

# Chronotopic Metric Theory: Emergent Geometry, Gauge Structure, and Quantum Linearity from Kernel Coherence

Matěj Rada

Email: MatejRada@email.cz

This work is licensed under a Creative Commons Attribution–NonCommercial–NoDerivatives 4.0 License.

**Abstract**—Chronotopic Metric Theory (CTMT) is a metric-free, kernel-based framework in which spacetime geometry, gauge structure, and quantum linearity emerge from coherence transport and Fisher-curvature stability. No background manifold, metric, or Hilbert space is assumed. Instead, dimensionality, null propagation, gauge redundancy, and linear quantum evolution arise as dynamically selected configurations of a chronotopic kernel.

In this focused contribution, we present explicit structural derivations showing that: (i) null propagation and Maxwell invariants arise from rank-one Fisher sectors, (ii) Newtonian gravity emerges from curvature minimization under coherence sourcing, (iii) Schrödinger-type dynamics appear at a rigidity threshold where kernel expectations linearize, and (iv)  $U(1)$ ,  $SU(2)$ , and  $SU(3)$  gauge groups arise from redundancy closures in kernel phase labeling. No adjustable dimensionless parameters are introduced; known constants appear only as empirical rigidity thresholds. The results place CTMT within the broader class of emergent spacetime and information-geometric approaches [1]–[5].

## I. FOUNDATIONAL AXIOMS

**Definition I.1** (Chronotopic kernel). Physical configurations are described by a complex-valued chronotopic kernel

$$\mathcal{K} : \text{Anch} \times \text{Topo} \times \mathbb{R} \rightarrow \mathbb{C},$$

where  $\text{Anch}$  encodes discrete anchors (sources, detectors, interaction loci) and  $\text{Topo}$  encodes purely topological neighborhoods of kernel support. No metric or dimensional structure is assumed on these sets.

**Axiom 1 (Kernel primacy).** There is no fundamental background spacetime. All observable geometry arises from coherence relations and curvature of  $\mathcal{K}$ .

**Axiom 2 (Action invariant).** All physical processes preserve the dimensionally closed invariant

$$S_* = \frac{E}{\nu},$$

where  $E$  is an effective energy scale and  $\nu$  a characteristic frequency. In particular, for photons  $E = h\nu$  gives  $S_* = h$ .

**Axiom 3 (No fundamental Hilbert space).** Linear state spaces, inner products, and superposition are not fundamental. They emerge only in a rigid-phase regime, once the kernel spectrum stabilizes under coherence flow.

## II. KERNEL CALCULUS AND COHERENCE

**Definition II.1** (Observable field). Given an anchor configuration  $\kappa$  and kernel  $\mathcal{K}$ , an induced observable field is

$$\psi(x, t) = \int \mathcal{K}(x, x'; t) \kappa(x') dx'.$$

**Definition II.2** (Coherence density). Let  $\phi$  denote kernel phases sampled under a reconstruction protocol. The local coherence density is

$$\rho_{\text{coh}} = |\mathbb{E}[e^{i\phi}]|.$$

**Definition II.3** (-volume). The global coherence measure

$$= \int \rho_{\text{coh}} d\mu$$

quantifies kernel stability, where  $d\mu$  is a topology-compatible measure on the support of  $\mathcal{K}$ .

No metric appears in these definitions; all effective geometry is encoded in coherence, curvature, and transport properties of  $\mathcal{K}$ .

## III. FISHER CURVATURE AND RANK ATTRACTION

**Definition III.1** (Fisher information Hessian). Let  $\Theta$  denote kernel parameters. The Fisher information Hessian of a likelihood functional  $\mathcal{L}$  is

$$H_{ij} \equiv \nabla_{\Theta}^2 \log \mathcal{L},$$

with indices  $i, j$  running over admissible kernel directions.

**Remark III.2.** Coherent eigenspaces of  $H$  define Fisher curvature sectors. Their rank controls the number of independent curvature directions that can be stably supported.

**Proposition III.3** (Rank attraction). *CTMT dynamics preferentially collapse onto low-rank coherent eigenspaces of  $H$ .*

*Sketch.* Sustaining coherence along  $N$  independent curvature directions requires simultaneous phase stability across  $N$  modes. Noise and rupture scale superlinearly with  $N$ , while  $\rho_{\text{coh}} \leq 1$ . Higher-rank sectors therefore lose coherence faster, causing eigenvalue collapse and dynamical reduction of effective rank.  $\square$

#### IV. NULL PROPAGATION AND MAXWELL STRUCTURE

We now show explicitly how null propagation and Maxwell invariants arise from a minimal kernel phase.

**Definition IV.1** (Minimal coherent phase). Consider a phase field

$$\Phi(x, t) = \omega t - kx,$$

with constant coherence density  $\rho_{\text{coh}} = \rho_0$ .

The Fisher information Hessian with respect to  $(t, x)$  is

$$H_{ij} = \mathbb{E}[\partial_i \Phi \partial_j \Phi] = \begin{pmatrix} \omega^2 & -\omega k \\ -\omega k & k^2 \end{pmatrix}.$$

Its determinant vanishes:

$$\det H = \omega^2 k^2 - (\omega k)^2 = 0.$$

**Proposition IV.2** (Null Fisher sector and signal speed). *The phase  $\Phi(x, t) = \omega t - kx$  defines a null Fisher-curvature direction and a characteristic propagation speed*

$$v \equiv \frac{\omega}{k}.$$

In the rigid-phase limit, this speed coincides operationally with the universal signal speed  $c$ . Interpreting  $\Phi$  as a kernel phase potential and  $A_\mu = \partial_\mu \Phi$  as a connection, we obtain a field strength

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$$

with Lorentz invariants

$$I_1 = F_{\mu\nu} F^{\mu\nu}, \quad I_2 = \epsilon^{\mu\nu\rho\sigma} F_{\mu\nu} F_{\rho\sigma}.$$

Null coherence propagation corresponds to  $I_1 = I_2 = 0$ , i.e. vacuum electromagnetic waves in emergent spacetime.

#### V. EMERGENT NEWTONIAN GRAVITY

We next derive the Poisson equation from Fisher-curvature stability.

**Definition V.1** (Fisher curvature scalar). The scalar Fisher curvature is

$$\mathcal{K} \equiv \text{Tr } H = \sum_i \partial_i^2 \Phi.$$

Consider a homogeneous, isotropic coherence background.

**Proposition V.2** (Laplacian from curvature stability). *Rank stability under fixed coherence density  $\rho_{\text{coh}}$  minimizes*

$$\int [(\nabla \Phi)^2 - \alpha \rho_{\text{coh}} \Phi] d^3x.$$

The Euler–Lagrange equation yields

$$\nabla^2 \Phi = \alpha \rho_{\text{coh}}.$$

Dimensional analysis of kernel curvature shows

$$[\rho_{\text{coh}}] = \frac{\text{action}}{\text{volume}}.$$

In the rigid regime where action localizes as  $E \times t$ , we obtain

$$\rho_{\text{coh}} \longrightarrow \frac{E}{V},$$

identifying  $\rho_{\text{coh}}$  with mass–energy density. In the static, weak-field limit,  $\Phi$  becomes operationally identifiable with the Newtonian potential  $\Phi_N$ , so

$$\nabla^2 \Phi_N = 4\pi G \rho$$

with  $\alpha = 4\pi G$ . Thus Newtonian gravity emerges as the curvature stability condition under coherence sourcing.

#### VI. RIGIDITY AND EMERGENT QUANTUM LINEARITY

We now formalize the rigidity threshold at which linear quantum mechanics emerges.

**Definition VI.1** (Rigidity functional). Define the dimensionless rigidity

$$\mathcal{R} = \frac{\text{Var}(\Phi)}{S_*^2},$$

where  $S_*$  is the action invariant of Axiom 2.

Three regimes follow:

- **Classical averaging:**  $\mathcal{R} \gg 1$ , destructive interference dominates.
- **Nonlinear coherence:**  $\mathcal{R} \sim 1$ , coherent but nonlinear transport.
- **Rigid phase:**  $\mathcal{R} \ll 1$ , phase fluctuations are small.

Consider the kernel expectation of an observable functional  $O$ ,

$$\langle O \rangle = \int \Xi(\Theta) e^{i\Phi(\Theta)/S_*} \mathcal{D}\Theta.$$

In the rigid-phase regime  $\text{Var}(\Phi) \ll S_*^2$ , expand

$$e^{i\Phi/S_*} \approx 1 + \frac{i}{S_*} \Phi.$$

This yields

$$\langle O \rangle \approx \langle \Xi \rangle + \frac{i}{S_*} \langle \Xi \Phi \rangle.$$

Let  $\hat{H}$  denote the generator of kernel phase flow. Time translation of  $\Phi$  induces

$$\partial_t \langle O \rangle = \frac{i}{S_*} \langle [\hat{H}, O] \rangle,$$

where the commutator expresses the action of phase flow as an infinitesimal generator on observables.

Rigidity enforces self-adjointness of  $\hat{H}$  with respect to the induced inner product, so evolution is norm-preserving. Identifying  $S_* = \hbar$  in the rigid regime yields

$$i\hbar \partial_t \psi = \hat{H} \psi,$$

a Schrödinger-type evolution law. Thus Hilbert space, linear superposition, and unitary time evolution arise as emergent structures at the rigidity threshold, in line with information-geometric and entropic approaches to quantum theory [1]–[3].

## VII. GAUGE STRUCTURE FROM KERNEL REDUNDANCY

Gauge freedom arises from redundancy in kernel phase labeling.

**Definition VII.1** (Redundant phase shift). A phase shift

$$\Phi \mapsto \Phi + \chi$$

with  $\chi \in \ker H$  lies in a null Fisher sector and leaves observables invariant:  $\delta \langle O \rangle = 0$ .

### A. $U(1)$ from Single-Phase Holonomy

A single dominant null-curvature sector (X-axis) supports a global phase redundancy

$$\Phi_X \mapsto \Phi_X + \theta(x, t).$$

Local closure of such redundancies yields a compact Abelian group with one generator, isomorphic to  $U(1)$ . This sector corresponds to charge-like holonomy in the X-type null-curvature direction.

### B. $SU(2)$ from X/Y Torsional Coupling

The X and Y sectors form a two-dimensional torsional subspace with antisymmetric coupling:

$$[\theta_X, \theta_Y] \neq 0.$$

Closure under local rotations that preserve the Fisher norm defines a compact, simply connected group with three generators: two transverse modes and one torsional generator. This Lie algebra is isomorphic to  $\mathfrak{su}(2)$ , yielding  $SU(2)$  as the minimal non-Abelian redundancy group stabilizing X/Y coherence.

### C. $SU(3)$ from Triple-Channel Redundancy

Adding the Z (compressive) sector, redundancy closures must preserve:

- total coherence density,
- null transport in X,
- torsional balance in Y,
- compressive stability in Z.

The minimal simple Lie group acting transitively on three coupled complex coherence channels under these constraints is  $SU(3)$ . Larger groups such as  $SU(4)$  introduce surplus generators that exceed the redundancy capacity allowed by coherence–redundancy stabilization, destabilizing rank coherence and triggering collapse.

## VIII. DIMENSIONAL ATTRACTION AND FAILURE MODES

**Proposition VIII.1** (Rank collapse beyond three spatial curvatures). Assume four independent spatial Fisher-curvature directions  $\{\theta_i\}_{i=1}^4$ . The integrated action scales as

$$\mathcal{A} \sim \rho_{\text{coh}} L^4.$$

Stationary-phase coherence requires  $\mathcal{A} \sim S_*$ , so

$$L \sim \left( \frac{S_*}{\rho_{\text{coh}}} \right)^{1/4}.$$

This scale is incompatible with null-sector transport and torsional closure: curvature eigenvalues split, coherence ruptures, and one direction dynamically collapses. The system relaxes to three spatial-like curvature directions plus a null ordering direction.

Thus 3 + 1-dimensional structure and Standard Model-like gauge groups appear as attractors of CTMT dynamics rather than arbitrary choices.

## IX. CONCLUSION

We have presented a focused subset of Chronotopic Metric Theory in which geometry, gauge structure, and quantum linearity emerge from kernel coherence and Fisher curvature stability. Key results include:

- explicit emergence of null propagation and Maxwell invariants from rank-one Fisher sectors;
- derivation of Newtonian gravity as the weak-field limit of kernel curvature stability under coherence sourcing;
- appearance of Schrödinger-type dynamics at a rigidity threshold where kernel expectations linearize;
- emergence of  $U(1)$ ,  $SU(2)$ , and  $SU(3)$  as minimal redundancy closures in kernel phase labeling; and
- dynamical instability of higher spatial rank, leading to 3 + 1-dimensional attraction.

No adjustable dimensionless parameters are introduced; known constants enter only as empirical rigidity thresholds. These structural correspondences place CTMT within the broader landscape of emergent spacetime and information-geometric theories, while providing a concrete, kernel-based realization.

## APPENDIX

This appendix collects the structural arguments underlying the gravitational limit, quantum rigidity, gauge-group emergence, and dimensional failure modes.

### A. Derivation of the Laplacian from Fisher Curvature Stability

In CTMT, the scalar phase functional  $\Phi(\Theta)$  generates observable structure through its Fisher information Hessian

$$H_{ij} = \partial_i \partial_j \Phi. \quad (1)$$

Define the Fisher curvature scalar as the trace over admissible kernel directions:

$$\mathcal{K} \equiv \text{Tr } H = \sum_i \partial_i^2 \Phi. \quad (2)$$

In a homogeneous, isotropic coherence background, rank stability requires minimizing curvature fluctuations under fixed coherence density  $\rho_{\text{coh}}$ . This yields the variational condition

$$\delta \int [(\nabla\Phi)^2 - \alpha \rho_{\text{coh}} \Phi] d^3x = 0. \quad (3)$$

Taking the Euler–Lagrange equation gives

$$\nabla^2\Phi = \alpha \rho_{\text{coh}}. \quad (4)$$

Thus the Laplacian form is not assumed but follows from Fisher-curvature stability under coherence sourcing. In the weak-field, static, rigid limit,  $\Phi$  becomes operationally identifiable with the Newtonian potential  $\Phi_N$ , and  $\rho_{\text{coh}}$  with mass–energy density, yielding  $\alpha = 4\pi G$ .

#### B. Explicit Kernel Toy Model and Null Curvature Sector

Consider again the minimal coherent kernel phase field

$$\Phi(x, t) = \omega t - kx, \quad (5)$$

with constant coherence density

$$\rho_{\text{coh}} = \rho_0. \quad (6)$$

The Fisher information Hessian for translations in  $(x, t)$  is

$$H_{ij} = \mathbb{E}[\partial_i\Phi \partial_j\Phi] = \begin{pmatrix} \omega^2 & -\omega k \\ -\omega k & k^2 \end{pmatrix}. \quad (7)$$

Its determinant vanishes identically:

$$\det H = \omega^2 k^2 - (\omega k)^2 = 0. \quad (8)$$

This demonstrates the existence of a null Fisher-curvature direction without assuming a metric. Defining the characteristic propagation speed

$$v \equiv \frac{\omega}{k}, \quad (9)$$

the null condition fixes  $v$  uniquely.

In the rigid-phase limit, this speed coincides operationally with the universal signal speed  $c$ . This explicitly reproduces:

- null propagation,
- vanishing electromagnetic invariants  $I_1 = F_{\mu\nu}F^{\mu\nu} = 0$  and  $I_2 = \epsilon^{\mu\nu\rho\sigma}F_{\mu\nu}F_{\rho\sigma} = 0$ ,
- Maxwell-like wave propagation without postulated space-time structure.

#### C. Coherence Density as Mass–Energy Density

In CTMT, curvature arises from the second variation of kernel phase under coherent perturbations:

$$\delta^2\Phi \sim \int \rho_{\text{coh}} \delta x^2. \quad (10)$$

Dimensional analysis yields

$$[\rho_{\text{coh}}] = \frac{\text{action}}{\text{volume}}. \quad (11)$$

In the rigid limit where action localizes as energy times time, coherence density reduces to

$$\rho_{\text{coh}} \longrightarrow \frac{E}{V}, \quad (12)$$

identifying  $\rho_{\text{coh}}$  with mass–energy density. This justifies the identification used in the main text when passing to the Newtonian limit.

#### D. Rigidity Linearization and Schrödinger Evolution

Consider the kernel expectation

$$\langle O \rangle = \int \Xi(\Theta) e^{i\Phi(\Theta)/S_*} \mathcal{D}\Theta. \quad (13)$$

In the rigid-phase regime  $\text{Var}(\Phi) \ll S_*^2$ , expand to first order:

$$e^{i\Phi/S_*} \approx 1 + \frac{i}{S_*} \Phi. \quad (14)$$

Time translation of  $\Phi$  generates

$$\partial_t \langle O \rangle = \frac{i}{S_*} \langle [\hat{H}, O] \rangle, \quad (15)$$

where  $\hat{H}$  is the generator of kernel phase flow and the commutator expresses its action as an infinitesimal generator on observables.

Rigidity enforces self-adjointness of  $\hat{H}$  with respect to the induced inner product, yielding unitary evolution. Identifying  $S_* = \hbar$ , this reduces to Schrödinger-type dynamics

$$i\hbar \partial_t \psi = \hat{H} \psi. \quad (16)$$

Thus linear quantum mechanics arises as the unique stable evolution law at the rigidity threshold.

#### E. Structural Origin of $SU(2)$ and $SU(3)$

Gauge structure in CTMT emerges from redundancy closure in kernel phase labeling.

a)  *$SU(2)$  from X/Y Torsional Coupling.*: The X and Y sectors form a two-dimensional torsional subspace with antisymmetric coupling:

$$[\theta_X, \theta_Y] \neq 0. \quad (17)$$

Closure under local rotations preserving the Fisher norm defines a compact, simply connected group with three generators: two transverse modes and one torsional generator. This algebra is isomorphic to  $\mathfrak{su}(2)$ , yielding  $SU(2)$  as the minimal non-Abelian redundancy group stabilizing X/Y coherence.

b)  *$SU(3)$  from Triple-Channel Redundancy.*: Including the Z (compressive) sector, redundancy closure must preserve:

- total coherence density,
- null transport in X,
- torsional balance in Y,
- compressive stability in Z.

The minimal simple Lie group acting transitively on three coupled complex coherence channels under these constraints is  $SU(3)$ . Larger groups introduce surplus generators that destabilize rank coherence.

#### F. Explicit Failure Mode: Rank Collapse Beyond Three Spatial Curvatures

Assume four independent spatial Fisher-curvature directions  $\{\theta_i\}_{i=1}^4$ . The integrated action scales as

$$\mathcal{A} \sim \rho_{\text{coh}} L^4. \quad (18)$$

Stationary-phase coherence requires  $\mathcal{A} \sim S_*$ , implying

$$L \sim \left( \frac{S_*}{\rho_{\text{coh}}} \right)^{1/4}. \quad (19)$$

This scale is incompatible with null-sector transport and torsional closure: curvature eigenvalues split, coherence ruptures, and one direction dynamically collapses. The system relaxes to three spatial-like curvature directions plus a null ordering direction.

This explicit instability demonstrates that the observed dimensionality and gauge structure are attractors of the CTMT dynamics, not arbitrary choices.

#### *G. Non-Arbitrariness and Falsifiability*

The framework satisfies the following constraints:

- 1) invariant tensor structure precedes interpretation;
- 2) no adjustable dimensionless parameters are introduced;
- 3) constants appear only as rigidity thresholds;
- 4) explicit instability modes (rank collapse, coherence loss) are identified.

Violation of any constraint destroys the structural correspondence and serves to falsify the construction.

#### REFERENCES

- [1] S. Amari and H. Nagaoka, *Methods of Information Geometry*, American Mathematical Society (2000).
- [2] A. Caticha, *Entropic Dynamics, Time, and Quantum Theory*, J. Phys. A **44**, 225303 (2011).
- [3] C. Rovelli, *Relational Quantum Mechanics*, Int. J. Theor. Phys. **35**, 1637–1678 (1996).
- [4] X.-G. Wen, *Quantum Field Theory of Many-body Systems*, Oxford University Press (2004).
- [5] J. Baez and J. P. Muniain, *Gauge Fields, Knots and Gravity*, World Scientific (1994).