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Spacecraft System Design

Tafelanschriften

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1 13.10.2014

$$m \cdot \ddot{\vec{r}} = F = -G \cdot \frac{m_1 m_2}{r^2}$$

G ist die *universal gravity constant*: $6.674 \cdot 10^{-11} \frac{Nm^2}{kg^2}$

$$\ddot{\vec{r}} = \frac{\mu}{\|\vec{r}\|^2}$$

μ	gravity constant of the specific gravity attractor.		$398600 \frac{km^3}{s^2}$ $5.97219 \cdot 10^{24} kg$
μ_E	gravity constant of earth:		
m_E	mass of earth:		

Die Lösungen dieser Gleichung sind Kegelschnitte. Außerdem gilt: $\mu_E = G \cdot m_E$

$$r = \frac{p}{1 + \varepsilon \cdot \cos \varphi}$$

r	radius towards focal point	
ε	eccentricity of the orbit	
p	parameters (spatial extensions of orbit)	
φ	true anomaly	

Interrelationships between parameters:

r_a apocenter, r_p pericenter

$$a = \frac{p}{1 - \varepsilon^2} \Rightarrow p = a(1 - \varepsilon^2) = r_p(1 + \varepsilon) = r_a(1 - \varepsilon)$$

$$a = \frac{r_a + r_p}{2}, r_a = a(1 + \varepsilon), r_p = a(1 - \varepsilon)$$

$$\varepsilon = \frac{r_a - r_p}{r_a + r_p}$$

For the specific case of circular orbits:

- gravitational acceleration in distance a to gravitational center: $\frac{\mu}{a^2}$
- centrifugal acceleration: $\omega^2 \cdot r$ with $\omega = \frac{2\pi}{T}$ angular velocity and T the orbital period

$$\left(\frac{2\pi}{T}\right)^2 \cdot a \Rightarrow \frac{\mu}{a^2} = \left(\frac{2\pi}{T}\right)^2 \cdot a \Rightarrow T = 2\pi \sqrt{\frac{a^3}{\mu}}$$

which also holds for general ellipses

$$v_{\text{circle}} = \frac{2\pi a}{T} = \sqrt{\frac{a^3}{\mu}}$$

Potential energy: $E_{\text{pot}} = \frac{-\mu m}{r} + C$

Kinetic energy: $E_{\text{kin}} = \frac{mv^2}{2}$

$$E = E_{\text{kin}} + E_{\text{pot}} = \frac{mv^2}{2} - \frac{\mu m}{r} \stackrel{\substack{\varepsilon \rightarrow 1, \\ v_a \rightarrow 0, \\ r_a \rightarrow 2a}}{=} -\frac{\mu m}{2a} (\text{constant})$$

v_a velocity at apocenter

Binet's equation, Vis-Viva equation:

$$v = \sqrt{\mu \left(\frac{2}{r} - \frac{1}{a} \right)}$$

For circles: $r = a$, $v = \sqrt{\frac{\mu}{a}}$

v first cosmic velocity: needed to bring a satellite on closed orbit (without perturbing forces, just gravity considered). In case of earth: $r = 6378 \text{ km}$, $\mu_E = 398600 \frac{\text{km}^3}{\text{s}^2}$
 $\Rightarrow 7.905 \frac{\text{km}}{\text{s}}$ required as minimum velocity from the rocket launch to insert into an earth orbit.

A rocket launch in equatorial direction and in rotation direction of earth, we receive a velocity component $0.463 \frac{\text{km}}{\text{s}}$ for free earth rotation rate.

For parabola: $a \rightarrow \infty$, $v = \sqrt{\frac{2\mu}{r}}$ v second cosmic velocity (escape velocity): is the minimum velocity applied to leave the gravitational field of the home planet. For earth: $11.179 \frac{\text{km}}{\text{s}}$

Derive position of the satellite in the orbit plane, e.g. we want the true anomaly φ as the function of time $\varphi = f(t)$. There is no explicit solution but only an algorithm using the support variable eccentric anomaly E can be derived.

Keplers equation:

$$E - \varepsilon \cdot \sin(E) = \frac{2\pi}{T}(t - t_{\text{perigee}}) = M(t)$$

$M(t)$ average anomaly at time t , can be solved by numerical methods. From E , you can calculate r and φ as follows:

$$r \sin \varphi = a \sin E \cdot \sqrt{1 - \varepsilon^2}$$

$$r \cos \varphi = a(\cos E - \varepsilon)$$

$$\Rightarrow \tan \frac{\varphi}{2} = \sqrt{\frac{1 + \varepsilon}{1 - \varepsilon}} \cdot \tan \frac{E}{2}$$

$$r = a(1 - \varepsilon \cos E)$$

Satellite Ground Tracks:

Analyse the triangle KFB in spherical coordinates.

$$\begin{array}{l|l} \theta & \text{latitude:} \\ \lambda & \text{longitude:} \end{array} \quad \left| \begin{array}{l} \sin \theta = \sin i \sin(\omega + \varphi) \\ \cos \lambda_1 = \frac{\cos(\omega + \varphi)}{\cos \theta} \end{array} \right.$$

$$\lambda = \lambda_1 + \Omega + \lambda_G$$

$$\lambda_G = \lambda_{G0} - 0.25068448^\circ / \text{min} \cdot t \quad (\text{Greenwich position versus } \gamma)$$

Draw the line from the satellite to earth center, the intersection with surface is specified by θ and λ (*subsatellite point*).

Perturbations:

gravitational potential of earth

$$u(r, \phi, \Lambda)$$

u gravitational at distance r and latitude ϕ and longitude Λ of the position.

$$u(r, \phi, \Lambda) = \frac{\mu}{r} \left(1 + \sum_{n=2}^{\infty} \left[\left(\frac{R_E}{r} \right)^n \cdot J_2 \cdot P_{n0} \cos \phi + \sum_{m=1}^n \left(\frac{R_e}{r} \right)^n \cdot (C_{nm} \cos(m\Lambda) + S_{nm} \sin(m\Lambda)) \cdot P_{nm}(\cos \phi) \right] \right)$$

spherical harmonic expansion with:

r	distance from the earth center	
P_{nm}	Legendre polynomials	
R_E	radius of earth	6378km

J_n, C_{nm}, S_{nm} are constants of body reflecting the mass distribution:

J_n zonal harmonic coefficient (mass distribution independent from longitude)

C_{nm}, S_{nm} earth's tesseral harmonic coefficients (for $n \neq m$) or earth's sectoral harmonic coefficients (for $n = m$)

By far, the largest coefficient is J_n (related to equatorial bulge) by 3 orders of magnitude.

Example: Molniya orbits of 63, 4° have minimum change in position of perigee and apogee.

Additional gravitational fields:

by other bodies in the solar system

$$a_d = \mu_d \sqrt{\vec{R} \cdot \vec{R}}$$

with $\vec{R} = \frac{\vec{r}_{sd}}{\|\vec{r}_d\|}$

$$\frac{a_d}{a_c} = \frac{M_d}{M_c} \left(\frac{\|\vec{r}_s\|}{\|\vec{r}_{sd}\|} \right)^3 \sqrt{1 + 3 \cos^2 \beta} = 2 \frac{M_d}{M_c} \left(\frac{\|\vec{r}_s\|}{\|\vec{r}_d\|} \right)^3$$

a_c | actual acceleration |

atmospheric drag:

Drag: $\vec{F}_0 = -\frac{1}{2} \rho S c_D \|\vec{v}_R\|^2 \frac{\vec{v}_r}{\|\vec{v}_r\|}$

v_r	velocity vector of the spacecraft relative to the atmosphere	
ρ	atmospheric density	
S	reference area of the spacecraft	
c_D	drag coefficient of the spacecraft (determined in wind channel)	default: $c_d \approx 2.5$

Main effect of altitude changes and changes of orbital period.

solar pressure:

1AU	distance earth-sun	$149.6 \cdot 10^6 \text{ km}$
	energy density	$1372 \frac{W}{m^2}$

from impulse conservation law:

$$p = \frac{1372W}{c \cdot m^2} = \frac{1372W \text{ s}}{299793km \text{ m}^2} = 0.46 \cdot 10^{-5} \frac{N}{m^2}$$

solar pressure is increased by reflection

$$p_{\text{eff}} = (1 + r)p$$

p_{eff}		effective pressure	
r		reflectivity factor	

At equatorial orbits the solar pressure affects a change of eccentricity ε .

orbit change manouvers:

Based on impule conservation law.

Propulsion systems:

Hydrazin	$2.2 \frac{km}{s}$	particle velocity
2-component	$3 \frac{km}{s}$	particle velocity
ion thrusters	$30 \frac{km}{s}$	particle velocity

$$\Delta v_s m_s = -\Delta m_s v_p$$

Δm_s		reduction of satellite mass due to ejected propellant	
v_p		velocity of propellant ejection	

$$\Delta v_s = -\frac{\Delta m_s}{m_s} v_p$$

$$\int_{v_0}^v dv = -v_p \int_{m_s}^{m_s - m_p} \frac{dm}{m}$$

$$\Delta v_s = v_p (\log_2 m_s - \log_2 (m_s - m_p))$$

Rocket equation, Ziolkowski equation:

$$m_{\text{after}} = m_{\text{before}} \cdot e^{\frac{-\Delta v_s}{v_p}}$$

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Keplerian Orbit Transfer

Calculate a transfer between two circular orbits with radius r_A to r_B . The velocity at pericenter of the transfer ellipse:

$$v_p^2 = 2\mu \left(\frac{1}{r_A} - \frac{1}{r_A + r_B} \right) = 2\mu \frac{r_B}{r_A(r_A + r_B)}$$

The required Δv_A to inject from the transfer orbit:

$$\Delta v_A = v_p - v_A = \sqrt{\frac{\mu}{r_A}} \left(\sqrt{\frac{2r_B}{r_A + r_B}} - 1 \right)$$

The required Δv to inject from the transfer orbit into orbit with r_B :

$$\Delta v_B = v_B - v_{apo} = \sqrt{\frac{\mu}{r_B}} \left(1 - \sqrt{\frac{2r_A}{r_A + r_B}} \right)$$

where v_B is the circular velocity at r_B .

The *Hohmann transfer* is the most energy-efficient transfer between two circular orbits.

$$\Delta v_{\text{total}} = \Delta v_A + \Delta v_B = \sqrt{\mu} \left[\sqrt{\left(\frac{2}{r_A} - \frac{2}{r_A + r_B} \right)} - \sqrt{\frac{1}{r_A}} + \sqrt{\frac{2}{r_B} - \frac{2}{r_A + r_B}} - \sqrt{\frac{1}{r_B}} \right]$$

2.3 Mission Analysis

Selection of orbit parameters for most appropriate orbits for the given task.

Earth synchronous orbits: after a given period of time the subsatellite's ground track repeats. *Rotation of Earth:* τ_E sidereal day (= 86161.10555 + 0.015C[s] where C is the number of centuries after 2000). As Earth rotation is in eastward direction, the satellite orbit with respect to Earth seems to have west direction.

$$\Delta \Phi_R = \frac{\tau}{\tau_E} - 2\pi \quad [\text{rad/rev}]$$

τ | satellite orbit period |

J_2 theory taking into account Earth oblateness, we can calculate the satellite orbit plane rotation:

$$\Delta \Omega = \frac{3\pi J_2 R_E^2 \cos i}{a^2(1 - \varepsilon^2)^2} \quad [\text{rad/rev}]$$

total displacement at equator: $\Delta \Phi = \Delta \Phi_R - \Delta \Omega$.

For Earth synchronous orbits, the following property needs to be satisfied:

$$n\Delta \Phi = m2\pi$$

n	number of orbits before identical ground track repeats	
m	number of Earth revolutions (day) before identical ground track repeats	

Sun synchronous orbits:

the sunlight incidents repeats

take into account the motion of Earth around Sun. Due to the Earth's orbit around Sun, the Earth/Sun-line intersect at points P moving westward.

$$\theta = \frac{360^\circ}{365.25 \text{ days}} \approx 1 \quad [^\circ/\text{day}]$$

orbital period of Earth around Sun τ_{ES} : $3.155815 \cdot 10^7 \text{ s} \approx 365.25 \text{ days}$.

angular velocity $\theta = 2\pi \frac{\tau_E}{\tau_{ES}} \quad [\text{rad/day}] = 2\pi \frac{\tau_E}{\tau_{ES}} \cdot \frac{\tau}{\tau_E} \quad [\text{rad/rev}]$

and for Sun synchronous orbits $\Leftrightarrow \Delta\Omega = \theta$.

Sun and Earth synchronous orbit:

$$n(\Delta\Phi_R - \Delta\Omega) = m2\pi \Rightarrow n\tau \left(1 - \frac{\tau_E}{\tau_{ES}}\right) = m\tau_E$$

displacement between subsequent orbits (\rightarrow gap in observation areas):

$$\Delta\Phi = 2\pi\tau \left(\frac{1}{\tau_E} - \frac{1}{\tau_{ES}}\right) = 7.27 \cdot 10^{-5} \cdot \tau \quad [\text{rad/rev}]$$

typical Earth observation orbits: $550 - 750 \text{ km}$ above surface $\Rightarrow \Delta\Phi \approx 4.3 \cdot 10^{-1} \text{ rads}$ and offset $\approx 2875 \text{ km}$.

Sun and earth synchronous satellites :

m	n	corresponding altitude
1	14	894km
1	13	567km
1	12	275km
2	22	

minimum drift orbits: $\frac{n\pm 1}{m} = k$, $k \in \mathbb{N}$, integer. orbits of successive days to fill the displacements between two successive orbits.

$k \mid$ orbits of successive days to fill the displacements between two successive orbits \mid

Example: Landsat 1/2

18 day repeat period minimum drift orbit

$\tau \approx 103.3\text{min}$, $i = 99^\circ$, $\varepsilon = 0.002$, apogee = 920km , descending node: 9 : 38local time

\Rightarrow orbit separation between successive orbits $\approx 2875\text{km}$.

$m = 18 \Rightarrow$ distance between adjacent ground tracks = 160km .

Eclipse calculation:

Angle between earth-sun direction \vec{s} and the orbit plane = β .

$$\beta = \sin^{-1}(\vec{s}\vec{n})$$

$\vec{n} \mid$ normal to orbit plane \mid

$$\sin \beta^* = \frac{r_E}{r_E + h} \Rightarrow \beta^* = \sin^{-1} \left(\frac{r_E}{r_E + h} \right)$$

$E_1 E_2 = 2\Delta v$ eclipse orbital arc

$$\Delta v = \cos^{-1} \left(\frac{\cos \beta^*}{\cos \beta} \right) \text{ in triangle } ACE_1$$

eclipse fraction of a circular orbit:

$$F_e = \frac{2\Delta v}{2\pi} = \frac{1}{\pi} \frac{\sqrt{h^2 + 2r_E h}}{(r_E + h) \cos \beta} \text{ [rad]}$$

h	altitude
β	angle sun-orbit-plane
r_E	earth radius

launch energy: $\varepsilon = \cos^{-1}(\cos^2 i + \sin^2 i - \cos \Delta\Omega)$

change angle related to plane (proportional):

$\Delta\Omega$: longitudinal separation

launch window: $2\frac{\Delta\Omega}{w_c} \Rightarrow \Delta\Omega = w_c t - w_c$ | earth rotation rate |

ground contact and coverage analysis

$$\sin \rho = \cos \lambda_0 = \frac{r_E}{r_E + h}$$

from known λ (from $\Lambda_t, \theta_t / \Lambda_s, \theta_s$) we derive η .

$$\tan \eta = \frac{\sin \rho \sin \lambda}{1 - \sin \rho \sin \lambda}$$

$$\cos \varepsilon = \frac{\sin \eta}{\sin \rho} \Rightarrow \sin \eta = \cos \varepsilon \sin \rho$$

$$\eta + \varepsilon + \lambda = 90^\circ$$

$$D = r_E \frac{\sin \lambda}{\sin \eta}$$

1) angular Radius of Earth φ : $\sin \varphi = \frac{R_E}{R_E + h}$

2) compute spacecraft viewing angles from subsatellite point (Λ_s, θ_s) and target point (Λ_t, θ_t) :

$$\cos \lambda = \sin \theta_s \sin \theta_t + \cos \theta_s \cos \theta_t \cos |\Lambda_s - \Lambda_t|$$

$$\cos A_z = \frac{\sin \theta_t - \cos \lambda \sin \theta_s}{\sin \lambda \cos \theta_s} \quad \text{Azimuth}$$

$$\tan \eta = \frac{\sin \varphi \sin \lambda}{1 - \sin \varphi \cos \lambda}$$

3) compute coordinates on Earth:

$$\cos \varepsilon = \frac{\sin \eta}{\sin \varphi}$$

$$\lambda + \eta + \varepsilon = 90^\circ$$

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ground station contacts

typically a minimum elevation angle is specified for each ground station ε_{\min} below contacts are limited. ($\varepsilon_{\min} \approx 5^\circ - 10^\circ$)

given a ε_{\min} , derive the maximum Earth central angle λ_{\max} and the max. nadir angle η_{\max} , and maximum range D_{\max} at which the satellite is still in the FOV of ground station.

$$\sin \eta_{\max} = \sin \varphi \cos \varepsilon_{\min}$$

$$\lambda_{\max} = 90^\circ - \varepsilon_{\min} - \eta_{\max}$$

$$D_{\max} = R_E \frac{\sin \lambda_{\max}}{\sin \eta_{\max}}$$

to calculate the timing of a flyover, we need λ_{\min} minimum Earth central angle between the satellite's ground track and the ground station.

$$\sin \lambda_{\min} = \sin(90^\circ - i) \sin \sigma_{gs} + \cos(90^\circ - i) \cos \sigma_{gs} \cos \Delta_{\text{long}}$$

where Δ_{long} is the longitude difference between ground station and the orbit pole ($\sigma_{\text{pole}} = 90^\circ - i, \Lambda_{\text{pole}} = \Lambda_{\text{node}} - 90^\circ$).

$$\tan \eta_{\min} = \frac{\sin \varphi \sin \lambda_{\min}}{1 - \sin \varphi \cos \lambda_{\min}}$$

$$\varepsilon_{\max} = 90^\circ - \lambda_{\min} - \eta_{\min}$$

$$D_{\min} = R_E \frac{\sin \lambda_{\min}}{\sin \eta_{\min}}$$

parameters at the satellites closest approach to ground station.

the maximum angular rate of the satellite, seen from the ground station

$$\sigma_{\max} = \frac{v_{\text{sat}}}{D_{\min}} = \frac{2\pi(R_E + h)}{TD_{\min}}$$

where T is the orbital period

$$AT = \frac{T}{180^\circ} \arccos \left(\frac{\cos \lambda_{\max}}{\cos \lambda_{\min}} \right)$$

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4.1 Onboard Data Handling (OBDH)

OBDH functions:

- to receive, validate, decode and distribute commands
- to gather, process and format the data onboard, spacecraft housekeeping and mission data
- spacecraft time synchronisation, health monitoring of computer itself (watchdogs), security interfaces
- autonomous reaction capabilities, redundancy

standards for command message formats CCSDS:

- synchronisation code
- address bit
- command message bits
- error check bits

Output: discrete commands (fixed amplitude, fixed pulse duration)

Data handling combines telemetry from multiple sources and provides it for downlink to earth.

- analog telemetry data (conversion needed)

4.2 Command & Data Handling System Sizing

Command processing

- number of command output channels
- requirements for stored commands
- requirements for computer commands of ACS functions

Telemetry processing

- spacecraft housekeeping data

- feedback for onboard control of spacecraft functions
- routing of onboard and subsystem data
 - to and from receivers and transmitters
 - to storage systems
 - to affected system controllers

Timing

- time granularity
- stability requirement
- acceptable uncertainty

Computer watchdog

constraints:

Spacecraft Bus constraints

- single unit systems
- multiple units, distributing system
- integrated systems

Reliability (→ redundancy/parts quality)

Radiation

Budgets (Size, Weight, Power)

⇒ OBDH is a tool to configure, control or program the payload and the spacecraft subsystems.

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