Part I: Linear ODEs Dynamical System

UNSW MATHEMATICS SOCIETY PRESENTS

MATH2121/2221 Revision Seminar

(Higher) Theory & Applications of Differential Equations

Seminar I / II

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Part I: Linear ODEs Dynamical Systems

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Part I: Linear ODEs

Differential Operators

Definition 1: Linear differential operators

Define the linear differential operator L of order m to be

$$Lu(x) = \sum_{j=0}^{m} a_j(x) \cdot D^j u(x)$$

= $a_m D^m u(x) + a_{m-1} D^{m-1} u(x) + \dots + a_0 u$,

where
$$D^j u = \frac{d^j u}{du^j}$$
 and $D^0 u = u$.

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$$D^j u = \frac{d^j u}{du^j}$$
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Definition 2: Singular ODEs

An ODE is said to be singular with respect to [a, b] if the leading coefficient vanishes for any $x \in [a, b]$.

Definition 3: Homogeneous ODEs

An ODE is said to be homogeneous if the right hand side is 0. That is, we have a differential equation of the form

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An ODE is said to be inhomogeneous if the right hand side is not 0. Then we can write the differential equation as

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$$Lu = f$$
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• **Example**: $u'' + u' + u = \cos(x)$.

Initial-valued problems (IVP)

Definition 5: Initial-valued problems

Consider an m-th order differential equation

$$Lu = f, \quad \text{on } [a, b] \tag{1}$$

along with the values

$$u(a) = v_0, u'(a) = v_1, ..., u^{(m-1)}(a) = v_{m-1}.$$
 (2)

The problem (1) with (2) is called an initial-valued problem.

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The problem (1) with (2) is called an initial-valued problem.

- **Example**: u' + u = x, u(0) = 0.
- **Solution**: $u(x) = x 1 + e^{-x}$.

Theorem 1: Unique solution for IVP

If f is continuous on [a, b] and the ODE Lu = f is non-singular on [a, b], then the IVP (1) and (2) has a unique solution.

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If L is a linear m-th order differential operator and non-singular on [a,b], then the set of all solutions to the homogenous equation Lu=0 on [a,b] forms a vector space of dimension m.

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• What does this mean?: The solution space has a basis of dimension m, with elements $u_1, ..., u_m$. And so every solution to the homogeneous equation can be written as a linear combination of this basis:

$$u(x) = c_1 u_1(x) + ... + c_m u_m(x), \quad \forall x \in [a, b].$$

This is called the general solution.

We can do the same for an inhomogeneous equation Lu=f by fixing a particular solution u_P . Then for any solution u, $L(u-u_p)=f-f=0$ and so we can write $u-u_p$ as a linear combination of the homogeneous equation basis:

$$u(x) - u_p(x) = c_1 u_1(x) + ... + c_m u_m(x), \quad \forall x \in [a, b].$$

Rearranging, we have the general solution for an inhomogeneous differential equation:

$$u(x) = u_P(x) + c_1 u_1(x) + ... + c_m u_m(x), \quad \forall x \in [a, b].$$

Example 1

Solve the second order ODE

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$$u(x) = c_1 e^{-x} + c_2 e^{x/2} - \cos(x) - 3\sin(x).$$

Linear Independence

Definition 6: Linear independence

Let $u_1(x),...,u_m(x)$ be functions on some interval $I \subset \mathbb{R}$. If there exists non-zero constants $a_1,...,a_m$ such that

$$a_1u_1(x) + ... + a_mu_m(x) = 0, \quad \forall x \in I,$$

then we say that $u_1, ..., u_m$ are linearly dependent.

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Part I: Linear ODEs

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then we say that $u_1, ..., u_m$ are **linearly dependent**. If the above equation only holds true for all constants zero, then we say that $u_1, ..., u_m$ are **linearly independent**.

Definition 7: Wronskian

The **Wronskian** of the functions $u_1, ..., u_m$ is the $m \times m$ determinant

$$W(x) = W(x; u_1, ..., u_m) = \det(D^{i-1}u_j).$$

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For example, a 3×3 Wronskian is

$$W(x) = W(x; u_1, ..., u_m) = \det \begin{pmatrix} u_1 & u_2 & u_3 \\ u'_1 & u'_2 & u'_3 \\ u''_1 & u''_2 & u''_3 \end{pmatrix}.$$

Lemmas

Lemma 1

If $u_1,...,u_m$ are linearly dependent over an interval $I \subset \mathbb{R}$ then $W(x;u_1,...,u_m)=0$ for all $x \in I$.

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Lemma 2

If $u_1, ..., u_m$ are solutions to the ODE

$$a_m(x)u^{(m)}(x) + a_{m-1}(x)u^{(m-1)}(x) + ... + a_0(x)u(x) = 0$$

on an interval $I \subset \mathbb{R}$, then the Wronskian satisfies

$$a_m(x)W'(x) + a_{m-1}(x)W(x) = 0 \quad \forall x \in I.$$

Example 2: MATH2221 2014 T2 2.iii).b

Given the functions u_1, u_2 , prove that if they are solutions to a second-order, homogeneous linear differential equation

$$a_2(x)u'' + a_1(x)u' + a_0(x)u = 0,$$

then the Wronskian W satisfies

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$$W = u_1 u_2' - u_1' u_2$$
, $W' = u_1 u_2'' - u_1'' u_2$.

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$$W = u_1 u_2' - u_1' u_2, \ W' = u_1 u_2'' - u_1'' u_2.$$

$$a_2 W' + a_1 W = a_2 (u_1 u_2'' - u_1'' u_2) + a_1 (u_1 u_2' - u_1' u_2)$$

$$= u_1 (a_2 u_2'' + a_1 u_2') - u_2 (a_2 u_1'' + a_1 u_1').$$

Add and subtract $a_0 u_1 u_2$:

$$a_2W' + a_1W = u_1(a_2u_2'' + a_1u_2' + a_0u_2) - u_2(a_2u_1'' + a_1u_1' + a_0u_1).$$

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Since u_1 and u_2 are solutions to the ODE, the RHS is 0.

Linear Independence of solutions

Theorem 3: Linear independence

Let $u_1, ..., u_m$ be solutions to the non-singular, linear, homogeneous m-th order ODE Lu = 0 on the interval [a, b]. Then either

$$W(x) = 0$$
 and the m solutions are linearly dependent,

or

 $W(x) \neq 0$ and the m solutions are linearly independent.

Polynomial solution guess

Theorem 4

Let L = p(D) be a linear differential operator of order m with constant coefficients. Assume that $p(0) \neq 0$. Then for any integer $r \geq 0$, there exists a unique polynomial u_P of degree r such that $Lu_P = x^r$.

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Let L = p(D) be a linear differential operator of order m with constant coefficients. Assume that $p(0) \neq 0$. Then for any integer $r \geq 0$, there exists a unique polynomial u_P of degree r such that $Lu_P = x^r$.

This means that if our linear ODE has a polynomial on the RHS, we should guess a polynomial of the same degree for our particular solution.

Exponential solution guess

Theorem 5

Let L = p(D) and $\mu \in \mathbb{C}$. If $p(\mu) \neq 0$, then the function

$$u_P(x) = \frac{e^{\mu x}}{p(\mu)}$$

satisfies $Lu_P = e^{\mu x}$.

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Let L = p(D) and $\mu \in \mathbb{C}$. If $p(\mu) \neq 0$, then the function

$$u_P(x) = \frac{e^{\mu x}}{p(\mu)}$$

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This means that we should guess a multiple of $e^{\mu x}$ when the RHS of a linear ODE is $e^{\mu x}$ if it is not already a solution of the homeogeneous solution.

Polynomial + exponential

Theorem 6

Let L = p(D) and assume $\mu \in \mathbb{C}$. If $p(\mu) \neq 0$ then for any integers $r \geq 0$, there exists a unique polynomial v of degree r such that

$$u_P(x) = v(x)e^{\mu x}$$

satisfies $Lu_p = x^r e^{\mu x}$.

Polynomial + exponential

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Let L = p(D) and assume $\mu \in \mathbb{C}$. If $p(\mu) \neq 0$ then for any integers $r \geq 0$, there exists a unique polynomial v of degree r such that

$$u_P(x) = v(x)e^{\mu x}$$

satisfies $Lu_p = x^r e^{\mu x}$.

So if the RHS of the inhomogeneous linear ODE is a polynomial times an exponential, and the exponential isnt in the homogeneous solution, then the guess for the particular solution should be a polynomial times exponential.

General solutions

Example 3: MATH2221 2014 T2 2.ii)

Let
$$p(z) = (z-1)(z+2)^2(z^2+1)$$
 and $D = \frac{d}{dx}$.

- Write down the general solution u_H of the 5-th order, linear homogeneous ODE p(D)u = 0.
- ullet Write down the form of a particular solution u_P to the inhomogeneous ODE

$$p(D)u = e^{-2x} + x^2 + \cos(x).$$

The zeros of p(z) are 1, -2 (multiplicity 2), i and -i. So the homogeneous ODE has general solution

$$u_H(x) = c_1 e^x + c_2 e^{-2x} + c_3 x e^{-2x} + c_4 \cos(x) + c_5 \sin(x).$$

General solutions

Now consider the inhomogeneous ODE. The solution form will need to contain a x^2e^{-2x} term (since all lower powers are in the homogeneous solution), as well a second order polynomial, as well as $x \cos(x)$ and $x \sin(x)$ terms (since $\cos(x)$ and $\sin(x)$ are in the homogeneous solution).

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$$u_P(x) = a_1 x^2 e^{-2x} + (a_2 x^2 + a_3 x + a_4) + x(a_5 \cos(x) + a_6 \sin(x)).$$

Theorem 7: Reduction of order

If we know a solution $u_1(x) \neq 0$ to the ODE

$$u'' + p(x)u' + q(x)u = 0$$

then we can find a second solution

$$u_2(x) = u_1(x) \int \frac{1}{u_1(x)^2 \exp\left(\int p(x) dx\right)} dx.$$

Example 4

Find the general solution to

$$x^2y'' + 2xy' - 2y = 0$$

given that $y_1(x) = x$ is a solution.

We need to rewrite the ODE in the form from Theorem 3:

$$y'' + \frac{2}{x}y' - \frac{2}{x^2}y = 0.$$

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Then $\exp\left(\int p(x)dx\right) = \exp\left(\int \frac{2}{x}dx\right) = \exp\left(2\ln(x)\right) = x^2$.

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Then $\exp\left(\int p(x)dx\right) = \exp\left(\int \frac{2}{x}dx\right) = \exp\left(2\ln(x)\right) = x^2$. Substituting into the reduction of order formula,

$$y_2(x) = x \int \frac{1}{x^4} dx = -\frac{1}{3x^2}.$$

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$$Lu = f(x)$$
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We find a differential operator A(D) which **annihilates** f(x), meaning A(D)f(x) = 0. For example D^3 annihilates x^2 .

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$$Lu = f(x)$$
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We find a differential operator A(D) which **annihilates** f(x), meaning A(D)f(x) = 0. For example D^3 annihilates x^2 . We apply this differential operator to both sides of our ODE:

$$A(D)Lu = 0.$$

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$$A(D)Lu=0.$$

A solution to this homogeneous differential equation will be a particular solution to the original differential equation.

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$$y''-2y'+y=e^x+\sin(x).$$

The LHS differential operator is $L(D) = (D-1)^2$. The function e^x is annihilated by (D-1). The function $\sin(x)$ is annihilated by (D^2+1) .

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$$A(D)L(D)y(x) = (D-1)^3(D^2+1)y(x) = 0.$$

Solutions to the characteristic equation here is 1 (multiplicity 3), i and -i.

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$$A(D)L(D)y(x) = (D-1)^3(D^2+1)y(x) = 0.$$

Solutions to the characteristic equation here is 1 (multiplicity 3), i and -i. The particular solution is of the form

$$y_P(x) = c_1 e^x + c_2 x e^x + c_3 x^2 e^x + c_4 \sin(x) + c_5 \cos(x).$$

The first two terms in this particular solution are contained in the homogeneous solution of the original ODE, so we discard them:

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Substitute into the original ODE, coefficients are $c_3 = \frac{1}{2}$, $c_4 = 0$ and $c_5 = \frac{1}{2}$.

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Substitute into the original ODE, coefficients are $c_3 = \frac{1}{2}$, $c_4 = 0$ and $c_5 = \frac{1}{2}$. So our particular solution is

$$y_P(x) = \frac{1}{2}x^2e^x + \frac{1}{2}\cos(x).$$

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If we have a linear, 2nd order, inhomogeneous ODE with leading coefficient 1,

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Let u_1, u_2 be a basis for the homogeneous solution space, and let W(x) be the Wronskian $W(x; u_1, u_2)$. Then a particular solution to the inhomogeneous equation is

$$u(x) = v_1(x)u_1(x) + v_2(x)u_2(x),$$

where

$$v_1'(x) = \frac{-u_2(x)f(x)}{W(x)}, \text{ and } v_2'(x) = \frac{u_1(x)f(x)}{W(x)}.$$

Example 6: MATH2121 2018 T2 1.i)

Use the variation of parameters method to solve

$$y'' - 2y' + y = e^x \cos(x).$$

Homogeneous solution has characteristic equation $t^2 - 2t + 1 = 0$ so $u_1(x) = e^x$ and $u_2(x) = xe^x$.

Example 6: MATH2121 2018 T2 1.i)

Use the variation of parameters method to solve

$$y''-2y'+y=e^x\cos(x).$$

Homogeneous solution has characteristic equation $t^2 - 2t + 1 = 0$ so $u_1(x) = e^x$ and $u_2(x) = xe^x$. The Wronskian is

$$W(x) = \det \begin{pmatrix} e^x & xe^x \ e^x & (x+1)e^x \end{pmatrix} = e^{2x}.$$

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Then

$$v_1'(x) = \frac{-xe^x e^x \cos(x)}{e^{2x}} = -x \cos(x),$$

and

$$v_2'(x) = \frac{e^x e^x \cos(x)}{e^{2x}} = \cos(x).$$

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Then a particular solution is

$$u(x) = (-x\sin(x) - \cos(x)) e^x + (\sin(x)) x e^x$$

= $-e^x \cos(x)$.

consider a general second-order, linear, homogeneous ODE

$$Lu = a_2(x)u'' + a_1(x)u' + a_0(x)u = 0,$$

or equivalently

$$Lu = u'' + p(x)u' + q(x)u = 0$$

with
$$p(x) = \frac{a_1(x)}{a_2(x)}$$
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with $p(x) = \frac{a_1(x)}{a_2(x)}$ and $q(x) = \frac{a_0(x)}{a_2(x)}$. Assume that a_0, a_1, a_2 are analytic at 0 (ie. they have local convergent power series), and $a_2(0) \neq 0$. Then p and q are analytic at 0,

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or equivalently

$$Lu = u'' + p(x)u' + q(x)u = 0$$

with $p(x) = \frac{a_1(x)}{a_2(x)}$ and $q(x) = \frac{a_0(x)}{a_2(x)}$. Assume that a_0, a_1, a_2 are analytic at 0 (ie. they have local convergent power series), and $a_2(0) \neq 0$. Then p and q are analytic at 0, so we can find power series expansions of both:

$$p(z) = \sum_{k=0}^{\infty} p_k z^k$$
, $q(z) = \sum_{k=0}^{\infty} q_k z^k$ for $|z| < \rho$,

where $\rho > 0$.

Theorem 8

If coefficients p(z) and q(z) are analytic for $|z| < \rho$, then the formal power series for the solution u(z) constructed in the previous slide, is also analytic for $|z| < \rho$.

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This means that if we find where p(z) and q(z) are both analytic, the power series solution u(z) is also analytic in this region.

Example 7: MATH2121 2018 T2 1.iii)

We aim to construct a series solution to the ODE about the ordinary point $x_0 = 0$:

$$(1-x^2)y''-2xy'+20y=0, \quad y(0)=1, y'(0)=0,$$

of the form

$$y(x) = \sum_{n=0}^{\infty} A_n x^n.$$

- Give the recurrence relation for the coefficients A_n .
- Explain from the recurrence relation that one of the series will terminate yielding a polynomial solution, and the other does not.
- Write down the polynomial solution.

Note that

$$y'(x) = \sum_{n=1}^{\infty} nA_n x^{n-1}, \quad y''(x) = \sum_{n=2}^{\infty} n(n-1)A_n x^{n-2}.$$

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Then we substitute into the ODE:

$$Ly = y'' + (-x^2y'' - 2xy' + 20y)$$

$$= \sum_{n=2}^{\infty} n(n-1)A_n x^{n-2} + \sum_{n=2}^{\infty} -n(n-1)A_n x^n$$

$$- \sum_{n=1}^{\infty} 2nA_n x^n + \sum_{n=0}^{\infty} 20A_n x^n$$

$$= \sum_{n=2}^{\infty} n(n-1)A_n x^{n-2} + \sum_{n=0}^{\infty} -n(n-1)A_n x^n$$

$$- \sum_{n=0}^{\infty} 2nA_n x^n + \sum_{n=0}^{\infty} 20A_n x^n.$$

We combine the last three sums as follows.

$$Ly = \sum_{n=2}^{\infty} n(n-1)A_n x^{n-2} + \sum_{n=0}^{\infty} (-n^2 - n + 20)A_n x^n.$$

In the first sum, we change n to n + 2.

$$Ly = \sum_{n=0}^{\infty} (n+1)(n+2)A_{n+2}x^n + \sum_{n=0}^{\infty} (-n^2 - n + 20)A_nx^n.$$

We can finally combine both sums.

$$Ly = \sum_{n=0}^{\infty} \left[(n+1)(n+2)A_{n+2} - (n+5)(n-4)A_n \right] x^n.$$

Since Ly = 0, we need $(n+1)(n+2)A_{n+2} - (n+5)(n-4)A_n = 0$. Rearranging,

$$A_{n+2} = \frac{(n+5)(n-4)}{(n+1)(n+2)}A_n.$$

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We have $A_0=1$ and $A_1=0$, so all odd terms are zero and the even terms terminate Since $A_6=0$. Hence the polynomial solution is

$$y(x) = 1 - 10x^2 + \frac{35}{3}x^4.$$

Example 8: MATH2221 2015 T2 1.iii)

Consider the ODE

$$(1+z^2)u'' - zu' - 3u = 0.$$

• Find the recurrence relation satisfied by the coefficients A_k in any power series solution:

$$u=\sum_{k=0}^{\infty}A_kz^k.$$

- Show that $A_5 = A_7 = A_9 = ... = 0$.
- Hence find the solution for which u(0) = 0, u'(0) = 6.

Note that

$$u'(z) = \sum_{k=1}^{\infty} k A_k z^{k-1}, \quad u''(z) = \sum_{k=2}^{\infty} k(k-1) A_k z^{k-2}.$$

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Substitute into the ODE:

$$Lu = u'' + (z^{2}u'' - zu' - 3u)$$

$$= \sum_{k=2}^{\infty} k(k-1)A_{k}z^{k-2} + \sum_{k=2}^{\infty} k(k-1)A_{k}z^{k}$$

$$- \sum_{k=1}^{\infty} kA_{k}z^{k} - \sum_{k=0}^{\infty} A_{k}z^{k}$$

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$$- \sum_{k=2}^{\infty} kA_{k}z^{k} - \sum_{k=0}^{\infty} A_{k}z^{k}.$$

Combine the last three sums.

$$Lu = \sum_{k=2}^{\infty} k(k-1)A_k z^{k-2} + \sum_{k=0}^{\infty} (k^2 - 2k - 3)A_k z^k.$$

In the first sum, change k to k+2.

$$Lu = \sum_{k=0}^{\infty} (k+1)(k+2)A_{k+2}z^k + \sum_{k=0}^{\infty} (k^2 - 2k - 3)A_kz^k.$$

Finally, we can combine the sums.

$$Lu = \sum_{k=0}^{\infty} (k+1) \left[(k+2)A_{k+2} + (k-3)A_k \right] z^k.$$

Since
$$Lu = 0$$
, we need $(k+2)A_{k+2} + (k-3)A_k = 0$. Rearranging,

$$A_{k+2} = -\frac{k-3}{k+2}A_k.$$

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$$A_{k+2} = -\frac{k-3}{k+2}A_k.$$

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The even terms start at $A_0=0$ so all even terms past this are zero. Hence

$$u(z)=6z+4z^3.$$

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$$Lu = ax^2u'' + bxu' + cu = f(x),$$

where a, b, c are constants with $a \neq 0$. This is singular at x = 0. Applying this differential operator L to x^r ,

$$Lx^{r} = [ar(r-1) + br + c]x^{r},$$

we can see that x^r is a solution to the homogeneous equation Lu = 0 iff

$$ar(r-1) + br + c = 0.$$

Lemma 3

Suppose there are distinct solutions r_1 , r_2 to the equation ar(r-1)+br+c=0. That is, $r_1 \neq r_2$. Then the general solution of the homogeneous equation Lu=0 is

$$u(x) = C_1 x^{r_1} + C_2 x^{r_2}, \quad x > 0.$$

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Lemma 4

Suppose there is one solution r_1 to ar(r-1)+br+c=0. Then the general solution to the homogeneous equation Lu=0 is

$$C_1 x^{r_1} + C_2 x^{r_1} \log(x), \quad x > 0.$$

For a particular solution to the inhomogeneous Cauchy-Euler equation

$$ax^2u'' + bxu' + cu = x^r$$
,

we can use the particular solution guess $u(x) = \alpha x^r$.

Example 9: MATH2121 2016 T2 2.i)

Find the general solution of the Cauchy-Euler ODE

$$2x^2y'' + 7xy' + 3y = 13x^{1/4}, \quad x > 0.$$

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$$y_H(x) = C_1 x^{-1} + C_2 x^{-3/2}$$
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A particular solution can be found by applying variation of parameters, giving us

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So
$$y(x) = C_1 x^{-1} + C_2 x^{-3/2} - 8x^{1/4}$$
.

Frobenious normal form

A frequent form of ODE that appears in many applications can be written in **Frobenious normal form**:

$$z^2u'' + zP(z)u' + Q(z)u = 0,$$

where P(z) and Q(z) are analytic at 0. Let $P_0 = P(0)$ and $Q_0 = Q(0)$, and define a series F as

$$F(z;r) = z^r \sum_{k=0}^{\infty} A_k(r) z^k.$$

Consider the equation $r(r-1) + P_0r + Q_0 = 0$ with solutions r_1 and r_2 .

Lemma 5

If $r_1 \neq r_2$, then $f(z; r_1)$ is a solution to the Frobenious normal form ODE. If $r_1 - r_2$ is **not** a whole number, then a second linearly independent solution is $F(z; r_2)$.

The **Bessel equation with parameter** ν is:

$$z^2u'' + zu' + (z^2 - \nu^2)u = 0.$$

This ODE is in Frobenious normal form, with indicial polynomial:

$$I(r) = (r + \nu)(r - \nu),$$

and we seek a series solution:

$$u(z) = \sum_{k=0}^{\infty} A_k z^{k+r}.$$

We assume $Re(\nu) \ge 0$, so $r_1 = \nu$ and $r_2 = -\nu$.

With the normalisation:

$$A_0 = \frac{1}{2^{\nu} \Gamma(1+\nu)}$$

the series solution is called the **Bessel function of order** ν and is denoted:

$$J_{
u}(z) = rac{(z/2)^{
u}}{\Gamma(1+
u)} \left[1 - rac{(z/2)^{
u}}{1+
u} + rac{(z/2)^4}{2!(1+
u)(2+
u)} - ...
ight].$$

And from the functional equation $\Gamma(1+z)=z\Gamma(z)$:

$$J_{\nu}(z) = \sum_{k=0}^{\infty} \frac{(-1)k(z/2)^{2k+\nu}}{k!\Gamma(k+1+\nu)}$$

If ν is not an integer, then a second linearly independent, solution is:

$$J_{-\nu}(z) = \sum_{k=0}^{\infty} \frac{(-1)^k (z/2)^{2k-\nu}}{k! \Gamma(k+1-\nu)}.$$

For an integer $\nu = n \in \mathbb{Z}$, since $\Gamma(n+1) = n!$, we have:

$$J_n(z) = \sum_{k=0}^{\infty} \frac{(-1)^k (z/2)^{2k+n}}{k!(k+n)!}.$$

Also, since $\frac{1}{\Gamma(z)} = 0$ for z = 0, -1, -2, ..., we find that J_n and J_{-n} are linearly independent; in fact:

$$J_{-n}(z) = (-1)^n J_n(z).$$

Example 10: MATH2221 2015 T2 2.ii)

① Use term-by-term differentiation to prove that for $\nu \in \mathbb{R}$ and x > 0:

$$\frac{d}{dx}(x^{\nu}J_{\nu}(x))=x^{\nu}J_{\nu-1}(x).$$

4 Hence evaluate the definite integral:

$$I = \int_{0}^{1} x^{\frac{7}{2}} J_{\frac{1}{2}}(x) dx.$$

Note that

$$x^{\nu}J_{\nu}(x) = \sum_{k=0}^{\infty} x^{\nu} \frac{(-1)^k (x/2)^{2k+\nu}}{k!\Gamma(k+1+\nu)}.$$

Moving the x^{ν} into the fraction,

$$x^{\nu}J_{\nu}(x) = \sum_{k=0}^{\infty} \frac{(-1)^{k}(x)^{2k+2\nu}}{2^{2k+\nu}k!\Gamma(k+1+\nu)}.$$

Take one term in this sum and differentiate with respect to x:

$$\frac{d}{dx}\left(\frac{(-1)^k(x)^{2k+2\nu}}{2^{2k+\nu}k!\Gamma(k+1+\nu)}\right) = \frac{(-1)^k(2k+2\nu)x^{2k+2\nu-1}}{2^{2k+\nu}k!\Gamma(k+1+\nu)}.$$

Since $\Gamma(k+1+\nu)=(k+\nu)\Gamma(k+1+\nu)$, then

$$\frac{d}{dx}\left(\frac{(-1)^k(x)^{2k+2\nu}}{2^{2k+\nu}k!\Gamma(k+1+\nu)}\right) = \frac{(-1)^k(2k+2\nu)x^{2k+2\nu-1}}{2^{2k+\nu}k!(k+\nu)\Gamma(k+\nu)}.$$

However there is a factor of $2(k + \nu)$ in both the numerator and denominator. So

$$\frac{d}{dx}\left(\frac{(-1)^k(x)^{2k+2\nu}}{2^{2k+\nu}k!\Gamma(k+1+\nu)}\right) = \frac{(-1)^k(x)^{2k+2\nu-1}}{2^{2k+\nu-1}k!\Gamma(k+\nu)}.$$

Finally, factor out a factor x^{ν} :

$$\frac{d}{dx}\left(\frac{(-1)^k(x)^{2k+2\nu}}{2^{2k+\nu}k!\Gamma(k+1+\nu)}\right) = x^{\nu}\frac{(-1)^k(x/2)^{2k+\nu-1}}{k!\Gamma(k+\nu)}.$$

Since the derivative is linear,

$$\frac{d}{dx}(x^{\nu}J_{\nu}(x)) = \sum_{k=0}^{\infty} \frac{d}{dx} \left(\frac{(-1)^{k}(x)^{2k+2\nu}}{2^{2k+\nu}k!\Gamma(k+1+\nu)} \right).$$

Substituting in the derivative we found,

$$\frac{d}{dx}(x^{\nu}J_{\nu}(x)) = x^{\nu}\sum_{k=0}^{\infty} \frac{(-1)^{k}(x/2)^{2k+\nu-1}}{k!\Gamma(k+\nu)}.$$

The RHS here is just $x^{\nu}J_{\nu-1}(x)$, so we are done.

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The RHS here is just $x^{\nu}J_{\nu-1}(x)$, so we are done. To find the integral $\int_{1}^{1} x^{\frac{7}{2}} J_{\frac{1}{2}}(x) dx$, we separate and use integration by parts:

$$\int_{0}^{1} x^{2} \cdot x^{\frac{3}{2}} J_{\frac{1}{2}}(x) dx$$

where $u' = x^{\frac{3}{2}} J_{\frac{1}{2}}(x)$ and $v = x^2$.

Applying integration by parts:

$$I = [x^{2}.x^{\frac{3}{2}}J_{\frac{3}{2}}(x)]_{0}^{1} - \int_{0}^{1} 2x \cdot x^{\frac{3}{2}}J_{\frac{3}{2}}(x)dx$$
$$= J_{\frac{3}{2}}(1) - 2\int_{0}^{1} x^{\frac{5}{2}}J_{\frac{3}{2}}(x)dx.$$

However by our previous result, $\frac{d}{dx}(x^{\nu}J_{\nu}(x))=x^{\nu}J_{\nu-1}(x)$. So the antiderivative of $x^{5/2}J_{3/2}$ is $x^{5/2}J_{5/2}$, hence

$$= J_{\frac{3}{2}}(1) - 2[x^{\frac{5}{2}}J_{\frac{5}{2}}(x)]_{0}^{1}$$

= $J_{\frac{3}{2}}(1) - 2J_{\frac{5}{2}}(1)$.

Legendre equation (2221 only)

The **Legendre equation** with paramater ν is:

$$(1-z^2)u'' - 2zu' + \nu(\nu+1)u = 0.$$

This ODE is not singular at z=0, so the solution has an ordinary Taylor series expansion:

$$u=\sum_{k=0}^{\infty}A_kz^k.$$

The A_k must satisfy:

$$(k+1)(k+2)A_{k+2}-[k(k+1)-\nu(\nu+1)]A_k=0.$$

The recurrence relation is:

$$A_{k+2} = \frac{(k-\nu)(k+\nu+1)}{(k+1)(k+2)} A_k$$
 for $k \ge 0$.

Legendre equation (2221 only)

We have:

$$u(z) = A_0 u_0(z) + A_1 u_1(z)$$

where:

$$u_0(z) = 1 - \frac{\nu(\nu+1)}{2!}z^2 + \frac{(\nu-2)\nu(\nu+1)(\nu+3)}{4!}z^-...$$

and:

$$u_1(z) = z - \frac{(\nu - 1)(\nu - 2)}{3!}z^3 + \frac{(\nu - 3)(\nu - 1)(\nu + 2)(\nu + 4)}{5!}z^5 - \dots$$

Suppose now that $\nu = n$ is a non-negative integer. If n is even, then the series for $u_0(z)$ terminates, whereas if n is odd, then the series for $u_1(z)$ terminates. The terminating solution is then called the **Legendre polynomial** of degree n and is denoted by $P_n(z)$ with the normalisation:

$$P_n(1) = 1.$$



Dynamical Systems

Dynamical Systems

State variables are natural variables which depending on a single independent variable. A **dynamical system** is a natural process described by these state variables. The state of a system at a given time is described by the values of the state variables at that instant.

Note that any *n*th order ODE can be writen as a system of **first order** ODEs (not vice versa):

$$\frac{d^n y}{dt^n} = g\left(y, \frac{dy}{dt}, ..., \frac{d^{n-1}y}{dt^{n-1}}\right)$$
$$\frac{dx}{dt} = f(x_1, x_2, ..., x_n)$$

Non-autonomous ODEs

Definition 8

A system of ODEs of the form:

$$\frac{d\mathbf{x}}{dt} = \mathbf{F}(\mathbf{x})$$

is said to be autonomous.

Definition 9

In a non-autonomous system, **F** may depend explicitly on t:

$$\frac{d\mathbf{x}}{dt} = \mathbf{F}(\mathbf{x}, t).$$

Definition 10

The number $L \in \mathbb{R}$ is a Lipschitz constant for a function $f: [a, b] \to \mathbb{R}$ if

$$|f(x)-f(y)| \leq L|x-y| \quad \forall x,y \in [a,b].$$

We say that the function f is Lipschitz if a Lipschitz constant for f exists.

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Theorem 9

If f is Lipschitz, then f is uniformly continuous.

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Theorem 9

If f is Lipschitz, then f is uniformly continuous.

Lemma 6

If $f: I \to \mathbb{R}$ is differentiable and f' is continuous on I, then f is Lipschitz.

Lipschitz Vector Field (2221 only)

We extend the definition of Lipschitz to vector fields.

Definition 11

A vector field $\mathbf{F}:S\subseteq\mathbb{R}^m\to\mathbb{R}^n$ is Lipschitz on S if

$$||\mathbf{F}(\mathbf{x}) - \mathbf{F}(\mathbf{y})|| \le L ||\mathbf{x} - \mathbf{y}|| \quad \forall \, \mathbf{x}, \mathbf{y} \in \mathcal{S}.$$

Here,

$$||\mathbf{x}|| = \left(\sum_{j=1}^{m} x_j^2\right)^{1/2}$$

denotes the Euclidean norm of the vector $\mathbf{x} \in \mathbb{R}^m$.

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Here,

$$||\mathbf{x}|| = \left(\sum_{j=1}^{m} x_j^2\right)^{1/2}$$

denotes the Euclidean norm of the vector $\mathbf{x} \in \mathbb{R}^m$.

We say that $\mathbf{F}(\mathbf{x}, t)$ is Lipschitz in \mathbf{x} if, for all t:

$$||\mathbf{F}(\mathbf{x},t) - \mathbf{F}(\mathbf{y},t)|| \le L ||\mathbf{x} - \mathbf{y}||, \quad \forall \mathbf{x}, \mathbf{y} \in \mathbb{R}^m.$$



Existence and Uniqueness Theorem (2221 only)

We want to find solutions to a non-autonomous system. The following theorem guarantees a unique solution for a non-autonomous system, under certain conditions.

Theorem 10

The initial value problem defined by:

$$\frac{d\mathbf{x}}{dt} = \mathbf{F}(\mathbf{x}, t), \quad \mathbf{x}(t_0) = \mathbf{x}_0$$

has a unique solution $\mathbf{x}(t)$ over a time interval $|t - t_0| < \alpha$ if $\mathbf{F}(\mathbf{x}, t)$ is continuous and Lipschitz.

In fact, $\mathbf{x}(t)$ is continuous and differentiable.

Notes

- The existence of solutions follow from continuity in x and t.
- The uniqueness of solutions follow from the Lipschitz condition in x.
- The theorem is a local existence theorem. It provides for the existence of solutions over a finite time interval.

Example 11

$$f(x) = 2\sqrt{x}, \quad x \in \mathbb{R}$$

Where is f Lipschitz?

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Example 11

$$f(x) = 2\sqrt{x}, \quad x \in \mathbb{R}$$

Where is *f* Lipschitz?

Since $f'(x) = \frac{1}{\sqrt{x}}$ is continuous on x > 0, then f is Lipschitz for x > 0 (Lemma 6). At 0, f is not Lipschitz since if it was,

$$|f(x) - f(y)| = 2|\sqrt{x} - \sqrt{y}|$$

$$\leq L|x - y|$$

$$\frac{2}{L} \leq |\sqrt{x} + \sqrt{y}|$$

which cannot hold true for x and y both 0.

Linear systems of ODEs

Definition 12

We say that the $n \times n$, first-order system of ODEs:

$$\frac{d\mathbf{x}}{dt} = \mathbf{F}(\mathbf{x}, t)$$

is linear if the RHS has the form:

$$\mathbf{F}(\mathbf{x},t) = A(t)\mathbf{x} + \mathbf{b}(t)$$

for some $n \times n$ matrix-valued function $A(t) = [a_{i,j}(t)]$ and a vector-valued function $\mathbf{b}(t) = [b_i(t)]$.

The linear first-order system is autonomous when A and \mathbf{b} are constant.

Global Existence and Uniqueness

We have a stronger existence result in the linear case:

Theorem 11

If A(t) and $\mathbf{b}(t)$ are continuous for $0 \le t \le T$, then the linear initial-value problem

$$\frac{d\mathbf{x}}{dt} = A(t)\mathbf{x} + \mathbf{b}(t) \quad \text{with } \mathbf{x}(0) = \mathbf{x}_0,$$

has a unique solution $\mathbf{x}(t)$ for $0 \le t \le T$.

A special case

It is much easier to work with the special case when A(t) = A is a constant $n \times n$ matrix and $\mathbf{b}(t) = \mathbf{0}$:

$$\frac{d\mathbf{x}}{dt} = A\mathbf{x}$$

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The general solution to this system is

$$\mathbf{x}(t) = \sum_{i=1}^{n} c_i e^{\lambda_i t} \mathbf{v}_i,$$

where λ_i is the *i*-th eigenvalue with corresponding eigenvector \mathbf{v}_i and $c_1, ..., c_n$ are constants.



Initial-valued system

Recall that for an $n \times n$ complex matrix A, we define the exponential of a matrix as

$$e^A = \sum_{k=0}^{\infty} \frac{A^k}{k!} = I + A + \frac{1}{2!}A^2 + \frac{1}{3!}A^3 + \cdots$$

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Consider the initial-valued system

$$\frac{d\mathbf{x}}{dt} = A\mathbf{x}, \quad \mathbf{x}(0) = \mathbf{x}_0.$$

Then we can write the solution as a matrix exponential,

$$\mathbf{x}(t)=e^{tA}\mathbf{x}_0.$$



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The issue is calculating e^{tA} , which requires finding A^k for $k \in \mathbb{N}$. To do this efficiently, we look to diagonalisation.



Diagonalising a matrix

Definition 13

An $n \times n$ complex matrix A is diagonalisable if there exists a non-singular matrix $n \times n$ matrix M such that $M^{-1}AM$ is diagonal.



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Theorem 12

An $n \times n$ complex matrix A is diagonalisable if and only if there exists a basis $\{\mathbf{v}_1,...,\mathbf{v}_n\}$ for \mathbb{C}^n , where $\mathbf{v}_1,...,\mathbf{v}_n$ are eigenvectors of A. In fact the columns of M are the eigenvectors of A, $M = (\mathbf{v}_1|\cdots|\mathbf{v}_n)$, and $M^{-1}AM = \Lambda$ where:

$$\Lambda = \begin{pmatrix} \lambda_1 & & \\ & \ddots & \\ & & \lambda_n \end{pmatrix}$$

for eigenvalues λ_i corresponding to eigenvector \mathbf{v}_i .



Matrix Powers

In general since $M^{-1}AM = \Lambda$, we can efficiently calculate A^k :

$$A^k = \overbrace{M \wedge M^{-1} \cdot M \wedge M^{-1} \cdots M \wedge M^{-1}}^{k \text{ times}} = M \wedge^k M^{-1}.$$



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This is better because

$$\Lambda^k = \begin{pmatrix} \lambda_1^k & & \\ & \ddots & \\ & & \lambda^n \end{pmatrix}.$$

In fact, due to the taylor series of the exponential function, we can simplify the solution to the initial-valued system further.

Exponential of a diagonalisable matrix

Theorem 13

If $A = M\Lambda M^{-1}$ is diagonalisable, then:

Exponential of a diagonalisable matrix

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If $A = M\Lambda M^{-1}$ is diagonalisable, then:

$$e^A = Me^{\Lambda}M^{-1}$$
 and $e^{\Lambda} = \begin{pmatrix} e^{\Lambda_1} & & \\ & \ddots & \\ & & e^{\lambda_n} \end{pmatrix}$.

We can now find the exponential of tA:



Example 12: MATH2121 2016 T2 2.iii)

For an $n \times n$ matrix A.

- State the definition of e^A .
- 2 Show that if $A\mathbf{v} = \lambda \mathbf{v}$, then $e^A \mathbf{v} = e^\lambda \mathbf{v}$.

The definition of the matrix exponential is

$$e^A = \sum_{k=0}^{\infty} \frac{A^k}{k!} = I + A + \frac{A^2}{2!} + \frac{A^3}{3!} + \cdots$$

Example 12: MATH2121 2016 T2 2.iii)

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- **1** State the definition of e^A .
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The definition of the matrix exponential is

$$e^A = \sum_{k=0}^{\infty} \frac{A^k}{k!} = I + A + \frac{A^2}{2!} + \frac{A^3}{3!} + \cdots$$

If $A\mathbf{v} = \lambda \mathbf{v}$ then

$$e^{A}\mathbf{v} = \left(\sum_{k=0}^{\infty} \frac{1}{k!} A^{k}\right) \mathbf{v}$$
$$= \sum_{k=0}^{\infty} \frac{1}{k!} \left(A^{k} \mathbf{v}\right).$$

But we can find $A^k \mathbf{v}$:

$$A^k \mathbf{v} = A^{k-1} \lambda \mathbf{v} = \cdots = \lambda^k \mathbf{v}.$$

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This is the exponential function Taylor series at the point λ , so

$$e^A \mathbf{v} = e^\lambda \mathbf{v}$$
.

Equilibrium points

Definition 14

We say that $\mathbf{a} \in \mathbb{R}^n$ is an equilibrium point for the dynamical system $\frac{d\mathbf{x}}{dt} = \mathbf{F}(\mathbf{x})$ if

$$F(a) = 0.$$

Equilibrium points

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We say that $\mathbf{a} \in \mathbb{R}^n$ is an equilibrium point for the dynamical system $\frac{d\mathbf{x}}{dt} = \mathbf{F}(\mathbf{x})$ if

$$F(a) = 0.$$

Suppose **a** is an equilibrium point for the system $\frac{d\mathbf{x}}{dt} = \mathbf{F}(\mathbf{x})$. Consider the initial-valued system

$$\frac{d\mathbf{x}}{dt} = \mathbf{F}(\mathbf{x}), \quad \mathbf{x}(0) = \mathbf{a}.$$

Then the solution is the constant function $\mathbf{x}(t) = \mathbf{a}$.

Stable Equilibrium

Definition 15

An equilibrium point ${\bf a}$ is stable if for every $\epsilon>0$, there exists $\delta>0$ such that whenever $||{\bf x}_0-{\bf a}||<\delta$, the solution of

$$\frac{d\mathbf{x}}{dt} = \mathbf{F}(\mathbf{x}), \quad \mathbf{x}(0) = \mathbf{x}_0$$

satisfies

$$||\mathbf{x}(t) - \mathbf{a}|| < \epsilon \quad \forall \ t > 0.$$

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satisfies

$$||\mathbf{x}(t) - \mathbf{a}|| < \epsilon \quad \forall \ t > 0.$$

Intuitively: if a solution starts close enough to the stable equilibrium point, then they will remain close to the stable equilibrium point.

Asymptotic Stability

This is a stronger form of stability, on a particular subset of \mathbb{R}^n .

Definition 16

Let N be an open subset of \mathbb{R}^n that contains an equilibrium point \mathbf{a} . We say that \mathbf{a} is asymptotically stable in N if \mathbf{a} is stable, and whenever $\mathbf{x}_0 \in N$ the solution of

$$\frac{d\mathbf{x}}{dt} = \mathbf{F}(\mathbf{x}) \quad \mathbf{x}(0) = \mathbf{x}_0$$

satisfies

$$\mathbf{x}(t) \to \mathbf{a}$$
 as $t \to \infty$.

We call N a domain of attraction for \mathbf{a} .

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 as $t \to \infty$.

We call N a domain of attraction for \mathbf{a} .

Intuitively: not only do the solutions stay close to the stable equilibrium point, but they also approach the equilibrium point as t goes to infinity.

Linear constant case

Consider the linear system

$$\frac{d\mathbf{x}}{dt} = A\mathbf{x} + \mathbf{b}, \quad \mathbf{x}(0) = \mathbf{x}_0,$$

where $det(A) \neq 0$.

Linear constant case

Consider the linear system

$$\frac{d\mathbf{x}}{dt} = A\mathbf{x} + \mathbf{b}, \quad \mathbf{x}(0) = \mathbf{x}_0,$$

where $det(A) \neq 0$. Then the unique equilibrium point of this system is

$$\mathbf{a} = -A^{-1}\mathbf{b},$$

and the solution to the system is

$$\mathbf{x}(t) = \mathbf{a} + e^{tA} (\mathbf{x}_0 - \mathbf{a}).$$

Linear constant case

Theorem 14

Consider the previous linear constant coefficient system. Let $\lambda_1, ..., \lambda_n$ be the eigenvalues of A. The equilibrium point $\mathbf{a} = -A^{-1}b$ is:

- **1** Stable if and only if $Re(\lambda_j) \leq 0$ for all j.
- **2** asymptotically stable if and only if $Re(\lambda_j) < 0$ for all j.

In the second case, the domain of attraction is the whole of \mathbb{R}^n .



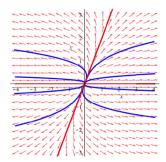
Classification of 2D Linear Systems

| Type | Eigenvalues | Eigenvectors | X(t) | Classification |
|---|---|--|--|----------------------------|
| 1: $\mathbf{B} = \begin{pmatrix} \lambda_1 & 0 \\ 0 & \lambda_2 \end{pmatrix}$ | $\lambda_1 \neq \lambda_2 \in \mathbb{R}$ | $v^{(1)}$, $v^{(2)}$ | $\left[\boldsymbol{v^{(1)}} e^{\lambda_1 t} \ \boldsymbol{v^{(2)}} e^{\lambda_2 t} \right]$ | Improper Node |
| | $\lambda_1 < 0 < \lambda_2$ | $v^{(1)}, v^{(2)}$ | $\left[v^{(1)}e^{\lambda_1 t}\ v^{(2)}e^{\lambda_2 t}\right]$ | Saddle Point |
| 2: $B = \begin{pmatrix} \lambda & \gamma \\ 0 & \lambda \end{pmatrix}$ $\lambda, \gamma \in \mathbb{R}$ | $\lambda_1 = \lambda_2 = \lambda \in \mathbb{R}$ (multiplicity 2) | $v^{(1)}, v^{(2)}$ $(v^{(2)}$ generalised eigenvector) | $\left[v^{(1)}e^{\lambda t} \ (v^{(2)}+tv^{(1)})e^{\lambda t}\right]$ | Deficient Node |
| | $\lambda_1 = \lambda_2 = \lambda$ (2D eigenspace) | $v^{(1)}, v^{(2)}$ any basis of \mathbb{R}^2 | $\left[\boldsymbol{v^{(1)}}e^{\lambda_1t}\ \boldsymbol{v^{(2)}}e^{\lambda_2t}\right]$ | Star (or proper) Node |
| 3: $\mathbf{B} = \begin{pmatrix} \alpha & -\omega \\ \omega & \alpha \end{pmatrix}$ $\alpha, \omega \in \mathbb{R}$ | $\lambda_1 = i\beta = \overline{\lambda_2}$ $(\beta \neq 0) \in \mathbb{R}$ | $v^{(2)} = \overline{v^{(1)}}$ | $\boxed{ \left[\mathbb{R} \mathrm{e} \big(\boldsymbol{v}^{(1)} e^{i\beta t} \big) \ \mathbb{I} \mathrm{m} \big(\boldsymbol{v}^{(2)} e^{i\beta t} \big) \right] }$ | Centre (or vortex) |
| | $\lambda_1 = \alpha + i\beta = \overline{\lambda_2}$ $(\alpha \neq 0, \beta \neq 0) \in \mathbb{R}$ | $v^{(2)} = \overline{v^{(1)}}$ | $\left[\mathbb{R}\mathrm{e}\big(\boldsymbol{v}^{(1)}e^{(\alpha+i\beta)t}\big)\ \mathbb{I}\mathrm{m}\big(\boldsymbol{v}^{(2)}e^{(\alpha+i\beta)t}\big)\right]$ | Spiral point (or focus) |



Improper Node

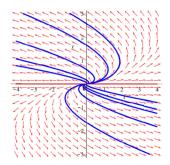
Here $\lambda_1 \neq \lambda_2 \in \mathbb{R}$. Which eigenline orbits tend to depends on the eigenvalues.





Deficient Node

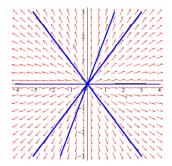
Here $\lambda_1 = \lambda_2 \in \mathbb{R}$, and the eigenspace is deficient (out of the scope of the course).





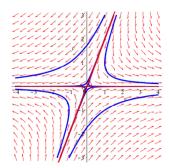
Star Node

Here $\lambda_1 = \lambda_2 \neq 0$ and all nonzero vectors are eigenvectors. Therefore, all orbits are either being attracted or repelled by the equilibrium points.



Saddle Point

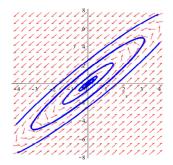
Here, $\lambda_1 < 0 < \lambda_2$. This means λ_1 is attracting, whilst λ_2 is repelling.





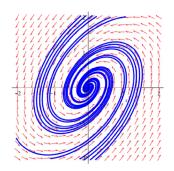
Centre

Here $\lambda_1=i\beta=\lambda_2$, where $\beta\in\mathbb{R}$. If eigenvalues are purely imaginary, orbits are given by ellipses around the eigen-plane as $e^{it}=\cos(t)+i\sin(t)$.



Spiral

Here $\lambda_1=\alpha+i\beta=\bar{\lambda_2}$. If eigenvalues are in conjugate pairs, real parts of the orbit will be defined by $e^{-\lambda t}=e^{-\operatorname{Re}(()\lambda)t}(c_1\cos(t)+c_2\sin(t))$, which will spiral inward or outward, clockwise or anticlockwise depending on values of $\operatorname{Re}(()\lambda)$ and $\operatorname{Im}(()\lambda)$.



Example 13: MATH2121 2018 T2 2.iii)

Solve for x(t) and y(t) and determine the type and stability of the equilibrium point of the following system of differential equations:

$$\frac{dx}{dt} = x + y;$$
$$\frac{dy}{dt} = 2x.$$

Example 13: MATH2121 2018 T2 2.iii)

Solve for x(t) and y(t) and determine the type and stability of the equilibrium point of the following system of differential equations:

$$\frac{dx}{dt} = x + y;$$
$$\frac{dy}{dt} = 2x.$$

To find the equilibrium point, we solve x + y = 0 and 2x = 0 simultaneously. The only solution is the point (0,0). The system can be written as

$$\frac{d}{dt} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} 1 & 1 \\ 2 & 0 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix}.$$

The eigenvalues of
$$\begin{pmatrix} 1 & 1 \\ 2 & 0 \end{pmatrix}$$
 are $\lambda_1 = -1$ and $\lambda_2 = 2$.

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The eigenvalues of $\begin{pmatrix} 1 & 1 \\ 2 & 0 \end{pmatrix}$ are $\lambda_1 = -1$ and $\lambda_2 = 2$.

Since $\lambda_1 < 0 < \lambda_2$, the equilibrium point (0,0) is a saddle point.

Since $Re(\lambda_2) > 0$ then (0,0) is an unstable equilibrium point.

Hence (0,0) is an unstable saddle point.

Definition 17

A function $G: \mathbb{R}^n \to \mathbb{R}$ is a **first integral** for a system of ODEs

$$\frac{d\mathbf{x}}{dt} = \mathbf{F}(\mathbf{x})$$

if $G(\mathbf{x}(t))$ is constant for every solution $\mathbf{x}(t)$.

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if $G(\mathbf{x}(t))$ is constant for every solution $\mathbf{x}(t)$.

Geometrically: G is a first integral iff

$$\nabla G(x) \perp \mathbf{F}(\mathbf{x})$$
 for all \mathbf{x} .

Example 14

Consider the system of ODEs

$$\frac{dx_1}{dt} = x_1 x_2,$$

$$\frac{dx_2}{dt} = -x_1^2.$$

Prove that $G(\mathbf{x}) = x_1^2 + x_2^2$ is a first integral.

Set the function F(x) as the RHS of the system:

$$\mathbf{F}(\mathbf{x}) = \begin{pmatrix} x_1 x_2 \\ -x_1^2 \end{pmatrix}.$$



We want to find the gradient of G:

$$\nabla G(\mathbf{x}) = \begin{pmatrix} 2x_1 \\ 2x_2 \end{pmatrix}.$$

So

$$\nabla G(\mathbf{x}) \cdot \mathbf{F}(\mathbf{x}) = \begin{pmatrix} 2x_1 \\ 2x_2 \end{pmatrix} \cdot \begin{pmatrix} x_1 x_2 \\ -x_1^2 \end{pmatrix} = 2x_1^2 x_2 - 2x_2 x_1^2 = 0.$$

This means that all solutions to the system of ODEs must be mapped to a constant under G. That is, if (x_1, x_2) is a solution then $x_1^2 + x_2^2 = C$ for some C. In other words, all the solutions to the system of ODEs lie on some circle around the origin. What this tells us is that if $x_1 = r\cos(t)$ then x_2 must be $r\sin(t)$ for whatever value of t > 0.