



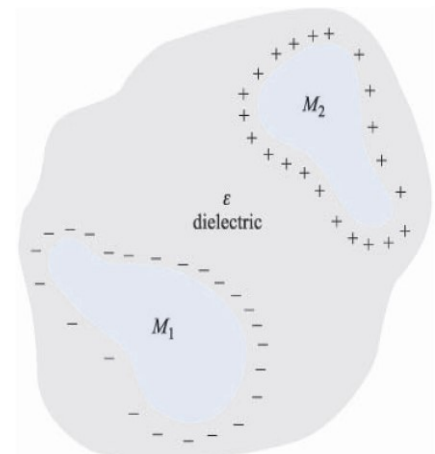
Lecture 7

1. The Electric Capacitor
2. Static Magnetic Fields

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Capacitance

- Consider two conductors embedded in a dielectric.
- Conductor M_2 carries a total positive charge Q , and M_1 carries an equal magnitude of negative charge.
- There are no other charges present, and the total charge of the system is zero.
- The charge is carried on the conductors surfaces as surface charge densities.
- The electric field is normal to the conductor surface
- The electric field is directed from M_2 ($+Q$) to M_1 ($-Q$), and M_2 is thus at higher potential.



Capacitance

- Work must be done to move a positive charge from M_1 to M_2 against the direction of the electric field.
- Let us denote the potential difference between M_2 and M_1 as $V_0 = V_2 - V_1$

The capacitance of this two-conductor system is defined as the ratio of the magnitude of the total charge on either conductor to the magnitude of the potential difference between conductors,

$$\boxed{C = \frac{Q}{V_0}} \longleftrightarrow \boxed{C = \frac{\oint_S \epsilon \mathbf{E} \cdot d\mathbf{S}}{-\int_-^+ \mathbf{E} \cdot d\mathbf{L}}}$$

Capacitance

$$C = \frac{\oint_S \epsilon \mathbf{E} \cdot d\mathbf{S}}{-\int_{-}^{+} \mathbf{E} \cdot d\mathbf{L}}$$

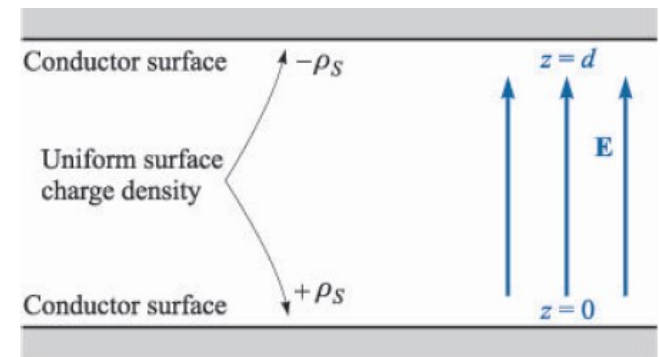
- In general terms,
 - we determine Q by a surface integral of the surface charge density over the positive conductor's area, and
 - we find V_0 by moving a unit positive charge from the negative to the positive surface i.e. against the electric field.
- The capacitance is independent of the potential and total charge, for their ratio is constant.
 - If the charge density is increased by a factor of N , the electric flux density, \mathbf{D} and electric field intensity, \mathbf{E} would be increased by a factor of N . It follows that the potential difference also would increase by the same factor.
 - Thus, the capacitance is a function only of the physical dimensions of the system of conductors and of the permittivity of the dielectric.

The Parallel Plate Capacitor

$$C = \frac{\oint_S \epsilon \mathbf{E} \cdot d\mathbf{S}}{-\int_{-}^{+} \mathbf{E} \cdot d\mathbf{L}}$$

- The conductors are two identical parallel plates with surface area S separated by a distance d ($S \gg d$).
- The plates are separated by a dielectric material and a uniform surface charge density (its magnitude is ρ_S) that exists on each conductor
- Choosing the lower conducting plate at $z = 0$, (ρ_S) and the upper one at $z = d$, ($-\rho_S$),
- We found that the electric field due to a uniformly charged infinite sheet, in free space, can be obtained from

$$\mathbf{E} = \frac{\rho_S}{2\epsilon_0} \mathbf{a}_N$$



Thus, within the capacitor,

$$\mathbf{E}_+ = \frac{\rho_S}{2\epsilon} \mathbf{a}_z \quad \mathbf{E}_- = \frac{\rho_S}{2\epsilon} \mathbf{a}_z$$

Therefore, the total \mathbf{E} within the capacitor

$$\mathbf{E} = \frac{\rho_S}{\epsilon} \mathbf{a}_z$$

The potential difference

$$V_0 = - \int_{\text{upper}}^{\text{lower}} \mathbf{E} \cdot d\mathbf{L} = - \int_d^0 \frac{\rho_S}{\epsilon} dz = \frac{\rho_S}{\epsilon} d$$

The magnitude of the total charge on each plate is

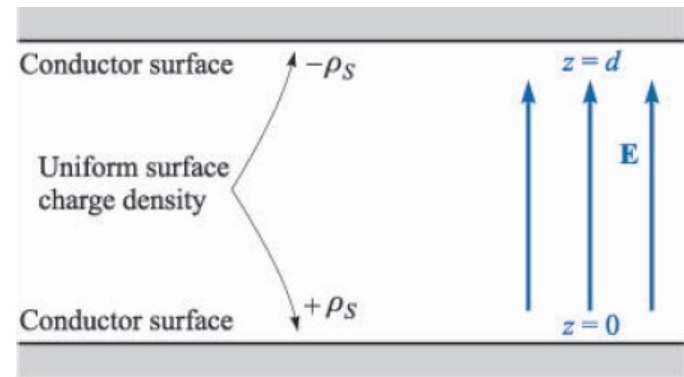
$$Q = \rho_S S$$

The capacitance of a parallel plate capacitor is

$$C = \frac{Q}{V_0} = \frac{\epsilon S}{d}$$

The Parallel Plate Capacitor

- The capacitance is independent of the potential and total charge, for their ratio is constant.
- The capacitance is a function only of the physical dimensions of the system of conductors and of the permittivity of the dielectric.



$$C = \frac{Q}{V_0} = \frac{\epsilon S}{d}$$

Example 1:

Calculate the capacitance of a parallel plate capacitor having a mica dielectric, $\epsilon_R = 6$, a plate area of 10 in^2 , and a separation of 0.01 in .

Solution. We may find that

$$S = 10 \times 0.0254^2 = 6.45 \times 10^{-3} \text{ m}^2$$

$$d = 0.01 \times 0.0254 = 2.54 \times 10^{-4} \text{ m}$$

$$C = \frac{\epsilon S}{d}$$


and therefore

$$C = \frac{6 \times 8.854 \times 10^{-12} \times 6.45 \times 10^{-3}}{2.54 \times 10^{-4}} = 1.349 \text{ nF}$$



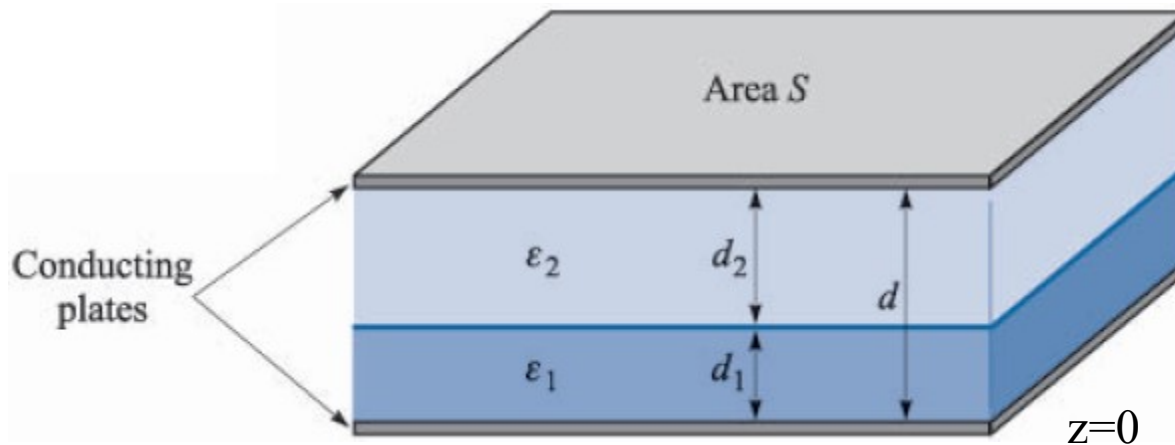
The Total Energy Stored in a Parallel Plate Capacitor

$$C = \frac{Q}{V_0} = \frac{\epsilon S}{d}$$


$$W_E = \frac{1}{2} C V_0^2$$

- The equation for the total energy stored in a parallel plate capacitor indicates that:
 - the energy stored in a capacitor with a fixed potential difference across it increases as the dielectric constant of the medium increases.

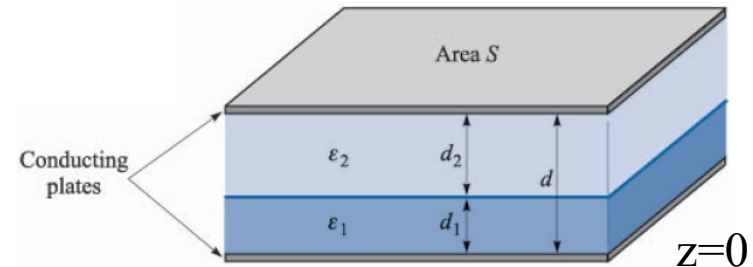
Example 2: Parallel Plate Capacitor containing two dielectrics



- Consider a parallel plate capacitor of area S and spacing d , with the usual assumption that d is small compared to the linear dimensions of the plates.
- The space between the two conducting plates is filled with two different dielectrics of permittivity ϵ_1 and ϵ_2 , and thicknesses d_1 and d_2 , respectively..

- Assume a uniform surface charge $\rho_s \text{ C/m}^2$ on lower plate and $-\rho_s \text{ C/m}^2$ on the upper plate, which leads to the uniform field,

$$\bar{E} = \begin{cases} \frac{\rho_s}{\epsilon_1} \hat{a}_z & 0 < z < d_1 \\ \frac{\rho_s}{\epsilon_2} \hat{a}_z & d_1 < z < d \end{cases}$$



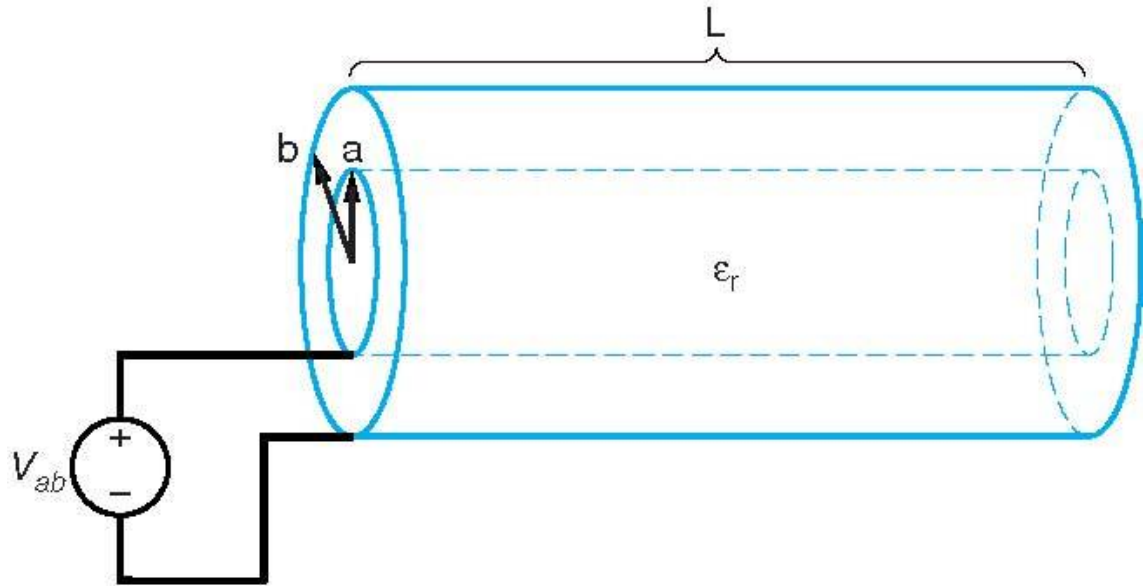
The potential difference is:

$$V_o = - \int_{upper}^{lower} \bar{E} \cdot d\bar{l} = - \int_d^{d_1} \frac{\rho_s}{\epsilon_2} dz - \int_{d_1}^0 \frac{\rho_s}{\epsilon_1} dz = \frac{\rho_s}{\epsilon_2} d_2 + \frac{\rho_s}{\epsilon_1} d_1$$

The surface area of each plate is “ S ” whose linear dimensions are much greater than the separation “ d ”. since $Q = \rho_s S$, then

$$C = \frac{Q}{V_o} = \frac{S}{\frac{d_2}{\epsilon_2} + \frac{d_1}{\epsilon_1}} \rightarrow C = \frac{1}{\frac{d_1}{\epsilon_1 S} + \frac{d_2}{\epsilon_2 S}} \quad \left\{ \begin{array}{l} C_1 = \frac{\epsilon_1 S}{d_1} \\ C_2 = \frac{\epsilon_2 S}{d_2} \end{array} \right.$$

Example 3: The Coaxial capacitor



- Consider a coaxial cable or coaxial capacitor of inner radius a , outer radius b , and length L .
- The inner surface carries a uniform charge per unit length $+\rho_l$ C/m and the outer one carries $-\rho_l$ C/m.
 - Find the Capacitance and the total energy stored in the capacitor



The electric field in the medium between the two cylindrical surfaces is given by:

$$\overline{E} = \frac{\rho_l}{2\pi\epsilon r} \hat{a}_r$$

The potential difference

$$V_o = -\int_b^a \frac{\rho_l}{2\pi\epsilon r} dr = \frac{\rho_l}{2\pi\epsilon} \ln(b/a)$$

Thus,

$$C = \frac{Q}{V_o} = \frac{\rho_l L}{\frac{\rho_l}{2\pi\epsilon} \ln(b/a)}$$

The capacitance of a Coaxial capacitor is

$$C = \frac{2\pi\epsilon L}{\ln(b/a)}$$

and

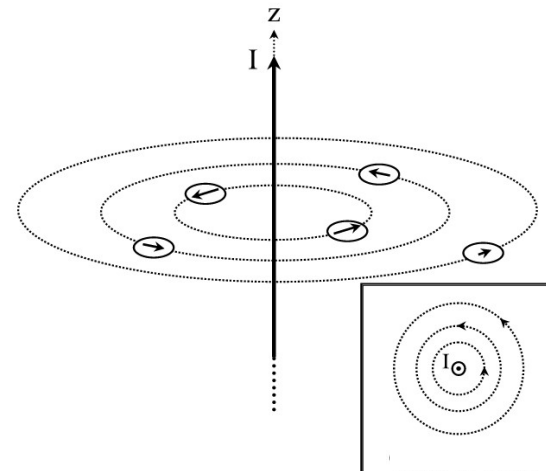
$$W_E = \frac{1}{2} C V_o^2$$

Magnetostatics – Oersted's Experiment

Magnetism and electricity were considered distinct phenomena until 1820 when *Hans Christian Oersted* (1777-1851) conducted an experiment that showed a compass needle deflecting when in proximity to a current carrying wire.

It was observed that moving away from the source of current, the field grows weaker.

Oersted's experiment



Magnetostatics – Biot-Savart's Law

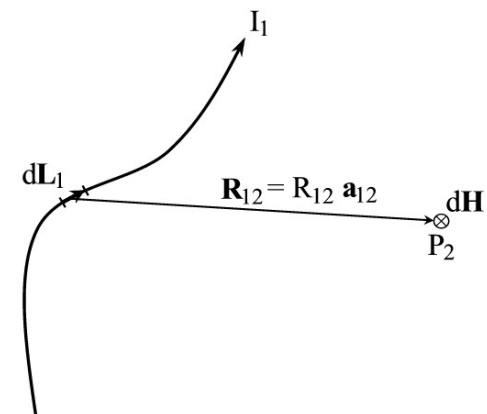
Jean Baptiste Biot (1774-1862) and Felix Savart (1791-1841) arrived at a mathematical relation between the magnetic field intensity and current.

The *Law of Biot-Savart* is

$$\boxed{d\mathbf{H}_2 = \frac{I_1 d\mathbf{L}_1 \times \mathbf{a}_{12}}{4\pi R_{12}^2}} \quad (\text{A/m})$$

To get the total magnetic field resulting from a current, one can sum the contributions from each segment by integrating

$$\boxed{\mathbf{H} = \int \frac{Id\mathbf{L} \times \mathbf{a}_R}{4\pi R^2}} \quad (\text{A/m})$$



$$d\mathbf{E}_2 = \frac{dQ_1 \mathbf{a}_{12}}{4\pi\epsilon R_{12}^2}$$

Note: Biot-Savart's law is analogous to Coulomb's law for the electric field resulting from a differential charge element.

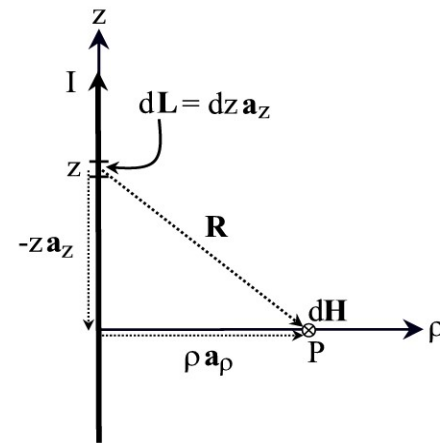
Magnetostatics – An Infinite Line current

Example 4: Consider an infinite length line along the z-axis conducting constant current I in the $+\mathbf{a}_z$ direction. We want to find the magnetic field at a point P, at distance ρ from the line along the ρ axis.

So we consider a differential length, $d\mathbf{L}$, at distance z from the origin.

The Biot-Savart Law

$$\mathbf{H} = \int \frac{I d\mathbf{L} \times \mathbf{a}_R}{4\pi R^2}$$



$I d\mathbf{L}$ is simply $I dz \mathbf{a}_z$, and the vector from the source to the test point is

$$R \mathbf{a}_R = -z \mathbf{a}_z + \rho \mathbf{a}_\rho$$

The Biot-Savart Law becomes

$$\mathbf{H} = \int_{-\infty}^{\infty} \frac{I dz \mathbf{a}_z \times (-z \mathbf{a}_z + \rho \mathbf{a}_\rho)}{4\pi (z^2 + \rho^2)^{3/2}}.$$

Pulling the constants to the left of the integral and realizing that $\mathbf{a}_z \times \mathbf{a}_z = 0$ and $\mathbf{a}_z \times \mathbf{a}_\rho = \mathbf{a}_\phi$, we have

$$\mathbf{H} = \frac{I\rho}{4\pi} \int_{-\infty}^{\infty} \frac{dz}{(z^2 + \rho^2)^{3/2}} \mathbf{a}_\phi$$

The integral can be evaluated using the following formula:

$$\int \frac{dx}{(x^2 + a^2)^{3/2}} = \frac{x}{a^2 \sqrt{x^2 + a^2}}$$

$$\int_{-\tau}^{\tau} \frac{dz}{(\rho^2 + z^2)^{3/2}} = \frac{z}{\rho^2 \sqrt{\rho^2 + z^2}} \Big|_{-\tau}^{\tau} = \left(\frac{\tau}{\rho^2 \sqrt{\rho^2 + \tau^2}} \right) - \left(\frac{-\tau}{\rho^2 \sqrt{\rho^2 + (-\tau)^2}} \right)$$

$$\left(\frac{\tau}{\rho^2 \sqrt{\rho^2 + \tau^2}} \right) + \left(\frac{\tau}{\rho^2 \sqrt{\rho^2 + \tau^2}} \right) = \frac{2\tau}{\rho^2 \sqrt{\rho^2 + \tau^2}}$$

When the limit $\tau \rightarrow \infty$

$$\int_{-\infty}^{\infty} \frac{dz}{(z^2 + \rho^2)^{3/2}} = \frac{2}{\rho^2 \sqrt{\left(\frac{\rho}{\tau}\right)^2 + 1}} = \frac{2}{\rho^2 \sqrt{\left(\frac{\rho}{\infty}\right)^2 + 1}} = \frac{2}{\rho^2}$$

$$\mathbf{H} = \frac{I \rho \mathbf{a}_\phi}{4\pi} \int_{-\infty}^{\infty} \frac{dz}{(z^2 + \rho^2)^{3/2}}$$

Substituting for $\int_{-\infty}^{\infty} \frac{dz}{(z^2 + \rho^2)^{3/2}} = \frac{2}{\rho^2}$

We find the magnetic field intensity resulting from an infinite length line of current is

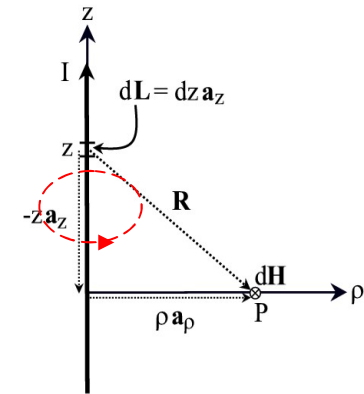
$$\mathbf{H} = \frac{I \mathbf{a}_\phi}{2\pi\rho}$$



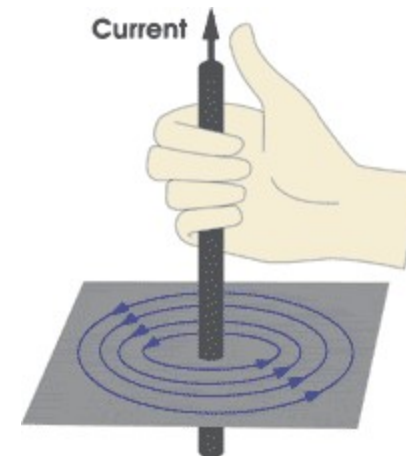
$$\mathbf{H} \propto \frac{1}{\rho} \mathbf{a}_\phi$$

Direction: The direction of the magnetic field can be found using the right hand rule.

Magnitude: The magnitude of the magnetic field is inversely proportional to radial distance.



An infinite length line of current



Magnetostatics – A Ring of Current

Example 5: Let us now consider a ring of current with radius “a” lying in the x-y plane with a current I in the $+\mathbf{a}_z$ direction.

The objective is to find an expression for the field at an arbitrary point P, at a height h on the z-axis.

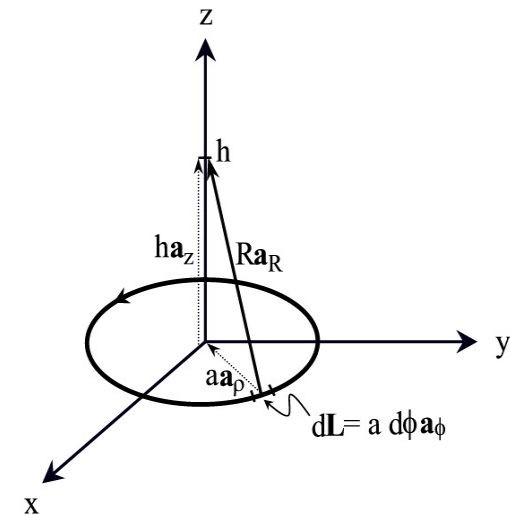
The *Biot-Savart Law*
$$\mathbf{H} = \int \frac{I d\mathbf{L} \times \mathbf{a}_R}{4\pi R^2}$$

The differential segment $dL = a d\phi \mathbf{a}_\phi$

The vector drawn from the source to the test point is

$$\mathbf{R} = R\mathbf{a}_R = h\mathbf{a}_z - a\mathbf{a}_\rho$$

Magnitude: $R = \sqrt{h^2 + a^2}$, Unit Vector: $\mathbf{a}_R = (h\mathbf{a}_z - a\mathbf{a}_\rho)/R$



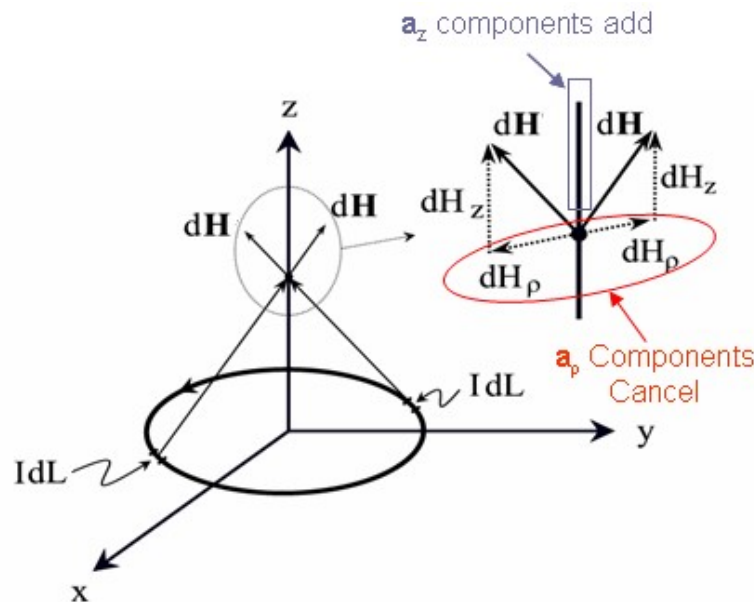
The Biot-Savart Law can be written as

$$\mathbf{H} = \int_{\phi=0}^{2\pi} \frac{I a d\phi \mathbf{a}_\phi \times (h\mathbf{a}_z - a\mathbf{a}_\rho)}{4\pi (h^2 + a^2)^{3/2}} = \frac{Ia}{4\pi} \int_{\phi=0}^{2\pi} \frac{d\phi \mathbf{a}_\phi \times (h\mathbf{a}_z - a\mathbf{a}_\rho)}{(h^2 + a^2)^{3/2}}$$

$$\mathbf{H} = \frac{Ia}{4\pi} \int_{\phi=0}^{2\pi} \frac{d\phi \mathbf{a}_\phi \times (h\mathbf{a}_z - a\mathbf{a}_\rho)}{(h^2 + a^2)^{3/2}} = \frac{Ia}{4\pi} \int_{\phi=0}^{2\pi} \frac{d\phi (h\mathbf{a}_\rho + a\mathbf{a}_z)}{(h^2 + a^2)^{3/2}}$$

A particular differential current element will give a field with an \mathbf{a}_ρ component (from $\mathbf{a}_\phi \times \mathbf{a}_z$) and an \mathbf{a}_z component (from $\mathbf{a}_\phi \times -\mathbf{a}_\rho$).

We can further simplify the previous expression for \mathbf{H} by considering the symmetry of the problem

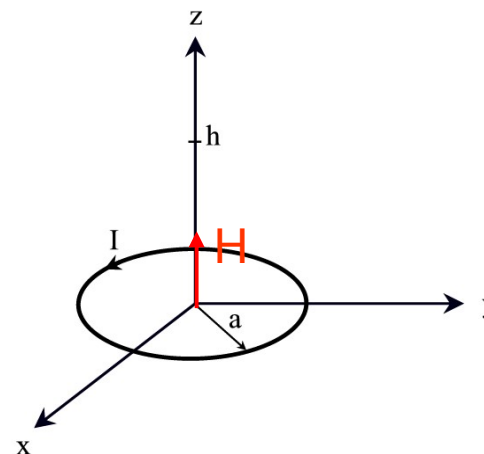


- Considering the field from a differential current element on the opposite side of the ring, it is apparent that the radial components cancel while the \mathbf{a}_z components add.

$$\mathbf{H} = \frac{Ia^2 \mathbf{a}_z}{4\pi (h^2 + a^2)^{3/2}} \int_0^{2\pi} d\phi \quad \longrightarrow \quad \boxed{\mathbf{H} = \frac{Ia^2}{2(h^2 + a^2)^{3/2}} \mathbf{a}_z}$$

At $h = 0$, the center of the loop, this equation reduces to

$$\boxed{\mathbf{H} = \frac{I}{2a} \mathbf{a}_z}$$

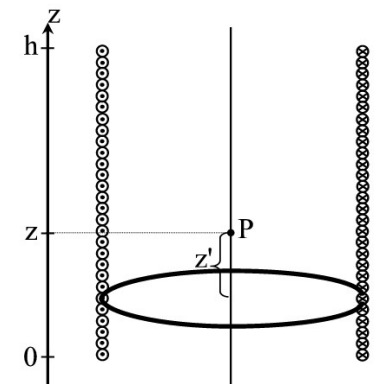
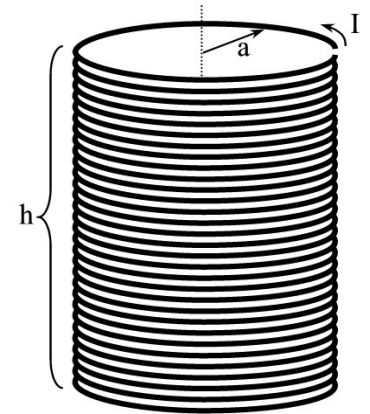


Magnetic Field Intensity for a Solenoid

- *Solenoids* (inductors) are many turns of insulated wire coiled in the shape of a cylinder.
- Suppose the solenoid has a length h , a radius a , and is made up of N turns of current carrying wire.
- For tight wrapping, we can consider the solenoid to be made up of N loops of current carrying wire.
- To find the magnetic field intensity from a single loop at a point P along the axis of the solenoid, we have
 - The differential amount of field resulting from a differential amount of current in one loop is given by

$$d\mathbf{H}_P = \frac{dIa^2}{2(z'^2 + a^2)^{3/2}} \mathbf{a}_z$$

A solenoid



$$\mathbf{H} = \frac{Ia^2}{2(h^2 + a^2)^{3/2}} \mathbf{a}_z$$

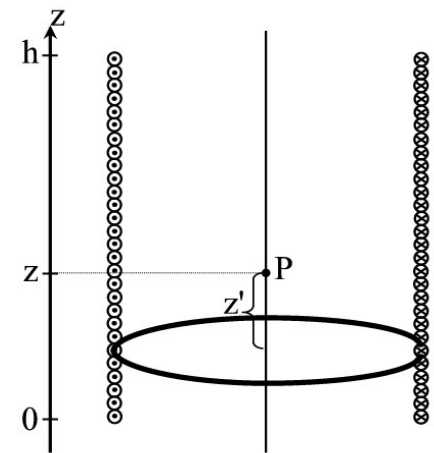
- The differential current element in one loop can be written as a function of the number of loops and the length of the solenoid as follows

$$dI = \frac{N}{h} Idz'$$

Fixing the point P where the field is desired, z' will range from $-z$ to $h-z$, or

$$\mathbf{H} = \int_{-z}^{h-z} \frac{NIa^2 dz'}{2h(z'^2 + a^2)^{3/2}} \mathbf{a}_z = \frac{NIa^2}{2h} \int_{-z}^{h-z} \frac{dz'}{(z'^2 + a^2)^{3/2}} \mathbf{a}_z.$$

$$\int \frac{dx}{(x^2 + a^2)^{3/2}} = \frac{x}{a^2 \sqrt{x^2 + a^2}}$$



$$\mathbf{H} = \frac{NI}{2h} \left[\frac{h-z}{\sqrt{(h-z)^2 + a^2}} + \frac{z}{\sqrt{z^2 + a^2}} \right] \mathbf{a}_z$$

At the center of the solenoid ($z = h/2$), with the assumption that the length is considerably larger than the loop radius ($h \gg a$), the equation reduces to

$$\mathbf{H} = \frac{NI}{h} \mathbf{a}_z$$



Announcements

- Midterm has been graded (Avg. 69%), check Webcourses.
 - [..\..\New folder\Midterm_Solution__EEL3470_Fall20.pdf](#)
- Homework 5 is assigned.