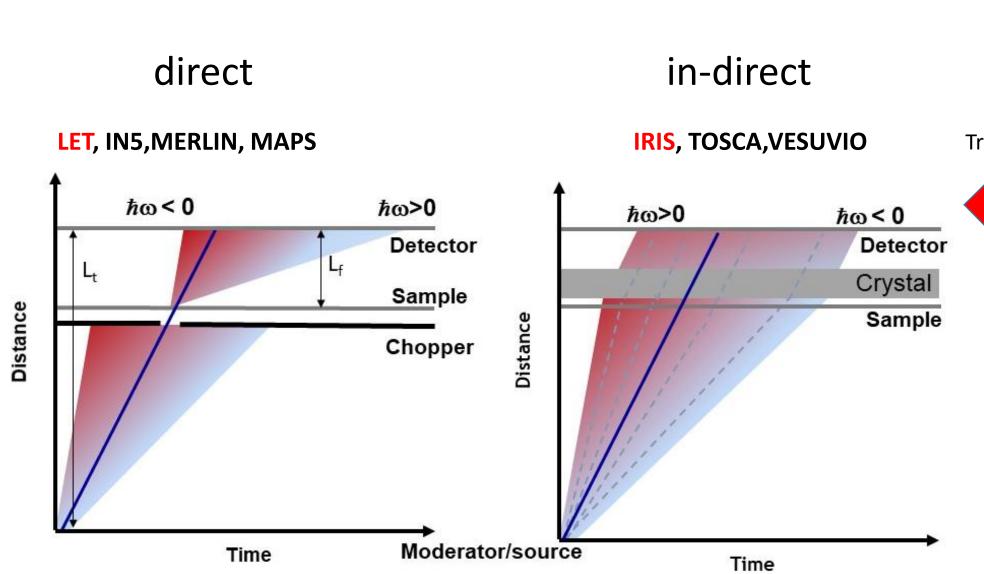
Spectrometer optimisation

McStas Workshop 13th-15th April 2021

Rob Bewley

What will we be focussing on **tof**





Where do you start designing your instrument?

- You don't start with McStas !!!
- Need the key drivers (resolution, sample environment, any limitations on design)
- Parameter space is large. To optimise its crucial to understand the 'Maths' of your instrument
 -how the resolution and flux depend on the instrument parameters
- What do we mean by an optimal design? need some sort of FOM

What do we mean by an optimal design?

Need a Figure Of Merit

$$FOM = I\sigma^2/w^2$$

For spectrometers FOM = Signal/Background

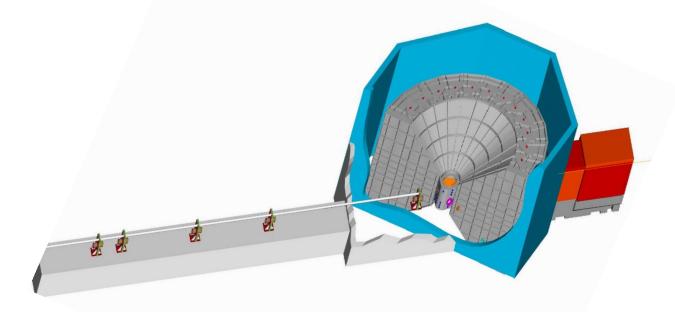
Optimal design will maximise the FOM for your key drivers

Background Cannot really calculate or simulate. Minimise with shielding and design

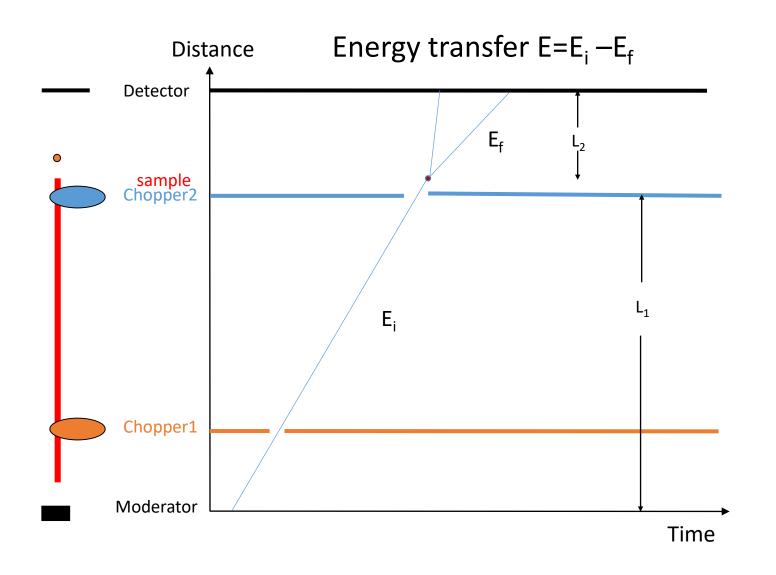
Rest of this talk is about the understanding the Maths of the instrument and how it relates to resolution and signal

'The Maths'

$$\Delta E^{2} = \left[\left(\frac{2E_{f}\Delta t}{t_{f}} \right)^{2} + \left(\frac{2E_{f}\Delta t_{ch} \left(L_{2} + L_{3} + L_{1} \right)}{t_{f}L_{1}} \right)^{2} + \left(\frac{2E_{f}\Delta t m_{od} L_{2}}{t_{f}L_{1}} \right)^{2} \right]$$



 $E=E_i-E_f$ where E_i fixed



$$\frac{\Delta E}{\Delta t} = -\frac{\Delta E_f}{\Delta t}$$

Where $E_f = 1/2 m_n (L_2/t_f)^2$

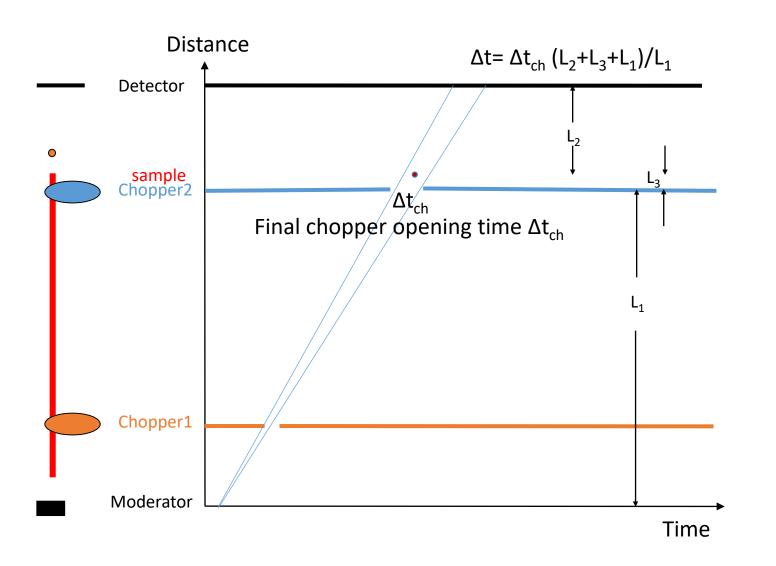
gives
$$\Delta E = \frac{2E_f \Delta t}{t_f}$$

Where

 $\begin{array}{lll} \Delta E & \text{energy uncertainty} \\ \Delta t & \text{time uncertainty at detector} \\ E_f & \text{energy of neutron after sample} \\ t_f & \text{flight time sample to detector} \end{array}$

Sources of Δt

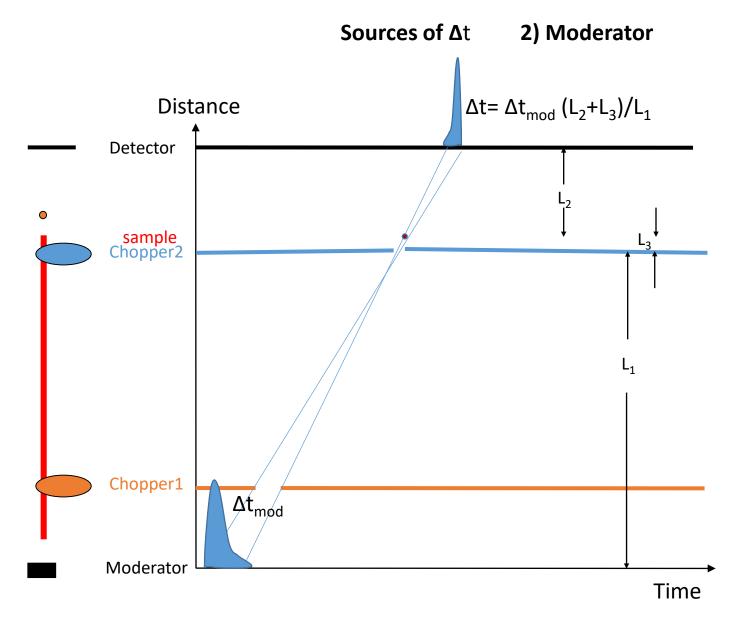
1) Chopper opening



$$\Delta \mathsf{E} = \frac{2E_f \Delta t}{t_f}$$

Chopper term

$$\Delta \mathsf{E} = \frac{2E_f \Delta t c_h \left(L_2 + L_3 + L_1\right)}{t_f L_1}$$



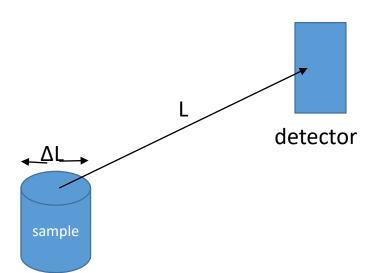
$$\Delta \mathsf{E} = \frac{2E_f \Delta t}{t_f}$$

Moderator term

$$\Delta \mathsf{E} = \frac{2E_f \Delta t m_{od} \left(L_2 + L_3 \right)}{t_f L_1}$$

Sources of Δt

3) Sample/detector size



$$\Delta \mathsf{E} = \frac{2E_f \Delta t}{t_f}$$

$$\frac{\Delta t}{t_f} = \frac{\Delta L}{L}$$

On LET $\Delta L = 2$ cm and L = 350 cm (sample to detector distance)

$$\frac{\Delta E}{E_f} = \frac{2x2}{350} \approx 1\%$$

Sample size only significant when energy resolution is close to 1%

Assuming no correlations all the terms convolute together. A good approximation is to add in quadrature

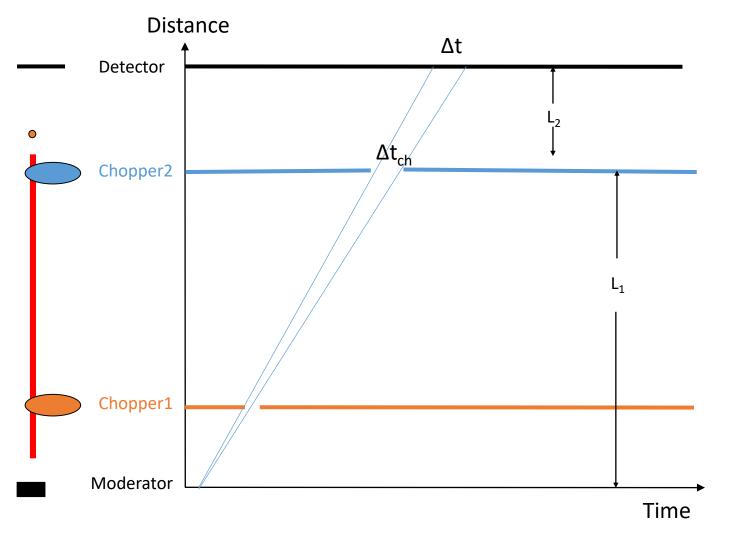
$$\Delta E^2 = \Delta E^2_{sam} + \Delta E^2_{chop} + \Delta E^2_{mod}$$

$$\Delta \mathsf{E}^2 = \left[\left(\frac{2E_f \Delta t_{sam}}{t_f} \right)^2 + \left(\frac{2E_f \Delta t_{chop} \left(L_2 + L_3 + L_1 \right)}{t_f L_1} \right)^2 + \left(\frac{2E_f \Delta t m_{od} L_2}{t_f L_1} \right)^2 \right]$$

What does it all mean?

Optimising Direct TOF spectrometer (LET) Optimising L₂

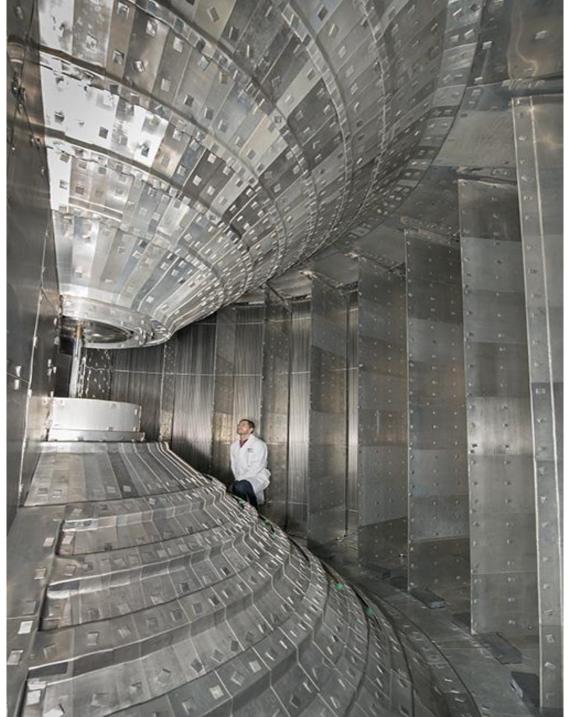
$$\Delta \mathsf{E} = \frac{2E_f \Delta t}{t_f}$$



$$\Delta \mathsf{E} = \frac{2E_f \Delta t}{252.8 \lambda_f L_2}$$

- Better resolution with larger L₂
- For fixed ΔE count rate goes as L₂

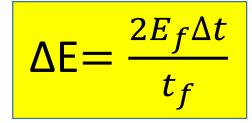
- Make L₂ as large as possible
- Gives better resolution
- Gives larger count rate
- Remember cost goes as L₂²

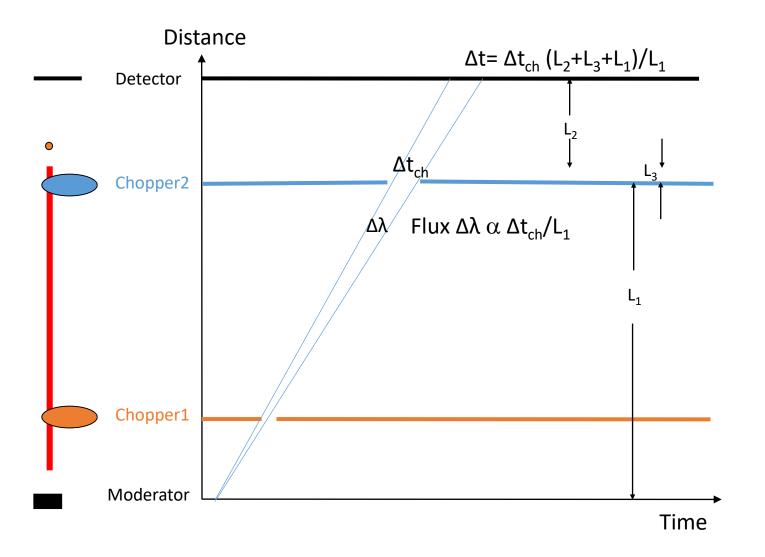


Main cost of Direct geometry spectrometers from secondary

• LET has 40 m² of detectors (approx. £20M today)

Optimising Direct TOF spectrometer (LET) Optimising L₁

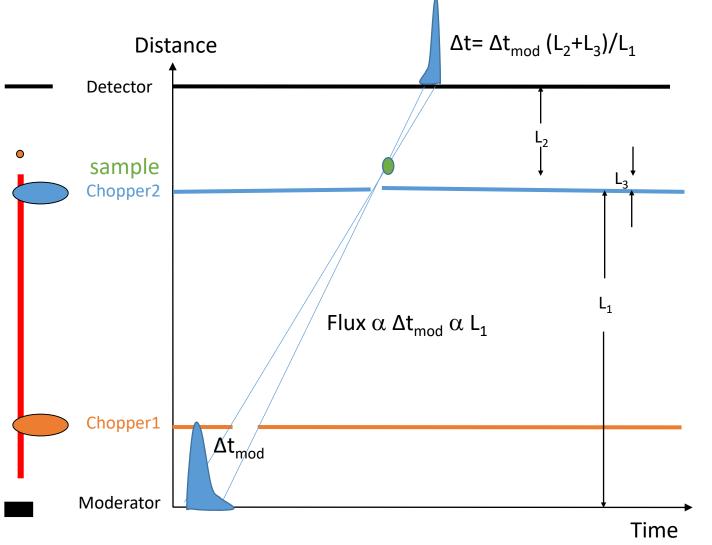




- keep ΔE constant just vary L₁
- So $\Delta t L_2$, L_3 are kept constant
- Flux $\alpha \Delta t_{ch}/L_1 = \Delta t/(L_2 + L_3 + L_1)$
- Flux α 1/ (a+L₁) where (a=L₂+L₃)
- When $L_1 >> L_2 + L_3$
- flux $\approx \alpha$ 1/L₁ for constant ΔE

Make L₁ as short as possible ????

Optimising L₁

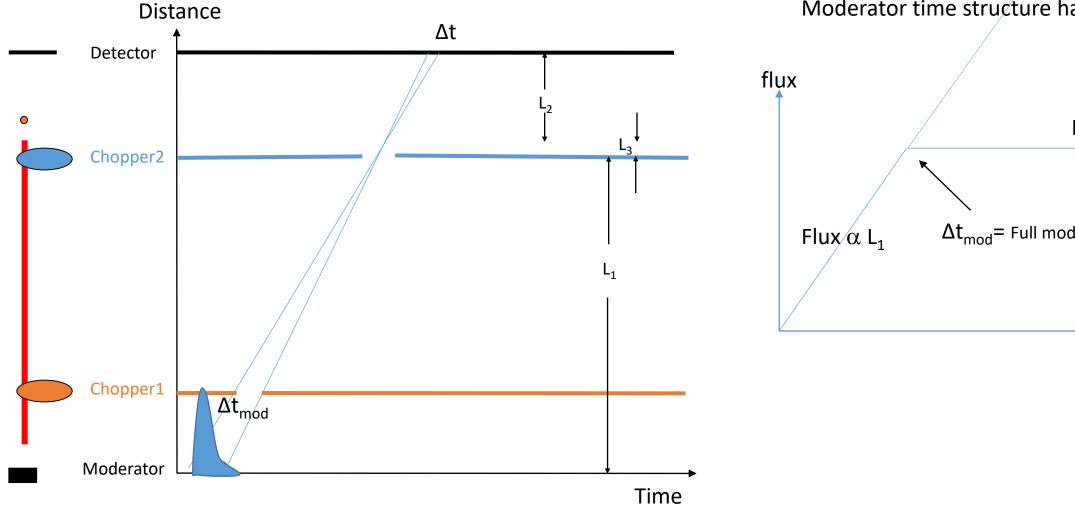


$$\Delta \mathsf{E} = \frac{2E_f \Delta t}{t_f}$$

- keep ΔE constant just vary L₁
- So Δt L₂, L₃ are kept constant
- flux $\alpha \Delta t_{mod} = L_1 \Delta t / (L_2 + L_3)$
- flux αL_1

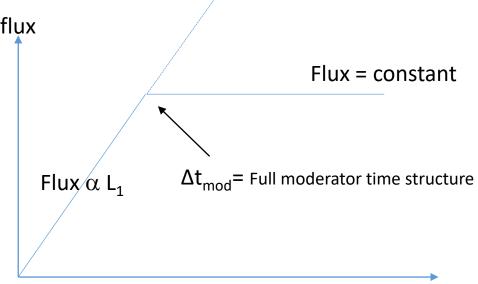
So flux increases with L₁ ???





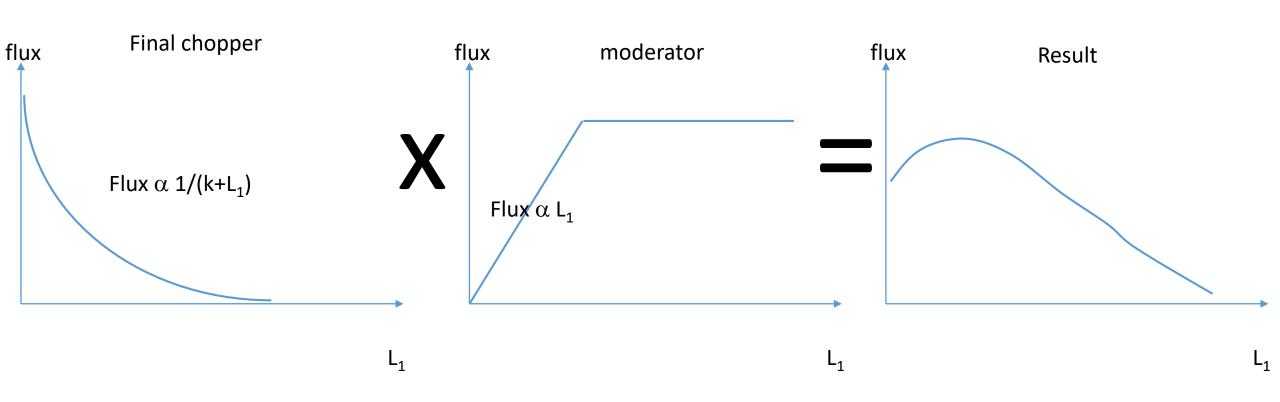
Moderator time structure has finite width

ESS



Optimising Direct TOF spectrometer (LET) Optimising L₁

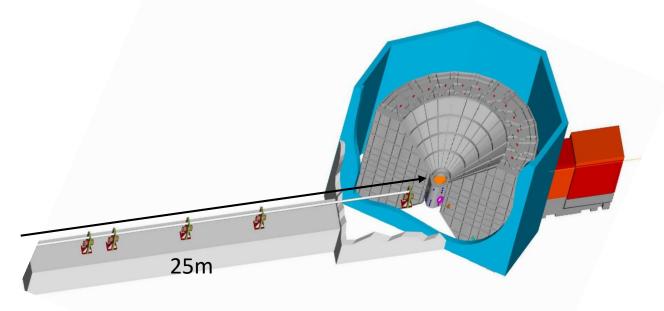
Total flux is chopper term x moderator term

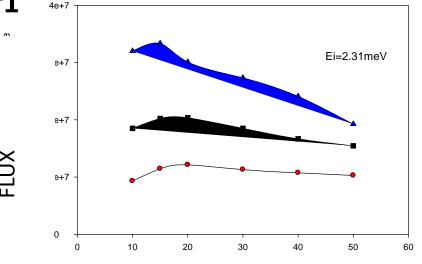


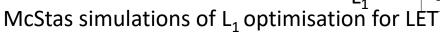
N.B This is nothing to do with guides. It's a consequence of the resolution equations

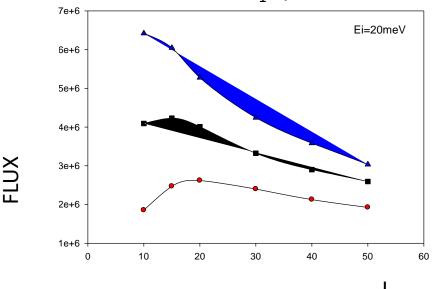
Optimising L₁

- Optimal L₁ is slightly less than 20m
- For LET set L₁=25m as need space for detectors

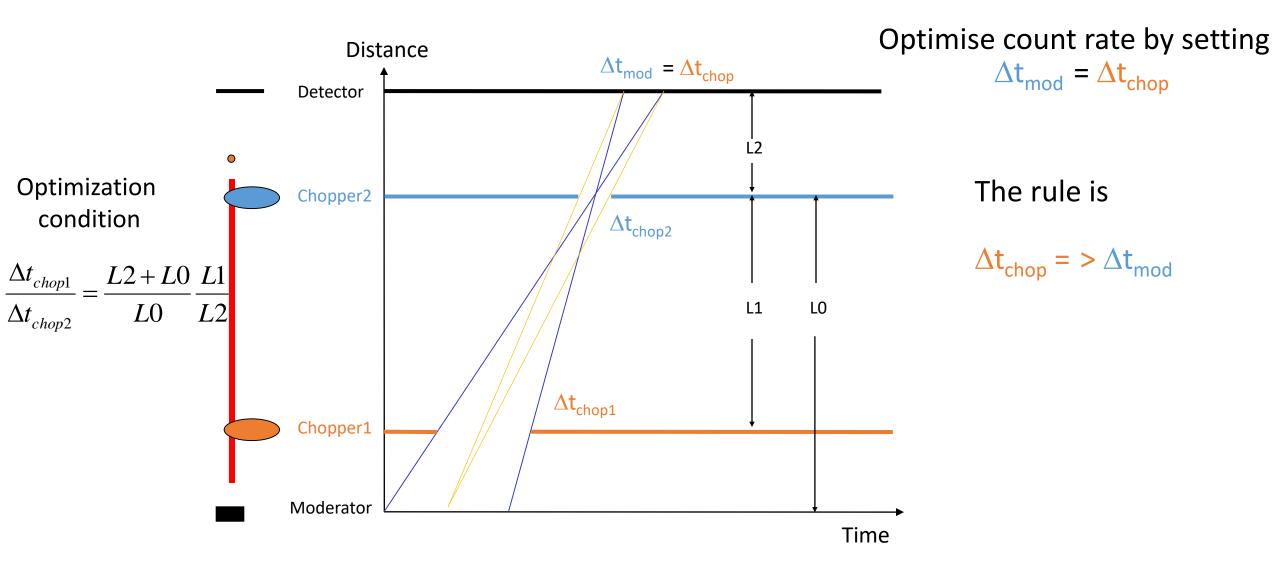








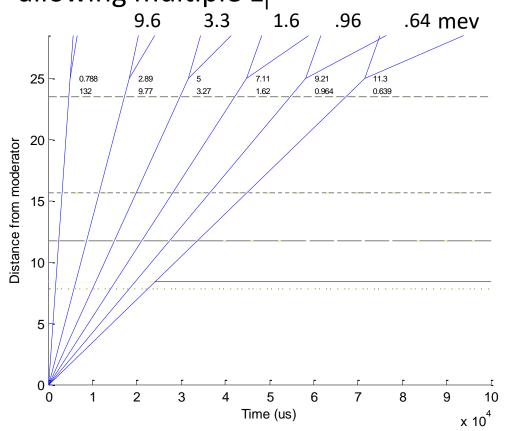
Optimising Direct TOF spectrometer (LET) Optimising chopper openings

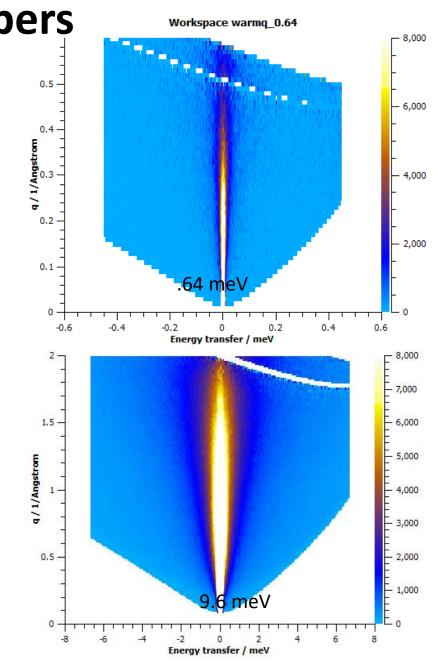


Optimising choppers

*Single measurement use fraction of time frame

•LET has a special chopper arrangement allowing multiple E_i





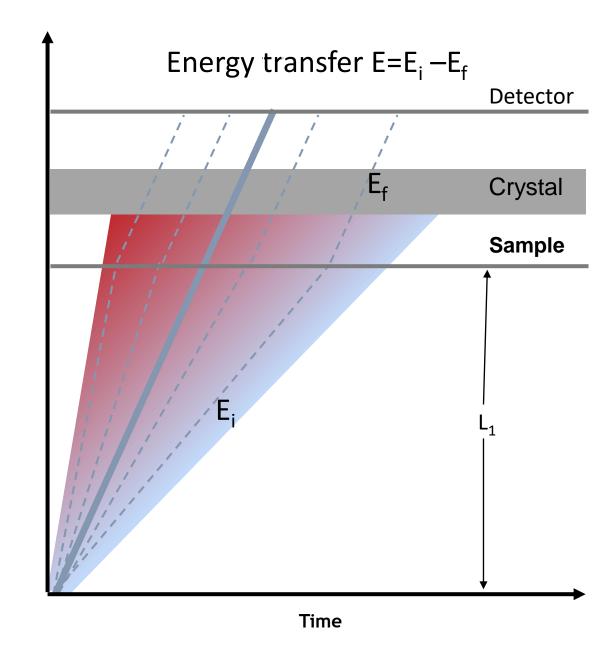
RULES to maximise your signal

- Make L2 as large as possible (better resolution and flux)
- There is an optimal L1
- Match moderator and chopper terms $\Delta E_{mod} = \Delta E_{chop}$ when possible
- Include RRM to fill the time frame

'The Maths'

$$\Delta \mathsf{E}^2 = \left[\left(\frac{2E_i \Delta t_{sam}}{t_i} \right)^2 + \left(2\left(\frac{\Delta d}{d} + cot\theta \Delta \theta \right) \left(\frac{t_f E_i}{t_i} + E_f \right) \right)^2 + \left(\frac{2E_i \Delta t m_{od}}{t_i} \right)^2 \right]$$

$$\begin{array}{c} \mathsf{Sample \ can} \\ \mathsf{Detector \ bank} \\ \mathsf{(mica\ analysed)} \\ \mathsf{Transmitted \ beam} \\ \mathsf{monitor} \\ \mathsf{Graphite \ analyser} \\ \mathsf{bank} \\ \mathsf{(graphite \ analysed)} \\ \end{array}$$



 $E=E_i-E_f$ where E_f fixed

$$\frac{\Delta E}{\Delta t} = -\frac{\Delta E_i}{\Delta t}$$

Where $E_i = \frac{1}{2} m_n (L_1/t_i)^2$

$$\Delta \mathsf{E} = \frac{2E_i \Delta t}{t_i}$$

 $\begin{array}{lll} \Delta \mathsf{E} & & \text{energy uncertainty} \\ \Delta t & & \text{time uncertainty at detector} \\ \mathsf{E_i} & & \text{energy of neutron on sample} \\ \mathsf{t_i} & & \text{flight time moderator to sample} \end{array}$

Distance

Comparison In-direct v direct

In-direct

$$\Delta \mathsf{E} = \frac{2E_i \Delta t}{t_i}$$

Direct

$$\Delta \mathsf{E} = \frac{2E_f \Delta t}{t_f}$$

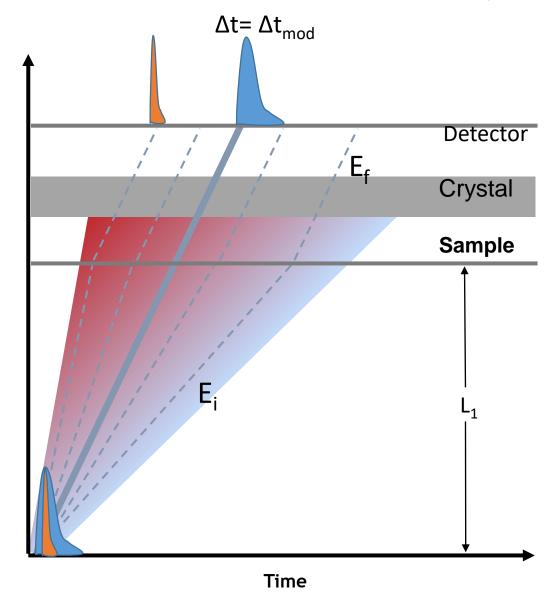
 $t_i >> t_f$ so ΔE smaller on In-direct for same Δt (In-direct for high resolution)

OR

For same ΔE can accept larger Δt and hence more flux (In-direct for higher count rate)

Sources of Δt

1) Moderator



Distance

$$\Delta \mathsf{E} = \frac{2E_i \Delta t}{t_i}$$

• Δt_{mod} doesn't alter from moderator to detector

$$\Delta \mathsf{E} = rac{2E_i \Delta t m_{od}}{t_i}$$

RESOLUTION

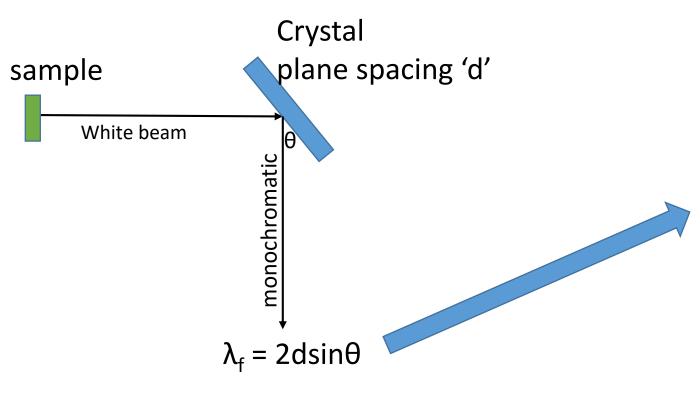
• Increasing L_1 reduces contribution of Δt_{mod} on ΔE

FLUX

• Flux α Δt_{mod} but is independent of L_1 or L_2

Sources of Δt

2) Analyser Crystal



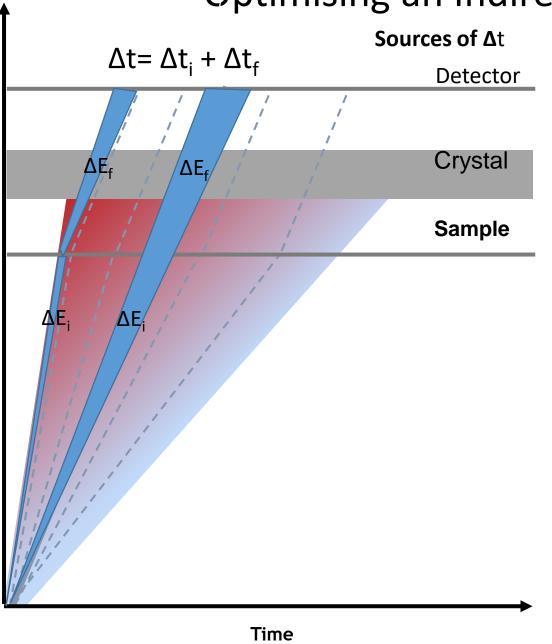
Differentiating Braggs law gives

$$\frac{\Delta \lambda_f}{\lambda_f} = \frac{\Delta d}{d} + \cot \theta \Delta \theta$$

Δλ is wavelength spread
Δd is spread in d spacing
Θ is Bragg angle
ΔΘ Bragg angle spread (mosaic + sample size)

$$\frac{\Delta E}{E} = \frac{2\Delta\lambda}{\lambda} \quad \text{thus} \quad \Delta E_f = 2E_f \left(\frac{\Delta d}{d} + cot\theta\Delta\theta\right)$$

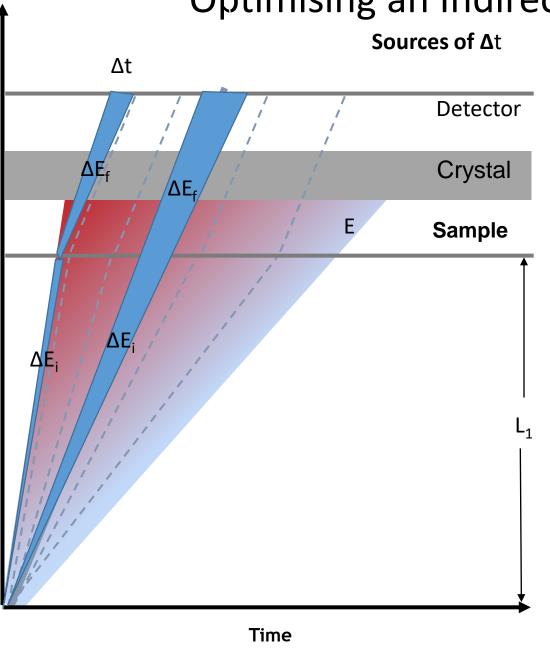
 ΔE_f is the energy spread of neutrons coming off the analyser



2) Analyser Crystal

$$\Delta \mathsf{E} = \frac{2E_i \Delta t}{t_i}$$

- 1) $\Delta E_f = 2E_f \left(\frac{\Delta d}{d} + cot\theta \Delta\theta\right)$ is a constant
- 2) Or $\Delta t_f = t_f \left(\frac{\Delta d}{d} + cot\theta \Delta \theta \right)$
- 3) BUT $\Delta E_i = \Delta E_f$ for any constant energy transfer
- 4) gives $\Delta t_i = t_i \frac{E_f}{E_i} (\frac{\Delta d}{d} + \cot \theta \Delta \theta)$ from 1) and 3)
- 5) $\Delta t = \Delta t_i + \Delta t_f$
- Thus $\Delta t = (\frac{\Delta d}{d} + \cot\theta \Delta\theta)(t_i \frac{E_f}{E_i} + tf)$ from 2) and 4)
- Or $\Delta E = 2 \left(\frac{\Delta d}{d} + \cot \theta \Delta \theta \right) \left(\frac{t_f E_i}{t_i} + E_f \right)$



$$\Delta \mathsf{E} = \frac{2E_i \Delta t}{t_i}$$

RESOLUTION

•
$$\Delta E = 2 \left(\frac{\Delta d}{d} + \cot \theta \Delta \theta \right) \left(\frac{t_f E_i}{t_i} + E_f \right)$$

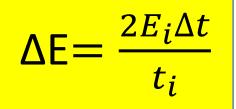
• As $\theta \rightarrow 90 \cot \theta \Delta \theta \rightarrow 0$ minimise ΔE

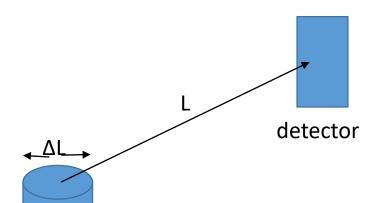
FLUX

- Flux $\alpha \Delta E_f = 2E_f \left(\frac{\Delta d}{d} + \cot\theta \Delta \theta\right)$
- Increase flux with decreasing θ but independent of L_1 and L_2

Sources of Δt

3) Sample/detector size



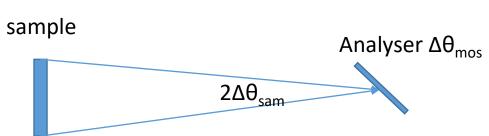


sample

$$\Delta \mathsf{E} = \frac{2E_i \Delta t s_{am}}{t_i}$$

Where
$$\frac{\Delta t_{sam}}{t_f} = \frac{\Delta L}{L}$$

More important is the angular uncertainty from sample size $\Delta\theta_{\text{sam}}$



The total uncertainty
$$\Delta\theta^2 = \Delta\theta^2_{sam} + \Delta\theta^2_{mos}$$

Assuming no correlations all the terms convolute together. A good approximation is to add in quadrature

$$\Delta E^2 = \Delta E^2_{sam} + \Delta E^2_{ana} + \Delta E^2_{mod}$$

$$\Delta E^{2} = \left[\left(\frac{2E_{i}\Delta t_{sam}}{t_{i}} \right)^{2} + \left(2\left(\frac{\Delta d}{d} + cot\theta\Delta\theta \right) \left(\frac{t_{f}E_{i}}{t_{i}} + E_{f} \right) \right)^{2} + \left(\frac{2E_{i}\Delta t_{mod}}{t_{i}} \right)^{2} \right]$$

Summary

- Flux from moderator or analyser does not depend on L1 or L2
- Only thing changing flux is energy spread accepted by analyser $\Delta E = E \cot \theta \Delta \theta$
- Flux increases as Bragg angle θ reduces but resolution worsens
- ΔE_{mod} reduces with L1 so make L1 large for high resolution (reduces bandwidth)

$$\Delta E_{\text{mod}} = \frac{2E_i \Delta t mod}{t_i}$$

- For a high resolution instrument like IRIS to maximise flux want $\Delta E_{ana} = \Delta E_{mod}$
- Matching resolution terms just like we did on LET

IRIS resolution

IRIS parameters

Ef=1.84 mev
$$\Delta t_{mod}$$
 =120 μs L1=36.41 m t_i =61371 μs $\Delta d/d$ =6E-4 $\Delta \theta$ =1.3°

$$\Delta E_{\text{mod}} = \frac{2E_i \Delta t mod}{t_i}$$

$$\Delta E_{\text{mod}} = 7.2 \, \mu \text{eV}$$

$$\Delta E_{\text{ana}} = 2 \left(\frac{\Delta d}{d} + \cot \theta \Delta \theta \right) \left(\frac{t_f E_i}{t_i} + Ef \right)$$

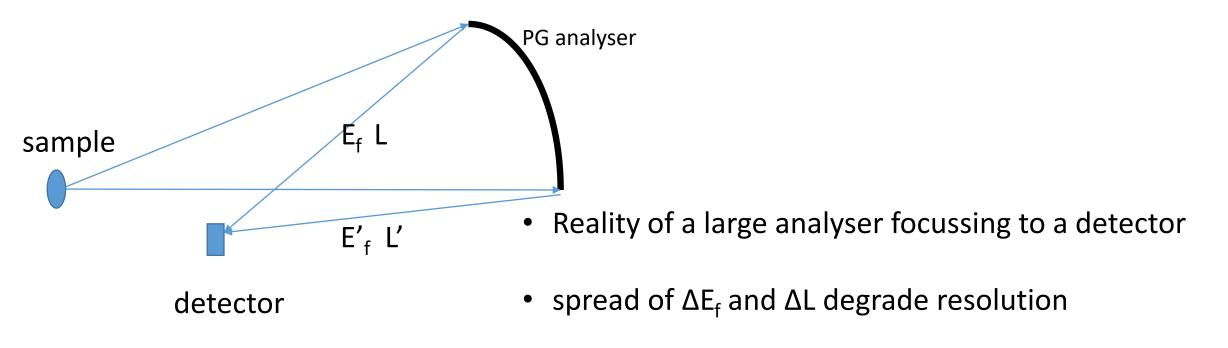
$$\Delta E_{ana} = 3.6 \mu eV$$

$$\Delta E = \operatorname{sqrt}(\Delta E_{\text{mod}}^2 + \Delta E_{\text{ana}}^2) = 8 \, \mu eV$$

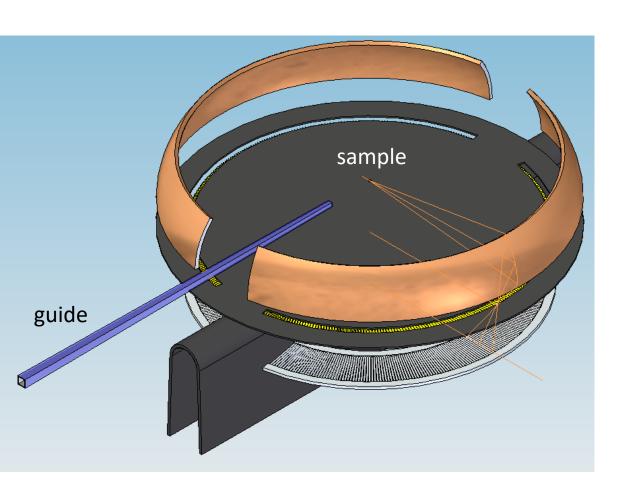
IRIS resolution = $17 \mu eV$ why the large difference?

Optimising an Indirect TOF spectrometer (IRIS) IRIS resolution

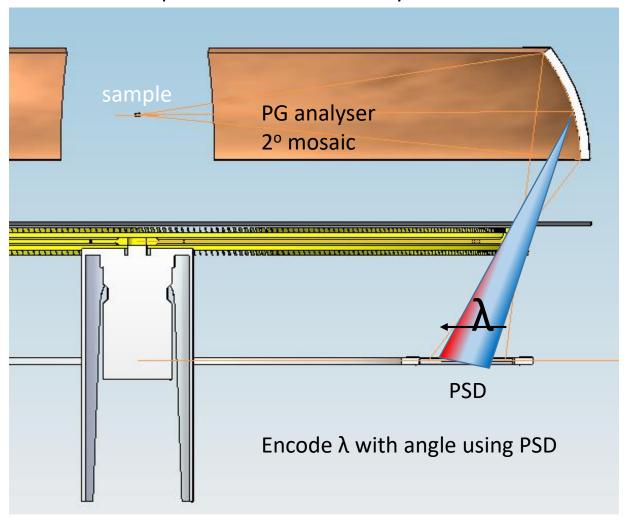
Real resolution worse than calculated due to analyser size

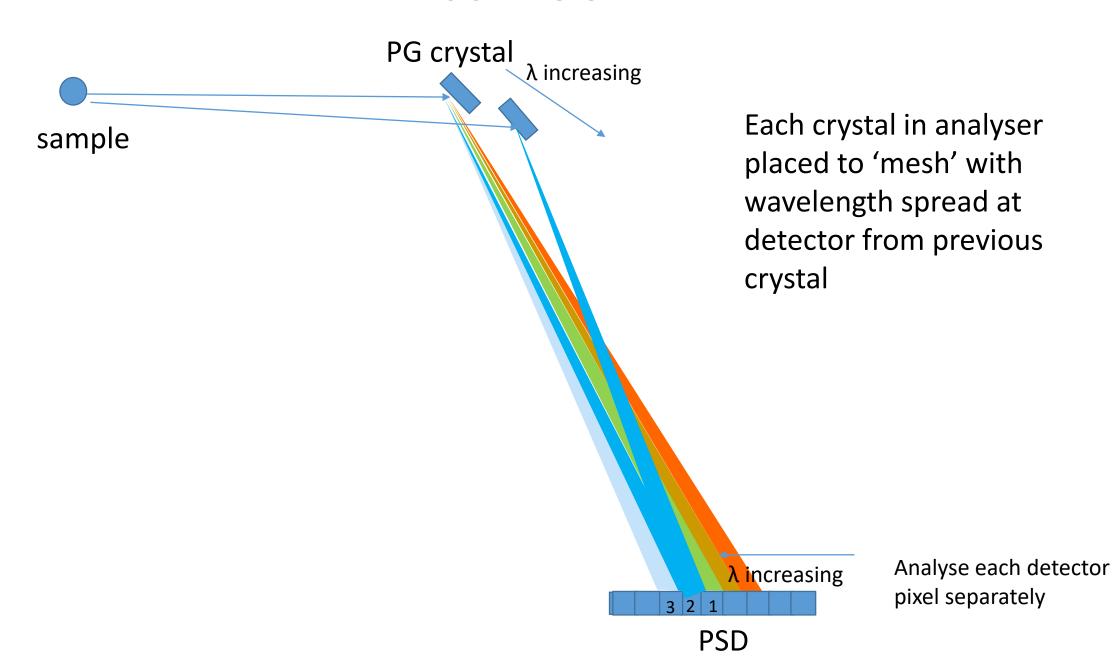


Nearly all the pain in designing in-direct is how to reduce spread of final energies Ef and path lengths L

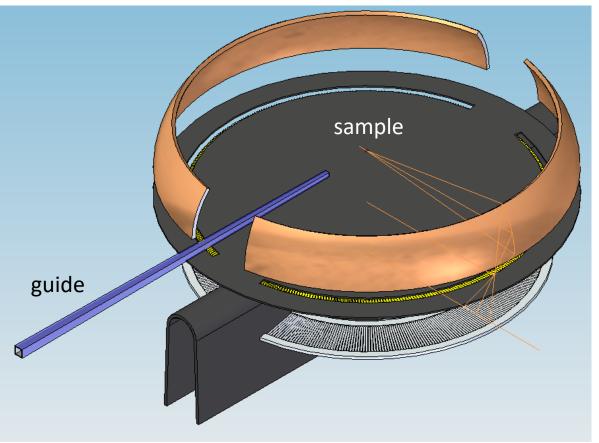


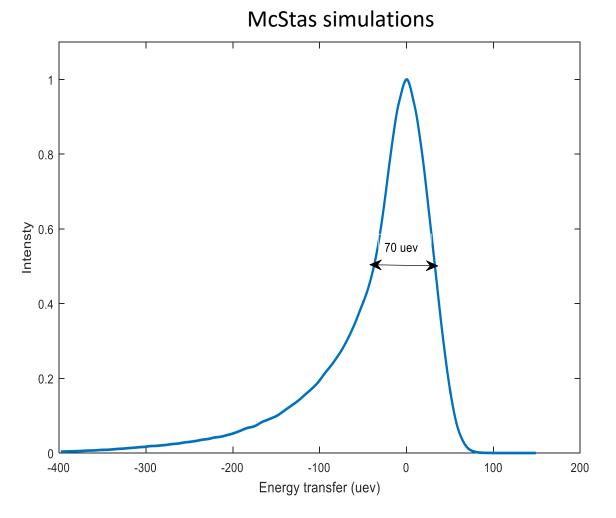
Use prismatic effect of PG crystals





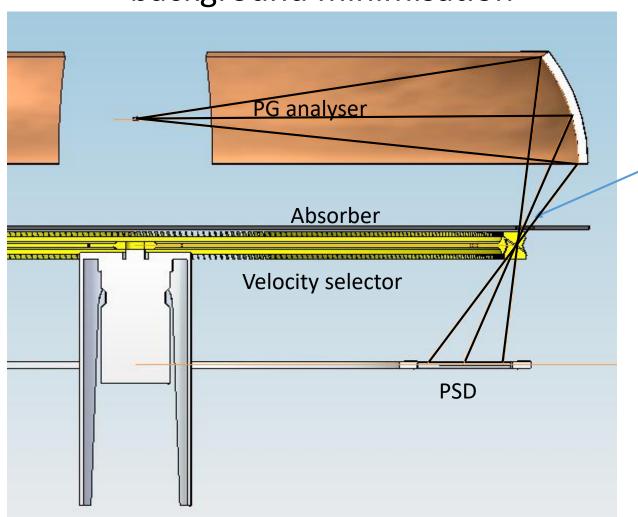
Mushroom has 70 μ eV resolution (elastic) Would be 260 μ eV without λ encoding





The large $\Delta\lambda$ range accepted by analyser gives 50 x flux of LET at same resolution

background minimisation



Neutron all focus though a point at velocity selector

Point on detector can only see point on analyser to minimised background from thermal scattering from PG crystals

Mushroom designed to maximise signal/background

Optimising a TOF spectrometer SUMMARY

- Don't start with McStas
- Know your key drivers
- FOM is signal/background
- Demonstrated how crucial it is to understand the 'Maths' of the instrument
 - -how it relates to resolution
 - -how it relates to count rate
- Without this information you WILL NOT build an optimal instrument

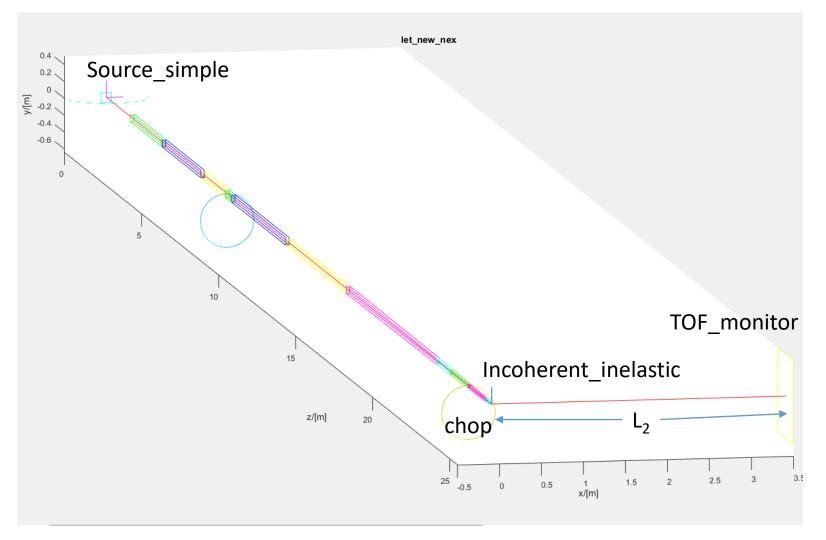
Have not even mentioned guides yet

.....but we will leave that for another day



Things to try

Direct geometry



- Keep Ei and chop opening the same
- Scatter to same solid angle
- Vary L₂
- Determine resolution and count rate v L₂

COMPONENT sample=Incoherent_inelastic(radius=0.005,yheight=0.01,focus_aw=5, focus_ah=5,target_z=0.0,target_y=0.0, target_x=1.0)
AT (0, 0, 24.999) RELATIVE Origin

Time-of-flight monitor C:\My_Docs\mc\LET_mcstas\let_new_nex_20210413_184840\detector.dat 0.01 Intensity Δt 0.005 2.915 2.9 2.905 2.91 2.92 2.925 2.93 Time-of-flight [\gms]

remember

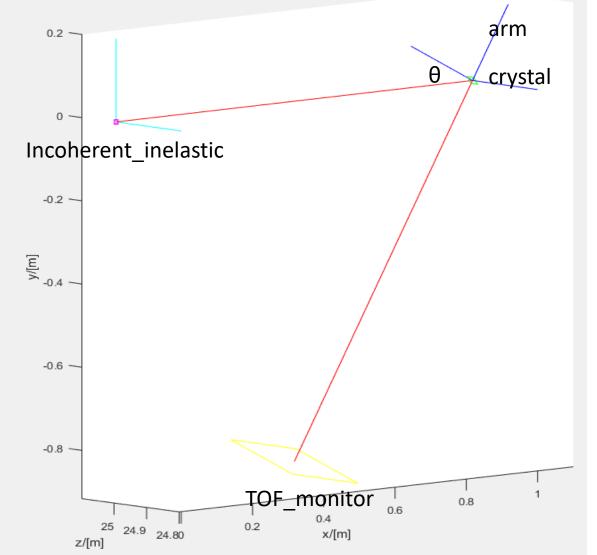
$$\Delta \mathsf{E} = \frac{2E_f \Delta t}{t_f}$$

- Elastic so $E_f = E_i$
- know L₂ so calculate t_f
- Get Δt from monitor (it's the FWHM)
- NOW calculate resolution ΔE
- Should be able to predict Δt from geometry and chopper opening
- How does resolution and flux vary with L₂should be able to predict
- If time. Vary L_2 but keep ΔE constant by varying chopper opening. Plot flux versus L_2

 $\times 10^4$

Things to try

In-direct geometry

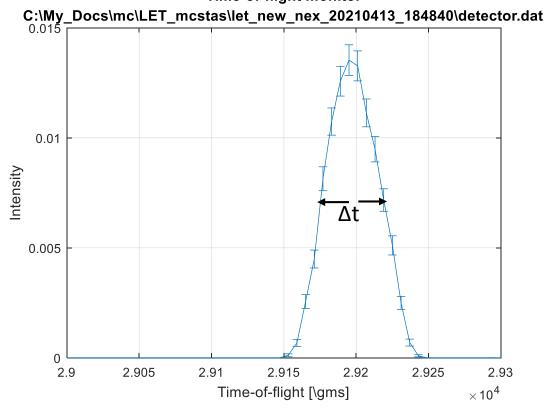


- Scatter from incoherent sample to analyser (remember to focus)
- Vary θ
- Determine resolution and count rate v θ

Trick is to use arm component to rotate crystal and TOF_monitor

COMPONENT analyser=Monochromator_flat(zwidth=0.02, yheight=0.02, mosaich=120.0, mosaicv=120.0,Q=1.8734) AT (0,0,1) RELATIVE detectorarm ROTATED (-rot, 90,0) RELATIVE detectorarm

Time-of-flight monitor



remember

$$\Delta \mathsf{E} = \frac{2E_i \Delta t}{t_i}$$

- Elastic so E_f = E_i
- know L₁ so calculate t_i
- Get Δt from monitor (it's the FWHM)
- NOW calculate resolution ΔΕ
- Should be able to predict Δt from geometry and crystal angle
- How does resolution and flux vary with θ should be able to predict

My solution Direct

```
DECLARE
%{
double v foc;
double lam min,lam max,den;
%}
// #include <math.h>
INITIALIZE
%{
v foc = SE2V*sqrt(E0);
TRACE
COMPONENT Origin = Arm()
 AT (0,0,0) ABSOLUTE
COMPONENT Source = Source_simple(radius=0.05, dist=1.7, focus_xw=.04, focus_yh=.09, E0=E0, dE=dE)
 AT (0, 0, 0) RELATIVE Origin
COMPONENT guide1 = Guide_channeled(
  w1 = 0.04, h1 = 0.09, w2 = 0.02, h2 = 0.04868, l = 22.68, alphax = 4.38, alphay = 4.38,
W=3e-3, mx = 4, my = 4)
AT (0, 0, 1.7) RELATIVE Origin
```

```
COMPONENT funnel = Guide channeled(
  w1 = 0.031, h1 = 0.05711, w2 = 0.02, h2 = 0.04868, l = 1.117, alphax = 4.38, alphay = 4.38,
W=3e-3, mx = 4, my = 4)
AT (0, 0, 22.373) RELATIVE Origin
COMPONENT Res5 = Chopper(
  R = 0.320, f = 2*PI*Res5, n=2, w=slit, pha=23.496/v foc)
AT (0, 0, 23.496) RELATIVE Origin
COMPONENT Res5 counter = Chopper(
  R = 0.320, f = -2*PI*Res5, n=2, w=slit, pha=-23.504/v foc)
AT (0, 0, 23.504) RELATIVE Origin
COMPONENT endguide = Guide channeled(
 w1 = 0.02, h1 = 0.0484, w2 = 0.020, h2 = 0.04, l = 1.1, alphax = 4.38, alphay = 4.38,
W=3e-3, mx = 4, my = 4)
AT (0, 0, 23.52) RELATIVE Origin
COMPONENT snot = Guide channeled(
 w1 = 0.02, h1 = 0.04, w2 = 0.02, h2 = 0.04, l = 0.23, alphax = 4.38, alphay = 4.38,
W=3e-3, mx = 4, my = 4)
AT (0, 0, 24.622) RELATIVE Origin
COMPONENT sample = Incoherent inelastic(radius=0.005,yheight=0.01,efmin=2.5,efmax=6.5,
   focus aw=5, focus ah=5,target z=0.0,target y=0.0,target x=1,sigma abs=0, sigma inc=5.08)
AT (0, 0, 24.999) RELATIVE Origin
COMPONENT detectorarm = Arm()
AT (0, 0,0) RELATIVE sample
ROTATED (0,90.0,0) RELATIVE sample
COMPONENT TOF target = TOF monitor(
  nt = 100, filename = "detector.dat",xmin = -0.5,
 xmax = 0.5, ymin = -0.5, ymax = 0.5, tmin = (1e6*(25+L2)/v foc) - 200, tmax = (1e6*(25+L2)/v foc) + 200
 AT (0, 0, L2) RELATIVE detectorarm
```

My solution Direct

```
DEFINE INSTRUMENT LET Mantid(I0=4,dl=3,rot=0)
```

TRACE

My solution to in-direct

```
COMPONENT Origin = Arm()
AT (0,0,0) ABSOLUTE
COMPONENT Source = Source simple(radius=0.1, dist=25, focus xw=.01, focus yh=.01, lambda0=l0, dlambda=dl)
AT (0, 0, 0) RELATIVE Origin
COMPONENT sampleMantid = Incoherent inelastic(radius=0.005, yheight=0.01, efmin=2.5, efmax=6.5, focus aw=1,
focus_ah=1,target_z=0.0,target_y=0.0,target_x=1,sigma_abs=0, sigma_inc=5.08)
AT (0, 0, 24.999) RELATIVE Origin
COMPONENT detectorarm = Arm()
AT (0, 0,0) RELATIVE sampleMantid
 ROTATED (0,90.0,0) RELATIVE sampleMantid
COMPONENT graphite analyser31=Monochromator flat(zwidth=0.02, yheight=0.02, mosaich=120.0,
mosaicv=120.0,Q=1.8734)
AT (0,0,1) RELATIVE detectorarm
ROTATED (-rot, 90,0) RELATIVE detectorarm
COMPONENT crystalarm = Arm()
AT (0, 0,1) RELATIVE detectorarm
ROTATED (-2*rot,0,0) RELATIVE detectorarm
COMPONENT TOF target = TOF monitor(
  nt = 1000, filename = "detector.dat",xmin = -0.1,
 xmax = 0.1, ymin = -0.1, ymax = 0.1, tmin=25000, tmax = 50000)
 AT (0, 0, -1) RELATIVE crystalarm
END
```