Equations

Schrödinger:

$$i\hbar\frac{\partial\left|\psi(t)\right\rangle}{\partial t}=H|\psi t\rangle \qquad \mathcal{L}=i\psi^{*}\dot{\psi}-\frac{1}{2m}\vec{\nabla}\psi^{*}\cdot\vec{\nabla}\psi$$

Klein-Gordon (Real Scalar Field):

$$(\partial \cdot \partial + m^2)\phi = 0 \qquad \mathcal{L} = \frac{1}{2} \left[\partial_\mu \phi \partial^\mu \phi - m^2 \phi^2 \right] = \frac{1}{2} \left[\dot{\phi}^2 - \left(\vec{\nabla} \phi \right)^2 - m^2 \phi^2 \right]$$

Dirac (Complex Scalar Field):

$$(i\partial \!\!\!/ - m)\varphi = 0$$
 $\overline{\varphi}(i\partial \!\!\!/ + m) = 0$ $\mathscr{L} = \partial_{\mu}\varphi\partial^{\mu}\varphi^* - m^2\varphi\varphi^*$

Pauli Matrices

$$\sigma^{\mu} = (\mathbf{1}, \vec{\boldsymbol{\sigma}}) \qquad \overline{\sigma}^{\mu} = (\mathbf{1}, -\vec{\boldsymbol{\sigma}})$$

$$\sigma^{1} = \sigma_{x} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \qquad \sigma^{2} = \sigma_{y} = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \qquad \sigma^{3} = \sigma_{z} = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \qquad \vec{\boldsymbol{\Sigma}} = \begin{pmatrix} \vec{\boldsymbol{\sigma}} & 0 \\ 0 & \vec{\boldsymbol{\sigma}} \end{pmatrix}$$

$$\{\sigma^{i}, \sigma^{j}\} = 2\delta^{ij} \qquad [\sigma^{i}, \sigma^{j}] = 2i\epsilon^{ijk}\sigma^{k} \qquad \sigma^{i}\sigma^{j} = \delta^{ij} + i\epsilon^{ijk}\sigma^{k} \qquad \sigma^{2}\sigma^{i}\sigma^{2} = -(\sigma^{i})^{*}$$

$$\vec{\boldsymbol{\sigma}} \cdot \vec{\mathbf{p}} = \begin{pmatrix} p_{z} & p_{x} - ip_{y} \\ p_{x} + ip_{y} & -p_{z} \end{pmatrix} \qquad (\vec{\boldsymbol{\sigma}} \cdot \vec{\mathbf{p}})^{2} = |\vec{\mathbf{p}}|^{2} \qquad (p \cdot \sigma)(p \cdot \vec{\sigma}) = p^{2}$$

$$p \cdot \vec{\boldsymbol{\sigma}} = \begin{pmatrix} p^{0} + p^{3} & p^{1} - ip^{2} \\ p^{1} + ip^{2} & p^{0} - p^{3} \end{pmatrix} \qquad p \cdot \sigma = \begin{pmatrix} p^{0} - p^{3} & -(p^{1} - ip^{2}) \\ -(p^{1} + ip^{2}) & p^{0} + p^{3} \end{pmatrix}$$

$$\sqrt{p \cdot \sigma} = \frac{E + m - \vec{\boldsymbol{\sigma}} \cdot \vec{\mathbf{p}}}{\sqrt{2(E + m)}} \qquad \sqrt{p \cdot \vec{\sigma}} = \frac{E + m + \vec{\boldsymbol{\sigma}} \cdot \vec{\mathbf{p}}}{\sqrt{2(E + m)}}$$

Dirac γ-Matrices

$$\begin{split} \{\gamma^{\mu},\gamma^{\nu}\} &= 2g^{\mu\nu} \quad \gamma_5 = +i\gamma^0\gamma^1\gamma^2\gamma^3 = -i\gamma_0\gamma_1\gamma_2\gamma_3 = -\frac{i}{4!}\epsilon^{\alpha\beta\gamma\delta}\gamma^{\alpha}\gamma^{\beta}\gamma^{\gamma}\gamma^{\delta} \\ \sigma^{\mu\nu} &= \frac{i}{2}[\gamma^{\mu},\gamma^{\nu}] \implies [\gamma^{\mu},\gamma^{\nu}] = -2i\sigma^{\mu\nu} \quad \{\gamma_5,\gamma^{\mu}\} = 0 \quad \gamma^0(\gamma^{\mu})^{\dagger}\gamma^0 = \gamma^{\mu} \\ \left(\gamma^0\right)^{\dagger} &= \gamma^0 \quad \left(\gamma^0\right)^2 = 1 \quad \left(\gamma^k\right)^{\dagger} = -\gamma^k \quad \left(\gamma^k\right)^2 = -1 \quad \left(\gamma^5\right)^{\dagger} = \gamma^5 \quad \left(\gamma^5\right)^2 = 1 \\ \overline{\Gamma} &= \gamma^0\Gamma^{\dagger}\gamma^0 \quad \overline{\gamma_5} = -\gamma_5 \quad \overline{\gamma^{\mu}} = \gamma^{\mu} \quad \overline{\gamma^{\mu}\gamma_5} = \gamma^{\mu}\gamma_5 \quad \overline{\sigma^{\mu\nu}} = \sigma^{\mu\nu} \end{split}$$

$$\begin{split} & \not A = A^{\mu} \gamma_{\mu} \qquad \gamma^{\mu} \gamma_{\mu} = 4 \qquad \gamma^{\mu} \gamma^{\nu} \gamma_{\mu} = -2 \gamma^{\nu} \qquad \gamma^{\mu} \gamma^{\nu} \gamma^{\rho} \gamma_{\mu} = 4 g^{\nu\rho} \\ & \gamma^{\mu} \gamma^{\nu} \gamma^{\rho} \gamma^{\sigma} \gamma_{\mu} = -2 \gamma^{\sigma} \gamma^{\rho} \gamma^{\nu} \qquad \gamma^{\mu} \gamma^{\nu} \gamma^{\rho} \gamma^{\sigma} \gamma^{\pi} \gamma_{\mu} = 2 [\gamma^{\pi} \gamma^{\nu} \gamma^{\rho} \gamma^{\sigma} + \gamma^{\sigma} \gamma^{\rho} \gamma^{\nu} \gamma^{\pi}] \\ & \gamma^{\mu} \gamma^{\nu} \gamma^{\rho} = i \epsilon^{\mu \nu \rho \lambda} \gamma_{\lambda} \gamma_{5} + g^{\mu \nu} \gamma^{\rho} - g^{\mu \rho} \gamma^{\nu} + g^{\nu \rho} \gamma^{\mu} \end{split}$$

Spin, Helicity and Chirality

$$\begin{split} \vec{\mathbf{S}} &= \frac{1}{2}\vec{\mathbf{\Sigma}} \qquad \tilde{h} = \frac{\vec{\mathbf{S}} \cdot \vec{\mathbf{p}}}{|\vec{\mathbf{p}}|} = \frac{\vec{\mathbf{\Sigma}} \cdot \vec{\mathbf{p}}}{2|\vec{\mathbf{p}}|} \qquad \tilde{h} = 2h = \frac{\vec{\mathbf{\Sigma}} \cdot \vec{\mathbf{p}}}{|\vec{\mathbf{p}}|} \qquad P_L = \frac{1-h}{2} \qquad P_R = \frac{1+h}{2} \\ \gamma_L &= \frac{1-\gamma_5}{2} \qquad \gamma_R = \frac{1+\gamma_5}{2} \qquad \gamma_{L,R}^2 = \gamma_{L,R} \qquad \gamma_{L,R}\gamma_{R,L} = 0 \qquad \gamma_L + \gamma_R = 1 \\ \gamma^\mu \gamma_{L,R} &= \gamma_{R,L} \gamma^\mu \qquad \gamma_5 \gamma_L = \gamma_L \gamma_5 = -\gamma_L \qquad \gamma_5 \gamma_R = \gamma_R \gamma_5 = \gamma_R \qquad \gamma_{L,R}^{\dagger} = \gamma_{L,R} \\ \overline{\gamma^\mu \gamma_{L,R}} &= \gamma^\mu \gamma_{L,R} \qquad \overline{\gamma_5 \gamma_L} = -\gamma_R \qquad \overline{\gamma_5 \gamma_R} = \gamma_L \qquad \overline{\gamma_{L,R}} = \gamma_0 \gamma_{L,R}^{\dagger} \gamma_0 = \gamma_{R,L} \end{split}$$

$$\begin{split} u_{\uparrow} &= N \begin{pmatrix} c \\ se^{i\phi} \\ kc \\ kse^{i\phi} \end{pmatrix} \approx N \begin{pmatrix} c \\ se^{i\phi} \\ c \\ se^{i\phi} \end{pmatrix} \qquad u_{\downarrow} = N \begin{pmatrix} c \\ ks \\ -kce^{i\phi} \end{pmatrix} \approx N \begin{pmatrix} c \\ s \\ -ce^{i\phi} \\ s \\ -ce^{i\phi} \end{pmatrix} \\ v_{\uparrow} &= N \begin{pmatrix} ks \\ -kce^{i\phi} \\ -s \\ ce^{i\phi} \end{pmatrix} \approx N \begin{pmatrix} c \\ -ce^{i\phi} \\ -s \\ ce^{i\phi} \end{pmatrix} \qquad v_{\downarrow} = N \begin{pmatrix} kc \\ kse^{i\phi} \\ c \\ se^{i\phi} \end{pmatrix} \approx N \begin{pmatrix} c \\ se^{i\phi} \\ c \\ se^{i\phi} \end{pmatrix} \\ \vec{\mathbf{p}} &= p(\sin\theta\cos\phi, \sin\theta\sin\phi, \cos\theta) \\ N &= \sqrt{E+m} \approx \sqrt{E} \qquad k = \frac{p}{E+m} \approx 1 \qquad s = \sin\left(\frac{\theta}{2}\right) \qquad c = \cos\left(\frac{\theta}{2}\right) \end{split}$$

Traces

$$\begin{split} \operatorname{Tr}[1] &= 4 & \operatorname{Tr}[\gamma^{\mu}\gamma^{\nu} \dots \gamma^{\rho}] = 0 & \operatorname{Tr}[\gamma^{\mu}\gamma^{\nu}] = 4g^{\mu\nu} \\ \operatorname{Odd} \operatorname{Number} & \operatorname{Tr}[\gamma_{5}] = 0 & \operatorname{Tr}[\gamma_{5}\gamma^{\mu}\gamma^{\nu} \dots \gamma^{\rho}] = 0 & \operatorname{Tr}[\gamma_{5}\gamma^{\mu}\gamma^{\nu}] = 0 \\ \operatorname{Odd} \operatorname{Number} & \operatorname{Tr}[\gamma_{5}\gamma^{\mu}\gamma^{\nu}\gamma^{\rho}\gamma^{\sigma}] = 0 & \operatorname{Tr}[\gamma_{5}\gamma^{\mu}\gamma^{\nu}] = 0 \\ \operatorname{Tr}[\gamma^{\mu}\gamma^{\nu}\gamma^{\rho}\gamma^{\sigma}] &= 4(g^{\mu\nu}g^{\rho\sigma} - g^{\mu\rho}g^{\nu\sigma} + g^{\mu\sigma}g^{\nu\rho}) & \operatorname{Tr}[\gamma_{5}\gamma^{\mu}\gamma^{\nu}\gamma^{\rho}\gamma^{\sigma}] = -4i\epsilon^{\mu\nu\rho\sigma} \\ \operatorname{Tr}[\gamma^{\mu}\gamma^{\nu}\gamma^{\rho}\gamma^{\sigma}\gamma_{L}] &= 2(g^{\mu\nu}g^{\rho\sigma} - g^{\mu\rho}g^{\nu\sigma} + g^{\mu\sigma}g^{\nu\rho} + i\epsilon^{\mu\nu\rho\sigma}) \\ \operatorname{Tr}[\gamma^{\mu}\gamma^{\nu}\gamma^{\rho}\gamma^{\sigma}\gamma_{R}] &= 2(g^{\mu\nu}g^{\rho\sigma} - g^{\mu\rho}g^{\nu\sigma} + g^{\mu\sigma}g^{\nu\rho} - i\epsilon^{\mu\nu\rho\sigma}) \\ \operatorname{Tr}[(g_{a} + m_{a})\gamma^{\mu}(g_{b} + m_{b})\gamma^{\nu}] &= 4\left[p_{a}^{\mu}p_{b}^{\nu} + p_{b}^{\mu}p_{a}^{\nu} + (m_{a}m_{b} - p_{a} \cdot p_{b})g^{\mu\nu}\right] \\ \operatorname{Tr}[\phi\gamma^{\mu}\beta\gamma^{\nu}\gamma_{L}] \operatorname{Tr}[\phi\gamma_{\mu}\delta\gamma_{\nu}\gamma_{L}] &= \operatorname{Tr}[\phi\gamma^{\mu}\beta\gamma^{\nu}\gamma_{L}] \operatorname{Tr}[\phi\gamma_{\mu}\delta\gamma_{\nu}\gamma_{L}] &= 16(a \cdot d)(b \cdot c) \end{split}$$

Spinors (Fermions)

$$\psi = ue^{+i(\vec{\mathbf{p}}\cdot\vec{\mathbf{x}}-E\cdot t)} = ve^{-i(\vec{\mathbf{p}}\cdot\vec{\mathbf{x}}-E\cdot t)} \qquad \overline{\psi} = \psi^{\dagger}\gamma^{0}$$

$$(\not{p}-m)u = \overline{u}(\not{p}-m) = 0 \qquad (\not{p}+m)v = \overline{v}(\not{p}+m) = 0$$

$$\overline{u}'(p)u^{s}(p) = +2m\delta^{rs} \qquad \overline{u}'(p)v^{s}(p) = 0 \qquad u^{r\dagger}(p)u^{s}(p) = 2E\delta^{rs} \qquad \sum_{s=1,2} u^{s}(p)\overline{u}^{s}(p) = \not{p}+m$$

$$\overline{v}'(p)v^{s}(p) = -2m\delta^{rs} \qquad \overline{v}'(p)u^{s}(p) = 0 \qquad v^{r\dagger}(p)v^{s}(p) = 2E\delta^{rs} \qquad \sum_{s=1,2} v^{s}(p)\overline{v}^{s}(p) = \not{p}-m$$

Polarization Vectors (Bosons)

External massless & Massive:

$$\sum_{\lambda=1}^{2} \epsilon^{\mu}(k,\lambda) \cdot \epsilon^{*\nu}(k,\lambda) = -g^{\mu\nu} \qquad \sum_{\lambda=1}^{3} \epsilon_{\mu}(q,\lambda) \epsilon_{\nu}^{*}(q,\lambda) = -g_{\mu\nu} + \frac{q_{\mu}q_{\nu}}{M^{2}}$$
Virtual massless:

$$\epsilon(k,\lambda) \cdot \epsilon^*(k,\lambda') = g^{\lambda \lambda'} \qquad \sum_{\lambda} g^{\lambda \lambda} \epsilon^{\mu}(k,\lambda) \cdot \epsilon^{*\nu}(k,\lambda) = g^{\mu \nu}$$

Quantization : Real Scalar Field (Bosons)

The "+" corresponds to positive frequency plane waves $e^{-ik \cdot x}$:

$$\begin{split} \mathcal{L} &= \frac{1}{2} \partial^{\mu} \phi \partial_{\mu} \phi - \frac{1}{2} m^{2} \phi \phi \qquad \pi = \dot{\phi} \qquad \left[a^{\dagger}(k,\lambda), a(k',\lambda') \right] = g^{\lambda \lambda'} \, \tilde{\delta}(k-k') \\ \left[\phi(\vec{\mathbf{x}},t), \phi(\vec{\mathbf{y}},t) \right] &= \left[\Pi(\vec{\mathbf{x}},t), \Pi(\vec{\mathbf{y}},t) \right] = 0 \qquad \left[\phi(\vec{\mathbf{x}},t), \Pi(\vec{\mathbf{y}},t) \right] = i \delta^{3} (\vec{\mathbf{x}} - \vec{\mathbf{y}}) \\ \phi(x) &= \int \widetilde{\mathrm{d}} k \left[a(k) e^{-ik \cdot x} + a^{\dagger}(k) e^{+ik \cdot x} \right] \\ |p\rangle &= a^{\dagger}(p) |0\rangle \qquad |p_{1}, p_{2}\rangle = a^{\dagger}(p_{2}) a^{\dagger}(p_{1}) |0\rangle = |p_{2}, p_{1}\rangle \end{split}$$

Quantization : Complex Scalar Field (Bosons)

$$\begin{split} \mathcal{L} &= : \partial^{\mu} \varphi^{\dagger} \partial_{\mu} \varphi - m^{2} \varphi^{\dagger} \varphi \colon \qquad \pi = \dot{\varphi}^{\dagger} \qquad \pi^{\dagger} = \dot{\varphi} \qquad \left[a_{\pm}(k), a_{\pm}^{\dagger}(k') \right] = \tilde{\delta}(k - k') \\ & \left[\varphi(\vec{\mathbf{x}}, t), \pi(\vec{\mathbf{y}}, t) \right] = \left[\varphi^{\dagger}(\vec{\mathbf{x}}, t), \pi^{\dagger}(\vec{\mathbf{y}}, t) \right] = i \, \delta^{3}(\vec{\mathbf{x}} - \vec{\mathbf{y}}) \\ & \varphi(x) = \varphi^{+}(x) + \varphi^{-}(x) = \int \widetilde{\mathrm{d}}k \Big[a_{+}(k) e^{-ik \cdot x} + a_{-}^{\dagger}(k) e^{+ik \cdot x} \Big] \\ & \varphi^{\dagger}(x) = \varphi^{+\dagger}(x) + \varphi^{-\dagger}(x) = \int \widetilde{\mathrm{d}}k \Big[a_{-}(k) e^{-ik \cdot x} + a_{+}^{\dagger}(k) e^{+ik \cdot x} \Big] \\ & \left| p^{+} \right\rangle = a_{+}^{\dagger}(p) \left| 0 \right\rangle \qquad \left| p^{-} \right\rangle = a_{-}^{\dagger}(p) \left| 0 \right\rangle \qquad \left| p^{+}_{1}, p^{-}_{2} \right\rangle = a_{-}^{\dagger}(p_{2}) a_{+}^{\dagger}(p_{1}) \left| 0 \right\rangle \end{split}$$

Quantization : Dirac Field (Fermions)

$$: a(x)b(x') := a(x)b(x') - b(x')a(x)$$

$$T(a(x)b(x')) = \theta(t-t')a(x)b(x') - \theta(t'-t)b(x')a(x)$$

$$\mathcal{L} = :i\overline{\psi}\gamma^{\mu}\partial_{\mu}\psi - m\overline{\psi}\psi: \qquad \pi_{\alpha} = \frac{\partial \mathcal{L}}{\partial \dot{\psi}_{\alpha}} = i\psi_{\alpha}^{\dagger} \qquad \pi_{\alpha}^{\dagger} = \frac{\partial \mathcal{L}}{\partial \dot{\psi}_{\alpha}^{\dagger}} = 0$$

$$\left\{b^{\dagger}(p,s), b(p',s')\right\} = \left\{d^{\dagger}(p,s), d(p',s')\right\} = \tilde{\delta}(p-p')\delta_{ss'}$$

$$\psi(x) = \psi^{+}(x) + \psi^{-}(x) = \int \widetilde{dp} \sum_{s} \left[b(p,s)u(p,s)e^{-ip\cdot x} + d^{\dagger}(p,s)v(p,s)e^{+ip\cdot x}\right]$$

$$\overline{\psi}(x) = \overline{\psi}^{\dagger}(x) + \overline{\psi}^{}(x) = \int \widetilde{dp} \sum_{s} \left[d(p,s)\overline{v}(p,s)e^{-ip\cdot x} + b^{\dagger}(p,s)\overline{u}(p,s)e^{+ip\cdot x}\right]$$

$$\left|e^{-}(p_{1},s_{1}), e^{-}(p_{2},s_{2})\right\rangle = b^{\dagger}(p_{2},s_{2})b^{\dagger}(p_{1},s_{1})\left|0\right\rangle = -\left|e^{-}(p_{2},s_{2}), e^{-}(p_{1},s_{1})\right\rangle$$

$$\psi^{+}(x)\left|p\right\rangle = \psi^{+}(x)b^{\dagger}(p)\left|0\right\rangle = \left|0\right\rangle u(p)e^{-ip\cdot x} \qquad \psi^{+}(b) \text{ destroy } e^{-}$$

$$\overline{\psi}^{\dagger}(x)\left|p\right\rangle = \overline{\psi}^{\dagger}(x)d^{\dagger}(p)\left|0\right\rangle = \left|0\right\rangle \overline{v}(p)e^{-ip\cdot x} \qquad \overline{\psi}^{\dagger}(d) \text{ destroy } e^{+}$$

$$\langle p|\psi^{-}(x) = \langle 0|d(p)\psi^{-}(x) = v(p)e^{+ip\cdot x}\langle 0| \qquad \psi^{-}(d^{\dagger}) \text{ creates } e^{+}$$

$$\langle p|\overline{\psi}^{}(x) = \langle 0|b(p)\overline{\psi}^{}(x) = \overline{u}(p)e^{+ip\cdot x}\langle 0| \qquad \overline{\psi}^{}(b^{\dagger}) \text{ creates } e^{-}$$

Quantization: Electromagnetic Field

$$\begin{split} \mathcal{L} &= -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} - \frac{1}{2\xi} (\partial \cdot A)^2 \qquad \pi^\mu = F^{\mu 0} - \frac{g^{\mu 0}}{\xi} (\partial \cdot A) \qquad \pi^0 = -\frac{1}{\xi} (\partial \cdot A) \qquad \pi^k = E^k \\ \left[A_\mu(\vec{\mathbf{x}},t), A_\nu(\vec{\mathbf{y}},t) \right] &= \left[\dot{A}_\mu(\vec{\mathbf{x}},t), \dot{A}_\nu(\vec{\mathbf{y}},t) \right] = 0 \qquad \left[\dot{A}_\mu(\vec{\mathbf{y}},t), A_\nu(\vec{\mathbf{x}},t) \right] = i g_{\mu\nu} \delta^3(\vec{\mathbf{x}} - \vec{\mathbf{y}}) \\ A_\mu(x) &= A_\mu^+(x) + A_\mu^-(x) = \int \widetilde{\mathrm{d}} k \sum_{\lambda=0}^3 \left[a(k,\lambda) \epsilon_\mu(k,\lambda) e^{-ik\cdot x} + a^\dagger(k,\lambda) \epsilon_\mu^*(k,\lambda) e^{+ik\cdot x} \right] \\ A_\mu^+(x) |k\rangle &= A_\mu^+(x) a^\dagger(k) |0\rangle = |0\rangle \epsilon_\mu(k) e^{-ik\cdot x} \qquad \left[\dot{a}(k,\lambda), a^\dagger(k',\lambda') \right] = -g^{\lambda\lambda'} \, \check{\delta}(k-k') \\ \langle k| \, A_\mu^-(x) &= \langle 0| \, a(k) A_\mu^-(x) = \epsilon_\mu^*(k) e^{+ik\cdot x} \, \langle 0| \qquad A^+(a) \, \mathrm{destroys} \, \gamma \quad A^-(a^\dagger) \, \mathrm{creates} \, \gamma \end{split}$$

Dyson Expansion, Propagators and Wick's Theorem

$$S_{fi} = \langle f | \Phi(\infty) \rangle = \langle f | S_{i} \rangle \qquad \sum_{f} \left| S_{fi} \right|^{2} = 1$$

$$S = \sum_{n=0}^{\infty} \frac{(-i)^{n}}{n!} \int d^{4}x_{1} d^{4}x_{2} \dots d^{4}x_{n} T(\mathcal{H}_{int}(x_{1}) \mathcal{H}_{int}(x_{2}) \dots \mathcal{H}_{int}(x_{n}))$$

$$\Delta_{F}(x_{1}-x_{2}) = \phi(x_{1})\phi(x_{2}) = \langle 0|T(\phi(x_{1})\phi(x_{2}))|0\rangle \qquad D_{F}^{\mu\nu} = A^{\mu}(x_{1})A^{\nu}(x_{2})$$

$$\Delta_{F}(x_{1}-x_{2}) = \phi(x_{1})\phi^{\dagger}(x_{2}) = \phi^{\dagger}(x_{2})\phi(x_{1}) \qquad S_{F\alpha\beta} = \psi_{\alpha}(x_{1})\overline{\psi_{\beta}}(x_{2}) = -\overline{\psi_{\beta}}(x_{2})\psi_{\alpha}(x_{1})$$

$$\begin{split} \mathbf{T}(ABCD...WXYZ) &= :ABCD...WXYZ: + \\ &+ : \stackrel{\square}{ABCD}...WXYZ: + :\stackrel{\square}{ABCD}...WXYZ: + \cdots + :ABCD...WXYZ: + \\ &+ : \stackrel{\square}{ABCD}...WXYZ: + \cdots + :ABCD...WXYZ: + \dots \end{split}$$

Decay: Decay Rates

$$\frac{\mathrm{d}\Gamma}{\mathrm{d}\Omega} = \frac{1}{32\pi^2} \frac{\left|\vec{\mathbf{p}}_{\mathrm{CM}}\right|}{m_0^2} |\vec{\mathcal{M}}|^2 \qquad \left|\vec{\mathbf{p}}_{\mathrm{CM}}\right| = \frac{1}{2m_0} \sqrt{\left[m_0^2 - (m_1 + m_2)^2\right] \left[m_0^2 - (m_1 - m_2)^2\right]}$$

Scattering: Cross Sections

$$\frac{\mathrm{d}\sigma}{\mathrm{d}\Omega} = \frac{1}{64\pi^2 s} \frac{|\vec{\mathbf{p}}_3|}{|\vec{\mathbf{p}}_1|} |\overline{\mathscr{M}}|^2 \qquad |\vec{\mathbf{p}}| = \frac{1}{2\sqrt{s}} \sqrt{\left[s - (m_a + m_b)^2\right] \left[s - (m_a - m_b)^2\right]}$$

$$s = (p_1 + p_2)^2 = (p_3 + p_4)^2 \approx +2p_1 \cdot p_2 \approx +2p_3 \cdot p_4 \quad \sqrt{s} = E_{CM}$$

$$t = (p_1 - p_3)^2 = (p_2 - p_4)^2 \approx -2p_1 \cdot p_3 \approx -2p_2 \cdot p_4 \quad s + t + u = m_1^2 + m_2^2 + m_3^2 + m_4^2$$

$$u = (p_1 - p_4)^2 = (p_2 - p_3)^2 \approx -2p_1 \cdot p_4 \approx -2p_2 \cdot p_3$$

$$s + t + u = m_1^2 + m_2^2 + m_3^2 + m_4^2$$

$$u = (p_1 - p_4)^2 = (p_2 - p_3)^2 \approx -2p_1 \cdot p_4 \approx -2p_2 \cdot p_3$$

$$f = \frac{sW}{8M_W^2} = \frac{sW}{8M_W^2} = \frac{sW}{8M_W^2} = \frac{1}{2}v^2 \quad M_W = M_Z \cos\theta_W = \frac{1}{2}g_W v \quad M_H = \sqrt{2\lambda}v$$

$$g_W = g_Z \cos\theta_W = \frac{e}{\sin\theta_W} \quad g_V^f = \frac{1}{2}T_f^3 - Q_f \sin^2(\theta_W) \quad g_A^f = \frac{1}{2}T_f^3 \quad g_{A,V} = \frac{1}{2}c_{A,V}$$

$$f + \overline{f} \longrightarrow g + \overline{g}:$$

$$\sqrt{s} \left(1 + 3 + 3 + 3\right) = \sqrt{s}$$

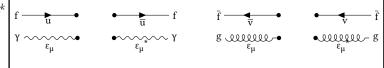
$$p_{1} = \frac{\sqrt{s}}{2} (1, 0, 0, +\beta_{f}) \qquad p_{2} = \frac{\sqrt{s}}{2} (1, 0, 0, -\beta_{f})$$

$$p_{3} = \frac{\sqrt{s}}{2} (1, +\beta_{g} \sin \theta, 0, +\beta_{g} \cos \theta) \qquad p_{4} = \frac{\sqrt{s}}{2} (1, -\beta_{g} \sin \theta, 0, -\beta_{g} \cos \theta)$$

$$m = 0$$
: $s = 2|\vec{\mathbf{p}}_{CM}|^2 = 2(p_{CM}^0)^2$ $t = -\frac{s}{2}(1 - \cos\theta)$ $u = -\frac{s}{2}(1 + \cos\theta)$

Feynman Rules for $i \mathcal{M}$

Goes in opposite way of arrows with the first one being adjoint, $\bar{\psi} = \psi^{\dagger} \gamma^0$:





$$\bullet - - \frac{\mathbf{h}}{i} - - \bullet$$

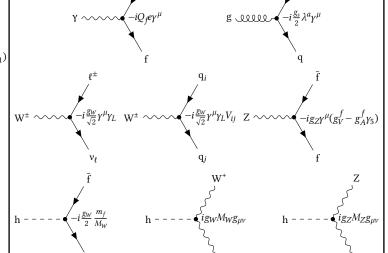
$$\frac{\mathbf{g}}{p^2 - M_h^2 + i\epsilon}$$

$$\bullet \underbrace{-i \frac{\mathbf{g}}{g_\mu N_h^{ab}}}_{k^2 + i\epsilon}$$

$$\bullet \underbrace{-i \frac{\mathbf{g}}{g_\mu N_h^{ab}}}_{k^2 + i\epsilon}$$

$$\bullet \underbrace{-i \frac{\mathbf{g}}{g_{\mu\nu} - k_\mu k_\nu / M_W^2}}_{k^2 - M_W^2 + i M_W \Gamma_W}$$

$$\bullet - \underbrace{-i \frac{\sigma^2}{p^2 - \xi_W M_W^2 + i\epsilon}}_{p^2 - \xi_W M_W^2 + i\epsilon}$$



Constants

$$\begin{split} \frac{G_F}{\sqrt{2}} &= \frac{g_W^2}{8M_W^2} = \frac{g_W^2}{8M_Z^2\cos^2(\theta_W)} = \frac{1}{2v^2} \qquad M_W = M_Z\cos\theta_W = \frac{1}{2}g_Wv \qquad M_H = \sqrt{2\lambda}v \\ g_W &= g_Z\cos\theta_W = \frac{e}{\sin\theta_W} \qquad g_V^f = \frac{1}{2}T_f^3 - Q_f\sin^2(\theta_W) \qquad g_A^f = \frac{1}{2}T_f^3 \qquad g_{A,V} = \frac{1}{2}c_{A,V} \\ c_L^f &= \frac{g_V^f + g_A^f}{2} \qquad c_R^f = \frac{g_V^f - g_A^f}{2} \qquad g_V^f = c_L^f + c_R^f \qquad g_A^f = c_L^f - c_R^f \end{split}$$

Fermions	Q_f	$I_{W}^{(3)}$	Y_L	Y_R	c_L	c_R	c_V	c_A
v_e, v_{μ}, v_{τ}	0	+1/2	-1	0	+1/2	0	+1/2	+1/2
e^-, μ^-, τ^-	-1	-1/2		-2	-0.27	+0.23	-0.04	-1/2
u, c, t	+2/3	+1/2	+1/3	+4/3	+0.35	-0.15	+0.19	+1/2
d, s, b	-1/3	-1/2	+1/3	-2/3	-0.42	+0.08	-0.35	-1/2

Relativity

$$\beta = \frac{v}{c} = \frac{p}{E} = \tanh(\eta) \qquad \gamma = \frac{1}{\sqrt{1 - \beta^2}} = \cosh(\eta) \qquad \gamma \beta = \sin(\eta)$$

$$E^2 = m^2 + |p|^2 \qquad \frac{\mathrm{d}^3 \vec{\mathbf{p}}'}{E'} = \frac{\mathrm{d}^3 \vec{\mathbf{p}}}{E} \qquad \widetilde{\mathrm{d}} p \equiv \frac{\mathrm{d}^3 \vec{\mathbf{p}}}{(2\pi)^3 2E_p} \qquad \widetilde{\delta}(p - q) \equiv (2\pi)^3 2E_p \, \delta^3(\vec{\mathbf{p}} - \vec{\mathbf{q}})$$

$$\begin{split} x^{\mu} &= (t, \vec{\mathbf{x}}) \qquad p^{\mu} &= (E, \vec{\mathbf{p}}) \qquad \partial^{\mu} &= \left(\partial_{t}, -\vec{\mathbf{V}}\right) \qquad A^{\mu} &= \left(\phi, \vec{\mathbf{A}}\right) \qquad J^{\mu} &= \left(\rho, \vec{\mathbf{J}}\right) \\ F^{\mu\nu} &\equiv \partial^{\mu}A^{\nu} - \partial^{\nu}A^{\mu} \qquad \mathcal{F}^{\mu\nu} &= \frac{1}{2} \epsilon^{\mu\nu\alpha\beta} F_{\alpha\beta} \qquad \partial_{\mu}F^{\mu\nu} &= J^{\nu} \qquad \partial_{\mu}\mathcal{F}^{\mu\nu} &= 0 \end{split}$$

$$\epsilon^{\mu\nu\rho\sigma} = \begin{cases} +1 & \text{If } (\mu\nu\rho\sigma) \text{ is an even permutation of (0123)} \\ -1 & \text{If } (\mu\nu\rho\sigma) \text{ is an odd permutation of (0123)} \\ 0 & \text{Otherwise} \end{cases}$$

$$\begin{split} \epsilon_{\alpha\beta_{1}\gamma_{1}\delta_{1}}\epsilon^{\alpha\beta_{2}\gamma_{2}\delta_{2}} &= +g_{\beta_{1}}^{\beta_{2}}g_{\gamma_{1}}^{\delta_{2}}g_{\delta_{1}}^{\gamma_{2}} - g_{\beta_{1}}^{\beta_{2}}g_{\gamma_{1}}^{\gamma_{2}}g_{\delta_{1}}^{\delta_{2}} + g_{\beta_{1}}^{\gamma_{2}}g_{\gamma_{1}}^{\beta_{2}}g_{\delta_{1}}^{\delta_{2}} - g_{\beta_{1}}^{\gamma_{2}}g_{\gamma_{1}}^{\beta_{2}}g_{\delta_{1}}^{\beta_{2}} \\ &\quad + g_{\beta_{1}}^{\delta_{2}}g_{\gamma_{1}}^{\gamma_{2}}g_{\delta_{1}}^{\beta_{2}} - g_{\beta_{1}}^{\delta_{2}}g_{\gamma_{1}}^{\gamma_{2}}g_{\delta_{1}}^{\beta_{2}} \\ &\quad \epsilon_{\alpha\beta\gamma_{1}\delta_{1}}\epsilon^{\alpha\beta\gamma_{2}\delta_{2}} = -2\Big(g_{\gamma_{1}}^{\gamma_{2}}g_{\delta_{1}}^{\delta_{2}} - g_{\gamma_{1}}^{\delta_{2}}g_{\delta_{1}}^{\gamma_{2}}\Big) \\ &\quad \epsilon_{\alpha\beta\gamma_{1}\delta_{1}}\epsilon^{\alpha\beta\gamma_{2}\delta_{2}} = -6g_{\delta_{1}}^{\delta_{2}} \end{split}$$

Quantum Mechanics

$$E \to i \frac{\partial}{\partial t} \qquad \vec{\mathbf{p}} \to -i \vec{\nabla} \cdot \implies p^{\mu} \to i \delta^{\mu}$$

Mathematics

Dirac Delta:

$$\left\{ i g_{Z} M_{Z} g_{\mu\nu} \middle| \delta^{n}(x'-x) = i \int \frac{\mathrm{d}^{n} p}{(2\pi)^{n}} e^{-ip(x'-x)} \qquad f(x_{i}) = 0 \implies \delta[f(x)] = \sum_{i} \left| \frac{\mathrm{d} f}{\mathrm{d} x} \right|^{-1} \cdot \delta(x-x_{i}) \right\}$$

Triangle Function:

$$\lambda(a, b, c) = \left[a - \left(\sqrt{b} + \sqrt{c}\right)^{2}\right] \left[a - \left(\sqrt{b} - \sqrt{c}\right)^{2}\right]$$

Residue Theorem:

$$\oint_{\Gamma} dz f(z) = \pm 2\pi i \sum_{k=1}^{n} \operatorname{Res}[f(a_k)] \qquad \operatorname{Res}[f(a_k)] = \lim_{z \to a_k} (z - a_k) f(z)$$

The "-" ("+") is used when Γ is oriented clockwise (counterclockwise).

$$f(z) = \frac{h(z)e^{-iz(t'-t)}}{(z-a_1)(z-a_2)\dots} \stackrel{t'-t<0}{\Longrightarrow} \int_{\Re(z)=-\infty}^{\Re(z)=+\infty} f(z) \, \mathrm{d}z = +2\pi i \sum_{a_k \text{ in UHP}} \mathrm{Res}[f(a_k)]$$

$$\stackrel{t'-t>0}{\Longrightarrow} \int_{\Re(z)=-\infty}^{\Re(z)=+\infty} f(z) \, \mathrm{d}z = -2\pi i \sum_{a_k \text{ in LHP}} \mathrm{Res}[f(a_k)]$$