

# CTMT: Full Elemental Computation

*Phase-Aware Kernel, Fisher Geometry, Rupture, and Groupwise Constraints with Dimensional Closure*

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## Abstract

This document presents the current consolidated formulation of the *Chronotopic Theory of Matter and Time* (CTMT) for elemental-scale computation across atomic number  $Z$ . We specify a bona fide phase channel within the kernel, define the likelihood and Fisher geometry under dimensional closure, and provide the inversion scheme (Gauss–Newton with damping) together with rupture analysis on phase-rich observables. Groupwise regularity (including lanthanide contraction with endpoint anchoring) and admissibility of mass modulation in the presence of phase sensitivity are formalized. Beyond formal definitions, an appendix documents the *execution trace*: step-by-step numerical operations where causality, phase, Fisher geometry, and rupture are computed. Citations are provided to standard references for Fisher information, robust statistics, and Gauss–Newton damping.

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## 1 Scope and Minimal Assumptions

We consider a collection of elements indexed by atomic number  $Z \in \mathcal{Z} \subset \mathbb{Z}_{>0}$  and a set of experiments producing observables  $\mathcal{O}(Z) \in \mathbb{R}^m$  under a forward operator determined by a kernel with amplitude and phase structure. CTMT imposes **dimensional closure**—all terms entering the objective are dimensionless or integrated with unit-canceling measures—and a **closed topology** under which admissible transformations preserve identifiability and causality.

**Definition 1** (Element index set). The elemental index set is  $\mathcal{Z} = \{Z_1, Z_2, \dots, Z_n\} \subset \mathbb{Z}_{>0}$ , equipped with the natural order and any physically relevant partition into chemical groups.

**Assumption 1** (Minimal structure). (A1) **Kernel regularity:** The kernel  $\mathcal{K}$  is measurable on a time domain  $\mathcal{T} \times \mathcal{T}$  with  $\mathcal{T} \subset \mathbb{R}$ , causal in the sense that it has negligible support for  $t < t'$ , and decomposes into amplitude and phase channels (Section 2).

(A2) **Noise model:** Measurement noise is zero-mean with known or estimable covariance  $\Sigma$ , and independent across experiments conditional on parameters (Section 3).

(A3) **Phase sensitivity:** The phase channel  $\Phi(t, t')$  satisfies  $\partial_{t'} \Phi \neq 0$  on a set of non-zero measure within  $\mathcal{T} \times \mathcal{T}$  (Section 2.1).

(A4) **Dimensional closure:** All objective terms are rendered dimensionless either by reference scaling, integration against appropriate measures, or normalization by Fisher-consistent factors (Section 9).

## 2 Kernel with a Bona Fide Phase Channel

Let  $\mathcal{K} : \mathcal{T} \times \mathcal{T} \rightarrow \mathbb{C}$  be written as

$$\mathcal{K}(t, t'; \vartheta) = A(t, t'; \vartheta) \exp(i \Phi(t, t'; \vartheta)), \quad (1)$$

where  $A \geq 0$  is a real-valued amplitude channel and  $\Phi$  is a real-valued phase channel. The parameter set  $\vartheta$  may include dispersion coefficients and couplings to element-level latent variables.

### 2.1 Phase Sensitivity and Measurability

We require that  $\Phi$  be *measurably sensitive* to its second argument in the sense that

$$\mu(\{(t, t') \in \mathcal{T}^2 : |\partial_{t'} \Phi(t, t'; \vartheta)| > \epsilon\}) > 0 \quad \text{for some } \epsilon > 0, \quad (2)$$

where  $\mu$  is the product measure on  $\mathcal{T}^2$ . This condition allows recovery of conjugate quantities (§6) and renders mass modulation admissible within CTMT (§6.1).

### 2.2 Frequency-Domain Representation

Assuming  $\mathcal{K}$  is  $L^2$  in both variables and sufficiently decaying, define the partial Fourier transform in  $(t - t')$  and denote frequency by  $\Omega$ :

$$\hat{\mathcal{K}}(\Omega; t') = \int_{\mathcal{T}} \mathcal{K}(t, t') e^{-i\Omega(t-t')} dt. \quad (3)$$

A dispersive mapping  $k(\Omega)$  may be embedded via  $\Phi(t, t') = \int_{t'}^t \Omega(\tau) d\tau$  with  $\Omega \mapsto k(\Omega)$  tied to an observable such as spectral phase or group delay.

### 3 Forward Model, Observables, and Likelihood

Let  $x(Z) \in \mathbb{R}^p$  denote element-resolved latent parameters (e.g., radii, masses, couplings). Observables  $\mathcal{O}(Z) \in \mathbb{R}^m$  are generated by

$$\mathcal{O}(Z) = H(x(Z); \mathcal{K}(\cdot, \cdot; \vartheta)) + \eta, \quad \eta \sim \mathcal{N}(0, \Sigma), \quad (4)$$

where  $H$  integrates the kernel against experiment-specific probes. Split  $\mathcal{O}$  into *monotone* channels (e.g.,  $\kappa \propto Z$ , diffusivities) and at least one *phase-linked* non-monotone observable,  $\mathcal{O}_{\text{phase}}$ , such as coherence length or spectral phase curvature.

Given independent experiments indexed by  $i$ , the negative log-likelihood (NLL) obeys dimensional closure by either scaling or normalizing each channel:

$$\mathcal{L}(x, \vartheta) = \frac{1}{2} \sum_{Z \in \mathcal{Z}} \sum_i \|W_i^{1/2} (\mathcal{O}_i^{\text{meas}}(Z) - \mathcal{O}_i^{\text{mod}}(Z; x, \vartheta))\|_2^2 + \mathcal{R}(x, \vartheta), \quad (5)$$

with  $W_i$  dimensionless weights (often  $\Sigma_i^{-1}$  post-normalization) and  $\mathcal{R}$  an admissible regularizer (§8).

### 4 Fisher Geometry and Identifiability

Let  $\theta$  collect all free parameters in  $\{x(Z)\}_{Z \in \mathcal{Z}}$  and  $\vartheta$ . The (expected) Fisher information matrix (FIM) is

$$\mathcal{F}(\theta) = \mathbb{E} \left[ \nabla_{\theta} \log p(\mathcal{D} | \theta) \nabla_{\theta} \log p(\mathcal{D} | \theta)^{\top} \right] \approx J(\theta)^{\top} W J(\theta), \quad (6)$$

where  $J$  is the Jacobian of residuals [1]. CTMT favors *isotropic* Fisher spectra along the inversion path; thus, we **rebalance** the observable set to avoid dominance by monotone channels, improving conditioning and Gauss–Newton performance.

**Proposition 1** (Phase channel improves local conditioning). If  $\mathcal{O}_{\text{phase}}$  is informative with non-collinear sensitivity to  $x$ , then the smallest eigenvalue of  $\mathcal{F}$  increases under its inclusion, reducing the condition number  $\kappa(\mathcal{F})$  provided  $W$  is scaled under closure.

#### 4.1 Elemental Fisher Trajectory

Across atomic number, CTMT monitors the Fisher spectrum as a function of  $Z$ :

$$\lambda_{\min}(Z), \quad \kappa(Z) = \frac{\lambda_{\max}(Z)}{\lambda_{\min}(Z)}. \quad (7)$$

Stable elemental computation requires bounded  $\kappa(Z)$  under (i) reweighting of monotone observables, (ii) inclusion/exclusion of the phase channel, and (iii) moderate regularization changes. Sharp increases in  $\kappa(Z)$  indicate loss of causal identifiability and precede observable rupture in  $Y(Z)$ .

### 5 Inversion: Gauss–Newton with Damping

Define residuals  $r(\theta) = W^{1/2}(\mathcal{O}^{\text{meas}} - \mathcal{O}^{\text{mod}})$  and  $J = \partial r / \partial \theta$ . One step of damped Gauss–Newton (GN) solves

$$(J^{\top} J + \lambda I) \Delta \theta = -J^{\top} r, \quad \theta_{k+1} = \theta_k + \Delta \theta, \quad (8)$$

with Levenberg–Marquardt damping  $\lambda \geq 0$  [2, 3]. Convergence is declared when  $\|\Delta \theta\|$  and the reduction in  $\mathcal{L}$  fall below tolerance. Success rates improve materially after rebalancing with  $\mathcal{O}_{\text{phase}}$ .

## 6 Conjugate Recovery and Mass Modulation

Consider a conjugate quantity  $S_*$  (e.g., an effective source) linked to the phase-bearing kernel via a linear operator  $\mathcal{A}$  on the same grid:

$$\mathcal{A}[S_*](t) = \int_{\mathcal{T}} \mathcal{K}(t, t') S_*(t') dt'. \quad (9)$$

**Proposition 2** ( $S^*$  recovery under phase sensitivity). If (2) holds and  $A(t, t')$  is non-degenerate on the support of  $S_*$ , then  $S_*$  is identifiable (up to null-space symmetries) from phase-bearing observables, with stability governed by the Fisher spectrum of the induced linearized map.

### 6.1 Admissibility of Mass Modulation (AAA)

Mass (or mass-like) modulation of a component AAA is admissible under CTMT only when the kernel carries a measurable phase channel per (2); otherwise the modulation is declined as non-identifiable under closure. With phase present, modulation parameters enter  $\theta$  and are regularized by Fisher-consistent priors.

## 7 Rupture Analysis on Phase-Rich Spectra

Let  $Y(Z)$  denote a scalar phase-linked observable across elements. We define a robust rupture score using standardized second differences:

$$R(Z_k) = \frac{\Delta^2 Y(Z_k) - \text{med}(\Delta^2 Y)}{1.4826 \text{ med} |\Delta^2 Y - \text{med}(\Delta^2 Y)|}, \quad \Delta^2 Y(Z_k) = Y(Z_{k+1}) - 2Y(Z_k) + Y(Z_{k-1}). \quad (10)$$

This statistic uses the median and MAD for robustness to outliers [4]. Kinks (change-points) satisfy  $|R(Z_k)| > \tau$  with  $\tau \approx 3.5$  and are validated by robustness to rebinning and mild regularization changes.

## 8 Groupwise Structure and Regularization

Partition  $\mathcal{Z}$  into chemical groups  $G_g$  (alkali, alkaline earth, transitions, pnictogens, chalcogens, halogens, noble gases, lanthanides, actinides). For a parameter field  $x(Z)$  we impose piecewise-smooth priors with optional endpoint anchoring:

$$\mathcal{R}(x) = \sum_g \alpha_g \sum_{Z \in G_g} (\nabla_Z x(Z))^2 + \sum_{(Z_a, Z_b) \in \mathcal{A}} \beta_{ab} (x(Z_a) - x_a^\circ)^2 + (x(Z_b) - x_b^\circ)^2, \quad (11)$$

where  $\mathcal{A}$  collects anchored endpoints  $(Z_a, Z_b)$  with targets  $x^\circ$ .

### 8.1 Lanthanide Contraction

For  $Z = 57 \dots 71$ , we enforce endpoint anchoring and monotone curvature consistent with CTMT rigidity:

$$\Delta_Z x(Z) < 0 \quad \text{for successive } Z, \quad x(57) = x_{\text{La}}^\circ, \quad x(71) = x_{\text{Lu}}^\circ. \quad (12)$$

External validation can be performed against Shannon ionic radii at coordination number 8; errors are summarized by

$$\text{MAE} = \frac{1}{n} \sum_k |x(Z_k) - x_{\text{ref}}(Z_k)|, \quad \text{RMSE} = \sqrt{\frac{1}{n} \sum_k (x(Z_k) - x_{\text{ref}}(Z_k))^2}. \quad (13)$$

For reference data, see [5].

## 9 Dimensional Closure

CTMT requires that each cost term be dimensionless or accompanied by a unit-canceling measure. Typical strategies include: (i) normalizing observables by reference scales, (ii) prewhitening with physically grounded  $W$ , (iii) scaling regularization by Fisher-consistent factors so that  $\mathcal{L}$  is unitless. This extends to all regimes and entities in the ontology, not only specific measurement protocols.

## 10 Computation Over Elements

For each  $Z \in \mathcal{Z}$ :

1. Assemble  $\mathcal{O}(Z)$  including at least one phase-linked observable  $O_{\text{phase}}$ .
2. Evaluate residuals and Jacobians under (5), enforcing closure.
3. Update  $\theta$  via (8) with trust adjustment of  $\lambda$ .
4. Monitor Fisher spectra and condition numbers; if anisotropy emerges, rebalance weights or augment observables.
5. Compute rupture score (10) on phase-rich series  $Y(Z)$ ; register kinks.
6. Apply groupwise constraints (11); in lanthanides, enforce (12).

## 11 Acceptance Criteria

A solution is accepted if:

- (C1) **Phase channel:**  $\partial_{t'}\Phi \neq 0$  on non-zero measure;  $S_*$  reconstruction on the kernel grid attains finite error bounded by Fisher conditioning.
- (C2) **Fisher isotropy:** Inclusion of  $O_{\text{phase}}$  reduces  $\kappa(\mathcal{F})$  relative to baseline and improves Gauss–Newton success rate.
- (C3) **Rupture:**  $R(Z)$  reveals stable kinks around physically expected regions (e.g., radioactives) with robustness to smoothing.
- (C4) **Lanthanides:** Contraction monotonicity holds under endpoint anchoring; external MAE/RMSE are within admissible bounds.
- (C5) **Closure:** The complete objective remains dimensionless; any mass modulation parameters are admissible only with the verified phase channel.

## A Execution Trace: Step-by-Step Guidance

This appendix translates the formal sections into a minimal execution trace for CTMT elemental computation.

### Step 1: Kernel Instantiation (Causality + Phase)

For fixed  $Z$ , define  $\mathcal{K}_Z(t, t') = A_Z(t, t') e^{i\Phi_Z(t, t')}$ . Numerically tabulate or parameterize  $A_Z$  and  $\Phi_Z$ . Verify phase sensitivity:  $\partial_{t'}\Phi_Z \neq 0$  on non-zero measure (Section 2.1). If this fails, CTMT halts and any mass modulation is forbidden (Section 6.1).

### Step 2: Observable Synthesis

Compute  $\mathcal{O}(Z) = H(x(Z); \mathcal{K}_Z)$ , ensuring at least one *phase-linked* observable depends non-linearly on  $\Phi_Z$ ; others may be monotone. This asymmetry is intentional to improve Fisher conditioning (Sections 3 and 4).

### Step 3: Dimensionally Closed Residuals

Form residuals  $r_i(Z) = W_i^{1/2}(\mathcal{O}_i^{\text{meas}}(Z) - \mathcal{O}_i^{\text{mod}}(Z))$  with  $W_i$  chosen to enforce dimensional closure (Section 9). Reject any term with hidden unit leakage.

### Step 4: Fisher Geometry Construction

Compute the Jacobian  $J_Z = \partial r(Z)/\partial \theta$  and the Fisher matrix  $\mathcal{F}_Z = J_Z^\top J_Z$  [1]. Test: full rank, acceptable isotropy, and whether adding the phase observable increases  $\lambda_{\min}(\mathcal{F}_Z)$ .

### Step 5: Gauss–Newton Update (Conditional)

Only if Fisher geometry passes, solve  $(J^\top J + \lambda I)\Delta\theta = -J^\top r$  with Levenberg–Marquardt damping  $\lambda \geq 0$  [2, 3]. GN is subordinate to Fisher geometry; degrade *or* damp updates if conditioning worsens.

### Step 6: Rupture Computation (Across $Z$ )

After sweeping  $Z$ , compute  $R(Z)$  per (10). Interpret large  $|R(Z)|$  as a break in transport coherence (structural change), not mere fit error. Validate kinks by robustness checks [4].

### Step 7: Groupwise Enforcement

Apply constraints such as  $\Delta_Z x(Z) < 0$  in lanthanides (Section 8). Treat these as CTMT closure conditions justified by Fisher rigidity, not arbitrary priors.

## B Minimal Elemental Computation Loop (Pseudo-code)

```
for Z in Z_set:
    K = build_kernel(A_Z, Phi_Z)           # Check causality; verify d/dt' Phi_Z != 0
    O_mod = forward_model(x[Z], K)         # Includes at least one phase-linked observable
    r = W_sqrt @ (O_meas[Z] - O_mod)       # Closure: r is dimensionless
    J = jacobian(r, theta)                 # Numeric differentiation or adjoint
    F = J.T @ J                           # Fisher geometry
    if passes_geometry(F):
        dtheta = solve((J.T@J+ lambda*I) * dtheta == -J.T * r) # LM damping
        theta = update(theta, dtheta)
    else:
        adjust_weights_or_observables()
# Across Z: rupture on phase-rich series
R = robust_second_diff(Y_by_Z)           # median/MAD standardized
kinks = { Z: abs(R[Z]) > 3.5 }
# Enforce groupwise constraints
apply_groupwise_constraints(x)
```

## References

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