

# Building a SUSY model II: using superfields

Seminar on Supersymmetry and its breaking

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# Main points of the talk

- Apply the superspace formalism and show why it is useful to build SUSY theories
- Principles to construct SUSY lagrangians
- SUSY gauge theories: QED and QCD
- SUSY predictions: particles, interaction, masses, ...

# D and F terms of the superfields

Chiral superfields

$$\begin{aligned}\Phi(x, \theta, \bar{\theta}) = & \phi(x) + i\bar{\theta}\bar{\sigma}^\mu\theta\partial_\mu\phi(x) + \frac{1}{4}\theta\theta\bar{\theta}\bar{\theta}\partial_\mu\partial^\mu\phi(x) + \sqrt{2}\theta\psi(x) \\ & - \frac{i}{\sqrt{2}}\theta\theta\bar{\theta}\bar{\sigma}^\mu\partial_\mu\psi(x) + \theta\theta F(x) \\ \Phi^* = & \text{blablabla}\end{aligned}$$

A priori  $\Phi$  and  $\Phi^*$  (hence the components) are independent.

Vector superfields

$$V(x, \theta, \bar{\theta}) = \text{longstuff}$$

Note that

$$\delta_\epsilon F = \delta_\epsilon F' = \delta_\epsilon D = 0$$

$\Rightarrow$  Use F terms of chiral superfields and D terms of vector superfields to build SUSY invariant lagrangians!

# Chiral and vector superfields from chiral superfields

Two ways to combine chiral superfields

1.  $\Phi^* \Phi \longrightarrow$  Forms a new vector field because  $(\Phi^* \Phi)^* = \Phi \Phi^* = \Phi^* \Phi$
2.  $\Phi \Phi \longrightarrow$  Forms a new chiral field because
$$\bar{D}_{\dot{\alpha}} \Phi \Phi = (\bar{D}_{\dot{\alpha}} \Phi) \Phi + \Phi (\bar{D}_{\dot{\alpha}} \Phi) = 0.$$

In general any analytic function of  $\theta$  (or  $\bar{\theta}$ ) would still be a chiral superfield.

# Naive combination of the fields

Let us try, very naively, to take a product of the first type

$$\begin{aligned}\Phi(x, \theta, \bar{\theta}) &= \phi(x) + i\bar{\theta}\bar{\sigma}^{\mu}\theta\partial_{\mu}\phi(x) + \frac{1}{4}\theta\theta\bar{\theta}\bar{\theta}\partial_{\mu}\partial^{\mu}\phi(x) + \sqrt{2}\theta\psi(x) \\ &\quad - \frac{i}{\sqrt{2}}\theta\theta\bar{\theta}\bar{\sigma}^{\mu}\partial_{\mu}\psi(x) + \theta\theta F(x) \\ \Phi^{*}\Phi &= \dots\end{aligned}$$

and "select" the SUSY invariant component

$$\begin{aligned}[\Phi^{*}\Phi]_D &= \int d^2\theta d^2\bar{\theta} \Phi^{*}(x, \theta, \bar{\theta})\Phi(x, \theta, \bar{\theta}) \\ &= -\partial^{\mu}\phi^{*}\partial_{\mu}\phi + i\psi^{\dagger}\bar{\sigma}^{\mu}\partial_{\mu}\psi + F^{*}F + \text{total derivative}\end{aligned}$$

Note how mentally easy it was to derive this lagrangian, and how the Grassman formalism allows us to choose the SUSY invariant component via simple integration.

# Full Wess-Zumino model

How can we introduce products of the second type?

Remember that analytic functions of a chiral superfield is in turn a chiral superfield  $\Rightarrow$  define a function via power series

$$W(\{\phi_i\}) \equiv a_i \phi_i + \frac{1}{2} m_{ij} \phi_i \phi_j + \frac{1}{3!} y_{ijk} \phi_i \phi_j \phi_k$$

Higher order terms lead to non-renormalisable theories!

$\rightarrow$  "select"  $F$  term  $\Rightarrow$

Qua ci va tutta la parte normale di derivazione di roba

## Another point of view on what we have done

$\Omega$  has  $8+8$  particles content  $\Phi, \Phi^+$  have 2 particles content each: a boson and a left (right) chiral fermion each,  $F$  is unphysical