TAREA A MECÁNICA ANÁLÍTICA. Saúl All Coodillo Rodigoca.

Pregonta 1

al Dadas las transformaciones:

$$P_1 = P_x \cos \mu - y \sin \mu$$
 $P_2 = P_y \cos \mu - y \sin \mu$

leorema: El conchele de Poisson es invariante bajo transformaciones canóniras.

Inversamente, cualquiri transformación que conserva la estructura del corchete de Poisson es canónica. Es decir, dado un conjunto de transformaciones Qi. P., con Qi = Qi + (q.p,t) y Pi = Pi | q.p.t |, donde p y q son las viejas coordinadas y momentos conjugados respectivamente, son conónicos si:

Entonces.

$$\cdot \left\{ P_1, P_2 \right\} = \frac{\partial P_1}{\partial P_2} \cdot \frac{\partial P_2}{\partial P_2} + \frac{\partial P_1}{\partial P_2} \cdot \frac{\partial P_2}{\partial P_2} - \frac{\partial P_2}{\partial P_2} - \frac{\partial P_2}{\partial P_2} \cdot \frac{\partial P_2}{\partial P_2} - \frac{\partial P_2}{\partial P_2} - \frac{\partial P_2}{\partial P_2} - \frac{\partial P_2}{\partial P_2} \cdot \frac{\partial P_2}{\partial P_2} - \frac{\partial P_2}$$

=
$$\cos m \cos m + 0 - 0 - \sin m (-\sin m) = \cos^2 m + \sin^2 m = 1$$

=
$$0 + \cos \mu \cos \mu - \sin \mu(-\sin \mu) - 0 = \cos^2 \mu + \sin^2 \mu = 1$$

 V mediante las propiedades de los avolheles de Paisson: se tiene:

$$\{q_2,q_1\}=-\{q_1,q_2\}=-0=0$$
 $\{P_2,P_1\}=-\{P_1,P_2\}=-0=0$ y Jado que $S_{ij}=\begin{pmatrix} 0&0\\0&1\end{pmatrix}$ se tiene entonces que las travejurmaciones q_1,q_2,P_1,P_2 son canónicas ya que praservan la estructura de l'corchete de Poisson.

b)
$$H = (q_1^2 + q_2^2 + p_1^2 + p_2^2)/2$$
 (Hamiltoniano original)
Dadas las transformaciones canónicas en a) obtenemos la tenatriz:

$$det(A) = cos M \begin{vmatrix} cos M sin M & 0 \\ -sin M cos M & 0 \end{vmatrix} - sin M & 0 - sin M cos M \\ 0 & 0 & cos M \end{vmatrix} - sin M & 0 - sin M & 0 & 0$$

$$\frac{\partial eL(A) = \cos \mu \left[\cos \mu \left[\cos \mu \left[-\sin \mu \left[-\sin \mu \left[-\sin \mu \left(-\cos \mu \left[-\cos \mu \left[-\sin \mu \left[\cos \mu \left[-\sin \mu \left[\cos \mu \left[-\sin \mu \left[\cos \mu \left[-\sin \mu \left[-\cos \mu \left[-\sin \mu \left[\cos \mu \left[-\sin \mu \left[-\cos \mu \left[\cos \mu \left$$

=
$$\alpha s_{M}(\alpha s_{M}^{3}\mu + s_{M}^{2}\mu \alpha s_{M}) - s_{M}\mu(-\alpha s_{M}^{2}m s_{M} - s_{M}^{3}m)$$

$$= \cos^3 \mu + \sin^2 \mu \cos^2 M + \sin^2 \mu \cos^2 \mu + \sin^4 \mu = \cos^2 \mu (\cos^2 M + \sin^2 \mu) + \sin^2 \mu (\cos^2 \mu + \sin^2 \mu)$$

=
$$crs^2 M + str^2 M = 1$$
 %. A as invertible y mediante el adcula de la malifiz adjunta se obtavo (inediante saftware) la inversa de $A: A':$

Inversas son:

$$X = q_1 \cos \mu - P_2 \sin \mu$$
 $y = q_2 \cos \mu - P_1 \sin \mu$

$$P_x = g_2 \sin \mu + P_1 \cos \mu$$

$$P_y = g_1 \sin \mu + P_2 \cos \mu$$

Hora encentrar la jurción generadora F tol que:

K=H +
$$\frac{\partial F_{i}}{\partial t}$$
, donde K es el vuevo Hamiltoniano.

l-bremes uso de las ecoaciones diferenciales: acopladas:

$$D_x = \frac{\partial F}{\partial x}$$
, $P_y = \frac{\partial F}{\partial y}$, $P_z = -\frac{\partial F}{\partial q_z}$, $P_z = -\frac{\partial F}{\partial q_z}$

* Note Exié, esto vo era vacesario.

$$H = \frac{q_1^2 + q_2^2 + P_1^2 + P_2^2}{7}$$

$$q_1^2 = (x \cos \mu + P_q \sin \mu)^2$$

= $x^2 \cos^2 \mu + 2 x \cos \mu P_q \sin \mu + P_q \sin^2 \mu$

$$q_{2}^{2} = (q_{70}s_{M} + V_{x}s_{1}^{2}n_{M})^{2} + q_{2}^{2}c_{3}s_{M}^{2} + 2q_{2}c_{3}s_{M} + V_{x}s_{1}^{2}n_{M} + V_{x}s_{1$$

$$P_{2}^{2} = (P_{y} \cos n - x \sin n)^{2} = (P_{y} \cos^{2} n - P_{y} \cos^{2} n + x^{2} \sin^{2} n)$$

$$= (P_{y} \cos n - x \sin n)^{2} + (P_{y}^{2} + Q_{z}^{2} + P_{y}^{2} + Q_{z}^{2} + Q_{z}^{2} + Q_{z}^{2})$$

$$= (P_{y} \cos n - x \sin n)^{2} + (P_{y}^{2} \cos^{2} n + P_{y}^{2} + Q_{z}^{2} + Q_{z}$$

c) De las ervadones:
$$0 = \frac{\partial H}{\partial P_1}$$
, $\dot{P}_1 = -\frac{\partial H}{\partial P_2}$ can la violarization $y = P_y = 0$

Se tiene $1 = \frac{x^2 + P_x^2}{2}$, por le tante: se tienen les signientes ecuariones de

movimento:

$$\dot{p}_1 = \frac{\partial H}{\partial x} = -\frac{\partial \left(x^2 + D_x^2\right)}{\partial x^2} = \left(x + \frac{D_x^2}{2}\right), \quad \dot{p}_2 = 0$$

$$4 = \frac{\partial H}{\partial P_x} = \frac{\partial \left(x^2 + P_x^2\right)}{\partial Q_x} = \frac{x^2}{2} + P_x$$

8.
$$\dot{p}_1 = -x - \frac{p_2^2}{2}$$
 $\dot{q}_1 = \frac{x^2}{2} + p_x$

Progonta 2.

a) Momentos de Inercia y gles principales:

Considerando la definición de tensor de momento de inevola:

$$I_{ij} = \int_{V} \rho(\bar{r}') \left(S_{ij} \sum_{K} x_{ik}^{2} - x_{i} x_{j} \right) dV \quad \text{an } dv = dx_{i} dx_{2} dx_{3}$$

$$dv = dx_{1} dx_{2} dx_{3}$$

Calculamos los momentos de inercia del disco:

 $I_{11} = \int_{V} \rho \left(x_{1}^{2} + x_{2}^{2} + x_{3}^{3} - x_{1}^{2} \right) dv = \rho \int_{V} \left(y_{1}^{2} + z_{2}^{2} \right) dx dy dz, dado que el disco el debado podemos despreciav la contribución de la integral en el eje z de modo que:$

•
$$I_{11} = P \int_{V} y^{2} dxdy$$
, en coordenadas polaves: $I_{11} = P \int_{V} r^{2} \sin^{2}\theta \, v \, dv \, d\theta$

$$\begin{split} & \left[\int_{0}^{1} - \rho \int_{0}^{A} \int_{0}^{2\pi} r^{3} \sin^{2}\theta \, d\theta \, dr = \rho \int_{0}^{A} r^{3} \left(\frac{1}{2} \theta - \frac{1}{4} \sin^{2}\theta \right) \right]_{0}^{2\pi} \, dv = \rho \int_{0}^{A} r^{3} \left(\frac{1}{2} 2\pi - \frac{1}{4} \sin^{2}\theta \right) \, dr \\ & = \rho \int_{0}^{A} r^{3} \pi \, dr = \rho \pi \left(\frac{1}{4} r^{4} \right)_{0}^{A} = \rho \pi \frac{1}{4} A^{4} = \frac{M}{\pi A^{2}} \frac{1}{4} A^{4} = \frac{MA^{2}}{4} \end{split}$$

$$\int_{22}^{2\pi} \left[\int_{V}^{2} dx dy \right] = \rho \int_{V}^{2} \cos^{2}\theta dy d\theta = \rho \int_{V}^{3} \cos^{2}\theta dy d\theta = \rho \int_{0}^{A} \int_{0}^{2\pi} v^{3} \cos^{2}\theta d\theta dy$$

$$= \rho \int_{0}^{A} r^{3} \left(\frac{1}{2} \theta + \frac{1}{4} \sin(4/2\theta) \right)_{0}^{2\pi} dr = \rho \int_{0}^{A} r^{3} \left[\frac{1}{2} 2\pi + \frac{1}{4} \sin(4\pi) - 0 - \frac{1}{4} \sin(0) \right] dr$$

$$= \rho \int_{0}^{A} r^{3} \pi = \rho \pi \frac{1}{4} A^{4} = \frac{MA^{2}}{4}$$

$$I_{33} = \rho \int_{V} (x^{2} + y^{2}) dx dy = \rho \int_{V} (y^{2} \cos^{2}\theta + \sin^{2}\theta y^{2}) r dy d\theta = \rho \int_{V} (y^{3} dx d\theta - \rho) \int_{0}^{A} \int_{0}^{2\pi} y^{3} d\theta dy$$

$$I_{38} = \rho \int_{0}^{A} r^{3} \theta \int_{0}^{2\pi} dv = \rho \int_{0}^{A} r^{3} 2\pi dr = \rho \frac{1}{4} r^{4} 2\pi \int_{0}^{A} = \rho \frac{1}{4} A^{4} 2\pi = \rho \frac{1}{2} A^{4} \pi = \frac{M}{\pi A^{2}} \frac{1}{2} A^{4} \pi^{2}$$

$$... \Lambda. I_{33} - \frac{1}{2}MA^2$$

Por la tanta las momentas de inercia son:

$$I_{11} = I_{22} = \frac{1}{4}MA^2$$
 $I_{33} = \frac{1}{2}MA^2$

Calculande les componentes del tensor de momente de inercia I r vestantes son:

•
$$I_{12} = \rho \int_{V} -x_1 x_2 dv = -\rho \int_{V} x_1 dx dy = -\rho \int_{V} r \cos\theta r \sin\theta r dr d\theta = -\rho \int_{0}^{A} \int_{0}^{2\pi} r^3 \cos\theta \sin\theta d\theta dr$$

$$= \int_{32}^{2} \left[-\rho \right]_{V}^{2} dx dy dz = -\rho \int_{V}^{2} y dx dy = -\rho \int_{V}^{A} \left[r^{2} \sin \theta d\theta dy \right]_{V}^{2} dx dy dz = -\rho \int_{V}^{A} \left[r^{2} \sin \theta d\theta dy \right]_{V}^{2} dx dy dz = -\rho \int_{V}^{A} \left[r^{2} \sin \theta d\theta dy \right]_{V}^{2} dx dy dz = -\rho \int_{V}^{A} \left[r^{2} \sin \theta d\theta dy \right]_{V}^{2} dx dy dz = -\rho \int_{V}^{A} \left[r^{2} \sin \theta d\theta dy \right]_{V}^{2} dx dy dz = -\rho \int_{V}^{A} \left[r^{2} \sin \theta d\theta dy \right]_{V}^{2} dx dy dz = -\rho \int_{V}^{A} \left[r^{2} \sin \theta d\theta dy \right]_{V}^{2} dx dy dz = -\rho \int_{V}^{A} \left[r^{2} \sin \theta d\theta dy \right]_{V}^{2} dx dy dz = -\rho \int_{V}^{A} \left[r^{2} \sin \theta d\theta dy \right]_{V}^{2} dx dy dz = -\rho \int_{V}^{A} \left[r^{2} \sin \theta d\theta dy \right]_{V}^{2} dx dy dz = -\rho \int_{V}^{A} \left[r^{2} \sin \theta d\theta dy \right]_{V}^{2} dx dy dz = -\rho \int_{V}^{A} \left[r^{2} \sin \theta d\theta dy \right]_{V}^{2} dx dy dz = -\rho \int_{V}^{A} \left[r^{2} \sin \theta d\theta dy \right]_{V}^{2} dx dy dy dz = -\rho \int_{V}^{A} \left[r^{2} \sin \theta d\theta dy \right]_{V}^{2} dx dy dz = -\rho \int_{V}^{A} \left[r^{2} \sin \theta d\theta dy \right]_{V}^{2} dx dy dz = -\rho \int_{V}^{A} \left[r^{2} \sin \theta d\theta dy \right]_{V}^{2} dx dy dz = -\rho \int_{V}^{A} \left[r^{2} \sin \theta d\theta dy \right]_{V}^{2} dx dy dz = -\rho \int_{V}^{A} \left[r^{2} \sin \theta d\theta dy \right]_{V}^{2} dx dy dz = -\rho \int_{V}^{A} \left[r^{2} \sin \theta d\theta dy \right]_{V}^{2} dx dy dz = -\rho \int_{V}^{A} \left[r^{2} \sin \theta d\theta dy \right]_{V}^{2} dx dy dz = -\rho \int_{V}^{A} \left[r^{2} \sin \theta d\theta dy \right]_{V}^{2} dx dy dz = -\rho \int_{V}^{A} \left[r^{2} \sin \theta d\theta dy \right]_{V}^{2} dx dy dz = -\rho \int_{V}^{A} \left[r^{2} \sin \theta d\theta dy \right]_{V}^{2} dx dy dz = -\rho \int_{V}^{A} \left[r^{2} \sin \theta d\theta dy \right]_{V}^{2} dx dy dz = -\rho \int_{V}^{A} \left[r^{2} \sin \theta d\theta dy \right]_{V}^{2} dx dy dz dy dz = -\rho \int_{V}^{A} \left[r^{2} \sin \theta d\theta dy \right]_{V}^{2} dx dy dz dy d$$

$$= -\rho \int_{0}^{A} r^{2} \cos \theta \int_{0}^{2\pi} dr = 0 \quad \text{s.} \quad I_{32} = I_{23} = 0$$

Por la tanta el tensor de mercia es:

$$I_{ij} = \begin{pmatrix} \frac{1}{4}MA^{2} & 0 & 0 \\ 0 & \frac{1}{4}MA^{2} & 0 \\ 0 & 0 & \frac{1}{2}MA^{2} \end{pmatrix} = \frac{1}{4}MA^{2} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 2 \end{pmatrix}$$

4 las eigenvalores de Ii; :

$$dd \begin{pmatrix} 1-2 & 0 & 0 \\ 0 & 1-2 & 0 \\ 0 & 0 & 2-2 \end{pmatrix} = (1-2)(1-2)(2-2) \implies \lambda_1 = 1, \ \lambda_2 = 1, \ \gamma_3 = 2$$

Evaluando la matriz en 2, = 2= 1.

$$\begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix} \Rightarrow Z=0 \Rightarrow \bar{X}=\begin{pmatrix} x \\ y \\ 0 \end{pmatrix}=x\begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}+y\begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix}, evaluanda en $I_3=Z$$$

$$\begin{pmatrix} -1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 0 \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{pmatrix} \Rightarrow \begin{array}{c} x = 0 \\ y = 0 \Rightarrow 0 \end{array} \quad \bar{x} = \begin{pmatrix} 0 \\ 0 \\ z \end{pmatrix} = z \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

ejes principales, as decir, les ejes principales son les ejes del sedema roordenado del disco.

1) Vector de momento angular.

El vertor de velocidad angular desde el sistema coordenado del disco está dada por:

$$\overline{\widetilde{w}} = \begin{pmatrix} \widetilde{w} \cos(\pi/2 - w) \\ \widetilde{w} \cos(\pi/2) \end{pmatrix} = \begin{pmatrix} \widetilde{w} \cos(\pi/2 - w) \\ 0 \\ w \cos(w) \end{pmatrix}, \text{ dende } ||\widetilde{w}|| = \omega, \quad \overline{\widetilde{w}} = \begin{pmatrix} \widetilde{w}_{r} \\ \widetilde{w}_{2} \\ w_{3} \end{pmatrix} = \begin{pmatrix} \widetilde{w}_{1} \\ \widetilde{w}_{2} \\ w_{3} \end{pmatrix}$$

Dado que
$$\overline{L} = \overline{I}\overline{w} \Rightarrow \overline{L}_a = \overline{I}_{ab}w_b = \overline{L}_b \overline{I}_{ab}w_b$$
, por lo tanto:

$$L_1 = L_{11}\omega_1 + I_{12}\omega_2 + I_{13}\omega_3 = I_{11}\omega_1 = \frac{1}{4}MA^2\omega \cos(\pi/2 - d)$$

$$L_{2} = \int_{21}^{7} w_{1} + I_{22} w_{2} + I_{23} w_{3} = I_{22} w_{2} = \frac{1}{4} MA^{2} w_{0} s(\pi/2) = 0$$

$$L_3 = I_{33}\omega_3 = \frac{1}{2}MA^2\omega \cos(\omega)$$

Por la tanta, el vertor de momenta angular es:

$$\begin{bmatrix}
\frac{1}{4}MA^{2}w\cos(\pi/2-\alpha) \\
0 \\
\frac{1}{2}MA^{2}w\cos(\alpha)
\end{bmatrix} = \frac{1}{4}MA^{2}\begin{pmatrix} w\cos(\pi/2-\alpha) \\
0 \\
2w\cos(\alpha) \end{pmatrix} = \frac{1}{4}MA^{2}w\begin{pmatrix} \sin(\alpha) \\
0 \\
2\cos(\alpha) \end{pmatrix}$$

Cuya magnitud a:

$$\begin{aligned} || \vec{L} || &= \sqrt{\left[\frac{1}{4} M A^2 w \cos(\pi/2 - \alpha)\right]^2 + \left(\frac{1}{2} M A^2 w \cos \alpha\right)^2} \\ &= \sqrt{\left(\frac{1}{4} M A^2 w\right)^2 \sin^2(\alpha) + \left(\frac{1}{4} M A^2 w\right)^2 (2 \cos \alpha)^2} \\ &= \frac{1}{4} M A^2 w \sqrt{\sin^2(\alpha) + 2^2 \cos^2(\alpha)} = \frac{1}{4} M A^2 w \sqrt{\sin^2(\alpha) + 4 \cos^2(\alpha)} = || \vec{L} || \end{aligned}$$

c) Magnitud y dirección de la torca velativa al srotema de referencia del discoc

$$51 \quad \frac{\partial \overline{L}}{\partial t} = \overrightarrow{\gamma} \quad \Rightarrow \overrightarrow{\gamma} = \frac{\partial \overline{L}}{\partial t} = \frac{\partial}{\partial t} \begin{pmatrix} \frac{1}{4} M A^2 w \sin d \\ 0 \\ \frac{1}{2} M A^2 w (\cos id) \end{pmatrix} = \frac{1}{4} M A^2 \dot{w} \begin{pmatrix} \sin (a) \\ 0 \\ 2 \cos (d) \end{pmatrix}$$

$$9 \quad ||\overrightarrow{\gamma}|| = \frac{1}{4} M A^2 \dot{w} \sqrt{3 \ln^2(d) + 4 \cos^2(d)}$$