## PLANES Documentation

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March 12, 2016

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# General overview

### 1.1 Presentation

PLANES (Porous LAum NumErical Simulator) project is a collection of Matlab/Fortran scripts to simulate the vibroacoustics response of coupled systems including acoustic, elastic, porous materials, PML...

# **Tables**

- interface $(n_1, n_2, e_1, e_2, n_{middle})$

# **Figures**

The number of the figure is xxxyyy. xxx is associated to the type of display and yyy corresponds to the type of field

#### Type of display xxx

- 1yyy: Modulus of the projection on x axis
- 2yyy: Modulus of the projection on y axis
- 3yyy: Angle of the projection on x axis
- 4yyy: Angle of the projection on y axis
- 5000: Shape of the displacement
- 10yyy: Map of the modulus for FEM
- 11yyy: Map of the angle for FEM
- 20yyy: Map of the modulus for DGM
- 21yyy: Map of the angle for DGM
- 30yyy: Map of the modulus in case of a FEM/DGM model
- 31yyy: Map of the angle in case of a FEM/DGM model

#### Type of fields xxx

- xxx1:  $v_x$  of the fluid or  $v_x^{eq}$  for an equivalent fluid or  $u_x^t$  for a Biot PEM.
- xxx2:  $v_y$  of the fluid or  $v_y^{eq}$  for an equivalent fluid or  $v_y^t$  for a Biot PEM
- xxx3: |v| of the fluid or  $|v^{eq}|$  for an equivalent fluid or  $u_y^t$  for a Biot PEM

• xxx4:  $v_x^s$  for a Biot PEM.

• xxx5:  $v_y^s$  for a Biot PEM

• xx10: p

## Weak forms

### 4.1 Finite-Elements

#### 4.1.1 Fluid

Unknown: the pressure field p.  $\rho$  is the density

$$\forall q, \quad \int_{\Omega} \frac{\nabla p \, \nabla q}{\rho \, \omega^2} - \frac{p \, q}{\rho \, c^2} \, d\Omega = \oint_{\partial \Omega} \frac{1}{\rho \, \omega^2} \frac{\partial p}{\partial n} \, d\Gamma \tag{4.1}$$

#### 4.1.2 Elastic solid

Unknown: the solid displacement  $\mathbf{u}$ .

$$\forall \mathbf{v}, \quad \int_{\Omega} \boldsymbol{\sigma}(\mathbf{u}) : \boldsymbol{\varepsilon}(\mathbf{v}) - \omega^2 \rho_s \mathbf{u}.\mathbf{v} \, d\Omega = \oint_{\partial \Omega} [\boldsymbol{\sigma}(\mathbf{u}).\mathbf{n}].\mathbf{v} \, d\Gamma$$
 (4.2)

### 4.2 Discontinuous Galerkin

#### 4.2.1 Abstract weak forms

$$\left(j\omega[\mathbf{M}] + [\mathbf{B}_x]\frac{\partial}{\partial x} + [\mathbf{B}_y]\frac{\partial}{\partial y}\right)\mathbf{S} = \mathbf{0} \text{ on } \Omega.$$
(4.3)

 ${f S}$  is the State Vector of the medium. In a conservative form:

$$\left(j\omega + [\mathbf{A}_x] \frac{\partial}{\partial x} + [\mathbf{A}_y] \frac{\partial}{\partial y}\right) \mathbf{S} = \mathbf{0} \text{ on } \Omega.$$
(4.4)

$$[\mathbf{A}_x] = [\mathbf{M}]^{-1}[\mathbf{B}_x], \quad [\mathbf{A}_y] = [\mathbf{M}]^{-1}[\mathbf{B}_y]. \tag{4.5}$$

 $\Omega$  is partitioned in n elements  $\Omega_e$  with e=1,...,n.

$$\sum_{e=1}^{n} \int_{\Omega_{e}} \mathbf{T}_{e}^{T} \left( j\omega[\mathbf{M}] \mathbf{S}_{e} + [\mathbf{B}_{x}] \frac{\partial \mathbf{S}_{e}}{\partial x} + [\mathbf{B}_{y}] \frac{\partial \mathbf{S}_{e}}{\partial y} \right) d\Omega = 0.$$
 (4.6)

 $\mathbf{T}_e$  is the test field on each element.

$$-\sum_{e=1}^{n} \int_{\Omega_{e}} \left( j\omega[\mathbf{M}]^{T} \mathbf{T}_{e} + [\mathbf{B}_{x}]^{T} \frac{\partial \mathbf{T}_{e}}{\partial x} + [\mathbf{B}_{y}]^{T} \frac{\partial \mathbf{T}_{e}}{\partial y} \right)^{T} \mathbf{S}_{e} d\Omega + \sum_{e=1}^{n} \int_{\partial \Omega_{e}} \mathbf{T}_{e}^{T} [\mathbf{G}_{e}] \mathbf{S}_{e} d\Gamma = 0.$$

where we have introduced the matrix  $[\mathbf{G}_e] = [\mathbf{B}_x]n_x + [\mathbf{B}_y]n_y$  which represents the normal fluxes across the boundary of the element  $\Omega_e$ . The unit normal  $\mathbf{n} = (n_x, n_y)$  on the element boundary  $\partial \Omega_e$  points out of element e. A key aspect of the wave-based DGM is to use test functions  $\mathbf{T}$  whose restrictions  $\mathbf{T}_e$  on each elements are solutions of the adjoint problem defined on each element:

$$\left(j\omega[\mathbf{M}]^T\mathbf{T}_e + [\mathbf{B}_x]^T\frac{\partial\mathbf{T}_e}{\partial x} + [\mathbf{B}_y]^T\frac{\partial\mathbf{T}_e}{\partial y}\right) = \mathbf{0} = 0.$$
(4.7)

With this choice of test functions the integral over each element  $\Omega_e$  vanishes and one is left with integrals on the interfaces between elements and on the boundary of the domain.

$$\sum_{e=1}^{n} \int_{\partial \Omega_e} \mathbf{T}_e^T [\mathbf{G}_e] \mathbf{S}_e \, \mathrm{d}\Gamma = 0. \tag{4.8}$$

$$[\mathbf{F}_e] = [\mathbf{A}_x] n_x + [\mathbf{A}_y] n_y \Longrightarrow [\mathbf{G}_e] = [\mathbf{M}_e] [\mathbf{F}_e]$$
(4.9)

$$\sum_{e=1}^{n} \int_{\partial \Omega_e} \mathbf{T}_e^T [\mathbf{M}_e] [\mathbf{F}_e] \mathbf{S}_e \, \mathrm{d}\Gamma = 0. \tag{4.10}$$

$$\sum_{e=1}^{n} \int_{\partial \Omega_e} \mathbf{V}_e[\mathbf{F}_e] \mathbf{S}_e \, d\Gamma = 0, \quad \mathbf{V}_e = [\mathbf{M}_e] \mathbf{T}_e. \tag{4.11}$$

#### 4.2.2 Characteristics

$$[\mathbf{F}_e][\mathbf{P}_e] = [\mathbf{P}_e][\mathbf{\Lambda}_e], \quad [\mathbf{F}_e] = [\mathbf{P}_e][\mathbf{\Lambda}_e][\mathbf{Q}_e], \quad [\mathbf{Q}_e] = [\mathbf{P}_e]^{-1}$$
 (4.12)

$$[\mathbf{\Lambda}_e] = \operatorname{diag}([\mathbf{\Lambda}_e^{in}], [\mathbf{\Lambda}_e^{out}], [\mathbf{0}]). \tag{4.13}$$

$$[\mathbf{F}_e] = [\mathbf{P}_e^{in}][\mathbf{\Lambda}_e^{in}][\mathbf{Q}_e^{in}] + [\mathbf{P}_e^{out}][\mathbf{\Lambda}_e^{out}][\mathbf{Q}_e^{out}]$$

$$(4.14)$$

#### 4.2.3 Boundary term

$$\mathbf{S}_e = [\mathbf{P}_e^{in}]\mathbf{S}_e^{in} + [\mathbf{P}_e^{out}]\mathbf{S}_e^{out} \tag{4.15}$$

$$[\mathbf{C}_e]\mathbf{S}_e = \mathbf{E}_e. \tag{4.16}$$

$$[\mathbf{C}_e][\mathbf{P}^{out}]\mathbf{S}_e^{out} = \mathbf{E}_e - [\mathbf{C}_e][\mathbf{P}^{in}]\mathbf{S}_e^{in}$$
(4.17)

$$\mathbf{S}_e^{out} = [\widetilde{\mathbf{R}}_e]\mathbf{S}_e^{in} + \widetilde{\mathbf{E}}_e, \tag{4.18}$$

$$[\widetilde{\mathbf{R}}_e] = -\left([\mathbf{C}_e][\mathbf{P}^{in}]\right)^{-1} [\mathbf{C}_e][\mathbf{P}^{in}], \quad \widetilde{\mathbf{E}}_e = \left([\mathbf{C}_e][\mathbf{P}^{in}]\right)^{-1} \mathbf{E}_e. \tag{4.19}$$

$$\mathbf{S}_{e} = \left( [\mathbf{P}^{in}] + [\mathbf{P}^{out}][\widetilde{\mathbf{R}}_{e}] \right) \mathbf{S}_{e}^{in} + [\mathbf{P}^{out}]\widetilde{\mathbf{E}}_{e}$$
(4.20)

$$\mathbf{S}_{e} = \underbrace{\left( [\mathbf{P}^{in}] + [\mathbf{P}^{out}][\widetilde{\mathbf{R}}_{e}] \right) [\mathbf{Q}_{e}^{in}]}_{[\widetilde{\mathbf{P}}_{e}]} \mathbf{S}_{e} + [\mathbf{P}^{out}]\widetilde{\mathbf{E}}_{e}$$
(4.21)

$$\int_{\partial\Omega_e} \mathbf{T}_e^T [\mathbf{M}_e] [\mathbf{F}_e] \mathbf{S}_e \, d\Gamma = \int_{\partial\Omega_e} \mathbf{T}_e^T [\mathbf{M}_e] [\mathbf{F}_e] \left[ [\widetilde{\mathbf{P}}_e] \mathbf{S}_e + [\mathbf{P}^{out}] \widetilde{\mathbf{E}}_e \right] \, d\Gamma$$
 (4.22)

$$= \int_{\partial\Omega_e} \mathbf{T}_e^T [\widetilde{\mathbf{F}}_e] \mathbf{S}_e \, \mathrm{d}\Gamma + \int_{\partial\Omega_e} \mathbf{T}_e^T \widetilde{\mathbf{S}}_e \, \mathrm{d}\Gamma$$
 (4.23)

$$[\widetilde{\mathbf{F}}_e] = [\mathbf{M}_e][\mathbf{F}_e][\widetilde{\mathbf{P}}_e], \quad \widetilde{\mathbf{S}}_e = [\mathbf{M}_e][\mathbf{F}_e][\mathbf{P}^{out}]\widetilde{\mathbf{E}}_e$$
 (4.24)

#### 4.2.4 Interfaces terms for DGM/DGM

$$I_{12} = \int_{\Gamma_{12}} \mathbf{T}_1[\mathbf{M}_1][\mathbf{F}_1] \mathbf{S}_1 \, \mathrm{d}\Gamma + \int_{\Gamma_{12}} \mathbf{T}_2[\mathbf{M}_2][\mathbf{F}_2] \mathbf{S}_2 \, \mathrm{d}\Gamma$$
 (4.25)

$$\mathbf{S}_1 = [\mathbf{P}_1^{in}]\mathbf{S}_1^{in} + [\mathbf{P}_1^{out}]\mathbf{S}_1^{out} \tag{4.26}$$

$$\mathbf{S}_2 = [\mathbf{P}_2^{in}]\mathbf{S}_2^{in} + [\mathbf{P}_2^{out}]\mathbf{S}_2^{out}. \tag{4.27}$$

$$[\mathbf{C}_1]\mathbf{S}_1 = [\mathbf{C}_2]\mathbf{S}_2. \tag{4.28}$$

$$[\mathbf{C}_1][\mathbf{P}_1^{out}]\mathbf{S}_1^{out} - [\mathbf{C}_2][\mathbf{P}_2^{out}]\mathbf{S}_2^{out} = [\mathbf{C}_2][\mathbf{P}_2^{in}]\mathbf{S}_2^{in} - [\mathbf{C}_1][\mathbf{P}_1^{in}]\mathbf{S}_1^{in}$$
(4.29)

$$\left\{\begin{array}{c} \mathbf{S}_{1}^{out} \\ \mathbf{S}_{2}^{out} \end{array}\right\} = \left[\widetilde{\mathbf{R}}\right] \left\{\begin{array}{c} \mathbf{S}_{1}^{in} \\ \mathbf{S}_{2}^{in} \end{array}\right\}.$$
(4.30)

$$[\widetilde{\mathbf{R}}] = [[\mathbf{C}_1][\mathbf{P}_1^{out}]] - [\mathbf{C}_2][\mathbf{P}_2^{out}]]^{-1} [-[\mathbf{C}_1][\mathbf{P}_1^{in}]][\mathbf{C}_2][\mathbf{P}_2^{in}]]$$
(4.31)

$$= \begin{bmatrix} \begin{bmatrix} \widetilde{\mathbf{R}}_{11} \end{bmatrix} & \begin{bmatrix} \widetilde{\mathbf{R}}_{12} \\ [\widetilde{\mathbf{R}}_{21}] & [\widetilde{\mathbf{R}}_{22}] \end{bmatrix}. \tag{4.32}$$

$$\mathbf{S}_{1} = \underbrace{\left( [\mathbf{P}_{1}^{in}] + [\mathbf{P}_{1}^{out}][\widetilde{\mathbf{R}}_{11}] \right) [\mathbf{Q}_{1}^{in}]}_{[\widetilde{\mathbf{P}}_{11}]} \mathbf{S}_{1} + \underbrace{[\mathbf{P}_{1}^{out}][\widetilde{\mathbf{R}}_{12}][\mathbf{Q}_{2}^{in}]}_{[\widetilde{\mathbf{P}}_{12}]} \mathbf{S}_{2}$$
(4.33)

$$\mathbf{S}_{2} = \underbrace{[\mathbf{P}_{2}^{out}][\widetilde{\mathbf{R}}_{21}][\mathbf{Q}_{1}^{in}]}_{[\widetilde{\mathbf{P}}_{21}]} \mathbf{S}_{1} + \underbrace{\left([\mathbf{P}_{2}^{in}] + [\mathbf{P}_{2}^{out}][\widetilde{\mathbf{R}}_{22}]\right)[\mathbf{Q}_{2}^{in}]}_{[\widetilde{\mathbf{P}}_{22}]} \mathbf{S}_{2}$$
(4.34)

$$I_{12} = \int_{\Gamma_{12}} \mathbf{T}_1[\widetilde{\mathbf{F}}_{11}] \mathbf{S}_1 \, \mathrm{d}\Gamma + \int_{\Gamma_{12}} \mathbf{T}_1[\widetilde{\mathbf{F}}_{12}] \mathbf{S}_2 \, \mathrm{d}\Gamma$$
 (4.35)

$$+ \int_{\Gamma_{12}} \mathbf{T}_2[\widetilde{\mathbf{F}}_{21}] \mathbf{S}_1 \, d\Gamma + \int_{\Gamma_{12}} \mathbf{T}_2[\widetilde{\mathbf{F}}_{22}] \mathbf{S}_2 \, d\Gamma$$
 (4.36)

$$[\widetilde{\mathbf{F}}_{11}] = [\mathbf{M}_1][\mathbf{F}_1][\widetilde{\mathbf{P}}_{11}], \quad [\widetilde{\mathbf{F}}_{12}] = [\mathbf{M}_1][\mathbf{F}_1][\widetilde{\mathbf{P}}_{12}]$$

$$(4.37)$$

$$[\widetilde{\mathbf{F}}_{21}] = [\mathbf{M}_2][\widetilde{\mathbf{P}}_{21}], \quad [\widetilde{\mathbf{F}}_{22}] = [\mathbf{M}_2][\widetilde{\mathbf{P}}_{22}]$$
 (4.38)

#### 4.2.5 Interfaces terms for FEM/FEM

$$I_{12} = -\int_{\Gamma_{12}} q_1 \frac{\mathbf{v}_1 \cdot \mathbf{n}_1}{j\omega} d\Gamma - \int_{\Gamma_{12}} q_2 \frac{\mathbf{v}_2 \cdot \mathbf{n}_2}{j\omega} d\Gamma$$
 (4.39)

- comes from the transposition of the right hand side integral.

Like in the previous section

$$\hat{\mathbf{S}}_1 = [\widetilde{\mathbf{P}}_{11}]\mathbf{S}_1 + [\widetilde{\mathbf{P}}_{12}]\hat{\mathbf{S}}_2 \tag{4.40}$$

$$\hat{\mathbf{S}}_2 = [\widetilde{\mathbf{P}}_{21}]\mathbf{S}_1 + [\widetilde{\mathbf{P}}_{22}]\mathbf{S}_2 \tag{4.41}$$

$$-\frac{\mathbf{v}_1 \cdot \mathbf{n}_1}{j\omega} = -\frac{1}{j\omega} \left[ \hat{\mathbf{S}}_1(1)\mathbf{n}_1(1) + \hat{\mathbf{S}}_1(2)\mathbf{n}_1(2) \right]$$
(4.42)

$$= [\widetilde{\mathbf{F}}_{11}]\mathbf{S}_1 + [\widetilde{\mathbf{F}}_{12}]\mathbf{S}_2 \tag{4.43}$$

with

$$\left[\widetilde{\mathbf{F}}_{11}\right] = -\frac{1}{i\omega} \left[ \left[\widetilde{\mathbf{P}}_{11}(1,:)\right] \mathbf{n}_{1}(1) + \left[\widetilde{\mathbf{P}}_{11}(2,:)\right] \mathbf{n}_{1}(2) \right]$$
(4.44)

$$[\widetilde{\mathbf{F}}_{12}] = -\frac{1}{i\omega} \left[ [\widetilde{\mathbf{P}}_{12}(1,:)] \mathbf{n}_1(1) + [\widetilde{\mathbf{P}}_{12}(2,:)] \mathbf{n}_1(2) \right]$$
 (4.45)

Concerning the second term

$$-\frac{\mathbf{v}_2 \cdot \mathbf{n}_2}{j\omega} = \frac{1}{j\omega} \left[ \hat{\mathbf{S}}_2(1) \mathbf{n}_2(1) + \hat{\mathbf{S}}_2(2) \mathbf{n}_2(2) \right]$$
(4.46)

$$= [\widetilde{\mathbf{F}}_{21}]\mathbf{S}_1 + [\widetilde{\mathbf{F}}_{22}]\mathbf{S}_2 \tag{4.47}$$

$$\left[\widetilde{\mathbf{F}}_{21}\right] = -\frac{1}{i\omega} \left[ \left[\widetilde{\mathbf{P}}_{21}(1,:)\right] \mathbf{n}_{2}(1) + \left[\widetilde{\mathbf{P}}_{21}(2,:)\right] \mathbf{n}_{2}(2) \right]$$
(4.48)

$$[\widetilde{\mathbf{F}}_{22}] = -\frac{1}{j\omega} \left[ [\widetilde{\mathbf{P}}_{22}(1,:)] \mathbf{n}_2(1) + [\widetilde{\mathbf{P}}_{22}(2,:)] \mathbf{n}_2(2) \right]$$
 (4.49)

$$p_1(x,y) \approx [\mathbf{N}_1(x,y)]\mathbf{p}_1, \quad \mathbf{p}_2(x,y) \approx [\mathbf{N}_2(x,y)]\mathbf{p}_2$$
 (4.50)

$$\mathbf{S}_{i} \approx \begin{bmatrix} [\mathbf{V}_{i}^{x}(x,y)] \\ [\mathbf{V}_{i}^{y}(x,y)] \\ [\mathbf{N}_{i}(x,y)] \end{bmatrix} \mathbf{p}_{i}$$
(4.51)

$$\left[\mathbf{V}_{i}^{x}(x,y)\right] = -\frac{1}{j\omega\rho} \left[ \frac{\partial \mathbf{N}_{i}(x,y)}{\partial x} \right], \quad \left[\mathbf{V}_{i}^{y}(x,y)\right] = -\frac{1}{j\omega\rho} \left[ \frac{\partial \mathbf{N}_{i}(x,y)}{\partial y} \right]$$
(4.52)

The boundary terms

$$\int_{\Gamma_{12}} q_i [\widetilde{\mathbf{F}}_{ij}] \mathbf{S}_j \, d\Gamma \approx \int_{\Gamma_{12}} \mathbf{q}_i^T \left[ \mathbf{N}_i(x, y) \right]^T \left[ \widetilde{\mathbf{I}}_{ij} \right] \mathbf{p}_j \, d\Gamma$$
 (4.53)

with

$$\left[\widetilde{\mathbf{I}}_{ij}\right] = \widetilde{\mathbf{F}}_{ij}(1) \left[ \mathbf{V}_{j}^{x}(x,y) \right] + \widetilde{\mathbf{F}}_{ij}(2) \left[ \mathbf{V}_{j}^{y}(x,y) \right] + \widetilde{\mathbf{F}}_{ij}(3) \left[ \mathbf{N}_{j}(x,y) \right]$$
(4.54)

#### 4.2.6 Interfaces terms for FEM/DGM

$$I_{12} = -\int_{\Gamma_{12}} q_1 \frac{\mathbf{v}_1 \cdot \mathbf{n}_1}{j\omega} d\Gamma + \int_{\Gamma_{12}} \mathbf{T}_2[\mathbf{M}_2][\mathbf{F}_2] \mathbf{S}_2 d\Gamma$$

$$(4.55)$$

- for the first comes from the transposition of the right hand side integral.

Like in the previous section

$$\hat{\mathbf{S}}_1 = [\widetilde{\mathbf{P}}_{11}]\mathbf{S}_1 + [\widetilde{\mathbf{P}}_{12}]\hat{\mathbf{S}}_2 \tag{4.56}$$

$$\hat{\mathbf{S}}_2 = [\widetilde{\mathbf{P}}_{21}]\mathbf{S}_1 + [\widetilde{\mathbf{P}}_{22}]\mathbf{S}_2 \tag{4.57}$$

$$-\frac{\mathbf{v}_1 \cdot \mathbf{n}_1}{j\omega} = \frac{1}{j\omega} \left[ \hat{\mathbf{S}}_1(1)\mathbf{n}_1(1) + \hat{\mathbf{S}}_1(2)\mathbf{n}_1(2) \right]$$
(4.58)

$$= [\widetilde{\mathbf{F}}_{11}]\mathbf{S}_1 + [\widetilde{\mathbf{F}}_{12}]\mathbf{S}_2 \tag{4.59}$$

$$[\widetilde{\mathbf{F}}_{11}] = -\frac{1}{j\omega} \left[ [\widetilde{\mathbf{P}}_{11}(1,:)] \mathbf{n}_1(1) + [\widetilde{\mathbf{P}}_{11}(2,:)] \mathbf{n}_1(2) \right]$$
(4.60)

$$\left[\widetilde{\mathbf{F}}_{12}\right] = -\frac{1}{j\omega} \left[ \left[\widetilde{\mathbf{P}}_{12}(1,:)\right] \mathbf{n}_{1}(1) + \left[\widetilde{\mathbf{P}}_{12}(2,:)\right] \mathbf{n}_{1}(2) \right]$$
(4.61)

$$[\widetilde{\mathbf{F}}_{21}] = [\mathbf{M}_2][\widetilde{\mathbf{P}}_{21}], \quad [\widetilde{\mathbf{F}}_{22}] = [\mathbf{M}_2][\widetilde{\mathbf{P}}_{21}][\widetilde{\mathbf{P}}_{22}]$$

$$(4.62)$$

$$I_{12} = \int_{\Gamma_{12}} q_1[\widetilde{\mathbf{F}}_{11}] \mathbf{S}_1 \, \mathrm{d}\Gamma + \int_{\Gamma_{12}} q_1[\widetilde{\mathbf{F}}_{12}] \mathbf{S}_2 \, \mathrm{d}\Gamma$$
 (4.63)

$$+ \int_{\Gamma_{12}} \mathbf{T}_{2}^{T} [\widetilde{\mathbf{F}}_{21}] \mathbf{S}_{1} d\Gamma + \int_{\Gamma_{12}} \mathbf{T}_{2}^{T} [\widetilde{\mathbf{F}}_{22}] \mathbf{S}_{2} d\Gamma$$

$$(4.64)$$

$$p_1(x,y) \approx [\mathbf{N}_1(x,y)]\mathbf{p}_1, \quad \mathbf{S}_2(x,y) \approx [\mathbf{N}_2(x,y)]\mathbf{X}_2$$
 (4.65)

$$\mathbf{S}_{1} \approx \begin{bmatrix} [\mathbf{V}_{1}^{x}(x,y)] \\ [\mathbf{V}_{1}^{y}(x,y)] \\ [\mathbf{N}_{1}(x,y)] \end{bmatrix} \mathbf{p}_{1}$$

$$(4.66)$$

$$\left[\mathbf{V}_{1}^{x}(x,y)\right] = -\frac{1}{j\omega\rho} \left[\frac{\partial \mathbf{N}_{1}(x,y)}{\partial x}\right], \quad \left[\mathbf{V}_{1}^{y}(x,y)\right] = -\frac{1}{j\omega\rho} \left[\frac{\partial \mathbf{N}_{1}(x,y)}{\partial y}\right]$$
(4.67)

$$\int_{\Gamma_{12}} q_1[\widetilde{\mathbf{F}}_{11}] \mathbf{S}_1 \, d\Gamma \approx \int_{\Gamma_{12}} \mathbf{q}_1^T \left[ \mathbf{N}_1(x, y) \right]^T \left[ \widetilde{\mathbf{I}}_{11} \right] \mathbf{p}_1 \, d\Gamma$$
(4.68)

with

$$[\widetilde{\mathbf{I}}_{11}] = \widetilde{\mathbf{F}}_{11}(1) \left[ \mathbf{V}_{1}^{x}(x,y) \right] + \widetilde{\mathbf{F}}_{11}(2) \left[ \mathbf{V}_{1}^{y}(x,y) \right] + \widetilde{\mathbf{F}}_{11}(3) \left[ \mathbf{N}_{1}(x,y) \right]$$
 (4.69)

$$\int_{\Gamma_{12}} q_1[\widetilde{\mathbf{F}}_{12}] \mathbf{S}_2 \, d\Gamma \approx \int_{\Gamma_{12}} \mathbf{q}_1^T [\widetilde{\mathbf{I}}_{12}] \mathbf{X}_2 \, d\Gamma$$
 (4.70)

with

$$[\widetilde{\mathbf{I}}_{12}] = [\mathbf{N}_1(x,y)]^T \,\widetilde{\mathbf{F}}_{12} \left[ \mathbf{N}_2(x,y) \right] \tag{4.71}$$

$$\int_{\Gamma_{12}} \mathbf{T}_{2}^{T} [\widetilde{\mathbf{F}}_{21}] \mathbf{S}_{1} d\Gamma \approx \int_{\Gamma_{12}} \mathbf{X}_{2}^{T} \left[ \mathbf{N}_{2}(x, y) \right]^{T} [\widetilde{\mathbf{I}}_{21}] \mathbf{p}_{1} d\Gamma$$
(4.72)

$$[\widetilde{\mathbf{I}}_{21}] = \widetilde{\mathbf{F}}_{11}(:,1) \left[ \mathbf{V}_{1}^{x}(x,y) \right] + \widetilde{\mathbf{F}}_{11}(:,2) \left[ \mathbf{V}_{1}^{y}(x,y) \right] + \widetilde{\mathbf{F}}_{11}(:,3) \left[ \mathbf{N}_{1}(x,y) \right]$$
 (4.73)

$$\int_{\Gamma_{12}} \mathbf{T}_{2}^{T} [\widetilde{\mathbf{F}}_{22}] \mathbf{S}_{2} \, d\Gamma \approx \int_{\Gamma_{12}} \mathbf{X}_{2}^{T} [\widetilde{\mathbf{I}}_{22}] \mathbf{X}_{2} \, d\Gamma$$

$$(4.74)$$

$$[\widetilde{\mathbf{I}}_{22}] = [\mathbf{N}_2(x,y)]^T [\widetilde{\mathbf{F}}_{22}] [\mathbf{N}_2(x,y)]$$

$$(4.75)$$

## Plane waves

### 5.1 Fluid

$$j\omega\rho v_x = -\frac{\partial p}{\partial x} \tag{5.1}$$

$$j\omega\rho v_y = -\frac{\partial p}{\partial y} \tag{5.2}$$

$$j\omega p = -\rho c^2 \left( \frac{\partial v_x}{\partial x} + \frac{\partial v_y}{\partial y} \right). \tag{5.3}$$

$$\mathbf{S} = \left\{ \begin{array}{c} v_x \\ v_y \\ p \end{array} \right\}. \tag{5.4}$$

 $[\mathbf{M}],\,[\mathbf{A}_x]$  and  $[\mathbf{A}_y]$  are square matrices of size m defined as

$$[\mathbf{M}] = \begin{bmatrix} \rho & 0 & 0 \\ 0 & \rho & 0 \\ 0 & 0 & \frac{1}{\rho c^2} \end{bmatrix} , \tag{5.5}$$

and

$$[\mathbf{B}_x] = \begin{bmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{bmatrix}, \quad [\mathbf{B}_y] = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{bmatrix}.$$
 (5.6)

## Finite-Element Matrices

### 6.1 Global system

#### 6.1.1 Unknown vector

- 1.  $X_{TR6}$
- 2.  $X_{H12}$
- 3.  $\mathbf{X}_{DGM}$
- 4.  $\mathbf{X}_{Bloch}$

$$\begin{bmatrix} [\mathbf{K}_{e}] - \omega^{2}[\mathbf{M}_{e}] & & & | \mathbf{I}_{R}^{e} & \mathbf{I}_{T}^{e} \\ & [\mathbf{K}] - \omega^{2} \widetilde{\rho}[\mathbf{M}] & -\widetilde{\gamma}[\mathbf{C}] & & | \mathbf{I}_{R}^{e} & \mathbf{I}_{T}^{e} \\ & -\widetilde{\gamma}[\mathbf{C}]^{t} & \frac{[\mathbf{H}]}{\omega^{2} \widetilde{\rho}_{eq}} - \frac{[\mathbf{Q}]}{\widetilde{K}_{eq}} & & | \mathbf{I}_{R}^{p} & \mathbf{I}_{T}^{p} \\ & & & \frac{[\mathbf{H}]}{\omega^{2} \rho_{0}} - \frac{[\mathbf{Q}]}{K_{0}} & | \mathbf{I}_{R}^{a} & \mathbf{I}_{T}^{a} \\ & & & | \mathbf{J}_{R}^{e} & \mathbf{I}_{T}^{a} \\ & \mathbf{J}_{T}^{e} & \mathbf{J}_{T}^{s} & \mathbf{J}_{T}^{p} & | \mathbf{J}_{T}^{a} \\ & & & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{a} \\ & & & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{a} \\ & & & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{a} \\ & & & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{a} \\ & & & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{a} & | \mathbf{J}_{T}^{a} \\ & & & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{a} & | \mathbf{J}_{T}^{a} \\ & & & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{a} & | \mathbf{J}_{T}^{a} \\ & & & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{a} & | \mathbf{J}_{T}^{a} \\ & & & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{a} & | \mathbf{J}_{T}^{a} \\ & & & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{a} & | \mathbf{J}_{T}^{a} \\ & & & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{a} & | \mathbf{J}_{T}^{a} \\ & & & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{a} \\ & & & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{a} & | \mathbf{J}_{T}^{a} \\ & & & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{a} \\ & & & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{a} \\ & & & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{a} & | \mathbf{J}_{T}^{a} \\ & & & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{e} \\ & & & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{e} \\ & & & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{e} \\ & & & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{e} \\ & & & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{e} \\ & & & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{e} \\ & & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{e} \\ & & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{e} \\ & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{e} \\ & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{e} \\ & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{e} \\ & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{e} \\ & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{e} \\ & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{e} \\ & | \mathbf{J}_{T}^{e} & | \mathbf{J}_{T}^{e$$

#### 6.2 Number of coefficient in matrices

### 6.3 DtN operators

#### 6.3.1 Fields properties in incident and transmission media

$$p^{inc}(x, y, t) = 1 e^{j(\omega t - k_x x - k_y y)} + \sum_{l \in \mathbb{Z}} e^{j(\omega t - k_x x + k_y (l)y)} R_l$$
(6.2)

$$k_x = k_0 \sin(\theta), \quad k_y = k_0 \cos(\theta), \quad \theta \in \left[0; \frac{\pi}{2}\right[, \quad k_x(l) = k_x^i + \frac{2\pi l}{D}\right]$$
 (6.3)

$$u^{inc}(x,y,t) = \frac{\nabla p^{inc}}{\rho_0 \omega^2} = \begin{vmatrix} -\frac{jk_x}{\rho \omega^2} \\ -\frac{jk_z}{\rho \omega^2} \\ -\frac{jk_z}{\rho \omega^2} \end{vmatrix} e^{j(\omega t - k_x x - k_y y)} + \sum_{l \in \mathbb{Z}} \begin{vmatrix} -\frac{jk_x}{\rho \omega^2} \\ \frac{jk_y(l)}{\rho \omega^2} \end{vmatrix} e^{j(\omega t - k_x x + k_y(l)y)} R_l \quad (6.4)$$

$$p^{tr}(x,y,t) = \sum_{l \in \mathbb{Z}} e^{j(\omega t - k_x x - k_y(l)y)} T_l$$
(6.5)

$$u^{tr}(x,y,t) = \frac{\nabla p^{tr}}{\rho_0 \omega^2} = \sum_{l \in \mathbb{Z}} \begin{vmatrix} -\frac{jk_x}{\rho \omega^2} \\ -\frac{jk_y(l)}{\rho \omega^2} \end{vmatrix} e^{j(\omega t - k_x x - k_y(l)y)} T_l$$
 (6.6)

#### 6.3.2 Formalisation of the fields

Let  $\mathbf{X}_l$  be an information vector associated to Bloch wave l.

$$\mathbf{Y} = \left\{ \begin{array}{c} 1 \\ \mathbf{X}_1 \\ \vdots \\ \mathbf{X}_n \end{array} \right\} \tag{6.7}$$

On each interface with the FEM model a physical field  $\varphi$  can be written

$$\varphi(x,y) = [\mathbf{\Omega}_F^{\varphi}]e^{-jk_x} + \sum_{l} [\mathbf{\Omega}_l^{\varphi}]\mathbf{X}_l e^{-jk_x(l)}.$$
 (6.8)

### 6.4 Case of a interface with fluid/Biot 1998

$$I_{FEM} = \int_{|\Gamma_R} \frac{1}{\rho \omega^2} \frac{\partial p^{inc}}{\partial n} q \, d\Gamma = \int_{|\Gamma_R} -u_y(x) q(x) \, d\Gamma$$
 (6.9a)

$$= \int_{|\Gamma_R} [\mathbf{\Omega}_F^{u_y}] e^{-jk_x} q \, d\Gamma - \sum_l \int_{|\Gamma_R} [\mathbf{\Omega}_l^{u_y}] \mathbf{X}_l e^{-jk_x(l)} q \, d\Gamma$$
 (6.9b)

- is highlighted as this term will end in the matrix system.

The information vector is scalar  $\mathbf{X}_l = R_l$  and correspond to the reflexion coefficient.

$$\left[\mathbf{\Omega}_{F}^{u_{y}}\right] = \frac{jk_{z}}{\rho_{0}\omega^{2}}, \quad \left[\mathbf{\Omega}_{l}^{u_{y}}\right] = \frac{jk_{z}(l)}{\rho_{0}\omega^{2}} \tag{6.9c}$$

The additional equation is the projection of the pressure

$$\int_{\Gamma_R} p(x)e^{jk_x(n)x} d\Gamma = \underbrace{\int_{\Gamma_R} e^{j\frac{2\pi nx}{D}} d\Gamma}_{D\delta_{n0}} + \underbrace{\int_{\Gamma_R} \sum_{l \in \mathbb{Z}} e^{-\frac{j2\pi(l-n)x}{D}} d\Gamma}_{D\delta_{nl}} d\Gamma R_l$$
 (6.10)

### 6.5 Case of a interface with an elastic medium

#### 6.5.1 Reflexion side

$$I_{FEM} = \int_{|\Gamma_R} \sigma_{xy} v_x - \sigma_{yy} v_y \, d\Gamma = \int_{|\Gamma_R} p v_y \, d\Gamma$$
 (6.11a)

$$= \int_{|\Gamma_R} [\mathbf{\Omega}_F^p] e^{-jk_x} q \, d\Gamma - \sum_l \int_{|\Gamma_R} [\mathbf{\Omega}_l^p] \mathbf{X}_l e^{-jk_x(l)} v_y \, d\Gamma$$
 (6.11b)

The information vector is scalar  $\mathbf{X}_l = R_l$  and correspond to the reflexion coefficient.

$$[\mathbf{\Omega}_F^p] = 1, \quad [\mathbf{\Omega}_I^p] = -1 \tag{6.11c}$$

The additional equation is the projection of the pressure

$$\int_{\Gamma_R} u_y(x)e^{jk_x(n)x} d\Gamma = \frac{-jk_y}{\rho_0\omega^2}D\delta_{n0} + \frac{jk_y(l)}{\rho_0\omega^2}D\delta_{nl}R_l$$
(6.12)

#### 6.5.2 Transmission side

$$I_{FEM} = \int_{|\Gamma_R} \sigma_{xy} v_x + \sigma_{yy} v_y \, d\Gamma = \int_{|\Gamma_R} p v_y \, d\Gamma$$
 (6.13a)

$$= \int_{|\Gamma_R} [\mathbf{\Omega}_F^p] e^{-jk_x} q \, d\Gamma - \sum_l \int_{|\Gamma_R} [\mathbf{\Omega}_l^p] \mathbf{X}_l e^{-jk_x(l)} v_y \, d\Gamma$$
 (6.13b)

The information vector is scalar  $\mathbf{X}_l = R_l$  and correspond to the reflexion coefficient.

$$[\mathbf{\Omega}_F^p] = -1, \quad [\mathbf{\Omega}_I^p] = 1 \tag{6.13c}$$

The additional equation is the projection of the pressure

$$\int_{\Gamma_T} u_y(x)e^{jk_x(n)x} d\Gamma = -\frac{jk_y(l)}{\rho_0\omega^2} D\delta_{nl}R_l$$
(6.14)

### 6.6 Case of integration of a plate medium

#### 6.6.1 Reflexion side

$$I_{FEM} = \int_{|\Gamma_R} -\sigma_{xz} v_x - \sigma_{zz} v_z \, d\Gamma \tag{6.15}$$

$$\left\{ \begin{array}{c}
 \sigma_{xz} \\
 u_z \\
 \sigma_{zz} \\
 u_x
 \end{array} \right\} = \begin{bmatrix}
 0 & 0 & 0 \\
 -\frac{jk_z}{\rho\omega^2} & \frac{jk_z}{\rho\omega^2} & 0 \\
 -1 & -1 & 0 \\
 0 & 0 & 1
 \end{bmatrix} \left\{ \begin{array}{c}
 1 \\
 R \\
 u_x
 \end{array} \right\} 
 \tag{6.16}$$

$$\left\{ \begin{array}{c} \sigma_{xz} \\ u_z \\ \sigma_{zz} \\ u_x \end{array} \right\}_{\mathbf{F}} = [\mathbf{T}] \left\{ \begin{array}{c} 0 \\ -\frac{jk_z}{\rho\omega^2} \\ -1 \\ 0 \end{array} \right\}, \quad \left\{ \begin{array}{c} \sigma_{xz} \\ u_z \\ \sigma_{zz} \\ u_x \end{array} \right\}_{\mathbf{R}} = [\mathbf{T}] \left[ \begin{array}{c} 0 & 0 \\ \frac{jk_z}{\rho\omega^2} & 0 \\ -1 & 0 \\ 0 & 1 \end{array} \right] \left\{ \begin{array}{c} R \\ u_x \end{array} \right\}$$
(6.17)

$$p_{|\Gamma_R}(x,z) = e^{-j(k_x^i x + k_z^i z)} + \sum_{l \in \mathbb{Z}} e^{-j(k_x(l)x - k_z(l)z)} R_l$$
(6.18)

$$\frac{\partial p}{\partial z|_{\Gamma_R}}(x,z) = -jk_z^i e^{-j(k_x^i x + k_z^i z)} + \sum_{l \in \mathbb{Z}} jk_y(l) e^{-j(k_x(l)x - k_z(l)z)} R_l$$
 (6.19)

$$p_{|\Gamma_R}(x,z) = \sum_{l \in \mathbb{Z}} e^{-j(k_x(l)x + k_z(l)z)} T_l$$
(6.20)

$$\frac{\partial p}{\partial z|_{\Gamma_T}}(x,z) = \sum_{l \in \mathbb{Z}} -jk_z(l)e^{-j(k_x(l)x + k_z(l)z)}T_l$$
(6.21)

Termes de surface

$$\frac{1}{\rho_a \omega^2} \int_{\Gamma_R} \frac{\partial p}{\partial n} q \, d\Gamma = \dots$$
 (6.22)

Equation supplmentaires

$$\int_{\Gamma_R} p(x)e^{jk_x(n)x} d\Gamma = \underbrace{\int_{\Gamma_R} e^{j\frac{2\pi nx}{D}} d\Gamma}_{D\delta_{n0}} + \underbrace{\int_{\Gamma_R} \sum_{l\in\mathbb{Z}} e^{-\frac{j2\pi(l-n)x}{D}} d\Gamma}_{D\delta_{nl}} R_l$$
(6.23)

## 6.7 Elementary matrix

#### 6.7.1 I

$$\hat{I} = \int_{a}^{b} (AX^{2} + BX + C)e^{-jk_{x}X} dX$$
 (6.24)

x = X - a

$$\hat{I} = e^{-jk_x a} \int_0^h \Phi(x) e^{-jk_x x} \, \mathrm{d}x$$
 (6.25)

$$\xi = -1 + \frac{2x}{h}, \quad x = \frac{h}{2}(\xi + 1)$$

$$\hat{I} = e^{-jk_x a} \frac{h}{2} \int_{-1}^{1} \Phi(\xi) e^{-jk_x \frac{h(\xi+1)}{2}} d\xi = e^{-jk_x \left(a + \frac{h}{2}\right)} \frac{h}{2} \int_{-1}^{1} \Phi(\xi) e^{-jk_x \frac{h\xi}{2}} d\xi$$
 (6.26)

# Benchmarks

### 7.1 Kundt Tube

### Sous projets

- 0: TR6
- 1: H12
- 2: TR6/H12
- $\bullet$  3: DGM on TR
- $\bullet$  4: DGM on H
- $\bullet~$  5: DGM on TR / DGM on H
- $\bullet$  6: H12 / DGM on H
- $\bullet$  7: TR6 / DGM on H

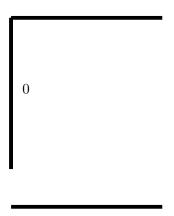


Figure 7.1: caption

## Hermite Finite-Elements

$$\Psi_1(x) = 1 - 3\left(\frac{x}{l_x}\right)^2 + 2\left(\frac{x}{l_x}\right)^3 \tag{8.1}$$

$$\Psi_2(x) = \left(\frac{x}{l_x}\right) - 2\left(\frac{x}{l_x}\right)^2 + \left(\frac{x}{l_x}\right)^3 \tag{8.2}$$

$$\Psi_3(x) = -\left(\frac{x}{l_x}\right)^2 + \left(\frac{x}{l_x}\right)^3 \tag{8.3}$$

$$\Psi_4(x) = 3\left(\frac{x}{l_x}\right)^2 - 2\left(\frac{x}{l_x}\right)^3 \tag{8.4}$$

$$\Psi_{ij}(x,y) = \Psi_i(x)\Psi_j(y) \tag{8.5}$$

Valeurs

	$\Psi_1$	$\Psi_2$	$\Psi_3$	$\Psi_4$
f(0)	1	0	0	0
$f(l_x)$	0	0	0	1
f'(0)	0	1/l	0	0
$f'(1_x)$	0	0	1/l	0

Ordre

	p	$\frac{\partial p}{\partial x}$	$\frac{\partial p}{\partial y}$	Ø
(0,0)	$\Psi_{11}$	$\Psi_{21}$	$\Psi_{12}$	$\Psi_{22}$
	1//1	2//2	3//3	$4//\emptyset$
(1,0)	$\Psi_{41}$	$\Psi_{31}$	$\Psi_{42}$	$\Psi_{23}$
	5//4	6//5	7//6	8//Ø
(1,1)	$\Psi_{44}$	$\Psi_{34}$	$\Psi_{43}$	$\Psi_{33}$
	9//7	10//8	11//9	$12//\emptyset$
(0,1)	$\Psi_{14}$	$\Psi_{24}$	$\Psi_{13}$	$\Psi_{32}$
	13//10	14//11	15//12	$16//\emptyset$

# Miscellaneous

#### 9.1 Integration of an exponential on one element

$$\int_{\Omega} e^{j\mathbf{k}\cdot(\mathbf{x}-\mathbf{x}_c)} d\Omega = \oint \nabla \cdot \left( - \begin{vmatrix} \frac{jk_x}{k_x^2 + k_y^2} \\ \frac{jk_y}{k_x^2 + k_y^2} \end{vmatrix} e^{j\mathbf{k}\cdot(\mathbf{x}-\mathbf{x}_c)} \right) d\Omega$$

$$= -\oint_{\partial\Omega} \left( \frac{jk_x n_x}{k_x^2 + k_y^2} + \frac{jk_y n_y}{k_x^2 + k_y^2} \right) e^{j\mathbf{k}\cdot(\mathbf{x}-\mathbf{x}_c)} d\Omega$$
(9.1)

$$= -\oint_{\partial\Omega} \left( \frac{jk_x n_x}{k_x^2 + k_y^2} + \frac{jk_y n_y}{k_x^2 + k_y^2} \right) e^{j\mathbf{k}\cdot(\mathbf{x} - \mathbf{x}_c)} d\Omega$$
 (9.2)

# Plane Waves 3D

### 10.1 Fluid medium

#### 10.1.1 Compressional wave

$$\delta_P^2 = k_x^2 + k_y^2 + k_z^2 \tag{10.1}$$

$$u^{\varepsilon}(x, y, z, t) = \begin{vmatrix} k_x \\ k_y \\ \varepsilon k_z \end{vmatrix} e^{j(\omega t - k_x x - k_y y - \varepsilon k_z z)}$$
(10.2)

$$p = (-K_0)(-j)(k_x^2 + k_y^2 + \varepsilon^2 k_z^2)$$
(10.3)

$$= jK_0\delta_P^2 \tag{10.4}$$

### 10.2 Elastic Medium

#### 10.2.1 Compressional wave

$$\delta_P^2 = k_x^2 + k_y^2 + k_z^2 \tag{10.5}$$

$$u^{+}(x, y, z, t) = \begin{vmatrix} k_x \\ k_y \\ \varepsilon k_z \end{vmatrix} e^{j(\omega t - k_x x - k_y y - \varepsilon k_z z)}$$
(10.6)

$$\sigma_{zz} = \lambda \left[ u_{x,x} + u_{y,y} + u_{z,z} \right] + 2\mu u_{z,z} \tag{10.7}$$

$$= -j\lambda(k_x^2 + k_y^2 + \varepsilon^2 k_z^2) + 2\mu(-j\varepsilon^2 k_z^2)$$
 (10.8)

$$= -j \left[ \lambda (k_x^2 + k_y^2) + (\lambda + 2\mu) k_z^2 \right]$$
 (10.9)

$$\sigma_{yz} = \mu \left[ u_{y,z} + u_{z,y} \right] = -j \left[ \varepsilon k_y k_z + \varepsilon k_z k_y \right] = -2j\mu \varepsilon k_y k_z \tag{10.10}$$

$$\sigma_{xz} = \mu \left[ u_{x,z} + u_{z,x} \right] = -j\mu \left[ \varepsilon k_x k_z + \varepsilon k_z k_x \right] = -2j\mu \varepsilon k_x k_z \tag{10.11}$$

#### 10.2.2 Shear wave 1

$$\delta_S^2 = k_x^2 + k_y^2 + k_z^2 \tag{10.12}$$

$$u^{\varepsilon}(x, y, z, t) = \begin{vmatrix} k_z \\ 0 \\ -\varepsilon k_x \end{vmatrix} e^{j(\omega t - k_x x - k_y y - \varepsilon k_z z)}$$
(10.13)

$$\sigma_{zz} = \lambda \left[ u_{x,x} + u_{y,y} + u_{z,z} \right] + 2\mu u_{z,z} \tag{10.14}$$

$$=2\mu jk_xk_z\tag{10.15}$$

$$\sigma_{yz} = \mu \left[ u_{y,z} + u_{z,y} \right] = -j\mu \varepsilon k_x k_y \tag{10.16}$$

$$\sigma_{xz} = \mu \left[ u_{x,z} + u_{z,x} \right] = -j\mu \left[ \varepsilon k_x k_z - \varepsilon k_z k_x \right] = -j\mu \varepsilon (k_z^2 - k_x^2) \tag{10.17}$$

### 10.2.3 Shear wave 2

$$\delta_S^2 = k_x^2 + k_y^2 + k_z^2 \tag{10.18}$$

$$u^{\varepsilon}(x, y, z, t) = \begin{vmatrix} 0 \\ k_z \\ -\varepsilon k_y \end{vmatrix} e^{j(\omega t - k_x x - k_y y - \varepsilon k_z z)}$$
(10.19)

$$\sigma_{zz} = \lambda \left[ u_{x,x} + u_{y,y} + u_{z,z} \right] + 2\mu u_{z,z} \tag{10.20}$$

$$=2\mu j k_y k_z \tag{10.21}$$

$$\sigma_{yz} = \mu \left[ u_{y,z} + u_{z,y} \right] = -j\mu \varepsilon (k_z^2 - k_y^2)$$
 (10.22)

$$\sigma_{xz} = \mu \left[ u_{x,z} + u_{z,x} \right] = -j\varepsilon k_x k_y \tag{10.23}$$

# Rotations

$$\frac{\partial f}{\partial x_1} = \frac{\partial F}{\partial X_1} \frac{\partial X_1}{\partial x_1} + \frac{\partial F}{\partial X_1} \frac{\partial X_1}{\partial x_2} + \frac{\partial F}{\partial X_1} \frac{\partial X_1}{\partial x_3}$$
(11.2)

$$= \mathbf{P}(1,:) \left\{ \frac{\frac{\partial}{\partial X_1}}{\frac{\partial}{\partial X_2}} \right\} F(X_1, X_2, X_3)$$
(11.3)

$$\left\{ \begin{array}{l} \frac{\partial}{\partial X_1} \\ \frac{\partial}{\partial X_2} \\ \frac{\partial}{\partial X_3} \end{array} \right\} F(X_1, X_2, X_3) = [\mathbf{Q}] \left\{ \begin{array}{l} \frac{\partial}{\partial x_1} \\ \frac{\partial}{\partial x_2} \\ \frac{\partial}{\partial x_3} \end{array} \right\} f(x_1, x_2, x_3) \tag{11.4}$$

$$\frac{\partial U_{i}}{\partial X_{j}} = \mathbf{Q}(i,:) \begin{cases} \frac{\partial}{\partial x_{1}} \\ \frac{\partial}{\partial x_{2}} \\ \frac{\partial}{\partial x_{3}} \end{cases} \mathbf{Q}(j,:) \begin{cases} u_{1} \\ u_{2} \\ u_{3} \end{cases}$$

$$= \mathbf{Q}(i,:) \begin{cases} \frac{\partial}{\partial x_{1}} \\ \frac{\partial}{\partial x_{2}} \\ \frac{\partial}{\partial x_{3}} \end{cases} (\mathbf{Q}_{j1}u_{1} + \mathbf{Q}_{j2}u_{2} + \mathbf{Q}_{j3}u_{3})$$

$$= \mathbf{Q}(i,:) \begin{cases} \mathbf{Q}_{j1} \frac{\partial u_{1}}{\partial x_{1}} + \mathbf{Q}_{j2} \frac{\partial u_{2}}{\partial x_{1}} + \mathbf{Q}_{j3} \frac{\partial u_{3}}{\partial x_{1}} \\ \mathbf{Q}_{j1} \frac{\partial u_{1}}{\partial x_{2}} + \mathbf{Q}_{j2} \frac{\partial u_{2}}{\partial x_{2}} + \mathbf{Q}_{j3} \frac{\partial u_{3}}{\partial x_{2}} \\ \mathbf{Q}_{j1} \frac{\partial u_{1}}{\partial x_{3}} + \mathbf{Q}_{j2} \frac{\partial u_{2}}{\partial x_{3}} + \mathbf{Q}_{j3} \frac{\partial u_{3}}{\partial x_{3}} \end{cases}$$

$$(11.5)$$

$$= \mathbf{Q}(i,:) \begin{cases} \frac{\partial}{\partial x_1} \\ \frac{\partial}{\partial x_2} \\ \frac{\partial}{\partial x_3} \end{cases} \left( \mathbf{Q}_{j1} u_1 + \mathbf{Q}_{j2} u_2 + \mathbf{Q}_{j3} u_3 \right)$$
(11.6)

$$= \mathbf{Q}(i,:) \left\{ \begin{aligned} \mathbf{Q}_{j1} \frac{\partial u_1}{\partial x_1} + \mathbf{Q}_{j2} \frac{\partial u_2}{\partial x_1} + \mathbf{Q}_{j3} \frac{\partial u_3}{\partial x_1} \\ \mathbf{Q}_{j1} \frac{\partial u_1}{\partial x_2} + \mathbf{Q}_{j2} \frac{\partial u_2}{\partial x_2} + \mathbf{Q}_{j3} \frac{\partial u_3}{\partial x_2} \\ \mathbf{Q}_{j1} \frac{\partial u_1}{\partial x_3} + \mathbf{Q}_{j2} \frac{\partial u_2}{\partial x_3} + \mathbf{Q}_{j3} \frac{\partial u_3}{\partial x_3} \end{aligned} \right\}$$

$$(11.7)$$

$$\frac{\partial U_i}{\partial X_j} = \sum_{k=1}^{3} \sum_{k'=1}^{3} Q_{ik} Q_{jk'} u_{k',k}$$
(11.8)

$$E_{ii} = Q_{i1}^2 \varepsilon_{11} + Q_{i2}^2 \varepsilon_{22} + Q_{i3}^2 \varepsilon_{33} + Q_{i1} Q_{i2} 2\varepsilon_{12} + Q_{i1} Q_{i3} 2\varepsilon_{13} + Q_{i2} Q_{i3} 2\varepsilon_{23}$$
(11.9)

$$2E_{ij} = (Q_{i1}Q_{j1} + Q_{j1}Q_{i1})\varepsilon_{11} + \dots (11.10)$$

$$+ (Q_{i1}Q_{j2} + Q_{j1}Q_{i2})2\varepsilon_{12} + \dots (11.11)$$