

Title

Max Varverakis

February 12, 2024

Contents

| | | |
|----------|--|-----------|
| 1 | Representation Theory Background | 2 |
| 2 | Topological Definitions | 7 |
| 3 | The Braid Group | 10 |
| 3.1 | Visualization of pure braids | 10 |
| 3.2 | General braids | 12 |
| 3.3 | Standard generators of the braid group | 12 |
| 3.4 | Automorphisms of the free group | 13 |
| 3.5 | The Burau Representation | 18 |
| 4 | tikz test | 23 |
| | References | 24 |

Chapter 3

The Braid Group

Definition 3.1. The *configuration space* of n ordered distinct points in the complex plane \mathbb{C} is defined as $M_n = \{(z_1, \dots, z_n) \in \mathbb{C}; z_i \neq z_j, \forall i \neq j\}$. Alternatively, consider \mathcal{D} to be the collection of all hyperplanes $H_{i,j} = \{z_i = z_j\} \in \mathbb{C}^n$ for $1 \leq i < j \leq n$. Then we can define $M_n = \mathbb{C}^n \setminus \mathcal{D}$.

Note that $(z_1, z_2, z_3, \dots, z_n)$ and $(z_2, z_1, z_3, \dots, z_n)$ are different points in the configuration space M_n . Before studying the various interpretations of the braid group, we first define the braid group itself.

Definition 3.2. The *pure braid group* on n strands, denoted PB_n , is the fundamental group of M_n . One can write $PB_n = \pi_1(M_n)$.

3.1 Visualization of pure braids

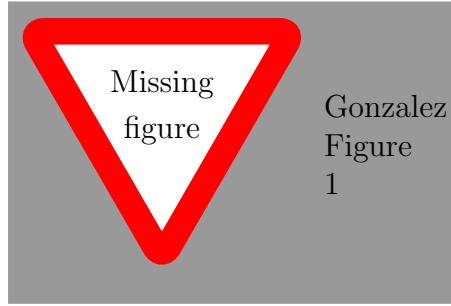
We can think of a pure braid as a loop in M_n :

$$\begin{aligned} \beta : [0, 1] &\rightarrow M_n \\ t &\mapsto \beta(t) = (\beta_1(t), \beta_2(t), \dots, \beta_n(t)), \end{aligned}$$

with some base point. Conventionally, we define the base point as the n -tuple of integers $(1, 2, 3, \dots, n) \in \mathbb{C}^n$. Then a pure braid can be thought of the motion of these points in the complex plane as t ranges from 0 to 1 in which $\beta_i(t)$ is defined and $\beta_i(t) \neq \beta_j(t)$ for every $t \in [0, 1]$ and $i \neq j \in \{1, 2, \dots, n\}$. Because each β_i is a loop, it must start and end at the point

i (e.g., $\beta_i(0) = \beta_i(1) = i$). Recall that the loops are actually equivalence classes of loops under homotopy. As a result, we can continuously deform the motion of the n points while maintaining the same pure braid (up to equivalence) so long as we preserve the pairwise distinction of the points for all time $t \in [0, 1]$.

A common visualization of pure braids is to plot the motion of the points in 3-dimensional space. For each $t \in [0, 1]$, we draw the points $(\beta_i(t), t)$ in $\mathbb{C} \times [0, 1]$ for every $i \in \{1, \dots, n\}$. The space $\mathbb{C} \times [0, 1]$ can be thought of as a spacetime diagram, where the motion of the points is plotted in the complex plane at each time t , with the time being the vertical axis. The convention is to have $\mathbb{C} \times \{0\}$ placed above $\mathbb{C} \times \{1\}$, so that the motion of the points is plotted from top to bottom.



For every $i \in \{1, \dots, n\}$, the motion of a single point starting at $(i, 0)$ and ending at $(i, 1)$ is known as the i -th *strand* of the pure braid. This can also be described by the i -th projection of the n -tuple $\beta(t)$. Thus, two braids are equivalent under homotopy if, for every moment of a continuous deformation of the n strands in $\mathbb{C} \times [0, 1]$, the (fixed) endpoints $((1, 0), (2, 0), \dots, (n, 0))$ and $((1, 1), (2, 1), \dots, (n, 1))$ are connected by strands that are pairwise disjoint where each strand intersects the plane $\mathbb{C} \times \{t\}$ exactly once for every $t \in [0, 1]$.

As pure braids are members of the pure braid group, multiplication is a well-defined operation. In the context of M_n , multiplication of pure braids involves the concatenation of loops. Visually, this is the process of stacking braids on top of each other, and then rescaling the time dimension so that t ranges from 0 to 1.

3.2 General braids

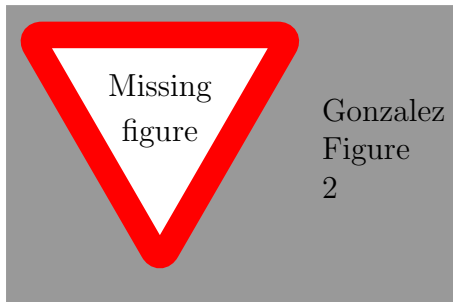
In the previous section, we defined pure braids in which the endpoints of each strand are identical at the beginning and end of the motion. This notion generalizes to define (non-pure) braids. First, we define a more general configuration space than M_n . The symmetric group S_n permutes the n distinct points in \mathbb{C} . Then the *configuration space of n unordered points in \mathbb{C}* is the quotient space $N_n = M_n/S_n$.

Definition 3.3. The braid group on n strands is the fundamental group of N_n , denoted $B_n = \pi_1(N_n)$.

The visualization of a braid is the same as in the case of pure braids, only now the endpoints of each strand do not necessarily match the starting points. For example, the i -th strand may start at the point $(i, 0)$ but end at the point $(j, 1)$ for $i, j \in \{1, \dots, n\}$. The equivalence of strands is still defined as before under the homotopy of loops. Loop concatenation defines the multiplication of braids, as before.

3.3 Standard generators of the braid group

Originally proposed by Artin [1], each braid can be decomposed into a product of *standard generators* of the braid group. When visualizing braids in $\mathbb{R} \times [0, 1]$, a crossing of two strands is clearly indicated by one going over the other. Suppose each crossing occurs at a different time $t \in [0, 1]$. Then by rescaling the time component of an arbitrary braid, we can deconstruct it into a stack of simple braids with only one crossing between neighboring strands per braid. Each single crossing of strands can be obtained by performing a transposition between neighboring endpoints of the strands.



For instance, swapping the endpoints of the i -th and $(i+1)$ -th strands can be written as applying σ_i to the identity braid (i.e., the braid that starts without any crossings of strands). It must be noted that there are two distinct ways to swap the endpoints of two strands. From a top-down perspective looking at the plane $\mathbb{C} \times \{t\}$ for some time t , σ_i swaps (i, t) and $(i+1, t)$ in a clockwise rotation. The reverse of this operation (i.e., twisting the endpoints around in the counterclockwise direction) is denoted σ_i^{-1} . The standard generators of the braid group B_n are defined as the set $\{\sigma_1, \sigma_2, \dots, \sigma_{n-1}\}$. An arbitrary braid can be constructed by concatenating (or stacking) the simple braids made from the standard generators before rescaling the time coordinate to $[0, 1]$.

3.4 Automorphisms of the free group

Consider the n -times punctured disk \mathbb{D}_n . The fundamental group of \mathbb{D}_n involves loops that start and end at the same (fixed) base point in $\partial\mathbb{D}_n$. Up to homotopy, a clockwise-directional loop that encompasses the i -th hole in \mathbb{D}_n corresponds to the i -th generator of the free group F_n of rank n , which is illustrated in Figure 3.1. In fact, $\pi_1(\mathbb{D}_n) = F_n$. This equality allows us to define a representation of the braid group on n strands as automorphisms of F_n .

Each braid $\beta \in B_n$ is realized as an automorphism of $\pi_1(\mathbb{D}_n) = F_n$ (up to isotopy) where each loop $\gamma \in \pi_1(\mathbb{D}_n)$ is sent to another loop $\beta(\gamma)$. In other words, we have a representation of the braid group defined by

$$\rho : B_n \rightarrow \text{Aut}(F_n) \tag{3.1}$$

$$\beta \mapsto \rho_\beta. \tag{3.2}$$

The action of β on a loop γ is defined by the rearrangements of the n holes in \mathbb{D}_n , similar to the action of the standard generators of B_n on the base points in $\mathbb{C} \times \{1\}$. In terms of the standard generators of B_n , each σ_i corresponds to switching the places of hole i and hole $i+1$ by means of a clockwise rotation, as seen in Figure 3.2. This is identical to viewing the action of σ_i on the base points (Section 3.3) from above, looking down on the $\mathbb{C} \times \{t\}$ -plane. As before, the inverse action σ_i^{-1} is a counterclockwise rotation of the two adjacent holes i and $i+1$ in \mathbb{D}_n . These actions respect the group operation of loop concatenation.

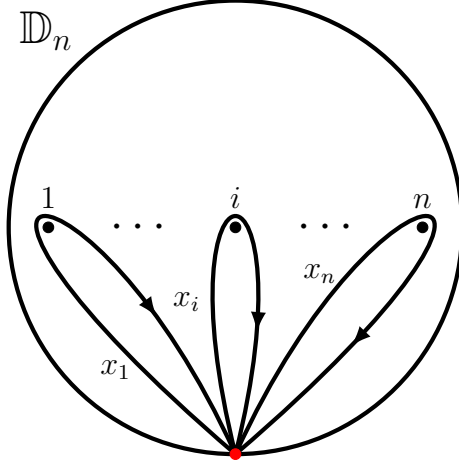


Figure 3.1: For each $i \in \{1, \dots, n\}$, the clockwise-directional loop encircling the i -th hole in \mathbb{D}_n corresponds to the i -th free generator of F_n (i.e., x_i). The red dot indicates the (arbitrary) base point for the loops in $\pi_1(\mathbb{D}_n)$.

A note on convention of directions/consistency up to inverses.

The automorphism ρ_β is most simply defined in terms of the action of the standard generators of B_n on the generators x_1, \dots, x_n of F_n (visualized as loops in \mathbb{D}_n). For each i , it follows that

$$\rho_{\sigma_i}(x_i) = x_i x_{i+1} x_i^{-1}, \quad (3.3)$$

$$\rho_{\sigma_i}(x_{i+1}) = x_i, \quad (3.4)$$

$$\rho_{\sigma_i}(x_j) = x_j, \text{ for } j \neq i, i-1. \quad (3.5)$$

Clearly, any two loops that are separated by at least one puncture will not interact while performing σ_i . The relations for adjacent loops can be verified graphically as illustrated in Figures 3.2 and 3.3.

For any σ_i , $\rho_{\sigma_i^{-1}}$ is well-defined. It follows that for any braid $\beta \in B_n$, we can decompose ρ_β the composition of the automorphisms of the standard generators $\sigma_1, \dots, \sigma_{n-1}$ and their inverses that make up β .

Notice that for any σ_i , $\rho_{\sigma_i}(x_1 \cdots x_n) = x_1 \cdots x_n$. This is because the loop $x_1 \cdots x_n$ in \mathbb{D}_n , encompassing all holes, is parallel to the boundary $\partial\mathbb{D}_n$. Thus, the action of σ_i on $x_1 \cdots x_n$ is trivial does not affect the structure of the loop up to isotopy. More generally, this implies that $\rho_\beta(x_1 \cdots x_n) = x_1 \cdots x_n$ for

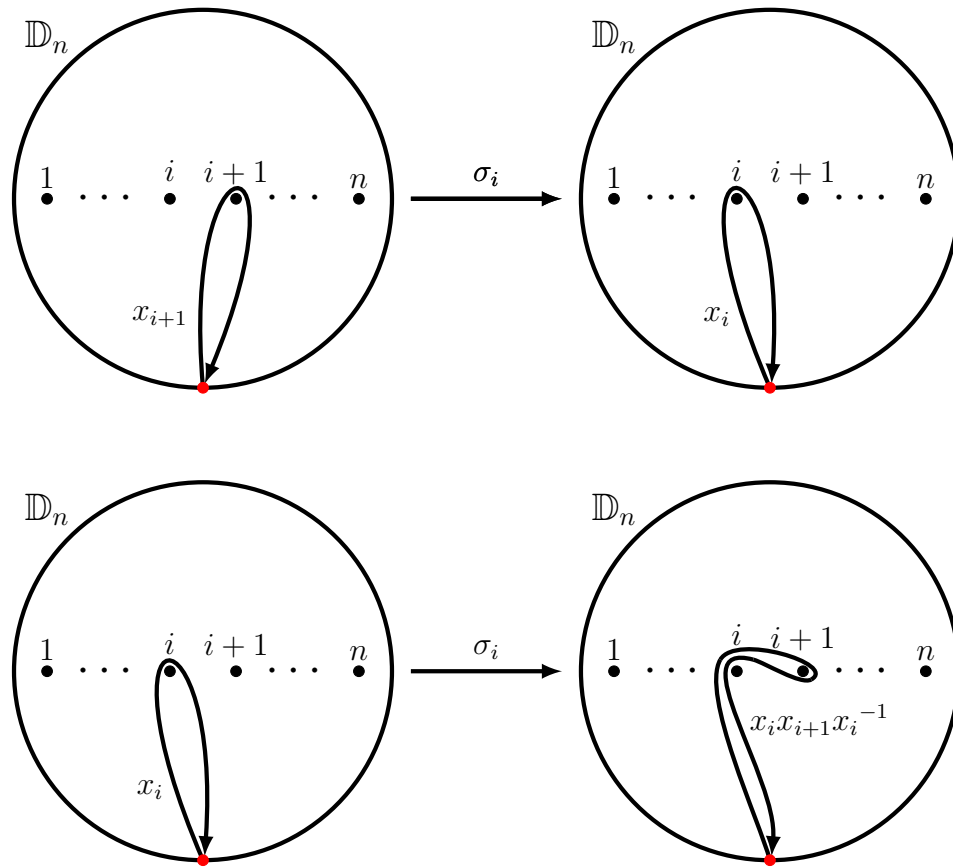


Figure 3.2: The action of σ_i on the generators x_i and x_{i+1} as described by Eqns. 3.3–3.5. The image of x_i under σ_i is verified visually in Figure 3.3

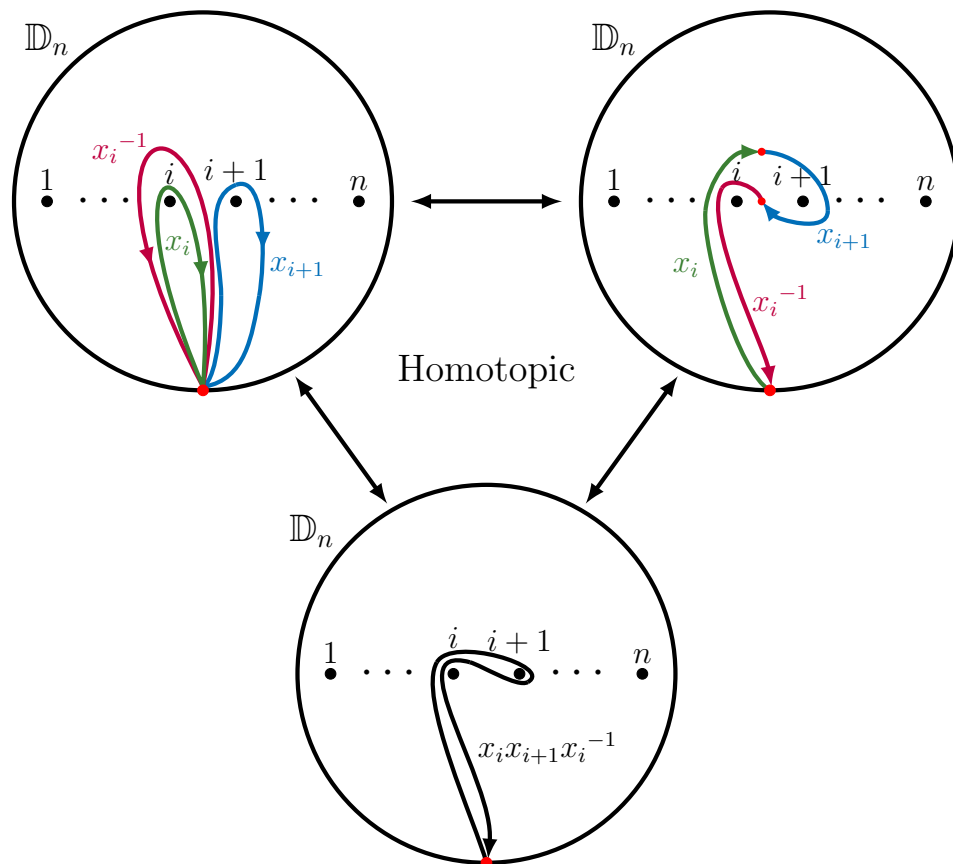


Figure 3.3: Up to homotopy, the product $x_i x_{i+1} x_i^{-1}$ is visualized as the concatenation of the loops $x_i, x_{i+1}, x_i^{-1} \in \pi_1(\mathbb{D}_n)$. In the top right diagram, small red dots are used indicate the (homotopically deformed) points of concatenation.

any $\beta \in B_n$. Paired with the observation that every generator is conjugate to another, Artin [1] showed that this is a necessary and sufficient condition for ρ_β to be an automorphism of F_n .

Theorem 3.1. *An automorphism $f \in \text{Aut}(F_n)$ is equal to ρ_β for some $\beta \in B_n$ if and only if*

1. $f(x_i)$ is a conjugate of some x_j for every $i \in \{1, \dots, n\}$, and
2. $f(x_1 \cdots x_n) = x_1 \cdots x_n$.

In this interpretation of the braid group, we can express B_n in terms of the standard generators:

$$B_n = \left\langle \sigma_1, \dots, \sigma_{n-1} \left| \begin{array}{ll} \sigma_i \sigma_j = \sigma_j \sigma_i, & |i - j| > 1 \\ \sigma_i \sigma_{i+1} \sigma_i = \sigma_{i+1} \sigma_i \sigma_{i+1}, & |i - j| = 1 \end{array} \right. \right\rangle. \quad (3.6)$$

The first condition that the standard generators commute if $|i - j| > 1$ is easily verified by looking at the graphic of Figure 3.3 in describing the property of the automorphism $\rho_{\sigma_i}(\sigma_j) = \sigma_j$. This follows by the fact that if two holes are non-adjacent, then the actions of σ_i and σ_j are mutually exclusive and thus commutative. The second condition on the standard generators is most easily verified in Figure 3.4 by looking at the corresponding braids in $\mathbb{R} \times [0, 1]$.

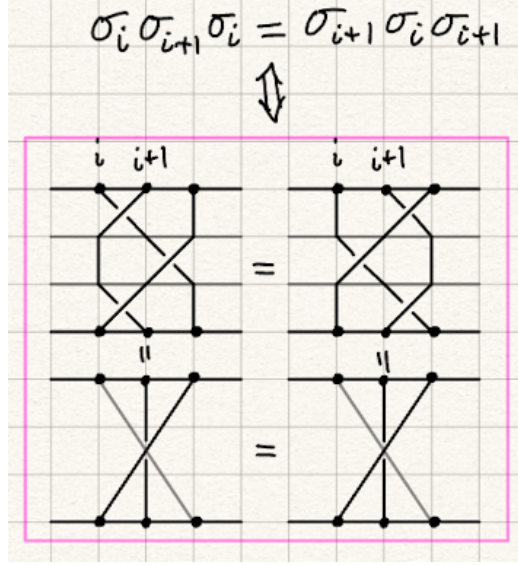


Figure 3.4: **Grayed-out strand indicates that it is behind all other strands.**

3.5 The Burau Representation

In the previous section, we defined a representation of the braid group B_n as automorphisms of the free group F_n . This representation is clearly non-abelian. Likewise, the Artin generators of B_n are nonabelian. Suppose we wish to abelianize the braid group. The details of the abelianization of B_n would require a quotient by the commutator $[a, b] = aba^{-1}b^{-1}$.

Sparing the details, let $B_{n,ab} = B_n / [B_n, B_n]$ be the abelianization of B_n , where $[B_n, B_n] = \{[\beta_1, \beta_2] \mid \beta_1, \beta_2 \in B_n\}$ is the commutator subgroup of B_n . Then, under the representation ρ from Section 3.4, the abelianization of Eqns. 3.3–3.5 become

$$x_i \xrightarrow{\sigma_i} x_i + x_{i+1} - x_i = x_{i+1} = \rho_{\sigma_i^{-1}}(x_i), \quad (3.7)$$

$$x_{i+1} \xrightarrow{\sigma_i} x_i, \quad (3.8)$$

$$x_j \xrightarrow{\sigma_i} x_j, \text{ for } j \neq i, i-1, \quad (3.9)$$

for each i . Thus, the generator $\sigma_i = \sigma_i^{-1}$, and corresponds to a transposition permutation in the symmetric group S_n . It follows that $B_{n,ab} \simeq S_n$. In this current construction, the abelianization of the braid group results in a loss of

complexity. This raises the question whether there exists such a reframing of the braid group that allows an abelian operation on the free generators while preserving the inequivalence of the Artin generators with their inverses.

First, we define a topological space that will aid in the desired construction.

Definition 3.4. Let X be a topological space. A *covering* of X is a space \tilde{X} together with a continuous map $p : \tilde{X} \rightarrow X$ such that, for every $x \in X$, there exists a path-connected open neighborhood U containing x such that $p^{-1}(U)$ is a disjoint union of open sets in \tilde{X} where each component of $p^{-1}(U)$ is mapped homeomorphically onto U by p . Each component of $p^{-1}(U)$ is called a *sheet* of the covering, where the i -th sheet is denoted by \tilde{X}_i , and total the number of sheets in $p^{-1}(U)$ is called the *degree* of the covering.

Example 3.1. One of the simplest examples of a covering space is the covering of the circle S^1 by the real line by the parameterization map $p : \mathbb{R} \rightarrow S^1$ defined by $p(t) = (\cos t, \sin t)$. Clearly, there are infinitely many sheets in this covering.

Example 3.2. A similar example is the covering of the circle S^1 through $p : [0, 1] \rightarrow S^1$ defined by $p(t) = e^{2\pi it}$. This defines a one-degree covering of S^1 . If we instead let our domain be $[0, 2]$, then we have a two-degree covering of S^1 .

With this topological tool, we construct a countably infinite-degree covering of the punctured disk \mathbb{D}_n , denoted $\tilde{\mathbb{D}}_n$, which can be visualized as an infinite stack of copies of \mathbb{D}_n , with a slight modification to be explained shortly. Let $\tilde{\mathbb{D}}_{n,i}$ denote the i -th sheet of $\tilde{\mathbb{D}}_n$, and consider the base sheet our covering to be $\tilde{\mathbb{D}}_{n,0}$.

We start this construction with a countably infinite stack of copies of \mathbb{D}_n . Then, for every $i \in \mathbb{Z}$, for each of the n punctures on $\tilde{\mathbb{D}}_{n,i}$, apply a cut from the hole to some point on the boundary of $\tilde{\mathbb{D}}_{n,i}$, as illustrated in Figure 3.5 for the case when $n = 3$. Each cut results in two edges, which will be referred to as the left edge and the right edge. Through a homeomorphic deformation, connect the left edge of $\tilde{\mathbb{D}}_{n,i}$ to the corresponding right edge of $\tilde{\mathbb{D}}_{n,i+1}$, and the right edge of $\tilde{\mathbb{D}}_{n,i}$ to the left edge of $\tilde{\mathbb{D}}_{n,i-1}$, for every cut on every sheet.

Now, viewing a single sheet, say $\tilde{\mathbb{D}}_{n,0}$, from above, when a loop with base

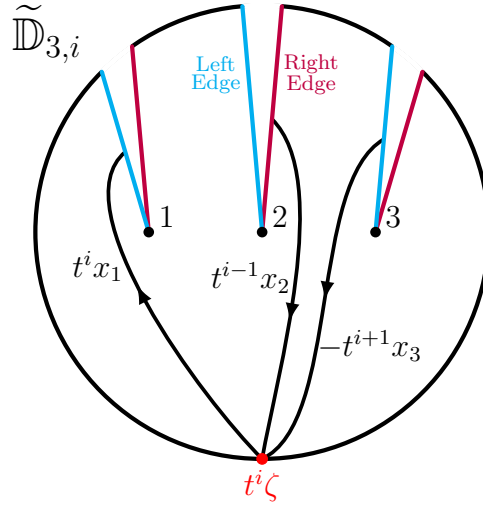


Figure 3.5: For the covering of \mathbb{D}_3 , we observe the i -th sheet of $\tilde{\mathbb{D}}_3$ with the cuts applied across each of the three punctures. The base point of the loop is indicated by the red dot, and labelled as $t^i \zeta$. The power of t indicates that we are on the i -th sheet of the covering. The portions of three different loops are drawn to illustrate the behavior of loops as they pass through various edges on $\tilde{\mathbb{D}}_{3,i}$. The loop that would traditionally be x_1 is labeled by $t^i x_1$ to indicate that it's starting on the i -th sheet. When it passes through the left edge corresponding to this puncture labeled with a 1, it traverses up to the sheet $\tilde{\mathbb{D}}_{3,i+1}$ and ends at base point $t^{i+1} \zeta$. Similarly, the loop that started on $\tilde{\mathbb{D}}_{3,i-1}$ and passed through the left edge of puncture 2 ends up coming out of the right edge of puncture 2 on $\tilde{\mathbb{D}}_{3,i}$ and ends at the base point $t^i \zeta$. This loop is labeled by the starting sheet, so it is $t^{i-1} x_2$. Finally, the loop that started on $\tilde{\mathbb{D}}_{3,i+1}$ and passed through the right edge of puncture 3 ends up coming out of the left edge of puncture 3 on $\tilde{\mathbb{D}}_{3,i}$ and is therefore labeled by $-t^{i+1} x_3$, with the negative sign indicating that the loop direction is reversed.

point $\tilde{\zeta}_0$ passes through a cut from the left, it traverses up to the next sheet, and ends at the base point $\tilde{\zeta}_1$. Similarly, a loop passing through a cut from the right ends at the base point $\tilde{\zeta}_{-1}$, on the sheet $\tilde{\mathbb{D}}_{n,-1}$ below $\tilde{\mathbb{D}}_{n,0}$. To keep track of the starting sheet of a loop, we use a free parameter t . For example, a loop γ that starts on $\tilde{\mathbb{D}}_{n,j}$ would be written $t^j\gamma$, for $j \in \mathbb{Z}$. Notice that the substitution of a complex number for the free parameter t results in a possibly finite degree covering. As an example, if we set t to an n -th root of unity, then we obtain an n -th degree covering of \mathbb{D}_n . For the purposes of this construction, we will keep t as a free parameter for now. Figure 3.5 demonstrates how loops interact with the cuts on different sheets. The following example describes the action of the standard generators of B_3 on the covering space $\tilde{\mathbb{D}}_3$.

Example 3.3. Consider the case when $n = 3$. Then we have the corresponding covering space $\tilde{\mathbb{D}}_3$ of \mathbb{D}_3 . See Figure 3.5 for the view of a single sheet with various loops interacting with the cuts on the sheet. The actions of the standard generators of B_3 in $\pi_1(\mathbb{D}_3)$ are known, and can be visually understood in Figure 3.2. In the context of the covering space $\tilde{\mathbb{D}}_3$, the action of the generators σ_1 and σ_2 is observed by reducing the visualization to only the base points on each sheet and the loops themselves. This can be seen in Figure 3.6 for the case of σ_1 .

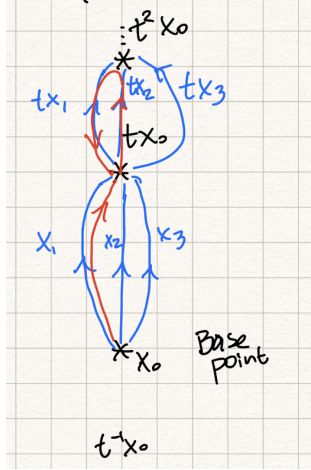


Figure 3.6: **CAPTION**

Now, we can express loop concatenation as an abelian operation, where

Eqs. 3.3–3.5 become

$$x_1 \xrightarrow{\sigma_1} x_1 + tx_2 - tx_1 = (1-t)x_1 + tx_2, \quad (3.10)$$

$$x_2 \xrightarrow{\sigma_1} x_1, \quad (3.11)$$

$$x_3 \xrightarrow{\sigma_1} x_3. \quad (3.12)$$

Consider the vector $\begin{bmatrix} x_1 \\ x_2 \\ x_3 \end{bmatrix}$. Then the action of σ_1 on the loops x_1, x_2, x_3 is realized by the matrix $\begin{bmatrix} 1-t & t & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 1 \end{bmatrix}$ since

$$\begin{bmatrix} 1-t & t & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \\ x_3 \end{bmatrix} = \begin{bmatrix} (1-t)x_1 + tx_2 \\ x_1 \\ x_3 \end{bmatrix}. \quad (3.13)$$

The action of σ_2 is obtained similarly, where

$$\sigma_2 \mapsto \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1-t & t \\ 0 & 1 & 0 \end{bmatrix}. \quad (3.14)$$

Notice that these matrices have entries in the ring of Laurent polynomials, $\Lambda = \mathbb{Z}[t, t^{-1}]$.

Clearly, the result from Example 3.3 generalizes to the case of braids on n strands. Fix $n > 1$. Let I_k denote the $k \times k$ dimensional identity matrix, and let

$$U = \begin{bmatrix} 1-t & t \\ 1 & 0 \end{bmatrix}. \quad (3.15)$$

For $i \in \{1, \dots, n-1\}$, the action of σ_i on $\pi_1(\widetilde{\mathbb{D}}_n)$ is realized as an $n \times n$ matrix with entries in $\Lambda = \mathbb{Z}[t, t^{-1}]$.

The Burau representation of B_n is then defined by:

$$\psi_n : B_n \rightarrow \text{GL}_n(\Lambda) \quad (3.16)$$

$$\sigma_i \mapsto \begin{bmatrix} I_{i-1} & 0 & 0 \\ 0 & U & 0 \\ 0 & 0 & I_{n-i-1} \end{bmatrix}. \quad (3.17)$$

Show invertible/inverse?

The Burau representation need only be defined on the standard generators, since any braid $\beta \in B_n$ decomposes into a product of $\sigma_1, \dots, \sigma_{n-1}$ and their inverses. Notice that if we set $t \rightarrow 1$, we recover the defining representation of S_n , as expected when we use a degree 1 covering space of \mathbb{D}_n and force the action of the generators to be abelian. Furthermore, by direct computation, it follows that

$$\psi_n(\sigma_i)\psi_n(\sigma_j) = \psi_n(\sigma_j)\psi_n(\sigma_i) \text{ for } |i - j| > 1, \quad (3.18)$$

$$\psi_n(\sigma_i)\psi_n(\sigma_{i+1})\psi_n(\sigma_i) = \psi_n(\sigma_{i+1})\psi_n(\sigma_i)\psi_n(\sigma_{i+1}) \text{ for } |i - j| = 1. \quad (3.19)$$

References

- [1] E. Artin. Theory of braids. *The Annals of Mathematics*, 48(1):101, January 1947.
- [2] W. Fulton. *Algebraic Topology: A First Course*. Graduate Texts in Mathematics. Springer New York, 1997.
- [3] Juan Gonzalez-Meneses. Basic results on braid groups, 2010.
- [4] Christian Kassel and Vladimir Turaev. *Homological Representations of the Braid Groups*, page 93–150. Springer New York, 2008.
- [5] Dale Rolfsen. Tutorial on the braid groups, 2010.
- [6] Jean-Luc Thiffeault. The bureau representation of the braid group and its application to dynamics. Presentation given at Topological Methods in Mathematical Physics 2022, Seminar GEOTOP-A, September 2022.
- [7] Wu-Ki Tung. *Group theory in physics: An introduction to symmetry principles, group representations, and special functions in classical and quantum physics*. World Scientific Publishing, Singapore, Singapore, January 1985.