www.afm-journal.de

# 2D Transition Metal Dichalcogenide Heterostructures for p- and n-Type Photovoltaic Self-Powered Gas Sensor

Youngjun Kim, Sangyoon Lee, Jeong-Gyu Song, Kyung Yong Ko, Whang Je Woo, Suk Woo Lee, Minwoo Park, Hoonkyung Lee, Zonghoon Lee, Hyunyong Choi, Woo-Hee Kim,\* Jusang Park,\* and Hyungjun Kim\*

2D transition metal dichalcogenides (TMDs) have attracted much attention for their gas sensing applications due to their superior responsivity at typical room temperature. However, low power consumption and reliable selectivity are the two main requirements for gas sensors to be applicable in future electronic devices. Herein, a p-type (WSe2/WS2) and n-type (MoS2/WSe2) photovoltaic self-powered gas sensor is demonstrated using 2D TMD heterostructures for the first time. The gas sensors are operated by the photovoltaic effect of 2D TMD heterostructures, which are uniformly synthesized by the vacuum-based synthesis. The gas sensing properties of the WSe<sub>2</sub>/WS<sub>2</sub> and MoS<sub>2</sub>/WSe<sub>2</sub> heterostructure gas sensors are investigated for NO<sub>2</sub> and NH<sub>3</sub> with changing gas concentration, and each sensor exhibits selectivity to NO2 and NH<sub>3</sub>. From the results, it is confirmed that the 2D TMD heterostructures can be a viable platform for highly sensitive, selective gas sensor applications without external bias due to their photovoltaic features. Further, this study contributes toward revealing the gas sensor mechanism in 2D TMD heterostructure.

Dr. Y. Kim, S. Lee, Dr. J.-G. Song, Dr. K. Y. Ko, W. J. Woo, Dr. J. Park, Prof. H. Kim

School of Electrical and Electronic Engineering

Yonsei University

50 Yonsei-Ro, Seodaemun-Gu, Seoul 03722, Republic of Korea

E-mail: jusang@yonsei.ac.kr; hyungjun@yonsei.ac.kr

S. W. Lee, Prof. Z. Lee

School of Materials Science and Engineering

Ulsan National Institute of Science and Technology (UNIST)

50 UNIST-gil, Ulsan 44919, Republic of Korea

M. Park, Prof. H. Lee

Department of Physics

Konkuk University

120 Neungdong-ro, Gwangjin-gu, Seoul 05029, Republic of Korea

Prof. H. Choi

Department of Physics and Astronomy

Seoul National University

1, Gwanak-ro, Gwanak-gu, Seoul 08826, Republic of Korea

Prof. W.-H. Kim

Department of Materials Science and Chemical Engineering

Hanyang University

55, Hanyangdaehak-ro, Sangnok-gu, Ansan-si

Gyeonggi-do 15588, Republic of Korea

E-mail: wooheekim@hanyang.ac.kr

The ORCID identification number(s) for the author(s) of this article can be found under https://doi.org/10.1002/adfm.202003360.

DOI: 10.1002/adfm.202003360

## 1. Introduction

Gas sensing technologies have recently drawn much attention due to the significantly increasing threats of industrial chemicals as well as daily air pollution caused by exhaust gases from automobiles, huge power plants, etc.[1-3] For daily monitoring and control of unpredictable types and concentrations of toxic gas, the need for highly selective and sensitive gas sensors has become increasingly important. In particular, semiconductor gas sensors have been considered as a prospective candidate for future gas sensor networks due to their simple sensing mechanism, which is based on resistance/current changes upon gas exposure.[4,5]

However, low power consumption and reliable selectivity are the two main requirements for gas sensors to be applicable in future electronic devices. Connected Internet of things (IoTs) devices

(e.g., smart phones and wireless sensor platforms) will facilitate personalized detection of environmental conditions, making them the basis of the future core technology of ubiquitous sensing, [6,7] Therefore, recent works have demonstrated promising concepts for realizing self-powered gas sensors that are capable of detecting gases without the need for external power sources to activate the sensor-gas interaction or to actively generate a read-out signal.[8,9] These sensors consume drastically less power than conventional semiconductor gas sensors, and they reduce the required space for integration.

Since 2D transition metal dichalcogenides (TMDs) are promising semiconducting materials as an electronic device for low power consumption and have sensitive gas detection properties at room temperature, [10-12] various studies on 2D TMD-based gas sensing devices have been reported.[13-16] Furthermore, 2D TMD heterostructure, such as MoS<sub>2</sub>/WSe<sub>2</sub>, has photovoltaic properties with open-circuit voltage  $(V_{oc})$  of 0.2-0.5 V, implying that atomically thin 2D TMD heterostructure could be used for photovoltaic self-powered gas sensors at typical room temperatures under white light illumination. [17,18]

Nevertheless, photovoltaic self-powered gas sensor using 2D TMD heterostructure has not yet been reported. This might be due to the difficulty in achieving a vertically stacked structure, which only one material is exposed to the surface. As 2D TMD www.advancedsciencenews.com



www.afm-journal.de

have different responsivity to gases, gas selectivity could be achieved only if they are vertically stacked. Note that most of 2D TMD heterostructures have been fabricated by the exfoliation or chemical vapor deposition (CVD), which inevitably exposes both TMD materials to air, which are not a suitable structure to secure gas selectivity.<sup>[17–22]</sup>

Herein, we have demonstrated a photovoltaic self-powered gas sensor through the 2D TMD heterostructures for the first time. Vertically stacked WSe<sub>2</sub>/WS<sub>2</sub> and MoS<sub>2</sub>/WSe<sub>2</sub> heterotructures were fabricated to the gas sensor and operated under white light illumination condition. The gas sensing experiments were performed using NO2 and NH3, which were representative electron acceptor and donor molecules, respectively.[12,14-16] Note that both 2D TMD heterostructures were synthesized by the vacuum-based synthesis for large-area uniformity. [15,23-25] For p-type surface, the WSe<sub>2</sub>/WS<sub>2</sub> was synthesized by in situ CVD and one transfer process was required. The n-type surface of MoS<sub>2</sub>/WSe<sub>2</sub> was synthesized by ex situ process and fabricated as a gas sensor using lithography to ensure its integration. In each gas sensor, only the WSe2 and MoS2 were exposed to the surface, respectively, and fabrication methods are suitable for demonstrating the photocurrent without external voltage due to photovoltaic characteristics. The WSe2/WS2 gas sensor presented the p-type behavior and selectivity to NO2 gas. On the other hand, MoS2/WSe2 gas sensor showed the n-type gas sensing characteristic and responsivity to NH3 gas was higher than the NO<sub>2</sub> gas. We found that the exposed surface plays the most dominant role in transferring the charges from the gas, even at monolayer thickness. Further, it was confirmed that the underneath layer could affect the responsivity depending on the type of the heterostructure material and exposed gas through computational research.

# 2. Results and Discussion

Two different methods were used to implement the photovoltaic self-powered gas sensors using 2D TMD heterostructures. The first process was for the fabrication of the WSe2/ WS<sub>2</sub> heterostructure gas sensor. WSe<sub>2</sub>, which is a p-type material, was exposed to the surface for p-type cell, and WS2, which is an n-type material, was located under the WSe<sub>2</sub> layer. [15,16,25] We attempted to form a heterostructure with few-layer thickness in each layers because the photovoltaic effect became too low in a monolayer/monolayer heterostructure due to the decreased light absorption. [17,19] On the contrary, if the thickness is too thick, it is difficult to use gas sensor at typical room temperatures because of a slow recovery rate.  $[^{26,27}]$  The uniform WSe2/WS2 was prepared by the in situ synthesis. It was composed of consecutive two-step CVD process for synthesizing centimeter-scale uniform WS<sub>2</sub> and WSe<sub>2</sub>, respectively.<sup>[15,25]</sup> Atomic force microscopy (AFM) height profile analysis of the as-deposited WS2 and the WSe2/WS2 heterostructure was implemented to measure the thickness and surface morphologies. Each sample was transferred onto a bare SiO<sub>2</sub>/Si substrate. The thickness of WS<sub>2</sub> was  $\approx$ 2.2 nm, indicating a trilayer (3L) (Figure 1a), and the overall thickness of the WSe<sub>2</sub>/WS<sub>2</sub> heterostructure was ≈4.5 nm (Figure 1b), indicating the additional 3L WSe2 on the 3L WS2. And, smooth surface could be identified

by the AFM images (Figure 1a,b). The synthesized WSe<sub>2</sub>, WS<sub>2</sub>, and WSe<sub>2</sub>/WS<sub>2</sub> heterostructure were thoroughly examined by Raman spectroscopy, as shown in Figure 1c. WSe<sub>2</sub> presented the Raman spectrum bands at  $\rm E^{1}_{2g}$  and  $\rm A_{1g}$  modes from in-plane and out-of-plane vibrations at  $\approx 250$  and 258 cm $^{-1}$ , respectively (yellow line), while WS $_{2}$  exhibited in-plane  $\rm E^{1}_{2g}$  and out-of-plane  $\rm A_{1g}$  modes at  $\approx 350$  and 419 cm $^{-1}$ , respectively (purple line). In the Raman spectra for the WSe<sub>2</sub>/WS $_{2}$  heterostructure, the main peaks for WSe $_{2}$  and WS $_{2}$  are represented as the red line. There was no spectrum of secondary phases such as WS $_{x}$ Se $_{2-x}$  (alloy formation) causing  $\rm A_{1g}$  mode at 380 cm $^{-1}$ .  $^{[29-31]}$  This confirmed that the WSe $_{2}$ /WS $_{2}$  heterostructure was successfully formed through the in situ CVD synthesis process.

To further investigate the structure and composition of the WSe<sub>2</sub>/WS<sub>2</sub> heterostructure, transmission electron microscopy (TEM) was conducted, as shown in Figure 1d. The cross-sectional TEM image clearly presents the formation of the vertically stacked heterostructure. A distinguishable layer structure evidently verifies the layer-by-layer growth for monolithic configuration with the precise control by the in situ CVD. In addition, the whole structure consisting of 3L for each WSe2 and WS<sub>2</sub> material was consistent with the AFM height analysis. Moreover, the energy-dispersive X-ray spectroscopy (EDS) mapping images in Figure 1e-h and high angle annular dark field (HAADF) image in Figure S1b in the Supporting Information further confirm the elemental composition of the vertical WSe<sub>2</sub>/WS<sub>2</sub> heterostructure. W atoms were well distributed across the entire heterostructure, while S and Se atoms were clearly distinguishable in the separated layered structure. From those measurements, the synthesized WSe2/WS2 showed the layer-by-layer CVD growth with highly controlled thickness, eventually achieving a vertically stacked heterostructure. As presented in Figure S1c,d in the Supporting Information, the highresolution top-view image TEM and the selected area electron diffraction pattern (SAED) exhibit the polycrystalline structures of the synthesized WSe<sub>2</sub>/WS<sub>2</sub> heterostructure. The grain size was measured to be ≈30 nm.

Based on the vertically stacked WSe <sub>2</sub>/WS<sub>2</sub> heterostructure, the gas sensor was fabricated with bottom (Au) and top (Pd) electrodes, as depicted in Figure 1i. This structure was designed to contact the electrodes with WSe<sub>2</sub> and WS<sub>2</sub>, respectively, and only expose the WSe<sub>2</sub> surface to the gases. A photograph of the completely fabricated WSe<sub>2</sub>/WS<sub>2</sub> gas sensor is shown in Figure 1j. The optical microscopy (OM) images in Figure S1e–g in the Supporting Information show the successive fabrication processes of uniform WSe<sub>2</sub>/WS<sub>2</sub> gas sensor from the bottom electrode patterning to the top electrode patterning.

Next to the p-type surface heterostructure, a self-powered gas sensor with n-type surface, MoS<sub>2</sub>/WSe<sub>2</sub> heterostructure, was fabricated. We synthesized the MoS<sub>2</sub> and WSe<sub>2</sub> as n-type and p-type materials using previously developed vacuum-based synthesis processes, respectively.<sup>[23,24]</sup> In this study, 3L WSe<sub>2</sub> on SiO<sub>2</sub> was synthesized by a self-limiting layer synthesis and monolayer (1L) MoS<sub>2</sub> was subsequently synthesized on 3L WSe<sub>2</sub>.<sup>[23,24]</sup> The 1L MoS<sub>2</sub> and 3L WSe<sub>2</sub> served sufficient light absorption thickness for photovoltaic effect.<sup>[17,32]</sup> Note that vertically stacked MoS<sub>2</sub> and WSe<sub>2</sub> layers with wafer-scale uniformity were confirmed to have n-type and p-type, respectively, in terms of their field-effect transistor (FET) characteristics.<sup>[23,24]</sup>

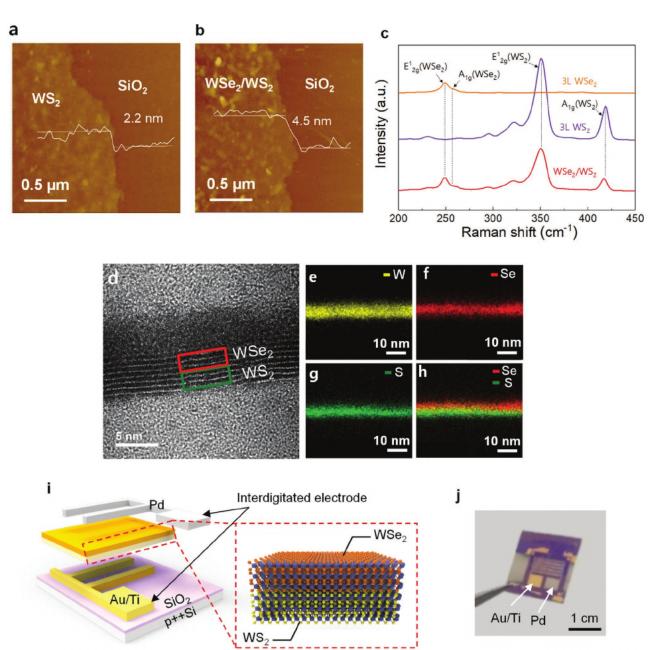


Figure 1. AFM image and height profile of a)  $WS_2$  and b)  $WSe_2/WS_2$  heterostructure. c) Raman spectra for  $WSe_2/WS_2$  heterostructure with  $WSe_2$  and  $WS_2$ . d) Cross-sectional TEM image of vertically stacked  $WSe_2/WS_2$  heterostructure. Cross-sectional EDS mapping images indicating e) W atoms, f) Se atoms, g) S atoms, and h) S and Se atoms. i) Schematic of the  $WSe_2/WS_2$  gas sensor. j) Photographic image of the fabricated  $WSe_2/WS_2$  gas sensor.

AFM was used for measuring the thickness of MoS<sub>2</sub>/WSe<sub>2</sub> and Raman spectroscopy was employed to confirm whether an alloy was formed between two materials during the synthesis of MoS<sub>2</sub> on WSe<sub>2</sub>. [32] **Figure 2a**,b shows AFM height profile analysis images of transferred 3L WSe<sub>2</sub> and 1L MoS<sub>2</sub>/3L WSe<sub>2</sub> on the SiO<sub>2</sub> substrate. The thickness was 3 nm, which corresponds to four layers (4L) of 2D TMDs. [33] Figure 2c shows the Raman spectra for 1L MoS<sub>2</sub>/SiO<sub>2</sub>, 3L WSe<sub>2</sub>/SiO<sub>2</sub> and 1L MoS<sub>2</sub>/3L WSe<sub>2</sub>/SiO<sub>2</sub>. E<sup>1</sup><sub>2g</sub> and A<sub>1g</sub> Raman peaks for 1L MoS<sub>2</sub>/SiO<sub>2</sub> were observed at 383.5 and 403.6 cm<sup>-1</sup>, respectively, with the peak distance of 20.1 cm<sup>-1</sup> indicating 1L MoS<sub>2</sub>/SiO<sub>2</sub>. [23,32] The sum

of the  $E^1_{2g}$  and  $A_{1g}$  peaks at 249.8 cm $^{-1}$  was observed in the 3L WSe $_2$ /SiO $_2$ . The height ( $\approx$ 2.3 nm) of 3L WSe $_2$  was measured by AFM, which corresponded to the Raman data. [24] For the 1L MoS $_2$ /3L WSe $_2$ , a 3L WSe $_2$  peak was observed with peaks for MoS $_2$  ( $E^1_{2g}$  at 378.7 cm $^{-1}$  and  $A_{1g}$  at 403.5 cm $^{-1}$ ), indicating that MoS $_2$  was synthesized on the 3L WSe $_2$ . Also, the absence of any Raman peaks related to MoSe $_2$  ( $E^1_{2g}$  at 286 cm $^{-1}$  and  $A_{1g}$  at 244 cm $^{-1}$ ) or WS $_2$  ( $E^1_{2g}$  at 356 cm $^{-1}$  and  $A_{1g}$  at 420 cm $^{-1}$ ) indicates that there was no significant mixing or alloying phenomena in the MoS $_2$ /WSe $_2$  heterostructure during the synthesis. [33,34] It should be noted that there was a noticeable 5 cm $^{-1}$  downshift in



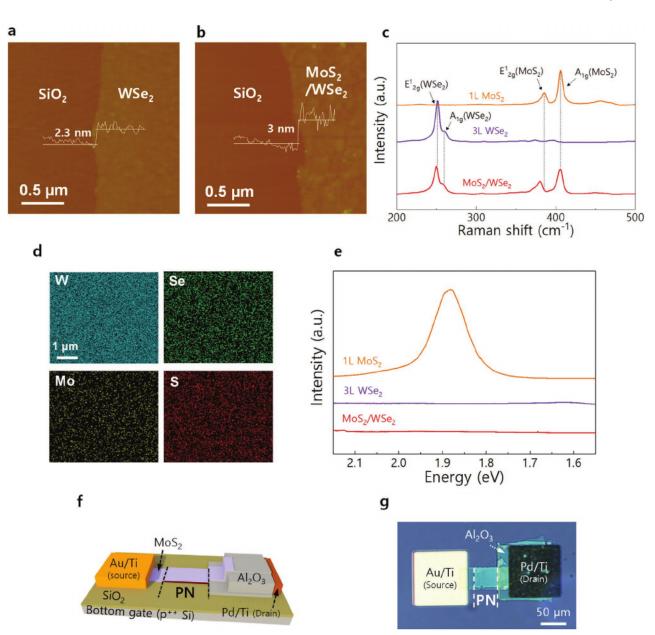


Figure 2. AFM image and height profile of a)  $WSe_2$  and b)  $MoS_2/WSe_2$  c) Raman spectra of 1L  $MoS_2$ , 3L  $WSe_2$  and 1L  $MoS_2/3L$   $WSe_2$  heterostructure. d) EDS mapping of  $MoS_2/WSe_2$  heterostructure. e) PL spectra of 1L  $MoS_2$ , 3L  $WSe_2$ , and 1L  $MoS_2/3L$   $WSe_2$  heterostructure. f) Schematic illustration and g) OM image of fabricated  $MoS_2/WSe_2$  gas sensor.

the  $\rm E^1_{2g}$  peak of MoS<sub>2</sub> in the 1L MoS <sub>2</sub>/3L WSe<sub>2</sub> region relative to the 1L MoS<sub>2</sub>. As reported previously, this is presumably attributed to the interaction between MoS<sub>2</sub> and WSe<sub>2</sub>. [<sup>32,35]</sup> Raman mapping was employed to examine the uniformity of 1L MoS<sub>2</sub> on 3L WSe<sub>2</sub>. Figure S2a,b in the Supporting Information shows the Raman mapping on the peak intensity of WSe<sub>2</sub>  $\rm E^1_{2g}$  and the intensity ratio of WSe<sub>2</sub>  $\rm E^1_{2g}$  and MoS<sub>2</sub> A<sub>1g</sub>, respectively. The results indicate the uniform synthesis of 1L MoS<sub>2</sub>/3L WSe<sub>2</sub>. Furthermore, we measured the EDS mapping for the vertically stacked MoS<sub>2</sub>/WSe<sub>2</sub>, which reveals the uniformly distributed W, Se, Mo, and S elements of WSe<sub>2</sub> and MoS<sub>2</sub> (Figure 2d). Figure 2e presents the photoluminescence (PL) spectra of 1L MoS<sub>2</sub>, 3L WSe<sub>2</sub>, and 1L MoS<sub>2</sub>/3L WSe<sub>2</sub>. It was evident that the

strong PL peak for the direct gap transition of 1L MoS<sub>2</sub> was highly suppressed by the WSe<sub>2</sub> layer in MoS<sub>2</sub>/WSe<sub>2</sub>, which is attributable to the rapid separation of charge carriers.<sup>[17,32,35]</sup> The schematic illustration and the OM image of the fabricated MoS<sub>2</sub>/WSe<sub>2</sub> gas sensor are shown in Figure 2f,g, respectively.

The PN diode characteristics of the fabricated MoS  $_2$ /WSe  $_2$  heterostructure were determined because it is closely related to the performance of photovoltaic behavior. The PN diode is dependent on the back-gate voltage, which is adjusted by varying carrier concentration through electrical doping. The calculated forward/reverse current ratio at  $V_{\rm ds} = |5 \text{ V}|$  clearly shows the gate-tunable PN diode characteristics (**Figure 3**a). Additionally, the calculated forward/reverse current ratio ( $\approx 560 \text{ at } V_{\rm g} = 10 \text{ V}$ )

www.advancedsciencenews.com

www.afm-journal.de

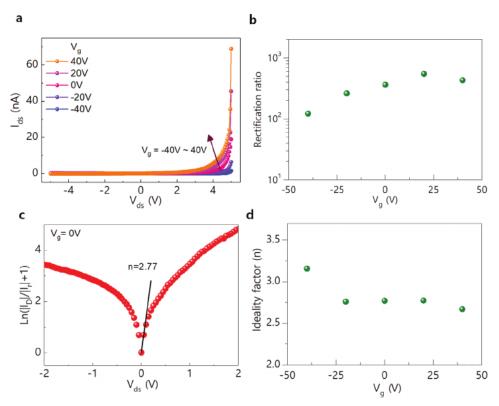


Figure 3. a) I-V characteristics of MoS<sub>2</sub>/WSe<sub>2</sub> PN diode with gate bias between -40 and 40 V. b) Rectification ratio of the MoS<sub>2</sub>/WSe<sub>2</sub> PN diode with gate bias. c) Ideality factor graph of MoS<sub>2</sub>/WSe<sub>2</sub> PN diode at  $V_g = 0$  V with  $V_{ds}$  between -2 and 2 V. d) Ideality factor depending on back-gate voltage.

was higher than that of a previously reported PN diode based on 1L MoS<sub>2</sub>/1L WSe<sub>2</sub> ( $\approx$ 50 at  $V_{\rm ds}$  = |8 V| and  $\approx$ 80 at  $V_{\rm ds}$  = |5 V|) (Figure 3b).[36] This gate-tunable characteristic of the PN diode could be explained by a variation in carrier density with electrical doping of the PN rectifying region. Initially, the majority carriers in MoS2 (electron) and WSe2 (hole) were depleted under a reverse bias.<sup>[36]</sup> Under a forward bias, the majority carriers were recombined at an interlayer that can be described by Shockley-Read-Hall and Langevin recombination, as well as the carriers were able to overcome the conduction and the valence band offset, thereby ensuring that a continuous current was maintained.[17,19,21,32,36] With increasing or decreasing gate voltage, p-type WSe2 was electrically doped to n-type or MoS2 was depleted, respectively. As a result, a continuous current was maintained regardless of the bias direction. In addition, to determine whether the diode was working ideally, we fitted the Shockley equation

$$|I_{\rm D}| = |I_{\rm rs}| \left( e^{\frac{eV_{\rm D}}{nk_{\rm B}T}} - 1 \right) \tag{1}$$

where  $I_{\rm D}$  is the drain current,  $V_{\rm D}$  is the drain bias,  $I_{\rm s}$  is the reverse saturation current, e is the electronic charge, n is the ideality factor,  $k_{\rm B}$  is the Boltzmann constant, and T is the operation temperature (300 K).<sup>[37]</sup> The ideal factor was 2.77 with a back-gate voltage of 0 V (Figure 3c), and the value was almost the same for back-gate voltages from -25 to 40 V (Figure 3d). The factor was similar to the PN diode heterostructure with monolayer MoS<sub>2</sub> and black phosphorus by exfoliation, which

we expected that the synthesized  $MoS_2/WSe_2$  heterostructure could produce the photovoltaic effect even if it is synthesized by an ex situ process. [17,38]

By using the above two types of 2D TMD heterostructures, we first checked the photovoltaic effect. Figure 4a,b shows schematics of two devices under white light illumination with an incident optical power density of 10 W m $^{-2}$ . The I-V curves of the WSe<sub>2</sub>/WS<sub>2</sub> (Figure 4c) and the MoS<sub>2</sub>/WSe<sub>2</sub> (Figure 4d) heterostructure under light illumination definitely indicate the photovoltaic characteristics. The short-circuit current density and the open-circuit voltage of the WSe2/WS2 heterostructure were 0.9 mA m<sup>-2</sup> and 0.74 mV, respectively. In case of the MoS<sub>2</sub>/WSe<sub>2</sub> heterostructure, the short-circuit current density and the open-circuit voltage were 0.45 mA m<sup>-2</sup> and 25 mV, respectively. Those current densities were lower than the MoS<sub>2</sub>/WSe<sub>2</sub> by exfoliation, where the differences between grain size, fabrication process and experimental conditions were assumed to be a reason.<sup>[17]</sup> In both devices, we found that the current increased with light illumination because of the generation of optically excited carriers. The open-circuit voltage directly indicated a photovoltaic effect. Without light illumination, both 2D TMD heterostructures showed no current at  $V_{ds} = 0$  V. When they were illuminated, electron-hole pairs in both MoS<sub>2</sub>/WSe<sub>2</sub> and WSe<sub>2</sub>/WS<sub>2</sub> layers were generated, as illustrated in Figure 4e. The schematic energy band diagram in Figure 4e depicts the separation processes of photoexcited charge carriers taking place within the WSe2/WS2 and the MoS<sub>2</sub>/WSe<sub>2</sub> heterostructures. Here, the electron-hole pairs in both the WSe2 and WS2 (MoS2) layers were photogenerated

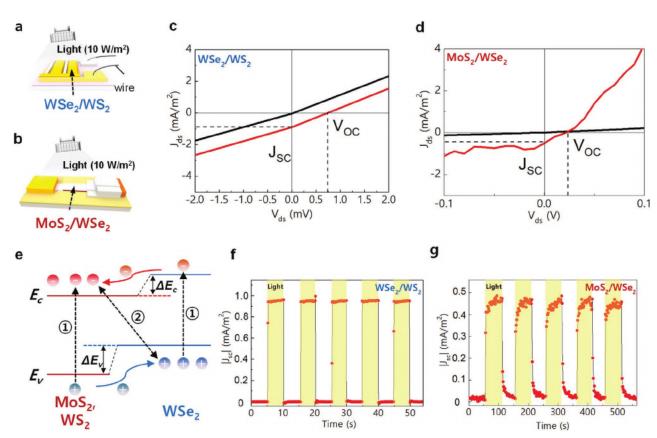


Figure 4. Schematic illustration of light illumination at fabricated a)  $WSe_2/WS_2$  and b)  $MoS_2/WSe_2$  gas sensor. *I–V* characteristics of c)  $WSe_2/WS_2$  and d)  $MoS_2/WSe_2$  gas sensor under each dark/light condition. e) Band energy diagram of 2D TMD heterostructures under light illumination. Photodynamics of f)  $WSe_2/WS_2$  and g)  $MoS_2/WSe_2$  gas sensor with 0 V bias.

because the electrons in the valence band were excited to the conduction band under light illumination. Then, the relaxation of the generated carriers occurred, which was driven by the band offsets  $\Delta E_{\rm C}$  and  $\Delta E_{\rm V}^{[36,39]}$  The electrons were diffused to the n-type WS<sub>2</sub> (MoS<sub>2</sub>) layers, while the holes were swept into the p-type WSe2 layers. This charge separation, through the aid of the built-in electric field across the junction, induced further electrical polarization of the heterostructure and eventually acted as a driving force for gas sensing for the WSe2/WS2 and MoS2/WSe2 heterostructures without the external bias voltage. In addition, the photodynamics of the WSe2/WS2 and MoS2/WSe, heterostructures at 0 V bias are presented in Figure 4f,g, respectively. The photocurrent switching under light illumination was operated for five repeats. From those experiments, we confirmed the photovoltaic characteristics of the fabricated 2D TMD heterostructure devices.

Based on the photovoltaic effect, the devices were tested to detect electron acceptor  $NO_2$  gas and electron donor  $NH_3$  gas, in the absence of external bias voltage. Accordingly, the photovoltaic self-powered gas sensing experiments were conducted on  $WSe_2/WS_2$  under light illumination with exposure to  $NO_2$  and  $NH_3$  gases, as illustrated in **Figure 5a**. Upon the exposure to both  $NO_2$  and  $NH_3$ , the changes in the photocurrent generated by each gas were observed respectively, as shown in Figure 5b,c. The gas sensing response (%) was calculated using the following equation  $^{[13,16]}$ 

Adv. Funct. Mater. 2020, 30, 2003360

Response (%) = 
$$\frac{I_g - I_0}{I_0} \times 100\% = \frac{\Delta I}{I_0} \times 100\%$$
 (2)

where  $I_0$  is the base current, and  $I_g$  refers to the corresponding current with the target gas exposure. Figure 5b shows the changes in photocurrent generated by the WSe $_2$ /WS $_2$  gas sensor upon the NO $_2$  exposure (concentrations of 10, 20, 50, 100, 200, and 500 ppm). The NO $_2$  exposure to the WSe $_2$ /WS $_2$  gas sensor was maintained for 10 min (blue color region in Figure 5b), followed by the inert N $_2$  gas for desorbing the NO $_2$  gas (white region in Figure 5b). The positive sensing response by NO $_2$  gas was observed with a gradual increase in response from 178% (10 ppm) to 322% (500 ppm) And, it showed good recovery at the room temperature compared to previous reports. [10,40]

In contrast, an opposite phenomenon was observed when exposed to NH<sub>3</sub>. Figure 5c shows the changes in photocurrent of the WSe<sub>2</sub>/WS<sub>2</sub> gas sensor upon the exposure to NH<sub>3</sub> (concentrations of 10, 20, 50, 100, 200, and 500 ppm). The exposure to NH<sub>3</sub> was also maintained for 10 min, followed by the N<sub>2</sub> purging. Then, the photocurrent level decreased with the NH<sub>3</sub> exposure. The increase in concentration of NH<sub>3</sub> gas up to 500 ppm led to the negative sensing response, gradually decreasing from -19% (10 ppm) to -23% (500 ppm) (Figure 5c). This indicates that the WSe<sub>2</sub>/WS<sub>2</sub> gas sensor is useful for the NO<sub>2</sub> gas, compared to the NH<sub>3</sub> gas.

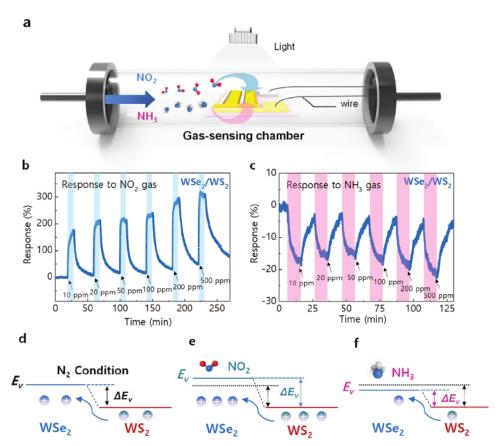


Figure 5. a) Schematic illustration of  $WSe_2/WS_2$  gas sensor for  $NO_2$  and  $NH_3$ . Consecutive measurement of b)  $NO_2$  and c)  $NH_3$  gas sensing. Charge transfer process at d)  $N_2$  condition, and band bending under e)  $NO_2$  and f)  $NH_3$  gas exposure at  $WSe_2/WS_2$  gas sensor.

The gas sensing characteristics of 2D TMD were based on the charge-transfer mechanism induced by the physisorption of gas molecules mostly over pristine surface areas of the 2D TMDs.[41,42] When the NO2 gas was injected, the adsorption of NO2 molecules acted as an electron acceptor and extracted electrons from the adjacent WSe2 layer, thereby increasing hole concentrations. From that, the valence band of the WSe2 layer with the adsorption of NO2 molecules increased compared to the N2 condition, as depicted in Figure 5d,e (dotted line: before gas adsorption; solid line: after gas adsorption). This eventually led to an increase in photocurrent compared with the case of the N<sub>2</sub> condition, as shown in Figure 5b. On the other hand, the adsorption of NH3 molecules on the WSe2 layer induced a decrease in hole concentrations due to the role as electron donors (Figure 5f).[16,43] Therefore, we observed a decrease in photocurrent in the WSe<sub>2</sub>/WS<sub>2</sub> gas sensor as in Figure 5c.

We further measured the  $V_{\rm oc}$  changes for each gas sensing (Figure S4, Supporting Information) and found a correspondence with the following ideal p–n heterostructure solar diode equation<sup>[44]</sup>

$$V_{\rm OC} = \frac{KT}{q} \ln \left( \frac{I_{\rm ph}}{I_{\rm rev,sat}} + 1 \right) \tag{3}$$

where K is the Boltzmann constant, q is charge,  $I_{\rm ph}$  is photocurrent generated by light illumination, and  $I_{\rm rev,sat}$  is reverse saturation current. Here, because  $I_{\rm rev,sat}$  can be considered as a

negligible factor with a constant value, the increasing hole concentrations from the  $\mathrm{NO_2}$  exposure induced the increment of  $V_{\mathrm{oc}}$  (Figure S5c, Supporting Information). On the other hand, the decreasing hole concentrations due to the NH<sub>3</sub> exposure led to a reduction in  $V_{\mathrm{oc}}$  (Figure S5d, Supporting Information). From these measurements, we were able to understand the changes in the photovoltaic characteristics in terms of the charge-transfer mechanism between the injected gas molecules and the surface of 2D TMD heterostructure under light illumination.

Similarly, we conducted the gas sensing experiments on the  $MoS_2/WSe_2$  for  $NO_2$  and  $NH_3$  gas, as shown in **Figure 6a**. Figure 6b presents the calculated gas sensing response of the  $MoS_2/WSe_2$  gas sensor upon the  $NO_2$  exposure (concentrations of 50, 100, 200, and 500 ppm). The  $MoS_2/WSe_2$  gas sensor was exposed to the  $NO_2$  for 5 min, and then it was purged in  $N_2$  gas for 5 min for recovery. In  $MoS_2/WSe_2$  gas sensor, we shortened the response time to 5 min because it is sufficient for  $NO_2$  gas response saturation. In this case,  $NO_2$  gas mainly adsorbed on the n-type  $MoS_2$  surface and acted as an electron acceptor (Figure 6e), and the responsivity range was  $\approx$ –15% (50 ppm) to –36% (500 ppm).

Figure 6c shows the current change of the  $MoS_2/WSe_2$  gas sensor upon the NH<sub>3</sub> exposure (concentrations of 10, 20, 50, 200, and 500 ppm). The  $MoS_2/WSe_2$  was exposed to NH<sub>3</sub> gas for 5 min and then purged with N<sub>2</sub> for 5 min, which are the same time for NO<sub>2</sub> gas. As NH<sub>3</sub> molecules acted as electron

www.advancedsciencenews.com

www.afm-journal.de

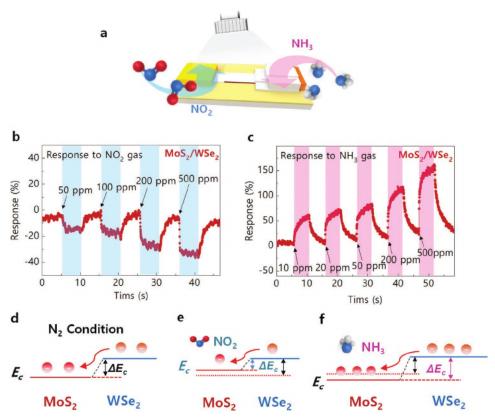


Figure 6. a) Schematic illustration of  $MoS_2/WSe_2$  gas sensor for  $NO_2$  and  $NH_3$ . Consecutive measurement of b)  $NO_2$  and c)  $NH_3$  gas sensing. Charge transfer process at d)  $N_2$  condition, and band bending under e)  $NO_2$  and f)  $NH_3$  gas exposure at  $MoS_2/WSe_2$  gas sensor.

donors, the electron concentration of  $MoS_2$  increased, as shown in Figure 6f.<sup>[14,42]</sup> Therefore, the electron concentration increased in the  $MoS_2$  layer, rising the photocurrent. Then, the photocurrent response of the  $MoS_2/WSe_2$  gas sensor resulted in a positive sensing response, increased from 62% (10 ppm) to 165% (500 ppm), which is much higher than that for  $NO_2$  gas. This tendency was entirely opposite to the case of the  $WSe_2/WS_2$  gas sensor.

The responsivities of the NO<sub>2</sub> and NH<sub>3</sub> gases for each gas sensor are presented in **Figure 7**a,b, respectively. Based on this comparative analysis of gas sensing responsivity between

WSe<sub>2</sub>/WS<sub>2</sub> and MoS<sub>2</sub>/WSe<sub>2</sub> heterostructures, we found that the exposed surface plays the most dominant role in transferring the charges from the gas, even at monolayer thickness.

To identify and interpret the observed difference in responsivity of surface material with gas molecules, we performed first-principles calculations based on the adsorption of gas molecules on 2D TMDs (**Figure 8**). According to the calculation results, NH<sub>3</sub> gives electrons to the 2D TMD layers, while NO<sub>2</sub> receives electrons. In other words, NH<sub>3</sub> acts like an n-type dopant and NO<sub>2</sub> acts like a p-type dopant, as shown in Figure 8, which agree with our experimental results. For a

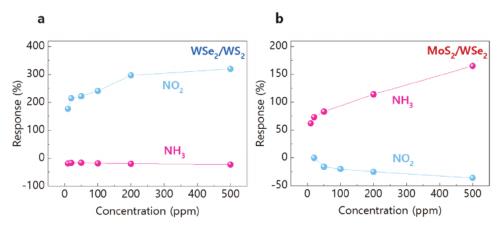


Figure 7. Calculated response of a) WSe<sub>2</sub>/WS<sub>2</sub> and b) MoS<sub>2</sub>/WSe<sub>2</sub> gas sensor corresponding to each gas concentration of NO<sub>2</sub> and NH<sub>3</sub>.

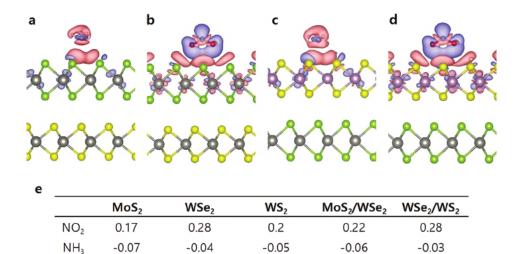


Figure 8. The difference in the charge densities of a)  $NH_3$  b)  $NO_2$  adsorbed on  $WSe_2/WS_2$  heterostructure, and c)  $NH_3$  d)  $NO_2$  adsorbed on  $MOS_2/WSe_2$  heterostructure, respectively, obtained from DFT calculations. Gray, yellow, green, and purple dots indicate the W, S, Se, and Mo atoms, respectively. Blue and red colors indicate electron accumulation and depletion. e) The calculated transferred charges between the gas molecules and 2D TMD. The unit of the number is the elementary charge, e, and the positive and negative signs indicate the accumulation and depletion of electrons.

single TMD materials in our calculation performed, NO2 had a larger amount of charge transfer than NH3, which are similar to previous reports.[42,43] Further, the adsorption energy of NH<sub>3</sub> was smaller than the NO<sub>2</sub> on WSe<sub>2</sub>.[43] In addition, even if the number of TMD layer increased, there was little difference in adsorption energy. <sup>[27]</sup> Therefore, in the 3L WSe<sub>2</sub>/3L WS<sub>2</sub> gas sensor, the lower responsivity of  $\text{NH}_3$  than that of NO  $_2$  gas is presumably due to the relatively lower adsorption energy and smaller charge transfer ability of NH3 gas with WSe2 surface.[15,16,43] On the other hand, it is interesting to note that the lower responsivity to NO<sub>2</sub> than NH<sub>3</sub> in MoS<sub>2</sub>/WS<sub>2</sub> gas sensor. [14] In case of MoS2, the adsorption energies of NH3 and NO2 were similar.[42,45] Therefore, the amount of charge transfer between the gas and MoS2/WSe2 was crucial for gas sensing responsivity. From the computational simulation results, NO2 gas had a larger charge transfer value in MoS<sub>2</sub>/WSe<sub>2</sub> than in MoS<sub>2</sub>. It means that there exists the charge transfer between the NO2 and the WSe2 layer. [46] As WSe2 and MoS2 behaves in an opposite manner to the NO<sub>2</sub> gas, it was considered that the response of MoS<sub>2</sub>/WSe<sub>2</sub> to NO<sub>2</sub> gas was low. In contrast, the degree of charge transfer between the NH3 and WSe2 was smaller compared to the case of NO<sub>2</sub>. Hence, the MoS<sub>2</sub>/WSe<sub>2</sub> gas sensor is more appropriate for detecting the NH<sub>3</sub> gas than NO<sub>2</sub> gas.

In addition, the response and recovery time of the NO<sub>2</sub> sensing (10 ppm) by WSe<sub>2</sub>/WS<sub>2</sub> (Figure S6a, Supporting Information) and the NH<sub>3</sub> sensing (10 ppm) by MoS<sub>2</sub>/WSe<sub>2</sub> (Figure S6b, Supporting Information) were measured. The response time was measured from the 10% response point to 90% of an end point (maximum), and the recovery time was measured at the point where the current hit the one-tenth point near the minimum from the one-ninth point near the maximum. Consequently, the response and recovery times for the NO<sub>2</sub> detection by the WSe<sub>2</sub>/WS<sub>2</sub> gas sensor were determined to be about 387 s ( $\approx$ 6 min) and 728 s ( $\approx$ 12 min), respectively. In the same manner, the time period for one cycle of NH<sub>3</sub> sensing (10 ppm) for MoS<sub>2</sub>/WSe<sub>2</sub> was characterized. For the NH<sub>3</sub> detection by the MoS<sub>2</sub>/WSe<sub>2</sub> gas sensor, it showed a response time

of  $\approx$ 205 s ( $\approx$ 3 min) and a recovery time of  $\approx$ 263 s ( $\approx$ 4 min). The ratio of recovery time to response time by our devices was relatively fast, compared to the previous reports. [14–16,47,48]

#### 3. Conclusion

In summary, we successfully demonstrated a photovoltaic selfpowered gas sensor using both WSe2/WS2 and MoS2/WSe2 heterostructures for the first time. The 2D TMD heterostructures effectively induced separation of photoexcited charge carriers with the built-in electric field under light illumination, which acted as a driving force for gas sensing. The fabricated 2D TMD heterostructure gas sensors showed strong abilities to detect NO2 and NH3 gases with high responsivity and selectivity without external bias voltage. The gas sensing mechanism was further elucidated through establishing the correlations of charge transfer with the type of adsorbed gas molecules onto the exposed surface of 2D TMD heterostructure gas sensor. We confirmed that the photovoltaic features of 2D TMD heterostructures could be a viable platform for highly sensitive gas sensor applications. Further, research related to practical gas sensing circumstances such as selective sensing of mixed gas will contribute to a significant advancement in environmental monitoring systems.

# 4. Experimental Section

Fabrication of WSe<sub>2</sub>/WS<sub>2</sub> Heterostructure Gas Sensor. The experimental procedure for the in situ synthesis of the WSe<sub>2</sub>/WS<sub>2</sub> heterostructure is illustrated in Figure S1a in the Supporting Information. Tube-type chamber equipment was used for in situ synthesis of the WSe<sub>2</sub>/WS<sub>2</sub> heterostructure, as depicted in the schematic image. The SiO<sub>2</sub> (285 nm) Si substrate was prepared through typical cleaning with acetone, IPA, and deionized water. The tube-type chamber was heated to 700 °C for 30 min. A WS<sub>2</sub> film was deposited on the SiO<sub>2</sub>/Si substrate through CVD using tungsten hexachloride (WCl<sub>6</sub>) and hydrogen sulfide (H<sub>2</sub>S)



ADVANCED FUNCTIONAL MATERIALS

gas. Then, the WSe $_2$  film was synthesized on a WS $_2$  continuously with WCl $_6$  and diethyl selenide (DESe) in the same chamber. During both the WS $_2$  and WSe $_2$  synthesis steps, the WCl $_6$  flowed into the chamber with the aid of Ar carrier gas at 50 sccm. Hydrogen (H $_2$ ) gas was inserted at 20 sccm during both synthetic processes to prevent unexpected carbon contamination.

The vertically stacked  $WSe_2/WS_2$  heterostructure was fabricated with bottom (Au) and top (Pd) electrodes. After pattering the bottom electrode on a bare  $SiO_2/Si$  substrate, the synthesized  $WSe_2/WS_2$  heterostructure was transferred onto the Au/Ti bottom electrode. Finally, Pd (40 nm) was deposited as a top contact. The overall pattern design of the electrodes was formed as an interdigitated electrode.

Fabrication of MoS<sub>2</sub>/WSe<sub>2</sub> Heterostructure Gas Sensor: The experimental procedure for the MoS<sub>2</sub>/WSe<sub>2</sub> heterostructure as a gas sensor is illustrated in Figure S2c in the Supporting Information. The overall pattern design of the MoS<sub>2</sub>/WSe<sub>2</sub>heterostructures and electrodes was formed by conventional lithography with large-area synthesis of 2D TMDs without a transfer method. The SiO<sub>2</sub> (285 nm)/Si substrate was prepatterned by a Pd/Ti electrode (Ti as an adhesion layer) through an evaporation and lift-off method. Then, the WSe2 was synthesized by a previously reported self-limiting layer (SLS) synthesis method.<sup>[24]</sup> The tube-type chamber was used to achieve a growth temperature of 700 °C. Tungsten hexachloride (WCl6) and DESe were used as a precursor and reactant, respectively. The SLS consisted of four steps, each using the same procedure of atomic layer deposition (ALD). A precursor-purgereactant-purge sequence was repeated for 150 cycles for 3L  $WSe_2$ synthesis. A canister containing the precursor was heated to 90 °C to ensure an adequate vapor pressure for the precursor molecules to be carried into the tube by a pure Ar carrier gas (30 sccm). DESe had sufficient vapor at a typical room temperature (25 °C), so there was no carrier gas injected. In addition, H2 (10 sccm) gas was injected with DESe during the reactant time. The synthesized WSe2 was patterned by reactive-ion etching with CHF3 and O2 gas for etching gas. Before synthesizing  $MoS_2$  on the  $WSe_2$ ,  $Al_2O_3$  was deposited by Al (2 nm) evaporation followed by oxidation. Al<sub>2</sub>O<sub>3</sub> was patterned by the HNO<sub>3</sub>.

The  $MoS_2$  was then synthesized by a catalytic chemical vapor deposition method. [23]  $MoS_2$  was synthesized on prepatterned  $WSe_2$  on the  $SiO_2$  (285 nm)/Si substrate. The tube-type chamber was used to achieve a growth temperature of 700 °C.  $MoCl_5$  and dimethyl sulfide (DMS) were used as a precursor and reactant, respectively. The temperature of the canister containing the precursor was heated to 90 °C to achieve adequate vapor pressure, and the vaporized precursor molecules were carried into the tube using Ar carrier gas (3 sccm). The DMS and  $H_2$  flow were fixed at 20 and 10 sccm, respectively. 10 mg of NaCl was added to the hot zone of the CVD reactor, which had a temperature of 700 °C.  $MoS_2$  was etched by reactive-ion etching with CHF<sub>3</sub> and  $O_2$  gas for channel patterning, and Au/Ti was deposited as an electrode to the  $MoS_2$  by evaporation.

Characterization of 2D TMD Heterostructures: OM images were taken using an Olympus DX51, and Raman spectroscopy was performed with a LabRam ARAMIS (HORIBA, 532 nm wavelength laser). AFM was conducted using a multimode scanning microscope (Veeco). Tecnai G2 F20 S-TWIN (accelerating voltage, 80 kV) was employed to characterize the TEM images of the vertically stacked WSe2/WS2 heterostructure. A source meter (Keithley 2400) was used to measure the *I*–V curves, photovoltaic characteristics, and gas sensing properties of the fabricated WSe2/WS2 gas sensor. A voltage/current meter (Keithley 4200, Keithley Instruments) was used to measure the *I*–V curves, photovoltaic characteristics, and gas sensing properties of the fabricated MoS2/WSe2 gas sensor.

Gas Sensing Measurement: The whole sensing was implemented at a typical room temperature (25 °C) under white light illumination. The concentration of each gas was set by controlling the flow rate of the target gas and purging gas ( $N_2$ ) using mass-flow controllers. Basically, the gas sensors were stabilized under  $N_2$  exposure for  $\approx 1$  h to blow off moisture, dust, and other particles that could influence the experimental outcome. No external bias was applied to the sensors, and 10 W m<sup>-2</sup> white light illuminated the sensor with the remainder of the chamber

Adv. Funct. Mater. 2020, 30, 2003360

surface covered to block natural light. The real-time output signal (current) was acquired using the source meter mentioned above.

Computational Details: In order to estimate a doping effect on the 2D TMD heterostructure by gas adsorption, the first-principle calculations were performed using density functional theory (DFT) implemented in the Vienna Ab-initio Simulation Package (VASP) with a projector augmented wave (PAW) method. [49] As the exchange–correlation energy functional, the local density approximation (LDA) within the Ceperley–Alder scheme was considered [50] 2  $\times$  2 primitive cells of 2D TMD were considered for a unit cell with one gas molecule to theoretical calculations. The kinetic energy cutoff was set to 400 eV. Geometrical optimization for the systems was performed until the Hellmann–Feynman force acting on each atom was smaller than 0.01 eV Å $^{-1}$ . The first Brillouin zone integration was carried out using 9  $\times$  9  $\times$  1 gamma centered grids. The transferred charges between the gas molecules and 2D TMD sheets were calculated using Bader analysis. [42]

# **Supporting Information**

Supporting Information is available from the Wiley Online Library or from the author.

## Acknowledgements

Y. Kim and S. Lee equally contributed to this work. This work was supported by the Materials and Components Technology Development Program of MOTIE/KEIT. [10080527, Development of commercialization technology of highly sensitive gas sensor based on chalcogenide 2D nanomaterial]; Samsung Display CO., LTD.; and by the Basic Science Research Program through the National Research Foundation of Korea (NRF) funded by the Ministry of Education, Science and Technology (NRF-2018R1D1A1A09084143 and NRF-2017R1C1B5076821), and by the National Research Foundation of Korea (NRF) grant funded by the Korea government(MSIT) (No. NRF-2018R1A2B6005289). This research was also supported by the MOTIE(Ministry of Trade, Industry & Energy, No. 20006504 and No. 20007000) and KSRC(Korea Semiconductor Research Consortium) support program for the development of the future semiconductor device.

#### **Conflict of Interest**

The authors declare no conflict of interest.

### Keywords

gas sensors, heterostructures, photovoltaics, self-powered sensors, transition metal dichalcogenide

Received: April 16, 2020 Revised: June 25, 2020 Published online: September 9, 2020

A. Das, R. Dost, T. Richardson, M. Grell, J. J. Morrison, M. L. Turner, Adv. Mater. 2007, 19, 4018.

<sup>[2]</sup> S. Ji, H. Wang, T. Wang, D. Yan, Adv. Mater. 2013, 25, 1755.

<sup>[3]</sup> J. F. Fernández-Sánchez, I. Fernández, R. Steiger, R. Beer, R. Cannas, U. E. Spichiger-Keller, Adv. Funct. Mater. 2007, 17, 1188.

<sup>[4]</sup> W. Tsujita, A. Yoshino, H. Ishida, T. Moriizumi, Sens. Actuators, B 2005, 110, 304.



·

- [5] S. V. Ryabtsev, A. V. Shaposhnick, A. N. Lukin, E. P. Domashevskaya, Sens. Actuators, B 1999, 59, 26.
- [6] Y. Mo, Y. Okawa, K. Inoue, K. Natukawa, Sens. Actuators, A 2002, 100, 94
- [7] K. Y. Choi, J. S. Park, K. B. Park, H. J. Kim, H. D. Park, S. D. Kim, Sens. Actuators, B 2010, 150, 65.
- [8] M. W. G. Hoffmann, L. Mayrhofer Dr, O. Casals Dr, L. Caccamo, F. Hernez-Ramirez, G. Lilienkamp, W. Daum, M. Moseler, A. Waag, H. Shen, J. D. Prades, Adv. Mater. 2014, 26, 8017.
- [9] M. W. G. Hoffmann, A. E. Gad, J. D. Prades, F. Hernandez-Ramirez, R. Fiz, H. Shen, S. Mathur, *Nano Energy* 2013, 2, 514.
- [10] H. Li, Z. Yin, Q. He, H. Li, X. Huang, G. Lu, D. W. H. Fam, A. I. Y. Tok, Q. Zhang, H. Zhang, Small 2012, 8, 63.
- [11] Q. He, Z. Zeng, Z. Yin, H. Li, S. Wu, X. Huang, H. Zhang, Small 2012, 8, 2994.
- [12] S. J. Choi, I. D. Kim, Electron. Mater. Lett. 2018, 14, 221
- [13] K. Y. Ko, J. G. Song, Y. Kim, T. Choi, S. Shin, C. W. Lee, K. Lee, J. Koo, H. Lee, J. Kim, T. Lee, J. Park, H. Kim, ACS Nano 2016, 10, 9287.
- [14] Y. Zhao, J. G. Song, G. H. Ryu, K. Y. Ko, W. J. Woo, Y. Kim, D. Kim, J. H. Lim, S. Lee, Z. Lee, J. Park, H. Kim, *Nanoscale* 2018, 10, 9338.
- [15] K. Y. Ko, S. Lee, K. Park, Y. Kim, W. J. Woo, D. Kim, J. G. Song, J. Park, J. H. Kim, Z. Lee, H. Kim, ACS Appl. Mater. Interfaces 2018, 10, 34163.
- [16] K. Y. Ko, K. Park, S. Lee, Y. Kim, W. J. Woo, D. Kim, J. G. Song, J. Park, H. Kim, ACS Appl. Mater. Interfaces 2018, 10, 23910.
- [17] C. H. Lee, G. H. Lee, A. M. Van Der Zande, W. Chen, Y. Li, M. Han, X. Cui, G. Arefe, C. Nuckolls, T. F. Heinz, J. Guo, J. Hone, P. Kim, Nat. Nanotechnol. 2014, 9, 676.
- [18] R. Cheng, D. Li, H. Zhou, C. Wang, A. Yin, S. Jiang, Y. Liu, Y. Chen, Y. Huang, X. Duan, *Nano Lett.* **2014**, *14*, 5590.
- [19] R. Frisenda, A. J. Molina-Mendoza, T. Mueller, A. Castellanos-Gomez, H. S. J. Van Der Zant, Chem. Soc. Rev. 2018, 47, 3339.
- [20] M. Sun, Q. Fang, D. Xie, Y. Sun, J. Xu, C. Teng, R. Dai, P. Yang, Z. Li, W. Li, Y. Zhang, Adv. Electron. Mater. 2017, 3, 1600502.
- [21] T. Roy, M. Tosun, X. Cao, H. Fang, D. H. Lien, P. Zhao, Y. Z. Chen, Y. L. Chueh, J. Guo, A. Javey, ACS Nano 2015, 9, 2071.
- [22] Y. Xue, Y. Zhang, Y. Liu, H. Liu, J. Song, J. Sophia, J. Liu, Z. Xu, Q. Xu, Z. Wang, J. Zheng, Y. Liu, S. Li, Q. Bao, ACS Nano 2016, 10, 573.
- [23] J. G. Song, G. Hee Ryu, Y. Kim, W. Je Woo, K. Yong Ko, Y. Kim, C. Lee, I. K. Oh, J. Park, Z. Lee, H. Kim, Nanotechnology 2017, 28, 465103.
- [24] K. Park, Y. Kim, J. G. Song, S. J. Kim, C. Wanlee, G. H. Ryu, Z. Lee, J. Park, H. Kim, 2D Mater. 2016, 3, 014004.
- [25] J. Park, W. Lee, T. Choi, S. H. Hwang, J. M. Myoung, J. H. Jung, S. H. Kim, H. Kim, *Nanoscale* 2015, 7, 1308.
- [26] Z. Qin, D. Zeng, J. Zhang, C. Wu, Y. Wen, B. Shan, C. Xie, Appl. Surf. Sci. 2017, 414, 244.
- [27] D. J. Late, Y. K. Huang, B. Liu, J. Acharya, S. N. Shirodkar, J. Luo, A. Yan, D. Charles, U. V. Waghmare, V. P. Dravid, C. N. R. Rao, ACS Nano 2013, 7, 4879.

- [28] W. Zhao, Z. Ghorannevis, K. K. Amara, J. R. Pang, M. Toh, X. Zhang, C. Kloc, P. H. Tan, G. Eda, Nanoscale 2013, 5, 9677.
- [29] M. Zhang, J. Wu, Y. Zhu, D. O. Dumcenco, J. Hong, N. Mao, S. Deng, Y. Chen, Y. Yang, C. Jin, S. H. Chaki, Y. S. Huang, J. Zhang, L. Xie, ACS Nano 2014, 8, 7130.
- [30] Y. Chen, J. Xi, D. O. Dumcenco, Z. Liu, K. Suenaga, D. Wang, Z. Shuai, Y. S. Huang, L. Xie, ACS Nano 2013, 7, 4610.
- [31] X. Duan, C. Wang, Z. Fan, G. Hao, L. Kou, U. Halim, H. Li, X. Wu, Y. Wang, J. Jiang, A. Pan, Y. Huang, R. Yu, X. Duan, *Nano Lett.* 2016, 16, 264.
- [32] Y. Kim, J. G. Song, Y. J. Park, G. H. Ryu, S. J. Lee, J. S. Kim, P. J. Jeon, C. W. Lee, W. J. Woo, T. Choi, H. Jung, H. B. R. Lee, J. M. Myoung, S. Im, Z. Lee, J. H. Ahn, J. Park, H. Kim, Sci. Rep. 2016, 6, 18754.
- [33] J. G. Song, J. Park, W. Lee, T. Choi, H. Jung, C. W. Lee, S. H. Hwang, J. M. Myoung, J. H. Jung, S. H. Kim, C. Lansalot-Matras, H. Kim, ACS Nano 2013, 7, 11333.
- [34] S. V. Bhatt, M. P. Deshpande, V. Sathe, R. Rao, S. H. Chaki, J. Raman Spectrosc. 2014, 45, 971.
- [35] M. H. Chiu, M. Y. Li, W. Zhang, W. T. Hsu, W. H. Chang, M. Terrones, H. Terrones, L. J. Li, ACS Nano 2014, 8, 9649.
- [36] M. M. Furchi, A. Pospischil, F. Libisch, J. Burgdörfer, T. Mueller, Nano Lett. 2014, 14, 4785.
- [37] D. Jariwala, V. K. Sangwan, C.-C. Wu, P. L. Prabhumirashi, M. L. Geier, T. J. Marks, L. J. Lauhon, M. C. Hersam, *Proc. Natl. Acad. Sci. USA* 2013, 110, 18076.
- [38] Y. Deng, Z. Luo, N. J. Conrad, H. Liu, Y. Gong, S. Najmaei, P. M. Ajayan, J. Lou, X. Xu, P. D. Ye, ACS Nano 2014, 8, 8292.
- [39] N. Huo, J. Yang, L. Huang, Z. Wei, S. S. Li, S. H. Wei, J. Li, Small 2015, 11, 5430.
- [40] S. Y. Cho, S. J. Kim, Y. Lee, J. S. Kim, W. Bin Jung, H. W. Yoo, J. Kim, H. T. Jung, ACS Nano 2015, 9, 9314.
- [41] C. Zhou, W. Yang, H. Zhu, J. Chem. Phys. 2015, 142, 214704.
- [42] B. Cho, M. G. Hahm, M. Choi, J. Yoon, A. R. Kim, Y. J. Lee, S. G. Park, J. D. Kwon, C. S. Kim, M. Song, Y. Jeong, K. S. Nam, S. Lee, T. J. Yoo, C. G. Kang, B. H. Lee, H. C. Ko, P. M. Ajayan, D. H. Kim, Sci. Rep. 2015, 5, 8052.
- [43] T. Wang, R. Zhao, X. Zhao, Y. An, X. Dai, C. Xia, RSC Adv. 2016, 6, 82793.
- [44] L. Liu, G. H. Li, Y. Wang, Y. Y. Wang, T. Li, T. Zhang, S. J. Qin, Nanoscale 2017, 9, 18579.
- [45] Q. Yue, Z. Shao, S. Chang, L. Jingbo, Nanoscale Res. Lett. 2013, 8, 425.
- [46] G. Liu, S. L. Rumyantsev, C. Jiang, M. S. Shur, A. A. Balandin, IEEE Electron Device Lett. 2015, 36, 1202.
- [47] B. Liu, L. Chen, G. Liu, A. N. Abbas, M. Fathi, C. Zhou, ACS Nano 2014, 8, 5304.
- [48] R. Kumar, N. Goel, M. Kumar, ACS Sens. 2017, 2, 1744.
- [49] D. Joubert, Phys. Rev. B: Condens. Matter Mater. Phys. 1999, 59, 1758.
- [50] D. M. Ceperley, B. J. Alder, Phys. Rev. Lett. 1980, 45, 566.