

Lecture-08

Gauss's Law:

$$\oint \mathbf{E} \cdot d\mathbf{s} = \frac{Q_{\text{enc}}}{\epsilon_0}$$

Integral form.

* Electric flux passing through any closed surface enclosing the charge is $\frac{Q}{\epsilon_0}$.

Differential form \Rightarrow

$$\oint \mathbf{E} \cdot d\mathbf{s} = \frac{Q}{\epsilon_0}$$

$$\Rightarrow \int_V (\nabla \cdot \mathbf{E}) dV = \frac{1}{\epsilon_0} \int_V \rho dV$$

$$\Rightarrow \boxed{\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0}}$$

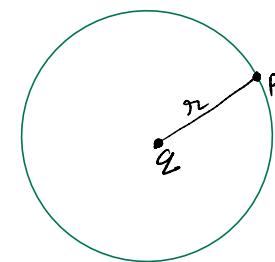
1st Maxwell's eqn.

Applications of Gauss's Law:

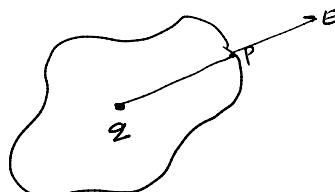
(a) \Rightarrow Point charge \Rightarrow

$$\oint \mathbf{E} \cdot d\mathbf{s} = \frac{Q}{\epsilon_0}$$

$$E \cdot 4\pi r^2 = \frac{q}{\epsilon_0} \Rightarrow E = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2}$$



Comment \Rightarrow * Gauss Law is always true but not always useful.



\Rightarrow * $E \cdot d\mathbf{s}$ are in different directions & E is not cont over the surface.

* Symmetry is crucial to the application of Gauss's law.

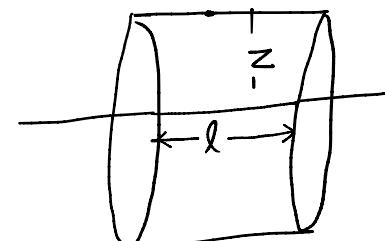
(ii) \Rightarrow Infinite Line charge \Rightarrow

* Field at distance z from the wire,

$$\oint \mathbf{E} \cdot d\mathbf{s} = \frac{Q}{\epsilon_0}$$

$$E \times 2\pi z l = \frac{\lambda l}{\epsilon_0}$$

$$E = \frac{\lambda}{2\pi\epsilon_0 z} \quad (\text{same as previously})$$



* Line charge density = λ

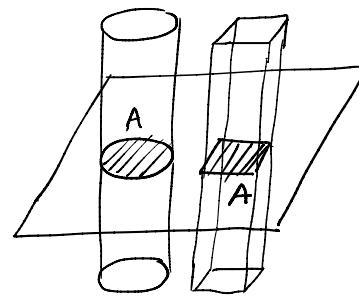
(ii) \Rightarrow Infinite Plane \Rightarrow

Surface charge density $= \sigma$

$$\int E \cdot dS = \frac{q}{\epsilon_0}$$

$$E \cdot 2A = \frac{\sigma A}{\epsilon_0}$$

$$E = \frac{\sigma}{2\epsilon_0} \quad \underline{\underline{}}$$



(iv) \Rightarrow Uniformly charged sphere \Rightarrow

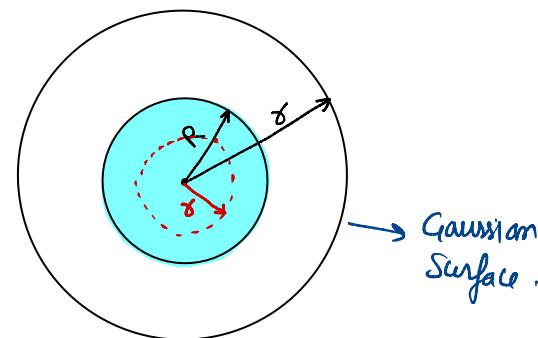
Radius R , Total charge $= q$

(a) \Rightarrow Outside \Rightarrow

$$\int E \cdot dS = \frac{q}{\epsilon_0}$$

$$E 4\pi r^2 = \frac{q}{\epsilon_0}$$

$$E = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \quad \underline{\underline{}}$$



(b) \Rightarrow Inside \Rightarrow

$$\int E \cdot dS = \frac{q}{\epsilon_0}$$

$$E 4\pi r^2 = \frac{q}{\epsilon_0} \times \frac{\frac{4}{3}\pi r^3}{R^3}$$

$$E = \frac{1}{4\pi\epsilon_0} \frac{qr}{R^3} \quad \underline{\underline{}}$$

Q \Rightarrow Suppose the electric field in some region is $E = k_2 r^3 \hat{r}$, in spherical coordinates.

(a) \Rightarrow Find the charge density ρ

(b) \Rightarrow Find the total charge contained in a sphere of radius R .

Ans (a) $\Rightarrow \nabla \cdot E = \frac{\rho}{\epsilon_0}$

$$\nabla \cdot E = \frac{1}{r^2 \sin\theta} \frac{\partial}{\partial r} (k_2 r^3 \cdot r \sin\theta) = \frac{\rho}{\epsilon_0}$$

$$= \frac{1}{r^2 \sin\theta} k_2 \sin\theta \cdot 5r^4 = \frac{\rho}{\epsilon_0}$$

$$\Rightarrow 5k_2 r^2 = \frac{\rho}{\epsilon_0} \Rightarrow \rho = 5\epsilon_0 k_2 r^2 \quad \underline{\underline{}}$$

$$\nabla \cdot A = \frac{1}{r_1 r_2 r_3} \left[\frac{\partial}{\partial r_1} (A_1 r_2 r_3) + \dots \right]$$

$$(b) \Rightarrow \text{Total charge } Q = \int \rho dV$$

$$\Rightarrow \int_0^R 5\epsilon_0 k_2 r^2 4\pi r^2 dr$$

$$\Rightarrow 20\pi \epsilon_0 k_2 \int_0^R r^4 dr$$

$$\Rightarrow 4\pi \epsilon_0 k_2 R^5 \cancel{\Delta}$$

Homework

Q.1 \Rightarrow Find the electric field inside a sphere which carries a charge density $\rho = k_2 r$.

$$\underline{\text{Ans}} \Rightarrow E = \frac{1}{4\pi\epsilon_0} k_2 r^2 \hat{r}.$$

Homework

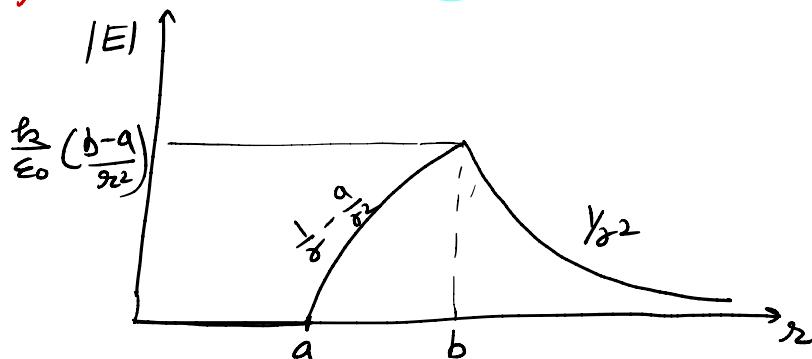
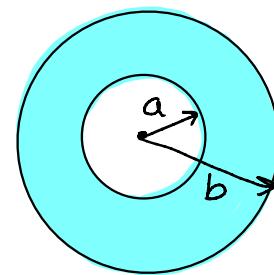
Q.2 \Rightarrow A hollow spherical shell carries charge density $\rho = \frac{k_2}{r^2}$ in the region $a \leq r \leq b$. Find the electric field in three regions:

$$(i) \Rightarrow r < a \quad E = 0$$

$$(ii) \Rightarrow a < r < b \quad E = \frac{k_2}{\epsilon_0} \left(\frac{r-a}{r^2} \right) \hat{r}$$

$$(iii) \Rightarrow r > b \quad E = \frac{k_2}{\epsilon_0} \left(\frac{b-a}{r^2} \right) \hat{r}$$

& Plot $|E|$ vs r .



Electric Potential \Rightarrow

* In electrostatic

$$\vec{E} = -y \hat{x} + x \hat{y}$$

or

$$\vec{E} = x \hat{y}$$

is not possible with any set of charge, their size or position.

\Rightarrow In electrostatic, the curl of Electric field is always zero.

Homework:

Q \Rightarrow One of these is an impossible electrostatic field. Which one?

$$(a) \Rightarrow \bar{E}_1 = k(x\hat{i} + 2yz\hat{j} + 3xz\hat{z}), \quad \nabla \times E_1 \neq 0$$

$$(b) \Rightarrow E_2 = k(y^2\hat{x} + (2xy + z^2)\hat{j} + 2yz\hat{z}) \quad \nabla \times E_2 = 0$$

$\therefore E_2$ is possible not E_1 .

#

$$\nabla \times E = 0$$

then $E = -\nabla V$

\downarrow
Electric Pot.

if $\nabla \times \bar{A} = 0$
then there is always
a φ such that

$$\bar{A} = -\nabla \varphi$$

* If we know V , we can easily determine E .

* E is a vector quantity (3 components) & V is scalar [1 component]

§ How one funⁿ can contain the information of three 'independent' ofⁿ?

Ans: The three components are not independent.

$$\nabla \times E = 0$$

$$\Rightarrow \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ E_x & E_y & E_z \end{vmatrix} = 0$$

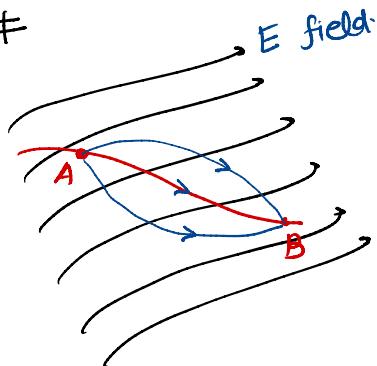
$$\Rightarrow \frac{\partial E_x}{\partial y} = \frac{\partial E_y}{\partial x},$$

$$\frac{\partial E_z}{\partial y} = \frac{\partial E_y}{\partial z}$$

$$\frac{\partial E_x}{\partial z} = \frac{\partial E_z}{\partial x}$$

* The advantage of pot. formalism is that it reduces the vector problem to a scalar one.

#



$$W = -\bar{F} \cdot d\ell$$

$$W = -q \int_A^B \mathbf{E} \cdot d\ell$$

* Pot. V is defined by Pot Energy per unit charge.

$$V = \frac{W}{q}$$

$$\therefore V = - \int_{O'}^R \mathbf{E} \cdot d\ell$$

* $\nabla \times \mathbf{E} = 0$

Stokes' theorem

$$\oint \mathbf{E} \cdot d\ell = \int (\nabla \times \mathbf{E}) \cdot ds$$

$$\Rightarrow \oint \mathbf{E} \cdot d\ell = 0$$

* V_{AB} is independent of the path taken.

Ex: Find the pot. inside & outside a uniformly charged solid sphere whose radius is R & whose total charge is q . Use infinity as your reference point.

Ans: $V(r) = - \int_{\infty}^r \mathbf{E} \cdot d\ell$

① Outside $\Rightarrow r > R$

$$V(r) = - \int_{\infty}^r \mathbf{E}_0 \cdot d\ell = - \int_{\infty}^r \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} dr$$

$$= \frac{1}{4\pi\epsilon_0} \frac{q}{r} \Big|_{\infty}^r = \frac{1}{4\pi\epsilon_0} \frac{q}{r} \quad \underline{\underline{}}$$

Outside	$\bar{E} = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{r}$
Inside	$\bar{E} = \frac{1}{4\pi\epsilon_0} \frac{qr}{R^3} \hat{r}$

② Inside $\Rightarrow r < R$

$$\begin{aligned} V(r) &= - \int_{\infty}^R E_0 \cdot dl - \int_R^r E_{in} \cdot dl \\ &= - \int_{\infty}^R \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} dr - \int_R^r \frac{1}{4\pi\epsilon_0} \frac{qr}{R^3} dr \\ &= \frac{1}{4\pi\epsilon_0} q \left[\frac{1}{R} - \frac{1}{R^3} \left(\frac{r^2 - R^2}{2} \right) \right] = \frac{1}{4\pi\epsilon_0} \frac{q}{2R} \left(3 - \frac{r^2}{R^2} \right) \end{aligned}$$

Ex \Rightarrow Find the pot. at a distance z from an infinitely long straight wire that carries a uniform line charge λ .

$$\text{Ans} \Rightarrow E = \frac{\lambda}{2\pi\epsilon_0 z}$$

$$V = - \int_a^z \frac{1}{2\pi\epsilon_0 z} dz = - \frac{1}{2\pi\epsilon_0} \ln z \Big|_a^z$$

$$V = - \frac{1}{2\pi\epsilon_0} \ln \frac{z}{a}$$

