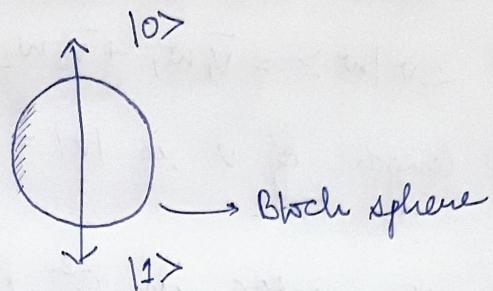


Quantum Computing

→ "Anyone who is not shocked by quantum theory has not understood it." - Niels Bohr.

→ A quantum bit, or qubit is the fundamental information unit of quantum computing.



$\vec{q} = a\vec{v} + b\vec{w}$, \vec{q} is a linear combination of vectors \vec{v} and \vec{w} and they are orthonormal basis (i.e. $\vec{v} \cdot \vec{w} = 0$). If a and b are non-zero the linear combination is non-trivial.

At any given time, q-bit is in a superposition state represented by linear combination of vectors $|0\rangle$ and $|1\rangle$ in C^2 :

$$a|0\rangle + b|1\rangle \text{ where } |a|^2 + |b|^2 = 1$$

probability
amplitudes

$\vec{v} = (v_1, v_2, v_3, \dots, v_n)$

$$= [v_1, v_2, v_3, \dots, v_n]$$

$$= \langle v | \quad \longrightarrow \text{"bra-}v\text{"}$$

transpose complex conjugate of $\langle v |$

$$= |w\rangle \quad \longrightarrow \text{"ket-}w\text{"}$$

$$\begin{bmatrix} w_1 \\ w_2 \\ w_3 \\ \vdots \\ w_n \end{bmatrix}$$

$|v><w|$ is the outer product,

$$|v><w| = \begin{bmatrix} v_1 \bar{w}_1 & v_1 \bar{w}_m \\ v_2 \bar{w}_2 & \vdots \\ \vdots & \vdots \\ v_n \bar{w}_1 & v_n \bar{w}_m \end{bmatrix}$$

$\langle v|w\rangle$ is the inner product,

$$\langle v|w\rangle = \bar{v}_1 w_1 + \bar{v}_2 w_2 + \dots + \bar{v}_n w_n$$

The length of v is $|v| = \sqrt{\langle v|v\rangle}$

$|0\rangle$ is the vector in C^2 that has coordinates $(1, 0)$ and $|1\rangle$ has coordinates $(0, 1)$ in the basis e_1 and e_2 .

They are frequently called computational basis.

$$\# |+\rangle = \left(\frac{\sqrt{2}}{2}, \frac{\sqrt{2}}{2} \right)$$

$$|-\rangle = \left(\frac{\sqrt{2}}{2}, -\frac{\sqrt{2}}{2} \right)$$

If L is a linear transformation on complex vector space, we write,

$$L|w\rangle,$$

If A is a matrix for L

$$A|w\rangle$$

Both A and w must be on the same basis.

property for $\langle v | L \rangle$ or $\langle v | A \rangle$

$$\langle v | L \times |w\rangle = \langle v | \times L |w\rangle = * \langle v | L |w\rangle$$
$$\langle v | A \times |w\rangle = \langle v | \times A |w\rangle = \langle v | A |w\rangle$$

A Hermitian matrix 'A' is a projection matrix if $A^2 = A$. The rank of the project is the dimension of its image vector space.

$$P(i) \rightarrow 1$$

$$P(i,j) \rightarrow 2$$

$$P(1, \dots, n) = I_n$$

Quantum state representation,

$$|\psi\rangle = a|0\rangle + b|1\rangle \xrightarrow{\text{quantum state}}$$

$$|\psi|^2 = |a|^2 + |b|^2 = 1$$

$$\boxed{\langle \psi | \psi \rangle = 1} \text{ same}$$

e.g. Let $a = \frac{1}{2}$ and $b = \frac{\sqrt{3}}{2}$. Then $|a|^2 = \frac{1}{4} = 0.25$

and $|b|^2 = \frac{3}{4} = 0.75$. This means that

the states of qubit has 25% chance of collapsing to $|0\rangle$ and 75% chance of collapsing to $|1\rangle$.

* If a complex number c in c is multiplied to $|\psi\rangle$, there's no change in probability if $|c|^2 = 1$.

* The state of single qubit can be represented by -

$$|\Psi\rangle = r_1 |0\rangle + r_2 e^{i\phi} |1\rangle$$

$$r_1^2 + r_2^2 = 1$$

$r_1, r_2 \in \mathbb{R}$

$0 \leq \phi < 2\pi$

$$r_1 = \cos\left(\frac{\theta}{2}\right)$$

$$r_2 = \sin\left(\frac{\theta}{2}\right)$$

$$0 \leq \theta \leq \pi$$

$$|\Psi\rangle = \cos\left(\frac{\theta}{2}\right) |0\rangle + \sin\left(\frac{\theta}{2}\right) e^{i\phi} |1\rangle$$

* Density Matrix.

for $|\Psi\rangle = a|0\rangle + b|1\rangle$, the density matrix of $|\Psi\rangle$ to be,

$$\begin{aligned} \rho &= |\Psi\rangle \langle \Psi| \\ &= \begin{bmatrix} a \\ b \end{bmatrix} \otimes \begin{bmatrix} \bar{a} & \bar{b} \end{bmatrix} \\ &= \begin{bmatrix} a\bar{a} & a\bar{b} \\ b\bar{a} & b\bar{b} \end{bmatrix} \\ &= \begin{bmatrix} |a|^2 & a\bar{b} \\ b\bar{a} & |b|^2 \end{bmatrix} \end{aligned}$$

$$\begin{aligned} \det|\rho| &= |a|^2 |b|^2 - a\bar{b} b\bar{a} \\ &= (ab)^2 - (ab)^2 \\ &= 0 \end{aligned}$$

* Observables and expectation.

$$M_0 = |0\rangle\langle 0| \quad \left. \begin{array}{l} \\ \end{array} \right\} \text{projectors}$$

$$M_1 = |1\rangle\langle 1|$$

$$|\Psi\rangle = a|0\rangle + b|1\rangle$$

$$\langle \Psi | M_0 | \Psi \rangle = \langle \Psi | (|0\rangle\langle 0|) | \Psi \rangle$$

$$= \cancel{\langle \Psi | 0 \rangle \langle 0 | \Psi \rangle}$$

$$= \langle \Psi | 0 \rangle \langle 0 | \Psi \rangle$$

$$= [\bar{a} \ \bar{b}] \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} a \\ b \end{bmatrix}$$

$$= [\bar{a} \ \bar{b}] \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} a \\ b \end{bmatrix}$$

$$= [\bar{a} \ \bar{b}] \begin{bmatrix} a \\ 0 \end{bmatrix}$$

$$= |a|^2$$

Similarly,

$$\langle \Psi | M_1 | \Psi \rangle = |b|^2$$

The eigenvalues of M_0 are 0 and 1

eigenvectors of M_0 are $|0\rangle$ and $|1\rangle$

The eigenvalues and eigenvectors of M_1 are same as M_0 .

$\therefore M_0$ and M_1 are examples of observables,
i.e. Hermitian matrices whose eigenvectors
form a basis for the quantum state
space.

Q. find the eigenvalues and eigenvectors of the matrix

$$A = \begin{bmatrix} 2 & 2 \\ 1 & 3 \end{bmatrix}$$

$$f_A(x) \Rightarrow |xI - A| = |A - xI| = 0$$

$$\Rightarrow (x-2)(x-3) - 2 = 0$$

$$\Rightarrow x^2 - 5x + 6 - 2 = 0$$

$$\Rightarrow x^2 - 5x + 4 = 0$$

$$x = 4, 1$$

$$\left. \begin{array}{l} \lambda_1 = 4 \\ \lambda_2 = 1 \end{array} \right\} \text{ eigenvalues}$$

$$[A - 4I][X_1] = 0$$

$$[A - I][X_2] = 0$$

$$\begin{bmatrix} -2 & 2 \\ 1 & -1 \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 1 & 2 \\ 1 & 2 \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{aligned} -2x + 2y &= 0 \\ x - y &= 0 \\ x &= y \\ -2x + 2x &= 0 \end{aligned}$$

Using elementary transformation

$$X_1 = \begin{bmatrix} 1 \\ 1 \end{bmatrix} \quad X_2 = \begin{bmatrix} -2 \\ 1 \end{bmatrix}$$

eigen vectors.

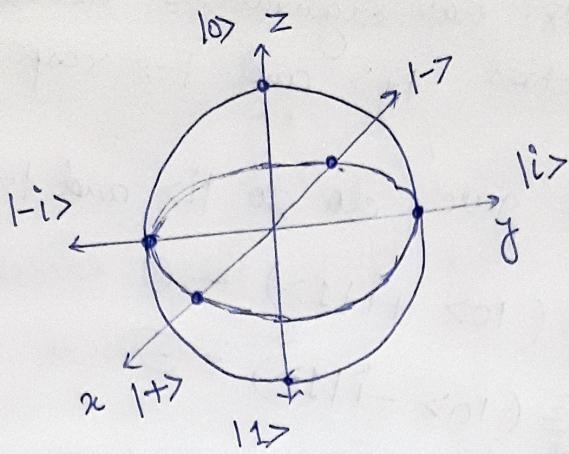
Measurable :-

$|\langle e_1 | \psi \rangle|^2 \rightarrow$ probability of getting e_1 ,

$|\langle e_2 | \psi \rangle|^2 \rightarrow$ probability of getting e_2

measurable .

* Bloch sphere



Basis

$$\{|0\rangle, |1\rangle\}$$

Common Name

Computational

$$\{|+\rangle, |-\rangle\}$$

Hadamard

$$\{|i\rangle, |-i\rangle\}$$

Circular

$$\hookrightarrow \{|u\rangle, |v\rangle\}$$

Bloch sphere
axis name

z

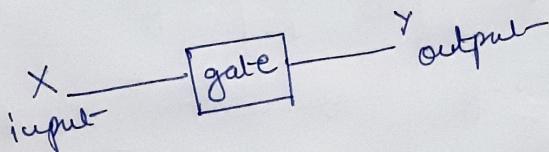
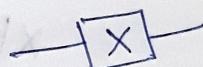
x

y

The Quantum X gate

$$I = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$

$$\sigma_x = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$$



$$\sigma_x |0\rangle = |1\rangle$$

$$\sigma_x |1\rangle = |0\rangle$$

$$\text{for, } |\psi\rangle = a|0\rangle + b|1\rangle$$

$$X|\psi\rangle = b|0\rangle + a|1\rangle$$

→ The X gate σ_x has eigenvalues +1 and -1 and eigenvectors $|+\rangle$ and $|-\rangle$ respectively.

Q. What does X gate do to $|i\rangle$ and $|-\bar{i}\rangle$?

$$|i\rangle = \frac{\sqrt{2}}{2} (|0\rangle + i|1\rangle) \quad \text{--- (1)}$$

$$|-\bar{i}\rangle = \frac{\sqrt{2}}{2} (|0\rangle - i|1\rangle) \quad \text{--- (2)}$$

$$A = X|i\rangle = \frac{\sqrt{2}}{2} (i|0\rangle + |1\rangle) \quad \text{--- (3)}$$

$$B = X|-\bar{i}\rangle = \frac{\sqrt{2}}{2} (-\bar{i}|0\rangle + |1\rangle) \quad \text{--- (4)}$$

Multiplying $-i$ in (3) and (4)

$$A = \frac{\sqrt{2}}{2} (|0\rangle - i|1\rangle)$$

$$B = \frac{\sqrt{2}}{2} (-|0\rangle - i|1\rangle)$$

$$\therefore X|i\rangle = -\frac{|-\bar{i}\rangle}{i}$$

$$X|-\bar{i}\rangle = \frac{|i\rangle}{i}$$

The quantum Z gate (phase flip gate) $I = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$

$$\sigma_z = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \quad \boxed{Z}$$

→ It rotates qubit states by π around the z -axis on Bloch sphere.

$$\begin{aligned}\sigma_z |+\rangle &= |-\rangle \\ \sigma_z |-\rangle &= |+\rangle \\ \sigma_z |i\rangle &= |-i\rangle \\ \sigma_z |-i\rangle &= |i\rangle\end{aligned}$$

swaps $|+\rangle, |-\rangle, |i\rangle, |-i\rangle$ but leaves $|0\rangle$ and $|1\rangle$.

for, $|\psi\rangle = a|0\rangle + b|1\rangle$

$$Z|\psi\rangle = a|0\rangle - b|1\rangle$$

→ σ_z has eigenvalues $+1$ and -1 and eigenvectors $|0\rangle$ and $|1\rangle$, respectively.

The quantum Y gate

$$\sigma_y = \begin{bmatrix} 0 & i \\ i & 0 \end{bmatrix} = i \begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix}$$

→ It rotates y-axis .

→ swaps $|0\rangle, |1\rangle, |+\rangle$ and ~~$|-\rangle$~~ $|-\rangle$ but leaves $|i\rangle$ and $|-i\rangle$ alone.



for, $|\psi\rangle = a|0\rangle + b|1\rangle$

$$Y|\psi\rangle = -bi|0\rangle + ai|1\rangle$$

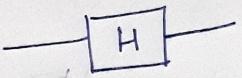
→ σ_y has eigenvalues $+1$ and -1 for eigen-vectors $|i\rangle$ and $|-i\rangle$, respectively.

The quantum ID gate



"does nothing"

The quantum H gate



$$H = \begin{bmatrix} \frac{\sqrt{2}}{2} & \frac{\sqrt{2}}{2} \\ \frac{\sqrt{2}}{2} & -\frac{\sqrt{2}}{2} \end{bmatrix} = \frac{\sqrt{2}}{2} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix}$$

$$H|0\rangle = \frac{\sqrt{2}}{2} (|0\rangle + |1\rangle) = |+\rangle$$

$$H|1\rangle = \frac{\sqrt{2}}{2} (|0\rangle - |1\rangle) = |- \rangle$$

$$H|+\rangle = H \left(\frac{\sqrt{2}}{2} (|0\rangle + |1\rangle) \right)$$

$$= \frac{\sqrt{2}}{2} (H|0\rangle + H|1\rangle)$$

$$= |0\rangle$$

$$H|-\rangle = |1\rangle$$

The quantum R_ϕ^z gate.

$$R_\phi^z = \begin{bmatrix} 1 & 0 \\ 0 & e^{i\phi} \end{bmatrix}$$

$$\begin{aligned} R_0^z &= ID \\ R_1^z &= Z \end{aligned} \quad \left. \right\}$$

The quantum S gate

$$S = \begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix}$$



The quantum $\sqrt{\text{NOT}}$ gate

$$= \frac{1}{2} \begin{bmatrix} 1+i & 1-i \\ 1-i & 1+i \end{bmatrix}$$



$$\begin{aligned} &+ |0\rangle\langle 0|_{AB} = |00\rangle\langle 00| \\ &+ |1\rangle\langle 1|_{AB} \\ &+ |01\rangle\langle 11|_{AB} \\ &+ |11\rangle\langle 01|_{AB} \end{aligned}$$

Quantum Entanglement

Quantum entanglement is the phenomenon that occurs when a group of particles are generated, interact or share spatial proximity in a way such that quantum state of each particle of group cannot be described independently of the state of others, including when the particles are separated by a large distance.

→ Let q_1 and q_2 be two qubits and let $\{|0\rangle, |1\rangle\}$, $\{|0\rangle_2, |1\rangle_2\}$ be standard orthonormal basis kets.

$$\begin{aligned} |\Psi\rangle &= a_1|0\rangle + b_1|1\rangle, & |a_1|^2 + |b_1|^2 &= 1 \\ |\Psi\rangle_2 &= a_2|0\rangle_2 + b_2|1\rangle_2 & |a_2|^2 + |b_2|^2 &= 1 \end{aligned}$$

The four kets,

$$\begin{aligned} |0\rangle, \otimes |0\rangle_2, |0\rangle, \otimes |1\rangle_2 \\ |1\rangle, \otimes |0\rangle_2, \cancel{|0\rangle}, |1\rangle, \otimes |1\rangle_2 \end{aligned}$$

are a basis for the combined set space.

$$\begin{aligned} |\Psi\rangle, \otimes |\Psi\rangle_2 &= a_1a_2|0\rangle, \otimes |0\rangle_2 + \\ &\quad a_1b_2|0\rangle, \otimes |1\rangle_2 + \\ &\quad b_1a_2|1\rangle, \otimes |0\rangle_2 + \\ &\quad b_1b_2|1\rangle, \otimes |1\rangle_2. \end{aligned}$$

Simplifying -

$$|\Psi\rangle_1 \otimes |\Psi\rangle_2 = a_1 a_2 |0\rangle_1 |0\rangle_2 + \\ a_1 b_2 |0\rangle_1 |1\rangle_2 + \\ b_1 a_2 |1\rangle_1 |0\rangle_2 + \\ b_1 b_2 |1\rangle_1 |1\rangle_2 .$$

$$|\Psi\rangle_1 \otimes |\Psi\rangle_2 = a_1 a_2 |0\rangle |0\rangle + \\ a_1 b_2 |0\rangle |1\rangle + \\ b_1 a_2 |1\rangle |0\rangle + \\ b_1 b_2 |1\rangle |1\rangle$$

$$|\Psi\rangle_1 \otimes |\Psi\rangle_2 = a_1 a_2 |00\rangle + \\ a_1 b_2 |01\rangle + \\ a_2 b_1 |10\rangle + \\ b_1 b_2 |11\rangle$$

$$|0\rangle = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$

$$|1\rangle = \begin{bmatrix} 0 \\ 1 \end{bmatrix}$$

$$|01\rangle = \begin{bmatrix} 1 \\ 0 \end{bmatrix} \otimes \begin{bmatrix} 0 \\ 1 \end{bmatrix}$$

$$= \begin{bmatrix} 1[0] \\ 0[1] \end{bmatrix}$$

$$= \begin{bmatrix} 0 \\ 1 \\ 0 \\ 0 \end{bmatrix}$$

$$|00\rangle = \begin{bmatrix} 1 \\ 0 \\ 0 \\ 0 \end{bmatrix} \quad |01\rangle = \begin{bmatrix} 0 \\ 1 \\ 0 \\ 0 \end{bmatrix} \quad |10\rangle = \begin{bmatrix} 0 \\ 0 \\ 1 \\ 0 \end{bmatrix}$$

$$|11\rangle = \begin{bmatrix} 0 \\ 0 \\ 0 \\ 1 \end{bmatrix}$$

$\rightarrow \langle 01|01\rangle = 1$

$\langle 01|11\rangle = 0$

same → 1
diff → 0

$$|\Psi\rangle_1 \otimes |\Psi\rangle_2 = a_1 a_2 |00\rangle + a_1 b_2 |01\rangle + b_1 a_2 |10\rangle + b_1 b_2 |11\rangle$$

$$|a_1 a_2|^2 + |a_1 b_2|^2 + |b_1 a_2|^2 + |b_1 b_2|^2 = 1$$

we can prove this

- Before measurement the qubits are in a state of entanglement.
- The entangled state is known as Bell state.

$$|\Phi^+\rangle = \frac{\sqrt{2}}{2} |00\rangle + \frac{\sqrt{2}}{2} |11\rangle$$

$$|\Phi^-\rangle = \frac{\sqrt{2}}{2} |00\rangle - \frac{\sqrt{2}}{2} |11\rangle$$

$$|\Psi^+\rangle = \frac{\sqrt{2}}{2} |01\rangle + \frac{\sqrt{2}}{2} |10\rangle$$

$$|\Psi^-\rangle = \frac{\sqrt{2}}{2} |01\rangle - \frac{\sqrt{2}}{2} |10\rangle$$

Q How to determine if $|\Psi\rangle$ is entangled?

Ans: Let $|\Psi\rangle$ be a 2-qubit quantum state in $C^2 \otimes C^2$.

$|\Psi\rangle$ is entangled if it cannot be written as the tensor products of two 1-qubit kets.

$$|\Psi^+\rangle = 0|00\rangle + \frac{\sqrt{2}}{2}|01\rangle + \frac{\sqrt{2}}{2}|10\rangle + 0|11\rangle \\ = a_1a_2|00\rangle + a_1b_2|01\rangle + a_2b_1|10\rangle + b_1b_2|11\rangle$$

$$a_1a_2 = 0 \quad a_1b_2 = \frac{\sqrt{2}}{2} \quad a_2b_1 = \frac{\sqrt{2}}{2} \quad b_1b_2 = 0$$

either $a_1 = 0, a_2 = 0$

which contradicts both a_1b_2 and a_2b_1

This means $|\Psi^+\rangle$ cannot be written as the tensor product of two 1-qubit kets and it is an entangled state.

→ If a quantum state is not entangled then it is separable.

Q Is $\frac{\sqrt{2}}{2}|00\rangle + \frac{\sqrt{2}}{2}|01\rangle$ an entangled state?

Ans: $a_1a_2 = \frac{\sqrt{2}}{2}$ $a_1b_2 = \frac{\sqrt{2}}{2}$ $a_2b_1 = 0$ $b_1b_2 = 0$
 $b_1 = 0$ $a_2 \neq 0$ $b_1 = 0$ $b_2 \neq 0$

satisfies 1-qubit state

∴ not - entangled
i.e separable

* The general case

n -bit quantum space

$2^n \rightarrow$ state space of dimension.

$$|\Psi\rangle = a_0|0\rangle_n + a_1|1\rangle_n + a_2|2\rangle_n + \dots + a_{2^n-1}|2^n-1\rangle_n$$

$$|a_0|^2 + |a_1|^2 + \dots + |a_{2^n-1}|^2 = 1$$

$$\underbrace{|00000\dots000\rangle}_n = |0\rangle^{\otimes n}$$

* Density Matrix

$$\rho = |\Psi\rangle\langle\Psi|$$

$$= \begin{bmatrix} |a_0|^2 & a_1\bar{a}_0 & a_2\bar{a}_0 & \dots & a_{2^n-1}\bar{a}_0 \\ a_0\bar{a}_1 & |a_1|^2 & & & \\ \vdots & & \ddots & & \\ a_0\bar{a}_{2^n-1} & & & \ddots & |a_{2^n-1}|^2 \end{bmatrix}$$

$$\text{tr}(\rho) = 1$$

The quantum $H^{\otimes n}$ gate

$$H = \frac{\sqrt{2}}{2} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix}$$

$$|\Psi_1\rangle = a_1|0\rangle + b_1|1\rangle, \quad |\Psi_2\rangle = a_2|0\rangle_2 + b_2|1\rangle_2$$

$$(H|\Psi_1\rangle) \otimes (H|\Psi_2\rangle)$$

$$= (H \otimes H)(|\Psi_1\rangle \otimes |\Psi_2\rangle)$$

$$= H^{\otimes 2} (|\Psi_1\rangle \otimes |\Psi_2\rangle)$$

$$H^{\otimes 2} = \begin{bmatrix} \frac{1}{2} & \frac{1}{2} & \frac{1}{2} & \frac{1}{2} \\ \frac{1}{2} & -\frac{1}{2} & \frac{1}{2} & -\frac{1}{2} \\ \frac{1}{2} & \frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \\ \frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} & \frac{1}{2} \end{bmatrix}$$

$$= \frac{1}{2} \begin{bmatrix} 1 & 1 & 1 & 1 \\ 1 & -1 & 1 & -1 \\ 1 & 1 & -1 & -1 \\ 1 & -1 & -1 & 1 \end{bmatrix}$$

$$= \frac{\sqrt{2}}{2} \begin{bmatrix} H & H \\ H & -H \end{bmatrix}$$

$$H^{\otimes 3} = \frac{\sqrt{2}}{2} \begin{bmatrix} H^{\otimes 2} & H^{\otimes 2} \\ H^{\otimes 2} & -H^{\otimes 2} \end{bmatrix}$$

Recursion? Yes.

$$H^{\otimes n} = \frac{\sqrt{2}}{2} \begin{bmatrix} H^{\otimes n-1} & H^{\otimes n-1} \\ H^{\otimes n-1} & -H^{\otimes n-1} \end{bmatrix}$$

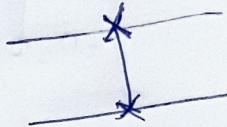
$$H^{\otimes n}|0\rangle_n = \frac{1}{\sqrt{2^n}} \sum_{j=0}^{n-1} |j\rangle_n$$

→ Dropping subscript.

$$H^{\otimes n}|0\rangle = \frac{1}{\sqrt{2^n}} \sum_{j=0}^{n-1} |j\rangle$$

Normalization constant.

The quantum SWAP gate



$$M(|\Psi_1\rangle \otimes |\Psi_2\rangle) = |\Psi_2\rangle \otimes |\Psi_1\rangle$$

$$M \begin{bmatrix} v_1 \\ v_2 \\ v_3 \\ v_4 \end{bmatrix} = \begin{bmatrix} v_1 \\ v_3 \\ v_2 \\ v_4 \end{bmatrix} ; \quad v_2, v_3 \rightarrow \text{interchange.}$$

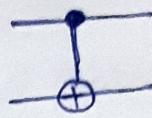
$$M = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$

The quantum CNOT gate | CX gate

"commonly used to create entangled qubits".

C → controlled

CNOT is conditional bit flip.



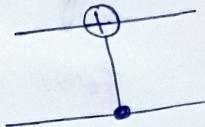
$\text{ID} |\Psi_1\rangle \text{ then } X |\Psi_2\rangle \text{ OR}$

$\text{ID} |\Psi_2\rangle \text{ then } X |\Psi_1\rangle$

$$\text{CNOT} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \\ v_4 \end{bmatrix} = \begin{bmatrix} v_1 \\ v_2 \\ v_4 \\ v_3 \end{bmatrix} \quad v_3, v_4 \rightarrow \text{interchange}$$

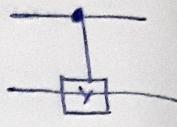
$$\text{CNOT} = \left[\begin{array}{cc|cc} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{array} \right] \xrightarrow{x}$$

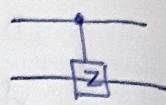
reverse CNOT



$v_1, v_2 \rightarrow \text{interchange?}$

#The quantum CY and CZ gate.

$$CY = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & -i \\ 0 & 0 & i & 0 \end{bmatrix}$$


$$CZ = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{bmatrix}$$


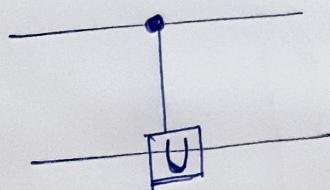
→ The CZ is conditional sign flip.

}

In general,

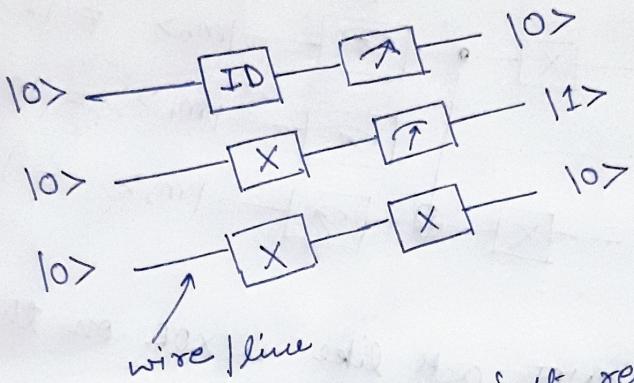
$$\text{if } U = \begin{bmatrix} a & b \\ c & d \end{bmatrix}$$

$$\text{controlled-}U(\text{CU}) = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & a & b \\ 0 & 0 & c & d \end{bmatrix}$$

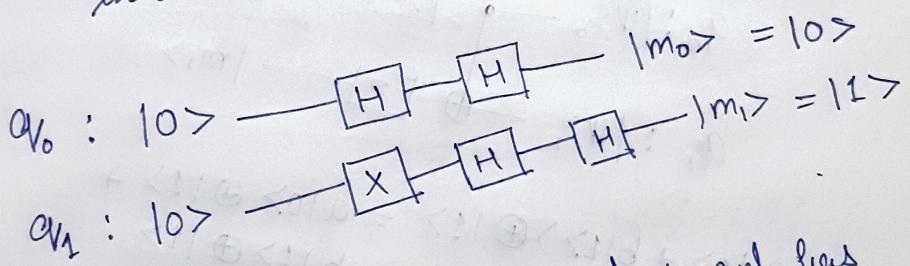


WIRING UP THE CIRCUITS

- quantum registers :- A quantum register is a collection of qubits we use for computation.
- A quantum circuit is a sequence of gates applied to one or more qubits in a quantum register.



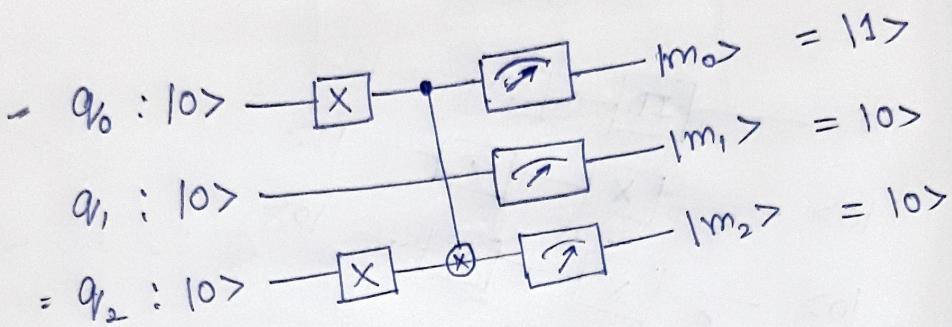
- If nothing is happening just remove the intermediate gates.



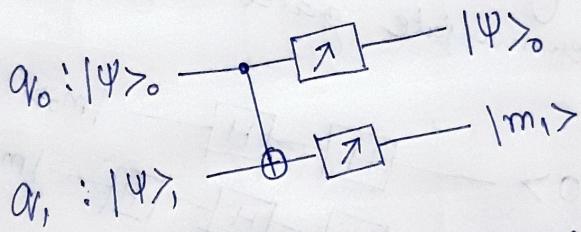
The first wire has depth 2 and second has depth 3.

$$* \text{CNOT } |100\rangle = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix} \begin{bmatrix} 1 \\ 0 \\ 0 \\ 0 \end{bmatrix} = \begin{bmatrix} 1 \\ 0 \\ 0 \\ 0 \end{bmatrix} = |100\rangle$$

$$\text{CNOT } |111\rangle = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix} \begin{bmatrix} 0 \\ 0 \\ 0 \\ 1 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 1 \\ 0 \end{bmatrix} = |110\rangle$$



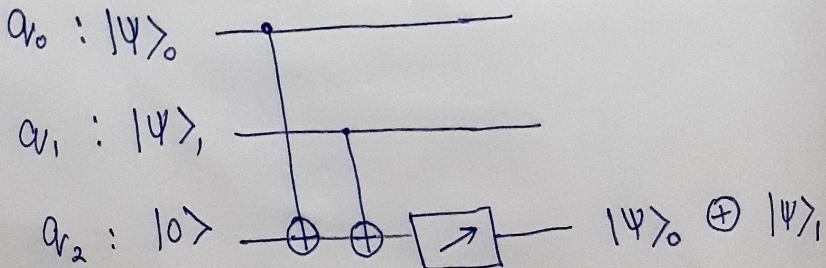
* The basic CNOT acts like a XOR on the standard basis.



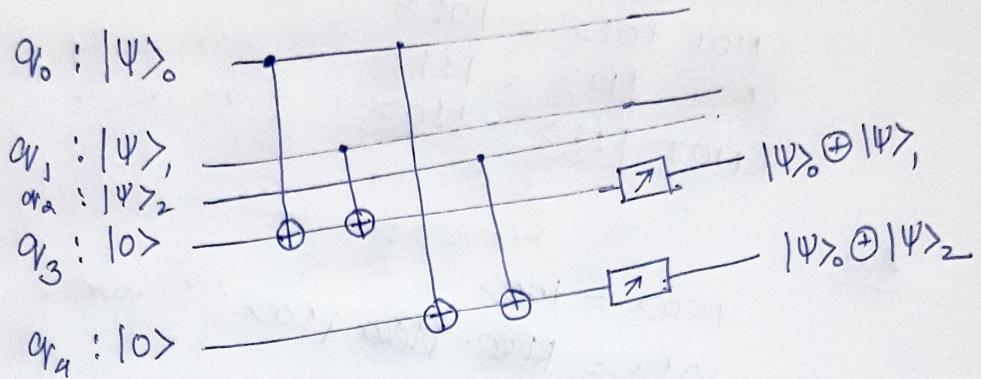
$$(a|0\rangle + b|1\rangle) \oplus |1\rangle = a|0\rangle \oplus |1\rangle + b|1\rangle \oplus |1\rangle$$

$$= a|1\rangle + b|0\rangle$$

\rightarrow If we want to modify one of the input qubits, we can put the value of the XOR in a third output, or ancilla qbit.

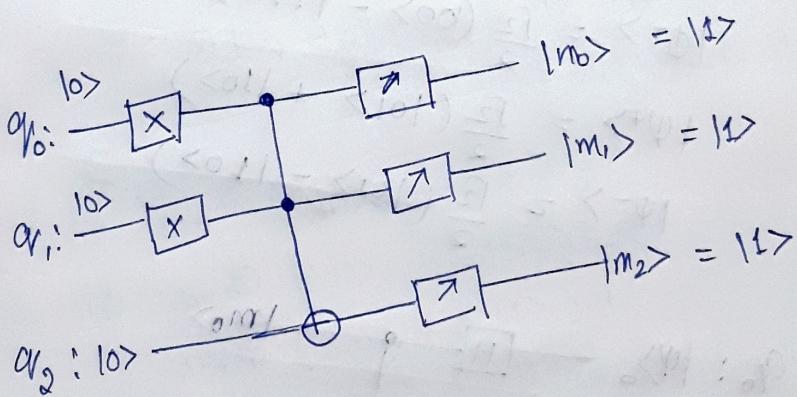


Q. Suppose we have three input qubits, q_0 , q_1 , and q_2 in states $|0\rangle$, $|1\rangle$, and $|1\rangle$ respectively. We want to put $|q_0\rangle \oplus |q_1\rangle$ in ancilla qubit q_3 and put $|q_0\rangle \oplus |q_2\rangle$ in ancilla qubit q_4 . Draw the circuit.



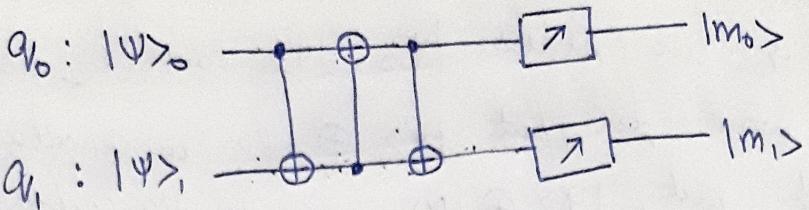
The quantum Toffoli gate

The quantum Toffoli gate operates on three qubits. If the first two qubits are $|1\rangle$ then it flips the third, otherwise does nothing.



- The toffoli gate can be used to create the quantum equivalent of a NAND gate.

Q.



$$CNOT|00\rangle = |00\rangle$$

$$CNOT|01\rangle = |01\rangle$$

$$CNOT|10\rangle = |11\rangle$$

$$CNOT|11\rangle = |10\rangle$$

Ans:

$$|00\rangle = |00\rangle$$

$$|01\rangle = \cancel{|00\rangle} \quad \cancel{|10\rangle} \quad |10\rangle$$

$$|10\rangle = |01\rangle$$

$$|11\rangle = |11\rangle$$

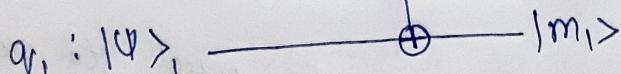
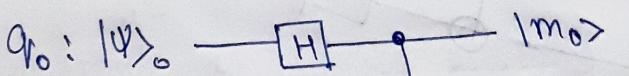
Q. Show that circuit creates bell states.

$$|\phi^+\rangle = \frac{\sqrt{2}}{2}(|00\rangle + |11\rangle)$$

$$|\phi^-\rangle = \frac{\sqrt{2}}{2}(|00\rangle - |11\rangle)$$

$$|\psi^+\rangle = \frac{\sqrt{2}}{2}(|01\rangle + |10\rangle)$$

$$|\psi^-\rangle = \frac{\sqrt{2}}{2}(|01\rangle - |10\rangle)$$



$$|00\rangle = CNOT(H|0\rangle \otimes |0\rangle)$$

$$= \frac{\sqrt{2}}{2}(|0\rangle + |1\rangle)$$

$$|PS\rangle = \frac{\sqrt{2}}{2}(|00\rangle + |10\rangle)$$

$$CNOT|\rho\rangle = \frac{\sqrt{2}}{2}(|00\rangle + |11\rangle)$$

Quantum Teleportation

→ My qubit Q is in some arbitrary quantum state $|Q\rangle_0 = a|0\rangle + b|1\rangle$.

When we are done, you will know this state but I will no longer have access to it.

M → my qubit

Y → your qubit

Q → where state I want to transfer from me to you.

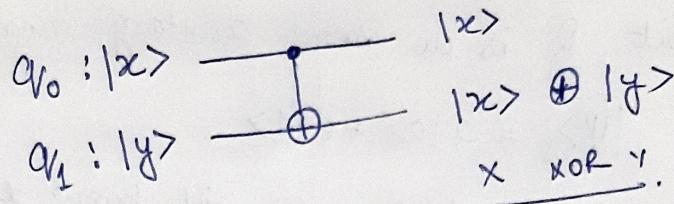
We begin by entangling M and Y
for that,

$$|\phi^+\rangle_{Mx} = \frac{\sqrt{2}}{2} (|00\rangle_{Mx} + |11\rangle_{Mx})$$

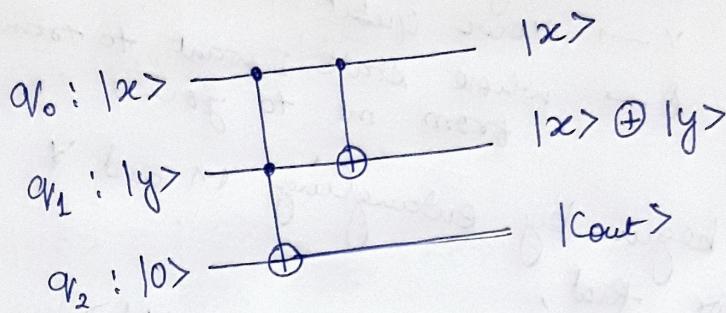
$$\begin{aligned} |\Psi\rangle_0 \otimes |\phi^+\rangle_{Mx} &= (a|0\rangle_0 + b|1\rangle_0) \otimes \left(\frac{\sqrt{2}}{2} (|00\rangle_{Mx} + |11\rangle_{Mx}) \right) \\ &= \frac{\sqrt{2}}{2} (a|000\rangle_{QMx} + a|011\rangle_{QMx} + b|100\rangle_{QMx} + b|111\rangle_{QMx}) \\ &= \frac{\sqrt{2}}{2} (a|00\rangle_{QM} \otimes |0\rangle_y + a|01\rangle_{QM} \otimes |1\rangle_y + b|10\rangle_{QM} \otimes |0\rangle_y \\ &\quad + b|11\rangle_{QM} \otimes |1\rangle_y) \\ &= \frac{1}{2} (|\phi^+\rangle_{QM} \otimes (a|0\rangle_y + b|1\rangle_y) + \\ &\quad |\phi^-\rangle_{QM} \otimes (a|0\rangle_y - b|1\rangle_y) + \\ &\quad |\psi^+\rangle_{QM} \otimes (b|0\rangle_y + a|1\rangle_y) + \\ &\quad |\psi^-\rangle_{QM} \otimes (-b|0\rangle_y + a|1\rangle_y)) \end{aligned}$$

Arithmetic

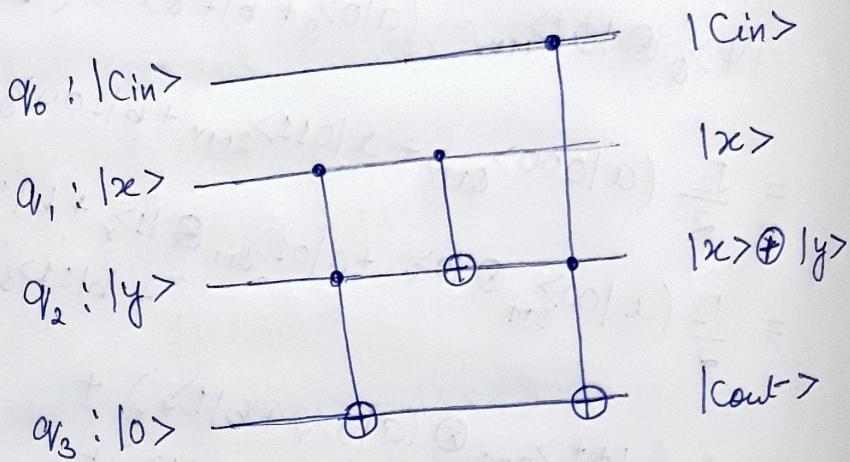
* XOR



"without carry out"



"with carry out"



"with carry in"

* SUM

