

# 第三次作业答案

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## 目录

### 第四题

解: (1)

$$N = \int_{-\infty}^{\infty} |\psi(x)|^2 dx = \int_{-\infty}^{\infty} e^{-2|x|} dx = \int_{-\infty}^0 e^{2x} dx + \int_0^{\infty} e^{-2x} dx = 1$$

得到:

$$\psi'(x) = \frac{1}{\sqrt{N}} \psi(x) = \psi(x) = e^{-|x|}$$
$$\phi(p) = \frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{\infty} e^{-|x|} e^{-\frac{i}{\hbar} px} dx = \frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^0 e^{(1-\frac{ip}{\hbar})x} dx + \frac{1}{\sqrt{2\pi\hbar}} \int_0^{\infty} e^{(-1-\frac{ip}{\hbar})x} dx = \frac{1}{\sqrt{2\pi\hbar}} \frac{2\hbar^2}{\hbar^2 + p^2}$$

### 第五题

解:

$$N = \int_{-\infty}^{\infty} |\psi(x)|^2 dx = \int_{-\infty}^{\infty} e^{-2\frac{x^2}{\sigma^2}} dx = \sigma \sqrt{\frac{\pi}{2}}$$

得到归一化后的波函数:

$$\psi(x) = \frac{1}{\sqrt{\sigma}} \left( \frac{2}{\pi} \right)^{\frac{1}{4}} e^{-\frac{x^2}{\sigma^2}}$$
$$\langle x \rangle = \int_{-\infty}^{\infty} \psi^*(x) x \psi(x) dx = \frac{1}{\sigma} \sqrt{\frac{2}{\pi}} \int_{-\infty}^{\infty} x e^{-\frac{2x^2}{\sigma^2}} dx = 0$$
$$\langle p \rangle = \int_{-\infty}^{\infty} \psi^*(x) (-i\hbar) \frac{d}{dx} \psi(x) dx = -\frac{2\sqrt{2}}{\pi\sigma^3} \int_{-\infty}^{\infty} x e^{-\frac{2x^2}{\sigma^2}} dx = 0$$

注: 奇函数积分性质。

$$\langle T \rangle = \int_{-\infty}^{\infty} \psi^*(x) \frac{-\hbar^2}{2m} \frac{d^2}{dx^2} \psi(x) dx = \frac{1}{\sigma} \sqrt{\frac{2}{\pi}} \frac{\hbar^2}{m\sigma^2} \int_{-\infty}^{\infty} \left( 1 - \frac{2x^2}{\sigma^2} \right) e^{-\frac{2x^2}{\sigma^2}} dx = \frac{\hbar^2}{2m\sigma^2}$$

## 第七题

解: 由不确定性关系:

$$\Delta x \Delta p \geq \frac{\hbar}{2}$$

$$\Delta p = \sqrt{\langle p^2 \rangle - \langle p \rangle^2}$$

因而有:

$$\langle T \rangle = \frac{\langle p^2 \rangle}{2m} \geq \frac{\Delta p^2}{2m} \geq \frac{\hbar^2}{8m\Delta x^2}$$

接下来带入计算即可。

## 第八题

解: 假设本征值为复数, 记作  $E = E_0 + i\Gamma$ , 则定态波函数可以写作:

$$\psi(\vec{x}, t) = \psi(\vec{x}, 0)e^{-i\frac{E_0 + i\Gamma}{\hbar}t}$$

由此我们可以验证, 在初始时刻归一化的波函数, 在  $t$  时刻是否依然满足归一化条件:

$$\int d^n \vec{x} |\psi(\vec{x}, t)|^2 = \int d^n \vec{x} |\psi(\vec{x}, 0)|^2 e^{\frac{2\Gamma t}{\hbar}} = e^{\frac{2\Gamma t}{\hbar}} = 1$$

可见, 为使得归一化条件一直成立, 就要求定态波函数的对应的本征值的虚部为  $\Gamma = 0$ , 即本征值为一个实数。

## 第十题

解:

$$-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \psi(x) + V(x)\psi(x) = E\psi(x)$$

给方程两边同时乘以  $\psi^*(x)$ , 并对整个方程在全空间积分, 得到:

$$\begin{aligned} E &= \int dx \psi^*(x) \left( -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + V(x) \right) \psi(x) = \\ &= -\frac{\hbar^2}{2m} \psi^*(x) \frac{\partial}{\partial x} \psi(x) \Big|_{\partial x} + \frac{\hbar^2}{2m} \int dx \frac{\partial}{\partial x} \psi^*(x) \frac{\partial}{\partial x} \psi(x) + \int dx |\psi(x)|^2 V(x) = \\ &= \frac{\hbar^2}{2m} \int dx \frac{\partial}{\partial x} \psi^*(x) \frac{\partial}{\partial x} \psi(x) + \int dx |\psi(x)|^2 V(x) \\ &\geq \int |\psi(x)|^2 V(x) \geq V_{\min} \end{aligned}$$

得证。

## 第十二题

a.

$$N = \int (\phi_1^*(x) - i\phi_2^*(x)) (\phi_1(x) + i\phi_2(x)) dx = \int dx |\phi_1(x)|^2 + |\phi_2(x)|^2 = 2$$

得到:

$$\psi(x, 0) = \frac{1}{\sqrt{2}} (\phi_1(x) + i\phi_2(x))$$

b.

$$\psi(x, t) = \frac{1}{\sqrt{2}} \left( \phi_1(x) e^{-i\frac{E_1}{\hbar}t} + i\phi_2(x) e^{-i\frac{E_2}{\hbar}t} \right)$$

$$E_n = \frac{(n\pi\hbar)^2}{2ma^2}$$

$$|\psi(x, t)|^2 = \frac{1}{2} \left( |\phi_1(x)|^2 + |\phi_2(x)|^2 + 2\phi_1(x)\phi_2(x) \sin(\omega t) \right)$$

$$\omega = \frac{E_2 - E_1}{\hbar} = \frac{3\pi^2\hbar}{2ma^2}$$

c.

$$\begin{aligned} \langle x \rangle &= \frac{1}{2} \int_0^a \left( |\phi_1(x)|^2 + |\phi_2(x)|^2 + 2\phi_1(x)\phi_2(x) \sin(\omega t) \right) x dx = \\ \frac{\langle x \rangle_1 + \langle x \rangle_2}{2} + \int_0^a \phi_1(x)\phi_2(x) \sin(\omega t) x dx &= \frac{a}{2} + \frac{2}{a} \sin(\omega t) \int_0^a \sin\left(\frac{\pi x}{a}\right) \sin\left(\frac{2\pi x}{a}\right) dx = \\ \frac{a}{2} - \frac{16a}{9\pi^2} \sin(\omega t) \\ \langle x^2 \rangle &= \frac{1}{2} \int_0^a \left( |\phi_1(x)|^2 + |\phi_2(x)|^2 + 2\phi_1(x)\phi_2(x) \sin(\omega t) \right) x^2 dx = \\ \frac{\langle x^2 \rangle_1 + \langle x^2 \rangle_2}{2} + \int_0^a \phi_1(x)\phi_2(x) \sin(\omega t) x^2 dx &= \\ \frac{a^2}{3} - \frac{5a^2}{16\pi^2} + \int_0^a \phi_1(x)\phi_2(x) \sin(\omega t) x^2 dx &= \end{aligned}$$

$$\begin{aligned}
& \frac{a^2}{3} - \frac{5a^2}{16\pi^2} - \frac{16D^2 \sin(\omega t)}{9\pi^2} \\
\langle p \rangle &= \frac{1}{2} \int_0^a \left( \phi_1(x) e^{i\frac{E_1}{\hbar}t} - i\phi_2(x) e^{i\frac{E_2}{\hbar}t} \right) p \left( \phi_1(x) e^{-i\frac{E_1}{\hbar}t} + i\phi_2(x) e^{-i\frac{E_2}{\hbar}t} \right) dx = \\
& \quad \frac{\langle p \rangle_1 + \langle p \rangle_2}{2} + \frac{\hbar}{2} \phi_1 \phi_2 e^{-i\omega t} \Big|_0^a - \hbar \int_0^a \phi_2 \frac{\partial}{\partial x} \phi_1 \cos(\omega t) dx = \\
& \quad - \frac{8\hbar \cos(\omega t)}{3D} \\
\langle p^2 \rangle &= \frac{1}{2} \int_0^a \left( \phi_1(x) e^{i\frac{E_1}{\hbar}t} - i\phi_2(x) e^{i\frac{E_2}{\hbar}t} \right) p^2 \left( \phi_1(x) e^{-i\frac{E_1}{\hbar}t} + i\phi_2(x) e^{-i\frac{E_2}{\hbar}t} \right) dx = \\
& \quad \frac{\langle p^2 \rangle_1 + \langle p^2 \rangle_2}{2} + i\frac{\hbar^2}{2} \left( \int_0^a \phi_2 \frac{\partial^2}{\partial x^2} \phi_1 e^{i\omega t} - \int_0^a \phi_1 \frac{\partial^2}{\partial x^2} \phi_2 e^{-i\omega t} \right) = \\
& \quad \frac{\langle p^2 \rangle_1 + \langle p^2 \rangle_2}{2} = \frac{5\hbar^2 \pi^2}{2D^2}
\end{aligned}$$

d.

$$\begin{aligned}
\langle E \rangle &= \frac{1}{2} \int_0^a \left( \phi_1(x) e^{i\frac{E_1}{\hbar}t} - i\phi_2(x) e^{i\frac{E_2}{\hbar}t} \right) H \left( \phi_1(x) e^{-i\frac{E_1}{\hbar}t} + i\phi_2(x) e^{-i\frac{E_2}{\hbar}t} \right) dx = \\
& \quad \frac{1}{2} \int_0^a \left( E_1 |\phi_1(x)|^2 + E_2 |\phi_2(x)|^2 + iE_1 \phi_1 \phi_2 e^{i\omega t} - iE_2 \phi_1 \phi_2 e^{-i\omega t} \right) dx = \\
& \quad \frac{E_1 + E_2}{2} = \frac{5\pi^2 \hbar^2}{4mD^2} \\
p(E_1) &= \left| \int_0^a \phi_1(x) \psi(x, t) dx \right|^2 = \frac{1}{2} \\
p(E_2) &= \left| \int_0^a \phi_2(x) \psi(x, t) dx \right|^2 = \frac{1}{2}
\end{aligned}$$

测得能量为  $E_1 = \frac{(\pi\hbar)^2}{2ma^2}$  及  $E_2 = \frac{(2\pi\hbar)^2}{2ma^2}$  的概率各为  $\frac{1}{2}$