#### Equations

Schrödinger, Klein-Gordon (Real Scalar Field), Dirac (Complex Scalar Field):

$$\begin{split} i\hbar\frac{\partial|\psi(t)\rangle}{\partial t} &= H|\psi(t)\rangle \qquad \mathcal{L} = i\psi^*\dot{\psi} - \frac{1}{2m}\vec{\nabla}\psi^*\cdot\vec{\nabla}\psi \\ (\partial\cdot\partial + m^2)\phi &= 0 \qquad \mathcal{L} = \frac{1}{2}\Big[\partial_\mu\phi\partial^\mu\phi - m^2\phi^2\Big] &= \frac{1}{2}\Big[\dot{\phi}^2 - \left(\vec{\nabla}\phi\right)^2 - m^2\phi^2\Big] \\ (i\not{\partial} - m)\phi &= 0 \qquad \overline{\phi}(i\not{\partial} + m) = 0 \qquad \mathcal{L} = \partial_\mu\phi\partial^\mu\phi^* - m^2\phi\phi^* \end{split}$$

#### Pauli Matrices

$$\begin{split} \sigma^1 &= \sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \quad \sigma^2 = \sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \quad \sigma^3 = \sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \quad \vec{\Sigma} = \begin{pmatrix} \vec{\sigma} & 0 \\ 0 & \vec{\sigma} \end{pmatrix} \\ \left\{ \sigma^i, \sigma^j \right\} &= 2\delta^{ij} \qquad \left[ \sigma^i, \sigma^j \right] = 2i\epsilon^{ijk}\sigma^k \qquad \sigma^i\sigma^j = \delta^{ij} + i\epsilon^{ijk}\sigma^k \qquad \sigma^2\sigma^i\sigma^2 = -(\sigma^i)^* \\ \sigma^\mu &= \begin{pmatrix} 1, \vec{\sigma} \end{pmatrix} \quad \vec{\sigma}^\mu = \begin{pmatrix} 1, -\vec{\sigma} \end{pmatrix} \qquad (p \cdot \sigma)(p \cdot \vec{\sigma}) = p^2 \qquad \left( \vec{\sigma} \cdot \vec{\mathbf{p}} \right)^2 = |\vec{\mathbf{p}}|^2 \\ p \cdot \vec{\sigma} &= \begin{pmatrix} p^0 + p^3 & p^1 - ip^2 \\ p^1 + ip^2 & p^0 - p^3 \end{pmatrix} \qquad p \cdot \sigma = \begin{pmatrix} p^0 - p^3 & -(p^1 - ip^2) \\ -(p^1 + ip^2) & p^0 + p^3 \end{pmatrix} \\ \sqrt{p \cdot \sigma} &= \frac{E + m - \vec{\sigma} \cdot \vec{\mathbf{p}}}{\sqrt{2(E + m)}} \qquad \sqrt{p \cdot \vec{\sigma}} &= \frac{E + m + \vec{\sigma} \cdot \vec{\mathbf{p}}}{\sqrt{2(E + m)}} \qquad \vec{\sigma} \cdot \vec{\mathbf{p}} &= \begin{pmatrix} p_z & p_x - ip_y \\ p_x + ip_y & -p_z \end{pmatrix} \end{split}$$

#### Dirac y-Matrices

$$\begin{split} \gamma^{\mu} &= \begin{pmatrix} 0 & \sigma^{\mu} \\ \overline{\sigma}^{\mu} & 0 \end{pmatrix} &<- \operatorname{Chiral} \mid \operatorname{Pauli-Dirac} -> & \gamma^{0} &= \begin{pmatrix} +1 & 0 \\ 0 & -1 \end{pmatrix} & \gamma^{k} &= \begin{pmatrix} 0 & \sigma^{k} \\ -\sigma^{k} & 0 \end{pmatrix} \\ & \{\gamma^{\mu}, \gamma^{\nu}\} = 2g^{\mu\nu} \quad \gamma_{5} = +i\gamma^{0}\gamma^{1}\gamma^{2}\gamma^{3} = -i\gamma_{0}\gamma_{1}\gamma_{2}\gamma_{3} = -\frac{i}{4!}\epsilon^{\alpha\beta\gamma\delta}\gamma^{\alpha}\gamma^{\beta}\gamma^{\gamma}\gamma^{\delta} \\ & \sigma^{\mu\nu} &= \frac{i}{2}[\gamma^{\mu}, \gamma^{\nu}] \implies [\gamma^{\mu}, \gamma^{\nu}] = -2i\sigma^{\mu\nu} & \{\gamma_{5}, \gamma^{\mu}\} = 0 & \gamma^{0}(\gamma^{\mu})^{\dagger}\gamma^{0} = \gamma^{\mu} \\ & (\gamma^{0})^{\dagger} = \gamma^{0} & (\gamma^{0})^{2} = 1 & (\gamma^{k})^{\dagger} = -\gamma^{k} & (\gamma^{k})^{2} = -1 & (\gamma^{5})^{\dagger} = \gamma^{5} & (\gamma^{5})^{2} = 1 \\ & \overline{\Gamma} = \gamma^{0}\Gamma^{\dagger}\gamma^{0} & \overline{\gamma_{5}} = -\gamma_{5} & \overline{\gamma^{\mu}} = \gamma^{\mu} & \overline{\gamma^{\mu}\gamma_{5}} = \gamma^{\mu}\gamma_{5} & \overline{\sigma^{\mu\nu}} = \sigma^{\mu\nu} \end{split}$$

$$\begin{split} & \not\! A = A^\mu \gamma_\mu \qquad \gamma^\mu \gamma_\mu = 4 \qquad \gamma^\mu \gamma^\nu \gamma_\mu = -2 \gamma^\nu \qquad \gamma^\mu \gamma^\nu \gamma^\rho \gamma_\mu = 4 g^{\nu\rho} \\ & \gamma^\mu \gamma^\nu \gamma^\rho \gamma^\sigma \gamma_\mu = -2 \gamma^\sigma \gamma^\rho \gamma^\nu \qquad \gamma^\mu \gamma^\nu \gamma^\rho \gamma^\sigma \gamma^\pi \gamma_\mu = 2 [\gamma^\pi \gamma^\nu \gamma^\rho \gamma^\sigma + \gamma^\sigma \gamma^\rho \gamma^\nu \gamma^\pi] \\ & \qquad \qquad \gamma^\mu \gamma^\nu \gamma^\rho = i \epsilon^{\mu\nu\rho\lambda} \gamma_\lambda \gamma_5 + g^{\mu\nu} \gamma^\rho - g^{\mu\rho} \gamma^\nu + g^{\nu\rho} \gamma^\mu \end{split}$$

### Spin, Helicity and Chirality

$$\vec{\mathbf{S}} = \frac{1}{2}\vec{\mathbf{\Sigma}} \qquad \tilde{h} = \frac{\vec{\mathbf{S}} \cdot \vec{\mathbf{p}}}{|\vec{\mathbf{p}}|} = \frac{\vec{\mathbf{\Sigma}} \cdot \vec{\mathbf{p}}}{2|\vec{\mathbf{p}}|} \qquad \tilde{h} = 2h = \frac{\vec{\mathbf{\Sigma}} \cdot \vec{\mathbf{p}}}{|\vec{\mathbf{p}}|} \qquad P_L = \frac{1-h}{2} \qquad P_R = \frac{1+h}{2}$$

$$\gamma_L = \frac{1-\gamma_5}{2} \qquad \gamma_R = \frac{1+\gamma_5}{2} \qquad \gamma_{L,R}^2 = \gamma_{L,R} \qquad \gamma_{L,R}\gamma_{R,L} = 0 \qquad \gamma_L + \gamma_R = 1$$

$$\gamma^{\mu}\gamma_{L,R} = \gamma_{R,L}\gamma^{\mu} \qquad \gamma_{5}\gamma_L = \gamma_L\gamma_5 = -\gamma_L \qquad \gamma_5\gamma_R = \gamma_R\gamma_5 = \gamma_R \qquad \gamma_{L,R}^{\dagger} = \gamma_L\gamma_R = \gamma_{L,R}$$

$$\overline{\gamma^{\mu}\gamma_{L,R}} = \gamma^{\mu}\gamma_{L,R} \qquad \overline{\gamma_5\gamma_L} = -\gamma_R \qquad \overline{\gamma_5\gamma_R} = \gamma_L \qquad \overline{\gamma_{L,R}} = \gamma_0\gamma_{L,R}^{\dagger}\gamma_0 = \gamma_{R,L}$$

$$\begin{split} u_{\uparrow}^T &= N \left( c \quad se^{i\phi} \quad kc \quad kse^{i\phi} \right) \approx N \left( c \quad se^{i\phi} \quad c \quad se^{i\phi} \right) & N = \sqrt{E+m} \approx \sqrt{E} \\ u_{\downarrow}^T &= N \left( -s \quad ce^{i\phi} \quad ks \quad -kce^{i\phi} \right) \approx N \left( -s \quad ce^{i\phi} \quad s \quad -ce^{i\phi} \right) & k = \frac{p}{E+m} \approx 1 \\ v_{\uparrow}^T &= N \left( ks \quad -kce^{i\phi} \quad -s \quad ce^{i\phi} \right) \approx N \left( s \quad -ce^{i\phi} \quad -s \quad ce^{i\phi} \right) & s = \sin \left( \frac{\theta}{2} \right) \\ v_{\downarrow}^T &= N \left( kc \quad kse^{i\phi} \quad c \quad se^{i\phi} \right) \approx N \left( c \quad se^{i\phi} \quad c \quad se^{i\phi} \right) & c = \cos \left( \frac{\theta}{2} \right) \\ \vec{\mathbf{p}} &= p (\sin \theta \cos \phi, \sin \theta \sin \phi, \cos \theta) \end{split}$$

#### Traces

Traces with an odd number of  $\gamma$  matrices are 0.

$$\begin{split} \operatorname{Tr}[1] &= 4 & \operatorname{Tr}[\gamma^{\mu}\gamma^{\nu}] &= 4g^{\mu\nu} & \operatorname{Tr}[\gamma^{\mu}\gamma^{\nu}\gamma^{\rho}\gamma^{\sigma}] &= 4(g^{\mu\nu}g^{\rho\sigma} - g^{\mu\rho}g^{\nu\sigma} + g^{\mu\sigma}g^{\nu\rho}) \\ \operatorname{Tr}[\gamma_{5}] &= 0 & \operatorname{Tr}[\gamma_{5}\gamma^{\mu}\gamma^{\nu}] &= 0 & \operatorname{Tr}[\gamma_{5}\gamma^{\mu}\gamma^{\nu}\gamma^{\rho}\gamma^{\sigma}] &= -4i\epsilon^{\mu\nu\rho\sigma} \\ & \operatorname{Tr}[\gamma^{\mu}\gamma^{\nu}\gamma^{\rho}\gamma^{\sigma}\chi_{L}] &= 2(g^{\mu\nu}g^{\rho\sigma} - g^{\mu\rho}g^{\nu\sigma} + g^{\mu\sigma}g^{\nu\rho} + i\epsilon^{\mu\nu\rho\sigma}) \\ & \operatorname{Tr}[\gamma^{\mu}\gamma^{\nu}\gamma^{\rho}\gamma^{\sigma}\gamma_{R}] &= 2(g^{\mu\nu}g^{\rho\sigma} - g^{\mu\rho}g^{\nu\sigma} + g^{\mu\sigma}g^{\nu\rho} - i\epsilon^{\mu\nu\rho\sigma}) \\ & \operatorname{Tr}[y^{\mu}\gamma^{\nu}\gamma^{\rho}\gamma^{\sigma}\gamma_{R}] &= 2(g^{\mu\nu}g^{\rho\sigma} - g^{\mu\rho}g^{\nu\sigma} + g^{\mu\sigma}g^{\nu\rho} - i\epsilon^{\mu\nu\rho\sigma}) \\ & \operatorname{Tr}[y^{\mu}\gamma^{\mu}\chi_{L}y^{\mu}\gamma^{\nu}(g_{V} - g_{A}\gamma_{5})y^{\mu}\chi_{L}y^{\mu}\chi_{L}y^{\mu}\gamma_{V}\chi_{L}] &= -16(g_{V} + g_{A})(p_{1} \cdot p_{4})(p_{2} \cdot p_{3}) \\ & \operatorname{Tr}[(y_{A} + m_{a})\gamma^{\mu}(y_{b} + m_{b})\gamma^{\nu}] &= 4\left[p_{A}^{\mu}p_{b}^{\nu} + p_{b}^{\mu}p_{a}^{\nu} + (m_{a}m_{b} - p_{a} \cdot p_{b})g^{\mu\nu}\right] \\ & \operatorname{Tr}[\phi\gamma^{\mu}\beta\gamma^{\nu}] \operatorname{Tr}[\phi\gamma_{\mu}d\gamma_{\nu}\gamma_{5}] &= 0 & \operatorname{Tr}[\phi\gamma^{\mu}\beta\gamma^{\nu}\gamma_{L}] \operatorname{Tr}[\phi\gamma_{\mu}d\gamma_{\nu}\gamma_{R}] &= 16(a \cdot d)(b \cdot c) \\ & \operatorname{Tr}[\phi\gamma^{\mu}\beta\gamma^{\nu}\gamma_{L}] \operatorname{Tr}[\phi\gamma_{\mu}d\gamma_{\nu}\gamma_{L}] &= \operatorname{Tr}[\phi\gamma^{\mu}\beta\gamma^{\nu}\gamma_{L}] \operatorname{Tr}[\phi\gamma_{\mu}d\gamma_{\nu}\gamma_{L}] &= 16(a \cdot c)(b \cdot d) \end{split}$$

## Spinors (Fermions)

$$\psi = ue^{+i(\vec{\mathbf{p}}\cdot\vec{\mathbf{x}}-Et)} = ve^{-i(\vec{\mathbf{p}}\cdot\vec{\mathbf{x}}-Et)} \qquad \overline{\psi} = \psi^{\dagger}\gamma^{0}$$

$$(\not{p}-m)u = \overline{u}(\not{p}-m) = 0 \qquad (\not{p}+m)v = \overline{v}(\not{p}+m) = 0$$

$$\overline{u}'(p)u^{s}(p) = +2m\delta^{rs} \qquad \overline{u}'(p)v^{s}(p) = 0 \qquad u^{r\dagger}(p)u^{s}(p) = 2E\delta^{rs} \qquad \sum_{s=1,2} u^{s}(p)\overline{u}^{s}(p) = \not{p}+m$$

$$\overline{v}'(p)v^{s}(p) = -2m\delta^{rs} \qquad \overline{v}'(p)u^{s}(p) = 0 \qquad v^{r\dagger}(p)v^{s}(p) = 2E\delta^{rs} \qquad \sum_{s=1,2} v^{s}(p)\overline{v}^{s}(p) = \not{p}-m$$

#### **Polarization Vectors (Bosons)**

External massless & Massive:

$$\sum_{\lambda=1}^{2} \epsilon^{\mu}(k,\lambda) \cdot \epsilon^{*\nu}(k,\lambda) = -g^{\mu\nu} \qquad \sum_{\lambda=1}^{3} \epsilon_{\mu}(q,\lambda) \epsilon_{\nu}^{*}(q,\lambda) = -g_{\mu\nu} + \frac{q_{\mu}q_{\nu}}{M^{2}}$$

### **Quantization: Real Scalar Field (Bosons)**

$$: \frac{1}{2} \left[ a^{\dagger}(k)a(k) + a(k)a^{\dagger}(k) \right] := a^{\dagger}(k)a(k) \qquad : a(x)b(x') := a(x)b(x') + b(x')a(x)$$

$$T(a(x)b(x')) = \theta(t - t')a(x)b(x') + \theta(t' - t)b(x')a(x)$$

The "+" corresponds to positive frequency plane waves  $e^{-ik \cdot x}$ :

$$\mathcal{L} = \frac{1}{2} \partial^{\mu} \phi \partial_{\mu} \phi - \frac{1}{2} m^{2} \phi \phi \qquad \pi = \dot{\phi} \qquad \left[ a^{\dagger}(k, \lambda), a(k', \lambda') \right] = g^{\lambda \lambda'} \tilde{\delta}(k - k')$$

$$\left[ \phi(\vec{\mathbf{x}}, t), \phi(\vec{\mathbf{y}}, t) \right] = \left[ \Pi(\vec{\mathbf{x}}, t), \Pi(\vec{\mathbf{y}}, t) \right] = 0 \qquad \left[ \phi(\vec{\mathbf{x}}, t), \Pi(\vec{\mathbf{y}}, t) \right] = i \delta^{3} (\vec{\mathbf{x}} - \vec{\mathbf{y}})$$

$$\phi(x) = \int \tilde{d}k \left[ a(k) e^{-ik \cdot x} + a^{\dagger}(k) e^{+ik \cdot x} \right]$$

$$|p\rangle = a^{\dagger}(p) |0\rangle \qquad |p_{1}, p_{2}\rangle = a^{\dagger}(p_{2}) a^{\dagger}(p_{1}) |0\rangle = |p_{2}, p_{1}\rangle$$

### **Quantization : Complex Scalar Field (Bosons)**

$$\begin{split} \mathcal{L} &= : \partial^{\mu} \varphi^{\dagger} \partial_{\mu} \varphi - m^{2} \varphi^{\dagger} \varphi \colon \qquad \pi = \dot{\varphi}^{\dagger} \qquad \pi^{\dagger} = \dot{\varphi} \qquad \left[ a_{\pm}(k), a_{\pm}^{\dagger}(k') \right] = \tilde{\delta}(k - k') \\ & \left[ \varphi(\vec{\mathbf{x}}, t), \pi(\vec{\mathbf{y}}, t) \right] = \left[ \varphi^{\dagger}(\vec{\mathbf{x}}, t), \pi^{\dagger}(\vec{\mathbf{y}}, t) \right] = i \, \delta^{3}(\vec{\mathbf{x}} - \vec{\mathbf{y}}) \\ & \varphi(x) = \varphi^{+}(x) + \varphi^{-}(x) = \int \widetilde{\mathrm{d}}k \left[ a_{+}(k) e^{-ik \cdot x} + a_{-}^{\dagger}(k) e^{+ik \cdot x} \right] \\ & \varphi^{\dagger}(x) = \varphi^{+\dagger}(x) + \varphi^{-\dagger}(x) = \int \widetilde{\mathrm{d}}k \left[ a_{-}(k) e^{-ik \cdot x} + a_{+}^{\dagger}(k) e^{+ik \cdot x} \right] \\ & | p^{+} \rangle = a_{+}^{\dagger}(p) | 0 \rangle \qquad | p^{-} \rangle = a_{-}^{\dagger}(p) | 0 \rangle \qquad | p_{1}^{+}, p_{2}^{-} \rangle = a_{-}^{\dagger}(p_{2}) a_{+}^{\dagger}(p_{1}) | 0 \rangle \end{split}$$

#### **Ouantization : Dirac Field (Fermions)**

$$: a(x)b(x') := a(x)b(x') - b(x')a(x)$$

$$T(a(x)b(x')) = \theta(t-t)a(x)b(x') - \theta(t'-t)b(x')a(x)$$

$$\mathcal{L} = i\bar{\psi}\gamma^{\mu}\partial_{\mu}\psi - m\bar{\psi}\psi : \qquad \pi_{\alpha} = \frac{\partial \mathcal{L}}{\partial\dot{\psi}_{\alpha}} = i\psi_{\alpha}^{\dagger} \qquad \pi_{\alpha}^{\dagger} = \frac{\partial \mathcal{L}}{\partial\dot{\psi}_{\alpha}^{\dagger}} = 0$$

$$\left\{ b^{\dagger}(p,s), b(p',s') \right\} = \left\{ d^{\dagger}(p,s), d(p',s') \right\} = \tilde{\delta}(p-p')\delta_{ss'}$$

$$\psi(x) = \psi^{+}(x) + \psi^{-}(x) = \int \widetilde{\mathrm{d}p} \sum_{s} \left[ b(p,s)u(p,s)e^{-ip\cdot x} + d^{\dagger}(p,s)v(p,s)e^{+ip\cdot x} \right]$$

$$\bar{\psi}(x) = \bar{\psi}^{\dagger}(x) + \bar{\psi}^{-}(x) = \int \widetilde{\mathrm{d}p} \sum_{s} \left[ d(p,s)\bar{v}(p,s)e^{-ip\cdot x} + b^{\dagger}(p,s)\bar{u}(p,s)e^{+ip\cdot x} \right]$$

$$|e^{-}(p_1,s_1), e^{-}(p_2,s_2)\rangle = b^{\dagger}(p_2,s_2)b^{\dagger}(p_1,s_1)|0\rangle = -|e^{-}(p_2,s_2), e^{-}(p_1,s_1)\rangle$$

$$\psi^{+}(x)|p\rangle = \psi^{+}(x)b^{\dagger}(p)|0\rangle = |0\rangle u(p)e^{-ip\cdot x} \qquad \psi^{+}(b) \text{ destroys } e^{-}$$

$$\bar{\psi}^{\dagger}(x)|p\rangle = \bar{\psi}^{\dagger}(x)d^{\dagger}(p)|0\rangle = |0\rangle \bar{v}(p)e^{-ip\cdot x} \qquad \bar{\psi}^{\dagger}(d) \text{ destroy } e^{+}$$

$$\langle p|\psi^{-}(x) = \langle 0|d(p)\psi^{-}(x) = v(p)e^{+ip\cdot x}\langle 0| \qquad \psi^{-}(d^{\dagger}) \text{ creates } e^{+}$$

$$\langle p|\bar{\psi}^{-}(x) = \langle 0|b(p)\bar{\psi}^{-}(x) = \bar{u}(p)e^{+ip\cdot x}\langle 0| \qquad \bar{\psi}^{-}(b^{\dagger}) \text{ creates } e^{-}$$

### **Quantization**: Electromagnetic Field

$$\begin{split} \mathcal{L} &= -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} - \frac{1}{2\xi}(\partial \cdot A)^2 \qquad \pi^{\mu} = F^{\mu 0} - \frac{g^{\mu 0}}{\xi}(\partial \cdot A) \qquad \pi^0 = -\frac{1}{\xi}(\partial \cdot A) \qquad \pi^k = E^k \\ & \left[A_{\mu}(\vec{\mathbf{x}},t),A_{\nu}(\vec{\mathbf{y}},t)\right] = \left[\dot{A}_{\mu}(\vec{\mathbf{x}},t),\dot{A}_{\nu}(\vec{\mathbf{y}},t)\right] = 0 \qquad \left[\dot{A}_{\mu}(\vec{\mathbf{y}},t),A_{\nu}(\vec{\mathbf{x}},t)\right] = ig_{\mu\nu}\,\delta^3(\vec{\mathbf{x}}-\vec{\mathbf{y}}) \\ & A_{\mu}(x) = A_{\mu}^+(x) + A_{\mu}^-(x) = \int \tilde{\mathrm{d}}k \sum_{\lambda=0}^3 \left[a(k,\lambda)\epsilon_{\mu}(k,\lambda)e^{-ik\cdot x} + a^{\dagger}(k,\lambda)\epsilon_{\mu}^*(k,\lambda)e^{+ik\cdot x}\right] \\ & A_{\mu}^+(x)|k\rangle = A_{\mu}^+(x)a^{\dagger}(k)|0\rangle = |0\rangle\,\epsilon_{\mu}(k)e^{-ik\cdot x} \qquad \left[\dot{a}(k,\lambda),a^{\dagger}(k',\lambda')\right] = -g^{\lambda\lambda'}\,\tilde{\delta}(k-k') \\ & \langle k|A_{\mu}^-(x) = \langle 0|a(k)A_{\mu}^-(x) = \epsilon_{\mu}^*(k)e^{+ik\cdot x}\langle 0| \qquad A^+(a) \text{ destroys } \gamma \quad A^-(a^{\dagger}) \text{ creates } \gamma \end{split}$$

#### **Classical Field Theory**

$$\begin{split} S &= \int_{t_i}^{t_f} L \, \mathrm{d}t = \int \mathrm{d}^4 x \, \mathscr{L} \qquad L = \int \mathrm{d}^3 \vec{\mathbf{x}} \, \mathscr{L}(\phi, \partial_\mu \phi) \qquad \frac{\partial \mathscr{L}}{\partial \phi_i} - \partial_\mu \frac{\partial \mathscr{L}}{\partial \left(\partial_\mu \phi_i\right)} = 0 \\ H &= \int \mathrm{d}^3 \vec{\mathbf{x}} \, \mathscr{H} \qquad \mathscr{H} = \sum_i \Pi_i(x) \partial_0 \phi_i(x) - \mathscr{L} \qquad \Pi_i(x) = \frac{\partial \mathscr{L}}{\partial \left(\partial_0 \phi_i(x)\right)} \\ \dot{\psi} &= \frac{\partial \mathscr{H}}{\partial \Pi} - \vec{\nabla} \cdot \frac{\partial \mathscr{H}}{\partial (\vec{\mathbf{V}}\Pi)} \qquad \dot{\Pi} = -\frac{\partial \mathscr{H}}{\partial \psi} + \vec{\nabla} \cdot \frac{\partial \mathscr{H}}{\partial (\vec{\mathbf{V}}\psi)} \\ \delta \mathscr{L} &= \mathscr{L} - \mathscr{L} = \partial_\mu C^\mu \qquad J^\mu = C^\mu - \frac{\partial \mathscr{L}}{\partial \partial_\mu \phi} \delta \phi \qquad \partial_\mu J^\mu = \frac{\partial J^0}{\partial t} + \vec{\nabla} \cdot \vec{\mathbf{J}} = 0 \\ Q &= \int \mathrm{d}^3 \vec{\mathbf{x}} \, J^0 \qquad \frac{\mathrm{d} Q}{\mathrm{d} t} = 0 \qquad A^\mu \to A^\mu + \partial^\mu \lambda \end{split}$$

Ward identity:  $k_{\mu} \mathcal{M}^{\mu} = 0$ , where  $\mathcal{M} = \epsilon_{\mu} \mathcal{M}^{\mu}$ . Photon polarizations  $\epsilon$  parallel to its direction of propagation don't contribute to the scattering amplitude.

# Dyson Expansion, Propagators and Wick's Theorem

$$S = \sum_{n=0}^{\infty} \frac{(-i)^n}{n!} \int d^4x_1 d^4x_2 \dots d^4x_n \, T(\mathcal{X}_{int}(x_1) \, \mathcal{H}_{int}(x_2) \dots \mathcal{H}_{int}(x_n)) \qquad \sum_{f} \left| S_{fi} \right|^2 = 1$$

$$\Delta_F(x_1 - x_2) = \phi(x_1) \phi(x_2) = \langle 0 | T(\phi(x_1) \phi(x_2)) | 0 \rangle \qquad D_F^{\mu\nu} = A^{\mu}(x_1) A^{\nu}(x_2)$$

$$\Delta_F(x_1 - x_2) = \phi(x_1) \phi^{\dagger}(x_2) = \phi^{\dagger}(x_2) \phi(x_1) \qquad S_{F\alpha\beta} = \psi_{\alpha}(x_1) \overline{\psi_{\beta}}(x_2) = -\overline{\psi_{\beta}}(x_2) \psi_{\alpha}(x_1)$$

$$T(ABCD...WXYZ) = :ABCD...WXYZ: + :ABCD...WXYZ: + :ABCD...WXYZ: + :ABCD...WXYZ: + :ABCD...WXYZ: + ...$$

#### **Decay: Decay Rates**

$$\frac{\mathrm{d}\Gamma}{\mathrm{d}\Omega} = \frac{1}{32\pi^2} \frac{\left|\vec{\mathbf{p}}_{\mathrm{CM}}\right|}{m_0^2} |\vec{\mathcal{M}}|^2 \qquad \left|\vec{\mathbf{p}}_{\mathrm{CM}}\right| = \frac{1}{2m_0} \sqrt{\left[m_0^2 - (m_1 + m_2)^2\right] \left[m_0^2 - (m_1 - m_2)^2\right]}$$

$$E_1 = \frac{m_0^2 + m_1^2 - m_2^2}{2m_0} \qquad E_2 = \frac{m_0^2 + m_2^2 - m_1^2}{2m_0}$$

### **Scattering: Cross Sections**

$$\begin{split} \frac{\mathrm{d}\sigma}{\mathrm{d}\Omega} &= \frac{1}{64\pi^2 s} \frac{|\vec{\mathbf{p}}_3|}{|\vec{\mathbf{p}}_1|} \overline{|\mathcal{M}|^2} \qquad |\vec{\mathbf{p}}_a| = |\vec{\mathbf{p}}_b| = \frac{1}{2\sqrt{s}} \sqrt{\left[s - (m_a + m_b)^2\right] \left[s - (m_a - m_b)^2\right]} \\ E_1 &= \frac{s + m_1^2 - m_2^2}{2\sqrt{s}} \quad E_2 &= \frac{s + m_2^2 - m_1^2}{2\sqrt{s}} \quad E_3 &= \frac{s + m_3^2 - m_4^2}{2\sqrt{s}} \quad E_4 &= \frac{s + m_4^2 - m_3^2}{2\sqrt{s}} \end{split}$$

$$s = (p_1 + p_2)^2 = (p_3 + p_4)^2 \approx +2p_1 \cdot p_2 \approx +2p_3 \cdot p_4 \qquad \sqrt{s} = E_{CM}$$

$$t = (p_1 - p_3)^2 = (p_2 - p_4)^2 \approx -2p_1 \cdot p_3 \approx -2p_2 \cdot p_4 \qquad s + t + u = m_1^2 + m_2^2 + m_3^2 + m_4^2$$

$$u = (p_1 - p_4)^2 = (p_2 - p_3)^2 \approx -2p_1 \cdot p_4 \approx -2p_2 \cdot p_3$$

$$f + \overline{f} \longrightarrow g + \overline{g} : \qquad p_1 = \frac{\sqrt{s}}{2} (1, 0, 0, +\beta_f) \qquad p_2 = \frac{\sqrt{s}}{2} (1, 0, 0, -\beta_f)$$
$$p_3 = \frac{\sqrt{s}}{2} (1, +\beta_g \sin \theta, 0, +\beta_g \cos \theta) \qquad p_4 = \frac{\sqrt{s}}{2} (1, -\beta_g \sin \theta, 0, -\beta_g \cos \theta)$$

$$m = 0$$
:  $s = 2|\vec{\mathbf{p}}_{\rm CM}|^2 = 2(p_{\rm CM}^0)^2$   $t = -\frac{s}{2}(1 - \cos\theta)$   $u = -\frac{s}{2}(1 + \cos\theta)$   $1 \, \text{kg} = 5.61 \times 10^{26} \, \text{GeV}$   $1 \, \text{m} = 55.07 \times 10^{15} \, \text{GeV}^{-1}$ 

#### Feynman Rules for i M

Goes in opposite way of arrows with the first one being adjoint,  $\bar{\psi} = \psi^{\dagger} \gamma^0$ :

#### Constants

$$\frac{G_F}{\sqrt{2}} = \frac{g_W^2}{8M_W^2} = \frac{g_W^2}{8M_Z^2 \cos^2(\theta_W)} = \frac{1}{2v^2} \qquad M_W = M_Z \cos \theta_W = \frac{1}{2} g_W v \qquad M_H = \sqrt{2\lambda} v$$

$$\begin{split} g_W &= g_Z \cos \theta_W = \frac{e}{\sin \theta_W} \qquad g_V^f = \frac{1}{2} T_f^3 - Q_f \sin^2(\theta_W) \qquad g_A^f = \frac{1}{2} T_f^3 \qquad g_{A,V} = \frac{1}{2} c_{A,V} \\ \alpha &= \frac{1}{4\pi\epsilon_0} \frac{e^2}{\hbar c} \qquad c_L^f = g_V^f + g_A^f \qquad c_R^f = g_V^f - g_A^f \qquad g_V^f = \frac{c_L^f + c_R^f}{2} \qquad g_A^f = \frac{c_L^f - c_R^f}{2} \end{split}$$

 $1 \text{ s} = 51.52 \times 10^{24} \,\text{GeV}^{-1}$ 

**Relativity & Quantum Mechanics** 

$$\beta = \frac{v}{c} = \frac{p}{E} = \tanh(\eta) \qquad \gamma = \frac{1}{\sqrt{1 - \beta^2}} = \cosh(\eta) \qquad \gamma \beta = \sin(\eta)$$

$$\cosh\left(\frac{\eta}{2}\right) = \sqrt{\frac{E + m}{2m}} \qquad \sinh\left(\frac{\eta}{2}\right) = \frac{|\vec{\mathbf{p}}|}{\sqrt{2m(E + m)}} \qquad p^{\mu} \to i\delta^{\mu}$$

$$E^2 = m^2 + |p|^2 \qquad \frac{\mathrm{d}^3\vec{\mathbf{p}}'}{E'} = \frac{\mathrm{d}^3\vec{\mathbf{p}}}{E} \qquad \widetilde{\mathrm{d}}p = \frac{\mathrm{d}^3\vec{\mathbf{p}}}{(2\pi)^3 2E_p} \qquad \widetilde{\delta}(p - q) = (2\pi)^3 2E_p \delta^3(\vec{\mathbf{p}} - \vec{\mathbf{q}})$$

$$x^{\mu} = (t, \vec{\mathbf{x}}) \qquad p^{\mu} = (E, \vec{\mathbf{p}}) \qquad \partial^{\mu} = \left(\partial_{t}, -\vec{\mathbf{V}}\right) \qquad A^{\mu} = \left(\phi, \vec{\mathbf{A}}\right) \qquad J^{\mu} = \left(\rho, \vec{\mathbf{J}}\right)$$

$$F^{\mu\nu} \equiv \partial^{\mu}A^{\nu} - \partial^{\nu}A^{\mu} \qquad \mathcal{F}^{\mu\nu} = \frac{1}{2}\epsilon^{\mu\nu\alpha\beta}F_{\alpha\beta} \qquad \partial_{\mu}F^{\mu\nu} = J^{\nu} \qquad \partial_{\mu}\mathcal{F}^{\mu\nu} = 0$$

$$\epsilon^{\mu\nu\rho\sigma} = \begin{cases} +1 & \text{If } (\mu\nu\rho\sigma) \text{ is an even permutation of } (0123) \\ -1 & \text{If } (\mu\nu\rho\sigma) \text{ is an odd permutation of } (0123) \\ 0 & \text{Otherwise} \end{cases}$$

$$\epsilon_{\alpha\beta_1\gamma_1\delta_1}\epsilon^{\alpha\beta_2\gamma_2\delta_2} = +g_{\beta_1}^{\beta_2}g_{\gamma_1}^{\delta_2}g_{\delta_1}^{\gamma_2} - g_{\beta_1}^{\beta_2}g_{\gamma_1}^{\gamma_2}g_{\delta_1}^{\delta_2} + g_{\beta_1}^{\gamma_2}g_{\gamma_1}^{\beta_2}g_{\delta_2}^{\delta_2} - g_{\beta_1}^{\gamma_2}g_{\gamma_1}^{\delta_2}g_{\delta_1}^{\beta_2} + g_{\beta_1}^{\delta_2}g_{\gamma_1}^{\gamma_2}g_{\delta_1}^{\delta_2} - g_{\beta_1}^{\gamma_2}g_{\gamma_1}^{\beta_2}g_{\delta_1}^{\gamma_2}g_{\delta_1}^{\delta_2} - g_{\beta_1}^{\gamma_2}g_{\gamma_1}^{\delta_2}g_{\delta_1}^{\gamma_2}g_{\delta_1}^{\delta_2} - g_{\beta_1}^{\gamma_2}g_{\delta_1}^{\delta_2} - g_{\beta$$

$$\begin{split} g_{\mu\nu}\Lambda^{\mu}_{\alpha}\Lambda^{\nu}_{,\beta} &= g_{\alpha\beta} \quad \Lambda^{\mu}_{\nu} = g^{\mu}_{\nu} + \omega^{\mu}_{\nu} = \delta^{\mu}_{\nu} + \omega^{\mu}_{\nu} \quad \omega_{\beta\alpha} = -\omega_{\alpha\beta} \quad \omega_{kl} = -\epsilon^{klm}\theta^{m} \quad \omega^{0k} = \eta^{klm}\theta^{m} \quad \omega^{0k}\theta^{m} \quad$$

$$\begin{aligned} (j_{-},j_{+}) &= (1/2,0) \qquad \vec{\mathbf{S}}_{-} &= \vec{\boldsymbol{\sigma}}/2 \quad \vec{\mathbf{S}}_{+} &= 0 \qquad \qquad \chi_{L} \to \Lambda_{L}\chi_{L} = \exp\{\left(-i\vec{\boldsymbol{\theta}} - \vec{\boldsymbol{\eta}}\right) \cdot \vec{\boldsymbol{\sigma}}/2\}\chi_{L} \\ (j_{-},j_{+}) &= (0,1/2) \qquad \vec{\mathbf{S}}_{+} &= \vec{\boldsymbol{\sigma}}/2 \quad \vec{\mathbf{S}}_{-} &= 0 \qquad \qquad \chi_{R} \to \Lambda_{R}\chi_{R} = \exp\{\left(-i\vec{\boldsymbol{\theta}} + \vec{\boldsymbol{\eta}}\right) \cdot \vec{\boldsymbol{\sigma}}/2\}\chi_{R} \\ \Lambda_{L}^{\dagger}\Lambda_{R} &= \Lambda_{R}^{\dagger}\Lambda_{L} &= 1 \qquad \sigma^{2}\Lambda_{L}^{*}\sigma^{2} = \Lambda_{R} \qquad \sigma^{2}\chi_{L}^{2} \to \Lambda_{R}(\sigma^{2}\chi_{L}^{*}) \end{aligned}$$

#### Mathematics

Dirac Delta:

$$\delta^{n}(x'-x) = i \int \frac{\mathrm{d}^{n}p}{(2\pi)^{n}} e^{-ip(x'-x)} \qquad f(x_{i}) = 0 \implies \delta[f(x)] = \sum_{i} \left| \frac{\mathrm{d}f}{\mathrm{d}x} \right|_{x_{i}}^{-1} \cdot \delta(x-x_{i})$$

Triangle Function:

$$\lambda(a,b,c) = \left[a - \left(\sqrt{b} + \sqrt{c}\right)^{2}\right] \left[a - \left(\sqrt{b} - \sqrt{c}\right)^{2}\right]$$

Residue Theorem:

$$\oint_{\Gamma} dz f(z) = \pm 2\pi i \sum_{k=1}^{n} \text{Res}[f(a_k)] \qquad \text{Res}[f(a_k)] = \lim_{z \to a_k} (z - a_k) f(z)$$

The "-" ("+") is used when  $\Gamma$  is oriented clockwise (counterclockwise).

$$f(z) = \frac{h(z)e^{-iz(t'-t)}}{(z-a_1)(z-a_2)\dots} \stackrel{t'-t<0}{\Longrightarrow} \int_{\Re(z)=-\infty}^{\Re(z)=+\infty} f(z) \, \mathrm{d}z = +2\pi i \sum_{a_k \text{ in UHP}} \mathrm{Res}[f(a_k)]$$

$$\stackrel{t'-t>0}{\Longrightarrow} \int_{\Re(z)=-\infty}^{\Re(z)=+\infty} f(z) \, \mathrm{d}z = -2\pi i \sum_{a_k \text{ in LHP}} \mathrm{Res}[f(a_k)]$$