

FAKULTÄT FÜR PHYSIK PHYSIKALISCHES PRAKTIKUM FÜR FORTGESCHRITTENE PRAKTIKUM MODERNE PHYSIK

Gruppe Nr	Kurs:	Mo	Mi	I	S 2016
Versuch:					
Namen:					
Assistent:					
durchgeführt am:					
Protokollabgabe am:					
Note gesamt +	0 -				
Anerkannt:	tum Unterschrift)			
Datum Rückgabe:					_
Bemerkung:					

1 Preparation

1.1 Theoretical Background

1.1.1 The Mößbauer Effect

When an atom at rest emits a photon, part of the released energy is used to satisfy momentum conservation and recoil the atom. Similarly when absorbing a photon, it has to carry the energy for the excitation and the recoil. Resonance absorption, where the absorber is of the same type as the emitter, can only occur, when the linewidth of the decay is larger than the recoil energy.

For visible light resonance absorption is possible, while it is impossible for nuclear γ -transitions.

In reality at finite Temperature of the samples the doppler effect increases the widths of the lines creating an overlap in emission and absorption spectra and enabling resonance absorption.

The Mößbauer effect occurs, when certain atoms (e.g. Fe^{57}) are built into a crystal structure. Atoms in the lattice can emit photons creating or consuming phonons, but there is a finite propability, that the internal energy of the crystal does not change.

The ratio of recoilless processes is given by the Debye-Waller-Factor:

$$f(T) = \exp(-\frac{p_{\gamma}^2}{2M} \cdot \frac{3}{2k_B\Theta}[1 + \frac{2\pi^2}{3}(\frac{T}{\Theta})^2])$$

where the first fraction gives the recoil the atom would receive, were it a free particle, the rest gives the expected value for the phonon spectrum in the Debye model with the Debye temperature Θ .

The Debye temperature of $\mathrm{Fe^{57}}$ is 480 K and the Debye-Waller factor at room temperature is 0.7.

1.1.2 Chemical Shift

Since the charge distribution of the nucleus is not pointlike, but has a radius dependant on excitation state, the coulomb interaction between nucleus and electron cloud changes the energy of the state and the transitions. The total shift is given as:

$$\nu = \frac{2\pi c}{3E_0} Ze^2 (|\psi_A(0)|^2 - |\psi_E(0)|^2) (R_e^2 - R_g^2)$$

Where A and E denote absorber and emitter and the ψ are the electronic wave function at the location of the nucleus. Since the wave function depends on the chemical environment of the atom the shift is only measured, when absorber and emitter nuclei are in different types of crystal.

1.1.3 Hyperfine Structure

Nuclear spin in the magnetic fields caused by the electron gas gives rise to zeeman splitting such that instead of finding a single resonance absorption peak there are multiple ones depending on the inner magnetic fields and the nuclear spin of the emitting and absorbing nuclei.

The Fe⁵⁷ emitter in this experiment is embedded in a crystal without inner magnetic fields in order to simplify the measurement.

Zeeman splitting of the absorber is given by:

$$\Delta E = -\frac{m}{I}\mu B$$

with magnetic quantum number m, nuclear spin quantum number I, magnetic dipole of the spin μ and inner magnetic field b.

For Fe⁵⁷ the ground state is split in two and the excited state in four levels. With the selection rules $\Delta m = \pm 1, 0$ there can be six different transitions.

1.1.4 Quadrupole Splitting

Nuclei with a spin I can also have an asymeetrical charge distribuation. In an electric field gradient this gives rise to another kind of splitting:

$$\Delta E_Q(m) = \frac{eQ}{4} \frac{\partial^2 V}{\partial z^2} \frac{3m^2 - I(I-1)}{3I^2 - I(I-1)}$$

Since the splitting is dependant on the square of the magnetic quantum number m, it only occurs for I>1 and for $\mathrm{Fe^{57}}$ there are only two lines instead of four. With a given quadrupole moment, the difference between the two lines can be used to determine the field gradient.

1.2 Assignments

1.2.1 Mößbauer Spectra of different Iron Compounds

The emitter is accelerated to produce a periodic pattern of velocities relative to the different absorbers and the transmission spectra are measured.

1.2.2 Lifetime of 14.4 keV State in Vacromium

The lines of γ -decays like the 14.4 keV line used in this experiment are orders of magnitude too narrow, to be resolved with interferometers. Instead the Mößbauer effect can be used to measure the linewidth and with this the life time of the excited state.

The ${\rm Fe^{57}}$ source is moved relative to a Vacromium absorber and the transmission is measured. At a certain relative velocity resonance absorption occurs.

Deviating from this velocity will quench the resonance, since the Doppler effect for small velocities gives:

$$\Delta E_{\gamma} = \frac{v}{c}E$$

Making it possible to sample the resonance spectrum.

The half width of the absorption peak is double the natural line width.

The 14.4 keV line of Fe⁵⁷ has a width of $5 \cdot 10^{-9}$ eV

1.2.3 Inner Magnetic Field and Magnetic Moment of the Excited State in Iron

Next the magnetic field and moment of iron will be measured. From the six absorption peaks one of the four cases for the relations outlined in V9 2.4.1 has to be selected.

1.2.4 Electric Field Gradient in Iron Compounds

The gradient can be determined from the distance in velocity of the two peaks:

$$\Delta v = \frac{eQ}{2} \frac{\partial^2 V}{\partial z^2} \frac{c}{E_0}$$

2 Analysis

2.1 Velocity Calibration

First the MCS has to be calibrated, in order to transfer channel numbers to velocities.

With the given formula:

$$v[i] = \frac{\lambda}{2} \frac{N[i]}{N[2] \cdot 10} \text{mm/s}$$

and the additional information, that the range is traversed from maximum positive to maximum negative velocity and back, the relation illustrated in fig(2.1) is found.

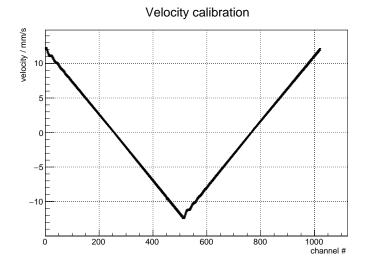


Abbildung 1: Velocities

2.2 Spectra

The following figures display the recorded spectra with gaussians fitted to the absorption peaks.

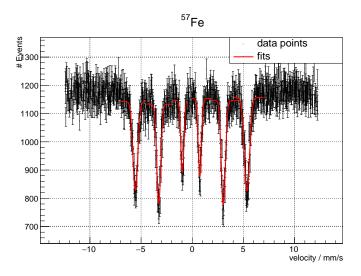


Abbildung 2: Velocities

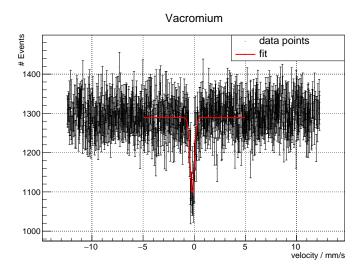


Abbildung 3: Velocities

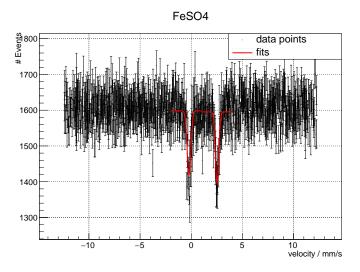


Abbildung 4: Velocities

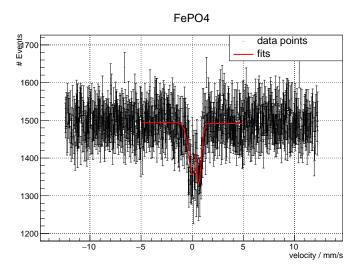


Abbildung 5: Velocities

2.3 Vacromium

The chemical shift for the vacromium peak computes to:

$$\nu = -0.1907 \pm 0.0129~\rm{mm/s}$$

The half width is:

$$v_{1/2} = 0.2529 \text{ mm/s}$$

which corresponds to an energy of $\Delta E = 1.215 \cdot 10^{-8}$ eV for 14.4 keV. The lifetime of the excited state follows:

$$\tau = \hbar/\Gamma_0 = 2 \cdot \hbar/\Delta E = 108.4 \text{ ns}$$

2.4 Iron

The fits give the following resonance velocities:

$$v_6 - 5.56 \pm 0.00789$$

 $v_5 - 3.27 \pm 0.00620$
 $v_4 - 0.97 \pm 0.00800$
 $v_3 0.74 \pm 0.00746$
 $v_2 3.04 \pm 0.00607$
 $v_1 5.34 \pm 0.00793$

and thus

$$I = -0.11125 \text{ mm/s}$$

for the chemical shift.

Since $v_1 = 2 \cdot v_2 + v_3 = 0.006305$ mm/s fulfills the criterium for the four possible parameter relations the best, it follows:

$$G = -v_2 - v_3 = -4.0066 \text{ mm/s}$$

and

$$A = v_1 - v_2 = 2.3044 \text{ mm/s}$$

from these the internal magnetic field and moment follow:

$$B = \frac{GI_g E_0}{\mu_g c} = -33.8 \text{ T}$$

$$\mu_e = \frac{AI_e E_0}{Bc} = -4.912 \cdot 10^{-9} \text{eV/T}$$

There were no values provided in the literature for this lab, but this source¹ gives -33 T and $-4.8768 \cdot 10^{-9} \text{eV/T}$, which is reasonably close to the values found here.

 $^{^{1}} https://www.chemie.uni-hamburg.de/bibliothek/2003/DissertationPickardt.pdf$

2.5 Iron Compounds

The field gradient of two different iron compounds is to be measured here.

2.5.1 FeSO4

The distance between the two peaks in the FeSO4 absorption spectrum is $\Delta v = 2.734 \text{ mm/s}$ and thus the gradient is:

$$\frac{\partial^2 V}{\partial z^2} = 1.251 \cdot 10^{22} \text{ V/m/m}$$

The chemical shift is $\nu = 1.954$ mm/s.

2.5.2 FePO4

For FePO4 the peaks are very close to each other. To resolve them, a single gaussian was fitted to both and the peak velocity of this fit used as the border between the fitting ranges of the single peaks.

With this the distance between peaks is $\Delta v = 0.5735$ mm/s and the gradient:

$$\frac{\partial^2 V}{\partial z^2} = 1.251 \cdot 10^{22} \text{ V/m/m}$$

The chemical shift is $\nu = 0.329$ mm/s.