

# 量统2015郭永期中

Deschain

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一、

1.(1)体系能量有确定值的状态。

(2)粒子在无穷远处出现的概率为0的状态。

(3)若 $\psi(x) = \psi(-x)$ ，则是偶（正）宇称；若 $\psi(x) = -\psi(-x)$ ，则是奇（负）宇称。

(4)能量为E的粒子有一定概率穿透高度为 $U_0 (U_0 > E)$ 的势垒。

2.

$$E = \hbar\omega = h\nu, \lambda = \frac{h}{p}$$

3.

$$[a, a^\dagger] = \frac{m\omega}{2\hbar}([x, x] + \frac{i}{m\omega}([\hat{p}_x, x] - [x, \hat{p}_x]) + \frac{i}{m^2\omega^2}[\hat{p}_x, \hat{p}_x]) = \frac{m\omega}{2\hbar} \times \frac{i}{m\omega} \times 2i\hbar = 1$$

4.

$$[\hat{p}_x, f(x)]\psi(x) = \hat{p}_x(f(x)\psi(x)) - f(x)\hat{p}_x\psi(x) = -i\hbar\frac{\partial}{\partial x}(f(x)\psi(x)) + i\hbar f(x)\frac{\partial}{\partial x}\psi(x) = -i\hbar(\frac{\partial}{\partial x}f(x))\psi(x)$$

$$[\hat{p}_x, f(x)] = -i\hbar\frac{\partial}{\partial x}f(x)$$

$$[\hat{L}_z, \hat{L}_+] = [\hat{L}_z, \hat{L}_x] + i[\hat{L}_z, \hat{L}_y] = i\hbar\hat{L}_y + \hbar\hat{L}_x = \hbar\hat{L}_+$$

$$[\hat{L}_z, \hat{L}_-] = [\hat{L}_z, \hat{L}_x] - i[\hat{L}_z, \hat{L}_y] = i\hbar\hat{L}_y - \hbar\hat{L}_x = -\hbar\hat{L}_-$$

5. 设 $a_n$ 对应的本征函数 $|\psi_{a_n}\rangle$ ,  $b_n$ 对应的本征函数 $|\psi_{b_n}\rangle$ ,

$$\langle \psi_{a_n} | \psi_{b_n} \rangle = \langle \psi_{b_n} | \psi_{a_n} \rangle$$

$$[\hat{A}, \hat{B}]\psi = \hat{A}\hat{B}\psi - \hat{B}\hat{A}\psi = \hat{A}\sum_i b_i\psi_{b_i} - \hat{B}\sum_k a_k\psi_{a_k} = \sum_i b_i \sum_j a_j\psi_{a_j, b_i} - \sum_k a_k \sum_l b_l\psi_{a_k, b_l}$$

$$\because P(a_n, b_n) = P(b_n, a_n) \therefore b_i = b_l, a_j = a_k (i = l, j = k) \therefore [\hat{A}, \hat{B}]\psi = 0, [\hat{A}, \hat{B}] = 0$$

6. 记 $\hat{L}_z$ 的本征态为 $|m\rangle$ ,

$$\langle m | \hat{L}_z | m \rangle = m\hbar | m \rangle$$

$$[\hat{L}_x, \hat{L}_y] | m \rangle = i\hbar \hat{L}_z | m \rangle = i\hbar^2 m | m \rangle$$

$\therefore$ 在 $\hat{L}_z$ 的本征态 $|m\rangle$ 中,  $[\hat{L}_x, \hat{L}_y] = i\hbar^2 m$

7. 原式 $= (\psi_n, (\hat{H}\hat{A} - \hat{A}\hat{H})\psi_n) = (\psi_n, \hat{H}\hat{A}\psi_n) - (\psi_n, \hat{A}\psi_n)E_n$

$\because \hat{H}$ 是线性Hermite算符  $\therefore (\psi_n, \hat{H}\hat{A}\psi_n) = (\hat{A}\psi_n, \hat{H}\psi_n) = E_n(\hat{A}\psi_n, \psi_n)$

$\because \hat{A}$ 是线性Hermite算符  $\therefore (\hat{A}\psi_n, \psi_n) = (\psi_n, \hat{A}\psi_n) \therefore (\psi_n, [\hat{H}, \hat{A}]\psi_n) = 0$

二、

1.

$$B^\dagger = A^\dagger(AA^\dagger)A = A^\dagger(1 - A^\dagger A)A = A^\dagger A - (A^\dagger)^\dagger A^\dagger = A^\dagger A = B$$

2.  $\because B^2 = B \therefore B$  有 2 个本征值  $\lambda_0 = 1, \lambda_1 = 0$

$$B = \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix}, \quad A = \begin{bmatrix} a_1 & a_2 \\ a_3 & a_4 \end{bmatrix}, \quad B^2 = A^\dagger A, \quad A^2 = 0$$

$$A = \begin{bmatrix} 0 & 0 \\ e^{i\delta} & 0 \end{bmatrix}$$

1 三、

1. 基态:

$$E_0 = \frac{\pi^2 \hbar^2}{ma^2}, \quad \psi_0 = \frac{2}{a} \sin\left(\frac{\pi x_1}{a}\right) \sin\left(\frac{\pi x_2}{a}\right), \quad f_0 = 1$$

2. 第一激发态:

$$E_1 = \frac{5\pi^2 \hbar^2}{2ma^2}$$

$$\psi_{11} = \frac{2}{a} \left[ \sin\left(\frac{\pi x_1}{a}\right) \sin\left(\frac{2\pi x_2}{a}\right) - \sin\left(\frac{\pi x_1}{a}\right) \sin\left(\frac{\pi x_2}{a}\right) \right] \times \begin{cases} |10\rangle \\ |11\rangle \\ |1-1\rangle \end{cases}$$

$$\psi_{12} = \frac{2}{a} \left[ \sin\left(\frac{\pi x_1}{a}\right) \sin\left(\frac{2\pi x_2}{a}\right) + \sin\left(\frac{\pi x_1}{a}\right) \sin\left(\frac{\pi x_2}{a}\right) \right] |00\rangle$$

$$f_1 = 4$$

四、

$$\psi(r, \theta, \varphi) = \sqrt{\frac{8\pi}{3}} R(r) Y_{11}(\theta, \varphi) + \sqrt{\frac{4\pi}{3}} R(r) Y_{10}(\theta, \varphi)$$

$$(1) L^2 = 2\hbar^2, \quad P = 1$$

$$(2) P(L_z = 0) = \frac{1}{3}, P(L_z = \hbar) = \frac{2}{3}, \overline{L_z} = \frac{2}{3}\hbar$$

五、

$$\hat{J} = \hat{S}_1 + \hat{S}_2, \quad \hat{H} = A\hat{J}_z + \frac{B}{2}(\hat{J}^2 - \hat{S}_1^2 - \hat{S}_2^2)$$

$$(1) j = 0, m = 0, E_1 = -\frac{3}{4}B\hbar^2, \psi_1 = |00\rangle$$

$$(2) j = 1, m = 1, E_2 = A\hbar + \frac{1}{4}B\hbar^2, \psi_2 = |11\rangle$$

$$(3) j = 1, m = 0, E_3 = \frac{1}{4}B\hbar^2, \psi_3 = |10\rangle$$

$$(4) j = 1, m = -1, E_4 = -A\hbar + \frac{1}{4}B\hbar^2, \psi_4 = |1-1\rangle$$

2 六、

$$H' = -q\varepsilon x, \quad E_n^{(0)} = (n + \frac{1}{2})\hbar\omega$$

$$E_n^{(1)} = \langle n | -q\varepsilon x | n \rangle = 0$$

$$H'_{n+1,n} = \langle n+1 | -q\varepsilon x | n \rangle = \sqrt{\frac{(n+1)\hbar}{2\mu\omega}}, \quad \text{else } H'_{m,n} = 0$$

$$E_n^{(2)} = \frac{|H'_{n+1,n}|^2}{E_n^{(0)} - E_{n+1}^{(0)}} = -\frac{n+1}{2\mu\omega^2}$$

$$\psi_n = \frac{H'_{n+1,n}\psi_{n+1}^{(0)}}{E_n^{(0)} - E_{n+1}^{(0)}} = -\sqrt{\frac{n+1}{2\mu\hbar\omega^3}} |n+1\rangle$$

$$\therefore \psi_n = |n\rangle - \sqrt{\frac{n+1}{2\mu\omega^3}} |n+1\rangle, \quad E_n = (n + \frac{1}{2})\hbar\omega - \frac{n+1}{2\mu\omega^2}$$