# COSC3500 2D Orbital Simulation Report

Maxwell Bo (43926871)

August 17, 2018

#### Description

The task was to create a stock-standard, 2-dimensional gravitational n-body simulator. All bodies were to be assumed to be point masses. The simulation was to be accurate, maintaining a constant total energy, and exhibiting phenomena such as apsidal precession. The simulation was to accept arguments specifying the granularity of the simulation (number of time steps, number of instances of data export), and a file specifying masses and their initial positions and velocities. The simulation was to produce its output as fast as possible, with minimal slowdown with an increasing n number of bodies.

The simulator did not need to handle collisions between bodies. The MS1 simulator was to be free of OpenMP multiprocessing, which would be implemented in preparation for the MS2 submission.

#### Implementation

At a high-level, the initial naive simulator:

- 1. Parsed input parameters and files
- 2. Constructed Body class instances, representing each point mass
- 3. Packed the Bodys into a std::vector<Body>, to maximize cache locality
- 4. Calculated forces between all pairwise combinations of n-bodies  $(O(n^2))$
- 5. Performed Euler's method to derive new velocities and positions
- 6. Output all necessary data
- 7. GOTO 4

By using a Quadtree, ('a tree datastructure in which each internal node has exactly four children') and the Barnes-Hut algorithm[1], the total cost of force calculation could be reduced to  $O(n \log n)$ , by grouping close-together bodies and approximating forces between the singular grouped pseudo-body, and distant bodies. New Quadtrees were constructed on each separate simulation step.

Dehen and Read note that the Euler method 'performs very poorly in practice', further noting that 'errors are proportional to  $\Delta t^2$ '. They contrast it with the second-order *Leapfrog* symplectic integrator, which is 'heavily used in collisionless N-body applications'.[2]

Leapfrog can be expressed many in forms[3] including a synchronised form:

$$x_i = x_{i-1} + v_{i-1/2} \Delta t$$
 $a_i = F(x_i)$ 
 $v_{i+1/2} = v_{i-1/2} + a_i \Delta t$ 

which only requires a single acceleration calculation per every two half timesteps (the timestep  $\Delta t$  must be constant to maintain stability), and a 'kick-drift-kick' form

$$v_{i+1/2} = v_i + a_i \frac{\Delta t}{2}$$

$$x_{i+1} = x_i + v_{i+1/2} \Delta t$$

$$v_{i+1} = v_{i+1/2} + a_{i+1} \frac{\Delta t}{2}$$

that is stable with variable timstepping, but incurs an additional acceleration calculation per every two half timesteps.

The synchronised form was implemented, but attempts to implement the kick-drift-kick form, and variable timestepping, were left unfinished.

Thus, the final implementation:

- 1. Parsed input parameters and files
- 2. Constructed Body class instances, representing each point mass
- 3. Packed the Bodys into a std::vector<Body>, to maximize cache locality
- 4. Inserted all Bodys into a fresh QuadTree on full timesteps, traversing the QuadTree with every Body to calculate forces  $(O(n \log n))$
- 5. Performed the appropriate Leapfrog step to derive new velocities or positions
- 6. Output all necessary data
- 7. GOTO 4

#### Correctness

I personally believe that the simulation is relatively accurate. By visualising the results with  $\mathtt{matplotlib}$ , we see something that resembles an n-body simulator. The total energy is flat, observing coefficients of variation as low as 0.001%. Euler's method consistently demonstrated coefficients of variation three times higher than that of Leapfrog. Due to recommendations by literature, and observed data, the use of Euler's method was gradually phased out during my testing to speed up the process.

Barnes-Hut caused a significant increase in observed coefficient of variation, averaging 0.08% across multiple runs. Furthermore, total energy was observed to step up and down at varying intervals1. I suspect that this was because certain force calculations were causing groups of bodies to be approximated as a single pseudobody, after other had strayed too far from the pseudobody's quadtree's node's centre of mass.

When bodies are in close proxmity, anomalous total energies are observed1. Furthermore, simulations with higher numbers of bodies produce more random low energy outliers (likely due to a greater number of close encounters). There seems to be no significant difference between Euler method and *Leapfrog*, with respect to observation of anomalies.

#### Performance & Scaling

Henceforth, the use of the 'recognizable' refers to eyeballing the output data, and making no significant effort to investigate the data more rigorously.

The Barnes-Hut algorithm appeared to be dominated by its constant factor.3 Generating input files with a high number of bodies that did not cause 'spontaneous combustion' (where bodies would fly away from each other in every direction), and unconstrained Quadtree growth (where bodies achieve escape velocity, rapidly expanding the size of the tree and causing a collapse in simulation accuracy) proved difficult. Further testing on goliath is needed to discover the n that causes Barnes-Hut to be more performant than the brute-force approach.

The addition of the <code>-march=native</code> compiler flag, which enables the use of all CPU specific instructions, provided no recognizable improvement in running time, but was left enabled in the instance that it improved performance on <code>goliath</code>.

The use of both th GCC and Clang Profile-Guided Optimisation features provided no recognizable improvement in running time.

Distressingly, -00, -01, -02, -03 showed no recognizable improvement in running time. -0fast led to an -4%-ish performance regression.

By profiling with callgrind, we saw that execution was dominated by only one incredibly costly user-defined method, with an exclusive cost of 26.33% of total running time.

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	Incl.		Self		Called	Function	Location
		100.00		0.00	(0)	■ 0x00000000001030	ld-2.17.so
		100.00		0.00	2	_dl_runtime_resolve_xsave	ld-2.17.so
-		100.00		0.00	1	0x0000000000401df6	nbody
- 1		100.00		0.00	(0)	(below main)	libc-2.17.so
		100.00	1	9.33	1	main	nbody: main.cpp, basic_string.h, string_conversio
		89.22		26.33	672 000 056	Body::exert_force_unidirec	nbody: Body.cpp
		62.90		2.93	672 033 712	distance(double, double,	nbody: utils.cpp
1		59.97	1	11.70	672 033 711	■ hypot	libm-2.17.so
		48.27		48.27	672 033 712	hypot_finite	libm-2.17.so
1		0.63		0.63	48 000 000	Body::frog(double)	nbody: Body.cpp
- 1		0.52		0.52	48 000 008	Body::leap(double)	nbody: Body.cpp
		0.21		0.21	48 000 000	Body::reset_force()	nbody: Body.cpp
		0.07		0.00	601	dump_timestep(double, st	nbody: main.cpp, stl_vector.h, stl_iterator.h

Performance fixes were divised.

```
void Body::exert_force_unidirectionally(const Body& there) {
   const double m1 = m;
   const double m2 = there.m;

   const double delta_x = there.x - x;
   const double delta_y = there.y - y;

- const double r = distance(x, y, there.x, there.y);
   const double r = hypot(delta_x, delta_y);
   const double r2 = r * r;
   const double r2 = pow(r, 2);

- const double F = (G * m1 * m2) / r2;

// turn the displacement vector between our two points into a force vector
// of the desired magnitude
   const double scale_factor = F / r;

Fx += delta_x * scale_factor;
```

```
Fy += dumpsdelta_y * scale_factor;
}
```

Instead of recalculating  $\Delta x$  and  $\Delta y$  twice (the second time in distance), we calculate them only once and make a direct call to hypot, rather than making a call to distance. We also used the specialized pow provided by cmath. This yielded a recognizable 15%ish performance improvement.

```
diff --git a/Assignment_1/src/Body.cpp b/Assignment_1/src/Body.cpp
+void Body::exert_force_bidirectionally(Body& there) {
     const double m1 = m;
     const double m2 = there.m;
     const double delta_x = there.x - x;
     const double delta_y = there.y - y;
     const double r = hypot(delta_x, delta_y);
     const double r2 = pow(r, 2);
     const double F = (G * m1 * m2) / r2;
    // turn the displacement vector between our two points into a force vector
     // of the desired magnitude
     const double scale_factor = F / r;
    Fx += delta_x * scale_factor;
+
    Fy += delta_y * scale_factor;
     there.Fx -= delta_x * scale_factor;
     there.Fy -= delta_y * scale_factor;
+}
diff --git a/Assignment_1/src/main.cpp b/Assignment_1/src/main.cpp
@@ -258,8 +259,7 @@ int main(int argc, char **argv) {
                 for (size_t j = i + 1; j < bodies.size(); j++) {
                     auto& y = bodies[j];
                     x.exert_force_unidirectionally(y);
                     y.exert_force_unidirectionally(x);
                     x.exert_force_bidirectionally(y);
                 }
             }
         }
  This fix halved execution time, for obvious reasons. \(^1\).
diff --git a/Assignment_1/src/Body.cpp b/Assignment_1/src/Body.cpp
@@ -83,11 +82,11 @@ double Body::kinetic_energy() const {
double Body::gravitational_potential_energy(const Body& there) const {
     const double R = distance(x, y, there.x, there.y); // final distance, aka, to edge
```

<sup>&</sup>lt;sup>1</sup>I had to throw out all my old profile data

```
return (-G * m * there.m) / R;

return (-Gm * there.m) / R;

}

@@ -101,7 +100,7 @@ void Body::exert_force_unidirectionally(const Body& there) {
    const double F = (G * m1 * m2) / r2;
    const double F = (Gm * m2) / r2;

# const double F = (G * m1 * m2) / r2;

@@ -113,7 +112,6 @@ void Body::exert_force_unidirectionally(const Body& there) {
    const double F = (G * m1 * m2) / r2;
    const double F = (Gm * m2) / r2;

# const double F = (Gm * m2) / r2;

@@ -189,6 +189,7 @@ std::vector<Body> parse_input_file(std::ifstream& input_fh) {
    for (size_t i = 0; i < bodies.size(); i++) {
        bodies[i].m = masses[i];
        bodies[i].Gm = G * masses[i];
}</pre>
```

Precomputing Gm yielded a 3%ish performance improvement.

### References

- [1] J. E. Barnes and P. Hut, "A hierarchical O(n-log-n) force calculation algorithm," *Nature*, vol. 324, p. 446, 1986.
- [2] W. Dehnen and J. I. Read, "N-body simulations of gravitational dynamics," *European Physical Journal Plus*, vol. 126, p. 55, May 2011.
- [3] R. D. Skeel, "Variable step size destabilizes the störmer/leapfrog/verlet method," BIT Numerical Mathematics, vol. 33, pp. 172–175, Mar 1993.

outputs/in\_8\_body\_counterclockwise-8b38ef5.data 100 50 0 -50 -100 50 -100 -50 ò 100 x (m) 1e18 -6.5595 -6.5600 -6.5605 -6.5610 2 6 4 8 10 Ó t (s)

Figure 1: Observed energy anomaly while bodies in close proxmity - Leapfrog

Figure 2: Barnes-Hut energy variation - Leapfrog outputs/in\_8\_body\_counterclockwise-8b38ef5.data 100 50 0 -50 -100 -100 -50 ò 50 100 x (m) 1e18 -6.545 -6.550-6.555 -6.560 2 4 6 8 10 t (s)

Figure 3: Scaling characteristics

## 60000 timestep running time



