(a)
$$[E] = kg \frac{m^2}{S^2}$$
 $[v] = \sqrt{\frac{[E]}{[m]}} = \frac{m}{S}$

$$= \sum_{w'} \left[w' \right] = \sqrt{\frac{\left[E' \right]}{\left[w \right] \alpha}} = \sqrt{\frac{\alpha}{\epsilon}} \left[w \right] = \sqrt{\frac{\alpha}{\epsilon}} \left[w \right] = \sqrt{\frac{\alpha}{\epsilon}} \left[w \right]$$

where a is the moss out and a the energy out

(i)
$$\langle x_i \rangle = k_B T$$
 $x_i \langle y_i \rangle$

$$H = E_{k} + V$$
 $\frac{\partial H}{\partial g^{i}} = \frac{\partial V}{\partial g^{i}}$ $\frac{\partial H}{\partial p_{i}} = \frac{\partial E_{k}}{\partial p_{i}}$

Summing over all particles and directions

$$\sum_{i=1}^{3N} k_B T = \frac{3N}{3N} \left\langle \frac{R^2}{m} \right\rangle = \sum_{i=1}^{N} \left\langle \frac{R^2}{m} \right\rangle + \left\langle \frac{R^2}{m} \right\rangle + \left\langle \frac{R^2}{m} \right\rangle$$

$$\frac{3}{2}Nk_{B}T = \frac{3}{2}\sum_{i=1}^{N}\frac{N}{N}\left\langle \frac{p_{i}^{2}}{m}\right\rangle = \left\langle \frac{E_{k}}{N}\right\rangle$$

Where I used that $N\left(\frac{1}{N}\sum_{i=1}^{N}3\langle p_{i}^{2}\rangle\right)=N\langle E_{k}^{Post}\rangle=\langle E_{k}\rangle$

Hence
$$C_{\sigma} = \frac{1}{10} \frac{\partial U}{\partial T} = \frac{3}{2} R$$
 only when $U \approx E_{\kappa}$

Numerically, we an adulate co as follows:

$$C_{\sigma} = \underbrace{\mathcal{E}}_{N} \underbrace{U(T+AT) - U(T)}_{\Delta T} = \underbrace{120K \cdot k_{B}}_{10K} \underbrace{N_{A}}_{N} \underbrace{\Delta U}_{N} = \underbrace{12\Delta U}_{N} R$$

We have $N=8^{\frac{3}{2}}$ 512 => $C_{U}/R=\frac{12\Delta U}{512} \approx 0.0234 \Delta U$ ΔU in natural units

$$\frac{C_{\sigma}\left(T=75\,\mathrm{K}\right)}{R} \simeq 1.55$$

$$\frac{C_{\sigma}\left(T=400\,\mathrm{K}\right)}{R} \simeq .1.51$$

As the temperature grows so it does the kinetic energy. Hence the supproximation $U_N \to K_K$ becomes more accurate (remember that $C_V = \frac{3}{2}K$ only if V = 0)

(i) The temperature is so low that the contribution to the total energy of the system comes almost only from the potential.

We expect the panticles to be at an average obstance equal to the equilibrium of the lennord-Jones potential, since temperature does not provide enough excitation to jump away temperature does not provide enough excitation to jump away we also expect (Cr. very different from 3 R