(a)
$$[E] = kg \frac{m^2}{S^2}$$
 $[v] = \sqrt{\frac{[E]}{[m]}} = \frac{m}{S}$

$$= \sum_{[v']} \left[\left[\frac{E'}{[w]} \right] \right] = \sqrt{\frac{[E][E]}{[w][a]}} = \sqrt{\frac{x}{E}} \left[\left[\frac{v}{v} \right] \right] = \sqrt{\frac{v}{E}} = \frac{[v']}{0}$$

where a is the moss out and a the energy out

(i)
$$\langle x_i \rangle = k_B T$$
 $x_i \langle \gamma_i \rangle$

$$H = E_k + V$$
 $\frac{\partial H}{\partial g^i} = \frac{\partial V}{\partial g^i}$ $\frac{\partial H}{\partial p_i} = \frac{\partial E_k}{\partial p_i}$

Summing over all partitles and directions

$$\sum_{i=1}^{3N} k_B T = \frac{3N}{3N} \left\langle \frac{R^2}{m} \right\rangle = \sum_{i=1}^{N} \left\langle \frac{R^2}{m} \right\rangle + \left\langle \frac{R^2}{m} \right\rangle + \left\langle \frac{R^2}{m} \right\rangle$$

$$\frac{3}{2}Nk_{B}T = \frac{3}{2}\sum_{i=1}^{N}\frac{N}{N}\left\langle \frac{p_{i}^{2}}{m}\right\rangle = \left\langle \frac{E_{K}}{N}\right\rangle$$

Where I used that $N\left(\frac{1}{N}\sum_{i=1}^{N}3\langle p_{i}^{2}\rangle\right)=N\langle E_{k}^{Post}\rangle=\langle E_{k}\rangle$

Hence
$$C_{\sigma} = \frac{1}{10} \frac{\partial U}{\partial T} = \frac{3}{2} R$$
 only when $U \approx E_{\kappa}$

n = number of moles

Numerically, we an adulate co as follows:

$$C_{\sigma} = \underbrace{\mathcal{E}}_{N} \underbrace{U(T+AT) - U(T)}_{AT}$$

$$= \underbrace{120 K \cdot k_{B}}_{AOK} \underbrace{N_{A}}_{N} \underbrace{AU}_{N} = \underbrace{12 AU}_{N} R$$

$$= \underbrace{12 \cdot (R_{A})}_{N} \underbrace{N_{A}}_{N} \underbrace{AU}_{N} = \underbrace{12 AU}_{N} R$$

We have $N=8^{\frac{3}{2}}$ 512 => $C_{U}/R=\frac{12\Delta U}{512} \approx 0.0234 \Delta U$ ΔU in natural units

$$\frac{C_{\sigma}\left(T=75\,\mathrm{K}\right)}{R} \simeq 1.55$$

$$\frac{C_{\sigma}\left(T=400\,\mathrm{K}\right)}{R} \simeq .1.51$$

As the temperature grows so it does the kinetic energy. Hence the supproximation $U_N \to K_K$ becomes more accurate (remember that $C_V = \frac{3}{2}K$ only if V = 0)

The temperature is so low that the contribution to the total energy of the system comes almost only from the potential.

We expect the panticles to be at an average obistance equal to the equilibrium of the lemnand-Jones potential, since temperature does not provide enough excitation to jump away temperature does not provide enough excitation to jump away we also expect (Cr. very different from 3 R































