

Brief Article

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This note gives some equations from the paper introducing the HLLD method by Miyoshi and Kusano [1].

Ideal MHD eigenvalues

$$\lambda_{2,6} = u \mp c_A, \quad \lambda_{1,7} = u \mp c_f, \quad \lambda_{3,5} = u \mp c_s, \quad \lambda_4 = u$$

where the Alfvén speed c_A , fast wave speed c_f and slow wave speed c_s are given by

$$c_A = \frac{|B_x|}{\sqrt{\rho}}, \quad c_{f,s} = \left[\frac{\gamma p + |\mathbf{B}|^2 \pm \sqrt{(\gamma p + |\mathbf{B}|^2)^2 - 4\gamma p B_x^2}}{2\rho} \right]^{1/2}$$

HLLD approximation. Entropy wave speed

$$S_M = \frac{(S_R - u_R)\rho_R u_R - (S_L - u_L)\rho_L u_L - p_{T_R} + p_{T_L}}{(S_R - u_R)\rho_R - (S_L - u_L)\rho_L} \quad (1)$$

We have $p_{T_L}^* = p_{T_L}^{**} = p_{T_R}^{**} = p_{T_R}^* = p_T^*$

$$p_T^* = \frac{(S_R - u_R)\rho_R p_{T_L} - (S_L - u_L)\rho_L p_{T_R} + \rho_L \rho_R (S_R - u_R)(S_L - u_L)(u_R - u_L)}{(S_R - u_R)\rho_R - (S_L - u_L)\rho_L} \quad (2)$$

With $\alpha = L, R$,

$$v_\alpha^* = v_\alpha - B_x B_{y\alpha} \frac{S_M - u_\alpha}{\rho_\alpha (S_\alpha - u_\alpha)(S_\alpha - S_M) - B_x^2} \quad (3)$$

$$B_{y\alpha}^* = B_{y\alpha} \frac{\rho_\alpha (S_\alpha - u_\alpha)^2 - B_x^2}{\rho_\alpha (S_\alpha - u_\alpha)(S_\alpha - S_M) - B_x^2} \quad (4)$$

$$w_\alpha^* = w_\alpha - B_x B_{z\alpha} \frac{S_M - u_\alpha}{\rho_\alpha (S_\alpha - u_\alpha)(S_\alpha - S_M) - B_x^2} \quad (5)$$

$$B_{z\alpha}^* = B_{z\alpha} \frac{\rho_\alpha (S_\alpha - u_\alpha)^2 - B_x^2}{\rho_\alpha (S_\alpha - u_\alpha)(S_\alpha - S_M) - B_x^2} \quad (6)$$

$$e_\alpha^* = \frac{(S_\alpha - u_\alpha)e_\alpha - p_{T\alpha}u_\alpha + p_T^*S_M + B_x(\mathbf{v}_\alpha \cdot \mathbf{B}_\alpha - \mathbf{v}_\alpha^* \cdot \mathbf{B}_\alpha^*)}{S_\alpha - S_M} \quad (7)$$

We have $\rho_\alpha^{**} = \rho_\alpha^*$ and $p_{T\alpha}^{**} = p_{T\alpha}^*$.

Propagation speed of Alfvén waves in the intermediate states:

$$S_L^* = S_M - \frac{|B_x|}{\sqrt{\rho_L^*}} \quad (8)$$

$$S_R^* = S_M + \frac{|B_x|}{\sqrt{\rho_R^*}} \quad (9)$$

Equalities:

$$v_L^{**} = v_R^{**} \equiv v^{**} \quad (10)$$

$$w_L^{**} = w_R^{**} \equiv w^{**} \quad (11)$$

$$B_{yL}^{**} = B_{yR}^{**} \equiv B_y^{**} \quad (12)$$

$$B_{zL}^{**} = B_{zR}^{**} \equiv B_z^{**} \quad (13)$$

$$v^{**} = \frac{\sqrt{\rho_L^*}v_L^* + \sqrt{\rho_R^*}v_R^* + (B_{yR}^* - B_{yL}^*)\text{sgn}(B_x)}{\sqrt{\rho_L^*} + \sqrt{\rho_R^*}} \quad (14)$$

$$w^{**} = \frac{\sqrt{\rho_L^*}w_L^* + \sqrt{\rho_R^*}w_R^* + (B_{yR}^* - B_{yL}^*)\text{sgn}(B_x)}{\sqrt{\rho_L^*} + \sqrt{\rho_R^*}} \quad (15)$$

$$B_y^{**} = \frac{\sqrt{\rho_L^*}B_{yR}^* + \sqrt{\rho_R^*}B_{yL}^* + \sqrt{\rho_L^*\rho_R^*}(v_R^* - v_L^*)\text{sgn}(B_x)}{\sqrt{\rho_L^*} + \sqrt{\rho_R^*}} \quad (16)$$

$$B_z^{**} = \frac{\sqrt{\rho_L^*}B_{yR}^* + \sqrt{\rho_R^*}B_{yL}^* + \sqrt{\rho_L^*\rho_R^*}(w_R^* - w_L^*)\text{sgn}(B_x)}{\sqrt{\rho_L^*} + \sqrt{\rho_R^*}} \quad (17)$$

The jump condition for the energy

$$e_L^{**} = e_L^* - \rho_L^*(\mathbf{v}_L^* \cdot \mathbf{B}_L^* - \mathbf{v}^{**} \cdot \mathbf{B}^{**}) \quad (18)$$

$$e_R^{**} = e_R^* + \rho_R^*(\mathbf{v}_R^* \cdot \mathbf{B}_R^* - \mathbf{v}^{**} \cdot \mathbf{B}^{**}) \quad (19)$$

The fluxes found from the jump conditions...

Signal speeds for the left and right sides of the Riemann fan are estimated as

$$S_L = \min(u_L, u_R) - \max(c_{fL}, c_{fR}) \quad (20a)$$

$$S_R = \max(u_L, u_R) + \max(c_{fL}, c_{fR}) \quad (20b)$$

References

- [1] Takahiro Miyoshi and Kanya Kusano. A multi-state HLL approximate Riemann solver for ideal magnetohydrodynamics. *Journal of Computational Physics*, 208(1):315–344, 2005.