

last time:

GUE random matrices. Fix $N \in \mathbb{N}$

$$M = (M_{ij})_{i,j=1}^N$$

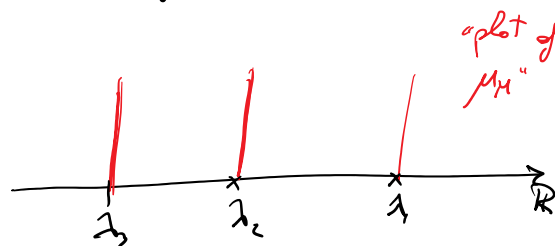
Random Hermitian Gaussian matrix

$$M_{ij} = \begin{cases} N(0, \frac{1}{2}) + i N(0, \frac{1}{2}) & i < j \\ N(0, 1) & i = j \\ \overline{M_{ji}} & i > j \end{cases}$$

Theorem (Semircircle law)

Define the empirical measure of the spectrum of M , $\lambda_1 \geq \lambda_2 \geq \dots \geq \lambda_N$
(the eigenvalues are real, by hermiticity)

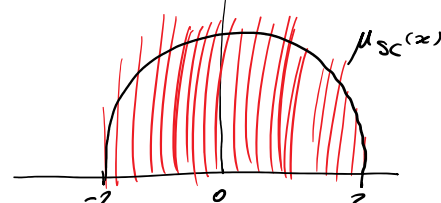
$$\mu_M(x) = \frac{1}{N} \sum_{i=1}^N \delta_{\lambda_i}(x)$$



Then, we have

$$\mu_M(\sqrt{N}x) \xrightarrow{\text{weakly a.s.}} \mu_{sc}(x) = \frac{1}{2\pi} \sqrt{(4-x^2)_+}$$

$$(z)_+ = \begin{cases} z & \text{if } z \geq 0 \\ 0 & \text{else} \end{cases}$$



Theorem (Hermite DPP)

Let $\lambda = (\lambda_1, \dots, \lambda_N)$ be the spectrum of a GUE matrix.

$$K_{N, \text{Hermite}}(x, y) = \phi_i(x) \phi_i(y) \left(= \frac{\phi_N(x) \phi_{N-1}(y) - \phi_{N-1}(x) \phi_N(y)}{x-y} \right)$$

Christoffel-Darboux identity

where $\phi_i(x)$ are the (L^2 -normalized) Hermite polynomials, which are an L^2 -orthogonal basis of $L^2(\mathbb{R}, dx)$.

This means that (informally)



$$\mathbb{P} \left[\begin{array}{l} \text{there is at least one eigenvalue} \\ \text{in the disjoint infinitesimal intervals} \\ (x_1, x_1+dx_1) \cup (x_2, x_2+dx_2) \cup \dots \cup (x_k, x_k+dx_k) \end{array} \right] = \det [K_{N, \text{Hermite}}(x_i, x_j)]_{i,j=1}^k dx_1 \dots dx_k$$

$\rightarrow 1, 0, \dots, 0$ limit

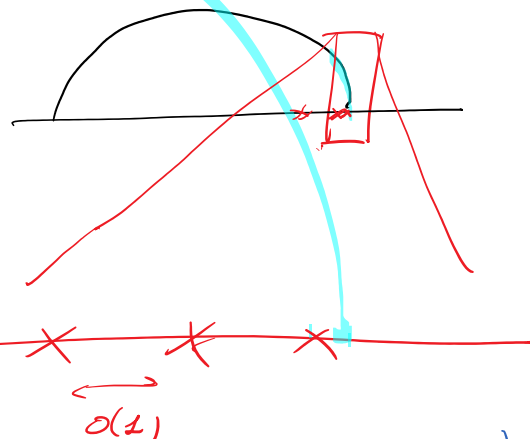
Edge Scaling limit

This can be determined from the DPP point of view.

We have

$$K_{N, \text{Hermite}} \left(z\sqrt{N} + \frac{x}{N^{1/6}}, z\sqrt{N}, \frac{y}{N^{1/6}} \right)$$

$$\xrightarrow{N \rightarrow \infty} K_{\text{Airy}}(x, y) = \int_0^\infty \text{Ai}(z+x) \text{Ai}(z+y) dz \quad \left(= \frac{\text{Ai}(x) \text{Ai}'(y) - \text{Ai}'(x) \text{Ai}(y)}{x-y} \right)$$



Dyson's Brownian Motion: We can add a dimension to the above analysis. (Idea: $N(0, \sigma^2) \mapsto B_t(0, \sigma^2)$)

Definition: Fix $N \in \mathbb{N}$

$$M(t) = (M_{ij}(t))_{i,j=1}^N$$

$$M_{ij}(t) = \begin{cases} B_t(0, \frac{1}{2}) + i \tilde{B}_t(0, \frac{1}{2}) & i < j \\ B_t(0, 1) & i = j \\ \overline{M_{ji}(t)} & i > j \end{cases}$$

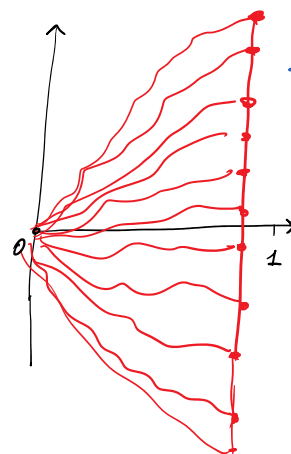
each entry i, j gets 2 indep. BM

Once again, since for each t $M(t)$ is Hermitian,

$\lambda(t) = (\lambda_1(t) \geq \dots \geq \lambda_N(t)) \in \mathbb{R}$. Then, the eigenvalues form a time dependent stochastic process, whose law at each fixed time is a rescaling of the law of the spectrum of a GUE matrix.

Picture

DBM



Theorem (Dyson's BM)

The process $\tilde{\lambda}(t) = (\lambda_1(t) \geq \dots \geq \lambda_N(t))$ evolve according to the SDE

$$d\lambda_i(t) = \sum_{j=1}^N \frac{dt}{\lambda_i(t) - \lambda_j(t)} + dB_t^{(i)}$$

$$d\lambda_i(t) = \sum_{1 \leq i < j \leq N} \frac{dt}{\lambda_i(t) - \lambda_j(t)} + dB_t$$

where $(B_t^{(i)})_{i=1, \dots, N}$ are independent BM.

Observation: We can write the above SDE as

$$d\vec{\lambda}(t) = \vec{V} \log \left(\prod_{i < j} (\lambda_i - \lambda_j) \right) dt + d\vec{B}_t$$

From the above expression one can recognize that λ is the "Dool's h -transform" of N independent Brownian motions, where $h = h(\lambda) = \prod_{i < j} (\lambda_i - \lambda_j)$ is the harmonic function encoding the infinitesimal probability of the N Brownian motions to never intersect.

Theorem $\vec{\lambda}(t)$ has the law of N Brownian motions started at 0 and conditioned to never intersect.

Corollary: The DBM has the Brownian Gibbs property

Theorem The DBM is an extended

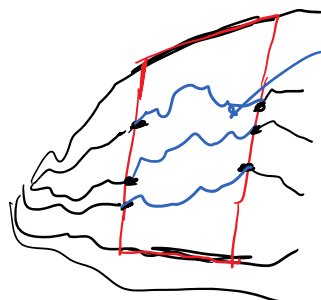
DPP. Fix t_1, \dots, t_k , then

$\mathbb{P} \left(\begin{array}{c} \text{there is a path crossing the folley} \\ \text{intervals} \end{array} \right)$

$$\{t_1, x(x_1, x_1 + dx_1), \dots, t_k, x(x_k, x_k + dx_k)\}$$

$$= \det \left[K_{N, \text{Hermite}}^{\text{ext.}}(x_i, t_i, x_j, t_j) \right]_{i,j=1}^k dx_1 \dots dx_k$$

↑
Explicit generalization of $K_{N, \text{Hermite}}$.



the DBM inside this box behaves like non-intersecting Brownian Bridges

