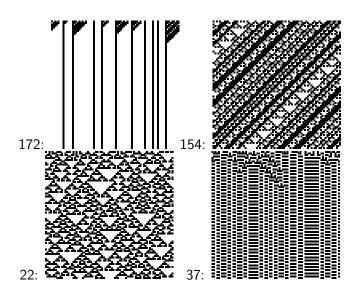
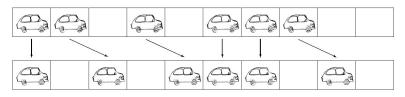
Some other elementary rules



Traffic model

Are CA useful in practice? Consider the following traffic model:



Suppose that $s_t(i)=1$ if there is a car at cell i, otherwise $s_t(i)=0$. Then we have CA with local rule

$$f(0,0,0) = 0, f(0,0,1) = 0, f(0,1,0) = 0, f(0,1,1) = 1,$$

 $f(1,0,0) = 1, f(1,0,1) = 1, f(1,1,0) = 0, f(1,1,1) = 1,$

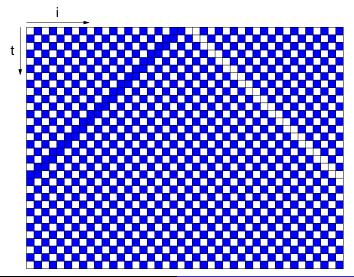
or equivalently

$$s_{t+1}(i) = s_t(i-1) + s_t(i)s_t(i+1) - s_t(i-1)s_t(i).$$

This is rule with W(f) = 184.

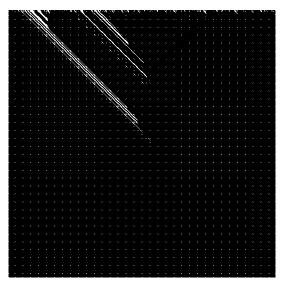
Spatio-temporal diagram for rule 184

If the initial configuration for rule 184 contains only one pair 11 followed by 00, we obtain:



Spatio-temporal diagram for rule 184, 500 sites

Larger simulation, M=500, 500 time steps.



Notation

Let $s_t(i)$ be the state of the cell i at time t. We will typically assume that the initial configuration is constructed randomly, and it is drawn from Bernoulli distribution, meaning that

$$Pr\left(s_0(i) = 1\right) = p \text{ for all } i,$$

$$Pr\left(s_0(i) = 0\right) = 1 - p \text{ for all } i.$$

If p=1/2, we will call this symmetric Bernoulli distribution. Let us also define probability of occurrence of block of symbols $a_1a_2\dots a_k$ at time t by

$$P_t(a_1 a_2 \dots a_k) = Pr(s_t(i) = a_1, s_t(i+1) = a_1, \dots, s_t(i+k) = a_k)$$

The above will be called block probability. Note that the block probability is independent of i, because both the initial configuration and the CA rule is are translationally-invariant.

Simulations

- In simulations, we will be typically using lattice of size M with periodic boundary condition, that is, we will take $i=0,1,\ldots M-1$, and assume that the left neighbour of $s_t(0)$ is $s_t(M-1)$, while the right neighbour of $s_t(0)$.
- In many cases, periodic boundary conditions approximate (in some respects) infinite system. Detailed discussion of this topic is unfortunately beyond the scope of this course.
- Moreover, if M is very large,

$$P_t(a_1, a_2, \dots a_k) \approx \frac{N_t(a_1, a_2, \dots a_k)}{M},$$

where $N_t(a_1, a_2, \dots a_k)$ is the number of occurrences of the block $a_1, a_2, \dots a_k$ after t iterations of the rule.

Now, let us go back to rule 184. Define density of 1's to be

$$\rho_t = P_t(1)$$

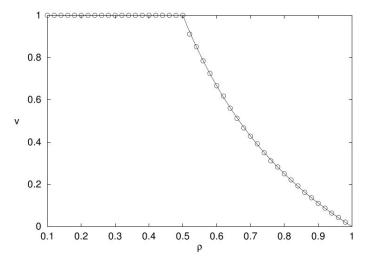
Note that $\rho_t = \rho_{t+1}$ for all t, because the number of cars does not change. We can thus drop the index t for ρ .

If we think of this rule as a system of moving cars, a car at site i can move only if site i+1 is empty. Such car has velocity 1 (one site per time step).

If, on the other hand, site i+1 is occupied, the velocity is zero. The expected value of the velocity of cars at time t is, therefore,

$$v_t = 1 \cdot \frac{P_t(10)}{P_t(1)} + 0 \cdot \frac{P_t(10)}{P_t(1)} = \frac{P_t(10)}{\rho}$$

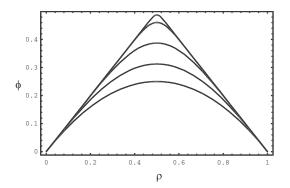
Let $v = \lim_{t \to \infty} v_t$. From computer experiments (taking large, but finite t):



Free moving phase, $\rho < 0.5$, v = 1 Jammed phase, $\rho > 0.5$, v < 1.

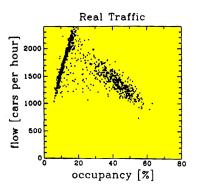
Define now the flow at time t as $\phi_t = \rho v_t$. Graph of the flow ϕ_t versus density ρ is called *fundamental diagram*.

Fundamental diagram for rule 184, for t=0, 1, 5, 50, 5000 is shown below. For large t, it becomes "tent shaped".



Real traffic data from a highway in Japan

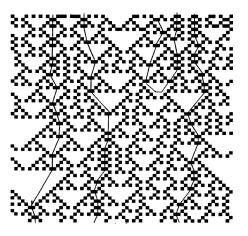
This is a fundamental diagram for a real traffic (on Japanese highway).



One can see that it is also tent-shaped, although somewhat slanted. Rule 184 reproduces the shape correctly (of course, better CA models exist).

Power laws in CA

An early example of a power law in CA has been found in rule 18 by P. Grassberger (Phys. Rev. A **28** 3666 1983). Grassberger noticed that in rule 18 number of pairs 11 decreases with time.



Plot from Grassberger's paper

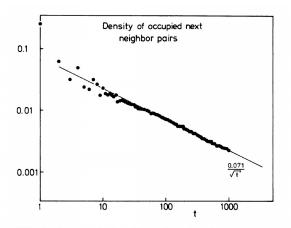


FIG. 3. Average density of occupied pairs after random start $(60\,000 \text{ sites})$. For large t, values averaged over t are plotted in order to suppress fluctuations.

$$P_t(11) \sim t^{-1/2}$$

The way to perform these calculations is to run simulation of rule 18 on a lattice with M sites, where $s_t(i)$ is the state of site i at time t. Then we simply compute the quantity

$$N_t(11) = \sum_{i=0}^{M} s_t(i)s_t(i+1)$$

We expect to obtain

$$N_t(11) = At^{-1/2}$$

where A is some constant. Taking \log of both sides of this we obtain

$$\ln N_t(11) = \ln A - \frac{1}{2} \ln t,$$

which means that if we fit linear function to data $(\ln t, \ln N_t(11))$, its slope should be -1/2.

Rule 18 is the simplest model in which power law of the type $x(t) \sim t^{-1/2}$ appears. Some other well know examples are

 Rule 184 (traffic rule discussed earlier). If one starts with symmetric Bernoulli initial configuration, then

$$P_t(11) \sim t^{-1/2}$$

 $P_t(00) \sim t^{-1/2}$

 Rule 14: If one starts with symmetric Bernoulli initial configuration, then,

$$P_0(1)=0.5$$

$$P_1(1)=0.375$$

$$P_0(1)-P_t(1)\sim t^{-1/2} \mbox{ for large } t$$

Rule 14

In rule 14, the origin of the power law is now fully understood. Local function of the elementary cellular automaton rule 14 is defined as f(0,0,1)=f(0,1,0)=f(0,1,1)=1, and $f(x_0,x_1,x_2)=0$ for all other triples $(x_0,x_1,x_2)\in\{0,1\}^3$. This is equivalent to

$$f(x_0, x_1, x_2) = x_1 + x_2 + x_1 x_0 x_2 - x_1 x_2 - x_0 x_2 - x_1 x_0.$$

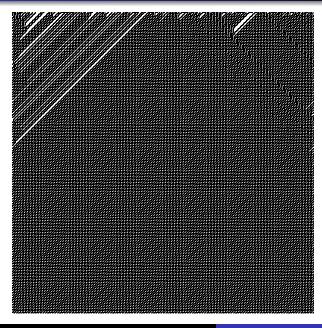
Theorem (H. Fukś and J. Haroutunian, J. of cellular automata, 4:99–110, 2009) If in the initial configuration all cells are independently in state 1 with probability 1/2 or in state 0 with probability 1/2, then the expected density of ones after t iterations of rule 14 is

$$\rho_t = \frac{1}{2} \left(1 - \frac{2t - 1}{4^t} C_{t-1} \right),$$

where C_n is the n-th Catalan number,

$$C_n = \frac{1}{n+1} {2n \choose n} = \frac{(2n)!}{n!(n+1)!}.$$

Rule 14 - spatiotemporal pattern for 500 sites



Stirling's formula for large n

$$n! \sim \sqrt{2\pi n} \left(\frac{n}{e}\right)^n$$

yields

$$C_n = \frac{1}{n+1} \frac{(2n)!}{(n!)^2} \approx \frac{1}{n+1} \frac{\sqrt{4\pi n} \left(\frac{2n}{e}\right)^{2n}}{\left(\sqrt{2\pi n} \left(\frac{n}{e}\right)^n\right)^2} = \frac{1}{n+1} \frac{2^{2n}}{\sqrt{\pi n}} = \frac{1}{n+1} \frac{4^n}{\sqrt{\pi n}}$$

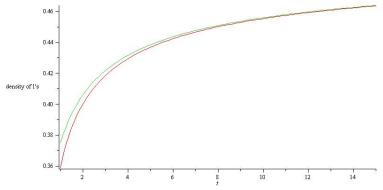
Now

$$\rho_t = \frac{1}{2} \left(1 - \frac{2t - 1}{4^t} C_{t-1} \right) \approx \frac{1}{2} \left(1 - \frac{2t - 1}{4^t} \cdot \frac{1}{t} \cdot \frac{4^{t-1}}{\sqrt{\pi(t-1)}} \right)$$

For large t we can replace t-1 by t and 2t-1 by 2t, thus

$$\rho_t \approx \frac{1}{2} \left(1 - \frac{2t}{4^t} \cdot \frac{1}{t} \cdot \frac{4^{t-1}}{\sqrt{\pi t}} \right) = \frac{1}{2} - \frac{1}{4\sqrt{\pi}} t^{-1/2}$$

Graph of the exact density (green) and the one obtained via Stirling's approximation (red) for rule 14.



For rule 14, therefore, the existence of the power law $x(t) \sim t^{-1/2}$ is fully understood. Similar explanation can be given for rule 184 (H. Fukś Phys. Rev. E, 60:197–202, 1999). No rigorous explanation is currently known for rule 18.

Other exponents in CA

Are there any rules with exponents different that -1/2? Consider rule 54, defined by

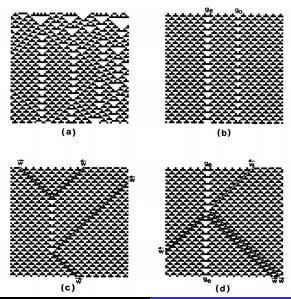
$$f(0,0,1)=f(1,0,0)=f(0,1,0)=f(1,0,1)=1, \text{ and } f(x_0,x_1,x_2)=0 \text{ for all other triples } (x_0,x_1,x_2)\in\{0,1\}^3.$$
 Alternatively,

$$f(x_0, x_1, x_2) = x_0 - 2x_0x_1 + x_2 + x_1 - x_0x_2 + 2x_0x_2x_1 - 2x_2x_1.$$

This is one of the most complex elementary CA rules. Its spatiotemporal patterns reveal existence of four types of "particles", shown in the next two slides and denoted as g_o , g_e , \overleftarrow{w} and \overrightarrow{w} .

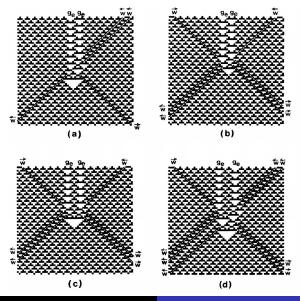
"Particles" in rule 54

Figure from N. Boccara, J. Nasser and M. Roger, Phys. Rev. A 44:2 866 (1991)



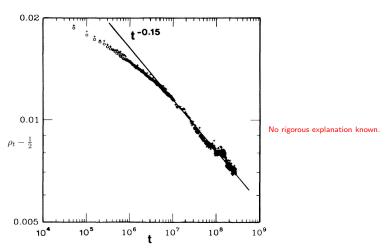
"Particles" in rule 54

Figure from N. Boccara, J. Nasser and M. Roger, Phys. Rev. A 44:2 866 (1991)



Power law in rule 54

Number of particles of type g_e as a function of time decreases as $t^{-0.15}$. Density of ones ρ_t approaches 1/2 also as $t^{-0.15}$, as shown below. (adapted from N. Boccara, J. Nasser and M. Roger, *Phys. Rev. A* 44:2 866 (1991))



Rule 54 – lattice of 1000 sites

