

El Lagrangiano en términos de Q y \dot{Q}

$$L = \frac{1}{2} m (\dot{Q}^2 - b^2 Q^2) + \frac{b^2 m Q^2}{4} - \frac{1}{2} k^2 Q^2$$

Como el Lagrangiano ya tenemos no aparece explícitamente el tiempo podemos sospechar que será la simetría que tenemos

$$\frac{dL}{dt} = 0$$

$$\frac{dL}{dt} = \frac{\partial L}{\partial Q} \dot{Q} + \frac{\partial L}{\partial \dot{Q}} \ddot{Q} \quad \frac{\partial L}{\partial \dot{Q}} = \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{Q}} \right)$$

$$\frac{dL}{dt} = \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{Q}} \dot{Q} + \frac{\partial L}{\partial Q} \right) = \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{Q}} \dot{Q} \right)$$

$$\frac{d}{dt} \left(L - \frac{\partial L}{\partial \dot{Q}} \dot{Q} \right) = 0$$

(conservación de la energía)

$$\frac{dL}{dt} = m \dot{Q} - \frac{b^2 m Q}{2} \quad \frac{\partial L}{\partial \dot{Q}} = m \dot{Q}^2 - \frac{b^2 m Q \dot{Q}}{2}$$

$$L - \frac{\partial L}{\partial \dot{Q}} \dot{Q} = \frac{1}{2} m \dot{Q}^2 - \frac{b^2 m Q \dot{Q}}{2} - \frac{1}{2} k^2 Q^2 - m \dot{Q}^2 + \frac{b^2 m Q \dot{Q}}{2}$$

$$= -\frac{1}{2} m \dot{Q}^2 + \frac{b^2 m Q}{2} - \frac{1}{2} k^2 Q^2$$

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Pregunta 4

Considere el siguiente Lagrangiano con un grado de libertad

$$L = e^{bt} \left(\frac{1}{2} m \dot{q}^2 - \frac{1}{2} k^2 q^2 \right) \quad k, b, m \text{ constantes positivas}$$

a) Encontrar las ecuaciones de Euler-Lagrange

$$\frac{\partial L}{\partial q} = -e^{bt} k^2 q \quad \frac{\partial L}{\partial \dot{q}} = e^{bt} m \dot{q}$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}} \right) = e^{bt} m \ddot{q} + b e^{bt} m \dot{q}$$

$$e^{bt} m \ddot{q} + b e^{bt} m \dot{q} + e^{bt} k^2 q = 0$$

$$m \ddot{q} + b m \dot{q} + k^2 q \longrightarrow \ddot{q} + b \dot{q} + \omega_0^2 k q = 0$$

→ El sistema a) que se parece es a un péndulo amortiguado.

b) Realiza un cambio de variable $Q = e^{bt/2} q$ y construye el Lagrangiano. Encuentra la simetría continua y deduce la cantidad conservada asociada a ella usando Teorema de Noether. Pre-escribe la cantidad en términos de q

Podemos sustituir al nivel del Lagrangiano

$$q = Q e^{-bt/2}$$

$$\dot{q} = \dot{Q} e^{-bt/2} - \frac{b}{2} e^{-bt/2} Q$$

$$L = e^{bt} \left(\frac{1}{2} m \dot{q}^2 - \frac{1}{2} k^2 q^2 \right)$$

$$q^2 = Q^2 e^{-bt}$$

$$\dot{q}^2 = e^{-bt} \left(\dot{Q} - \frac{b}{2} Q \right)^2 = e^{-bt} \left(\dot{Q}^2 - b \dot{Q} Q + \frac{b^2}{4} Q^2 \right)$$

$$L = \frac{1}{2} m \left(\dot{Q}^2 - b \dot{Q} Q + \frac{b^2}{4} Q^2 \right) - \frac{1}{2} k^2 Q^2$$

$$\theta = 0 \quad \dot{\theta} = 0 \quad \text{Puntos de equilibrio}$$

$$-g + \frac{k}{m} r = 0 \quad g + \frac{k}{m} r = 0$$

$$r_1 = \frac{g m}{k} \quad r_2 = -\frac{g m}{k} \quad \theta_1 = 0 \quad \theta_2 = \pi$$

c) Haz una expansión alrededor de los puntos de equilibrio y resuelve el sistema.

Para $\theta = 0$ el punto de equilibrio será

$$P_e = r + \frac{g m}{k}$$

Para $\theta = \pi$ punto de equilibrio será

$$P_e' = r - \frac{g m}{k}$$

Las ecuaciones de movimiento que obtenemos serán

$$\ddot{r} - (r + l)\dot{\theta}^2 - g + \frac{k}{m} r = 0$$

$$(r + l)\ddot{\theta} + 2\dot{\theta}\dot{r} + g \sin \theta = 0$$

linealizando $\cos \theta \approx 1$ y $\sin \theta \approx \theta$

$$\ddot{r} - (r + l)\dot{\theta}^2 - g + \frac{k}{m} r = 0$$

$$(r + l)\ddot{\theta} + 2\dot{\theta}\dot{r} + g\theta = 0$$

Tomamos punto de equilibrio

$$\left(-r + \frac{g m}{k}\right)\ddot{\theta} + 2\dot{\theta}\left(r + \frac{g m}{k}\right) + g\theta = 0$$

$$\left(r + \frac{g m}{k}\right)\ddot{\theta} + g\theta = 0 \rightarrow \text{Ecuación de un péndulo}$$

$$\ddot{\theta} = -\omega^2 \theta \quad \theta = A \cos(\omega t - \delta) \quad \text{donde } \omega = \sqrt{\frac{g k}{k r + g m}}$$

obtenemos las ecuaciones de movimiento

$$\frac{\partial \mathcal{L}}{\partial \dot{r}} = m(\dot{r} + \dot{\theta}^2 r) + mg \cos \theta - kr = 0$$

$$\frac{\partial \mathcal{L}}{\partial \dot{\theta}} = m\dot{r} \frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{r}} \right) = m\ddot{r}$$

$$m\ddot{r} - m(\dot{r} + \dot{\theta}^2 r) + mg \cos \theta - kr = 0$$

$$\ddot{r} - (\dot{r} + \dot{\theta}^2 r) - g \cos \theta + \omega_0^2 r = 0$$

$$\frac{\partial \mathcal{L}}{\partial \theta} = -mg(\dot{r} + \dot{\theta}^2 r) \sin \theta$$

$$\frac{\partial \mathcal{L}}{\partial \dot{\theta}} = m(\dot{r} + \dot{\theta}^2 r) \frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{\theta}} \right) = m(\dot{r} + \dot{\theta}^2 r) \ddot{\theta} + 2m(\dot{r} + \dot{\theta}^2 r) \dot{\theta}$$

$$m(\dot{r} + \dot{\theta}^2 r) \ddot{\theta} + 2m(\dot{r} + \dot{\theta}^2 r) \dot{\theta} + mg(\dot{r} + \dot{\theta}^2 r) \sin \theta = 0$$

$$(\dot{r} + \dot{\theta}^2 r) \ddot{\theta} + 2\dot{\theta} \dot{r} + g \sin \theta = 0$$

las ecuaciones de movimiento son

$$\ddot{r} - (\dot{r} + \dot{\theta}^2 r) - g \cos \theta + \omega_0^2 r = 0$$

$$(\dot{r} + \dot{\theta}^2 r) \ddot{\theta} + 2\dot{\theta} \dot{r} + g \sin \theta = 0$$

b) Encuentre los puntos de equilibrio y describa su estabilidad para encontrar los puntos de equilibrio

$$\ddot{r} = \dot{r} = \dot{\theta} = \ddot{\theta} = 0$$

$$-g \cos \theta + \frac{kr}{m} = 0$$

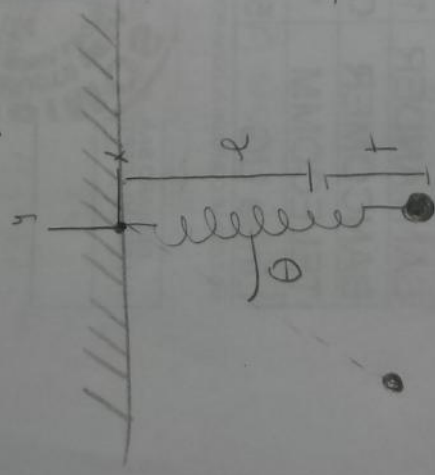
$$g \sin \theta = 0 \implies \theta = n\pi \quad n=0,1,\dots \quad \theta=0 \text{ y } \theta=\pi$$

sustituimos para obtener los puntos en r

Pregunta 3

A un péndulo simple se le reemplaza por un resorte de longitud en reposo l y constante K

a) Construye el Lagrangiano del sistema y deduce las ecuaciones de Euler-Lagrange.



$$\mathcal{L} = T - V$$

$$\text{donde } T = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2)$$

→ 2 grados de libertad

→ hacemos cambio de variable

$$x = R \sin \theta$$

$$y = -R \cos \theta$$

$$\text{donde } R = (l + r)$$

$$\dot{x} = (l + r) \sin \theta$$

$$\dot{y} = -(l + r) \cos \theta$$

$$\dot{x}^2 = (l + r)^2 \sin^2 \theta \dot{\theta}^2$$

$$\dot{y}^2 = (l + r)^2 \cos^2 \theta \dot{\theta}^2$$

$$\dot{x}^2 + \dot{y}^2 = (l + r)^2 (\sin^2 \theta + \cos^2 \theta) \dot{\theta}^2 = (l + r)^2 \dot{\theta}^2$$

$$T = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2) = \frac{1}{2} m (l + r)^2 \dot{\theta}^2$$

$$V = mgy + \frac{1}{2} K r^2$$

Para el potencial sea debido al potencial gravitacional + el del resorte.

$$V = mgy + \frac{1}{2} K r^2$$

$$V(\theta, r) = -mg(l + r) \cos \theta + \frac{1}{2} K r^2$$

$$\mathcal{L}(r, \theta) = \frac{1}{2} m ((l + r)^2 \dot{\theta}^2 + \dot{r}^2) + mg(l + r) \cos \theta - \frac{1}{2} K r^2$$

Jugamos, con índices Contrarios

$$\delta S = \frac{1}{2} \int (g_{pa} \{^p_{bc}\} + g_{pa} \{^p_{bc}\}) \dot{q}^a \dot{q}^b \dot{q}^c + 2 g_{ab} \delta \dot{q}^a \dot{q}^b dt$$

$$\delta S = \int g_{pa} \{^p_{bc}\} \dot{q}^a \dot{q}^b \dot{q}^c + g_{ab} \delta \dot{q}^a \dot{q}^b dt$$

$$\delta S = \int g_{ab} \{^b_{pc}\} \dot{q}^a \dot{q}^p \dot{q}^c + g_{ab} \delta \dot{q}^a \dot{q}^b dt$$

$$\delta S = \int g_{ap} \{^p_{bc}\} \dot{q}^a \dot{q}^b \dot{q}^c + g_{ap} \delta \dot{q}^a \dot{q}^b dt$$

$$\delta S = \int g_{ap} \{^a_{bc}\} \dot{q}^p \dot{q}^b \dot{q}^c + g_{ap} \delta \dot{q}^a \dot{q}^b dt = 0$$

Como vemos podemos sacar factor común, la métrica y la velocidad de q^p

$$\{^a_{bc}\} \dot{q}^a \dot{q}^b \dot{q}^c + \delta \dot{q}^a \dot{q}^b \dot{q}^c = 0$$

Con un pequeño análisis tensorial.

Una conexión métrica de 1º especie es la siguiente

$$[\chi_j, \kappa] = \frac{1}{2} \left(\frac{\partial g_{\kappa j}}{\partial x^j} + \frac{\partial g_{j\kappa}}{\partial x^\kappa} - \frac{\partial g_{\chi j}}{\partial x^\chi} \right)$$

Tomamos otra conexión en índices cambiados

$$[\chi_{\kappa, j}] = \frac{1}{2} \left(\frac{\partial g_{\chi j}}{\partial x^\kappa} + \frac{\partial g_{\kappa j}}{\partial x^j} - \frac{\partial g_{\chi \kappa}}{\partial x^j} \right)$$

Si sumamos ambas

$$[\chi_j, \kappa] + [\chi_{\kappa, j}] = \frac{1}{2} \left(\frac{\partial g_{\chi \kappa}}{\partial x^j} + \frac{\partial g_{j\kappa}}{\partial x^\chi} - \frac{\partial g_{\chi j}}{\partial x^\kappa} + \frac{\partial g_{\kappa j}}{\partial x^j} + \frac{\partial g_{\chi j}}{\partial x^\kappa} - \frac{\partial g_{\chi \kappa}}{\partial x^j} \right)$$

$$[\chi_j, \kappa] + [\chi_{\kappa, j}] = \frac{\partial g_{\kappa j}}{\partial x^\chi}$$

En términos de conexiones de 2º especie

$$\left\{ \begin{matrix} e \\ \chi_j \end{matrix} \right\} = g^{\kappa j} [\chi_j, \kappa] \longrightarrow g^{\kappa m} \left\{ \begin{matrix} e \\ \chi_j \end{matrix} \right\} = [\chi_j, m]$$

$$[\chi_j, \kappa] + [\chi_{\kappa, j}] = \frac{\partial g_{\kappa j}}{\partial x^\chi}$$

$$\frac{\partial g_{\kappa j}}{\partial x^\chi} = g^{\epsilon i} \left\{ \begin{matrix} e \\ \kappa \chi \end{matrix} \right\} + g^{\epsilon \kappa} \left\{ \begin{matrix} e \\ \chi j \end{matrix} \right\}$$

Retomamos

$$\delta S = \int \frac{\partial g_{ab}}{\partial q^a} \frac{\partial q^a}{\partial q^c} \delta q^c \dot{q}^a \dot{q}^b + 2 g_{ab} \delta \dot{q}^a \dot{q}^b \int dt$$

donde

$$\frac{\partial g_{ab}}{\partial q^a} = g_{pa} \left\{ \begin{matrix} p \\ b a \end{matrix} \right\} + g_{pb} \left\{ \begin{matrix} p \\ a b \end{matrix} \right\}$$

$$\delta S = \int (g_{pa} \left\{ \begin{matrix} p \\ b a \end{matrix} \right\} + g_{pb} \left\{ \begin{matrix} p \\ a b \end{matrix} \right\}) \delta \dot{q}^a \dot{q}^b \delta q^c + 2 g_{ab} \delta \dot{q}^a \dot{q}^b \int dt$$

$$\delta S = \int_{t_0}^t \frac{d^2}{dt^2} \left(\frac{\partial \mathcal{L}}{\partial \ddot{q}_i} \right) \delta q_i dt + \int_{t_0}^t \frac{\partial \mathcal{L}}{\partial \dot{q}_i} \delta \dot{q}_i dt + \int_{t_0}^t \frac{\partial \mathcal{L}}{\partial q_i} \delta q_i dt$$

$$\int_{t_0}^t \frac{\partial \mathcal{L}}{\partial \dot{q}_i} \delta \dot{q}_i dt \quad u = \frac{\partial \mathcal{L}}{\partial \dot{q}_i} \quad dv = \delta \dot{q}_i dt$$

$$dv = \frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{q}_i} \right) dt \quad v = \delta q_i$$

$$\int_{t_0}^t \frac{\partial \mathcal{L}}{\partial \dot{q}_i} \delta \dot{q}_i dt = \left. \frac{\partial \mathcal{L}}{\partial \dot{q}_i} \delta q_i \right|_{t_0}^t - \int_{t_0}^t \frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{q}_i} \right) \delta q_i dt$$

$$\delta S = \int_{t_0}^t \left\{ \frac{d^2}{dt^2} \left(\frac{\partial \mathcal{L}}{\partial \ddot{q}_i} \right) \delta q_i - \frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{q}_i} \right) \delta q_i + \frac{\partial \mathcal{L}}{\partial q_i} \delta q_i \right\} dt$$

$$\delta S = 0$$

$$\int_{t_0}^t \left\{ \frac{d^2}{dt^2} \left(\frac{\partial \mathcal{L}}{\partial \ddot{q}_i} \right) - \frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{q}_i} \right) + \frac{\partial \mathcal{L}}{\partial q_i} \right\} \delta q_i dt = 0$$

$$\frac{d^2}{dt^2} \left(\frac{\partial \mathcal{L}}{\partial \ddot{q}_i} \right) - \frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{q}_i} \right) + \frac{\partial \mathcal{L}}{\partial q_i} = 0$$

Pregunta 2

Enuncie el modelo sigma las Ecuaciones de Euler-Lagrange

$$L(\dot{q}, q, t) = \frac{1}{2} g_{ab}(q) \dot{q}^a \dot{q}^b$$

$$S = \int L(\dot{q}, q, t) dt \quad \delta S = \int \delta L(\dot{q}, q, t) dt = 0$$

$$\delta S = \frac{1}{2} \int (\delta g_{ab}(q) \dot{q}^a \dot{q}^b + g_{ab}(q) \delta \dot{q}^a \dot{q}^b + g_{ab}(q) \dot{q}^a \delta \dot{q}^b) dt$$

$$\delta S = \frac{1}{2} \int \left(\frac{\partial g_{ab}}{\partial q^c} \delta q^c \dot{q}^a \dot{q}^b + 2 g_{ab}(q) \dot{q}^a \delta \dot{q}^b \right) dt$$

Tarea #3 Juan Manuel López Vega

Pregunta 1

Deduce paso a paso las Ecuaciones de Euler - Lagrange para un lagrangiano que dependa de la aceleración.

$$L = L(\ddot{q}_x, \dot{q}_x, q_x, t)$$

Partiendo de obtener la acción y extremizarla

$$S = \int_{t_0}^t L(\ddot{q}_x, \dot{q}_x, q_x, t) dt \quad \text{por el principio de mínima acción}$$
$$\delta S [q_x] = 0$$

$$\delta S = \int_{t_0}^t \delta L(\ddot{q}_x, \dot{q}_x, q_x, t) dt$$

$$\delta S = \int_{t_0}^t \left[\frac{\partial L}{\partial \ddot{q}_x} \delta \ddot{q}_x + \frac{\partial L}{\partial \dot{q}_x} \delta \dot{q}_x + \frac{\partial L}{\partial q_x} \delta q_x \right] dt$$

$$\delta S = \int_{t_0}^t \frac{\partial L}{\partial \ddot{q}_x} \delta \ddot{q}_x dt + \int_{t_0}^t \frac{\partial L}{\partial \dot{q}_x} \delta \dot{q}_x dt + \int_{t_0}^t \frac{\partial L}{\partial q_x} \delta q_x dt$$

integrando por partes los 2 primeros miembros

$$\int_{t_0}^t \frac{\partial L}{\partial \ddot{q}_x} \delta \ddot{q}_x dt \quad u = \frac{\partial L}{\partial \ddot{q}_x} \quad dv = \delta \ddot{q}_x dt$$

$$du = \frac{d}{dt} \left(\frac{\partial L}{\partial \ddot{q}_x} \right) \quad v = \delta \dot{q}_x$$

$$\int_{t_0}^t \frac{\partial L}{\partial \ddot{q}_x} \delta \ddot{q}_x dt = \left. \frac{\partial L}{\partial \ddot{q}_x} \delta \dot{q}_x \right|_{t_0}^t - \int_{t_0}^t \frac{d}{dt} \left(\frac{\partial L}{\partial \ddot{q}_x} \right) \delta \dot{q}_x dt$$

$$\cancel{\delta \dot{q}_x} = \frac{d}{dt} \cancel{\delta \dot{q}_x} \Big|_{t_0}^t \quad \text{como } q_x \text{ este fijo en los extremos}$$

$$u = \frac{d}{dt} \left(\frac{\partial L}{\partial \ddot{q}_x} \right) \quad dv = \delta \dot{q}_x dt$$

$$\int_{t_0}^t \frac{\partial L}{\partial \ddot{q}_x} \delta \ddot{q}_x dt = - \int_{t_0}^t \frac{d}{dt} \left(\frac{\partial L}{\partial \ddot{q}_x} \right) \delta \dot{q}_x dt$$

$$du = \frac{d^2}{dt^2} \left(\frac{\partial L}{\partial \ddot{q}_x} \right) dt \quad v = \delta q_x$$

$$\int_{t_0}^t \frac{\partial L}{\partial \ddot{q}_x} \delta \ddot{q}_x dt = - \left. \frac{d}{dt} \left(\frac{\partial L}{\partial \ddot{q}_x} \right) \delta q_x \right|_{t_0}^t + \int_{t_0}^t \frac{d^2}{dt^2} \left(\frac{\partial L}{\partial \ddot{q}_x} \right) \delta q_x dt$$