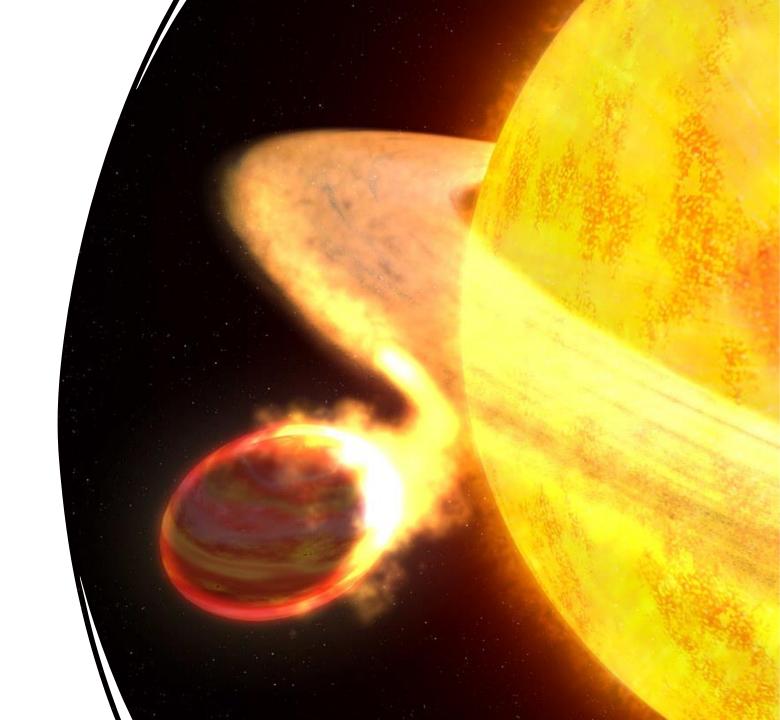
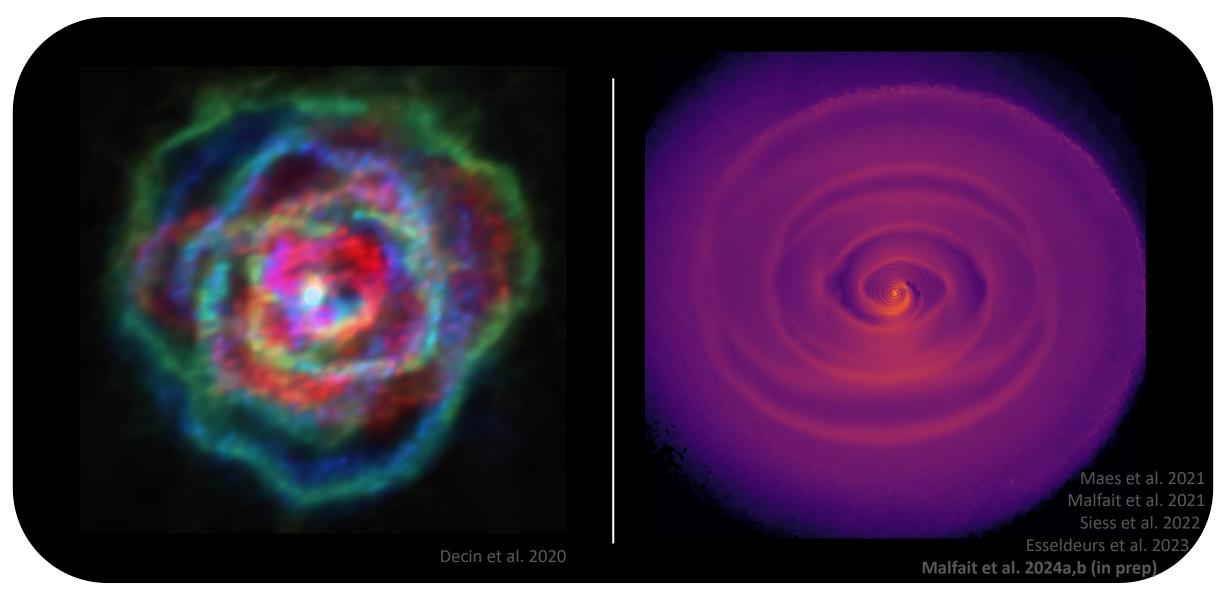
Tidal Dissipation in Cool Evolved Stars

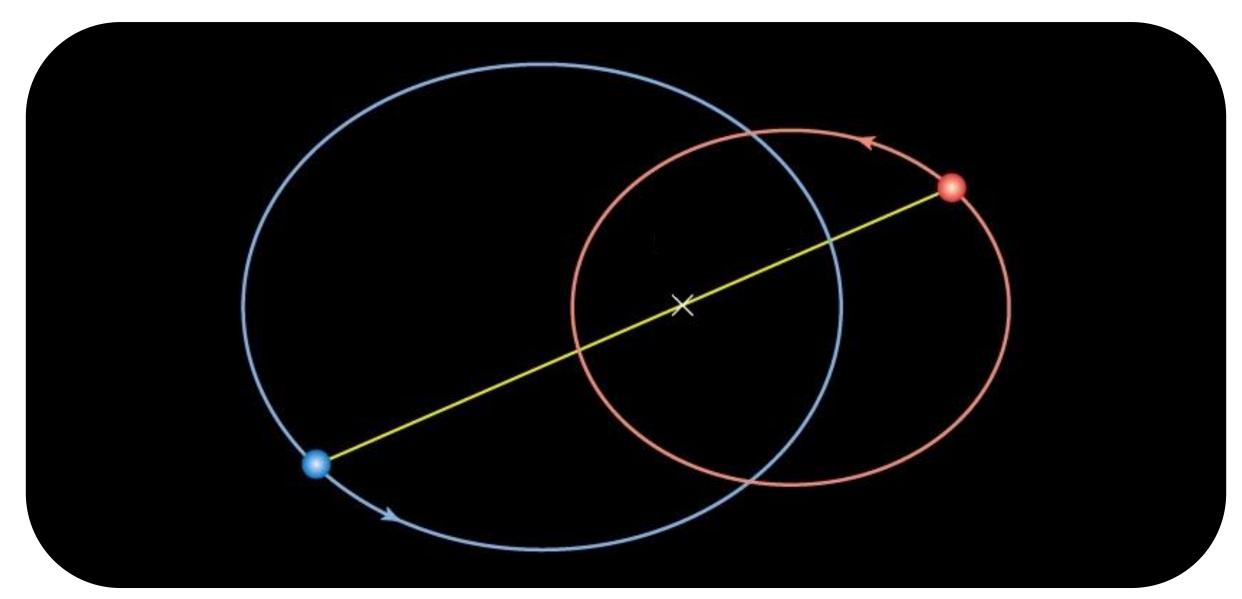
Mats Esseldeurs Stéphane Mathis Leen Decin



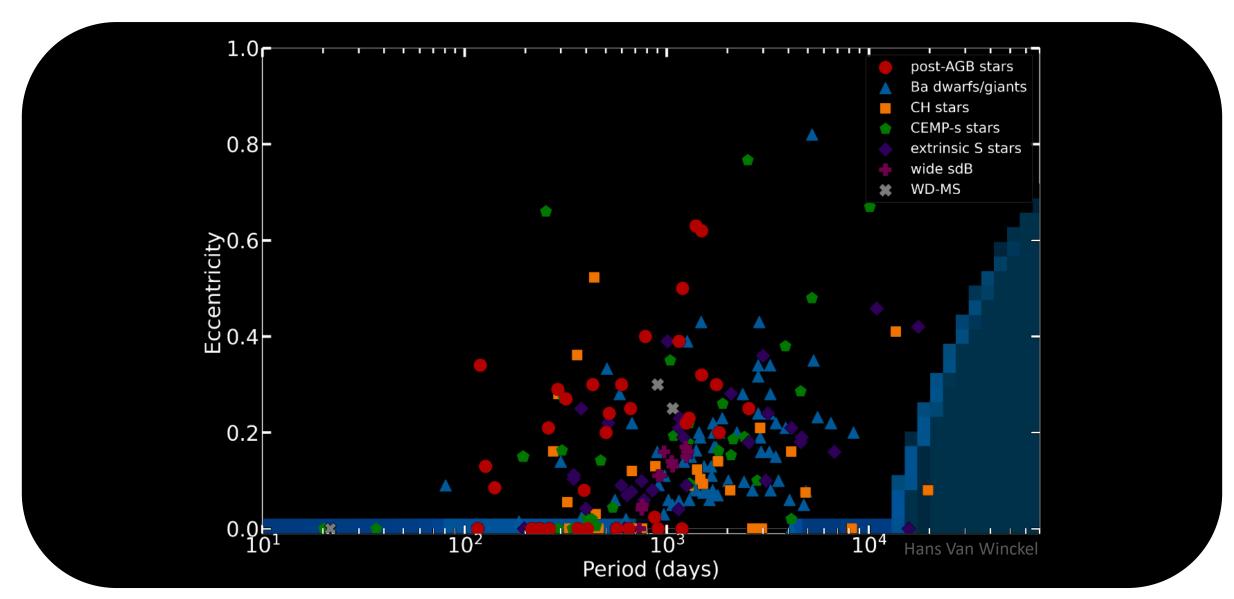
Observations and Simulations



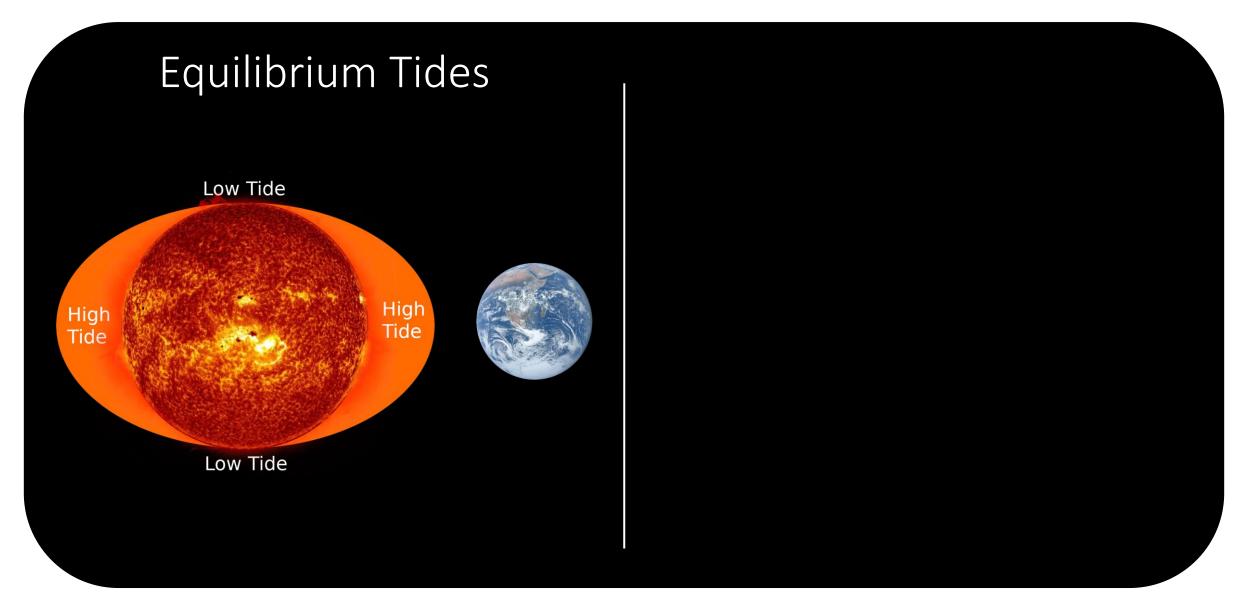
Orbital properties of binary systems



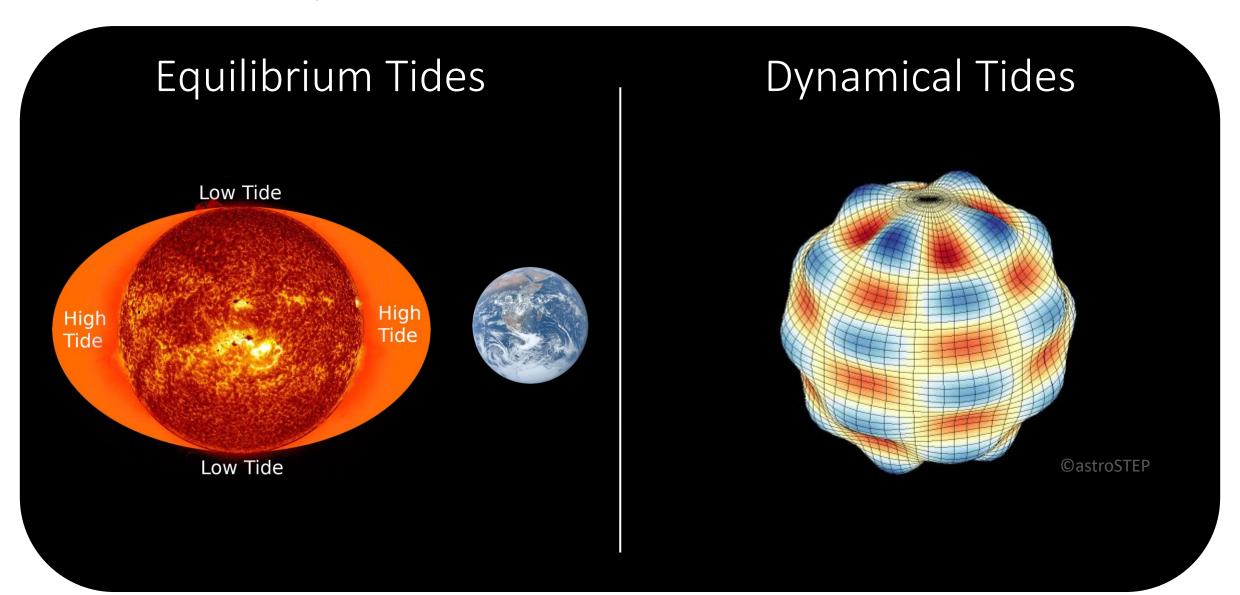
Statistics of orbital properties



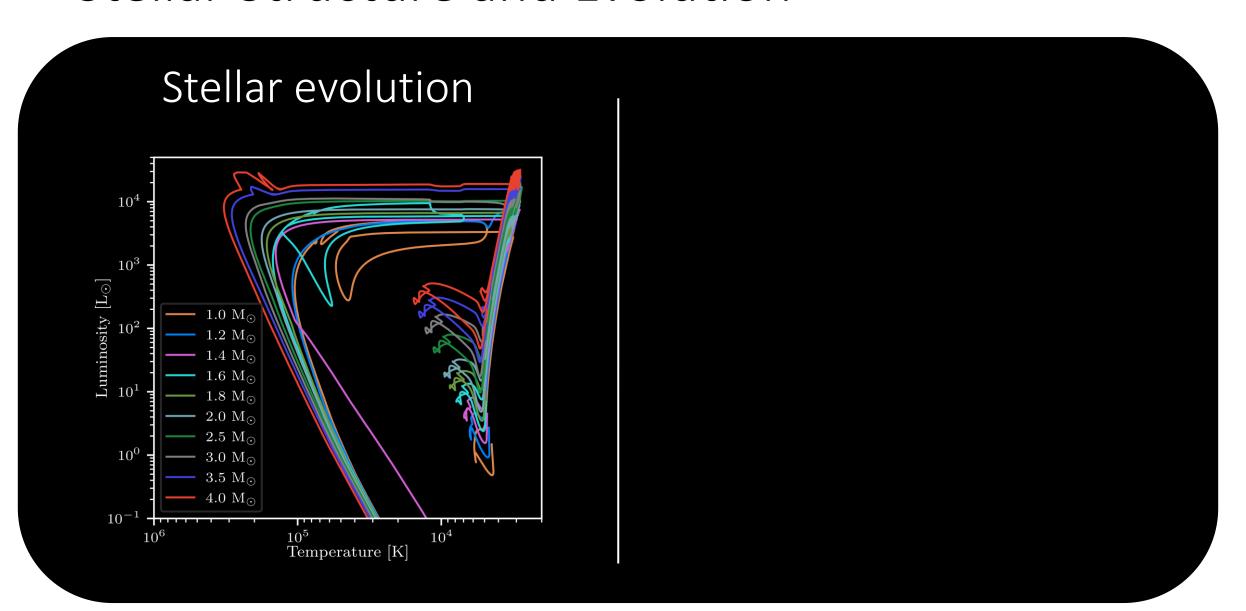
Tidal Dissipation mechanisms



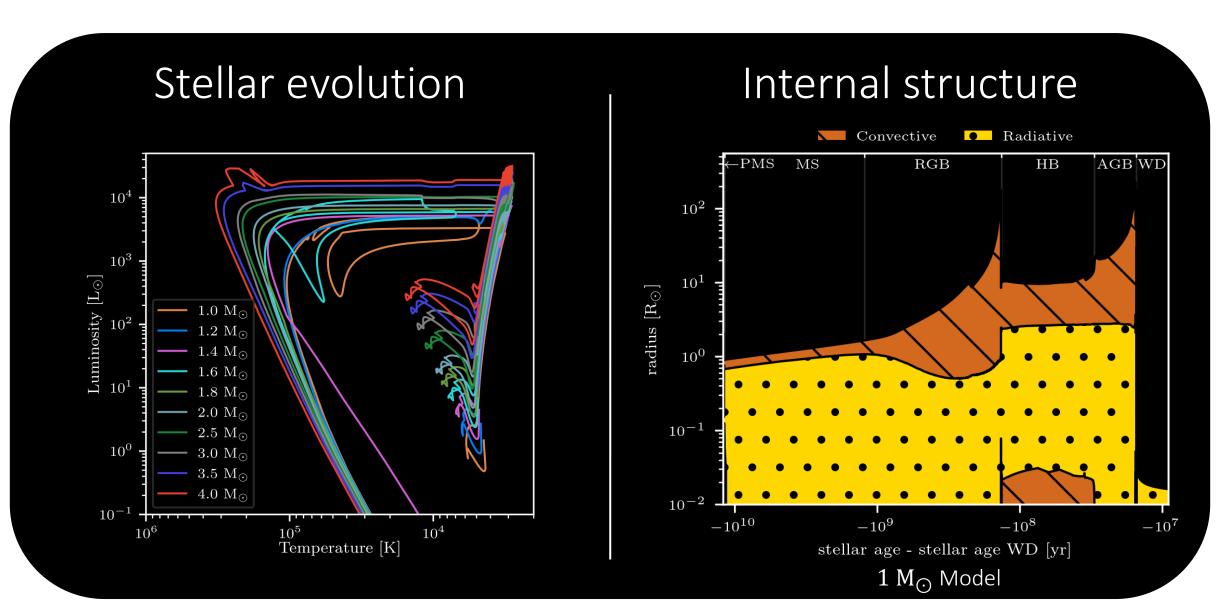
Tidal Dissipation mechanisms



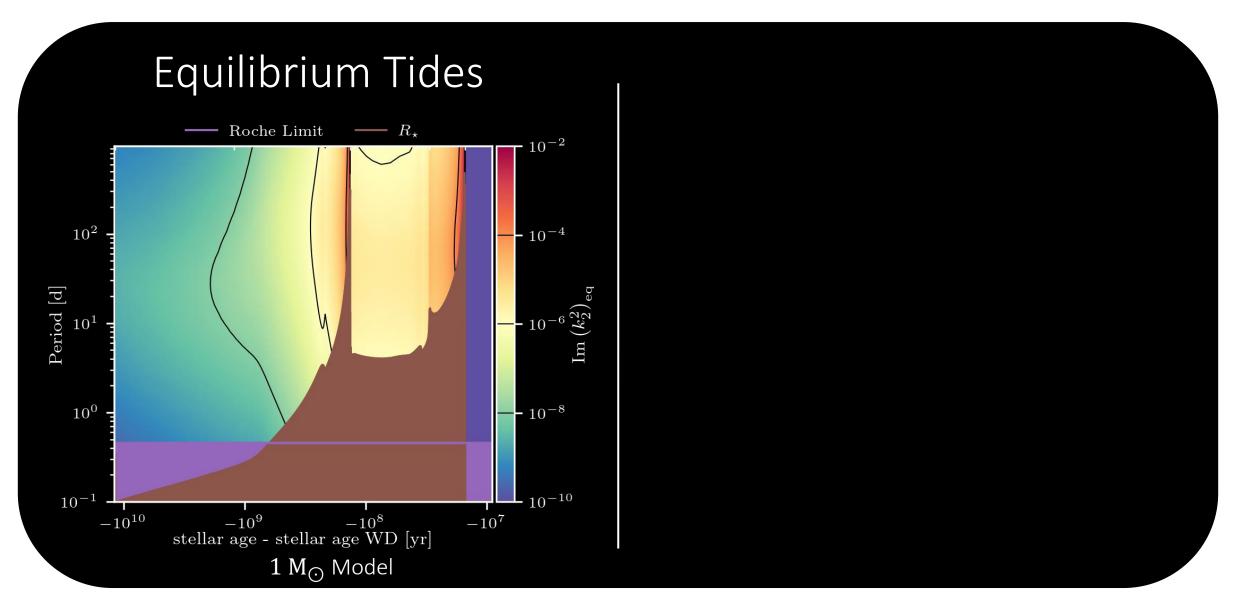
Stellar Structure and Evolution



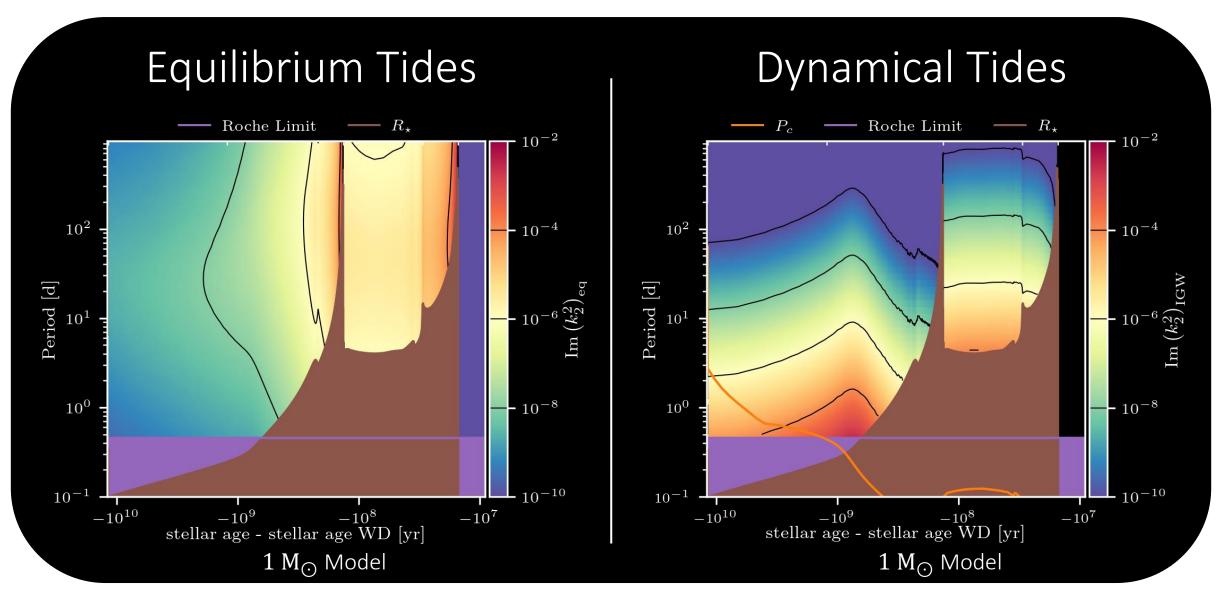
Stellar Structure and Evolution



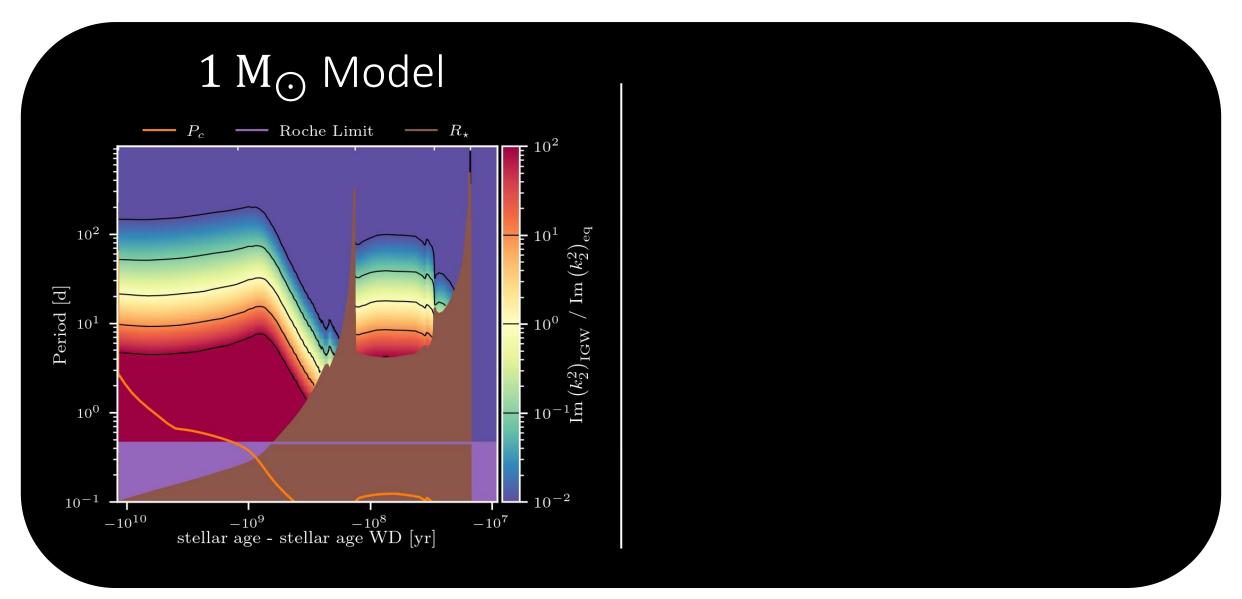
Tidal Dissipation in Cool Evolved Stars



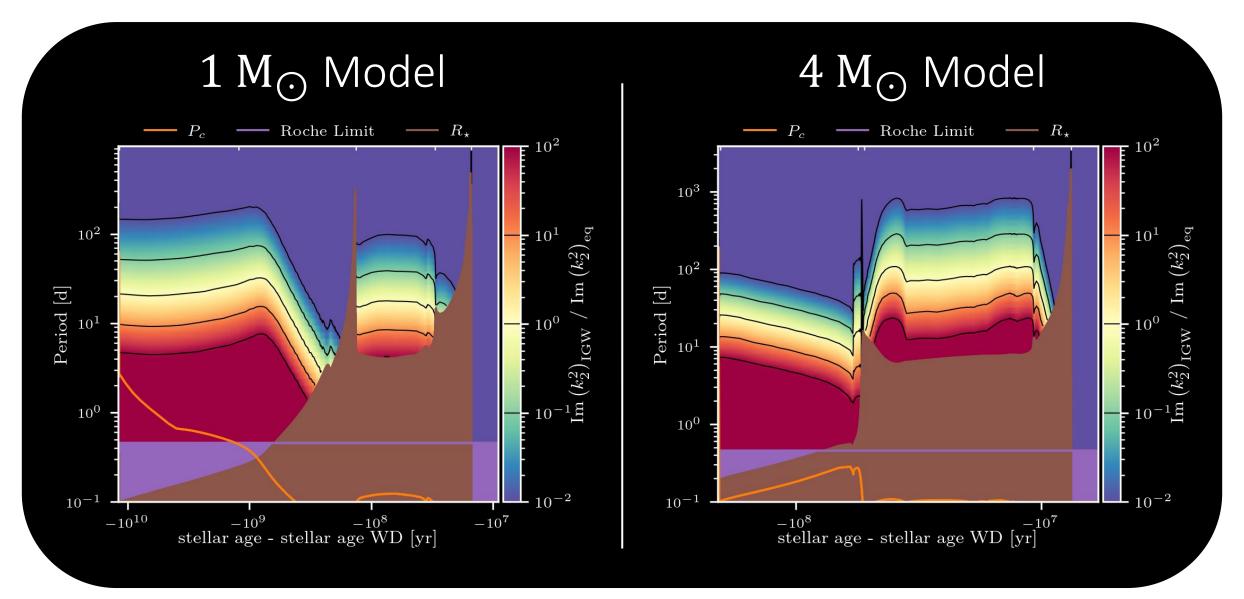
Tidal Dissipation in Cool Evolved Stars



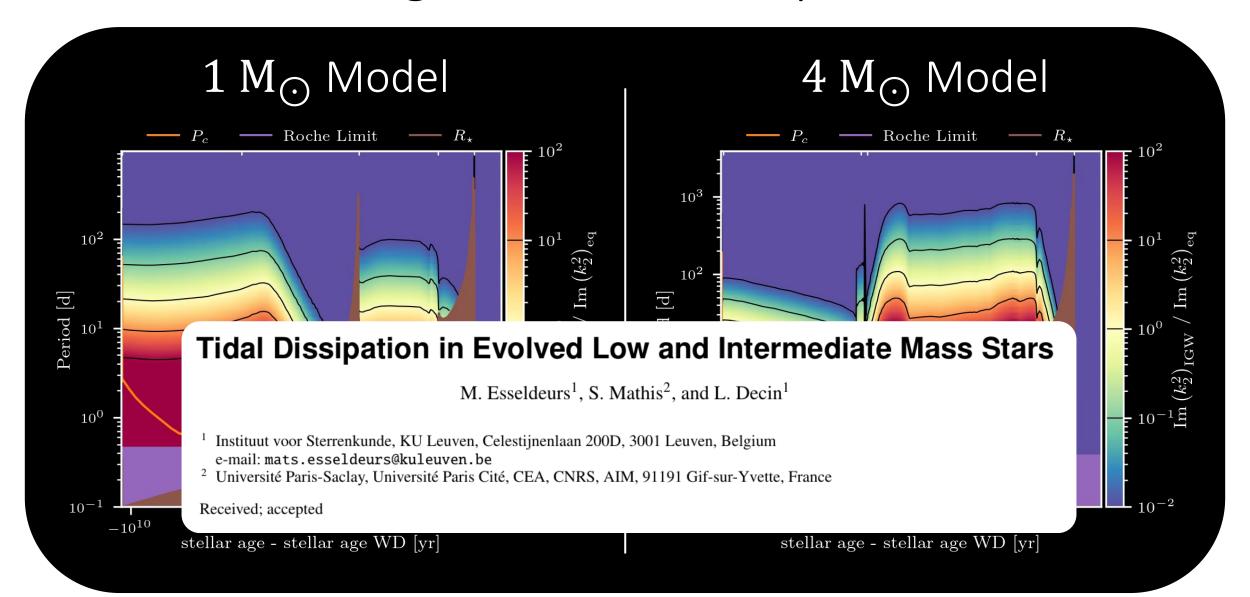
Relative strengths of tidal dissipation



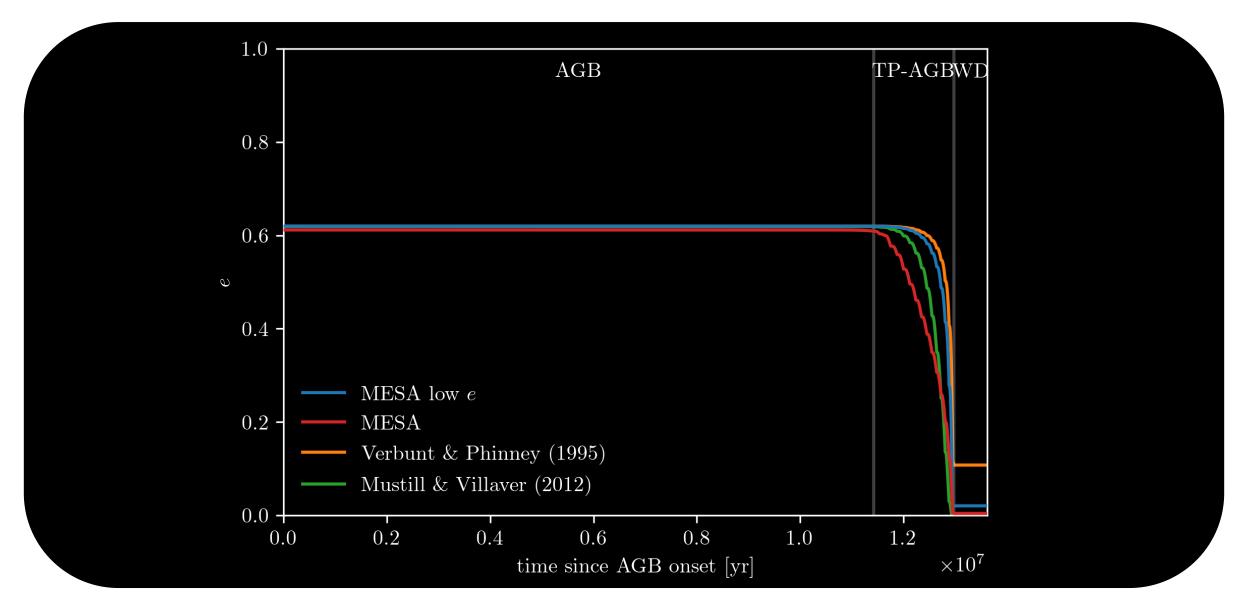
Relative strengths of tidal dissipation



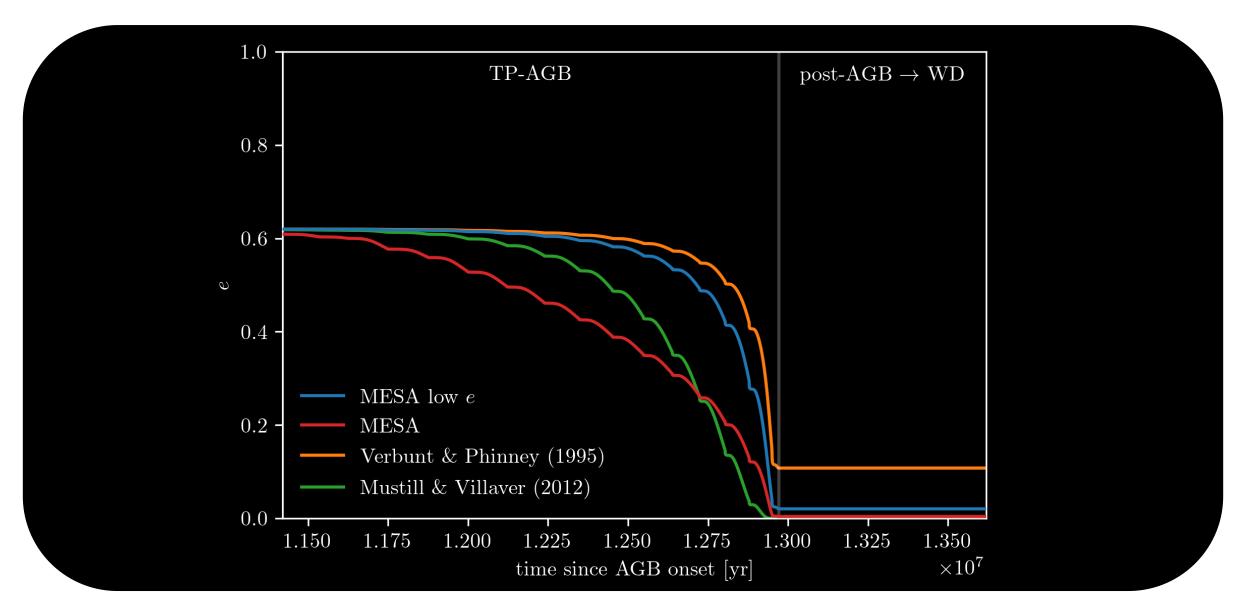
Relative strengths of tidal dissipation



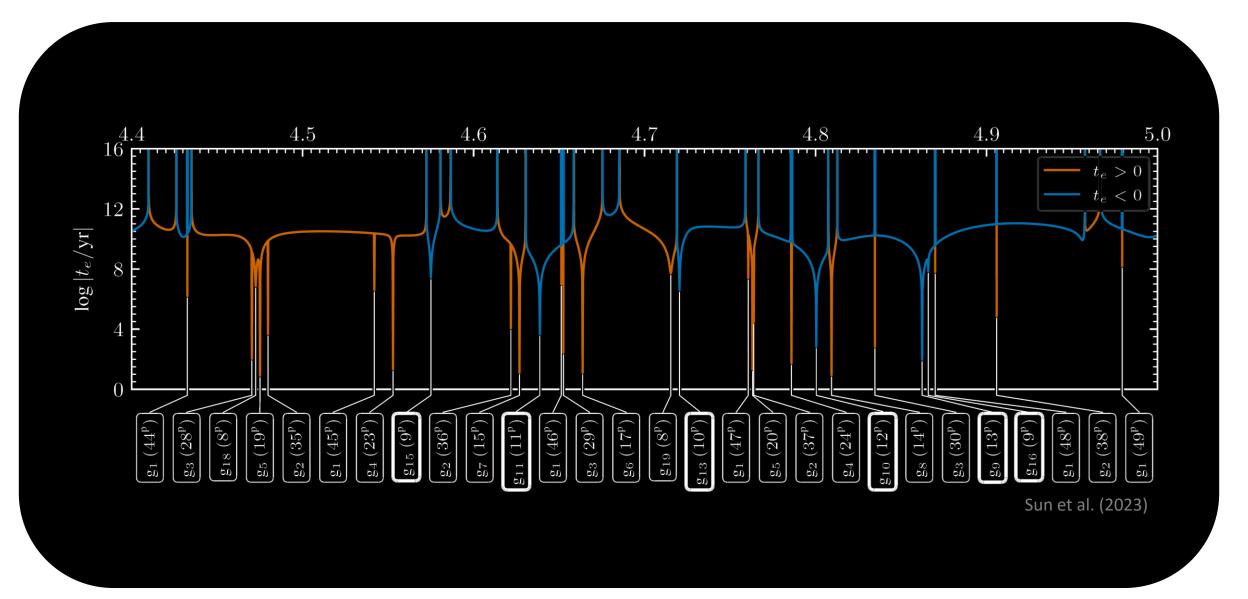
Eccentricity change during the AGB phase



Eccentricity change during the AGB phase



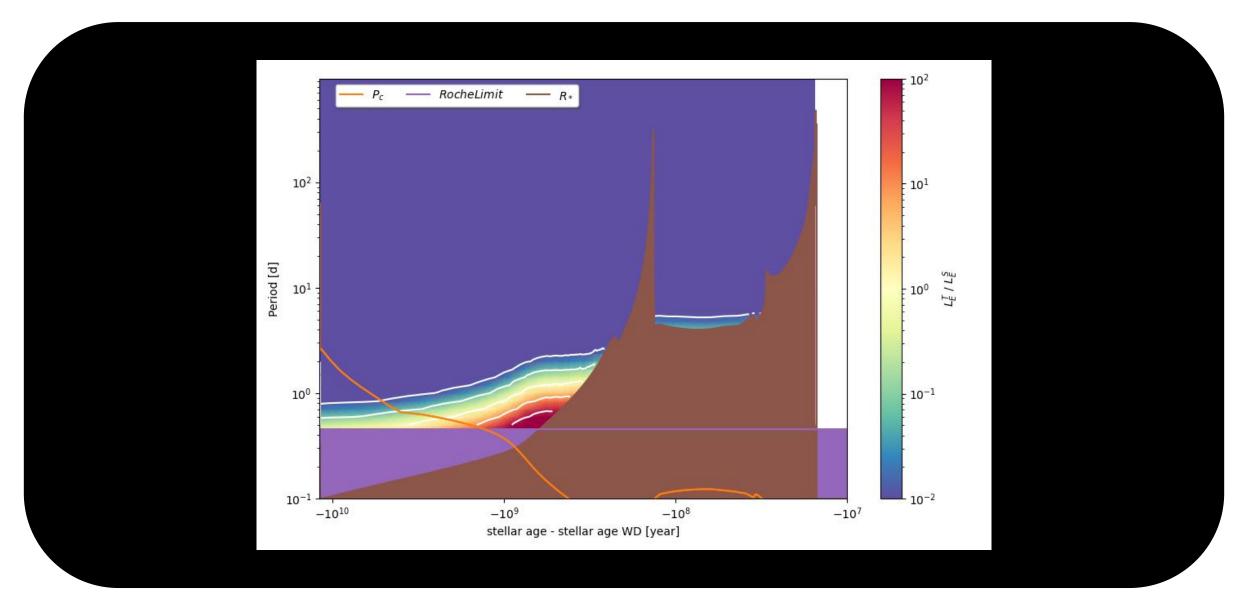
Tidal eccentricity pumping through resonances



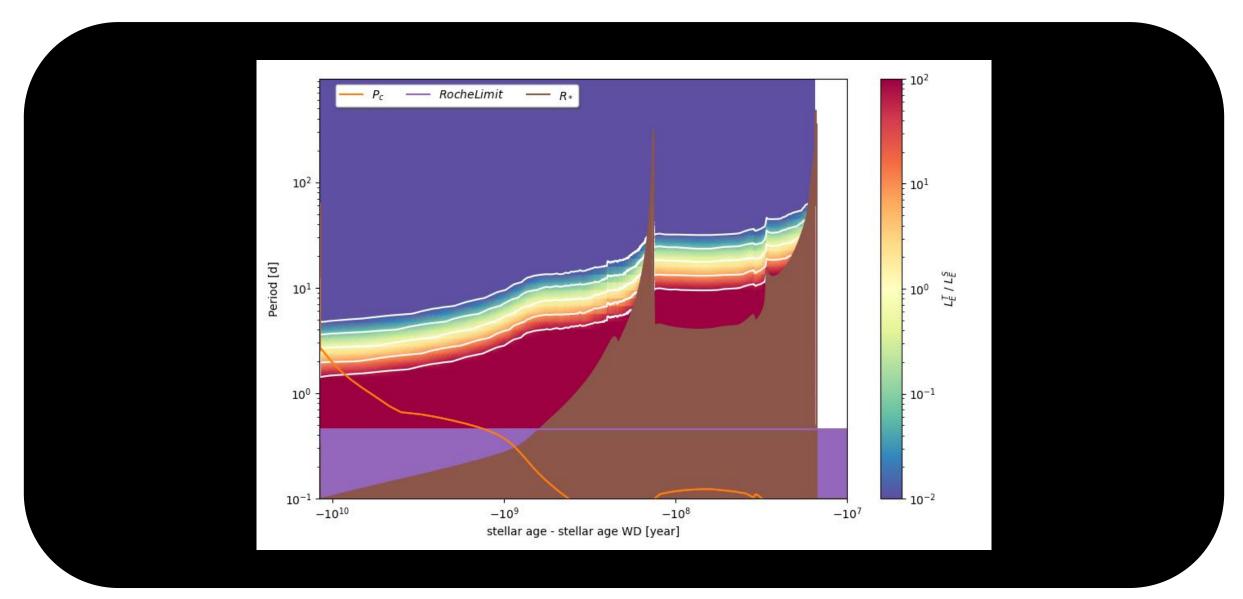
Conclusions

- Tidal dissipation can be calculated ab-initio throughout the entire lifetime of a star
- The dynamical tide of gravity waves remains moderate during the giant phases
- The eccentricity problem is not yet solved
- The dynamical tide connecting with pressure modes remains to be studied

Mixing due to tidal waves



Mixing due to tidal waves



Tidal Dissipation in Cool Evolved Stars

Equilibrium Tides

$$\begin{split} &\frac{1}{r^2} \frac{\mathrm{d}}{\mathrm{d}r} \left(r^2 \frac{\mathrm{d}\Phi_l^{\mathrm{nw}}}{\mathrm{d}r} \right) - \frac{l(l+1)}{r^2} \Phi_l^{\mathrm{nw}} - 4\pi G \frac{\mathrm{d}\rho_0}{\mathrm{d}r} \frac{1}{g_0} \left(\Phi_l^{\mathrm{nw}} + \Psi_l \right) = 0 \\ &\left\{ \frac{\mathrm{d}\ln\Phi_l^{\mathrm{nw}}}{\mathrm{d}\ln r} = l & \text{at } r = \eta R_\star \text{ for } \eta \to 0 \\ \frac{\mathrm{d}\ln\Phi_l^{\mathrm{nw}}}{\mathrm{d}\ln r} = -(l+1) & \text{at } r = R_\star \\ &\xi_{r,l}^{\mathrm{nw}} = -\frac{\Phi_l^{\mathrm{nw}} + \Psi_l}{g_0} , & \xi_{h,l} = \frac{1}{l(l+1)} \left(2\xi_{r,l}^{\mathrm{nw}} + r \frac{\mathrm{d}\xi_{r,l}^{\mathrm{nw}}}{\mathrm{d}r} \right) . \\ &D_l(r) = \frac{1}{3} \left(3 \frac{\mathrm{d}\xi_{r,l}^{\mathrm{nw}}}{\mathrm{d}r} - \frac{1}{r^2} \frac{\mathrm{d} \left(r^2 \xi_{r,l}^{\mathrm{nw}} \right)}{\mathrm{d}r} + l(l+1) \frac{\xi_{h,l}^{\mathrm{nw}}}{r} \right)^2 \\ &+ l(l+1) \left(\frac{\xi_{r,l}^{\mathrm{nw}}}{r} + r \frac{\mathrm{d} \left(\xi_{h,l}^{\mathrm{nw}} / r \right)}{\mathrm{d}r} \right)^2 \\ &+ l(l+1) l(l+1) (l+2) \left(\frac{\xi_{h,l}^{\mathrm{nw}}}{r} \right)^2 , \end{split}$$

$$\mathrm{Im} \left(k_2^2 \right)_{\mathrm{eq}} = \frac{16\pi G \omega_t}{4(2l+1)R_\star |\varphi_T(R_\star)|^2} \int_0^{R_\star} r^2 \rho \nu_t D_l(r) \mathrm{d}r , \end{split}$$

Dynamical Tides

$$\begin{split} \mathcal{F}_{\text{in}} &= \int_{0}^{r_{\text{in}}} \left[\left(\frac{r^{2} \varphi_{T}}{g_{0}} \right)'' - \frac{l(l+1)}{r^{2}} \left(\frac{r^{2} \varphi_{T}}{g_{0}} \right) \right] \frac{X_{1,\text{in}}}{X_{1,\text{in}} (r_{\text{in}})} \mathrm{d}r \\ \mathcal{F}_{\text{out}} &= \int_{r_{\text{out}}}^{R_{\star}} \left[\left(\frac{r^{2} \varphi_{T}}{g_{0}} \right)'' - \frac{l(l+1)}{r^{2}} \left(\frac{r^{2} \varphi_{T}}{g_{0}} \right) \right] \frac{X_{1,\text{out}}}{X_{1,\text{out}} (r_{\text{out}})} \mathrm{d}r \\ \begin{cases} X_{1,\text{out}}'' - \frac{\partial_{r} \rho_{0}}{\rho_{0}} X_{1,\text{out}}' - \frac{l(l+1)}{r^{2}} X_{1,\text{out}} = 0 \\ X_{1,\text{out}}(r)_{r \to 0} \propto \left(\frac{1}{2} + \sqrt{\frac{1}{4} + l(l+1)}}{r^{2}} X_{1,\text{in}} = 0 \\ X_{1,\text{out}}(r)_{r \to R_{\star}} \propto \rho_{0} \left(r - R_{\star} - \frac{\varphi_{T}(R_{\star})}{g_{0}(R_{\star})} \right) \\ X_{1,\text{out}}'(r)_{r \to R_{\star}} \propto \rho_{0}(R_{\star}) \end{split}$$

Tidal Dissipation in Cool Evolved Stars

Equilibrium Tides

$$\begin{split} &\frac{1}{r^2} \frac{\mathrm{d}}{\mathrm{d}r} \left(r^2 \frac{\mathrm{d}\Phi_l^{\mathrm{nw}}}{\mathrm{d}r} \right) - \frac{l(l+1)}{r^2} \Phi_l^{\mathrm{nw}} - 4\pi G \frac{\mathrm{d}\rho_0}{\mathrm{d}r} \frac{1}{g_0} \left(\Phi_l^{\mathrm{nw}} + \Psi_l \right) = 0 \\ &\left\{ \frac{\mathrm{d} \ln \Phi_l^{\mathrm{nw}}}{\mathrm{d} \ln r} = l & \text{at } r = \eta R_{\star} \text{ for } \eta \to 0 \\ \frac{\mathrm{d} \ln \Phi_l^{\mathrm{nw}}}{\mathrm{d} \ln r} = -(l+1) & \text{at } r = R_{\star} \\ &\xi_{r,l}^{\mathrm{nw}} = -\frac{\Phi_l^{\mathrm{nw}} + \Psi_l}{g_0} , & \xi_{h,l} = \frac{1}{l(l+1)} \left(2\xi_{r,l}^{\mathrm{nw}} + r \frac{\mathrm{d}\xi_{r,l}^{\mathrm{nw}}}{\mathrm{d}r} \right) . \\ &D_l(r) = \frac{1}{3} \left(3 \frac{\mathrm{d}\xi_{r,l}^{\mathrm{nw}}}{\mathrm{d}r} - \frac{1}{r^2} \frac{\mathrm{d} \left(r^2 \xi_{r,l}^{\mathrm{nw}} \right)}{\mathrm{d}r} + l(l+1) \frac{\xi_{h,l}^{\mathrm{nw}}}{r} \right)^2 \\ &+ l(l+1) \left(\frac{\xi_{r,l}^{\mathrm{nw}}}{r} + r \frac{\mathrm{d} \left(\xi_{h,l}^{\mathrm{nw}} / r \right)}{\mathrm{d}r} \right)^2 \\ &+ (l-1)l(l+1)(l+2) \left(\frac{\xi_{h,l}^{\mathrm{nw}}}{r} \right)^2 , \end{split}$$

$$\mathrm{Im} \left(k_2^2 \right)_{\mathrm{eq}} = \frac{16\pi G \omega_t}{4(2l+1)R_{\star} |\varphi_T(R_{\star})|^2} \int_0^{R_{\star}} r^2 \rho \nu_t D_l(r) \mathrm{d}r , \end{split}$$

Dynamical Tides

$$\mathcal{F}_{\text{in}} = \int_{0}^{r_{\text{in}}} \left[\left(\frac{r^{2} \varphi_{T}}{g_{0}} \right)'' - \frac{l(l+1)}{r^{2}} \left(\frac{r^{2} \varphi_{T}}{g_{0}} \right) \right] \frac{X_{1,\text{in}}}{X_{1,\text{in}}(r_{\text{in}})} dr$$

$$\mathcal{F}_{\text{out}} = \int_{r_{\text{out}}}^{R_{\star}} \left[\left(\frac{r^{2} \varphi_{T}}{g_{0}} \right)'' - \frac{l(l+1)}{r^{2}} \left(\frac{r^{2} \varphi_{T}}{g_{0}} \right) \right] \frac{X_{1,\text{out}}}{X_{1,\text{out}}(r_{\text{out}})} dr$$

$$\text{Im} \left(k_{2}^{2} \right)_{\text{IGW}} = \frac{3^{-\frac{1}{3}} \Gamma^{2} \left(\frac{1}{3} \right)}{2\pi} [l(l+1)]^{-\frac{4}{3}} \omega_{t}^{\frac{8}{3}} \frac{a^{6}}{GM_{2}^{2} R_{\star}^{5}}$$

$$\times \left(\rho_{0} (r_{\text{in}}) r_{\text{in}} \left| \frac{dN^{2}}{d \ln r} \right|_{r_{\text{in}}}^{-\frac{1}{3}} \mathcal{F}_{\text{in}}^{2} \right)$$

$$+ \rho_{0} (r_{\text{out}}) r_{\text{out}} \left| \frac{dN^{2}}{d \ln r} \right|_{r_{\text{out}}}^{-\frac{1}{3}} \mathcal{F}_{\text{out}}^{2} \right)$$