# Boosting ground states preparation with double-bracket quantum algorithms

PART 1: Introducing and optimizing double-bracket quantum algorithms

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# A couple of references

### Double-bracket quantum algorithms for diagonalization

Marek Gluza

This work proposes double-bracket terations as a famework for obtaining diagonation; quantum circuits. Their implementation or a quantum computer consists of interliance evolutions generated by the input Hamiltonian with diagonal evolutions which can be closen variationally. No qubit overheads or controlled-unitary operations are needed but the method is recursive which makes the circuit deful prove exponentially with the number of cerusion steps. In other lands here are deem implementations valide, the proposal includes optimization of diagonal evolution generations and of recursion step durations. Indeed, thanks to this numerical examples show that the expressive power of double-bracket actions suffices is approximate eigenstates or relevant quantum models with fever recursion steps. Compared to brute-force epitimization of uncircuitared circuits obtain-bracket interliance and on stuffer from the same transhibly limitations. Moreover, with an implementation cost lower than required for quantum phase estimation they are more suitable for near-term quantum compiling operiments. Now box books, this work opens a pathway for constructing purposally quantum algorithms based on so called double-bracket flows also for tasks different from diagonalization and thus enlarges the quantum compiling to obtain each office.

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Subjects: Quantum Physics (quant-ph)
Cite as: arXiv:2206.11772 [quant-ph]

## Check out this because:

- 1. Marek is a very nice and smart guy;
- 2. it contains math foundations of this talk;

#### Strategies for optimizing double-bracket quantum algorithms

Li Xiaoyue, Matteo Robbiati, Andrea Pasquale, Edoardo Pedicillo, Andrew Wright, Stefano Carrazza, Marek Gluza

Recently double-bracket quantum algorithms have been proposed as a way to comple circuits for approximating eigenstates. Physically, they consist of appropriately composing evolutions under an input Hamiltonian together with diagonal evolutions. Here, we present strategies to optimize the choice of the double-bracket evolutions to enhance the diagonalization efficiency. This can be done by Indiag optimal generators and durations of the evolutions. We present numerical results regarding the preparation of double-bracket iterations, both in ideal cases where the algorithm's setup provides analytical convergence quaranteres and in more heuristic cases, where we use an adaptive and variational approach to optimize the generators of the evolutions. As an example, we discuss the efficiency of these optimization strategies when considering a spin-fall in-intelliginal as the target. To propose algorithms that can be executed starting today, fully aware of the limitations of the quantum technologies at our disposal, we finally present a selection of diagonal evolution parametrizations that can be directly complete into CNDTs and single-quilit criation gates. We discuss the advantages and limitations of this compilation and propose a way to take advantage of this approach when used in senserow with other existion methods.

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Cite as: arXiv:2400.07431 [guant-ph].

## And this because:

- 1. We delve deeper into the optimization of DBQAs;
- 2. we base the optimization of the next talk on these studies.

# Ingredients

- 1. Input Hamiltonian  $\hat{H}_0$  (here 1D XXZ);
- 2. an anti-hermitian  $((i\hat{W})^{\dagger} = -i\hat{W})$  rotation generator  $\hat{W}_0 = [\hat{D}_0, \hat{H}_0]$ ;
- 3. point 2. is used to build a unitary operation:

$$\hat{R}_0 \equiv e^{s\hat{W_0}}$$

with s stepsize or flow duration;

4.  $\hat{R}_0$  applied to  $\hat{H}_0$  as double-bracket rotation:

$$\hat{H}_1 = e^{s\hat{W}_0} \hat{H}_0 e^{-s\hat{W}_0}.$$

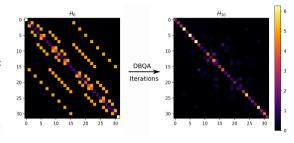


Figure 1: Ten DBQA rotations to a 5 qubit 1D XXZ.

## Nomenclature reason

The presented framework satisfies a Heisenberg equation involving two, not one, brackets:

$$\partial_s \hat{H}_0(s) = [\hat{H}_0(s), [\hat{H}_0(0), \hat{D}_0]].$$

1. We care of diagonal  $\hat{D}_k$  in  $\hat{W}_k = [\hat{D}_k, \hat{H}_k]$ . A natural choice is the canonical generator:

$$\hat{W}_k^{\mathrm{can}} = [\hat{H}_k, \hat{\Delta}(\hat{H}_k)], \qquad \text{with} \qquad \begin{cases} \hat{\Delta}(\hat{H}_k) = \mathrm{diag}(\hat{H}_k), \\ \hat{H}_k = \hat{\Delta}(\hat{H}_k) + \hat{\sigma}(\hat{H}_k), \end{cases}$$

but - as we will see - many choices can be done here.

2. **Lemma**<sup>1</sup>: given  $\sigma(\hat{H}_0(s))$  off-diagonal restriction of  $\hat{H}_0(s)$  and s small enough:

$$\partial_s \|\sigma(\hat{H}_0(s))\|_{\mathrm{HS}}^2 = -2 \, \langle \hat{\mathcal{W}}_0, [\hat{H}_0, \sigma(\hat{H}_0(s))] \rangle_{\mathrm{HS}} \,.$$

- 3. If  $\hat{D}_0 = ... = \hat{D}^*$ , as long as  $s_0 = ... = s^*$  are sufficiently small and  $\hat{D}^*$  non degenerate, the recursion converges to a fixed point  $\hat{H}_{\infty}$  with  $[\hat{H}_{\infty}, \hat{D}^*]^2$ ;
- $2. o ext{ If } \hat{D}_0$  is diagonal, then  $\hat{H}_1$  will be more diagonal than  $\hat{H}_0$ ,  $3. o ext{ the discrete DBI can converge arbitrary well,}$   $3. o ext{ analytically motivated}$  methods work in some setups but they can be slow.

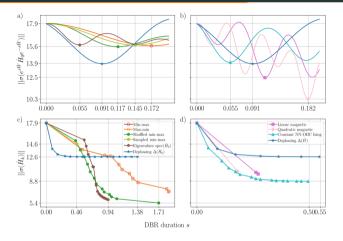
<sup>&</sup>lt;sup>1</sup>Definition and proof of Lemma 2 in Double-bracket quantum algorithms for diagonalization, M. Gluza, 2024

<sup>&</sup>lt;sup>2</sup> If you are really brave: Optimization and Dynamical Systems [Ch. 2.3], Uwe Helmke . John B. Moore, 1994





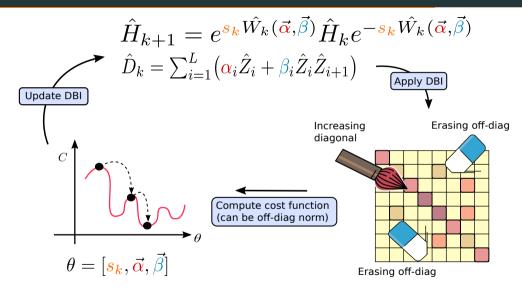
strategies







Function	Off-diagonal norm $(f_1)$	Least-squares $(f_2)$	Energy $(f_3)$	Energy fluctuation $(f_4)$
Formula	$  \sigma(\hat{H}_k)  _{ ext{HS}}$	$\frac{1}{2}  \hat{D}_k  _{\mathrm{HS}}^2 - \mathrm{Tr}(\hat{H}_k\hat{D}_k)$	$\langle \psi   \hat{H}   \psi \rangle$	$\sqrt{(\langle \psi   \hat{H}_k^2   \psi \rangle - \langle \psi   \hat{H}_k   \psi \rangle^2)}$
Utility	Diagonalization	Diagonalization	Ground state preparation	Eigenstate preparation





The heart of the compilation of DBQAs is one of the **group commutator formulas** introduced by Dawson and Nielsen in their Solovay-Kitaev algorithm:

$$\hat{V}^{\text{GC}}(\hat{A},\hat{B}) = e^{i\hat{A}}e^{i\hat{B}}e^{-i\hat{A}}e^{-i\hat{B}},$$
  $\hat{A},\hat{B}$  hermitian,

by means of which:

$$e^{-[\hat{A},\hat{B}]} = \hat{V}^{\mathrm{GC}}(\hat{A},\hat{B}) + \hat{E}^{\mathrm{GC}},$$

where

$$\hat{E}^{\text{GC}} \leq \|[\hat{A}, [\hat{A}, \hat{B}]]\| + \|[\hat{B}, [\hat{B}, \hat{A}]]\|$$

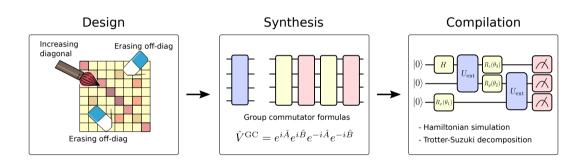
in any unitarily invariant norm (e.g. Hilbert-Schmidt norm).

## How about DBI?

$$e^{-s[\hat{D}_0,\hat{H}_0]} = e^{i\sqrt{s_0}\hat{H}_0}e^{-i\sqrt{s_0}\hat{D}_0}e^{-i\sqrt{s_0}\hat{H}_0}e^{i\sqrt{s_0}\hat{D}_0} + O(s_0^{3/2}).$$

This is compilable into a sequence of local Hamiltonian evolutions.

• Higher-order formulas can be implemented to reduce the approximation error.



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