## **Equations**

Schrödinger, Klein-Gordon (Real Scalar Field), Dirac (Complex Scalar Field):

$$\begin{split} i\hbar\frac{\partial|\psi(t)\rangle}{\partial t} &= H|\psi(t)\rangle \qquad \mathcal{L} = i\psi^*\dot{\psi} - \frac{1}{2m}\vec{\nabla}\psi^*\cdot\vec{\nabla}\psi \\ (\partial\cdot\partial + m^2)\phi &= 0 \qquad \mathcal{L} = \frac{1}{2}\Big[\partial_\mu\phi\partial^\mu\phi - m^2\phi^2\Big] = \frac{1}{2}\Big[\dot{\phi}^2 - \left(\vec{\nabla}\phi\right)^2 - m^2\phi^2\Big] \\ (i\not{\!\partial} - m)\varphi &= 0 \qquad \overline{\varphi}(i\not{\!\partial} + m) = 0 \qquad \mathcal{L} = \partial_\mu\varphi\partial^\mu\varphi^* - m^2\varphi\varphi^* \end{split}$$

#### Pauli Matrices

$$\sigma^{1} = \sigma_{x} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \quad \sigma^{2} = \sigma_{y} = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \quad \sigma^{3} = \sigma_{z} = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \quad \vec{\Sigma} = \begin{pmatrix} \vec{\sigma} & 0 \\ 0 & \vec{\sigma} \end{pmatrix}$$

$$\{\sigma^{i}, \sigma^{j}\} = 2\delta^{ij} \quad [\sigma^{i}, \sigma^{j}] = 2i\epsilon^{ijk}\sigma^{k} \quad \sigma^{i}\sigma^{j} = \delta^{ij} + i\epsilon^{ijk}\sigma^{k} \quad \sigma^{2}\sigma^{i}\sigma^{2} = -(\sigma^{i})^{*}$$

$$\sigma^{\mu} = (1, \vec{\sigma}) \quad \overline{\sigma}^{\mu} = (1, -\vec{\sigma}) \quad (p \cdot \sigma)(p \cdot \overline{\sigma}) = p^{2} \quad (\vec{\sigma} \cdot \vec{\mathbf{p}})^{2} = |\vec{\mathbf{p}}|^{2}$$

$$p \cdot \overline{\sigma} = \begin{pmatrix} p^{0} + p^{3} & p^{1} - ip^{2} \\ p^{1} + ip^{2} & p^{0} - p^{3} \end{pmatrix} \quad p \cdot \sigma = \begin{pmatrix} p^{0} - p^{3} & -(p^{1} - ip^{2}) \\ -(p^{1} + ip^{2}) & p^{0} + p^{3} \end{pmatrix}$$

$$\sqrt{p \cdot \sigma} = \frac{E + m - \vec{\sigma} \cdot \vec{\mathbf{p}}}{\sqrt{2(E + m)}} \quad \sqrt{p \cdot \overline{\sigma}} = \frac{E + m + \vec{\sigma} \cdot \vec{\mathbf{p}}}{\sqrt{2(E + m)}} \quad \vec{\sigma} \cdot \vec{\mathbf{p}} = \begin{pmatrix} p_{z} & p_{x} - ip_{y} \\ p_{x} + ip_{y} & -p_{z} \end{pmatrix}$$

## Dirac y-Matrices

$$\begin{split} \gamma^{\mu} &= \begin{pmatrix} 0 & \sigma^{\mu} \\ \overline{\sigma}^{\mu} & 0 \end{pmatrix} \quad \text{<- Chiral } | \, \text{Pauli-Dirac } \text{->} \quad \gamma^0 = \begin{pmatrix} +1 & 0 \\ 0 & -1 \end{pmatrix} \quad \gamma^k = \begin{pmatrix} 0 & \sigma^k \\ -\sigma^k & 0 \end{pmatrix} \\ \{\gamma^{\mu}, \gamma^{\nu}\} &= 2g^{\mu\nu} \quad \gamma_5 = +i\gamma^0\gamma^1\gamma^2\gamma^3 = -i\gamma_0\gamma_1\gamma_2\gamma_3 = -\frac{i}{4!}\epsilon^{\alpha\beta\gamma\delta}\gamma^{\alpha}\gamma^{\beta}\gamma^{\gamma}\gamma^{\delta} \\ \sigma^{\mu\nu} &= \frac{i}{2}[\gamma^{\mu}, \gamma^{\nu}] \implies [\gamma^{\mu}, \gamma^{\nu}] = -2i\sigma^{\mu\nu} \quad \{\gamma_5, \gamma^{\mu}\} = 0 \qquad \gamma^0(\gamma^{\mu})^{\dagger}\gamma^0 = \gamma^{\mu} \\ (\gamma^0)^{\dagger} &= \gamma^0 \quad (\gamma^0)^2 = 1 \quad \left(\gamma^k\right)^{\dagger} = -\gamma^k \quad \left(\gamma^k\right)^2 = -1 \quad \left(\gamma^5\right)^{\dagger} = \gamma^5 \quad \left(\gamma^5\right)^2 = 1 \\ \overline{\Gamma} &= \gamma^0\Gamma^{\dagger}\gamma^0 \quad \overline{\gamma_5} = -\gamma_5 \quad \overline{\gamma^{\mu}} = \gamma^{\mu} \quad \overline{\gamma^{\mu}\gamma_5} = \gamma^{\mu}\gamma_5 \quad \overline{\sigma^{\mu\nu}} = \sigma^{\mu\nu} \end{split}$$

$$\begin{split} & \not\! A = A^\mu \gamma_\mu \qquad \gamma^\mu \gamma_\mu = 4 \qquad \gamma^\mu \gamma^\nu \gamma_\mu = -2 \gamma^\nu \qquad \gamma^\mu \gamma^\nu \gamma^\rho \gamma_\mu = 4 g^{\nu\rho} \\ & \gamma^\mu \gamma^\nu \gamma^\rho \gamma^\sigma \gamma_\mu = -2 \gamma^\sigma \gamma^\rho \gamma^\nu \qquad \gamma^\mu \gamma^\nu \gamma^\rho \gamma^\sigma \gamma^\pi \gamma_\mu = 2 [\gamma^\pi \gamma^\nu \gamma^\rho \gamma^\sigma + \gamma^\sigma \gamma^\rho \gamma^\nu \gamma^\pi] \\ & \qquad \qquad \gamma^\mu \gamma^\nu \gamma^\rho = i \epsilon^{\mu\nu\rho\lambda} \gamma_\lambda \gamma_5 + g^{\mu\nu} \gamma^\rho - g^{\mu\rho} \gamma^\nu + g^{\nu\rho} \gamma^\mu \end{split}$$

## Spin, Helicity and Chirality

$$\vec{\mathbf{S}} = \frac{1}{2}\vec{\mathbf{\Sigma}} \qquad \tilde{h} = \frac{\vec{\mathbf{S}} \cdot \vec{\mathbf{p}}}{|\vec{\mathbf{p}}|} = \frac{\vec{\mathbf{Z}} \cdot \vec{\mathbf{p}}}{2|\vec{\mathbf{p}}|} \qquad \tilde{h} = 2h = \frac{\vec{\mathbf{Z}} \cdot \vec{\mathbf{p}}}{|\vec{\mathbf{p}}|} \qquad P_L = \frac{1-h}{2} \qquad P_R = \frac{1+h}{2}$$

$$\gamma_L = \frac{1-\gamma_5}{2} \qquad \gamma_R = \frac{1+\gamma_5}{2} \qquad \gamma_{L,R}^2 = \gamma_{L,R} \qquad \gamma_{L,R}\gamma_{R,L} = 0 \qquad \gamma_L + \gamma_R = 1$$

$$\gamma^\mu \gamma_{L,R} = \gamma_{R,L} \gamma^\mu \qquad \gamma_{SYL} = \gamma_L \gamma_S = -\gamma_L \qquad \gamma_{SYR} = \gamma_R \gamma_S = \gamma_R \qquad \gamma_{L,R}^{\dagger} = \gamma_{L,R}$$

$$\overline{\gamma^\mu \gamma_{L,R}} = \gamma^\mu \gamma_{L,R} \qquad \overline{\gamma_{SYL}} = -\gamma_R \qquad \overline{\gamma_{SYR}} = \gamma_L \qquad \overline{\gamma_{L,R}} = \gamma_0 \gamma_{L,R}^{\dagger} \gamma_0 = \gamma_{R,L}$$

$$\begin{aligned} u_L &= u_\downarrow & u_R &= u_\uparrow & v_L &= v_\uparrow & v_R &= v_\downarrow \\ P_L u_\downarrow &= u_\downarrow & P_L u_\uparrow &= 0 & P_R u_\downarrow &= 0 & P_R u_\uparrow &= u_\uparrow & \overline{u_\downarrow} P_L &= 0 & \overline{u_\uparrow} P_L &= u_\uparrow & \overline{u_\downarrow} P_R &= u_\downarrow & \overline{u_\uparrow} P_R &= 0 \\ \gamma_L u_L &= u_L & \gamma_L u_R &= 0 & \gamma_R u_L &= 0 & \gamma_R u_R &= u_R & \overline{u_L} \gamma_L &= 0 & \overline{u_R} \gamma_L &= u_R & \overline{u_L} \gamma_R &= u_L & \overline{u_R} \gamma_R &= 0 \\ P_L v_\downarrow &= 0 & P_L v_\uparrow &= v_\uparrow & P_R v_\downarrow &= v_\downarrow & P_R v_\uparrow &= 0 & \overline{v_\downarrow} P_L &= v_\downarrow & \overline{v_\uparrow} P_L &= 0 & \overline{v_\downarrow} P_R &= 0 & \overline{v_\uparrow} P_R &= v_\uparrow \\ \gamma_L v_L &= v_L & \gamma_L v_R &= 0 & \gamma_R v_L &= 0 & \gamma_R v_R &= v_R & \overline{v_L} \gamma_L &= 0 & \overline{v_R} \gamma_L &= v_R & \overline{v_L} \gamma_R &= v_L & \overline{v_R} \gamma_R &= 0 \end{aligned}$$

$$\begin{aligned} u_{\uparrow}^T &= N \left( c \quad s e^{i \phi} \quad k c \quad k s e^{i \phi} \right) \approx N \left( c \quad s e^{i \phi} \quad c \quad s e^{i \phi} \right) & N = \sqrt{E + m} \approx \sqrt{E} \\ u_{\downarrow}^T &= N \left( -s \quad c e^{i \phi} \quad k s \quad -k c e^{i \phi} \right) \approx N \left( -s \quad c e^{i \phi} \quad s \quad -c e^{i \phi} \right) & k = \frac{p}{E + m} \approx 1 \\ v_{\uparrow}^T &= N \left( k s \quad -k c e^{i \phi} \quad -s \quad c e^{i \phi} \right) \approx N \left( s \quad -c e^{i \phi} \quad -s \quad c e^{i \phi} \right) & s = \sin \left( \frac{\theta}{2} \right) \\ v_{\downarrow}^T &= N \left( k c \quad k s e^{i \phi} \quad c \quad s e^{i \phi} \right) \approx N \left( c \quad s e^{i \phi} \quad c \quad s e^{i \phi} \right) & c = \cos \left( \frac{\theta}{2} \right) \\ \vec{\mathbf{p}} &= p (\sin \theta \cos \phi, \sin \theta \sin \phi, \cos \theta) \end{aligned}$$

#### Traces

Traces with an odd number of  $\gamma$  matrices are 0.

$$\begin{split} \operatorname{Tr}[1] &= 4 & \operatorname{Tr}[\gamma^{\mu}\gamma^{\nu}] = 4g^{\mu\nu} & \operatorname{Tr}[\gamma^{\mu}\gamma^{\nu}\gamma^{\rho}\gamma^{\sigma}] = 4(g^{\mu\nu}g^{\rho\sigma} - g^{\mu\rho}g^{\nu\sigma} + g^{\mu\sigma}g^{\nu\rho}) \\ \operatorname{Tr}[\gamma_{5}] &= 0 & \operatorname{Tr}[\gamma_{5}\gamma^{\mu}\gamma^{\nu}] = 0 & \operatorname{Tr}[\gamma_{5}\gamma^{\mu}\gamma^{\nu}\gamma^{\rho}\gamma^{\sigma}] = -4i\epsilon^{\mu\nu\rho\sigma} \\ & \operatorname{Tr}[\gamma^{\mu}\gamma^{\nu}\gamma^{\rho}\gamma^{\sigma}\chi_{L}] = 2(g^{\mu\nu}g^{\rho\sigma} - g^{\mu\rho}g^{\nu\sigma} + g^{\mu\sigma}g^{\nu\rho} + i\epsilon^{\mu\nu\rho\sigma}) \\ & \operatorname{Tr}[\gamma^{\mu}\gamma^{\nu}\gamma^{\rho}\gamma^{\sigma}\gamma_{R}] = 2(g^{\mu\nu}g^{\rho\sigma} - g^{\mu\rho}g^{\nu\sigma} + g^{\mu\sigma}g^{\nu\rho} - i\epsilon^{\mu\nu\rho\sigma}) \\ & \operatorname{Tr}[y^{\mu}\gamma_{L}y^{\mu}\gamma_{L}y^{\nu}\gamma^{\nu}(g_{V} - g_{A}\gamma_{5})y_{A}\gamma_{\mu}\gamma_{L}y_{3}\gamma_{\nu}\gamma_{L}] = -16(g_{V} + g_{A})(p_{1} \cdot p_{4})(p_{2} \cdot p_{3}) \\ & \operatorname{Tr}[(y_{A} + m_{a})\gamma^{\mu}(y_{b} + m_{b})\gamma^{\nu}] = 4\left[p_{A}^{\mu}p_{b}^{\nu} + p_{b}^{\mu}p_{a}^{\nu} + (m_{a}m_{b} - p_{a} \cdot p_{b})g^{\mu\nu}\right] \\ & \operatorname{Tr}[\phi\gamma^{\mu}\beta\gamma^{\nu}\gamma_{L}]\operatorname{Tr}[\phi\gamma_{\mu}\phi\gamma_{\nu}\gamma_{5}] = 0 & \operatorname{Tr}[\phi\gamma^{\mu}\beta\gamma^{\nu}\gamma_{L}]\operatorname{Tr}[\phi\gamma_{\mu}\phi\gamma_{\nu}\gamma_{R}] = 16(a \cdot d)(b \cdot c) \\ & \operatorname{Tr}[\phi\gamma^{\mu}\beta\gamma^{\nu}\gamma_{L}]\operatorname{Tr}[\phi\gamma_{\mu}\phi\gamma_{\nu}\gamma_{L}]\operatorname{Tr}[\phi\gamma^{\mu}\phi\gamma^{\nu}\gamma_{L}]\operatorname{Tr}[\phi\gamma^{\mu}\phi\gamma_{\nu}\gamma_{L}] = 16(a \cdot c)(b \cdot d) \end{split}$$

# Spinors (Fermions)

$$\psi = ue^{+i(\vec{\mathbf{p}}\cdot\vec{\mathbf{x}}-Et)} = ve^{-i(\vec{\mathbf{p}}\cdot\vec{\mathbf{x}}-Et)} \qquad \overline{\psi} = \psi^{\dagger}\gamma^{0}$$

$$(\not{p}-m)u = \overline{u}(\not{p}-m) = 0 \qquad (\not{p}+m)v = \overline{v}(\not{p}+m) = 0$$

$$\overline{u}'(p)u^{s}(p) = +2m\delta^{rs} \qquad \overline{u}'(p)v^{s}(p) = 0 \qquad u^{r\dagger}(p)u^{s}(p) = 2E\delta^{rs} \qquad \sum_{s=1,2} u^{s}(p)\overline{u}^{s}(p) = \not{p}+m$$

$$\overline{v}'(p)v^{s}(p) = -2m\delta^{rs} \qquad \overline{v}'(p)u^{s}(p) = 0 \qquad v^{r\dagger}(p)v^{s}(p) = 2E\delta^{rs} \qquad \sum_{s=1,2} v^{s}(p)\overline{v}^{s}(p) = \not{p}-m$$

## **Polarization Vectors (Bosons)**

External massless & Massive:

$$\sum_{\lambda=1}^{2} \epsilon^{\mu}(k,\lambda) \cdot \epsilon^{*\nu}(k,\lambda) = -g^{\mu\nu} \qquad \sum_{\lambda=1}^{3} \epsilon_{\mu}(q,\lambda) \epsilon_{\nu}^{*}(q,\lambda) = -g_{\mu\nu} + \frac{q_{\mu}q_{\nu}}{M^{2}}$$

## Quantization: Real Scalar Field (Bosons)

$$: \frac{1}{2} \left[ a^{\dagger}(k)a(k) + a(k)a^{\dagger}(k) \right] := a^{\dagger}(k)a(k) \qquad : a(x)b(x') := a(x)b(x') + b(x')a(x)$$

$$T(a(x)b(x')) = \theta(t - t')a(x)b(x') + \theta(t' - t)b(x')a(x)$$

The "+" corresponds to positive frequency plane waves  $e^{-ik \cdot x}$ :

$$\begin{split} \mathcal{L} &= \frac{1}{2} \partial^{\mu} \phi \partial_{\mu} \phi - \frac{1}{2} m^{2} \phi \phi \qquad \pi = \dot{\phi} \qquad \left[ a^{\dagger}(k,\lambda), a(k',\lambda') \right] = g^{\lambda \lambda'} \, \tilde{\delta}(k-k') \\ \left[ \phi(\vec{\mathbf{x}},t), \phi(\vec{\mathbf{y}},t) \right] &= \left[ \Pi(\vec{\mathbf{x}},t), \Pi(\vec{\mathbf{y}},t) \right] = 0 \qquad \left[ \phi(\vec{\mathbf{x}},t), \Pi(\vec{\mathbf{y}},t) \right] = i \delta^{3} (\vec{\mathbf{x}} - \vec{\mathbf{y}}) \\ \phi(x) &= \int \widetilde{dk} \left[ a(k) e^{-ik \cdot x} + a^{\dagger}(k) e^{+ik \cdot x} \right] \\ |p\rangle &= a^{\dagger}(p) |0\rangle \qquad |p_{1}, p_{2}\rangle = a^{\dagger}(p_{2}) a^{\dagger}(p_{1}) |0\rangle = |p_{2}, p_{1}\rangle \end{split}$$

## **Quantization : Complex Scalar Field (Bosons)**

$$\begin{split} \mathcal{L} &= : \partial^{\mu} \varphi^{\dagger} \partial_{\mu} \varphi - m^{2} \varphi^{\dagger} \varphi \colon \qquad \pi = \dot{\varphi}^{\dagger} \qquad \pi^{\dagger} = \dot{\varphi} \qquad \left[ a_{\pm}(k), a_{\pm}^{\dagger}(k') \right] = \tilde{\delta}(k - k') \\ & \left[ \varphi(\vec{\mathbf{x}}, t), \pi(\vec{\mathbf{y}}, t) \right] = \left[ \varphi^{\dagger}(\vec{\mathbf{x}}, t), \pi^{\dagger}(\vec{\mathbf{y}}, t) \right] = i \, \delta^{3}(\vec{\mathbf{x}} - \vec{\mathbf{y}}) \\ & \varphi(x) = \varphi^{+}(x) + \varphi^{-}(x) = \int \widetilde{\mathrm{d}} k \Big[ a_{+}(k) e^{-ik \cdot x} + a_{-}^{\dagger}(k) e^{+ik \cdot x} \Big] \\ & \varphi^{\dagger}(x) = \varphi^{+\dagger}(x) + \varphi^{-\dagger}(x) = \int \widetilde{\mathrm{d}} k \Big[ a_{-}(k) e^{-ik \cdot x} + a_{+}^{\dagger}(k) e^{+ik \cdot x} \Big] \\ & \left| p^{+} \right\rangle = a_{+}^{\dagger}(p) \left| 0 \right\rangle \qquad \left| p^{-} \right\rangle = a_{-}^{\dagger}(p) \left| 0 \right\rangle \qquad \left| p^{+}_{1}, p^{-}_{2} \right\rangle = a_{+}^{\dagger}(p_{2}) a_{+}^{\dagger}(p_{1}) \left| 0 \right\rangle \end{split}$$

### **Quantization**: Dirac Field (Fermions)

$$: a(x)b(x') := a(x)b(x') - b(x')a(x)$$

$$T(a(x)b(x')) = \theta(t-t)a(x)b(x') - \theta(t'-t)b(x')a(x)$$

$$\mathcal{L} = i\bar{\psi}\gamma^{\mu}\partial_{\mu}\psi - m\bar{\psi}\psi: \qquad \pi_{\alpha} = \frac{\partial \mathcal{L}}{\partial \dot{\psi}_{\alpha}} = i\psi_{\alpha}^{\dagger} \qquad \pi_{\alpha}^{\dagger} = \frac{\partial \mathcal{L}}{\partial \dot{\psi}_{\alpha}^{\dagger}} = 0$$

$$\left\{b^{\dagger}(p,s), b(p',s')\right\} = \left\{d^{\dagger}(p,s), d(p',s')\right\} = \tilde{\delta}(p-p')\delta_{ss'}$$

$$\psi(x) = \psi^{+}(x) + \psi^{-}(x) = \int \widetilde{dp} \sum_{s} \left[b(p,s)u(p,s)e^{-ip\cdot x} + d^{\dagger}(p,s)v(p,s)e^{+ip\cdot x}\right]$$

$$\bar{\psi}(x) = \bar{\psi}^{\dagger}(x) + \bar{\psi}^{-}(x) = \int \widetilde{dp} \sum_{s} \left[d(p,s)\bar{v}(p,s)e^{-ip\cdot x} + b^{\dagger}(p,s)\bar{u}(p,s)e^{+ip\cdot x}\right]$$

$$|e^{-}(p_{1},s_{1}), e^{-}(p_{2},s_{2})\rangle = b^{\dagger}(p_{2},s_{2})b^{\dagger}(p_{1},s_{1})|0\rangle = -|e^{-}(p_{2},s_{2}), e^{-}(p_{1},s_{1})\rangle$$

$$\psi^{+}(x)|p\rangle = \psi^{+}(x)b^{\dagger}(p)|0\rangle = |0\rangle u(p)e^{-ip\cdot x} \qquad \psi^{+}(b) \text{ destroys } e^{-}$$

$$\bar{\psi}^{\dagger}(x)|p\rangle = \bar{\psi}^{\dagger}(x)d^{\dagger}(p)|0\rangle = |0\rangle \bar{v}(p)e^{-ip\cdot x} \qquad \bar{\psi}^{\dagger}(d) \text{ destroy } e^{+}$$

$$\langle p|\psi^{-}(x) = \langle 0|d(p)\psi^{-}(x) = v(p)e^{+ip\cdot x}\langle 0| \qquad \psi^{-}(d^{\dagger}) \text{ creates } e^{+}$$

$$\langle p|\bar{\psi}^{-}(x) = \langle 0|b(p)\bar{\psi}^{-}(x) = \bar{u}(p)e^{+ip\cdot x}\langle 0| \qquad \bar{\psi}^{-}(b^{\dagger}) \text{ creates } e^{-}$$

## **Quantization : Electromagnetic Field**

Ward Identity and U(1) gauge invariante:  $A_{\mu} \to A_{\mu} + \partial_{\mu} \Lambda \implies \epsilon_{\mu} \to \epsilon_{\mu} + ck_{\mu}$   $\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} - \frac{1}{2\xi} (\partial \cdot A)^2 \qquad \pi^{\mu} = F^{\mu 0} - \frac{g^{\mu 0}}{\xi} (\partial \cdot A) \qquad \pi^{0} = -\frac{1}{\xi} (\partial \cdot A) \qquad \pi^{k} = E^{k}$   $\left[ A_{\mu}(\vec{\mathbf{x}}, t), A_{\nu}(\vec{\mathbf{y}}, t) \right] = \left[ \dot{A}_{\mu}(\vec{\mathbf{x}}, t), \dot{A}_{\nu}(\vec{\mathbf{y}}, t) \right] = 0 \qquad \left[ \dot{A}_{\mu}(\vec{\mathbf{y}}, t), A_{\nu}(\vec{\mathbf{x}}, t) \right] = ig_{\mu\nu} \delta^{3}(\vec{\mathbf{x}} - \vec{\mathbf{y}})$   $A_{\mu}(x) = A_{\mu}^{+}(x) + A_{\mu}^{-}(x) = \int \widetilde{dk} \sum_{\lambda=0}^{3} \left[ a(k, \lambda) \epsilon_{\mu}(k, \lambda) e^{-ik \cdot x} + a^{\dagger}(k, \lambda) \epsilon_{\mu}^{*}(k, \lambda) e^{+ik \cdot x} \right]$   $A_{\mu}^{+}(x) |k\rangle = A_{\mu}^{+}(x) a^{\dagger}(k) |0\rangle = |0\rangle \epsilon_{\mu}(k) e^{-ik \cdot x} \qquad \left[ \dot{a}(k, \lambda), a^{\dagger}(k', \lambda') \right] = -g^{\lambda\lambda'} \, \tilde{\delta}(k - k')$   $\langle k|A_{\mu}^{-}(x) = \langle 0|a(k)A_{\mu}^{-}(x) = \epsilon_{\mu}^{*}(k) e^{+ik \cdot x} \langle 0| \qquad A^{+}(a) \text{ destroys } \gamma \qquad A^{-}(a^{\dagger}) \text{ creates } \gamma$ 

## **Classical Field Theory**

$$\begin{split} S &= \int_{t_i}^{t_f} L \, \mathrm{d}t = \int \mathrm{d}^4 x \, \mathscr{L} \qquad L = \int \mathrm{d}^3 \vec{\mathbf{x}} \, \mathscr{L}(\phi, \partial_\mu \phi) \qquad \frac{\partial \mathscr{L}}{\partial \phi_i} - \partial_\mu \frac{\partial \mathscr{L}}{\partial \left(\partial_\mu \phi_i\right)} = 0 \\ H &= \int \mathrm{d}^3 \vec{\mathbf{x}} \, \mathscr{H} \qquad \mathscr{H} = \sum_i \Pi_i(x) \partial_0 \phi_i(x) - \mathscr{L} \qquad \Pi_i(x) = \frac{\partial \mathscr{L}}{\partial \left(\partial_0 \phi_i(x)\right)} \\ \dot{\psi} &= \frac{\partial \mathscr{H}}{\partial \Pi} - \vec{\nabla} \cdot \frac{\partial \mathscr{H}}{\partial (\vec{\mathbf{V}}\Pi)} \qquad \dot{\Pi} = -\frac{\partial \mathscr{H}}{\partial \psi} + \vec{\nabla} \cdot \frac{\partial \mathscr{H}}{\partial (\vec{\mathbf{V}}\psi)} \\ \delta \mathscr{L} &= \mathscr{L} - \mathscr{L} = \partial_\mu C^\mu \qquad J^\mu = C^\mu - \frac{\partial \mathscr{L}}{\partial \partial_\mu \phi} \delta \phi \qquad \partial_\mu J^\mu = \frac{\partial J^0}{\partial t} + \vec{\nabla} \cdot \vec{\mathbf{J}} = 0 \\ Q &= \int \mathrm{d}^3 \vec{\mathbf{x}} \, J^0 \qquad \frac{\mathrm{d} Q}{\mathrm{d} t} = 0 \qquad A^\mu \to A^\mu + \partial^\mu \lambda \end{split}$$

Ward identity:  $k_{\mu} \mathcal{M}^{\mu} = 0$ , where  $\mathcal{M} = \epsilon_{\mu} \mathcal{M}^{\mu}$ . Photon polarizations  $\epsilon$  parallel to its direction of propagation don't contribute to the scattering amplitude.

# Dyson Expansion, Propagators and Wick's Theorem

$$S = \sum_{n=0}^{\infty} \frac{(-i)^n}{n!} \int d^4x_1 d^4x_2 \dots d^4x_n \operatorname{T}(\mathcal{X}_{int}(x_1) \mathcal{X}_{int}(x_2) \dots \mathcal{X}_{int}(x_n)) \qquad \sum_{f} \left| S_{fi} \right|^2 = 1$$

$$\Delta_F(x_1 - x_2) = \overrightarrow{\phi(x_1)} \phi(x_2) = \langle 0|\operatorname{T}(\phi(x_1)\phi(x_2))|0 \rangle \qquad D_F^{\mu\nu} = \overrightarrow{A^{\mu}(x_1)} \overrightarrow{A^{\nu}}(x_2)$$

$$\Delta_F(x_1 - x_2) = \overrightarrow{\phi(x_1)} \overrightarrow{\phi^{\dagger}}(x_2) = \overrightarrow{\phi^{\dagger}}(x_2) \overrightarrow{\phi}(x_1) \qquad S_{F\alpha\beta} = \overrightarrow{\psi_{\alpha}(x_1)} \overrightarrow{\psi_{\beta}}(x_2) = -\overrightarrow{\psi_{\beta}}(x_2) \overrightarrow{\psi_{\alpha}}(x_1)$$

$$T(ABCD...WXYZ) = :ABCD...WXYZ: + :ABCD...WXYZ: + :ABCD...WXYZ: + :ABCD...WXYZ: + ...$$

$$+ \cdots + :ABCD...WXYZ: + :ABCD...WXYZ: + \cdots + :ABCD...WXYZ: + ...$$

## **Decay: Decay Rates**

$$\frac{\mathrm{d}\Gamma}{\mathrm{d}\Omega} = \frac{1}{32\pi^2} \frac{|\vec{\mathbf{p}}_{\mathrm{CM}}|}{m_0^2} |\vec{\mathcal{M}}|^2 \qquad |\vec{\mathbf{p}}_{\mathrm{CM}}| = \frac{1}{2m_0} \sqrt{\left[m_0^2 - (m_1 + m_2)^2\right] \left[m_0^2 - (m_1 - m_2)^2\right]}$$

$$E_1 = \frac{m_0^2 + m_1^2 - m_2^2}{2m_0} \qquad E_2 = \frac{m_0^2 + m_2^2 - m_1^2}{2m_0}$$

# **Scattering: Cross Sections**

$$\frac{d\sigma}{d\Omega} = \frac{1}{64\pi^2 s} \frac{|\mathbf{p}_3|}{|\mathbf{p}_1|} |\mathcal{M}|^2 \qquad |\mathbf{p}_a| = |\mathbf{p}_b| = \frac{1}{2\sqrt{s}} \sqrt{\left[s - (m_a + m_b)^2\right] \left[s - (m_a - m_b)^2\right]}$$

$$E_1 = \frac{s + m_1^2 - m_2^2}{2\sqrt{s}} \quad E_2 = \frac{s + m_2^2 - m_1^2}{2\sqrt{s}} \quad E_3 = \frac{s + m_3^2 - m_4^2}{2\sqrt{s}} \quad E_4 = \frac{s + m_4^2 - m_3^2}{2\sqrt{s}}$$

$$s = (p_1 + p_2)^2 = (p_3 + p_4)^2 \approx +2p_1 \cdot p_2 \approx +2p_3 \cdot p_4 \qquad \sqrt{s} = E_{CM}$$

$$t = (p_1 - p_3)^2 = (p_2 - p_4)^2 \approx -2p_1 \cdot p_3 \approx -2p_2 \cdot p_4 \qquad s + t + u = m_1^2 + m_2^2 + m_3^2 + m_4^2$$

$$u = (p_1 - p_4)^2 = (p_2 - p_3)^2 \approx -2p_1 \cdot p_4 \approx -2p_2 \cdot p_3$$

$$f + \overline{f} \longrightarrow g + \overline{g} : \qquad p_1 = \frac{\sqrt{s}}{2} (1, 0, 0, +\beta_f) \qquad p_2 = \frac{\sqrt{s}}{2} (1, 0, 0, -\beta_f)$$
$$p_3 = \frac{\sqrt{s}}{2} (1, +\beta_g \sin \theta, 0, +\beta_g \cos \theta) \qquad p_4 = \frac{\sqrt{s}}{2} (1, -\beta_g \sin \theta, 0, -\beta_g \cos \theta)$$

$$m = 0$$
:  $s = 2|\vec{\mathbf{p}}_{\rm CM}|^2 = 2(p_{\rm CM}^0)^2$   $t = -\frac{s}{2}(1 - \cos\theta)$   $u = -\frac{s}{2}(1 + \cos\theta)$   $1 \text{ kg} = 5.61 \times 10^{26} \text{ GeV}$ 

## Feynman Rules for i M

Goes in opposite way of arrows with the first one being adjoint,  $\bar{\psi} = \psi^{\dagger} \gamma^0$ :

$$\begin{array}{c} \mathbf{f} & \mathbf{u} & \bullet & \mathbf{v} & \mathbf{f} \\ \mathbf{f} & \mathbf{v} & \bullet & \mathbf{v} & \mathbf{f} \\ \mathbf{f} & \mathbf{v} & \bullet & \mathbf{v} & \mathbf{f} \\ \mathbf{f} & \mathbf{v} & \bullet & \mathbf{v} & \mathbf{f} \\ \mathbf{f} & \mathbf{v} & \bullet & \mathbf{v} & \mathbf{f} \\ \mathbf{f} & \mathbf{v} & \bullet & \mathbf{v} & \mathbf{f} \\ \mathbf{f} & \mathbf{v} & \bullet & \mathbf{v} & \mathbf{f} \\ \mathbf{f} & \mathbf{f} & \mathbf{v} & \bullet & \mathbf{v} & \mathbf{f} \\ \mathbf{f} & \mathbf{f} & \mathbf{f} & \mathbf{f} & \mathbf{f} \\ \mathbf{f} & \mathbf{f} & \mathbf{f} & \mathbf{f} & \mathbf{f} \\ \mathbf{f} & \mathbf{f} & \mathbf{f} & \mathbf{f} & \mathbf{f} \\ \mathbf{f} & \mathbf{f} & \mathbf{f} & \mathbf{f} & \mathbf{f} \\ \mathbf{f} & \mathbf{f}$$

 $\int ig_Z M_Z g_{\mu\nu}$ 

$$\begin{split} \frac{G_F}{\sqrt{2}} &= \frac{g_W^2}{8M_W^2} = \frac{g_W^2}{8M_Z^2 \cos^2(\theta_W)} = \frac{1}{2v^2} \qquad M_W = M_Z \cos\theta_W = \frac{1}{2}g_W v \qquad M_H = \sqrt{2\lambda}v \\ g_W &= g_Z \cos\theta_W = \frac{e}{\sin\theta_W} \qquad g_V^f = \frac{1}{2}T_f^3 - Q_f \sin^2(\theta_W) \qquad g_A^f = \frac{1}{2}T_f^3 \qquad g_{A,V} = \frac{1}{2}c_{A,V} \\ \alpha &= \frac{1}{4\pi\epsilon_0} \frac{e^2}{\hbar c} \qquad c_L^f = g_V^f + g_A^f \qquad c_R^f = g_V^f - g_A^f \qquad g_V^f = \frac{c_L^f + c_R^f}{2} \qquad g_A^f = \frac{c_L^f - c_R^f}{2} \end{split}$$

Constants

Fermions	$Q_f$		$I_{W}^{(3)}$	$Y_L$	$Y_{j}$	?	$c_L$		$c_R$ $c$		'V	$c_A$
$\nu_e, \nu_\mu, \nu_\tau$		0	+1/2	-1	0		+1/2	0		+1/2		+1/2
e <sup>-</sup> ,μ <sup>-</sup> ,τ <sup>-</sup>	-	-1	-1/2	-1	-2		-0.27	+0.23		-0.04		-1/2
u, c, t	+2	2/3	+1/2	+1/3	+4/3		+0.35	-0.15		+0.19		+1/2
d, s, b	-:	1/3	-1/2	+1/3	-2/3		-0.42	+0.08		-0.35		-1/2
Particle	$W^{\pm}$	$Z^0$	$H^0$	$e^{\pm}$	$\mu^{\pm}$	τ	u	d	c	s	t	b
Mass (MeV)	80 379	91 188	125 100	0.511	105.7	1777	2.16	4.67	1270	93	172 760	4180
41 5 (4 40 <sup>2</sup> 6 C V 4 55 CF 40 <sup>1</sup> 5 C V - 1 4 54 F0 40 <sup>2</sup> 4 C V - 1												

**Relativity & Quantum Mechanics** 

$$\beta = \frac{v}{c} = \frac{p}{E} = \tanh(\eta) \qquad \gamma = \frac{1}{\sqrt{1 - \beta^2}} = \cosh(\eta) \qquad \gamma \beta = \sin(\eta)$$

$$\cosh\left(\frac{\eta}{2}\right) = \sqrt{\frac{E + m}{2m}} \qquad \sinh\left(\frac{\eta}{2}\right) = \frac{|\vec{\mathbf{p}}|}{\sqrt{2m(E + m)}} \qquad p^{\mu} \to i\delta^{\mu}$$

$$E^2 = m^2 + |p|^2 \qquad \frac{\mathrm{d}^3\vec{\mathbf{p}}'}{E'} = \frac{\mathrm{d}^3\vec{\mathbf{p}}}{E} \qquad \widetilde{\mathrm{d}}p \equiv \frac{\mathrm{d}^3\vec{\mathbf{p}}}{(2\pi)^3 2E_p} \qquad \widetilde{\delta}(p - q) \equiv (2\pi)^3 2E_p \delta^3(\vec{\mathbf{p}} - \vec{\mathbf{q}})$$

$$e^{\mu} = (t, \vec{\mathbf{x}})$$
  $p^{\mu} = (E, \vec{\mathbf{p}})$   $\partial^{\mu} = (\partial_t, -\vec{\mathbf{V}})$   $A^{\mu} = (\phi, \vec{\mathbf{A}})$   $J^{\mu} = (\rho, \vec{\mathbf{J}})$   
 $F^{\mu\nu} \equiv \partial^{\mu}A^{\nu} - \partial^{\nu}A^{\mu}$   $\mathcal{E}^{\mu\nu} = \frac{1}{2}\epsilon^{\mu\nu\alpha\beta}F_{\alpha\beta}$   $\partial_{\mu}F^{\mu\nu} = J^{\nu}$   $\partial_{\mu}\mathcal{E}^{\mu\nu} = 0$ 

$$\epsilon^{\mu\nu\rho\sigma} = \begin{cases} +1 & \text{If } (\mu\nu\rho\sigma) \text{ is an even permutation of } (0123) \\ -1 & \text{If } (\mu\nu\rho\sigma) \text{ is an odd permutation of } (0123) \\ 0 & \text{Otherwise} \end{cases}$$

$$\epsilon_{\alpha\beta_1\gamma_1\delta_1}\epsilon^{\alpha\beta_2\gamma_2\delta_2} = +g_{\beta_1}^{\beta_2}g_{\gamma_1}^{\delta_2}g_{\delta_1}^{\gamma_2} - g_{\beta_1}^{\beta_2}g_{\gamma_1}^{\gamma_2}g_{\delta_1}^{\delta_2} + g_{\beta_1}^{\gamma_2}g_{\delta_1}^{\beta_2} - g_{\beta_1}^{\gamma_2}g_{\gamma_1}^{\delta_2}g_{\delta_1}^{\delta_2} + g_{\beta_1}^{\delta_2}g_{\gamma_1}^{\gamma_2}g_{\delta_1}^{\delta_2} - g_{\beta_1}^{\delta_2}g_{\gamma_1}^{\delta_2}g_{\delta_1}^{\delta_2} - g_{\beta_1}^{\delta_2}g_{\gamma_1}^{\delta_2}g_{\gamma_1}^{\delta_2}g_{\delta_1}^{\delta_2} - g_{\beta_1}^{\delta_2}g_{\gamma_1}^{\delta_2}g_{\delta_1}^{\delta_2} - g_{\beta_1}^{\delta_2}$$

$$[J^{\mu\nu}, J^{\rho\sigma}] = i(g^{\nu\rho}J^{\mu\sigma} - g^{\nu\sigma}J^{\mu\rho} - g^{\mu\rho}J^{\nu\sigma} + g^{\mu\sigma}J^{\nu\rho}) \quad \vec{\mathbf{J}} = (J^{23}, J^{31}, J^{12})$$

$$J^{k} = \frac{1}{2}\epsilon^{klm}J^{lm} \quad K^{k} = J^{k0} \qquad J^{\pm k} = \frac{1}{2}(J^{k} \pm iK^{k}) \quad \vec{\mathbf{K}} = (J^{10}, J^{20}, J^{30})$$

$$[J^{i}, J^{j}] = i\epsilon^{ijk}J^{k} \quad [K^{i}, K^{j}] = -i\epsilon^{ijk}J^{k} \quad [K^{i}, J^{j}] = i\epsilon^{ijk}K^{k}$$

$$[J^{\pm i}, J^{\pm j}] = i\epsilon^{ijk}J^{\pm k} \quad [J^{i}_{+}, J^{j}_{-}] = 0$$

 $\delta V^{\alpha} \sim \omega_{.\beta}^{\alpha} V^{\beta} = -\frac{i}{2} \omega_{\mu\nu} (J^{\mu\nu})_{.\beta}^{\alpha} V^{\beta} \quad (J^{\mu\nu})_{.\beta}^{\alpha} = i (g^{\mu\alpha} g \beta^{\nu} - g^{\nu\alpha} g \beta^{\mu})$ 

$$\begin{aligned} (j_{-}, j_{+}) &= (1/2, 0) \qquad \vec{\mathbf{S}}_{-} &= \vec{\boldsymbol{\sigma}}/2 \quad \vec{\mathbf{S}}_{+} &= 0 \\ (j_{-}, j_{+}) &= (0, 1/2) \qquad \vec{\mathbf{S}}_{+} &= \vec{\boldsymbol{\sigma}}/2 \quad \vec{\mathbf{S}}_{-} &= 0 \\ \Lambda_{L}^{\dagger} \Lambda_{R} &= \Lambda_{R}^{\dagger} \Lambda_{L} &= 1 \end{aligned} \qquad \begin{aligned} \chi_{L} &\to \Lambda_{L} \chi_{L} &= \exp\left\{\left(-i\vec{\boldsymbol{\theta}} - \vec{\boldsymbol{\eta}}\right) \cdot \vec{\boldsymbol{\sigma}}/2\right\} \chi_{L} \\ \chi_{R} &\to \Lambda_{R} \chi_{R} &= \exp\left\{\left(-i\vec{\boldsymbol{\theta}} + \vec{\boldsymbol{\eta}}\right) \cdot \vec{\boldsymbol{\sigma}}/2\right\} \chi_{R} \end{aligned}$$

## Mathematics

Dirac Delta:

Direct Delta: 
$$\delta^{n}(x'-x) = i \int \frac{\mathrm{d}^{n} p}{\left(2\pi\right)^{n}} e^{-ip(x'-x)} \qquad f(x_{i}) = 0 \implies \delta[f(x)] = \sum_{i} \left| \frac{\mathrm{d} f}{\mathrm{d}x} \right|_{x_{i}}^{-1} \cdot \delta(x-x_{i})$$

Triangle Function:

$$\lambda(a, b, c) = \left[a - \left(\sqrt{b} + \sqrt{c}\right)^{2}\right] \left[a - \left(\sqrt{b} - \sqrt{c}\right)^{2}\right]$$

Residue Theorem:

$$\oint_{\Gamma} dz f(z) = \pm 2\pi i \sum_{k=1}^{n} \text{Res}[f(a_k)] \qquad \text{Res}[f(a_k)] = \lim_{z \to a_k} (z - a_k) f(z)$$

The "-" ("+") is used when  $\Gamma$  is oriented clockwise (counterclockwise).

$$f(z) = \frac{h(z)e^{-iz(t'-t)}}{(z-a_1)(z-a_2)\dots} \stackrel{t'-t<0}{\Longrightarrow} \int_{\Re(z)=-\infty}^{\Re(z)=+\infty} f(z) \, \mathrm{d}z = +2\pi i \sum_{a_k \text{ in UHP}} \mathrm{Res}[f(a_k)]$$

$$\stackrel{t'-t>0}{\Longrightarrow} \int_{\Re(z)=-\infty}^{\Re(z)=+\infty} f(z) \, \mathrm{d}z = -2\pi i \sum_{a_k \text{ in LHP}} \mathrm{Res}[f(a_k)]$$