

EXERCISES WITH THEIR CORRECTIONS

EXERCICE 1 ADJOINT OPERATOR OF A SYMMETRICAL SECOND-ORDER OPERATOR

Let V be a subspace of $H^1(\Omega)$, let A be the linear operator defined from V into V' as follows:

$$A(y) = -\operatorname{div}(\lambda \nabla y) + c y$$

with λ and c given in $L^\infty(\Omega)$.

Question I) Write the expression of A^* in weak form, in the following cases.

Case 1) With Dirichlet B.C.. Homogeneous Dirichlet conditions are imposed on the whole boundary $\partial\Omega$. In this case, we set $V = H_0^1(\Omega)$.

Show that in this case the operator is self-adjoint, i.e. $A^* = A$.

Case 2) With mixed B.C..

Considering mixed B.C., we set $\partial\Omega = \Gamma_0 \cup \Gamma_1$ and:

$$V = H_{\Gamma_0}^1(\Omega) = \{z \in H^1(\Omega), z = 0 \text{ on } \Gamma_0\}$$

Question II)

Write the expression of A^* in weak form, next in classical form, in the following cases.

We consider the mixed BC as previously, with:

Case 1) $(-\lambda \nabla y \cdot n) = \varphi$ given on Γ_1 ,

Case 2) $(-\lambda \nabla y \cdot n) = \varphi(y)$ given on Γ_1 .

***.

Correction. Recall that by definition, given A a linear operator, its adjoint operator A^* satisfies:

$$\langle Ay, z \rangle_{V' \times V} = \langle y, A^*z \rangle_{V \times V'} \quad \forall (y, z) \in V \times V$$

Question I)

Case 1)

Let $(y, z) \in V \times V$; $V = H_0^1(\Omega)$.

The operator $A(y) = -\operatorname{div}(\lambda \nabla y) + c y$ is linear.

By definition, we have: $\langle A(y), z \rangle_{V' \times V} = \int_{\Omega} [-\operatorname{div}(\lambda \nabla y) + c y] z \, dx$

By applying Green's formulae, we obtain:

$$\langle A(y), z \rangle_{V' \times V} = \int_{\Omega} (\lambda \nabla y \nabla z + c y z) \, dx \quad \text{since } z \in H_0^1(\Omega)$$

By applying Green's formulae again, we obtain:

$$\begin{aligned} \langle A(y), z \rangle_{V' \times V} &= \int_{\Omega} [-\operatorname{div}(\lambda \nabla z) + c z] y \, dx + \int_{\partial\Omega} \lambda \nabla z \cdot n y \, ds \\ &= \int_{\Omega} [-\operatorname{div}(\lambda \nabla z) + c z] y \, dx \quad \text{since } y \in H_0^1(\Omega) \\ &= \langle y, A^*z \rangle_{V \times V'} \end{aligned}$$

Therefore: $A(z) \equiv Az = A^*z$ for all $z \in V$. In other words: $A = A^*$.

Case 2)

Case of mixed BCs with $V = H_{\Gamma_0}^1(\Omega)$.

Let $(y, z) \in V \times V$. By applying Green's formulae, we obtain:

$$(0.1) \quad \langle A(y), z \rangle_{V' \times V} = \int_{\Omega} (\lambda \nabla y \nabla z + c y z) \, dx - \int_{\Gamma_1} \lambda \nabla y \cdot n z \, ds$$

By applying Green's formulae again, it follows:

$$\begin{aligned}
\langle A(y), z \rangle_{V' \times V} &= \int_{\Omega} [-\operatorname{div}(\lambda \nabla z) + c z] y \, dx + \int_{\Gamma_1} \lambda \nabla z \cdot n y \, ds - \int_{\Gamma_1} \lambda \nabla y \cdot n z \, ds \\
&= \langle y, A^* z \rangle_{V \times V'} \text{ by definition.}
\end{aligned}$$

We obtain the weak form of the adjoint operator: for $(p, z) \in V \times V$,

$$(0.2) \quad \langle A^* p, z \rangle_{V' \times V} = \int_{\Omega} (\lambda \nabla p \nabla z + c p z) \, dx - \int_{\Gamma_1} \lambda \nabla z \cdot n p \, ds$$

By combining (0.1) and (0.2), we get:

$$\langle (Ay - A^* y), z \rangle_{V' \times V} = - \int_{\Gamma_1} \lambda \nabla y \cdot n z \, ds + \int_{\Gamma_1} \lambda \nabla z \cdot n y \, ds$$

As a consequence, because of the mixed BCs, the operator A is *a-priori* not self-adjoint in $V = H_{\Gamma_0}^1(\Omega)$. However... see next question!...

Note that we have:

$$\langle A^* p, z \rangle_{V' \times V} = \int_{\Omega} [-\operatorname{div}(\lambda \nabla p) + c p] z \, dx + \int_{\Gamma_1} \lambda \nabla p \cdot n z \, ds - \int_{\Gamma_1} \lambda \nabla z \cdot n p \, ds$$

Question II)

Case 1) We have: $-(\lambda \nabla y \cdot n) = \varphi$ given on Γ_1 . Therefore, for all $(y, z) \in V \times V$,

$$(0.3) \quad \langle Ay, z \rangle_{V' \times V} = \int_{\Omega} (\lambda \nabla y \nabla z + c y z) \, dx + \int_{\Gamma_1} \varphi z \, ds$$

In the case $\varphi \neq 0$, the linear part of the operator is the map $y \mapsto \int_{\Omega} (\lambda \nabla y \nabla z + c y z) \, dx$ only.

Indeed the term $\int_{\Gamma_1} \varphi z \, ds$ does not depend on y (it is constant wrt y).

As a consequence, the linear part of the operator is self-adjoint again.

In the case $\varphi = 0$, the operator A is self-adjoint.

Case 2) We have: $-(\lambda \nabla y \cdot n) = \varphi(u)$.

In the case $\varphi(\cdot)$ not linear, the question does not apply. Indeed, the operator must be linear to define its adjoint.

In the case $\varphi(\cdot)$ linear, we have: $-(\lambda \nabla y \cdot n) = \alpha y$ on Γ_1 , α given.

In this case, we have for all $(y, z) \in V \times V$,

$$(0.4) \quad \langle Ay, z \rangle_{V' \times V} = \int_{\Omega} (\lambda \nabla y \nabla z + c y z) \, dx + \int_{\Gamma_1} \alpha y z \, ds$$

$$(0.5) \quad = \langle y, A^* z \rangle_{V \times V'} \text{ by definition.}$$

The (linear) operator A is self-adjoint.

For $(p, z) \in V \times V$,

$$\langle A^* p, z \rangle_{V' \times V} = \int_{\Omega} (\lambda \nabla p \nabla z + c p z) \, dx + \int_{\Gamma_1} \alpha p z \, ds$$

In the classical form, the operators read in Ω : $Ay = A^* y = -\operatorname{div}(\lambda \nabla y) + c y$.

Recall the mixed BC on y : $y = 0$ on Γ_0 , $-(\lambda \nabla y \cdot n) = \alpha y$ on Γ_1 .

These BCs are the same for the adjoint operator: $p = 0$ on Γ_0 , $-(\lambda \nabla p \cdot n) = \alpha p$ on Γ_1 .

EXERCISE 2: ADJOINT OPERATOR OF THE ADVECTION-DIFFUSION EQUATION

Let V be a subspace of $H^1(\Omega)$. Let $y(x)$ be a scalar function defined in V .

We consider the PDE operator A defined from V into V' as:

$$A(y) = -\operatorname{div}(\lambda \nabla y) + \mathbf{w} \cdot \nabla y$$

with λ given in $L^\infty(\Omega)$, $\lambda > 0$, the vector field \mathbf{w} given in $(L^\infty(\Omega))^d$.

Assumption. The vector field \mathbf{w} satisfies the following properties: $\operatorname{div}(\mathbf{w}) = 0$ in Ω and $\mathbf{w} \cdot \mathbf{n} \geq 0$ on Γ_1 .

In the case \mathbf{w} represents a velocity field, the velocity field is incompressible and incoming where the solution value is given (i.e. on Γ_0).

Case 1) Homogeneous Dirichlet conditions. Homogeneous Dirichlet conditions are imposed everywhere on $\partial\Omega$. In this case, we have $V = H_0^1(\Omega)$.

a) *Existence and uniqueness.*

Let $f \in L^2(\Omega)$ be given. Prove that the BVP,

$A(y) = f$ accompanied with homogeneous Dirichlet B.C. on $\partial\Omega$, admits a unique weak solution u in V .

b) *Adjoint operator.*

Write an expression of A^* in weak form, next in its classical form.

Correction.

a) *Existence and uniqueness.*

The weak form of the model reads as follows.

Find $y \in V = H_0^1(\Omega)$ such that $a(y, z) = l(z)$ with

$$a(y, z) = \int_{\Omega} \lambda \nabla y \nabla z \, dx + \int_{\Omega} \mathbf{w} \cdot \nabla y \, z \, dx \text{ and } l(z) = \int_{\Omega} f \, z \, dx$$

When applying the Lax-Milgram theorem, the most critical property to verify is the coercivity in $V = H_0^1(\Omega)$. Here, for all $y \in V = H_0^1(\Omega)$,

$$\int_{\Omega} \mathbf{w} \cdot \nabla y \, y \, dx = \frac{1}{2} \sum_i \int_{\Omega} \mathbf{w}_i \partial_i (y^2) \, dx = -\frac{1}{2} \sum_i \int_{\Omega} (y^2) \partial_i \mathbf{w}_i \, dx + \int_{\partial\Omega} \mathbf{w}_i n_i \, y^2 \, dx$$

Therefore:

$$\int_{\Omega} \mathbf{w} \cdot \nabla y \, y \, dx = 0$$

Hence the bilinear form $a(y, z)$ is coercive in $V = H_0^1(\Omega)$.

The existence and uniqueness of the weak solution follows from the Lax-Milgram theorem.

b) *Adjoint operator.*

By integrating by part, we get:

$$\int_{\Omega} \mathbf{w} \cdot \nabla y \, z \, dx = - \int_{\Omega} \operatorname{div}(z \mathbf{w}) \, y \, dx + \int_{\partial\Omega} \mathbf{w} \cdot \mathbf{n} \, z \, y \, dx = - \int_{\Omega} \operatorname{div}(z \mathbf{w}) \, y \, dx$$

For A linear, we have by definition:

$$a(y, z) = \langle Ay, z \rangle_{V' \times V} \stackrel{\text{def}}{=} \langle A^*z, y \rangle_{V' \times V}$$

with here $V = H_0^1(\Omega)$.

Using $\operatorname{div}(\mathbf{w})=0$ and the Dirichlet B.C., we get:

$$\langle A^*z, y \rangle_{V' \times V} = \int_{\Omega} \lambda \nabla z \nabla y \, dx - \int_{\Omega} \mathbf{w} \cdot \nabla z \, y \, dx$$

Therefore the adjoint of the advection term $(\mathbf{w} \cdot \nabla z)$ is $-(\mathbf{w} \cdot \nabla z)$.

PDE model operators containing the 1st order non symmetrical advective term $(\mathbf{w} \cdot \nabla y)$ are not self-adjoint.

However, the advection term $(\mathbf{w} \cdot \nabla y)$ simply transforms as $-(\mathbf{w} \cdot \nabla p)$.

Case 2) Mixed boundary conditions.

Let us set $\partial\Omega = \Gamma_0 \cup \Gamma_1$ and $V = H_{\Gamma_0}^1(\Omega) = \{z \in H^1(\Omega), z = 0 \text{ on } \Gamma_0\}$.

a) *Existence and uniqueness.*

Let $f \in L^2(\Omega)$ be given. Prove that the BVP

$A(y) = f$ with mixed homogeneous BCs (Dirichlet / Neumann) on Γ_0, Γ_1 respectively, admits an unique weak solution in V .

b) *Adjoint operator.*

Write an expression of A^* in weak form.

Correction.

a) *Existence and uniqueness using the Lax-Milgram theorem.*

The analysis is similar to the one in the previous case. We write:

$$\int_{\Omega} \mathbf{w} \cdot \nabla y \, y \, dx = \frac{1}{2} \sum_i \int_{\Omega} \mathbf{w}_i \partial_i (y^2) \, dx = -\frac{1}{2} \sum_i \int_{\Omega} (y^2) \partial_i \mathbf{w}_i \, dx + \int_{\partial\Omega} \mathbf{w}_i n_i \, y^2 \, dx$$

Therefore:

$$\int_{\Omega} \mathbf{w} \cdot \nabla y \, y \, dx = \int_{\Gamma_1} \mathbf{w}_i n_i \, y^2 \, dx \geq 0 \quad \text{by assumption.}$$

Therefore the bilinear form $a(y, z)$ is coercitive in V .

The existence and uniqueness of the weak solution follows from the Lax-Milgram theorem.

In other words, in the advection-diffusion equation with an incompressible velocity field \mathbf{w} , the flow must be outgoing on the boundary part where the solution is not imposed.

b) *Adjoint operator.*

We have for all $(y, z) \in V \times V$,

$$a(y, z) = \langle Ay, z \rangle_{V' \times V} = \int_{\Omega} \lambda \nabla y \nabla z \, dx - \int_{\Omega} \text{div}(z\mathbf{w}) \, y \, dx + \int_{\Gamma_1} \mathbf{w} \cdot \mathbf{n} \, z \, y \, ds - \int_{\Gamma_1} \lambda \nabla y \cdot \mathbf{n} \, z \, ds$$

Using the incompressibility assumption $\text{div}(\mathbf{w}) = 0$, we get: $\int_{\Omega} \text{div}(z\mathbf{w}) \, y \, dx = \int_{\Omega} \mathbf{w} \cdot \nabla z \, y \, dx$.

Recall that by definition: $\langle Ay, z \rangle_{V' \times V} = \langle A^*z, y \rangle_{V' \times V}$.

Let us write the adjoint equation LHS term $\langle A^*p, z \rangle_{V' \times V}$.

For all $(p, z) \in V \times V$,

$$\langle A^*p, z \rangle_{V' \times V} = \int_{\Omega} \lambda \nabla p \nabla z \, dx - \int_{\Omega} \mathbf{w} \cdot \nabla p \, z \, dx + \int_{\Gamma_1} \mathbf{w} \cdot \mathbf{n} \, p \, z \, ds - \int_{\Gamma_1} \lambda \nabla z \cdot \mathbf{n} \, p \, ds$$

Because of the non-symmetrical 1st order term, A is not self-adjoint.

The boundary terms on Γ_1 have to be clarified from Neumann type BCs.