

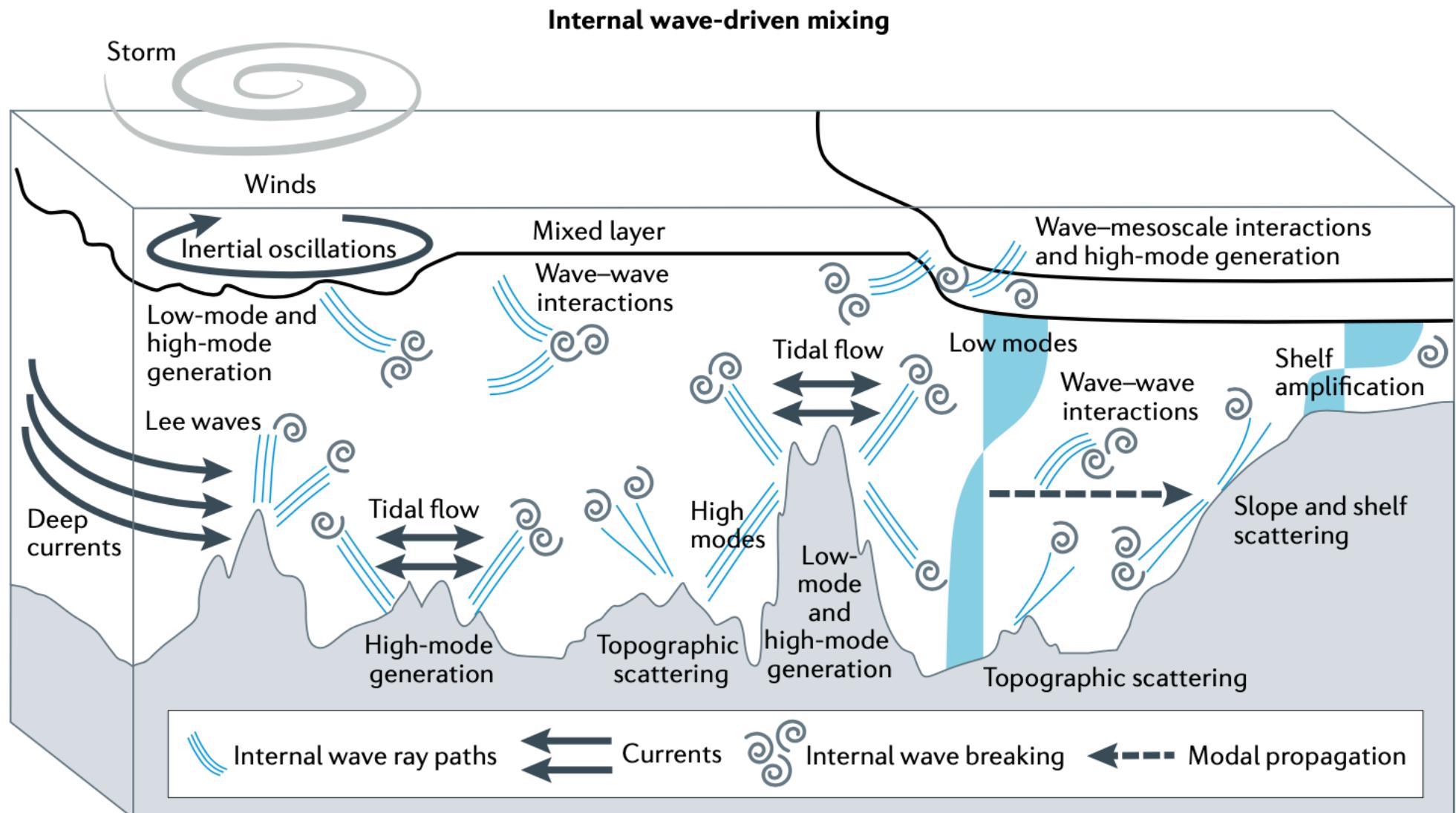
# INTERNAL WAVES

## 3. GENERATION AND DISSIPATION

---

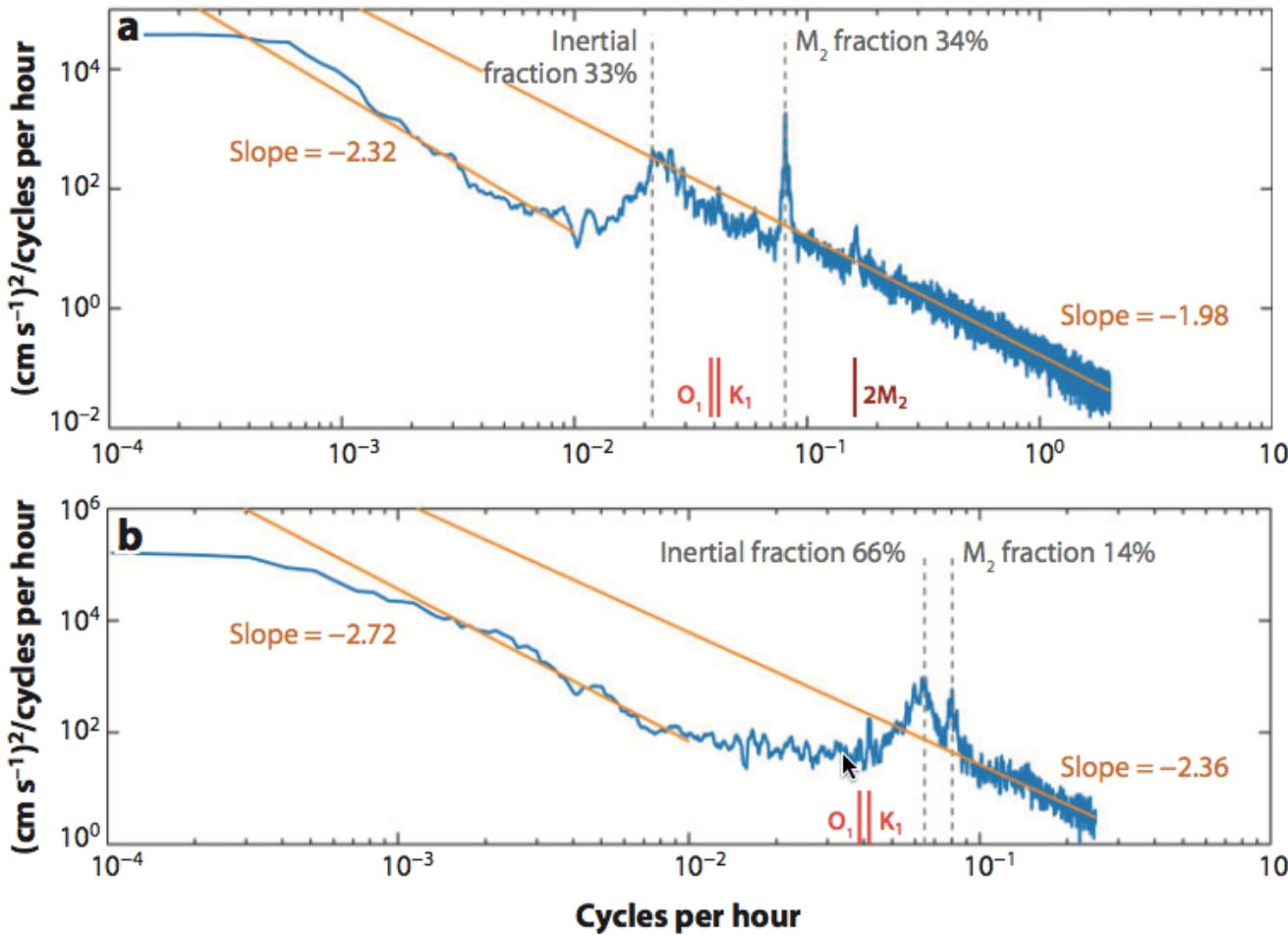
- **Today : Generation and propagation of Internal waves**
  - Internal tides
  - Lee waves
  - Near-Inertial waves
- **Next lesson : Propagation and dissipation of internal waves**

# 1. Internal waves generation



# 1. Internal waves generation

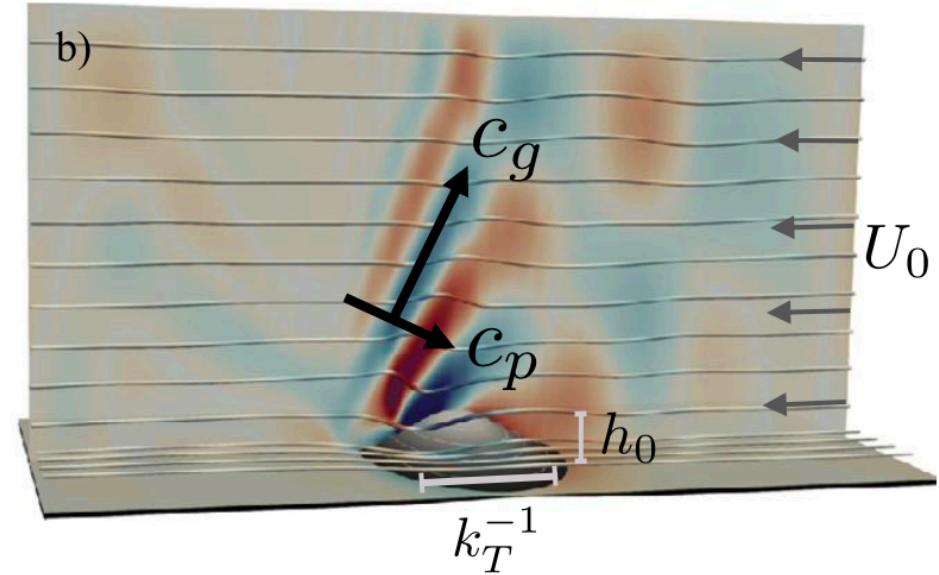
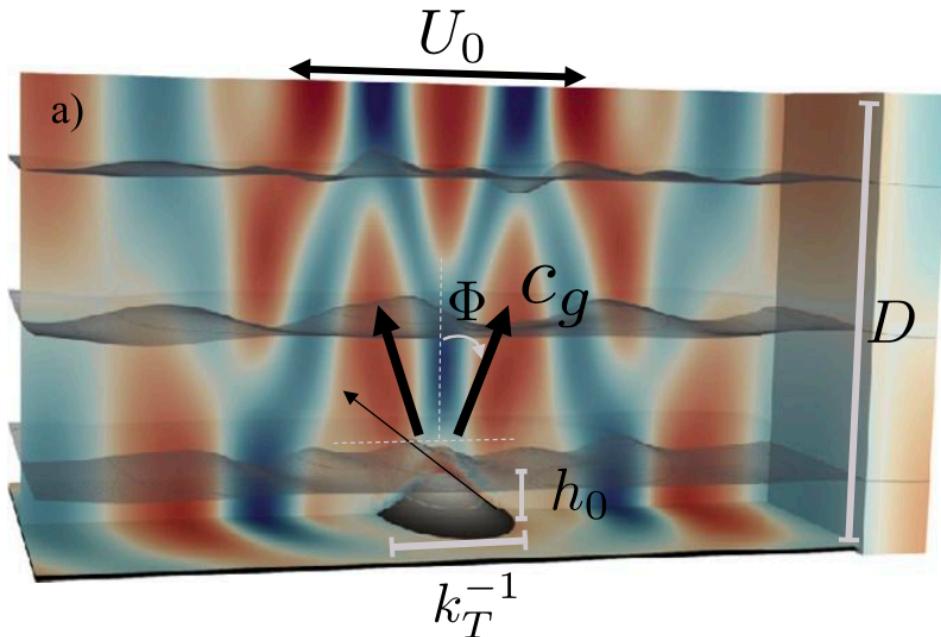
- Winds generate mostly near-inertial waves (frequency close to  $f$ )
- Barotropic tides generate internal waves at the frequency of tides that are called internal tides (or baroclinic tides)



(a) Kinetic energy estimate for an instrument in the western North Atlantic near  $15^\circ\text{N}$  at 500 m. (b) Power density spectral estimate from a record at 1000 m at  $50.7^\circ\text{S}$ ,  $143^\circ\text{W}$ , south of Tasmania in the Southern Ocean

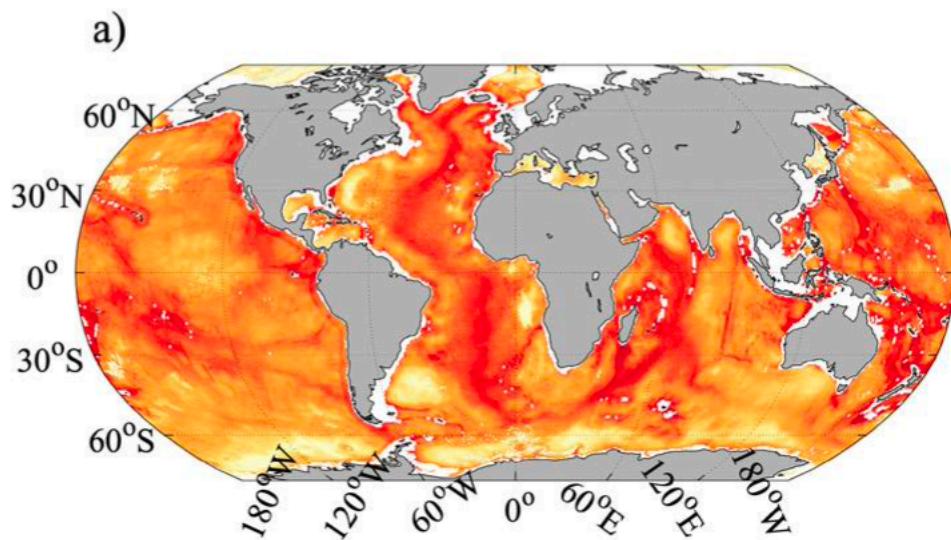
# 1. Internal waves generation

Idealised experiments of the generation of internal tides (left) and lee waves (right)

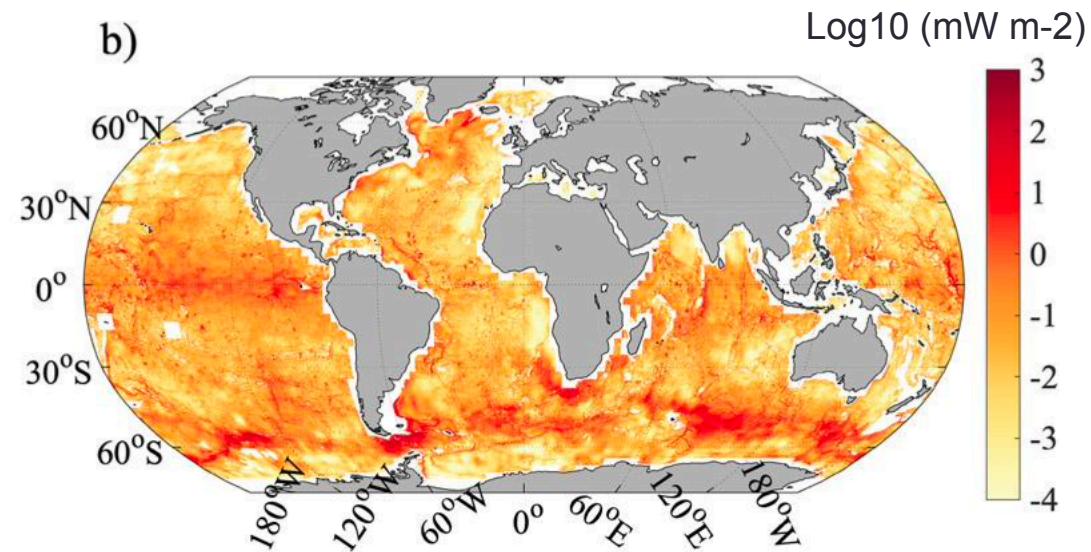


# 1. Internal waves generation

Geography of the generation of internal tides (left) and lee waves (right)



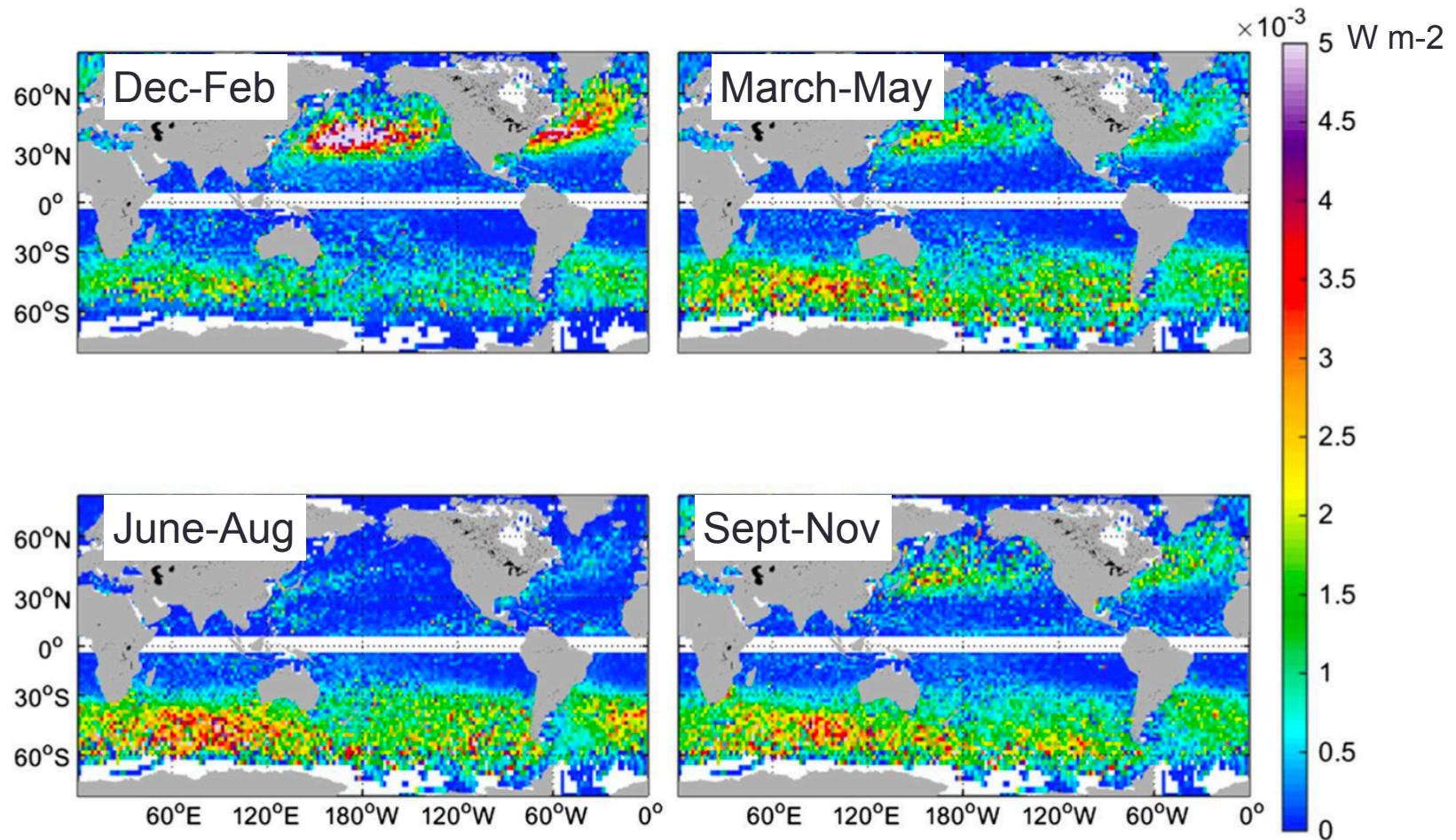
Global contribution ~ 1 TW



Global contribution ~ 0.15-0.75 TW

# 1. Internal waves generation

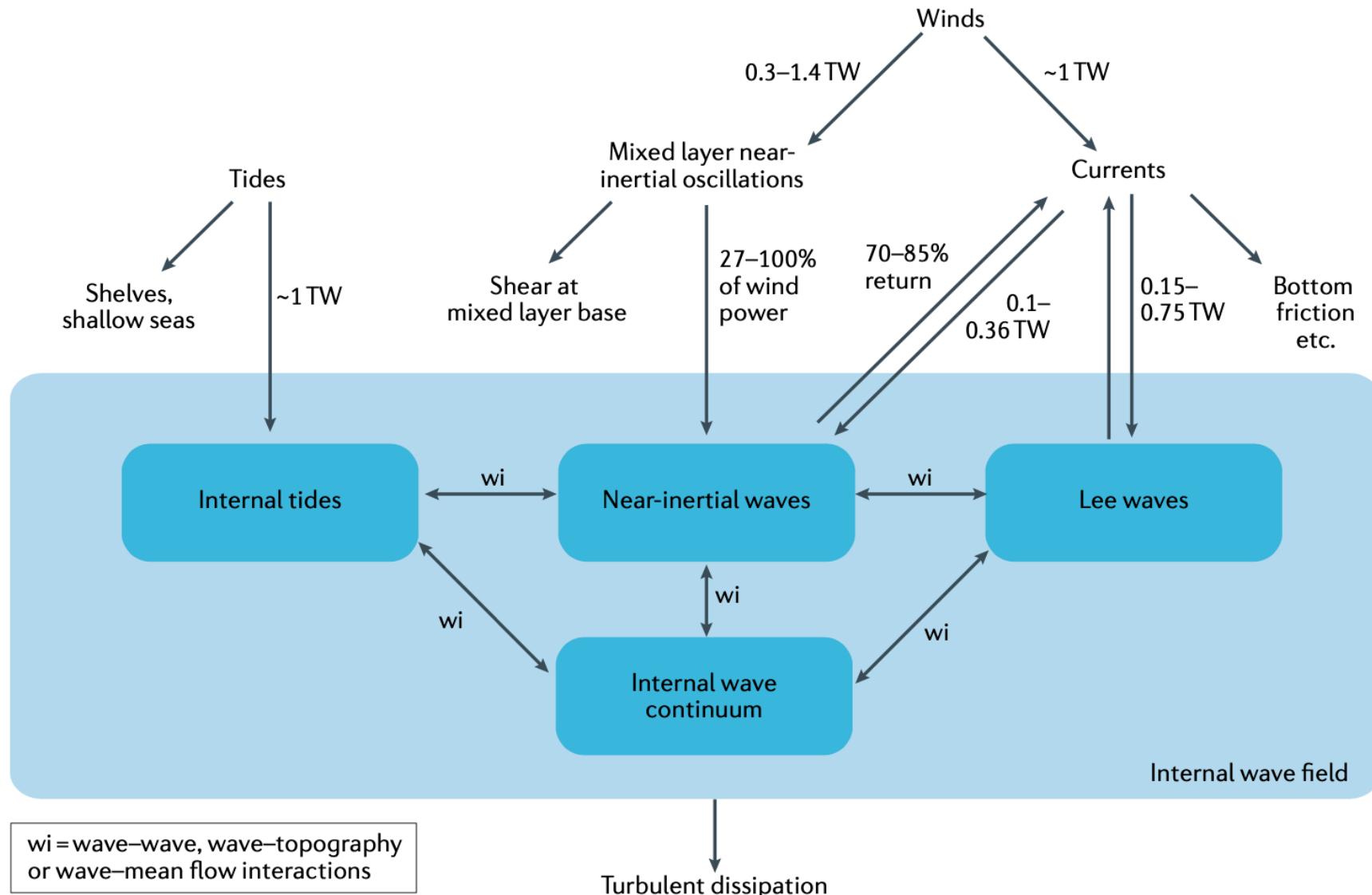
Geography of the generation of near-inertial waves by the wind



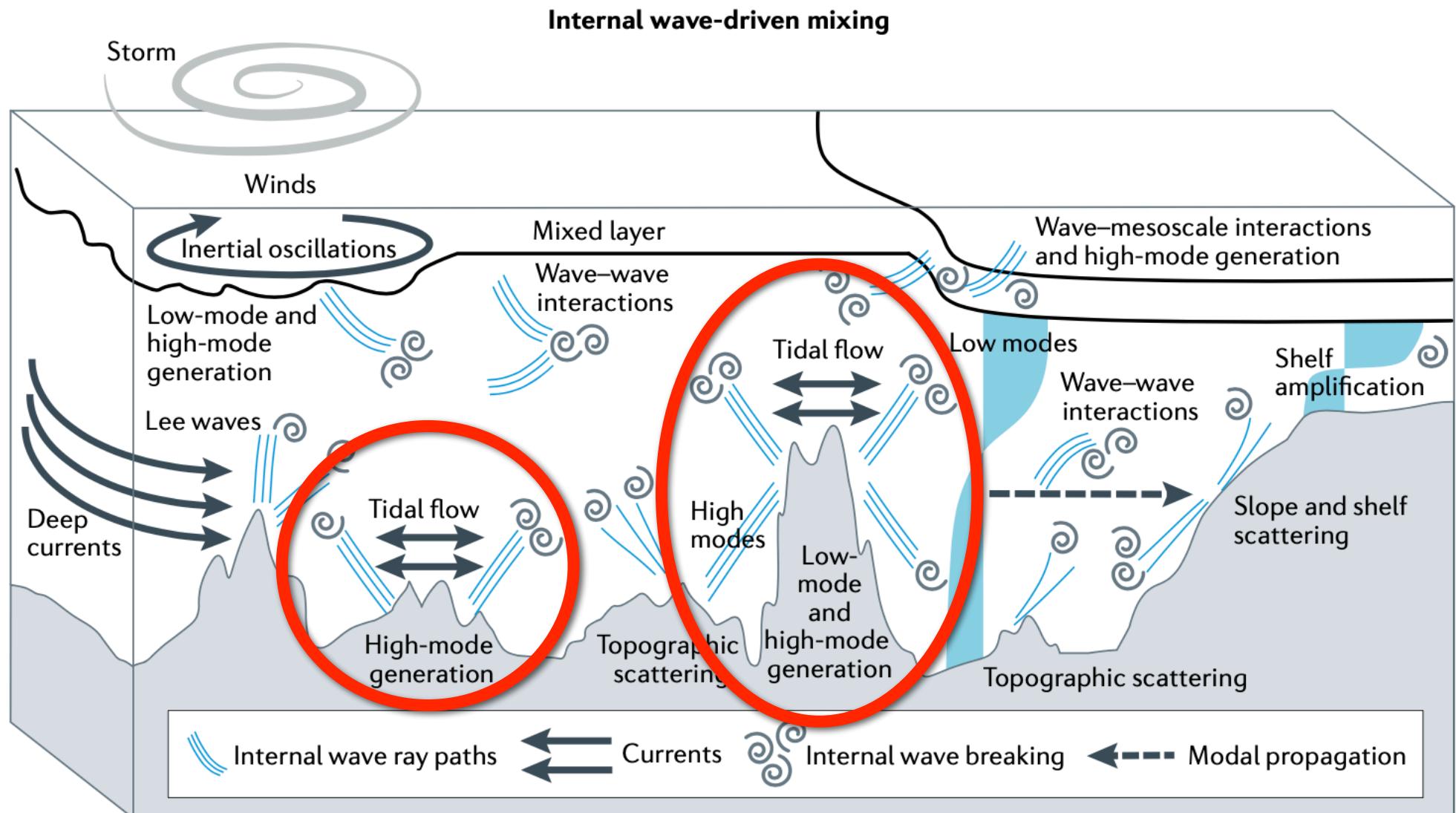
Global contribution  $\sim 0.3\text{-}1.5$  TW

Thomas and Zhai — mixing book 2022

# 1. Internal waves generation

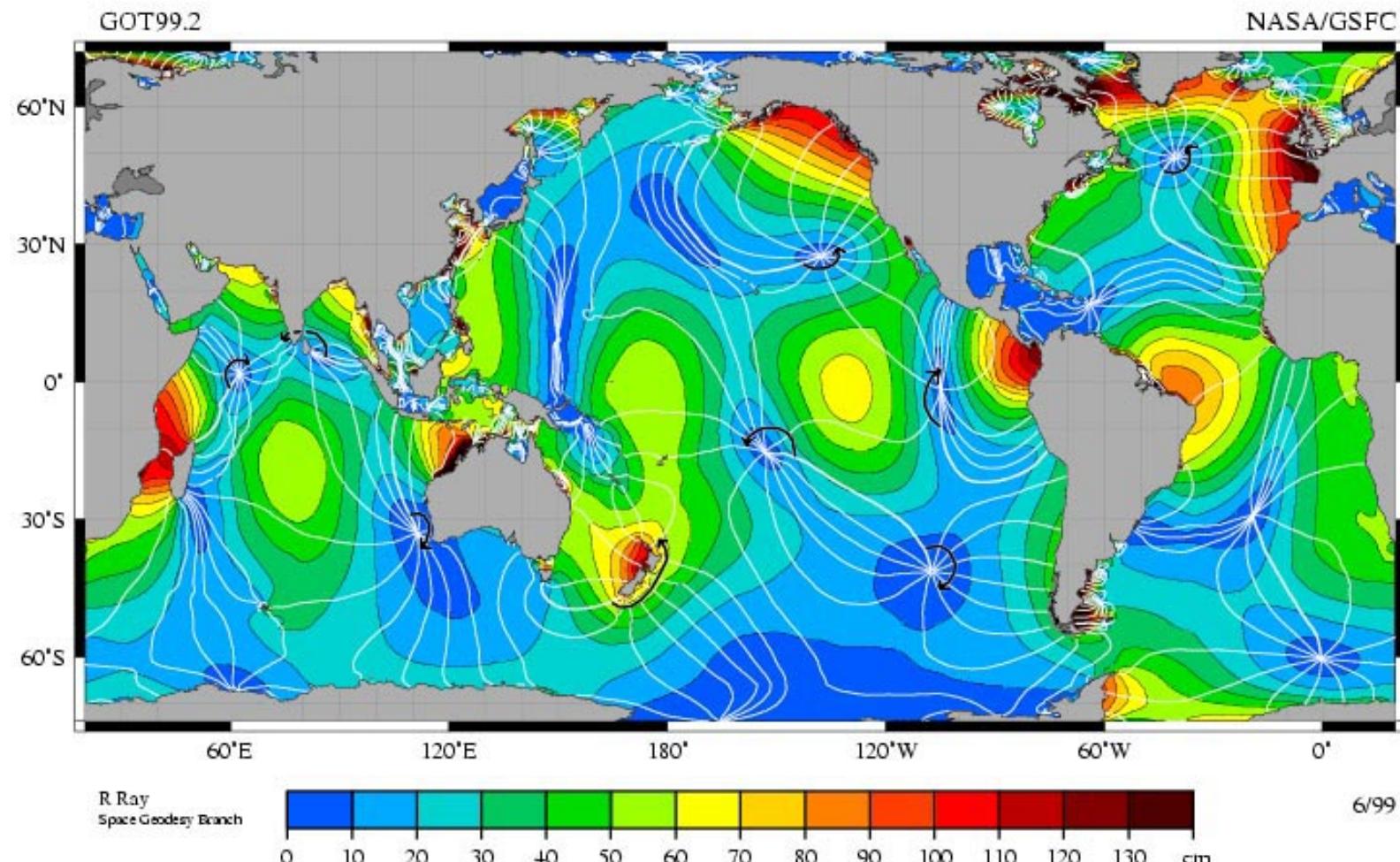


# 1.1 Generation of internal tides



# 1.1 Generation of internal tides

- Barotropic tides

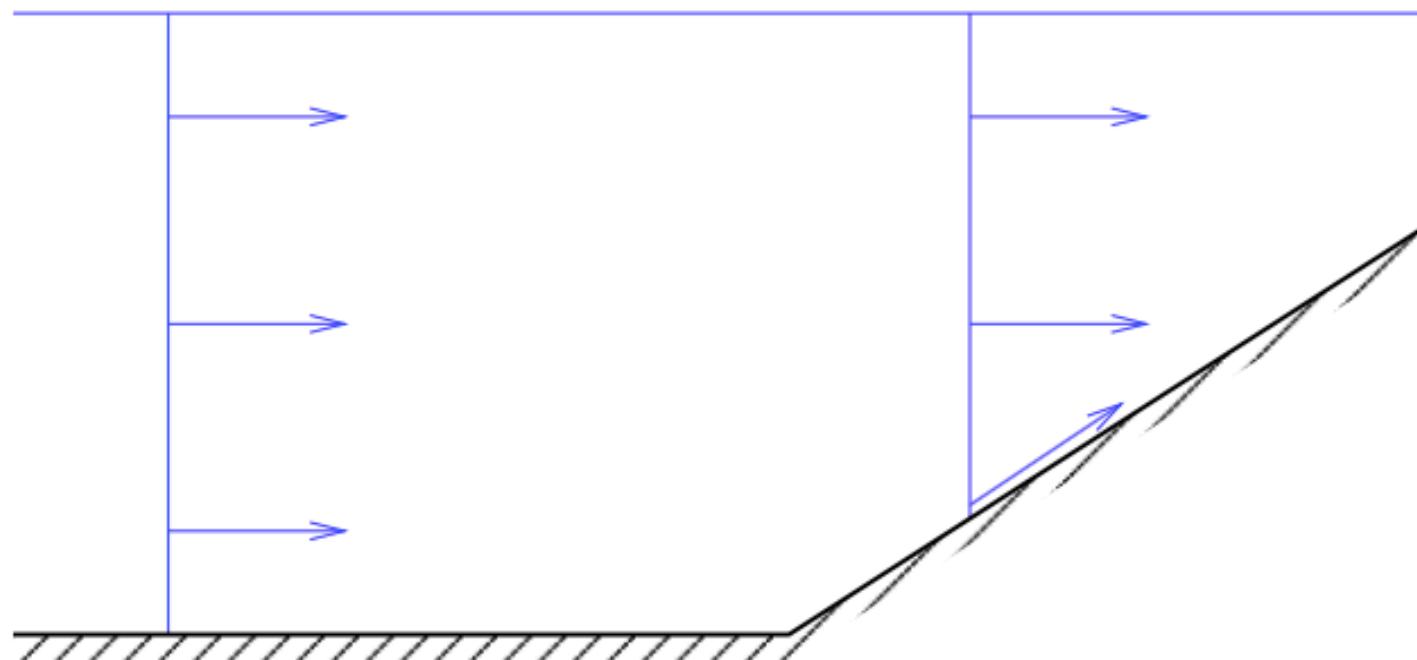


The  $M_2$  tidal constituent. Amplitude is indicated by color, and the white lines are cotidal differing by 1 hour. The colors indicate where tides are most extreme (highest highs, lowest lows), with blues being least extreme.

# 1.1 Generation of internal tides

- Generation of Internal Tide on a Slope:

A barotropic velocity profile impinging on a sloping boundary. On the boundary, the velocity vector can only have a component tangential to it and hence over the slope, the velocity profile is no longer barotropic. Energy must go into generating internal waves, and propagate away along characteristics (WHOI course, Hendershott and Garrett)



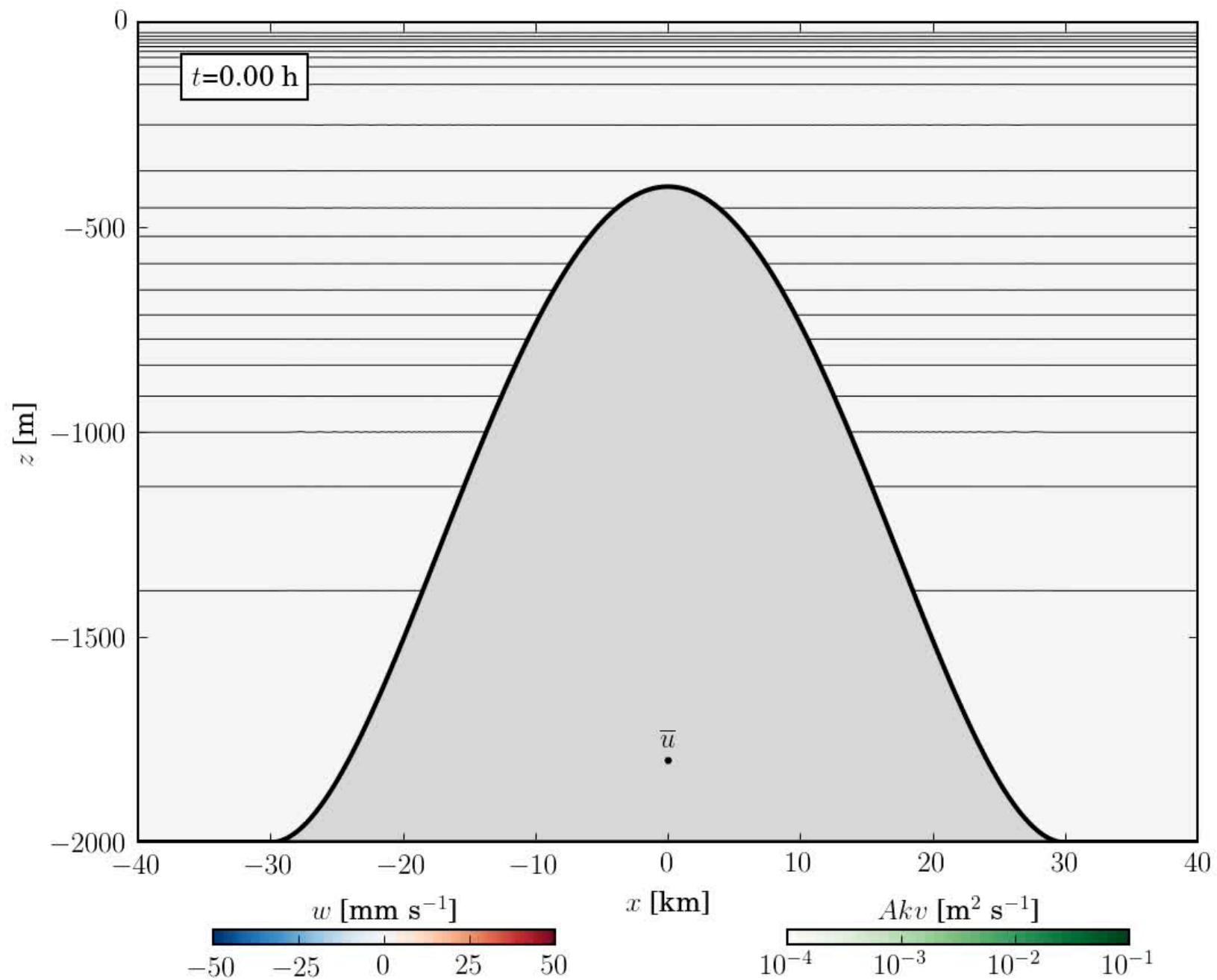
# 1.1 Generation of internal tides

## Internal Tide Generation in the Deep Ocean

Chris Garrett and Eric Kunze

Department of Physics and Astronomy, University of Victoria, Victoria,  
British Columbia V8W 2Y2; Canada

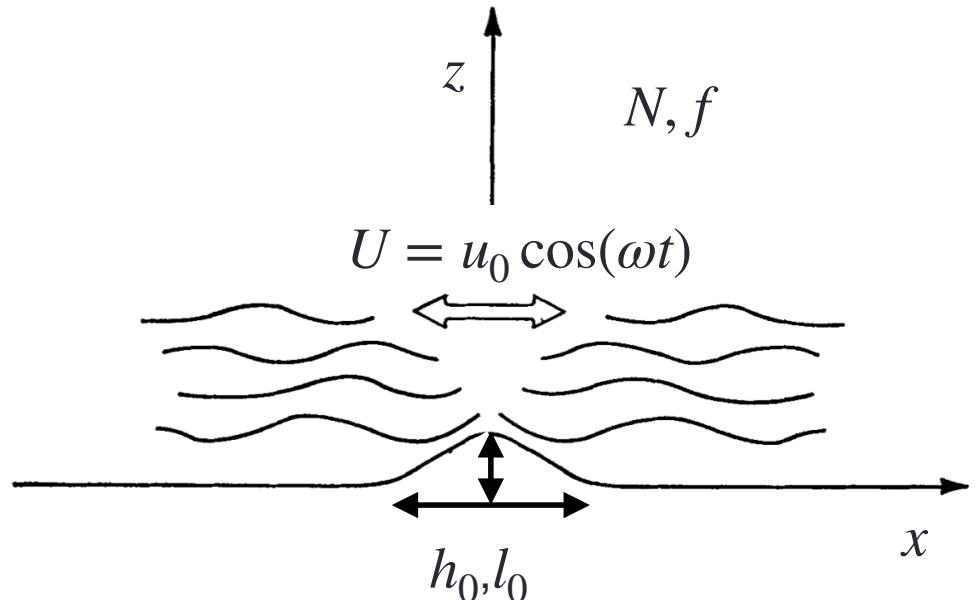
# 1.1 Generation of internal tides



# 1.1 Generation of internal tides

6 dimensional parameters :

- Coriolis frequency  $f$
- Brunt-Vaisala frequency  $N$
- Forcing (tidal) frequency  $\omega$
- Topographic height  $h_0$
- Topographic wavenumber  $k = 2\pi/l_0$
- Magnitude of the tidal current  $u_0$



4 dimensionless parameters :

- $\omega/N$
- $\omega/f$
- $ku_0/\omega$ , ratio of the “tidal excursion” to the topographic length scale
- $\epsilon = \frac{kh_0}{\alpha} = \frac{\text{topographic slope}}{\text{beam slope}}$ , with  $\alpha = \sqrt{\frac{\omega^2 - f^2}{N^2 - \omega^2}}$

# 1.1 Generation of internal tides

Assumptions to derive a linear expression for  $w$  :

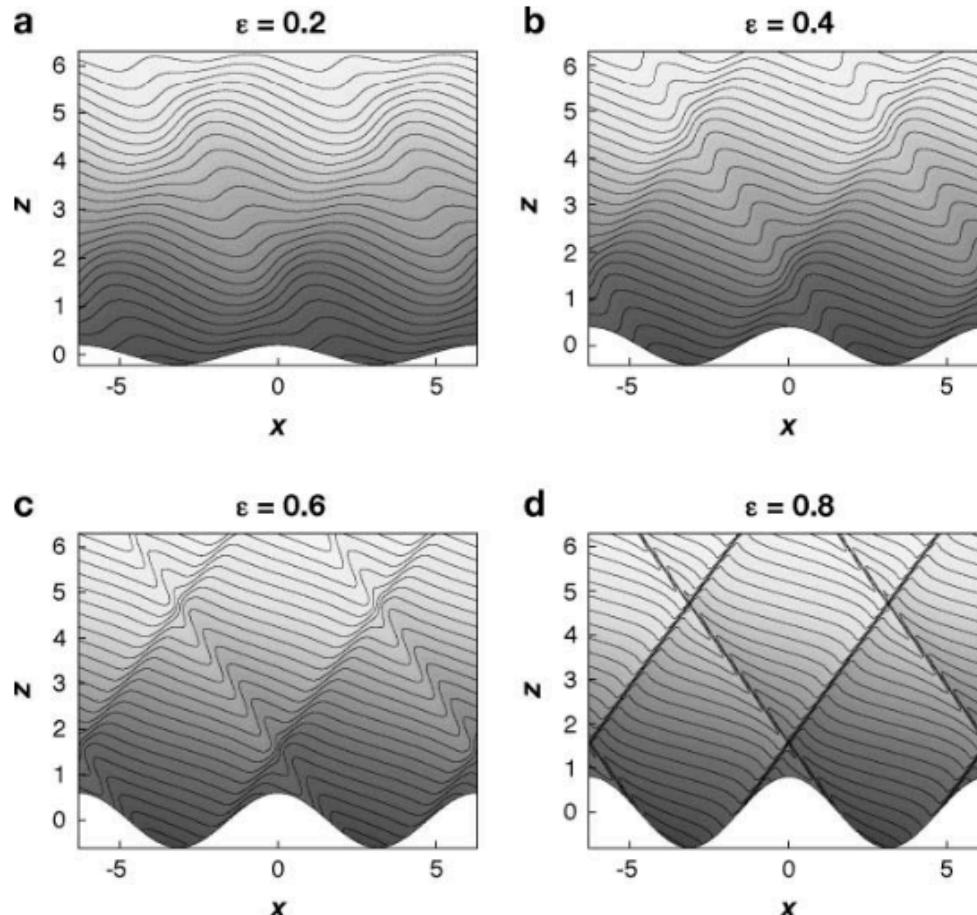
$$(1) - f < \omega < N$$

# 1.1 Generation of internal tides

Assumptions to derive a linear expression for  $w$  :

(1) -  $f < \omega < N$

(2) - The topographic slopes are subcritical :  $\epsilon = \frac{kh_0}{\alpha} = \frac{\text{topographic slope}}{\text{beam slope}} < 1$



Beam slope in the deep ocean :  
 $N^2 = 10^{-6} - 10^{-7} (\text{rad s}^{-1})^2$   
 $f = 10^{-4} \text{ rad s}^{-1}$   
 $\omega = 1.4 \times 10^{-4} \text{ rad s}^{-1}$  (M2 tide)  
Give  
 $\alpha = 0.09 - 0.34 = 9 - 34\%$  (steep !)

Vs typical topographic slopes are ~1-2%

# 1.1 Generation of internal tides

Assumptions to derive a linear expression for  $w$  :

(1) -  $f < \omega < N$

(2) - The topographic slopes are subcritical :  $\epsilon = \frac{kh_0}{\alpha} = \frac{\text{topographic slope}}{\text{beam slope}} < 1$

(3) - The tidal excursion is small :  $ku_0/\omega < 1$  :

—> the distance traveled by a particle in half a tidal cycle =  $u_0 \times \frac{1}{2} \times \frac{2\pi}{\omega} = \frac{\pi u_0}{\omega}$

should be smaller than half the topographic wavelength  $\frac{1}{2}l_0 = \frac{\pi}{k}$

—> watch again the movie slide 11 !

For a semi-diurnal tide with  $\omega = 1.4 \times 10^{-4} \text{ rad s}^{-1}$  and  $u_0 = 1 - 4 \text{ cm s}^{-1}$ ,

The theory is valid for  $l_0 > \frac{2\pi u_0}{\omega} > 400 - 1800 \text{ m}$ ,

—> OK for most topographic structures over mid-ocean ridges

# 1.1 Generation of internal tides

For the 2D (x-z) problem with  $h = h_0 \cos(kx)$  and a tidal current  $u = u_0 \cos(\omega t)$ , the response is:

$$u = mh_0u_0 \cos(kx)\sin(mz + \omega t)$$

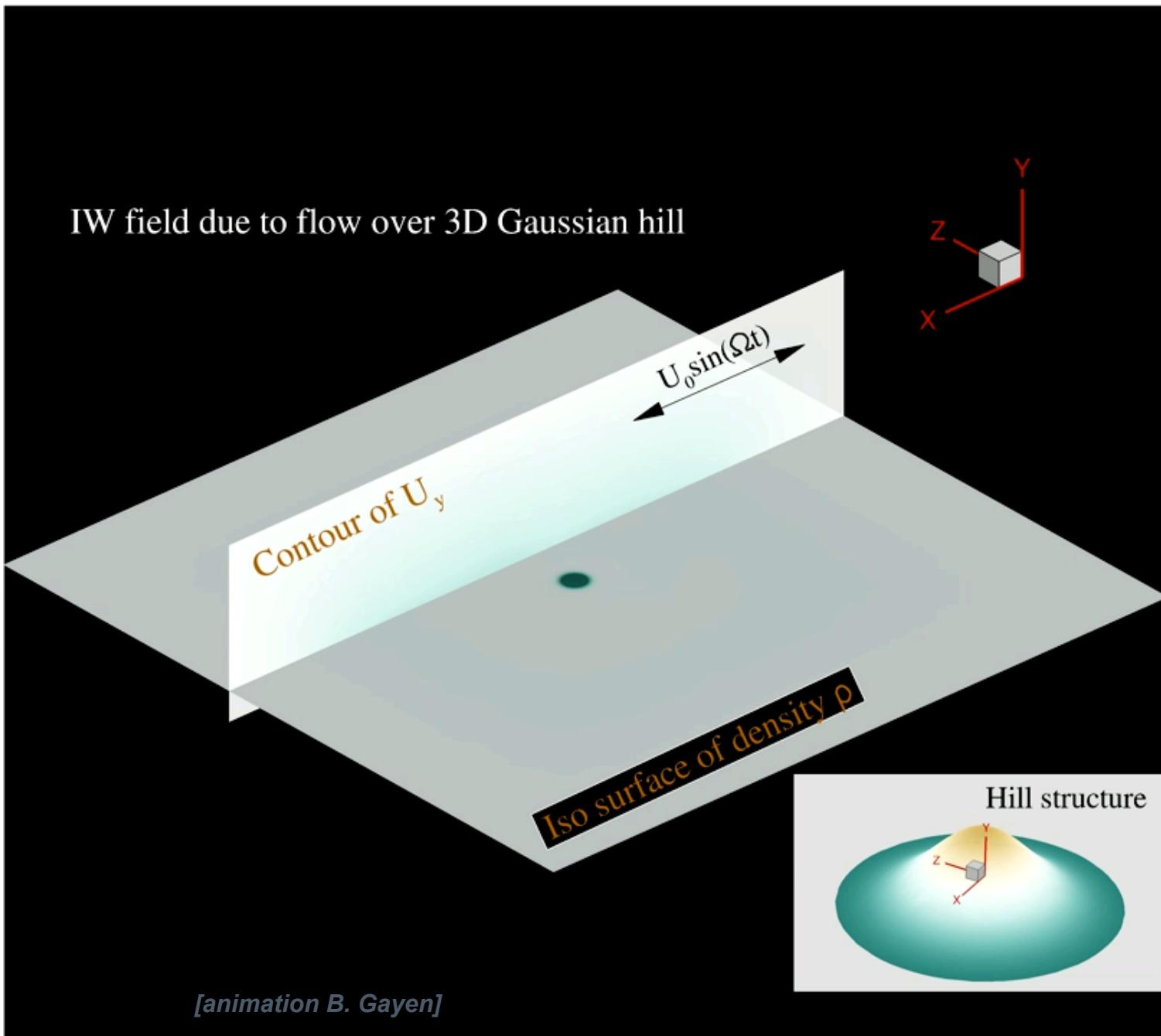
$$w = -kh_0u_0 \sin(kx)\cos(mz + \omega t)$$

The vertical energy flux is  $F = \frac{1}{4}\rho_0\omega^{-1} [(N^2 - \omega^2)(\omega^2 - f^2)]^{1/2} ku_0^2 h_0^2$

**Take-home message** : the energy flux (conversion) depends on :

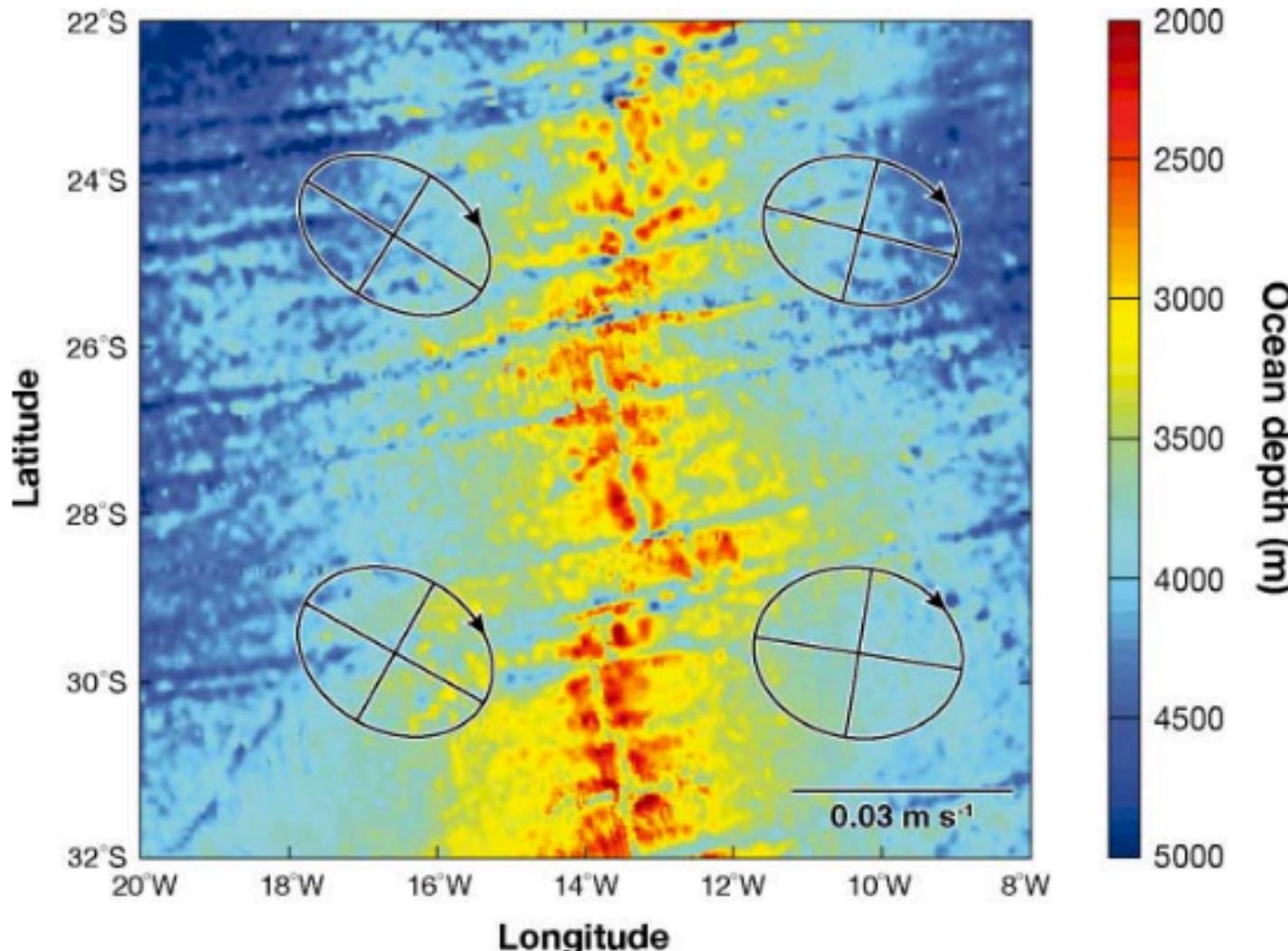
- The shape of the seafloor topography  $k, h_0^2$
- The frequencies of the systems  $\omega, f, N$
- The squared amplitude of the barotropic tide current  $u_0^2$

# 1.1 Generation of internal tides



# 1.1 Generation of internal tides

This can be generalised to a 3D problem with random bathymetry

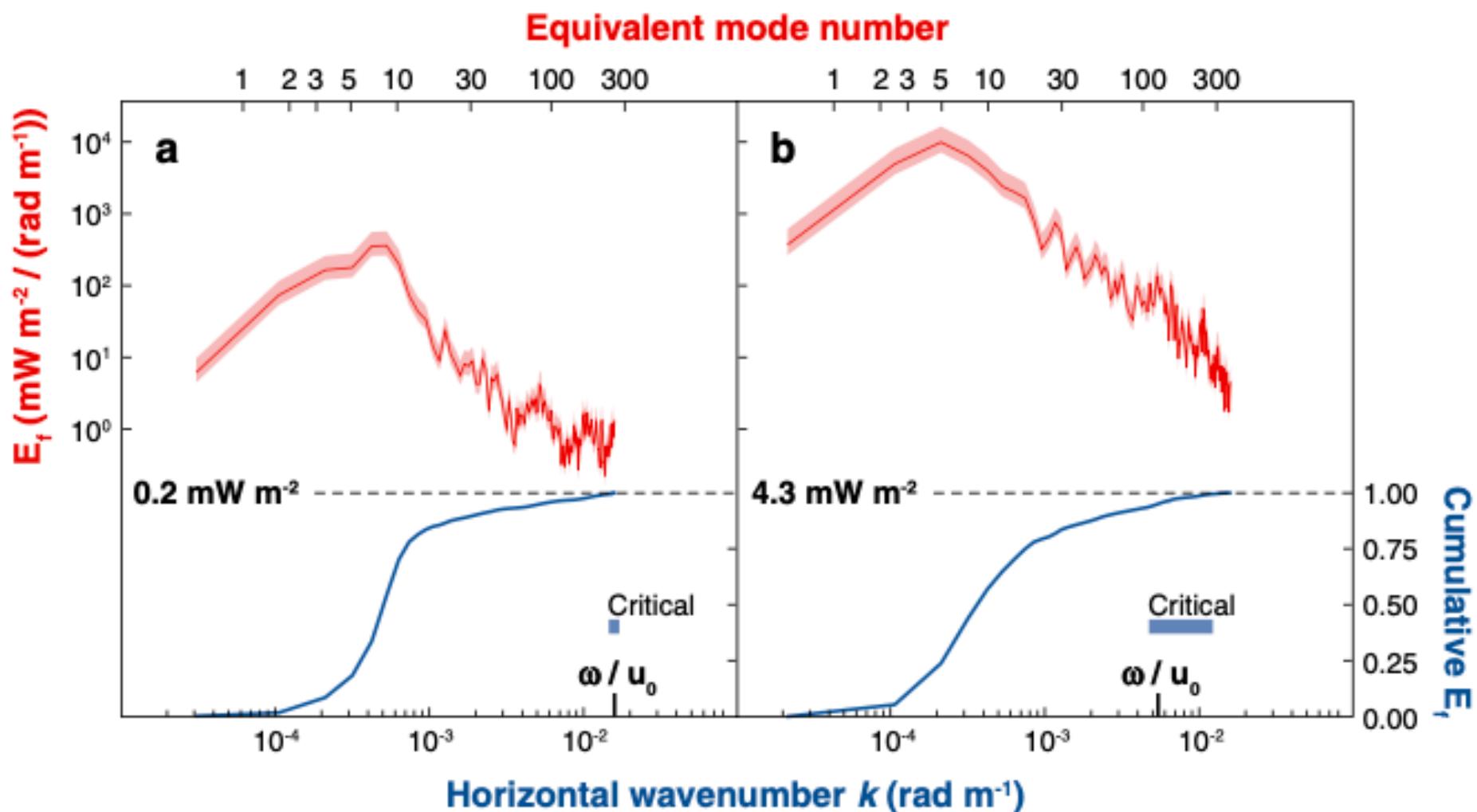


You must know :

- bathymetry properties (power density spectrum)
- tidal currents

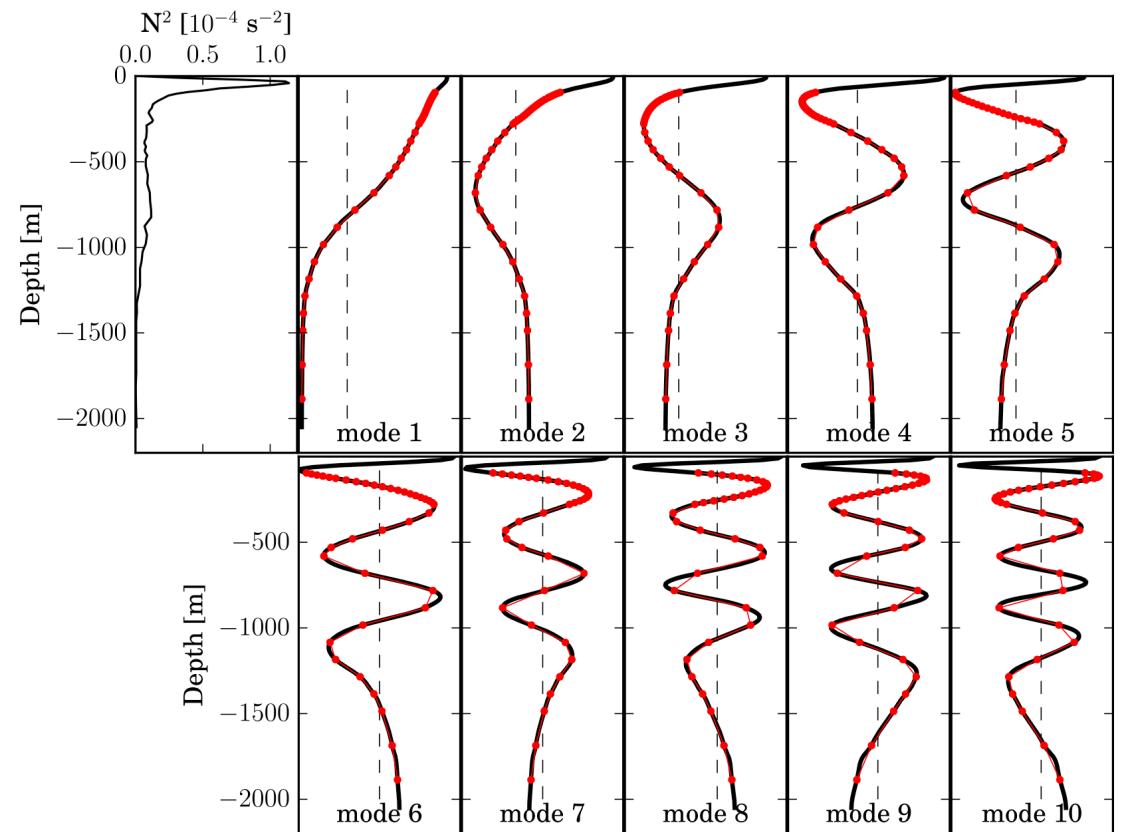
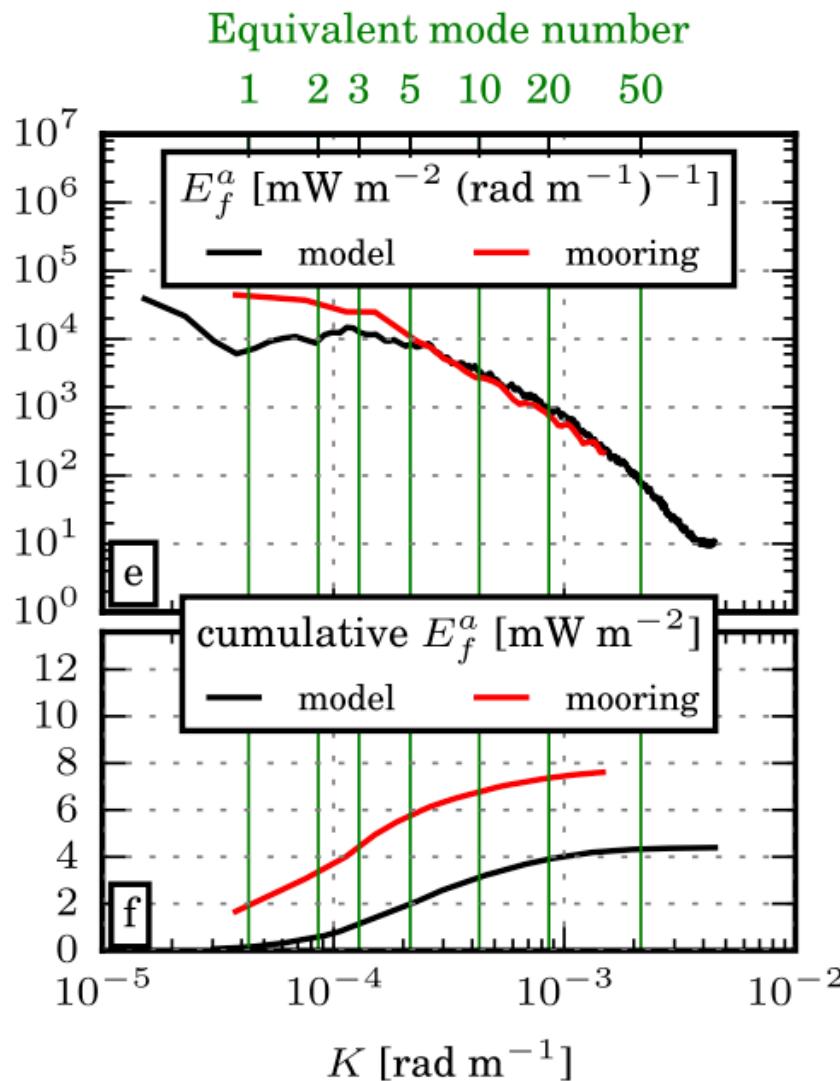
# 1.1 Generation of internal tides

This can be generalised to a 3D problem with random bathymetry



# 1.1 Generation of internal tides

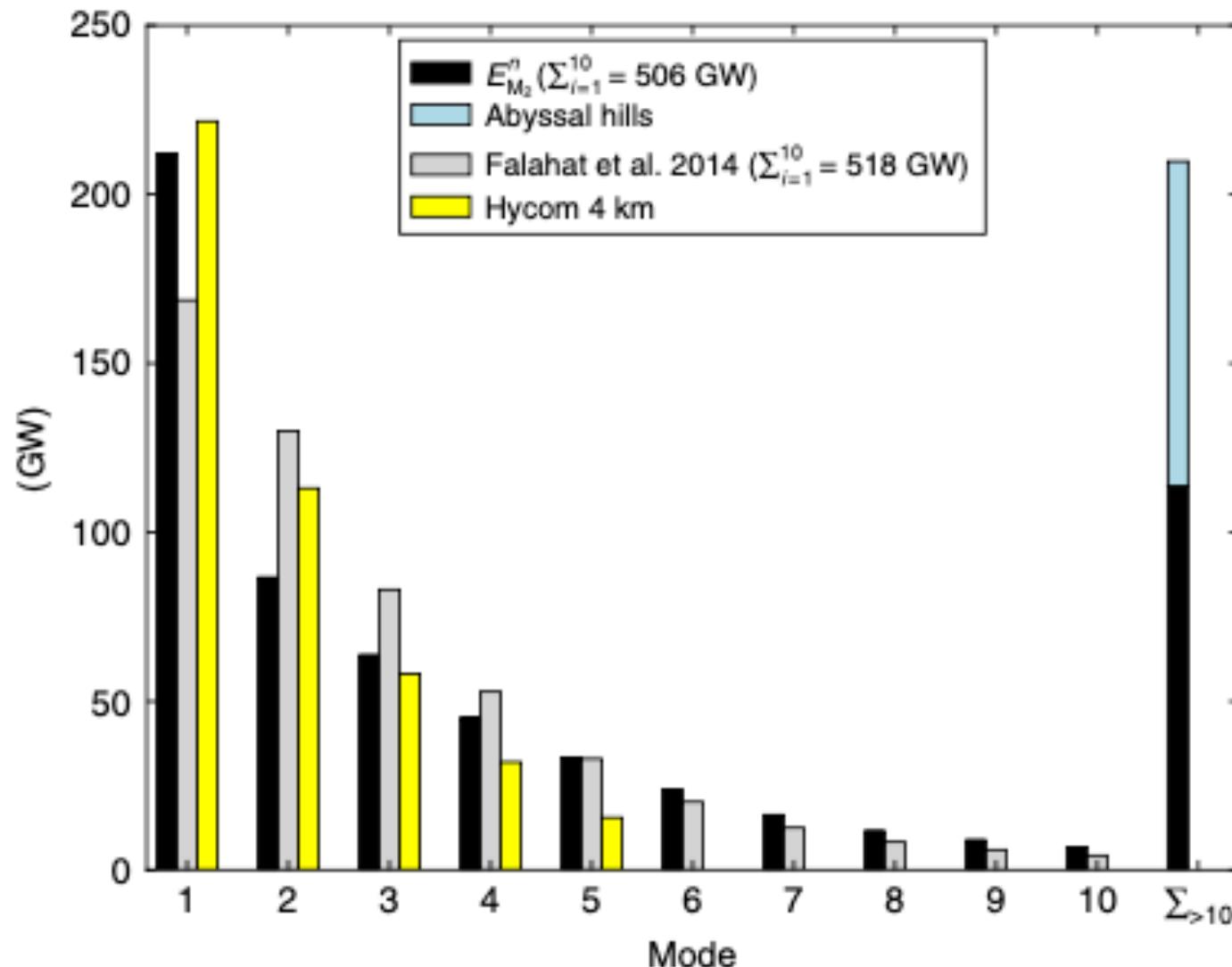
This can be generalised to a 3D problem with random bathymetry



Vertical modes reconstructed with instruments on a mooring line

# 1.1 Generation of internal tides

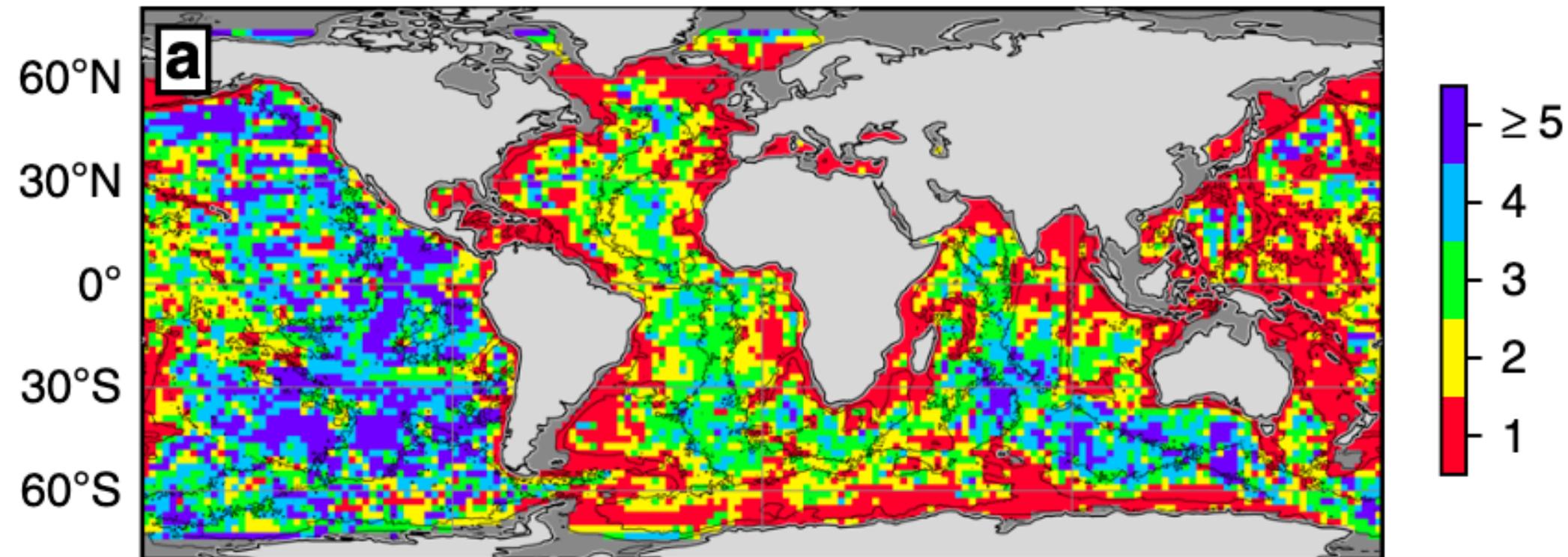
Mapping the internal tide generation mode by mode



# 1.1 Generation of internal tides

Mapping the internal tide generation mode by mode

Most energetic mode



Interest : high modes are more unstable hence generate more mixing (next lesson)

# 1.1 Generation of internal tides

## Internal Tide Generation in the Deep Ocean

Chris Garrett and Eric Kunze

Department of Physics and Astronomy, University of Victoria, Victoria,  
British Columbia V8W 2Y2; Canada

On continental shelf breaks : assumptions to derive a linear theory do not hold !

# 1.1 Generation of internal tides

On continental shelf breaks : assumptions to derive a linear theory do not hold ! **Why ?**

# 1.1 Generation of internal tides

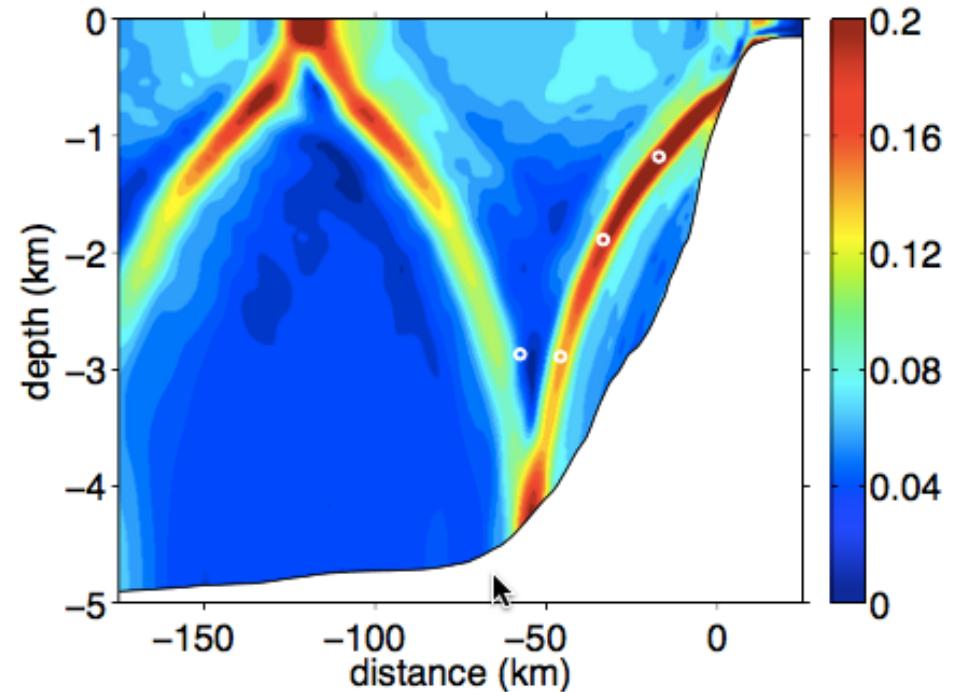
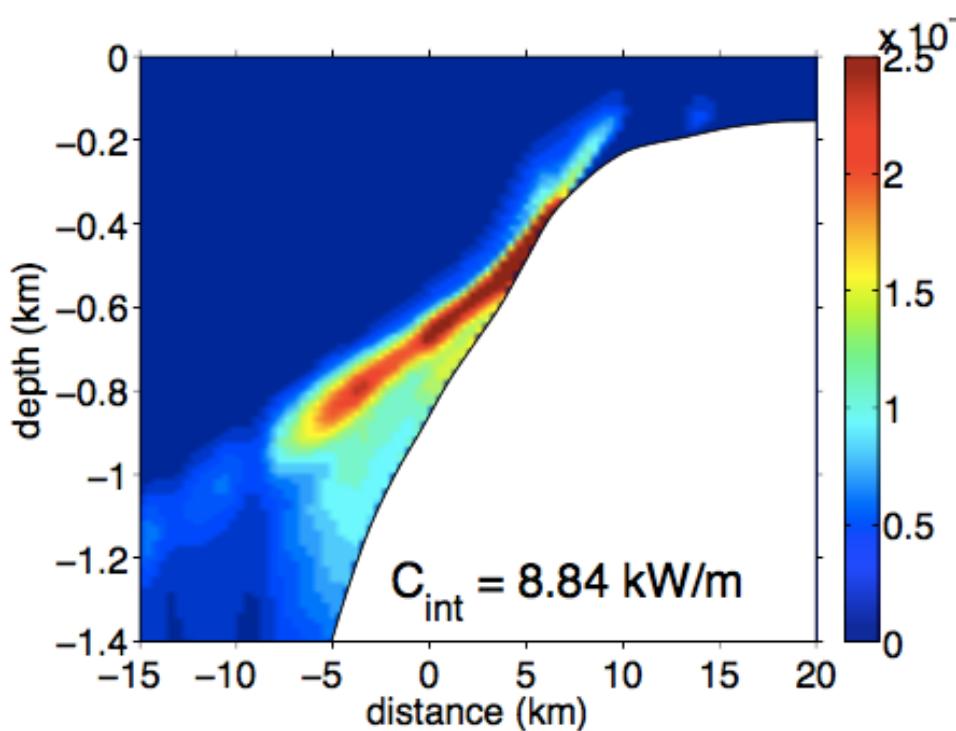
On continental shelf breaks : assumptions to derive a linear theory do not hold ! **Why ?**

- Shallower depths :  $N^2$  is much larger —> beam is more horizontal
- Topographic slopes are steeper

$$\rightarrow \epsilon = \frac{kh_0}{\alpha} = \frac{\text{topographic slope}}{\text{beam slope}} \text{ can be no longer } < 1$$

—> critical slopes  $\epsilon = 1$  and supercritical slopes  $\epsilon > 1$

# 1.1 Generation of internal tides

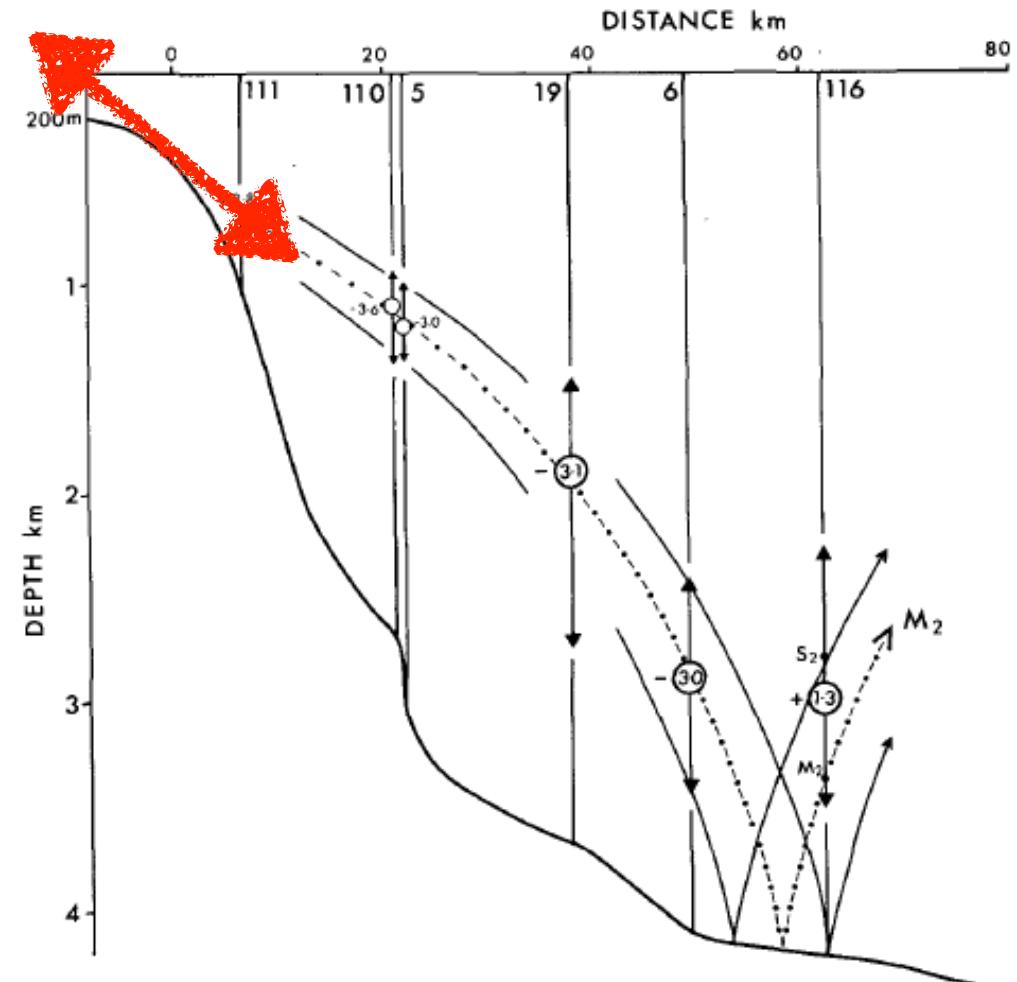


Results from a numerical model, here applied to the Bay of Biscay. Left: The spatial distribution tidally-averaged conversion rate,  $C = -\rho * \langle bW \rangle$ , in W/m<sup>3</sup>; Right: The internal tide emanating from the continental slope, here depicted in terms of the amplitude of the cross-slope velocity  $u$ , in m/s.

# 1.1 Generation of internal tides

Internal beams are preferentially excited where the slope is critical

- One ray goes toward the coast (difficult to see because of mixing).
- A second ray goes downward, is reflected at the bottom,



# 1.1 Generation of internal tides

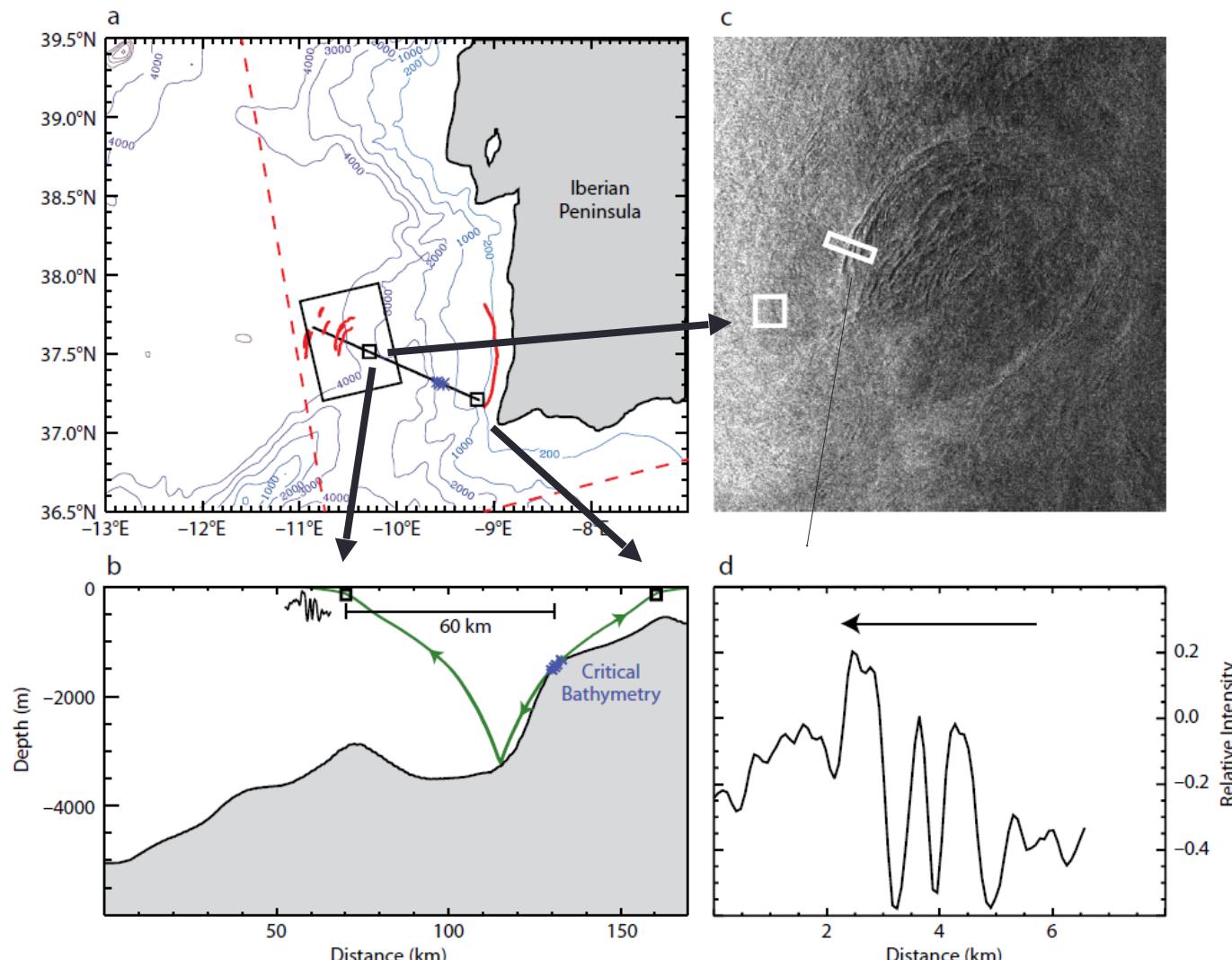
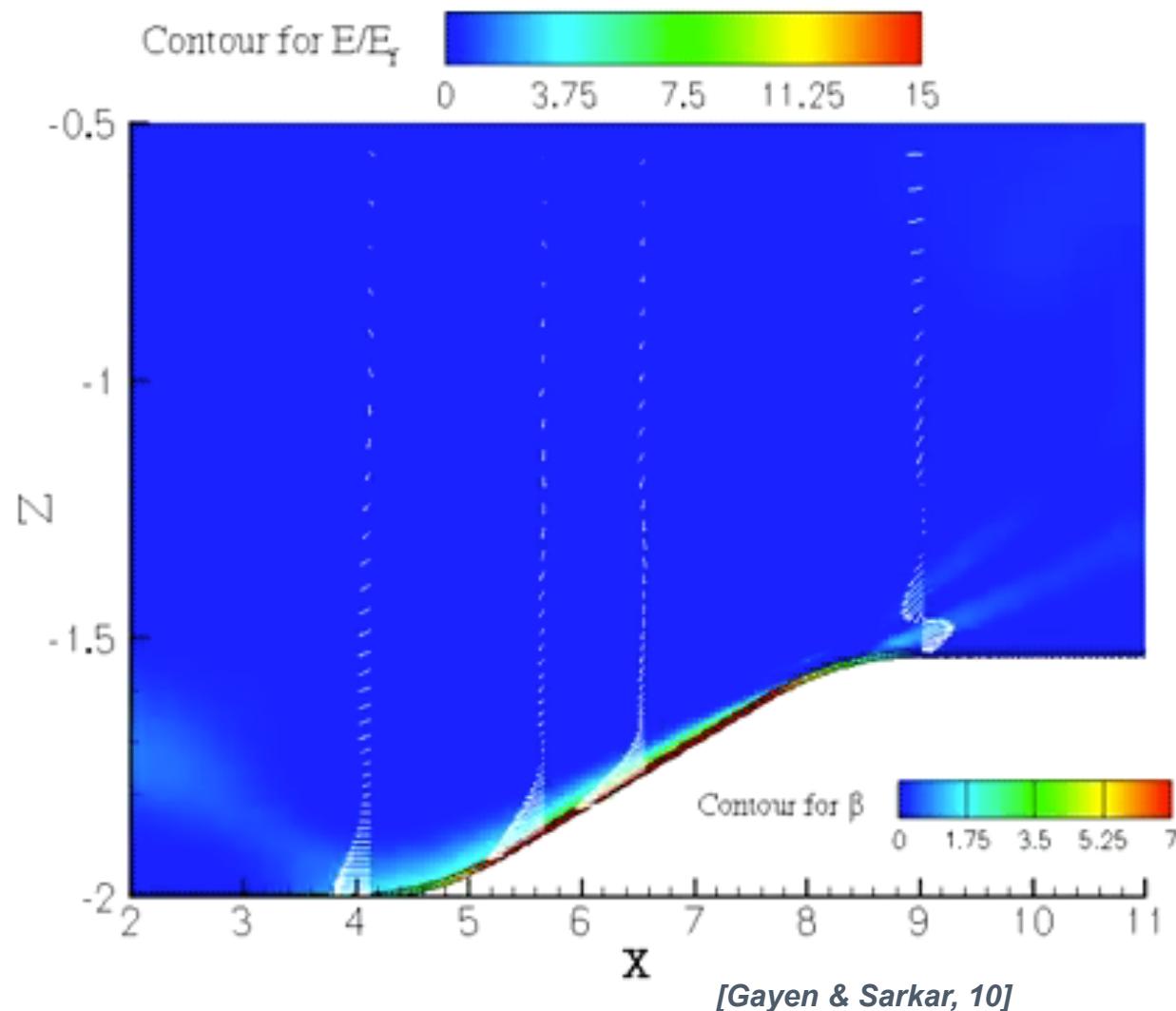


Figure 5. (a) Map of area southwest of the Iberian Peninsula with internal solitary wave crests marked in red based on one Envisat Advanced Synthetic Aperture Radar (ASAR) image in Wide-Swath Mode dated August 4, 2004 (22:27 UTC). (b) Ray-tracing diagram showing internal tide ray paths (in green; emanating from critical topography, in blue) along the black line in part (a). The small black squares in part (a) show where the ray path crosses the near-surface thermocline (taken at a depth of 50 m), and are also marked in part (b). (c) Full-resolution detail of a nonlinear internal wave train observed to propagate toward the west-northwest and believed to be generated by the tidal beam in part (b). (d) Radar backscatter profile showing the wavelength cross section of a nonlinear internal wave train generated by the tidal beam. The narrow rectangle in part (c) represents the cross section from which part (d) was obtained. The square in part (c) is a background backscatter reference used to normalize the radar profile in part (d). Zero in part (d) represents the average unperturbed backscatter of the SAR image in part (c).

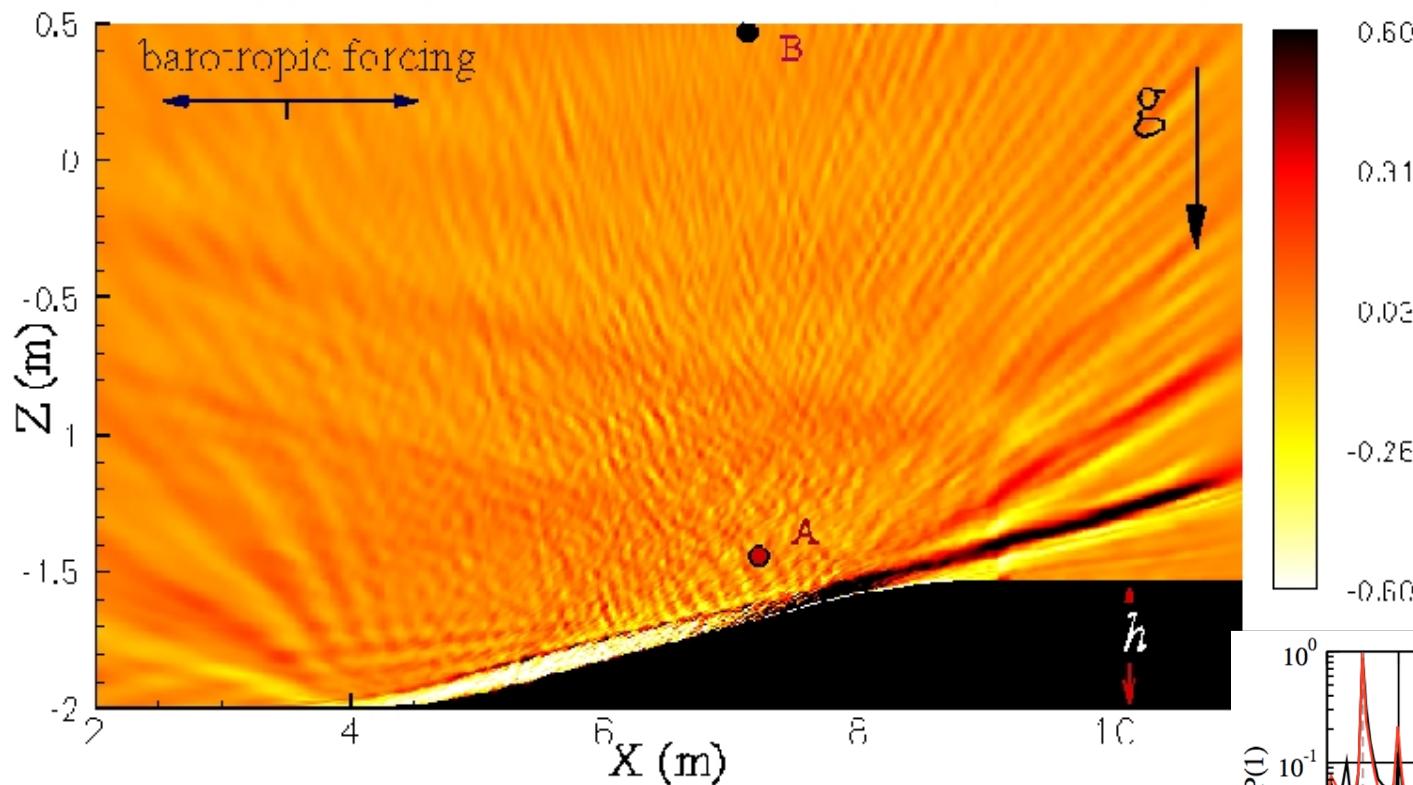
# 1.1 Generation of internal tides

- Generation of Internal Tide on a Critical Slope:

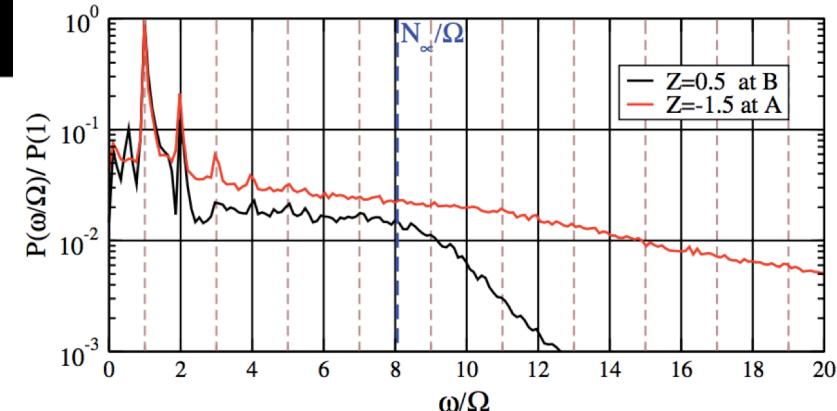


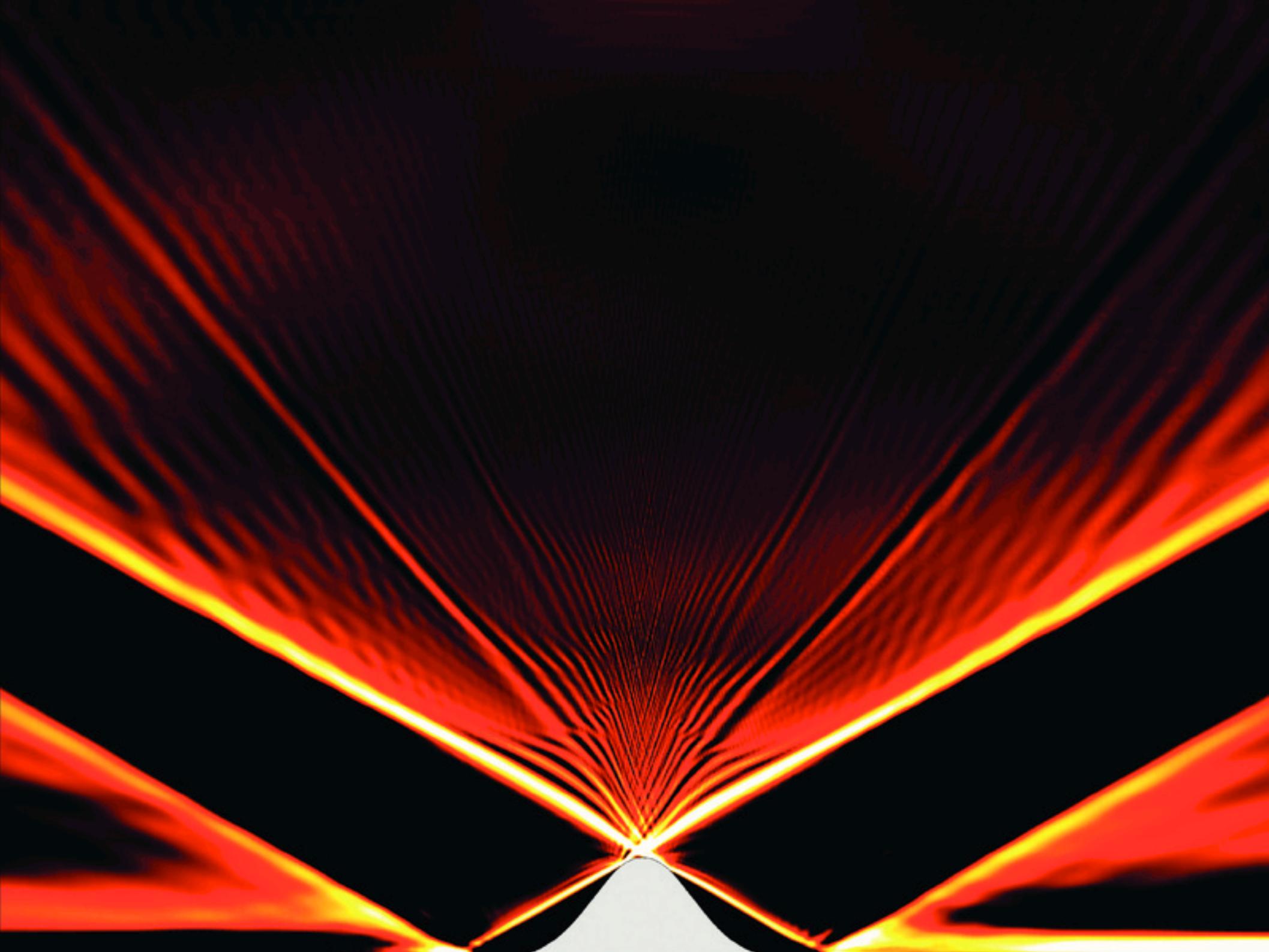
# 1.1 Generation of internal tides

- Generation of Internal Tide on a Critical Slope:

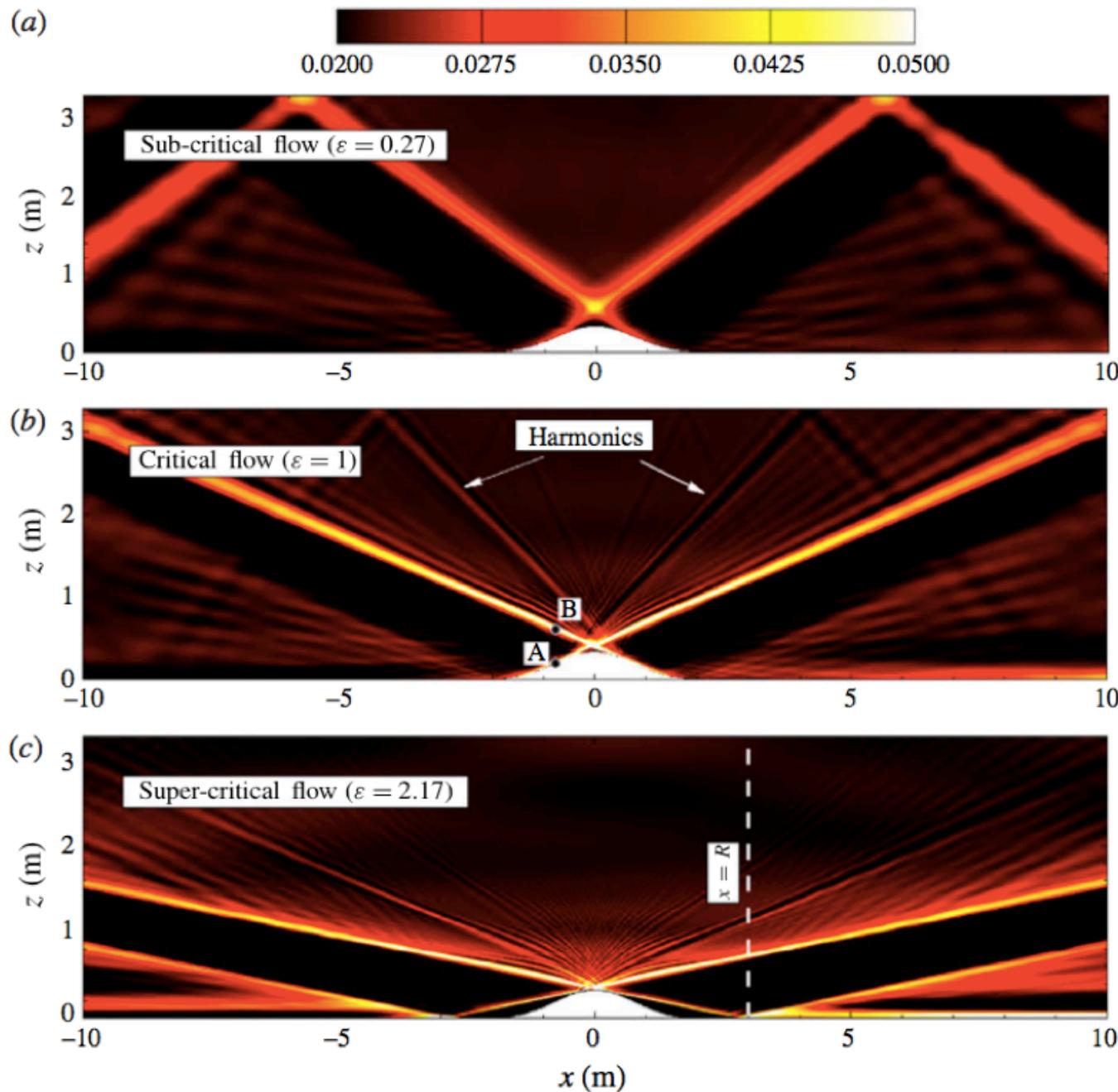


Internal wave field from DNS visualized by a slice of  $dw/dz$  field in  $x$ - $z$  plane  
[Gayen & Sarkar, 10]





# 1.1 Generation of internal tides



# 1.1 Generation of internal tides

Impact of tidally-driven mixing on erosion and sediment transport

## The Shaping of Continental Slopes by Internal Tides

D. A. Cacchione,<sup>1</sup> L. F. Pratson,<sup>2</sup> A. S. Ogston<sup>3</sup>

The angles of energy propagation of semidiurnal internal tides may determine the average gradient of continental slopes in ocean basins ( $\sim 2$  to 4 degrees). Intensification of near-bottom water velocities and bottom shear stresses caused by reflection of semi-diurnal internal tides affects sedimentation patterns and bottom gradients, as indicated by recent studies of continental slopes off northern California and New Jersey. Estimates of bottom shear velocities caused by semi-diurnal internal tides are high enough to inhibit deposition of fine-grained sediment onto the slopes.

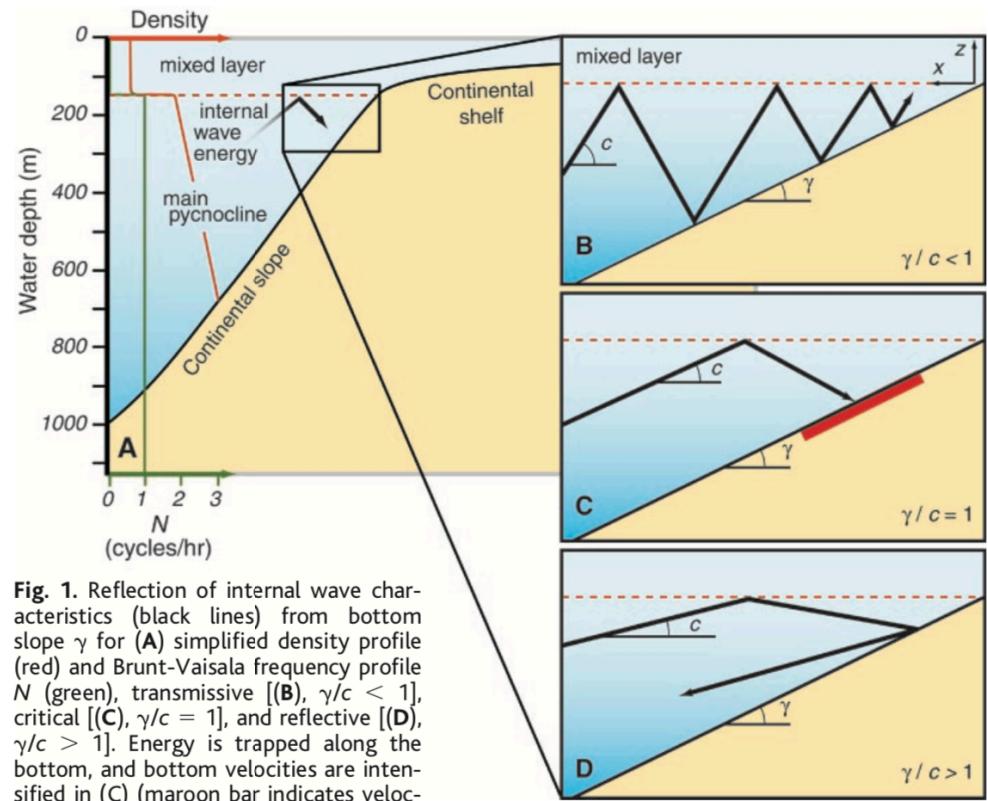
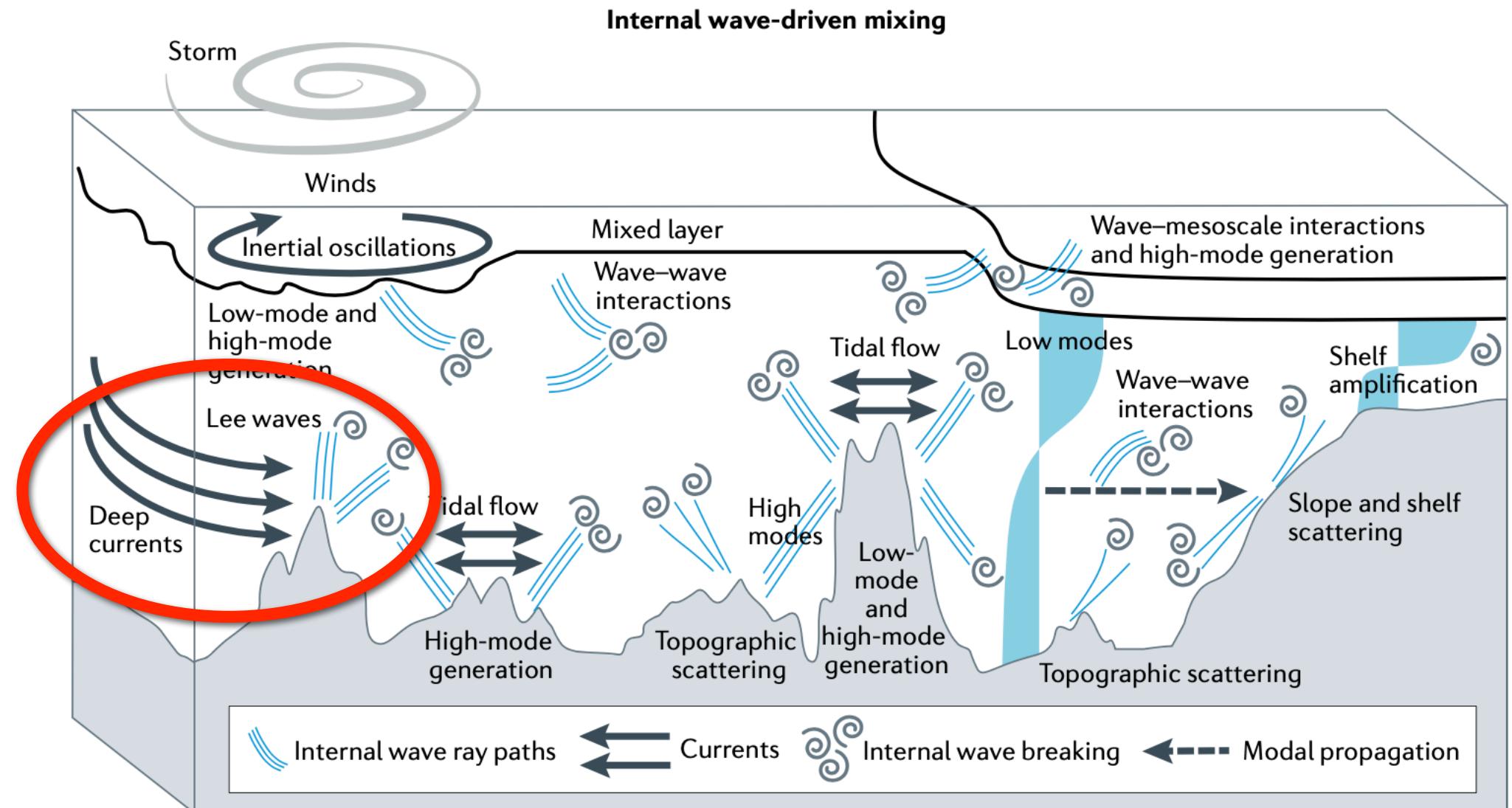


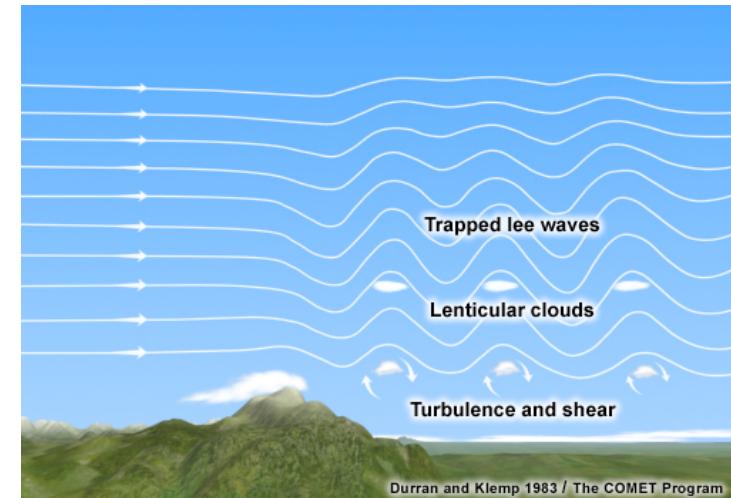
Fig. 1. Reflection of internal wave characteristics (black lines) from bottom slope  $\gamma$  for (A) simplified density profile (red) and Brunt-Vaisala frequency profile  $N$  (green), transmissive [(B),  $\gamma/c < 1$ ], critical [(C),  $\gamma/c = 1$ ], and reflective [(D),  $\gamma/c > 1$ ]. Energy is trapped along the bottom, and bottom velocities are intensified in (C) (maroon bar indicates velocity intensification); bottom velocities also increase upslope in (B) (10).

# 1.2 Generation of lee waves



# 1.2 Generation of Lee waves

- **Lee waves** (or mountain waves) can be seen in ocean and atmosphere

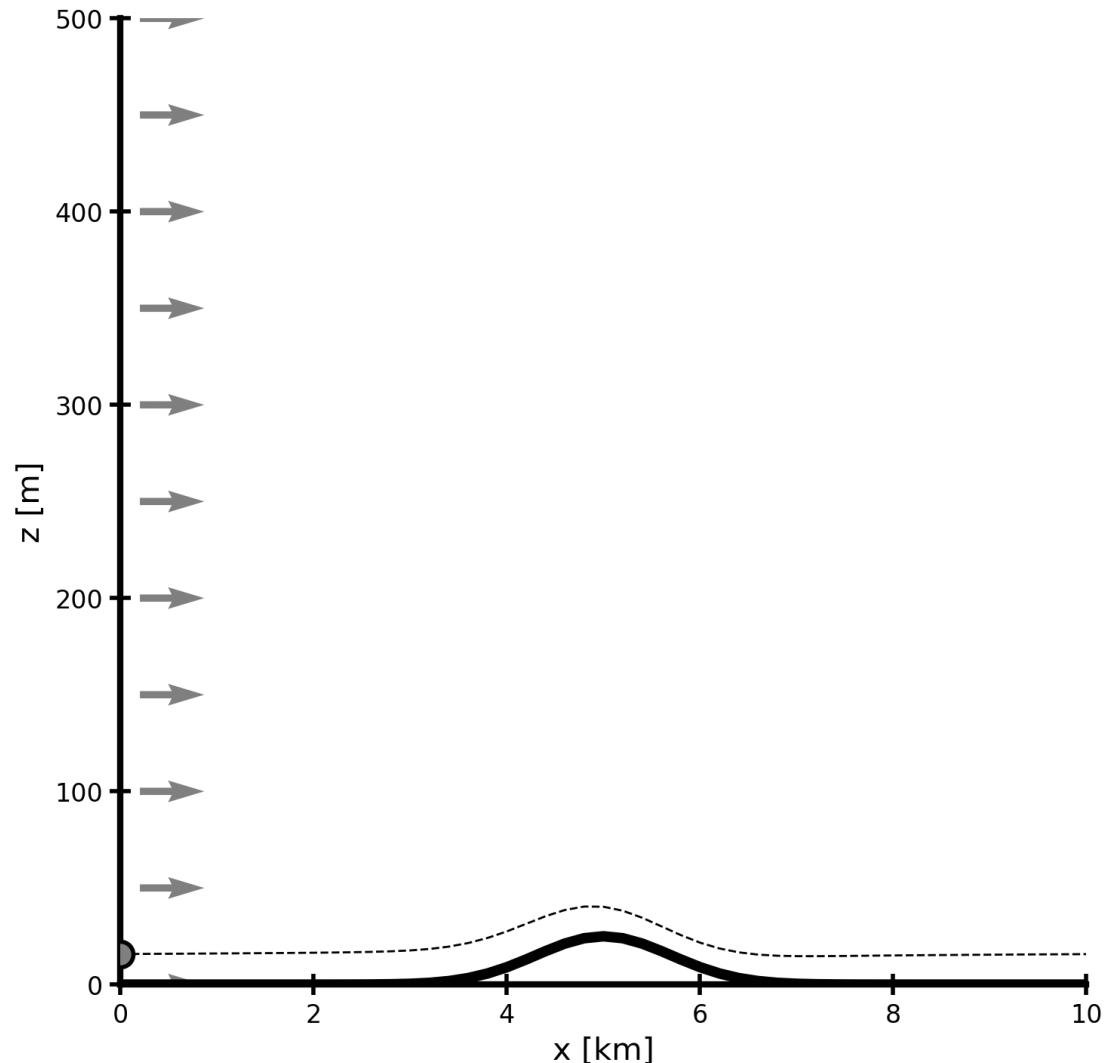


Modis, 23/11/2009 - South Atlantic – Sandwich Islands  
The lower atmosphere is drier – Downwind of the islands the waves are seen when the air goes up, condense and form clouds

## 1.2 Generation of Lee waves

Ingredients for the propagation of a lee wave :

- ✓ A current...
- ✓ ...flowing over a topography

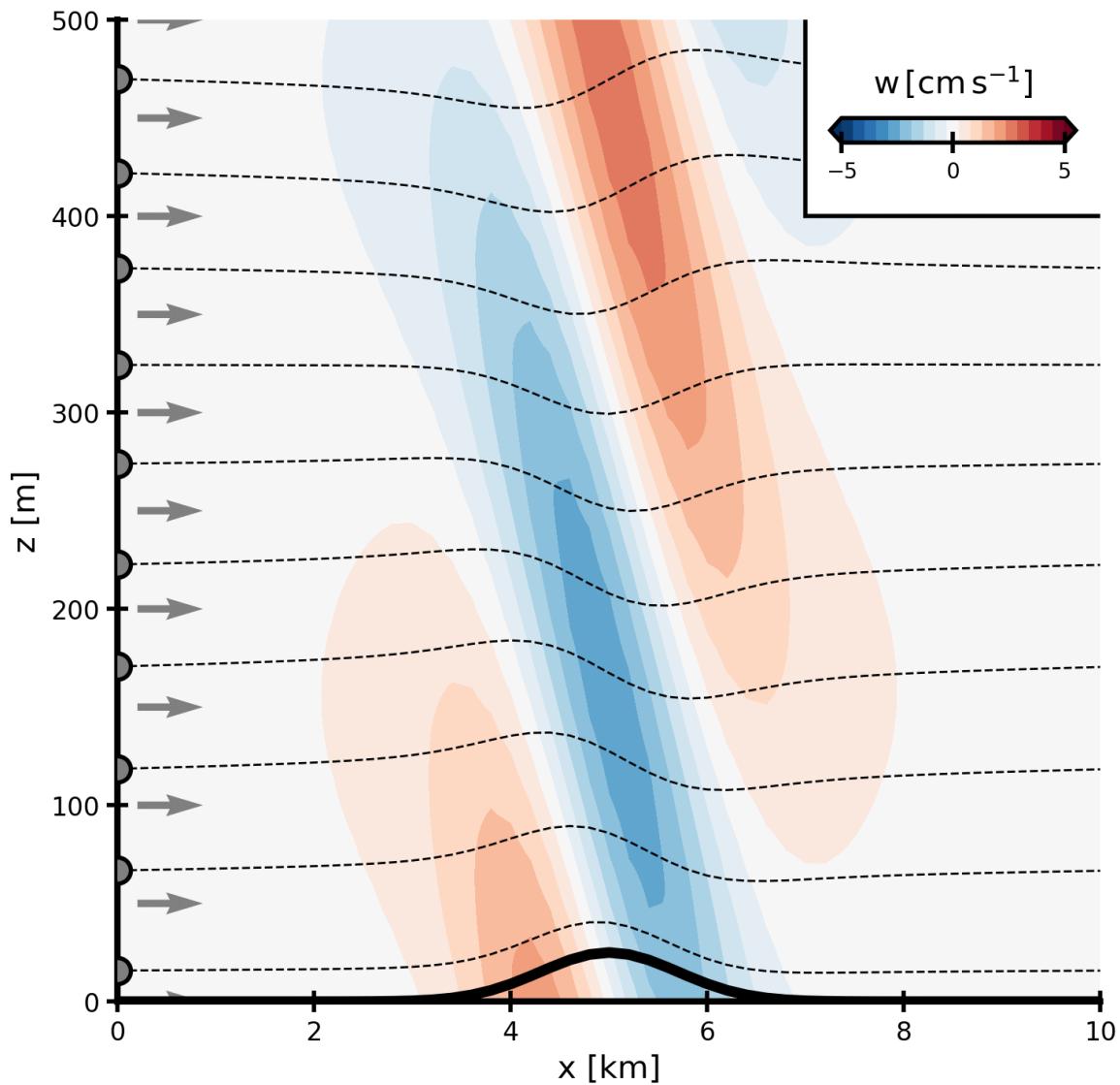


## 1.2 Generation of Lee waves

Ingredients for the propagation of a lee wave :

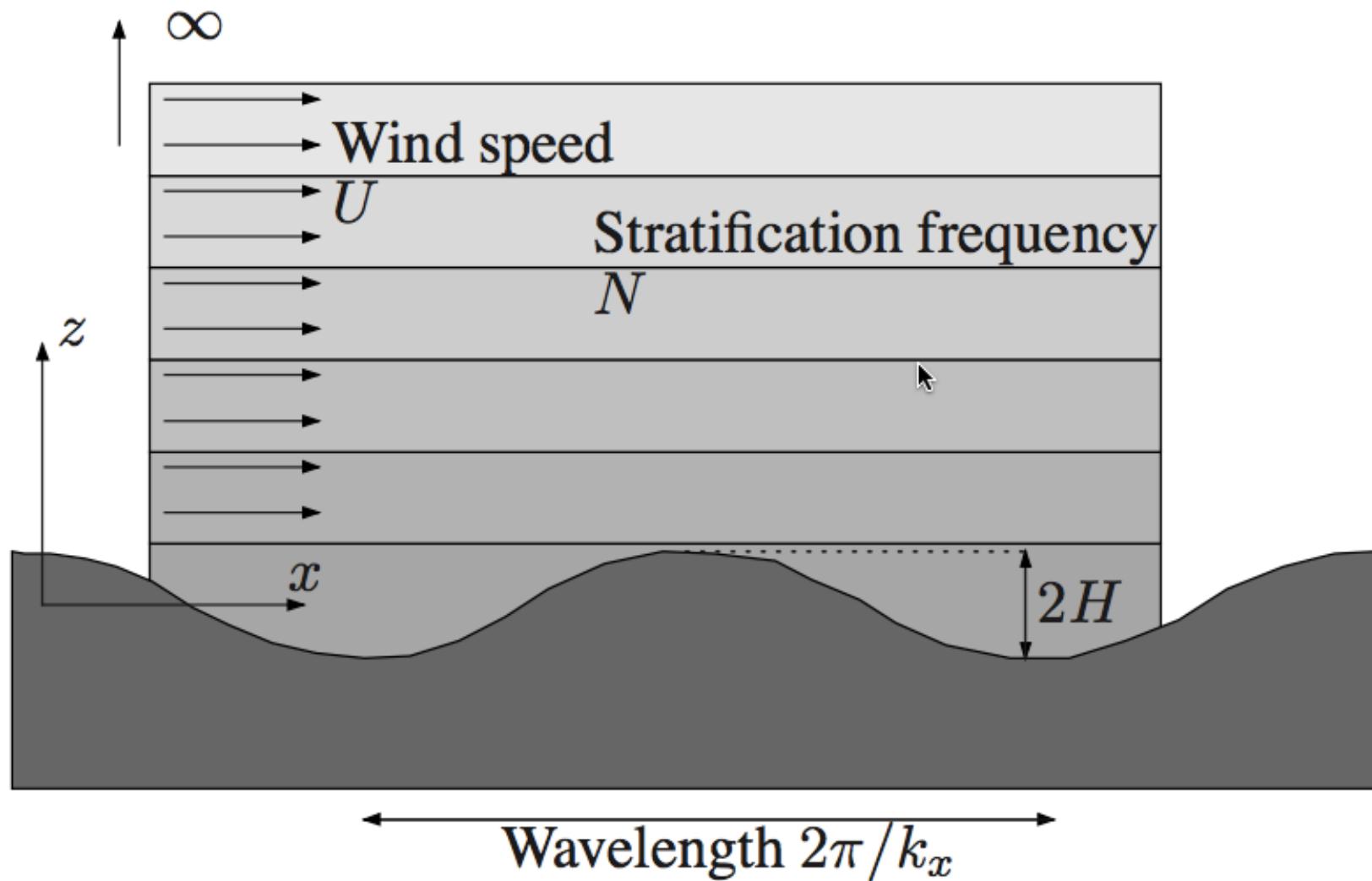
- ✓ A current...
- ✓ ...flowing over a topography
- ✓ A stratification

→ *Propagation of an internal wave forced by the topography*



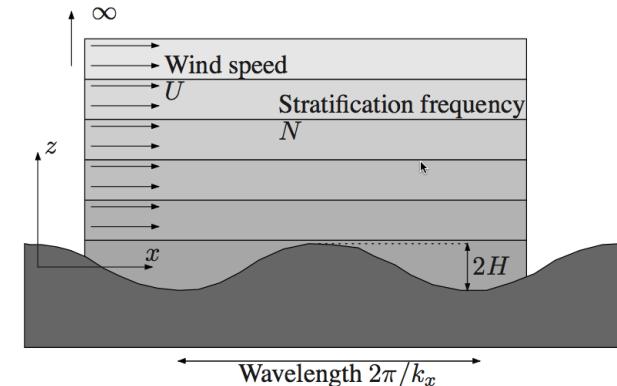
## 1.2 Generation of Lee waves

- Internal waves generated at horizontal boundary:



# 1.2 Generation of Lee waves

- Topography  $z_b = H \sin(k_x x)$
- In the frame moving at constant speed  $U$ ,  $z_b = H \sin(k_x(x + Ut))$
- We can identify the frequency of the motion:  $\omega = -k_x U$
- Recall the general equation for internal waves:  $(w_{xx} + w_{yy} + w_{zz})_{tt} + f^2 w_{zz} + N^2(w_{xx} + w_{yy}) = 0$



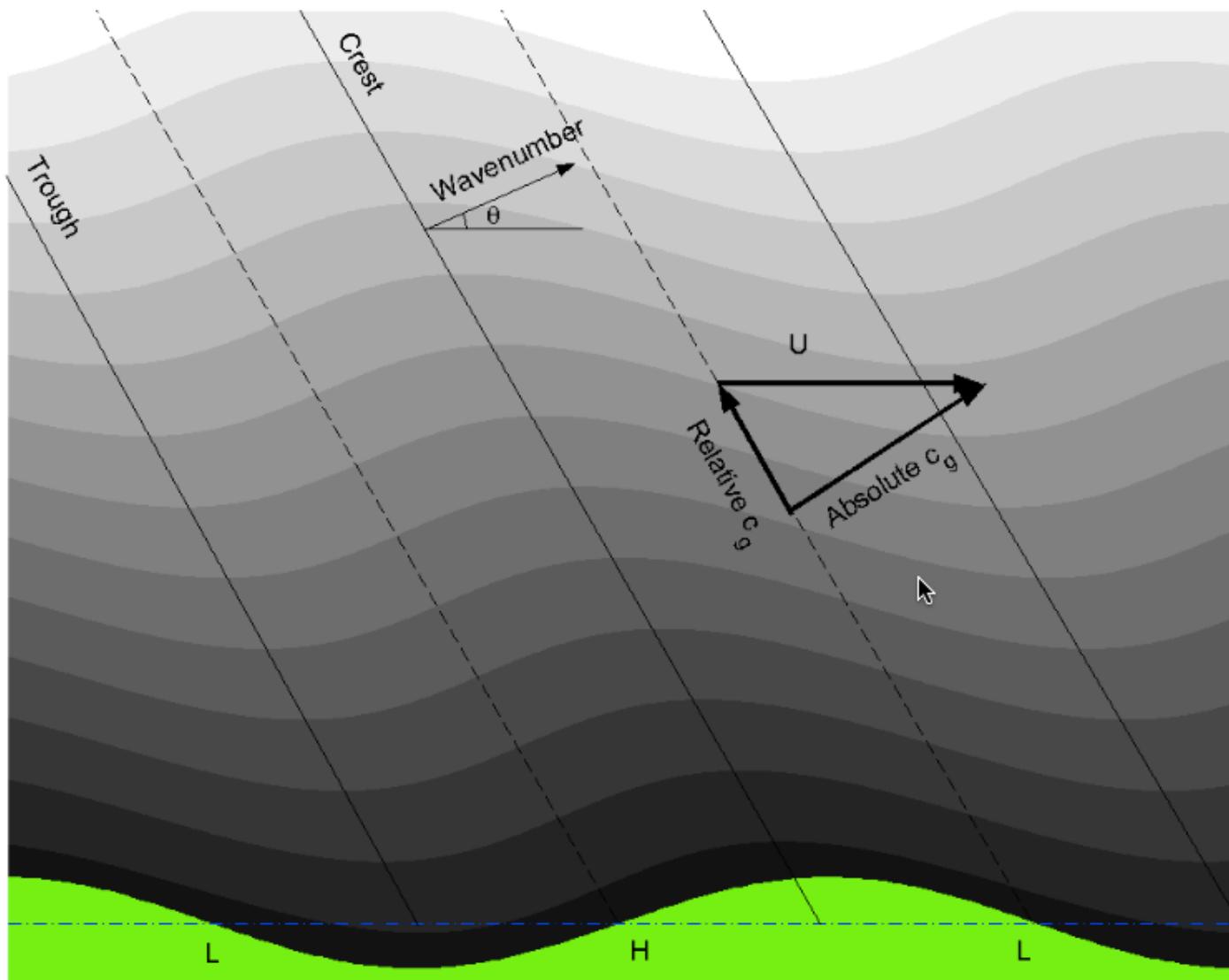
Activity: Solve the equation for linear internal waves in the  $x - z$  plane assuming:

- Earth's rotation can be neglected  $f = 0$
- The topography is small compared to typical scales of stratification  $H \ll U/N$

1. Write the bottom boundary condition for the vertical velocity  $w(z_b) \approx w(0)$
2. Solve for a solution in the form  $w = w_0 \exp(i(k_x x + k_z z - \omega t))$

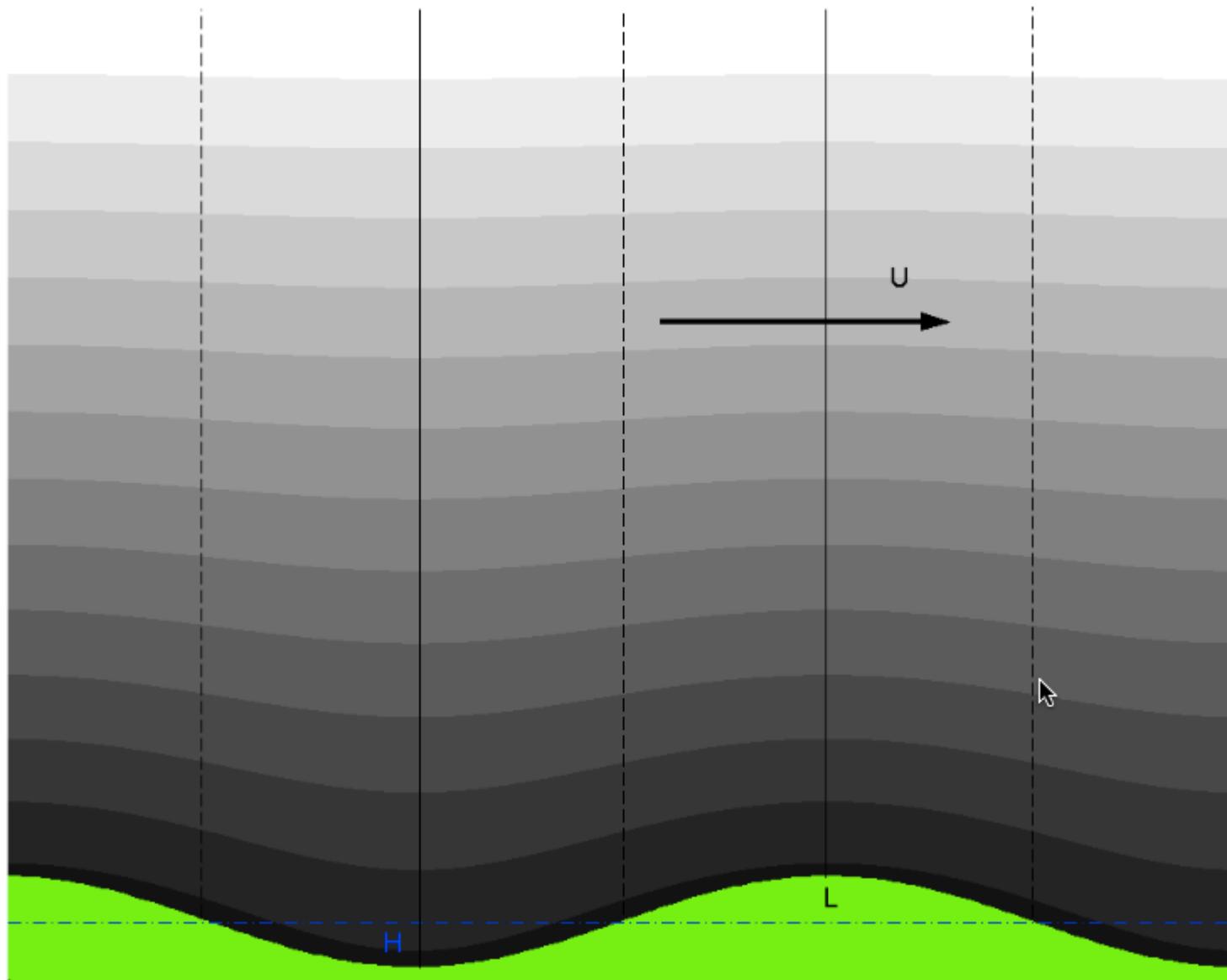
# 1.2 Generation of Lee waves

- Case 1 = Radiating waves



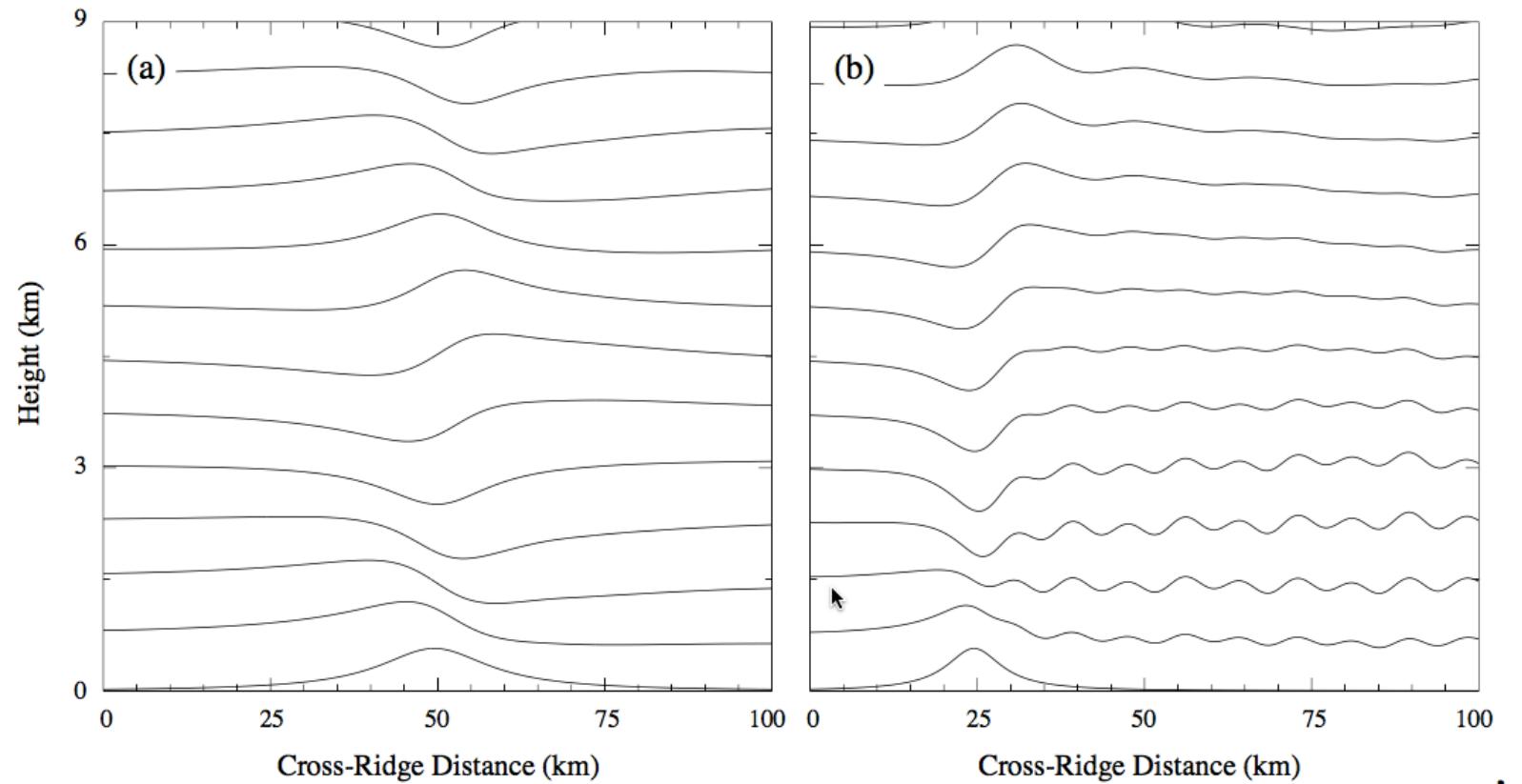
## 1.2 Generation of Lee waves

- Case 2 = Trapped waves



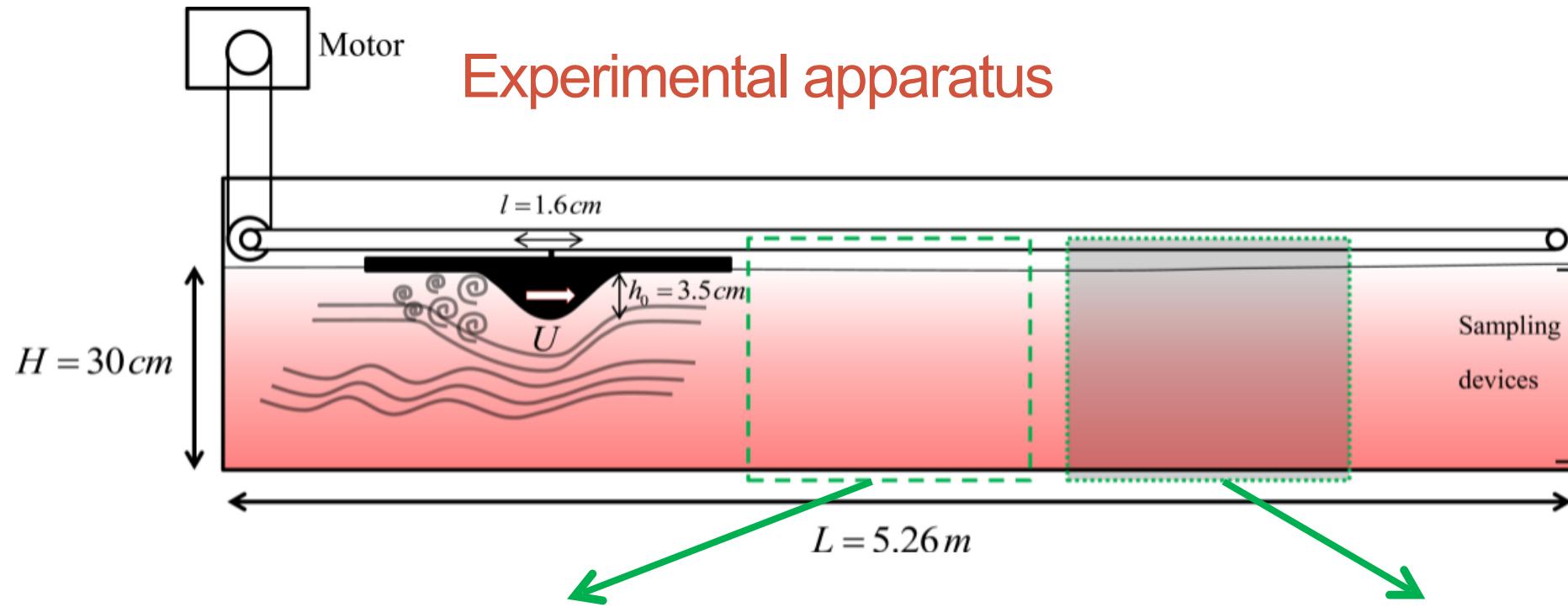
## 1.2 Generation of Lee waves

- General topography

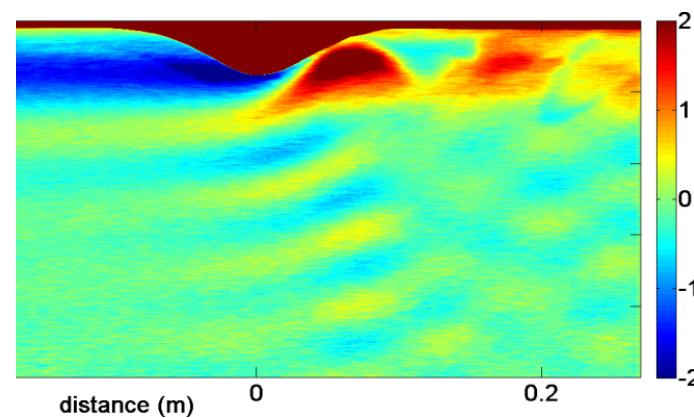
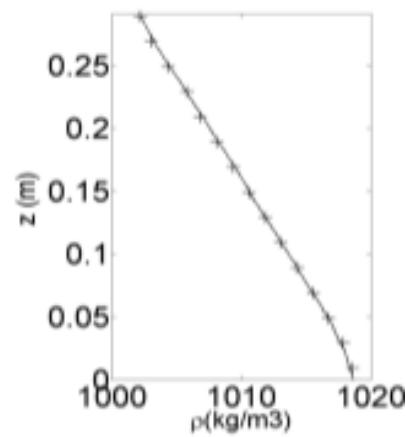


Streamlines in steady airflow over an isolated ridge as predicted by linear theory

# 1.2 Generation of Lee waves - Experiments



**Light attenuation technique**  
High resolution density measurements

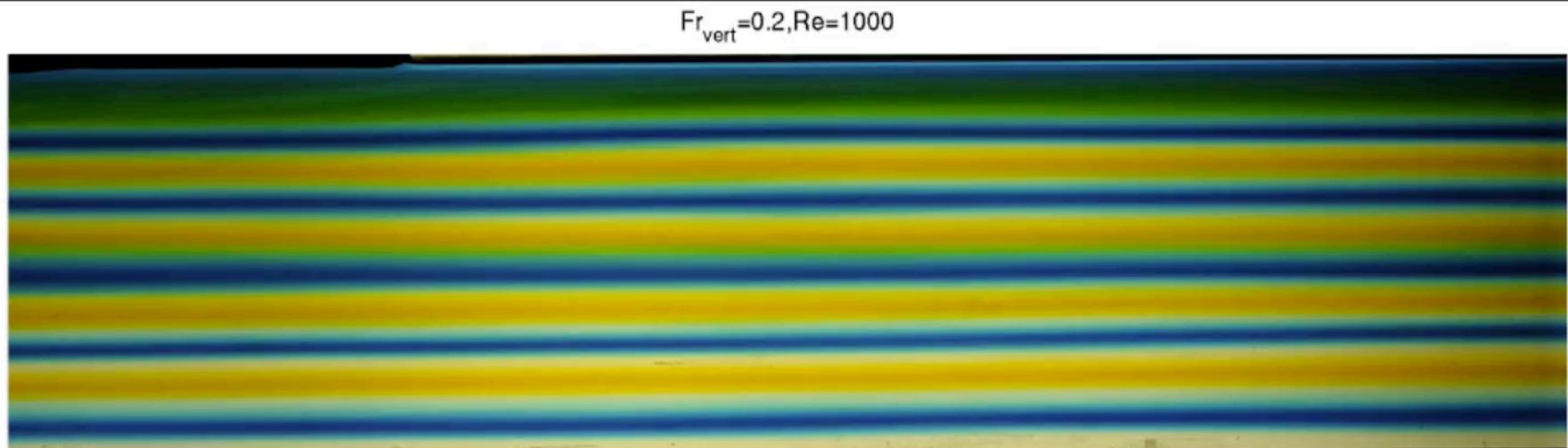


**Shadowgraph pictures**  
Direct observation of mixing



[Y. Dossmann]

## 1.2. Generation of Lee waves - Experiments

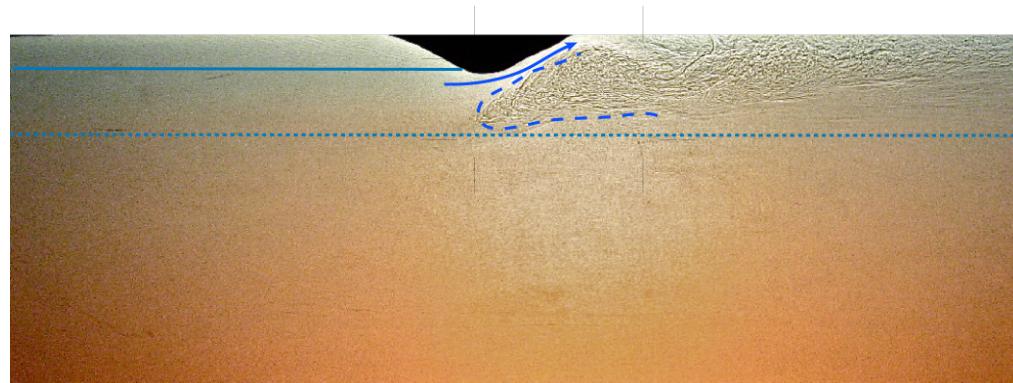


Froude number : Flow regime

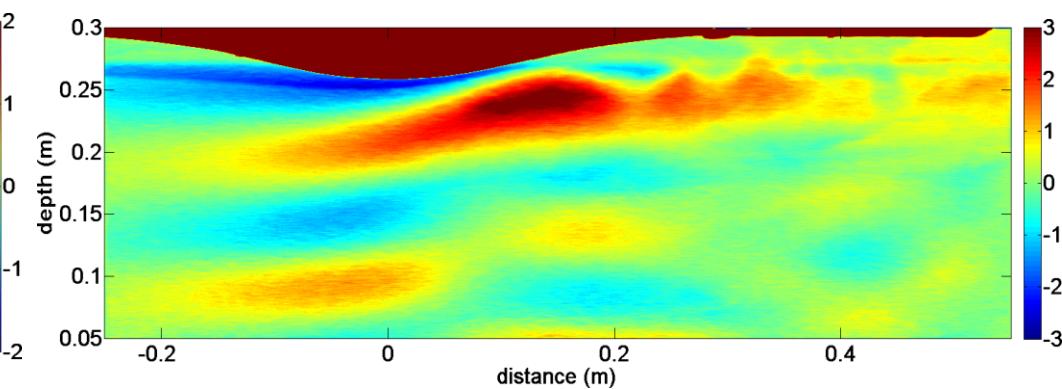
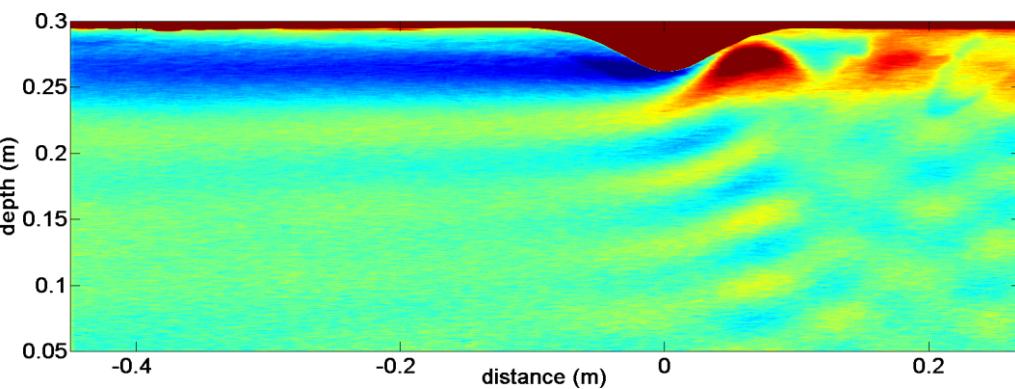
Reynolds number : Turbulence regime

# 1.2. Generation of Lee waves - Experiments

Fr=0.2; Re=500



Fr=0.4; Re=1700



## 1.2. Generation of Lee waves - Experiments

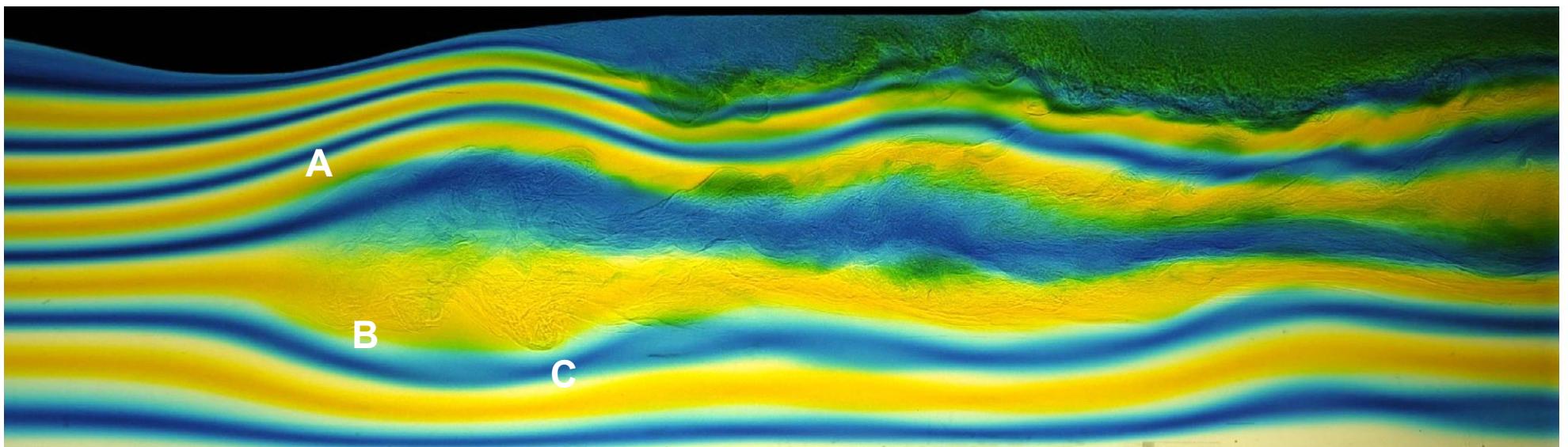
$Fr_{vert} = 1, Re = 5000$



[Y. Dossmann]

## 1.2. Generation of Lee waves - Experiments

$Fr=1$  ;  $Re=3500$

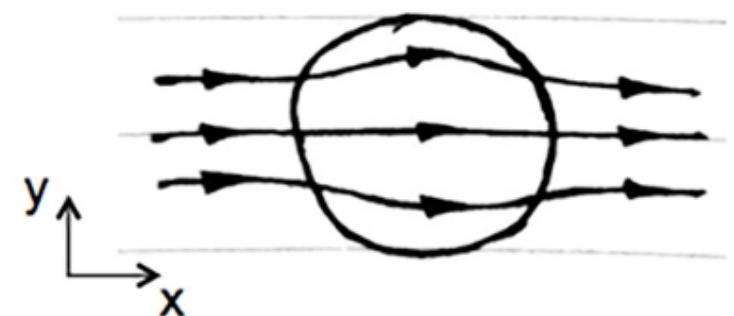
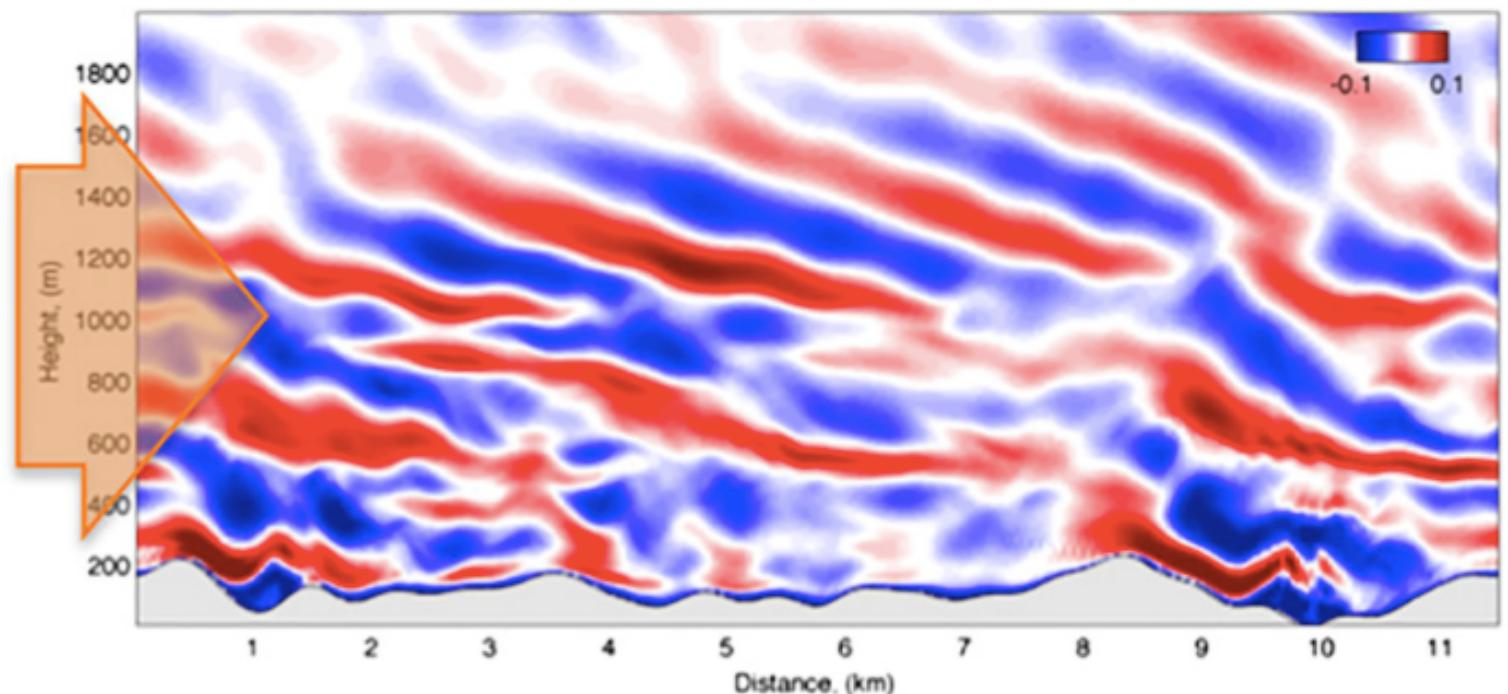


- A: Unstable Lee waves
- B: Narrow accelerating jet
- C: Periodic turbulent patches

[Y. Dössmann]

## 1.2. Generation of Lee waves

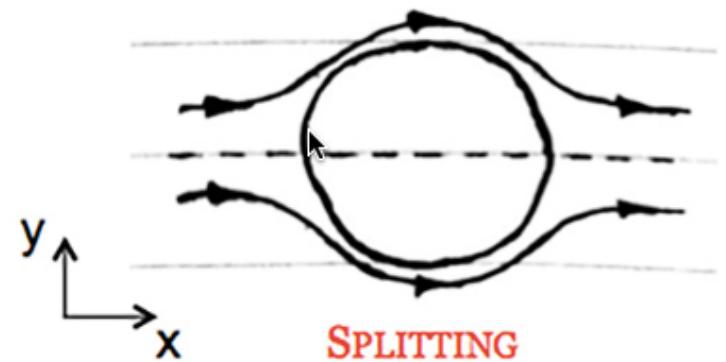
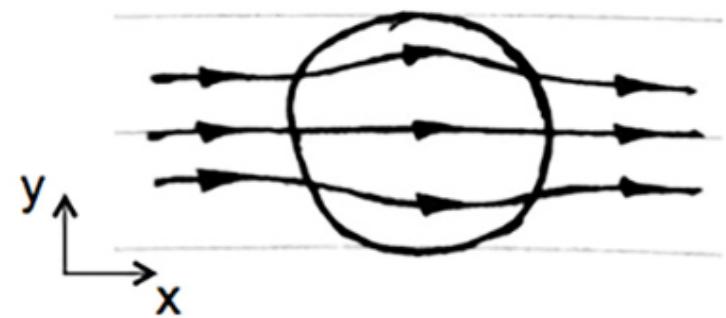
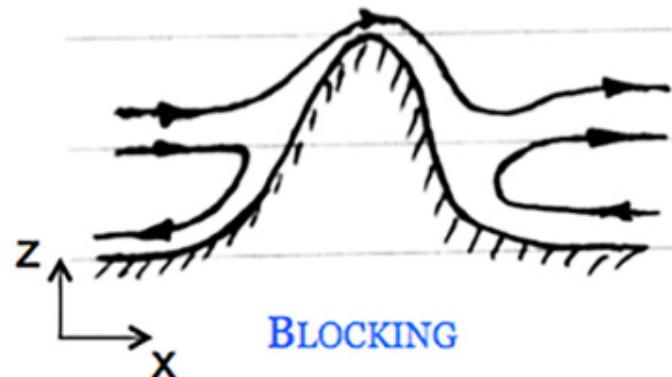
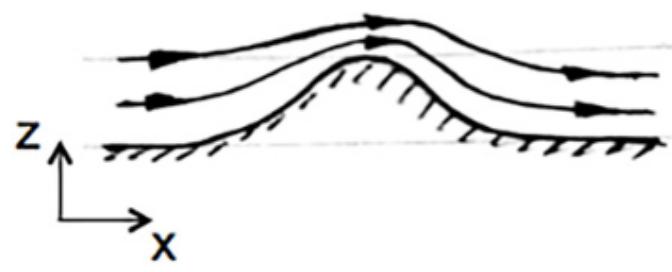
Limitation of the 2d linear theory:



[Nikurashin et al., 2014]

## 1.2. Generation of Lee waves

Limitation of the 2d theory



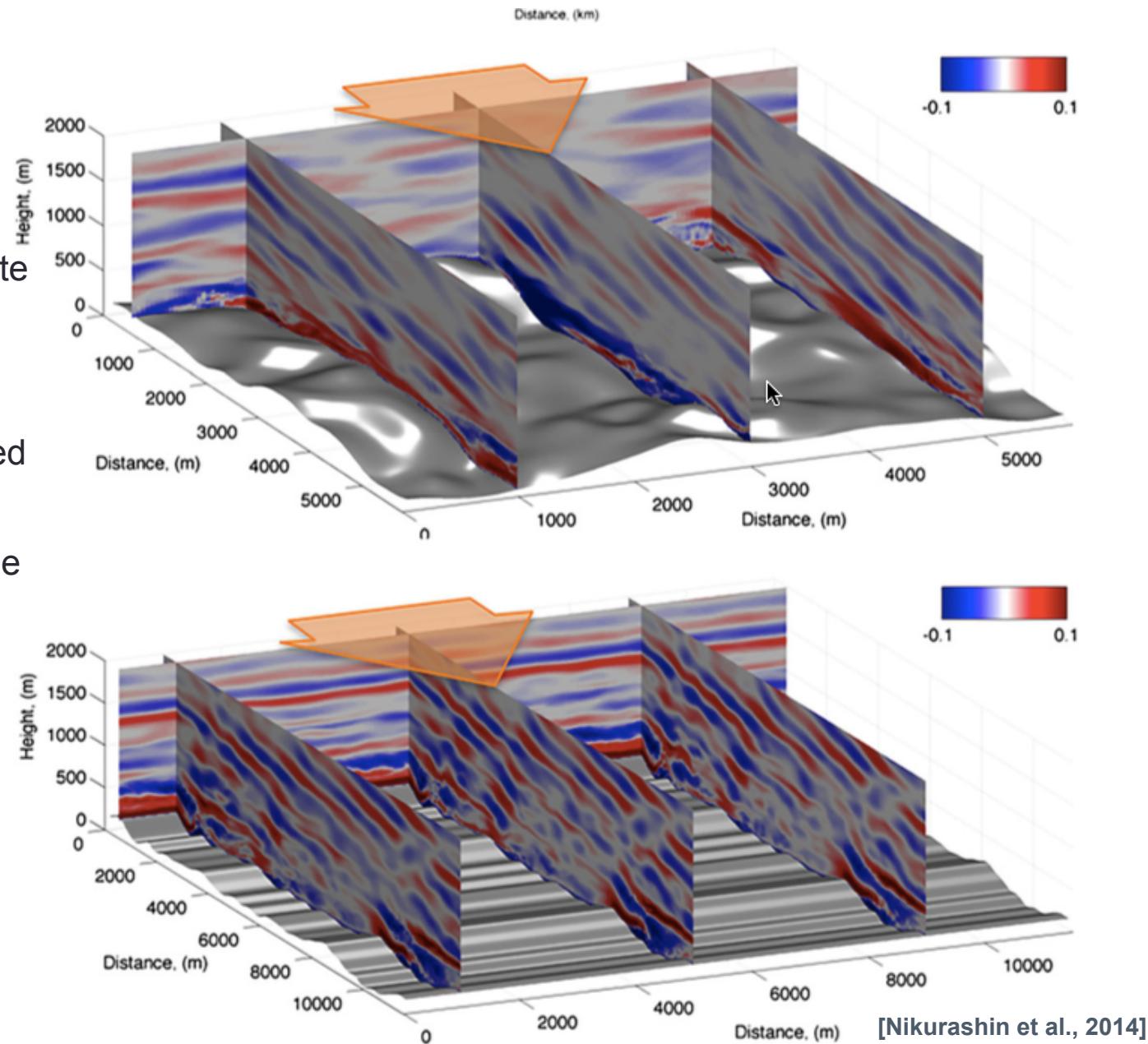
[Nikurashin et al., 2014]

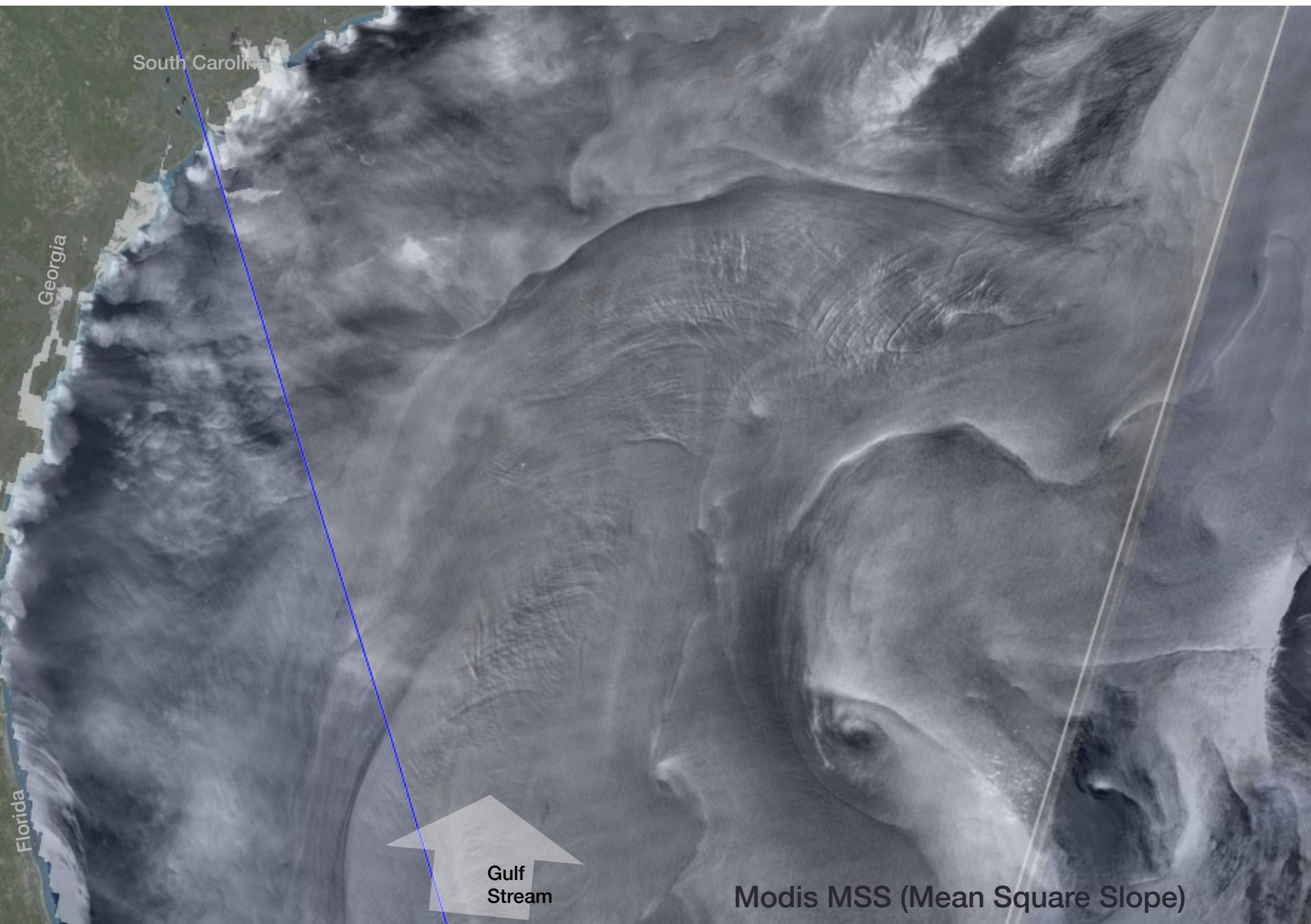
# 1.2. Generation of Lee waves

## Limitation of the 2d theory

Internal wave generation at 3D, finite bottom topography is reduced compared to the two-dimensional case.

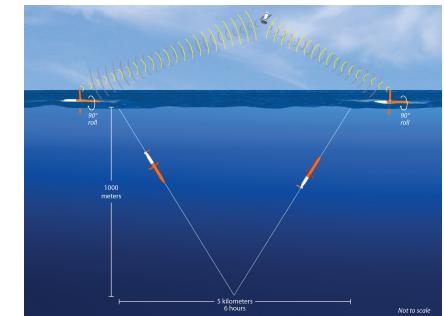
The reduction is primarily associated with finite-amplitude bottom topography effects that suppress vertical motions and thus reduce the amplitude of the internal waves radiated from topography.





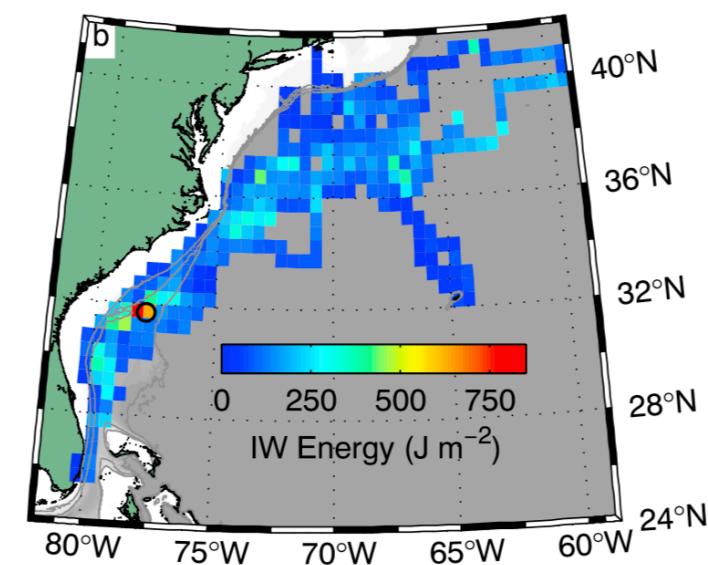
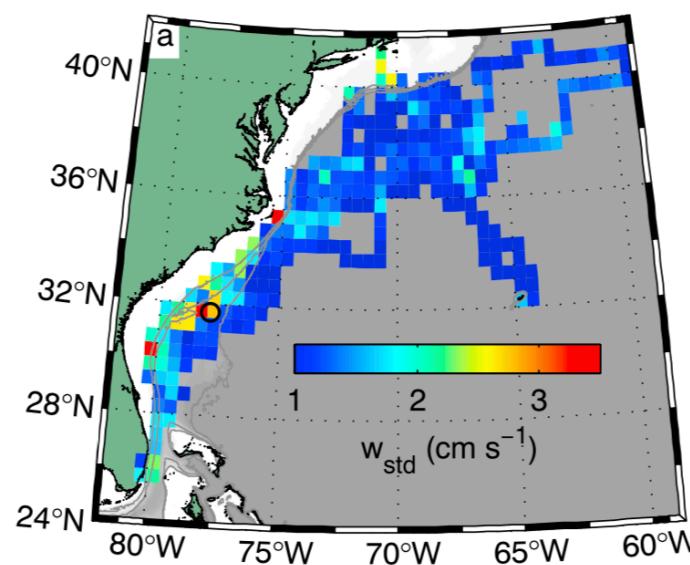
## 1.2. Generation of Lee waves - Real example

- Internal waves with vertical velocities exceeding  $0.1 \text{ m s}^{-1}$  are found over the Charleston Bump.



internal wave activity in the GS

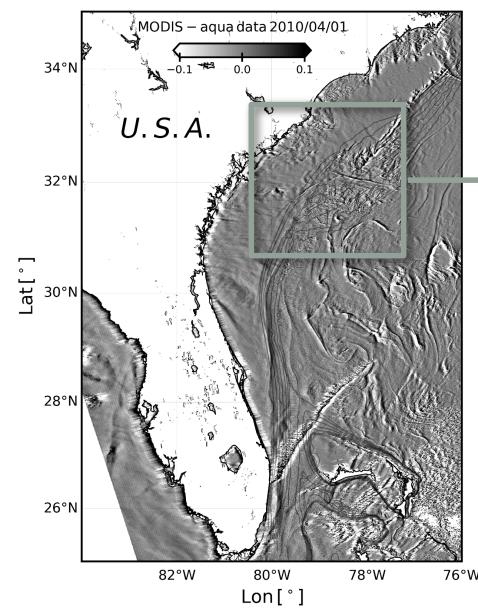
Glider data



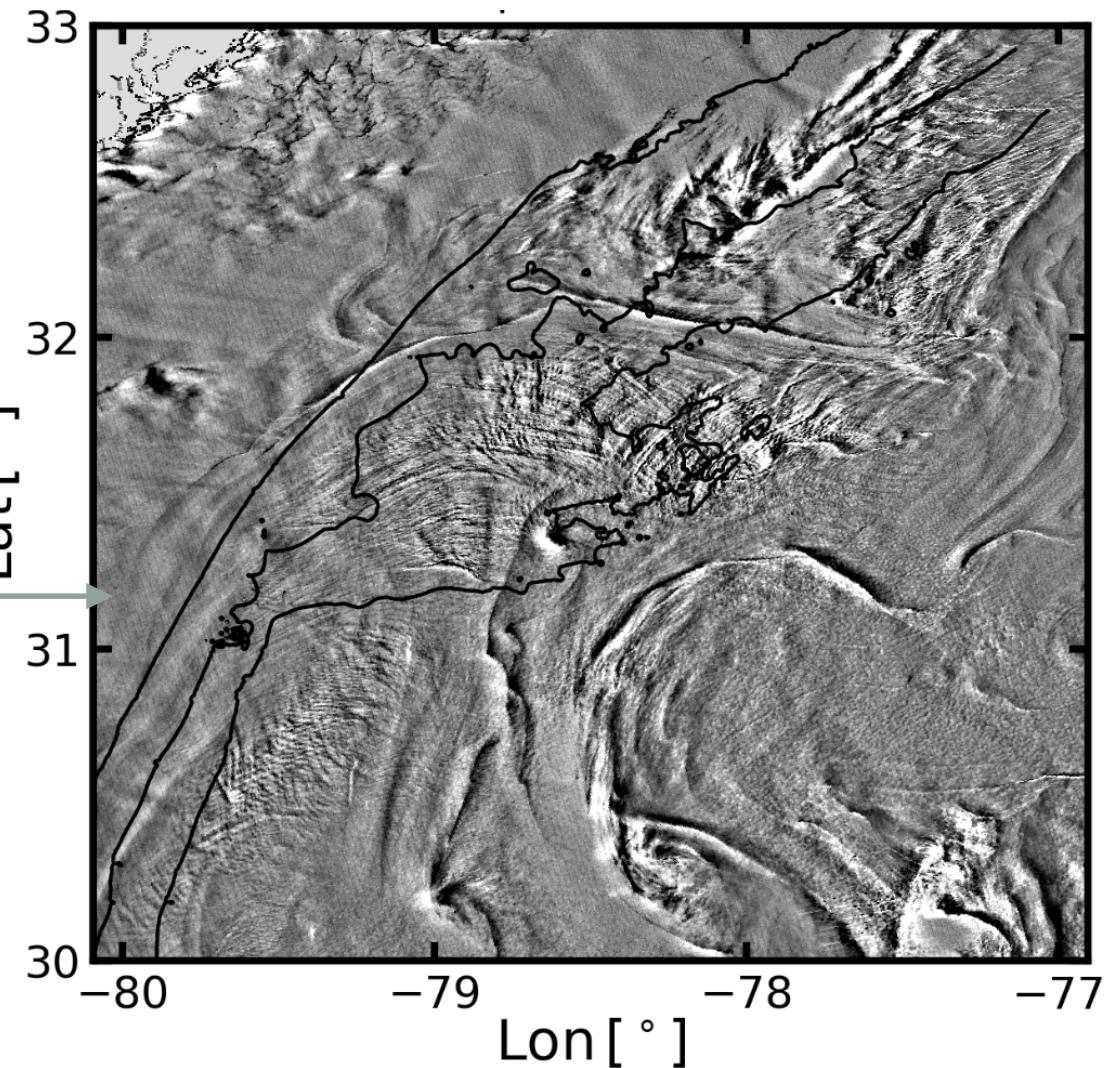
(a) Standard deviations of vertical velocities from individual glider dives averaged in  $0.5^\circ \times 0.5^\circ$  boxes. (b) Vertically integrated internal wave energy [Todd, 17]

## 1.2. Generation of Lee waves - Real example

- Internal waves can be observed using synthetic aperture radars (SAR) or Sun-glitter images through their surface roughness signature.



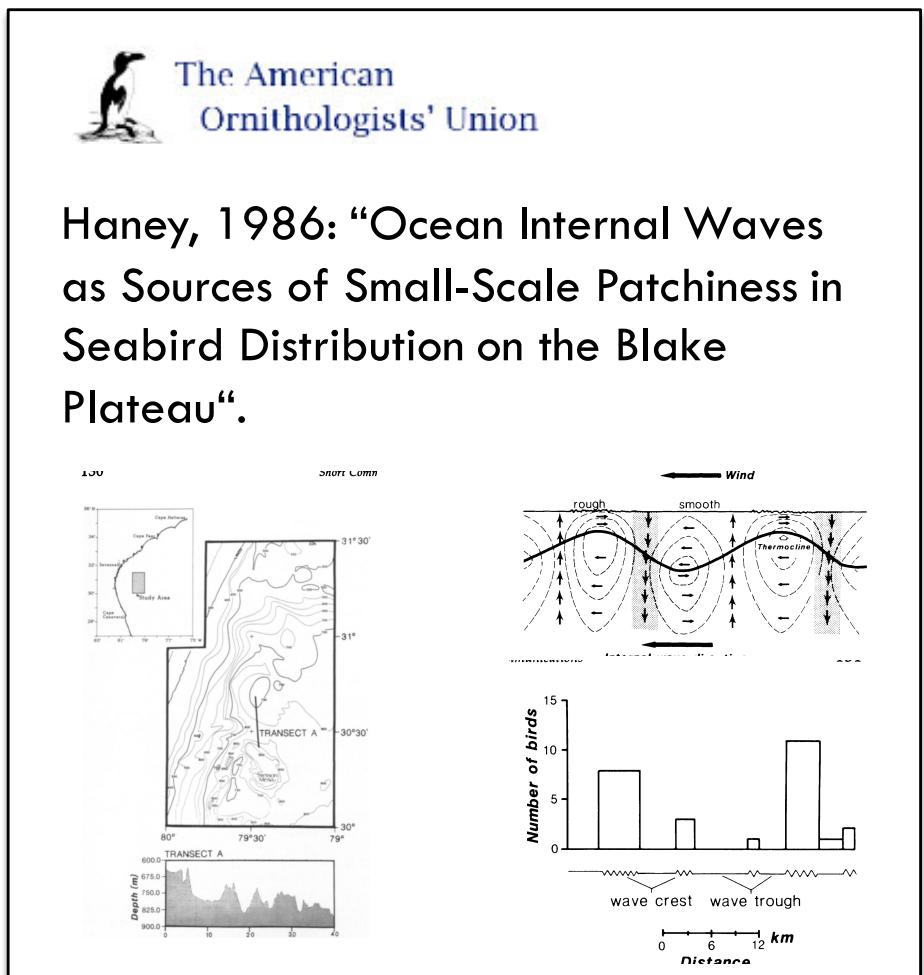
Observations of Lee Waves in the Gulf Stream



## 1.2. Generation of Lee waves - Real example

- They were also previously inferred from patchiness in seabirds distribution:

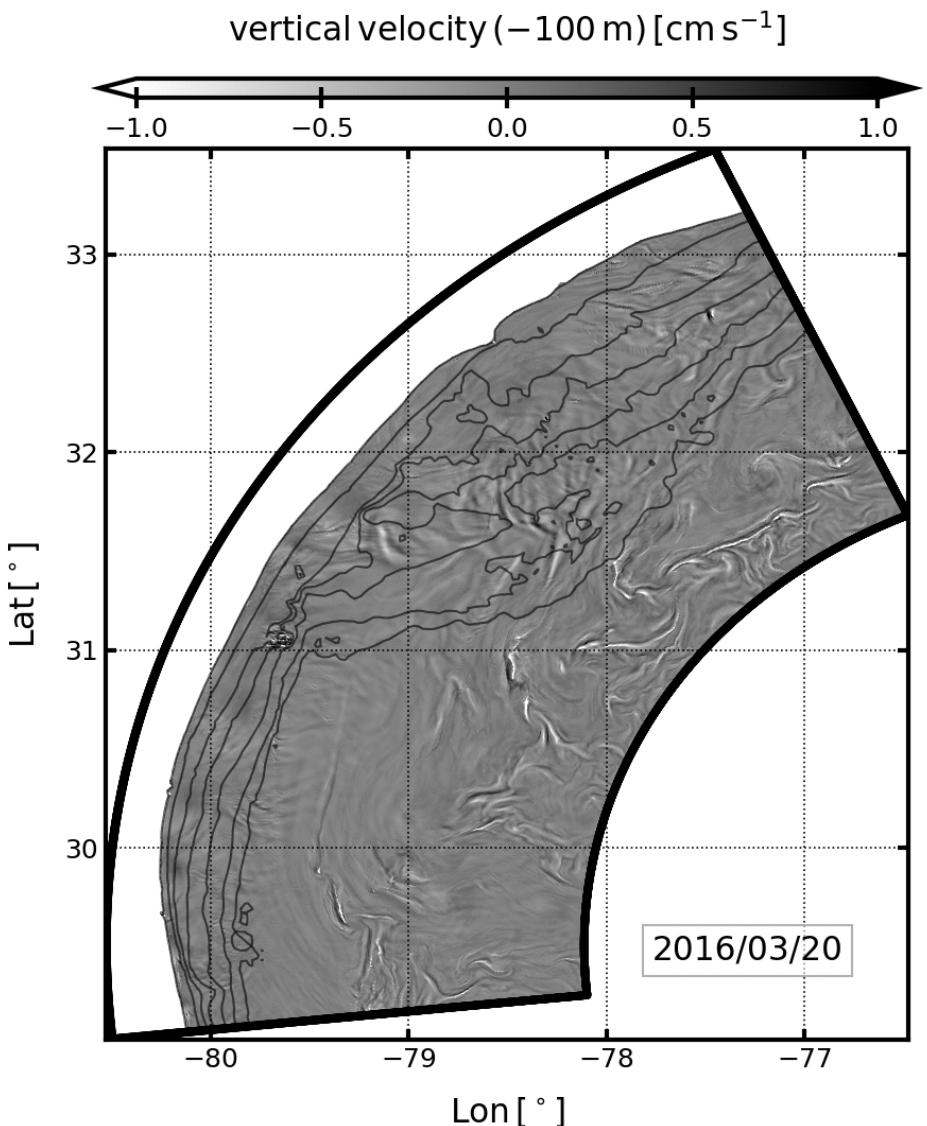
### Observations of Lee Waves in the Gulf Stream



## 1.2. Generation of Lee waves - Real example

- We performed a ROMS/CROCO nest  $dx = 200$  m,
- $Nx \times Ny \times Nz = 1024 \times 2048 \times 128$
- Climatological forcings (monthly winds, no tides)
- 

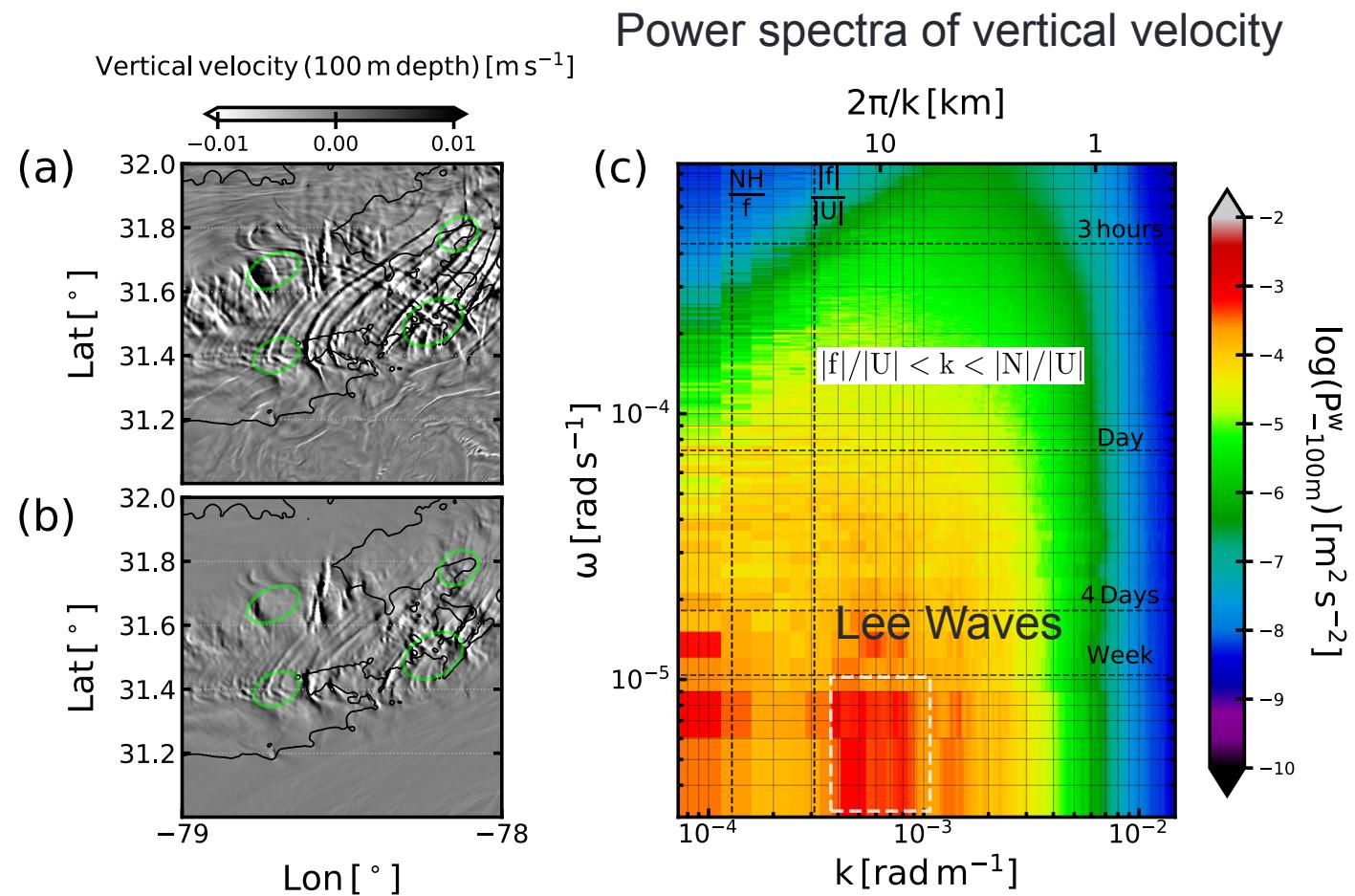
Simulation of the Gulf Stream



# 1.2. Generation of Lee waves - Real example

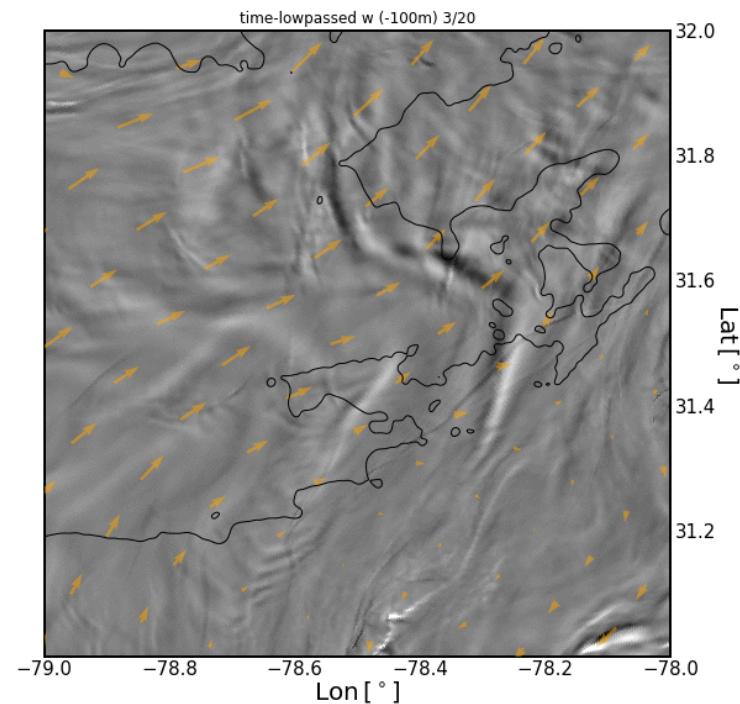
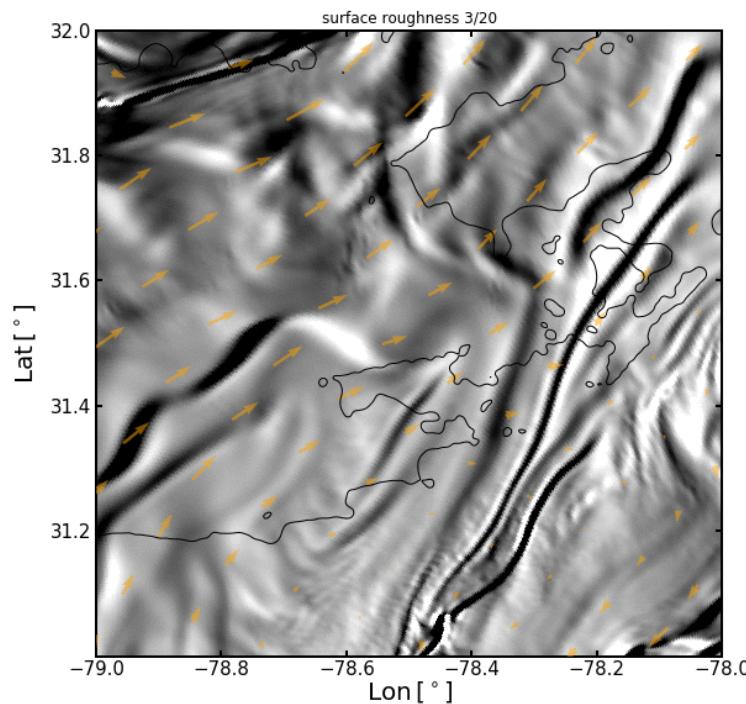
- We can isolate patterns related to Lee Waves by looking at time-low passed vertical velocities below the thermocline:

Lee Waves in the  
Gulf Stream



## 1.2. Generation of Lee waves - Real example

- We can isolate patterns related to Lee Waves by looking at time-low passed vertical velocities below the thermocline:



# 1.2. Generation of Lee waves - Real example

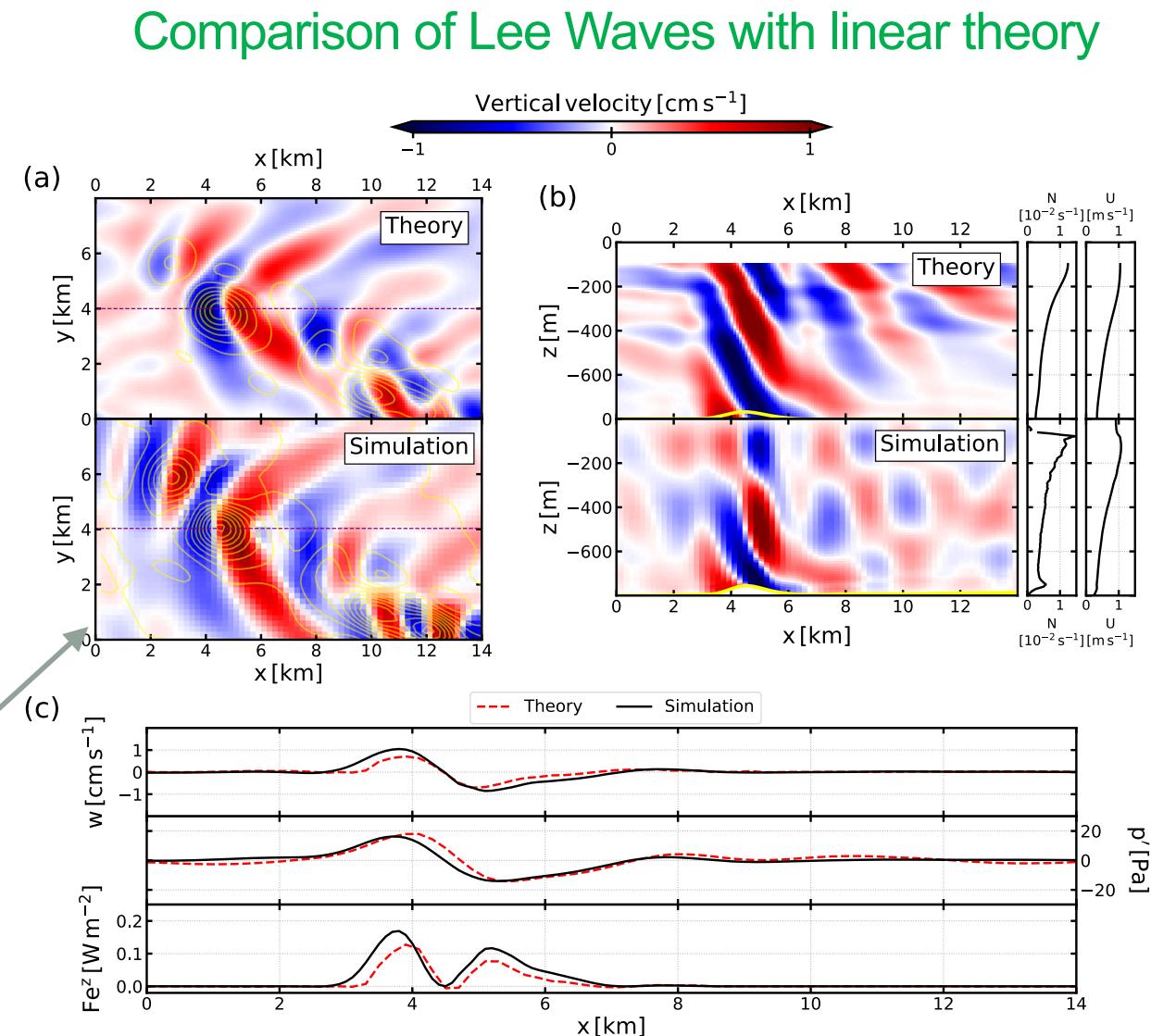
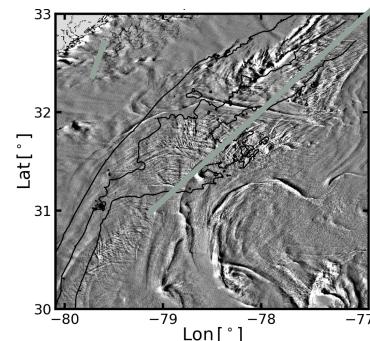
- 3D theoretical prediction of hydrostatic lee waves is obtained by numerically solving:

$$\partial_{zz} \tilde{\eta}(k, m, z) + n^2 \tilde{\eta}(k, m, z) = 0$$

- under the WKBJ approximation with the dispersion relation:

$$n^2(k, m, z) = k^{-2} (k^2 + m^2) \left( \frac{N^2}{U^2} + \frac{\partial_{zz} U}{U} \right)$$

- and radiation condition at the top of the domain and a Dirichlet condition at the bottom.



# 1.2. Generation of Lee waves - Real example

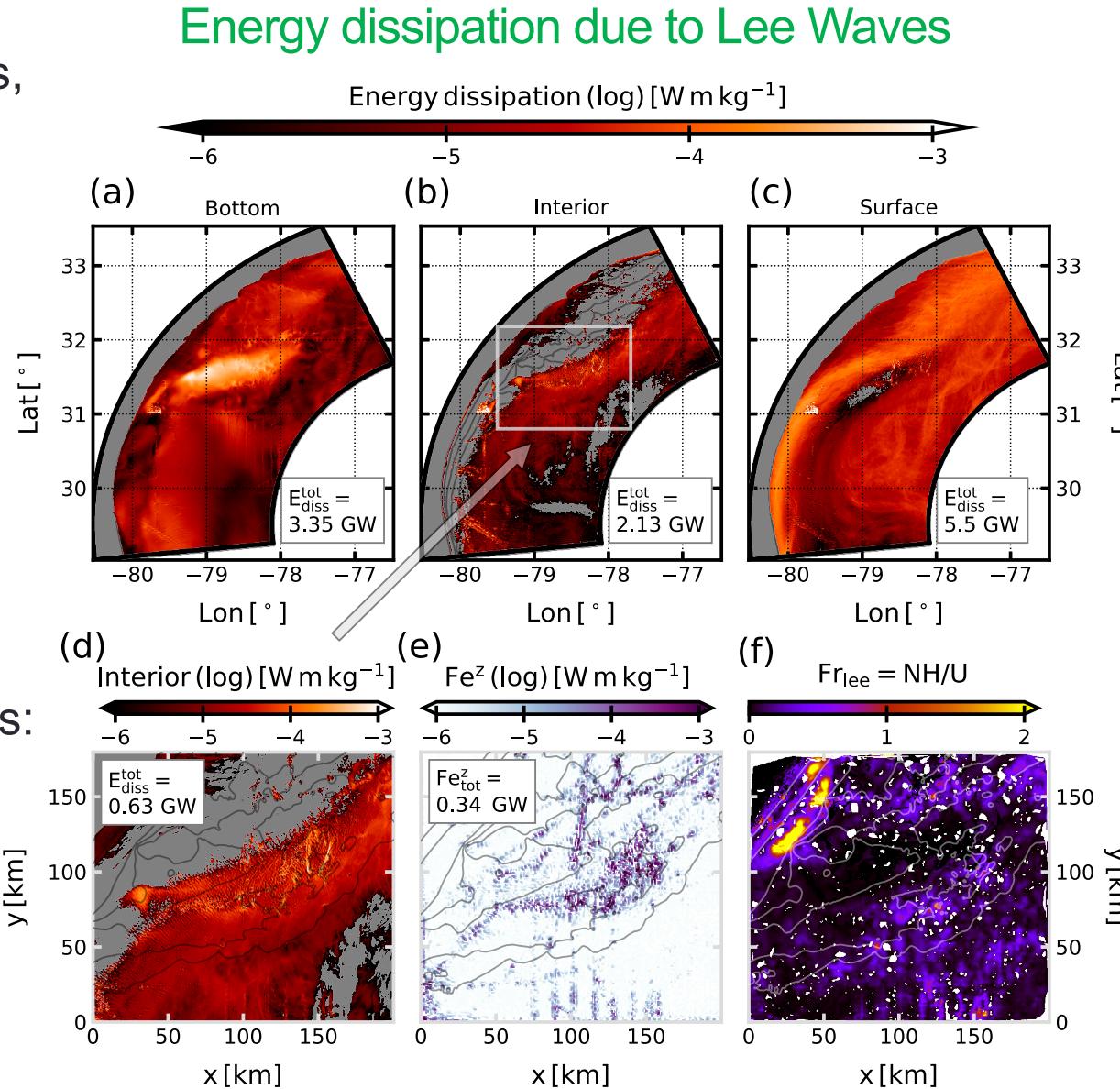
- A significant amount of energy ( $O(1)$  GW) is being dissipated in the interior of the fluid by the lee waves, including at locations where these waves may exhibit nonlinearities.
- The lee wave energy flux ( $Fe^z = p'w$ ) computed using analytical formulation of Nikurashin and Ferrari (2011) yields a total conversion of  $O(0.3)$  GW.

$$E = \frac{\rho_0 |\mathbf{U}|}{2\pi} \int_{|f|/|\mathbf{U}|}^{N/|\mathbf{U}|} P_*(k) \sqrt{N^2 - |\mathbf{U}|^2 k^2} \sqrt{|\mathbf{U}|^2 k^2 - f^2} dk,$$

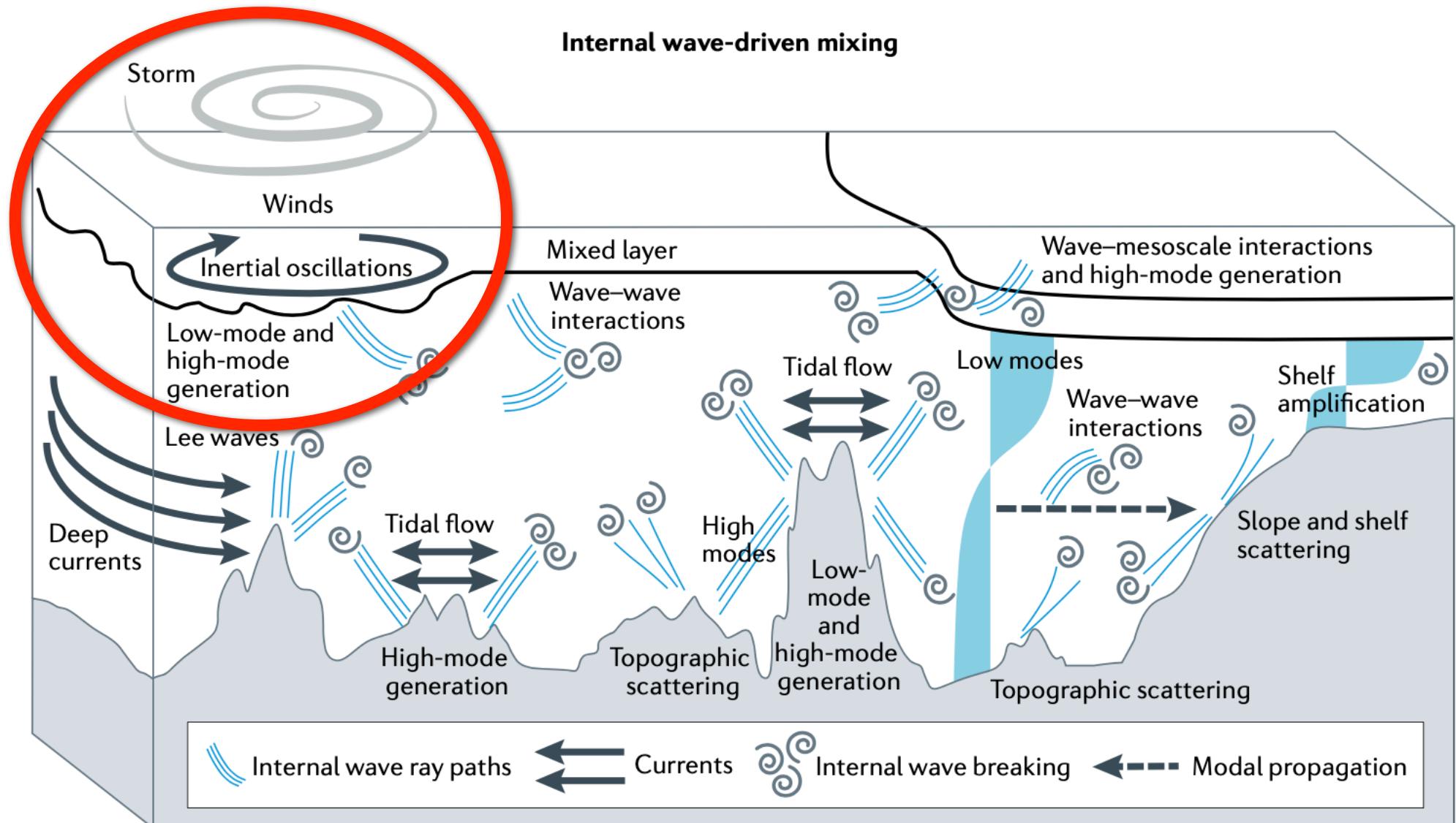
- The wave drag can be computed as:

$$\text{Form Drag} = \rho_0 \sqrt{(u'w)^2 + (v'w)^2}$$

- It amounts to  $O(1)$  GN and represents about 20% of the form drag exerted by the whole Charleston Bump.

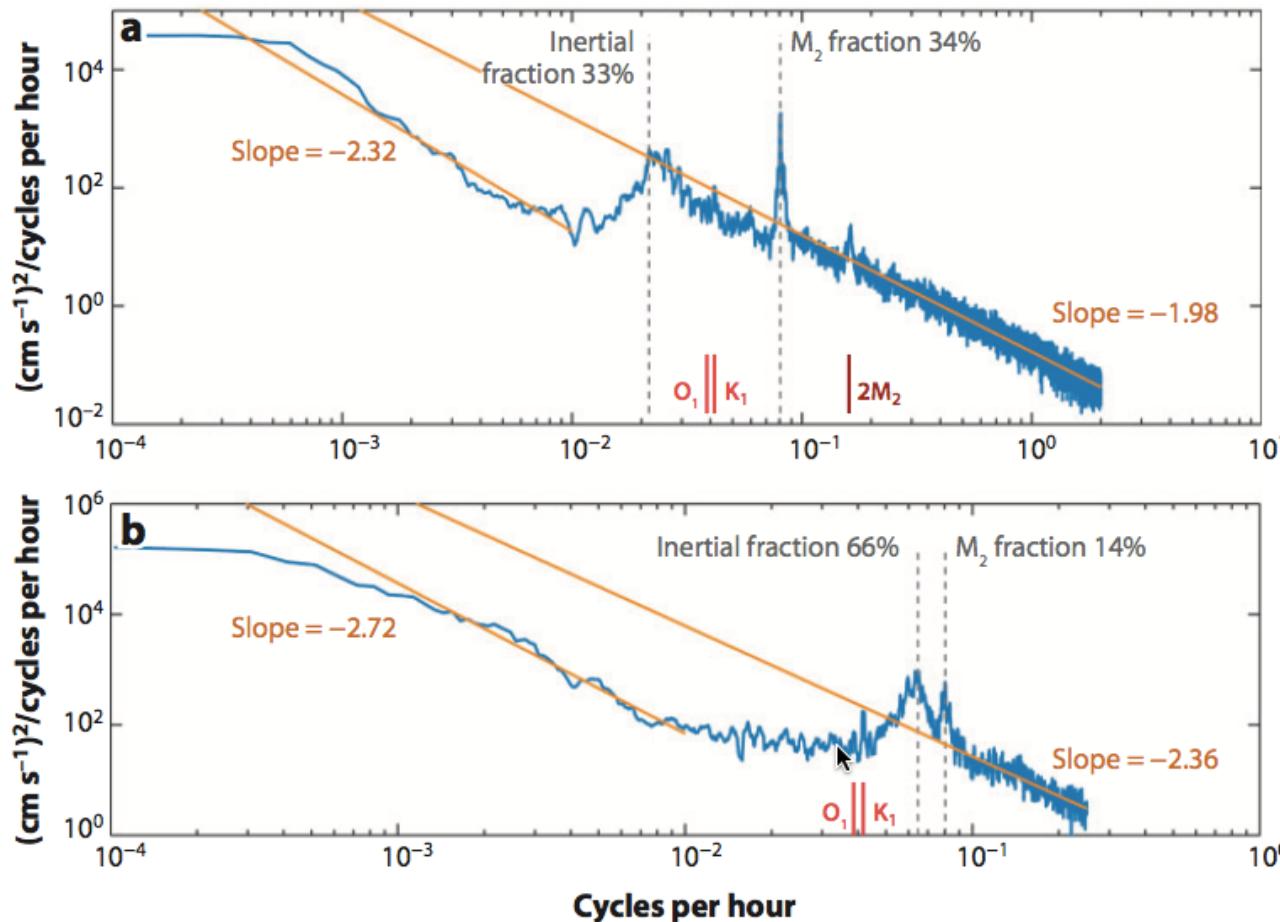


# 1.3 Generation of near-inertial waves



# 1. Internal waves generation

- Winds generate mostly near-inertial waves (frequency close to  $f$ )
  - = The near-inertial peak contains about half the kinetic energy in the internal wave spectrum



(a) Kinetic energy estimate for an instrument in the western North Atlantic near  $15^\circ\text{N}$  at 500 m. (b) Power density spectral estimate from a record at 1000 m at  $50.7^\circ\text{S}$ ,  $143^\circ\text{W}$ , south of Tasmania in the Southern Ocean

## 1.3. Generation of Near-Inertial waves (NIW)

- Near-inertial frequencies are associated primarily with atmospheric forcing.
- The ocean selects near-inertial frequencies for amplification because the aspect ratio of the forcing is exceedingly small.
  - Storms tends to have horizontal scales of  $O(100-1000 \text{ km})$
  - The ocean mixed layer is typically no deeper than 100 m
  - Storms move rapidly so that the ocean does not have time to adjust
  -

## 1.3. Generation of Near-Inertial waves (NIW)

- We can integrate the linearized hydrostatic Boussinesq equations on the  $f$ -plane over the mixed layer (Horizontal gradients and vertical velocities can be neglected at the lowest order):

$$\frac{\partial u}{\partial t} = -\frac{1}{\rho_0} \frac{\partial p}{\partial x} + fv + \frac{1}{\rho_0} \frac{\partial}{\partial z} \tau_x,$$

$$\frac{\partial v}{\partial t} = -\frac{1}{\rho_0} \frac{\partial p}{\partial y} - fu + \frac{1}{\rho_0} \frac{\partial}{\partial z} \tau_y,$$

$$\frac{\partial p}{\partial z} = -g\rho',$$

$$\frac{\partial}{\partial z} \mathbf{u} = 0,$$

$$\frac{\partial \rho'}{\partial t} = \frac{\rho_0}{g} N^2 w,$$



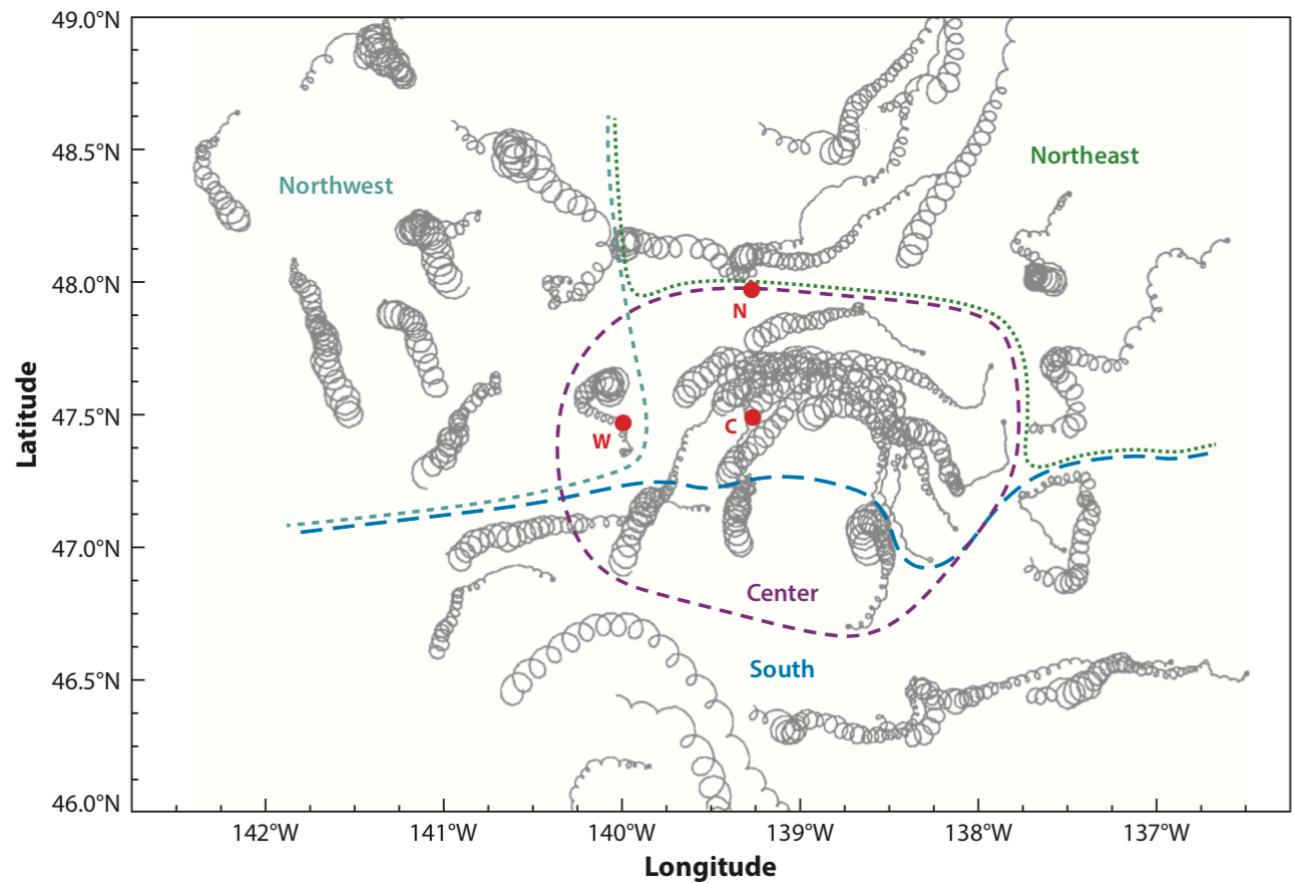
$$\frac{\partial U}{\partial t} = +fV + \frac{1}{\rho_0} \frac{\partial}{\partial z} T_x,$$

$$\frac{\partial V}{\partial t} = -fU + \frac{1}{\rho_0} \frac{\partial}{\partial z} T_y,$$

- Solutions are the sum of a steady Ekman transport to the right of the wind and anticyclonic circular motions at the local inertial frequency,  $\mathbf{f}$

# 1.3. Generation of Near-Inertial waves (NIW)

- NIW are recognizable by their characteristic circularly polarized velocities



Twenty-five days of surface drifter trajectories after a storm in the eastern north Pacific. The drifters trajectories represent a combination of decaying inertial motions (circular oscillations) and weak geostrophic flow (the time-averaged drift). [D'Asaro et al, 1995]

## 1.3. Generation of Near-Inertial waves (NIW)

- Dispersion relation for internal waves:

$$\omega^2 = \frac{N^2 k^2 + f^2 m^2}{k^2 + m^2}$$

.

## 1.3. Generation of Near-Inertial waves (NIW)

- Dispersion relation for internal waves:

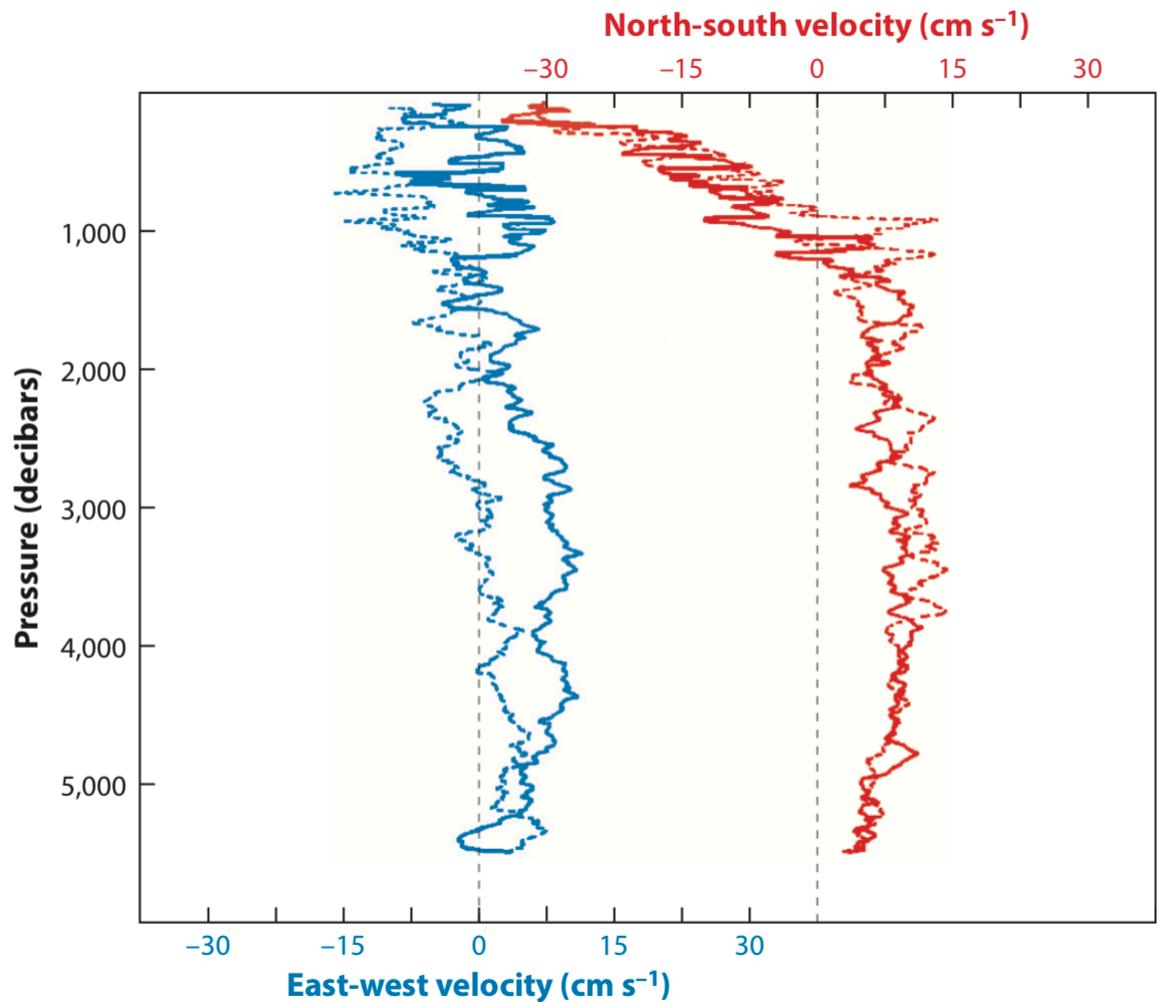
$$\omega^2 = \frac{N^2 k^2 + f^2 m^2}{X^2 + m^2}$$

- For hydrostatic NIW:

$$\omega^2 = f^2 + \frac{N^2 k^2}{m^2} \quad k \ll m$$

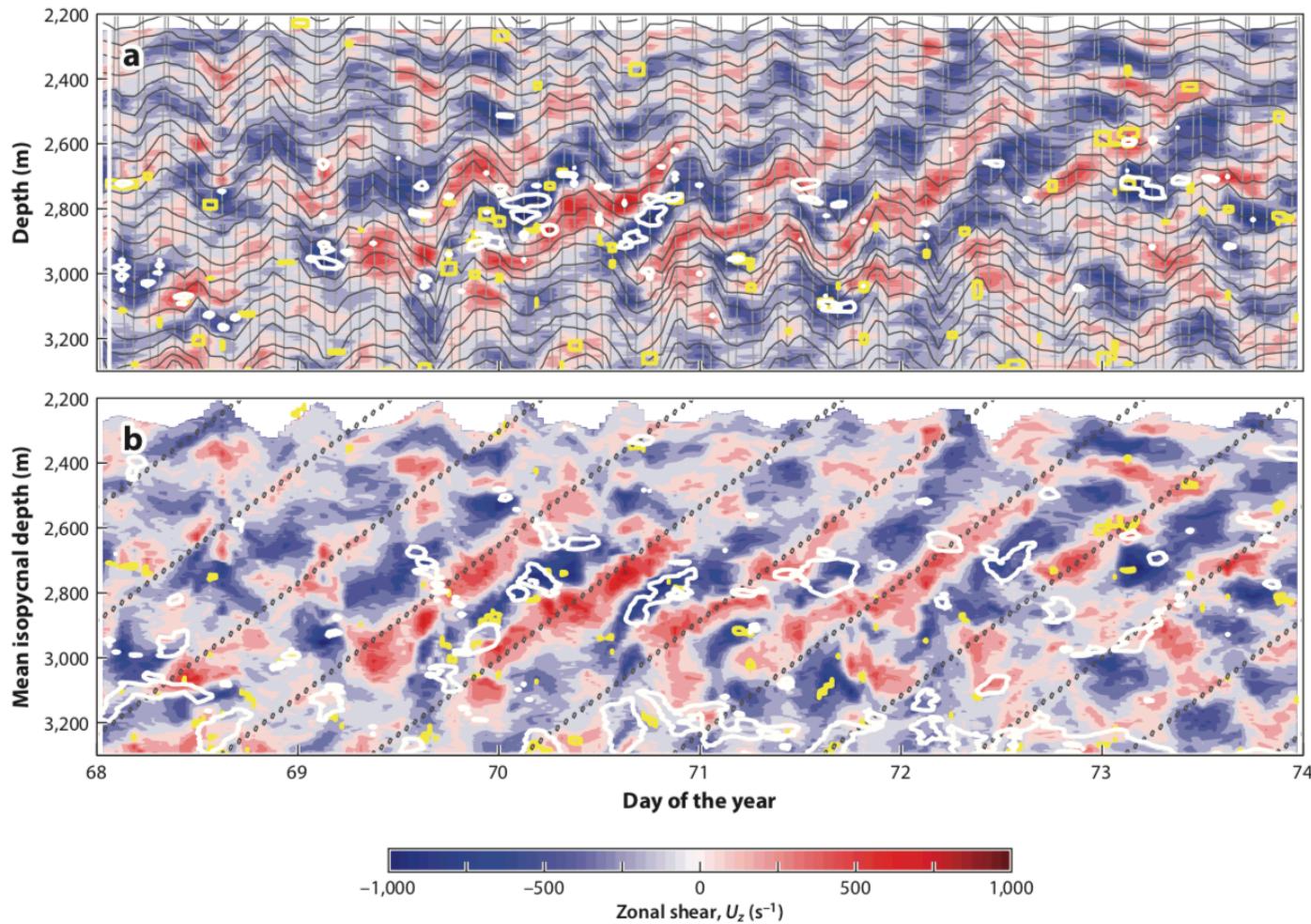
## 1.3. Generation of Near-Inertial waves (NIW)

- NIW are also recognizable by their a strong vertical shears



Two profiles (dashed and solid lines) of east-west (blue) and north-south (red ) velocity taken at an interval of half an inertial period, showing high-wavenumber near-inertial motions. [Leaman & Sanford, 1975]

# 1.3. Generation of Near-Inertial waves (NIW)



For vertical wavelength of 350 m and  $\omega = 1.03f$ , the horizontal wavelength would be 80 km and the vertical group velocity  $10^{-3} \text{ m s}^{-1}$

Zonal shear plotted versus (a) depth and (b) density (expressed as the mean depth of each isopycnal), showing the upward phase propagation of near-inertial shear layers. Shear layers are aligned with isopycnals (dark gray lines in panel a), which are heaved vertically by internal waves of other frequencies (primarily tidal; light gray lines); this heaving is removed in the isopycnal-following frame (panel b). The dissipation rate from overturns is contoured in yellow (contour value  $1.5 \times 10^{-8} \text{ W kg}^{-1}$ ); the Richardson number (white contours) drops below 0.8 once per wave period. In panel b, the dotted lines represent the wave phase for a signal with  $\omega = 1.03 f$  and  $\lambda = 350$  m.

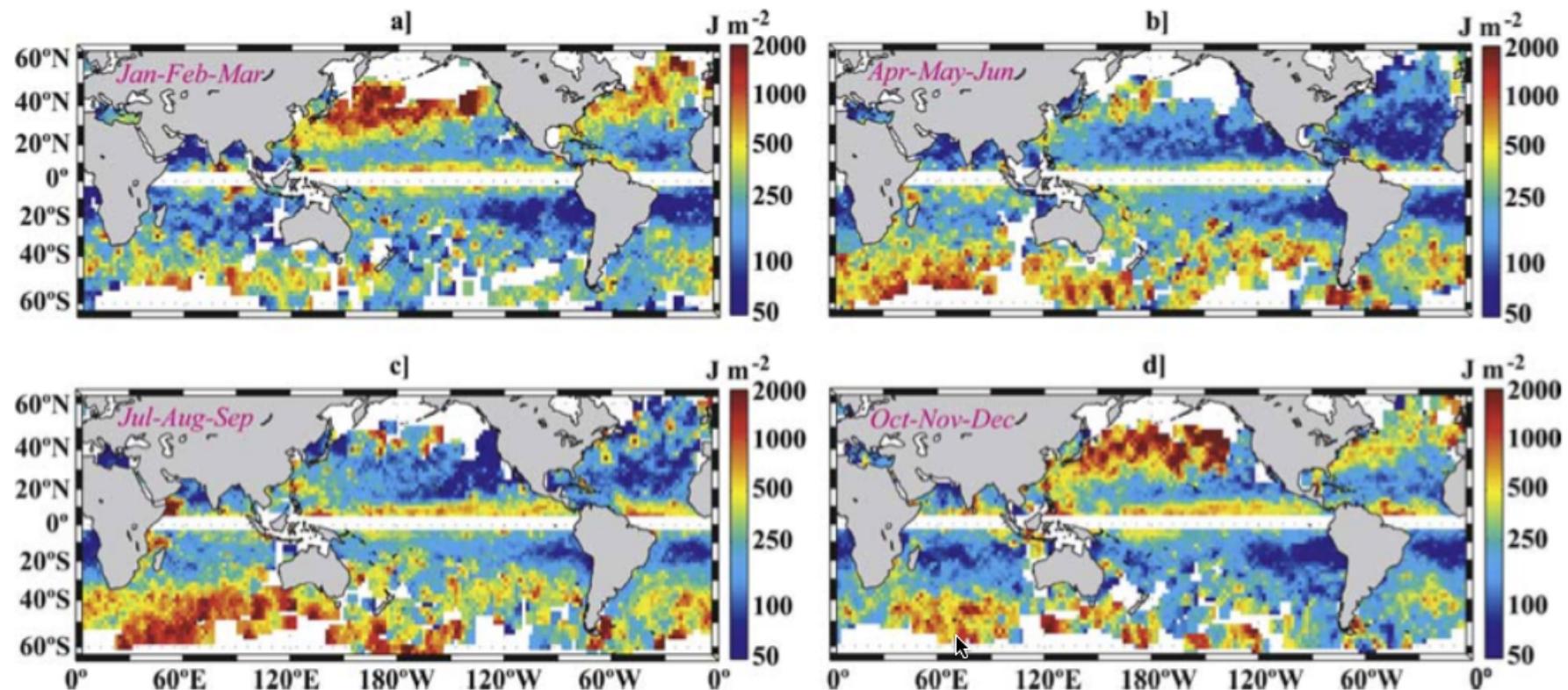
# 1.3. Generation of Near-Inertial waves (NIW)

**The general characteristics of NIWs are as follows:**

- Clockwise and counterclockwise polarization in the Northern and Southern Hemispheres, respectively
- Frequency of  $1\text{--}1.2f$
- Vertical wavelength of 100–400 m
- Lateral scale of 10–500 km

# 1.3. Generation of Near-Inertial waves (NIW)

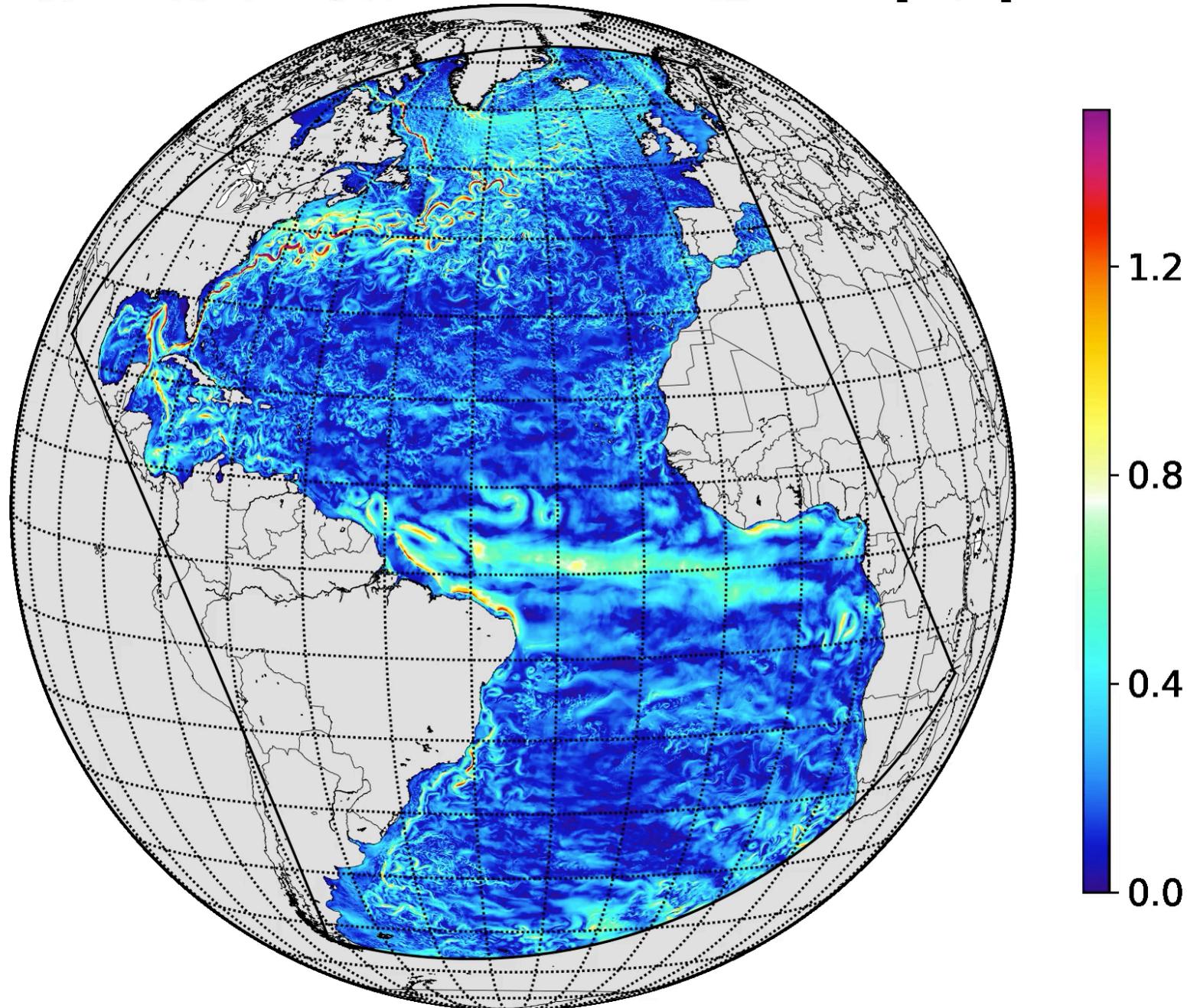
Seasonal band-passed near-inertial energy density in the boundary layer is elevated under storm tracks.  
Results computed from surface drifter trajectories. [Chaigneau et al., 2008]



1999 - Feb 25 - 16:00

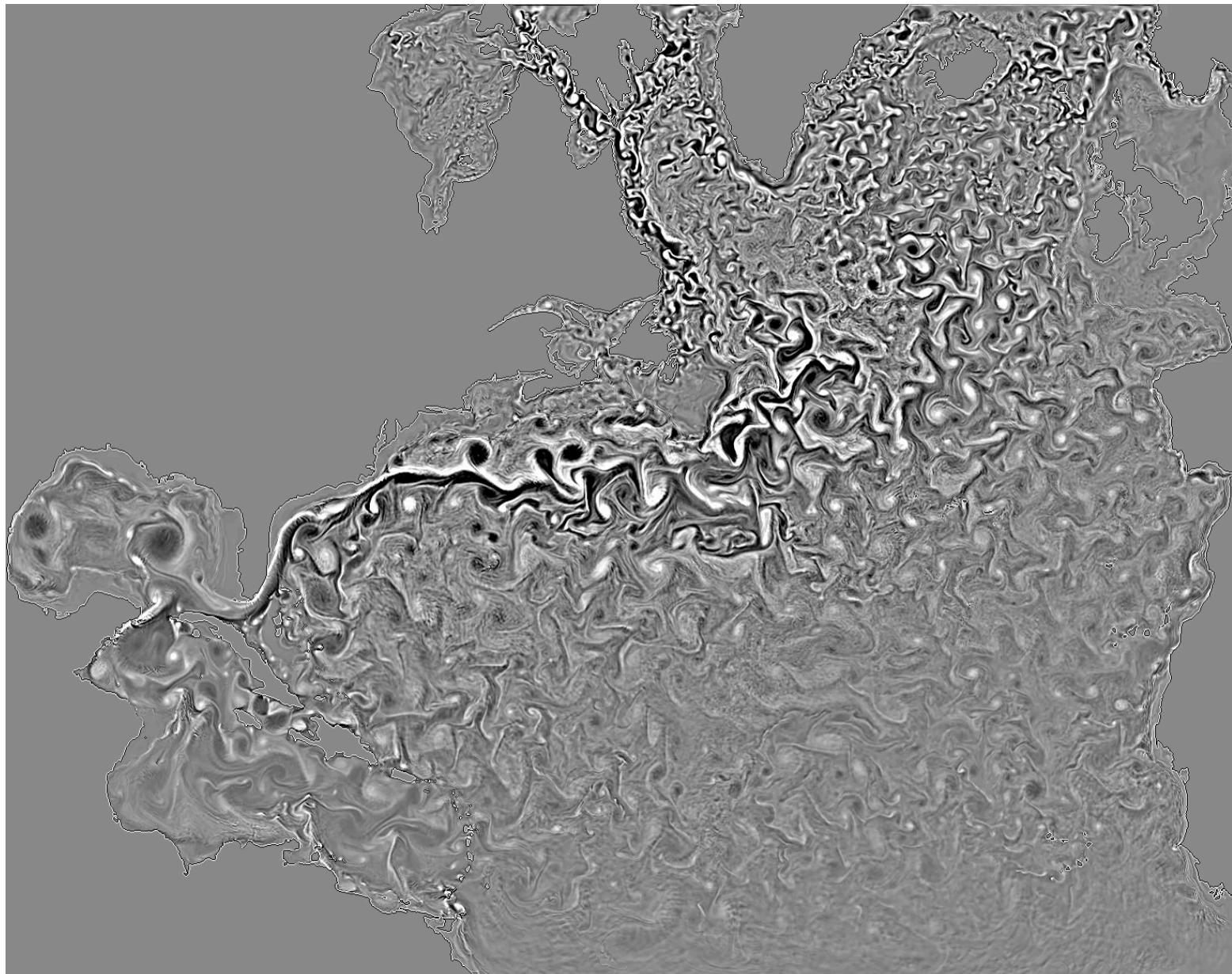
currents [m/s]

- High-frequency surface forcings (using hourly CFSR data and bulk flux formulae)

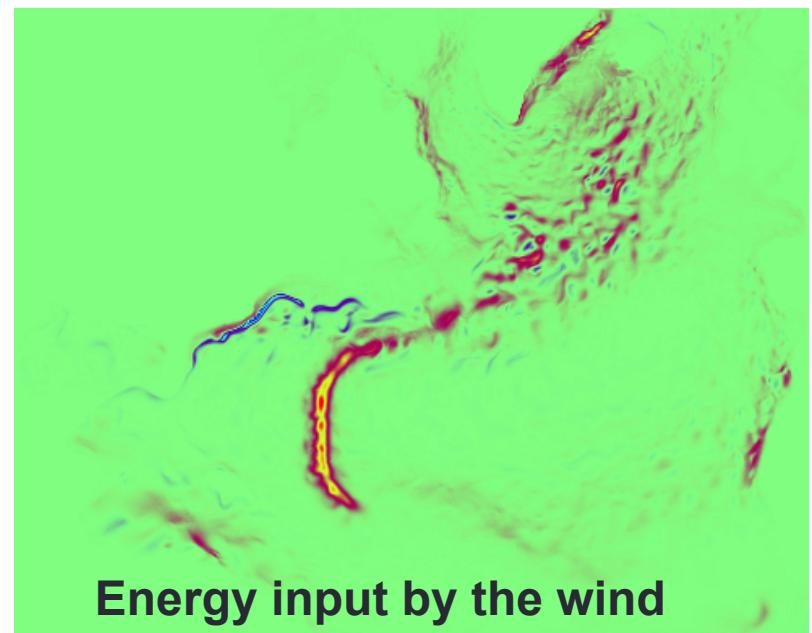
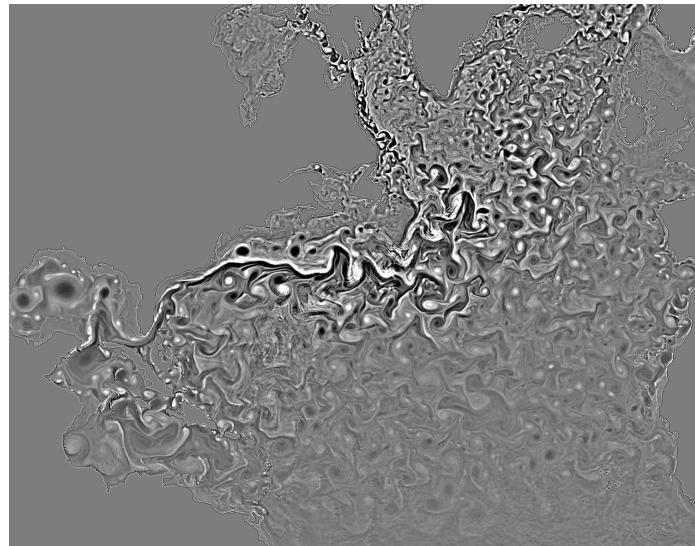
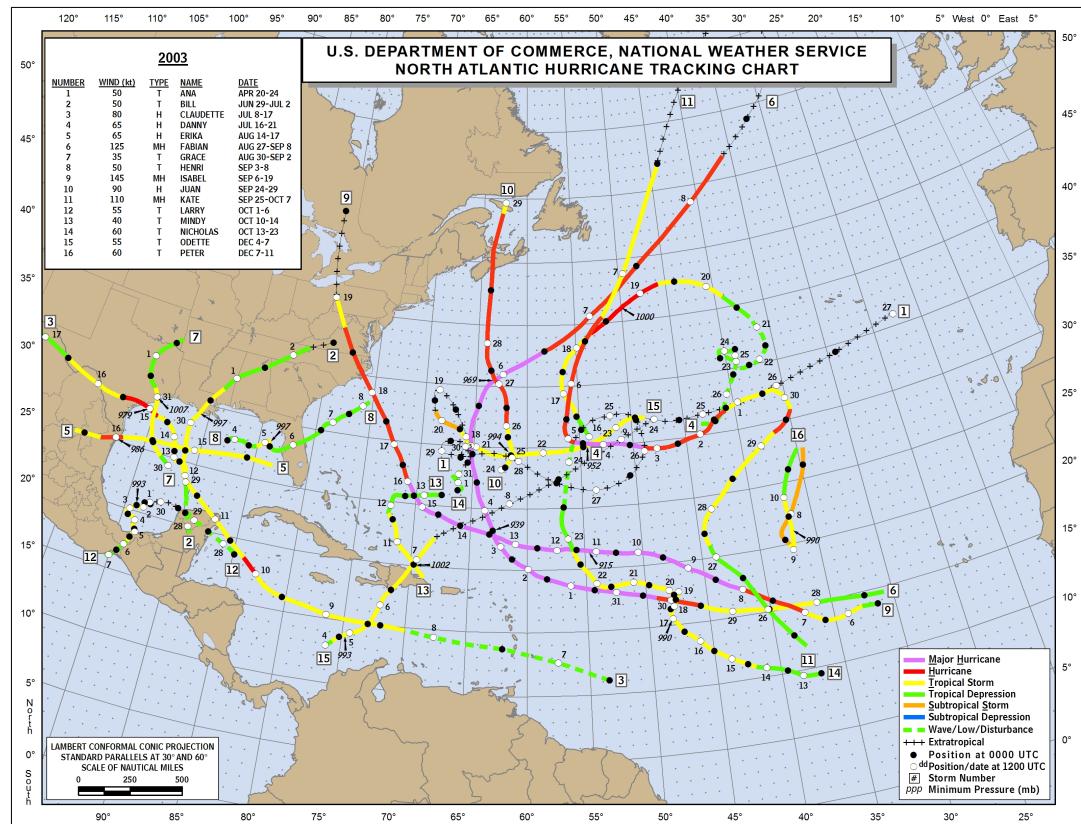


## 1.3. Generation of Near-Inertial waves (NIW)

- High-frequency surface forcings (using hourly CFSR data and bulk flux formulae)

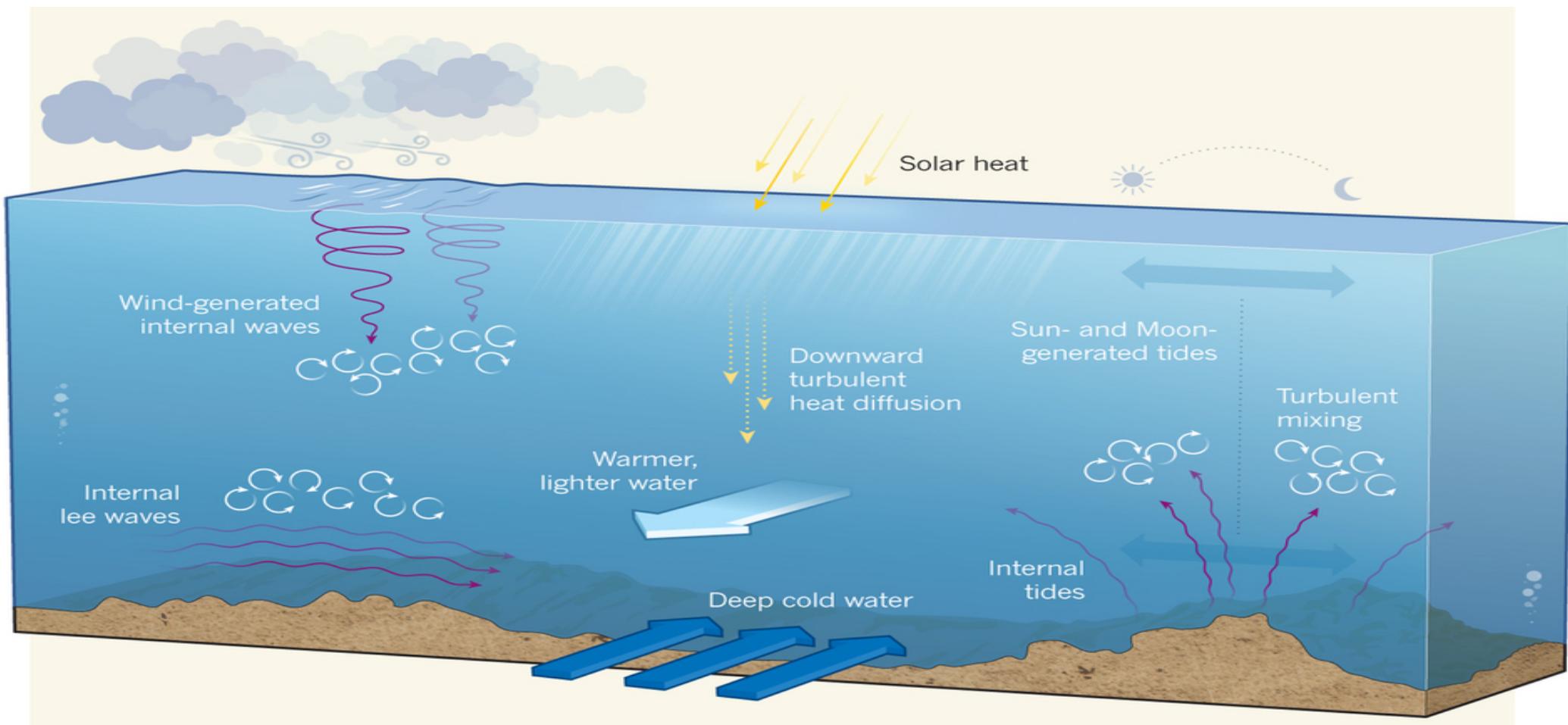


# 1.3. Generation of Near-Inertial waves (NIW)



# 1. Internal waves generation

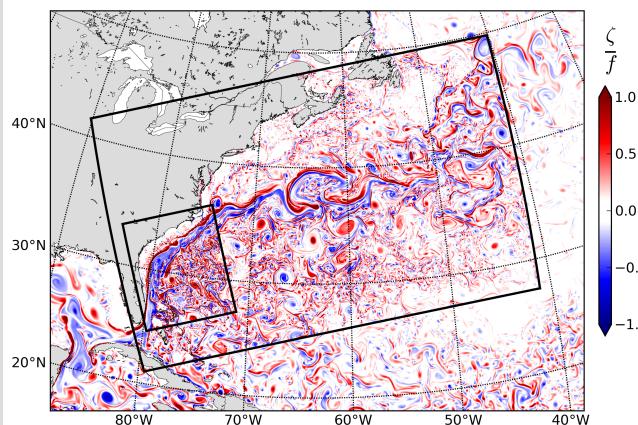
- Mechanisms:



From Mackinnon 2013

# Realistic modelling including tides and NIW

- **Gulf Stream / Sargasso Sea**

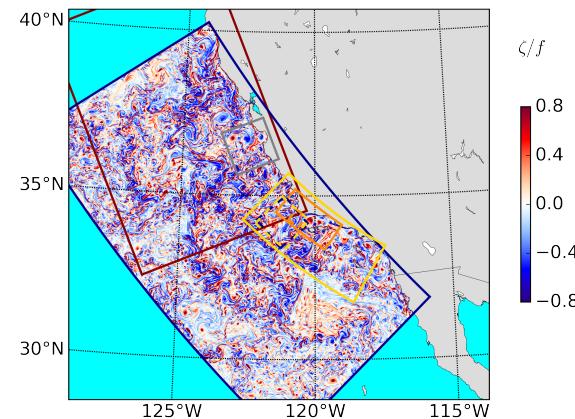


750 m res. forced with:

- Monthly winds
- Hourly winds
- Hourly winds + Tides

(+200 m with HF forcings )

- **California Current**

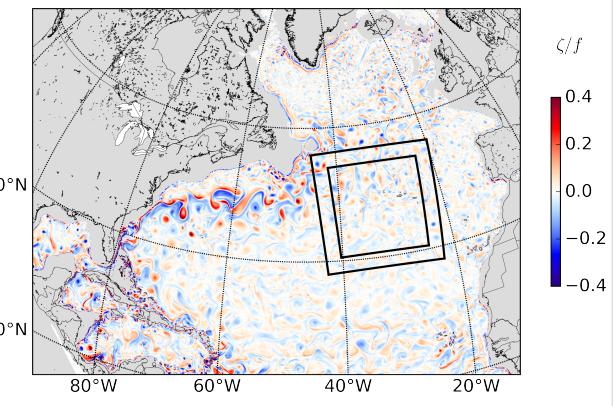


1 km res. forced with:

- Daily winds
- Hourly winds
- Hourly winds + Tides

(+300 m, 100 m with HF forcings and tides)

- **Mid-Atlantic ridge**

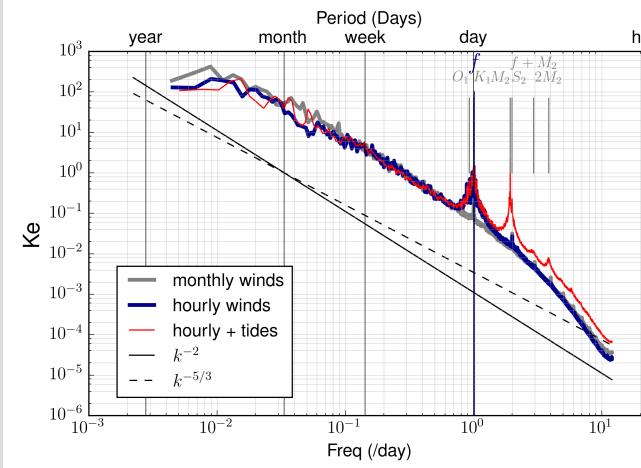
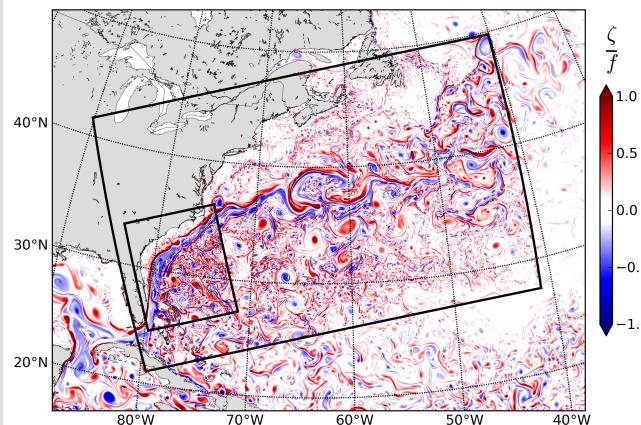


750 m res. forced with:

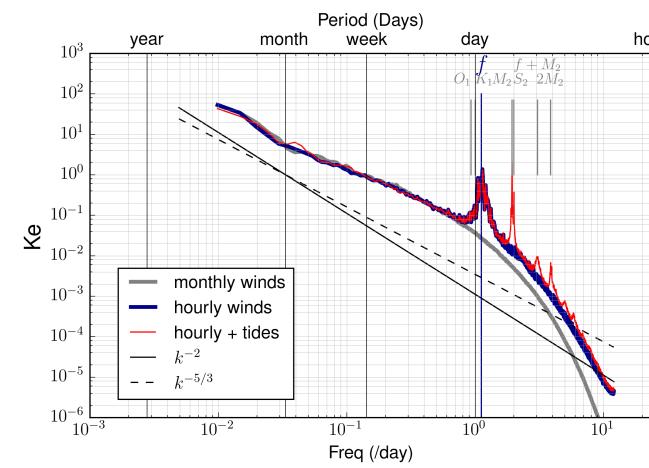
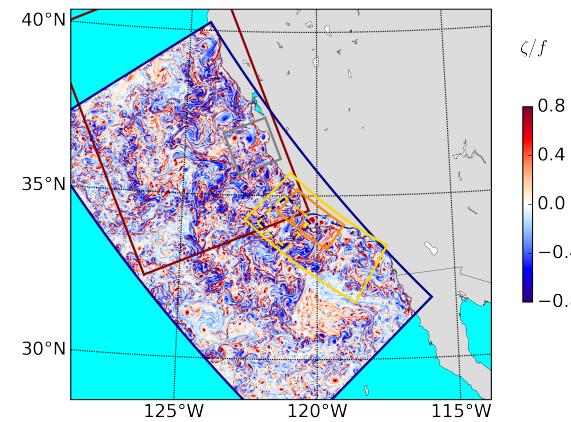
- daily winds
- daily winds + Tides
- Only tides

# Realistic modelling including tides and NIW

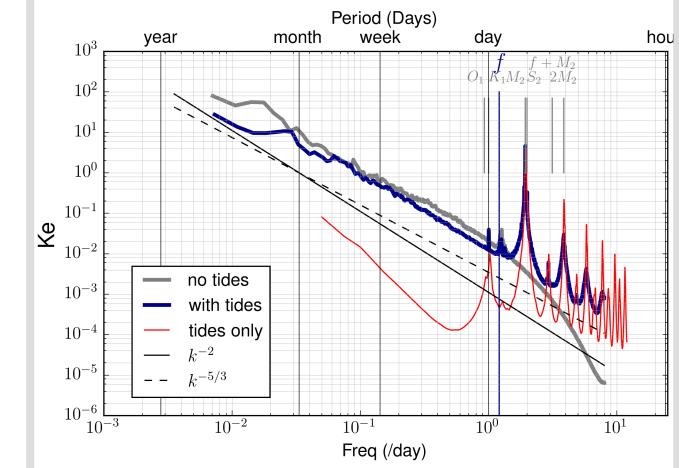
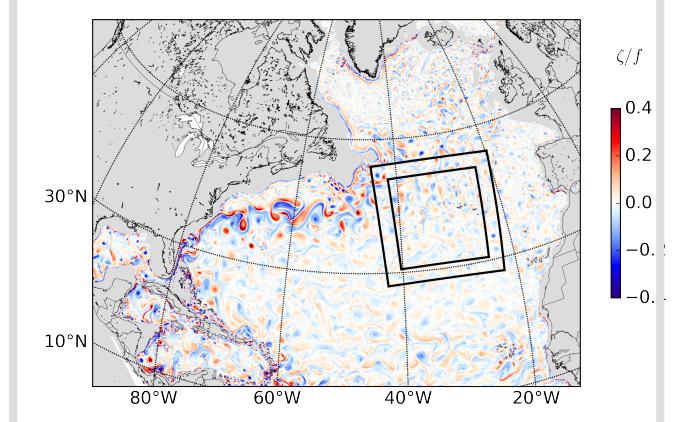
- **Gulf Stream / Sargasso Sea**



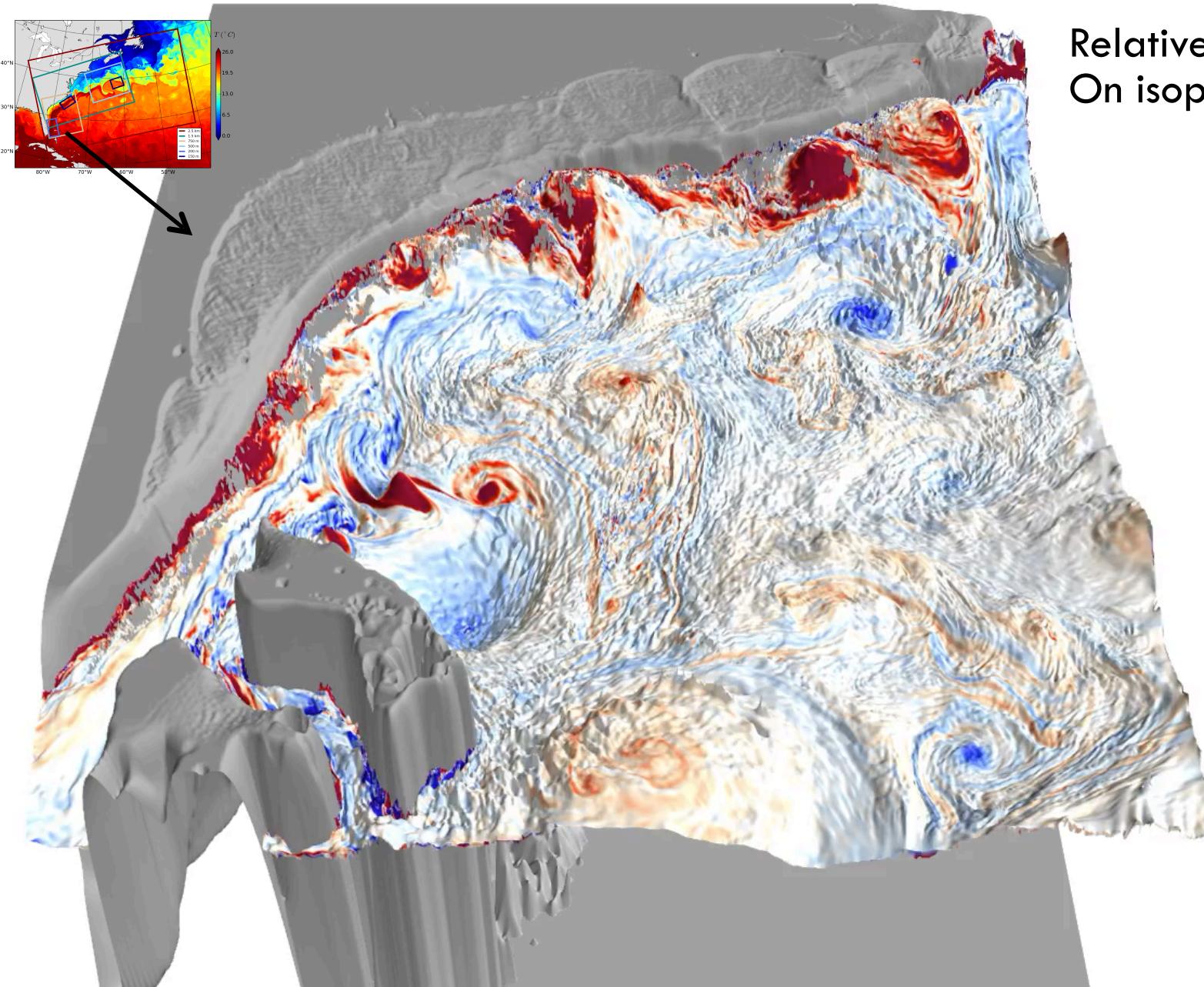
- **California Current**



- **Mid-Atlantic ridge**



# Method: Realistic modelling including tides

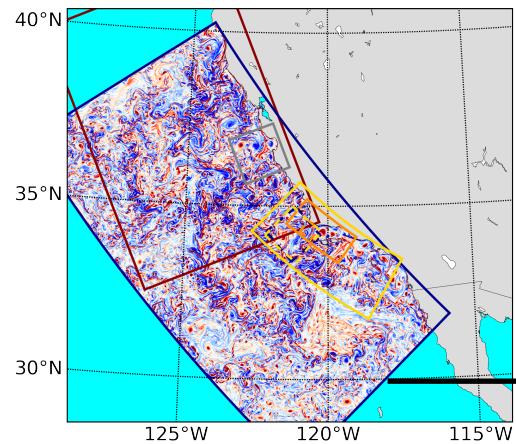


Relative vorticity  $(\pm f)$   
On isopycnal  $\sigma = 27 \text{ kg m}^{-3}$

$$\Delta x = 750 \text{ m}$$

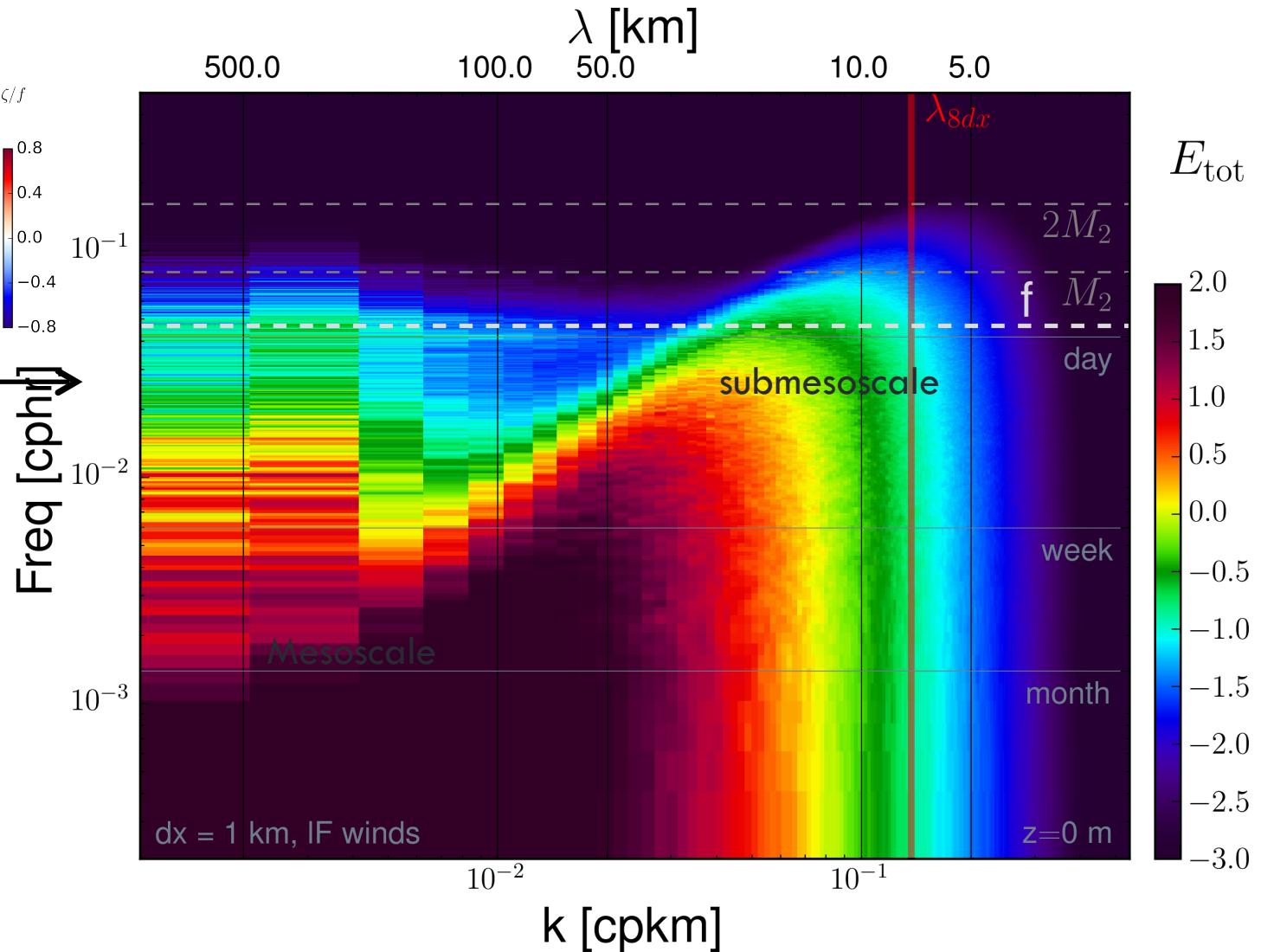
HF wind  
+  
Tides

# Realistic modelling including tides and NIW



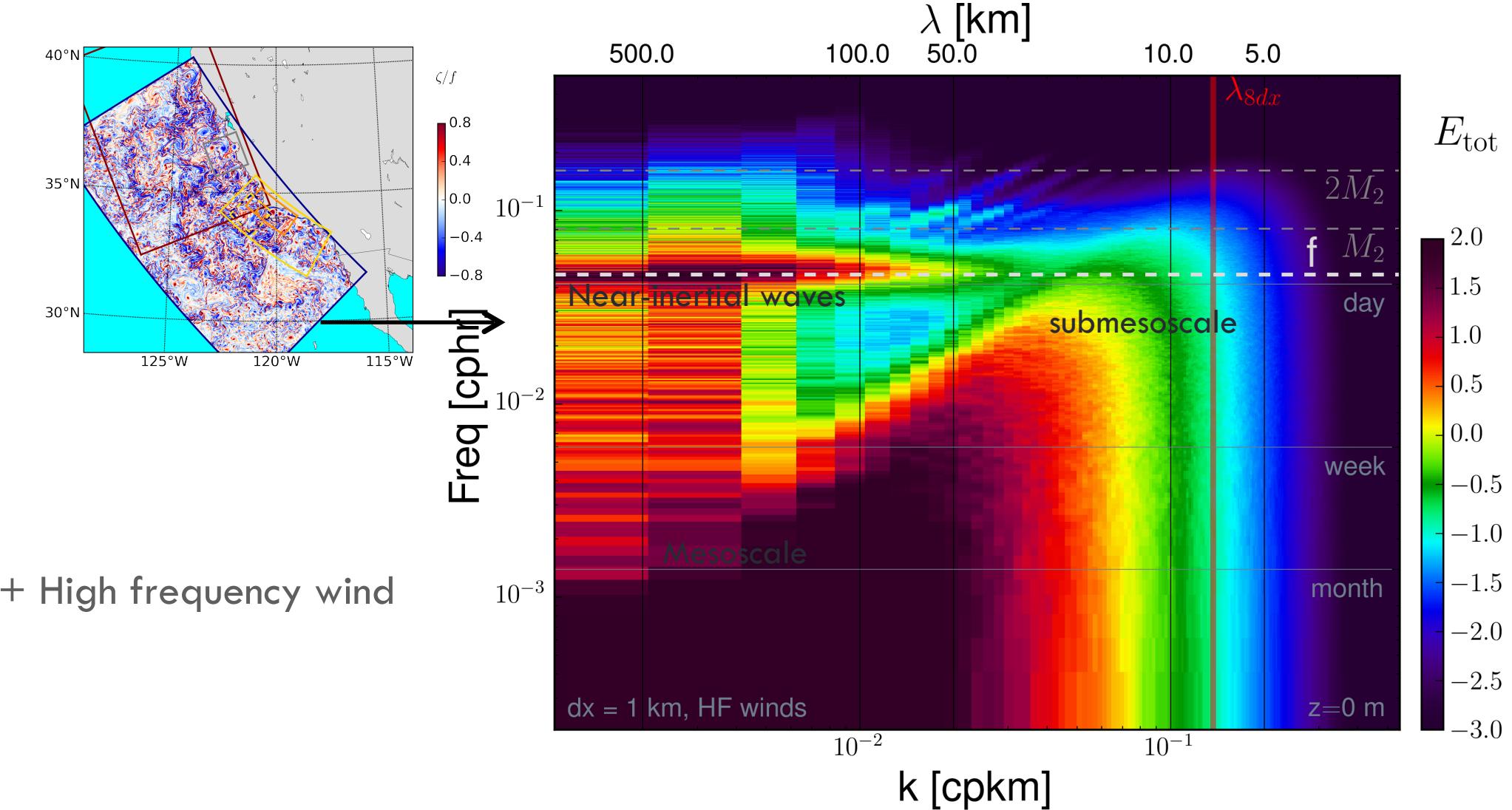
Daily winds No tides

= deficient in internal waves



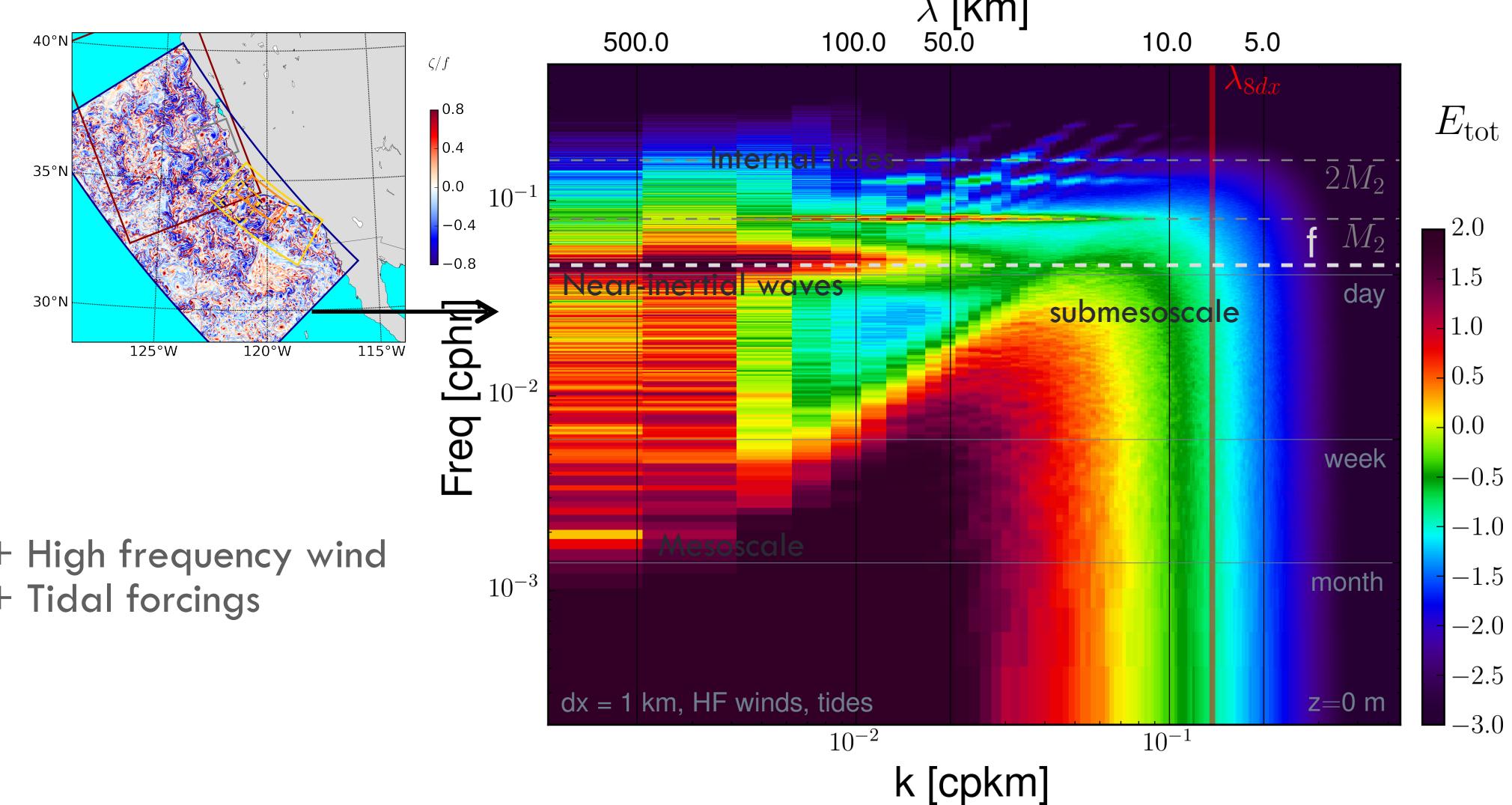
Azimuthally-averaged 2D frequency-wavenumber kinetic energy spectra in the California Current.

# Realistic modelling including tides and NIW



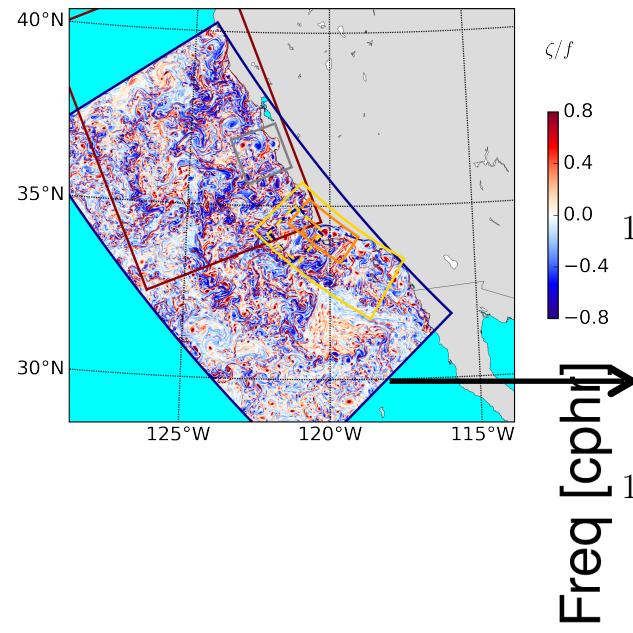
Azimuthally-averaged 2D frequency-wavenumber kinetic energy spectra  
in the California Current.

# Realistic modelling including tides and NIW

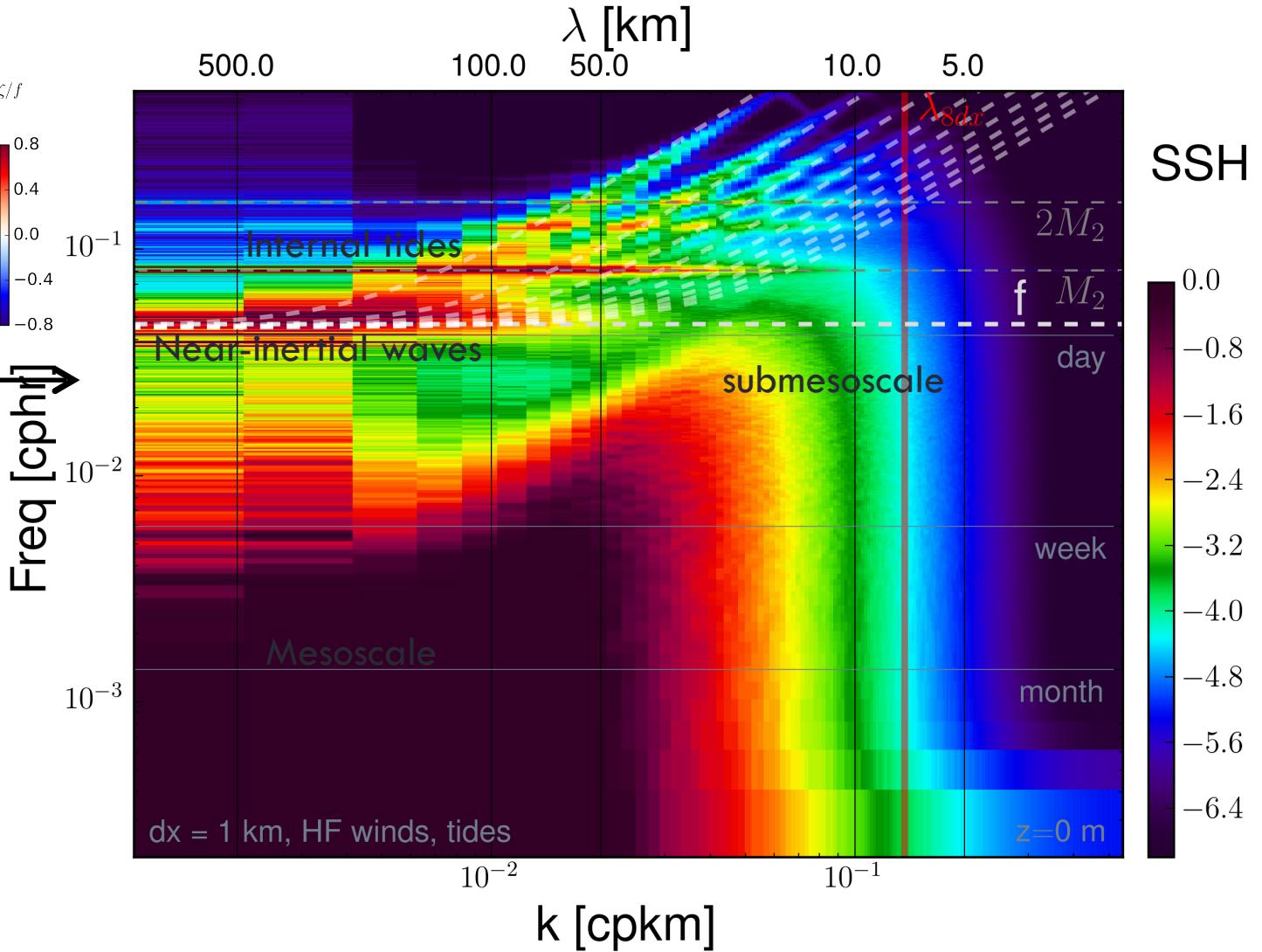


Azimuthally-averaged 2D frequency-wavenumber kinetic energy spectra in the California Current.

# Realistic modelling including tides and NIW



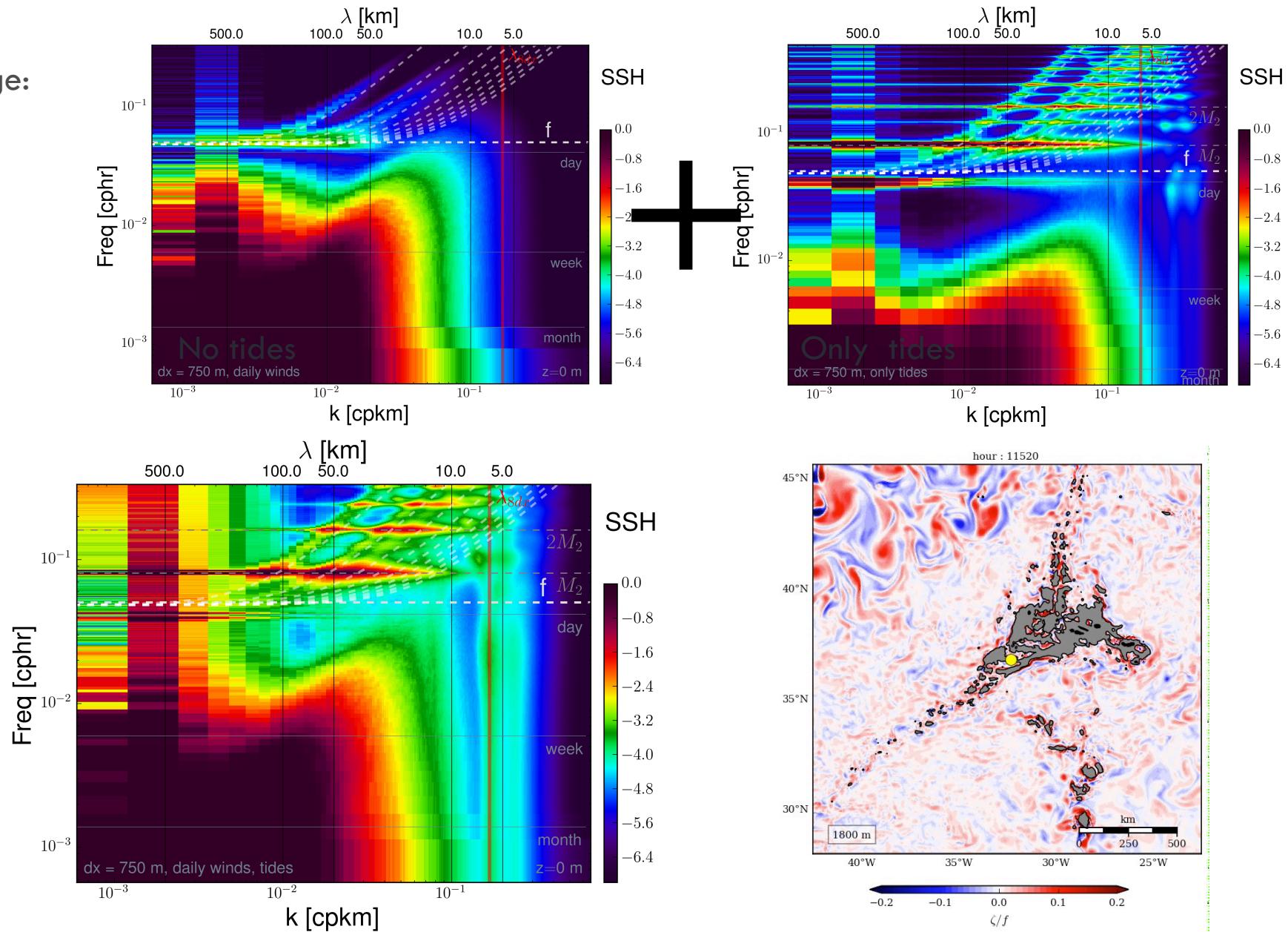
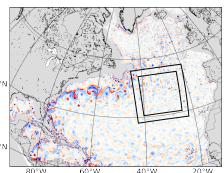
Signature of high-frequency internal waves are amplified on **Sea Surface Height**.



Azimuthally-averaged 2D frequency-wavenumber spectra for SSH in California Current

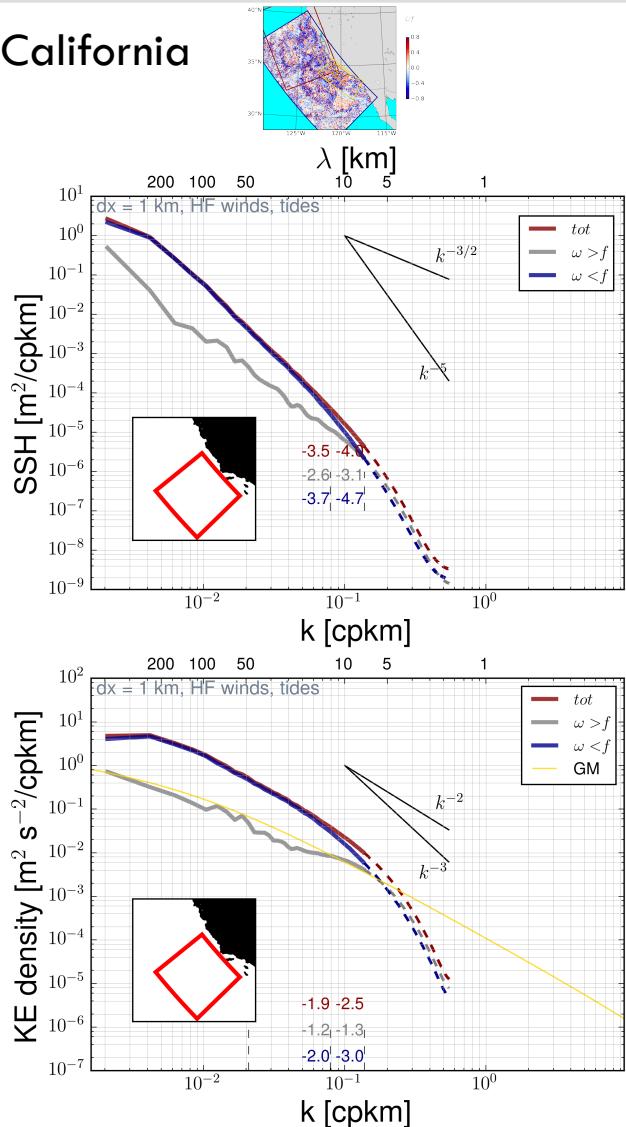
# Realistic modelling including tides and NIW

Mid-Atlantic ridge:

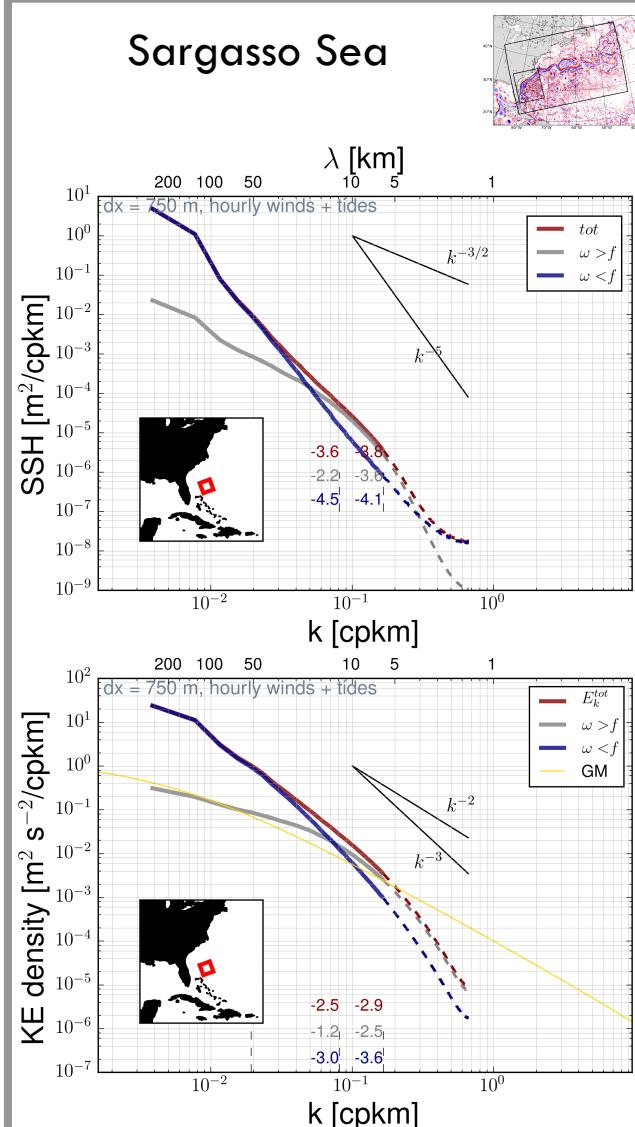


# Where is the flow balanced?

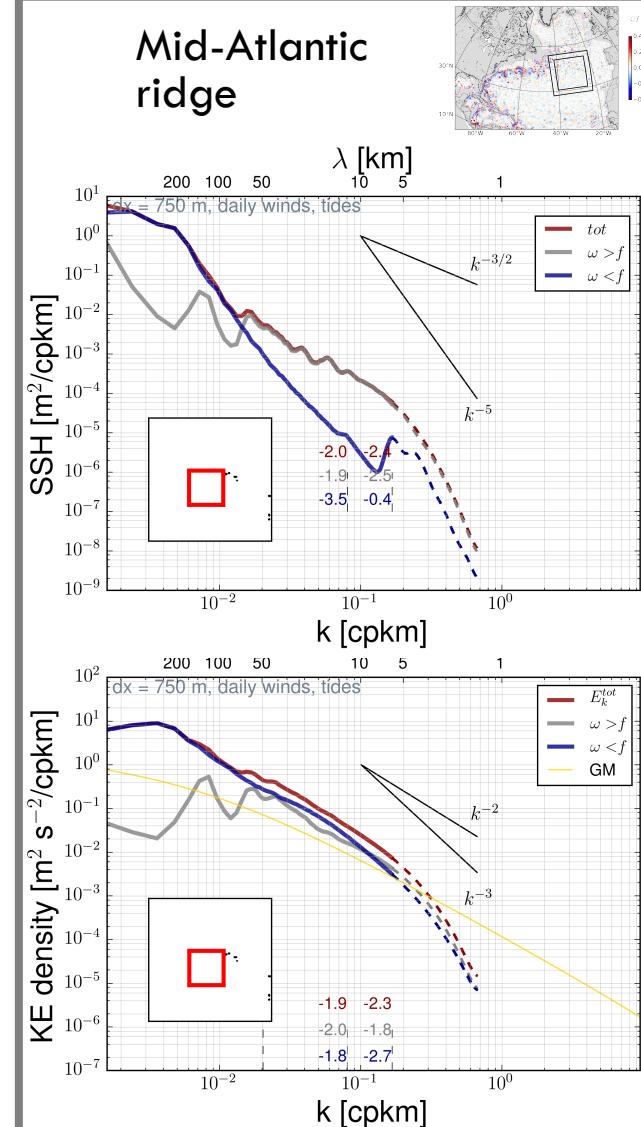
**California**



**Sargasso Sea**

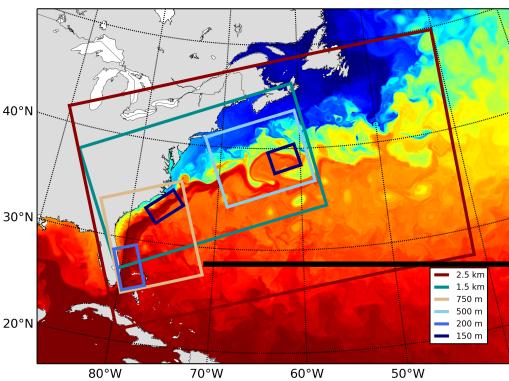


**Mid-Atlantic ridge**

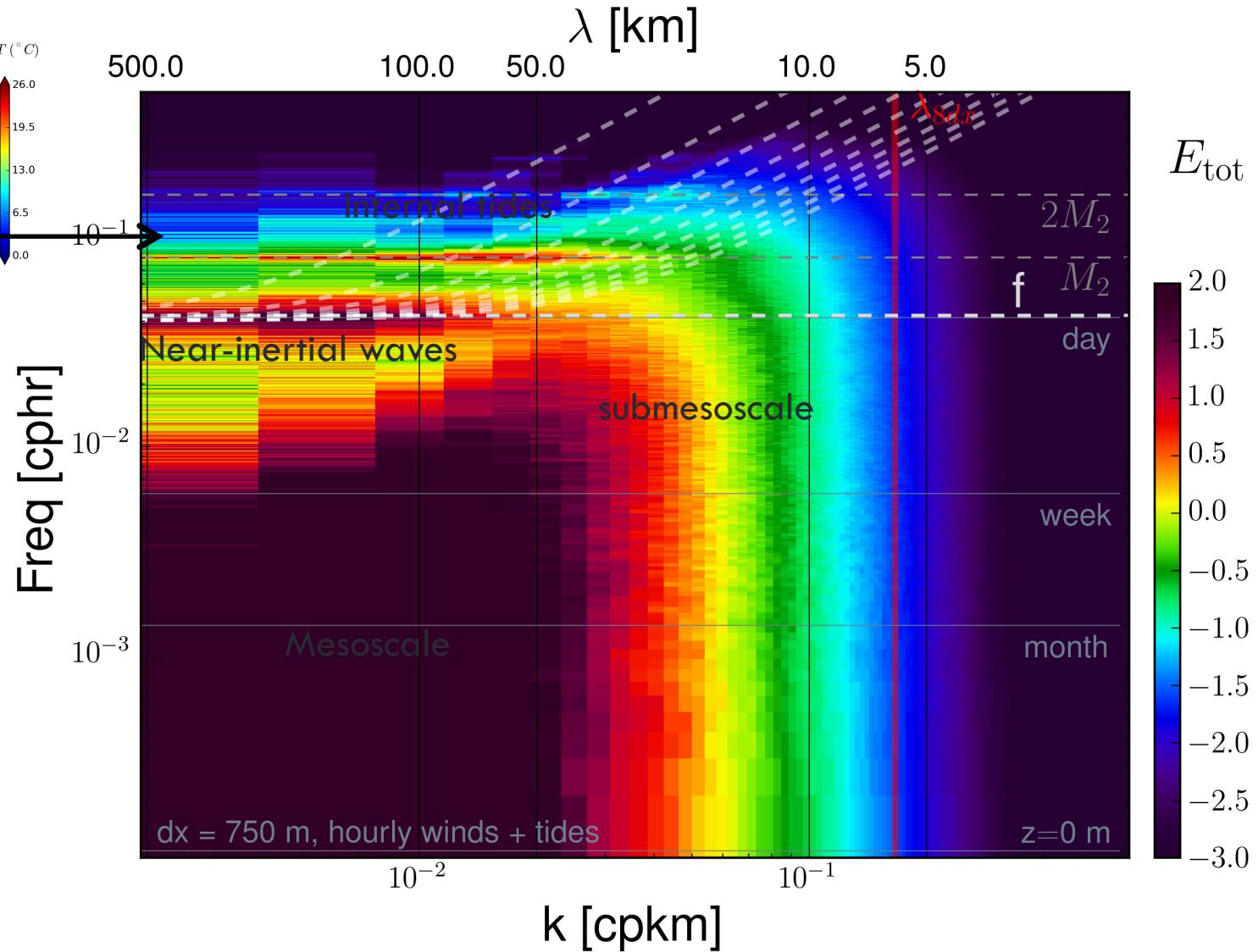


Signature of internal waves are amplified on **Sea Surface Height** – Leading to more dramatic breaks in slopes than in Kinetic energy.

# Where is the flow balanced?



- + High frequency wind
- + Tidal forcings

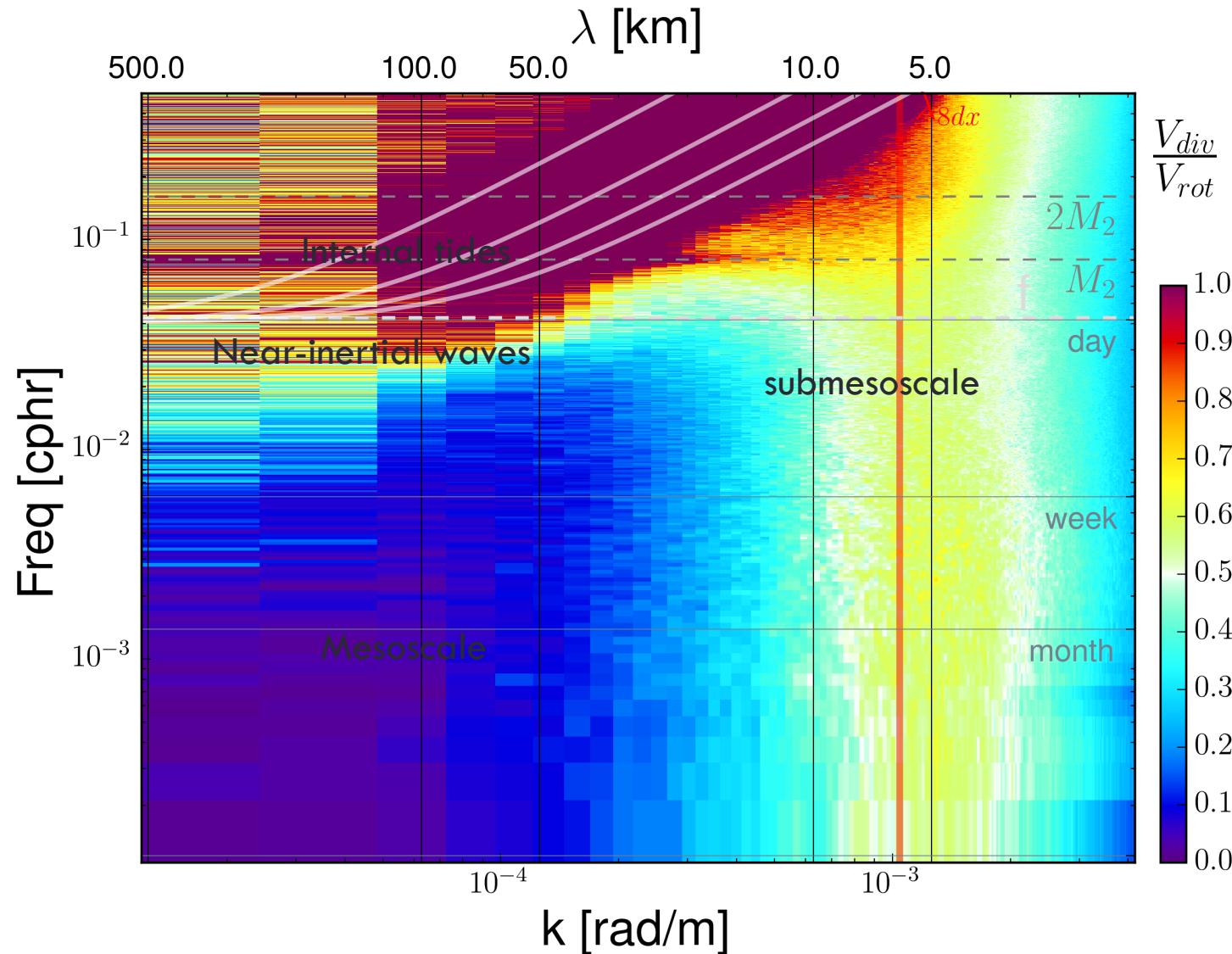


*Azimuthally-averaged 2D frequency-wavenumber kinetic energy spectra in the Sargasso Sea.*

# Where is the flow balanced?

Ratio of divergent /  
rotational part of the  
kinetic energy  
  
using Helmholtz  
decomposition of a 3d  
incompressible flow

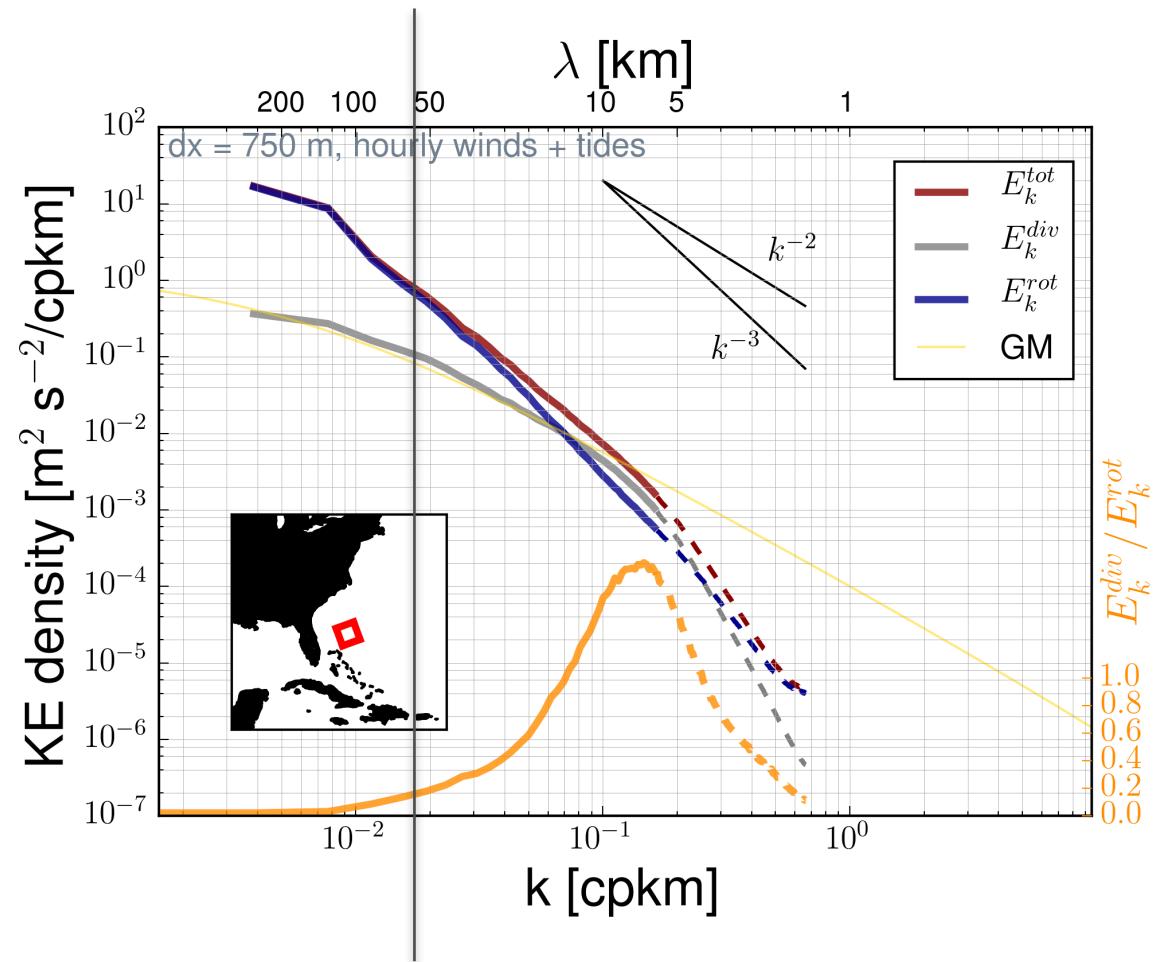
$$\mathbf{u}_h = \mathbf{u}_r + \mathbf{u}_d, \\ \nabla_h \cdot \mathbf{u}_r = 0 \\ \hat{\mathbf{z}} \cdot \nabla_h \times \mathbf{u}_d = 0$$



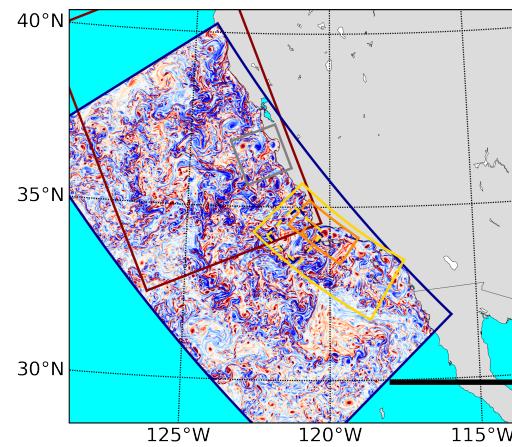
# Where is the flow balanced?

Transition between balanced and unbalanced at scales between 10 – 100 km depending on the dynamical regime (geographical region + season).

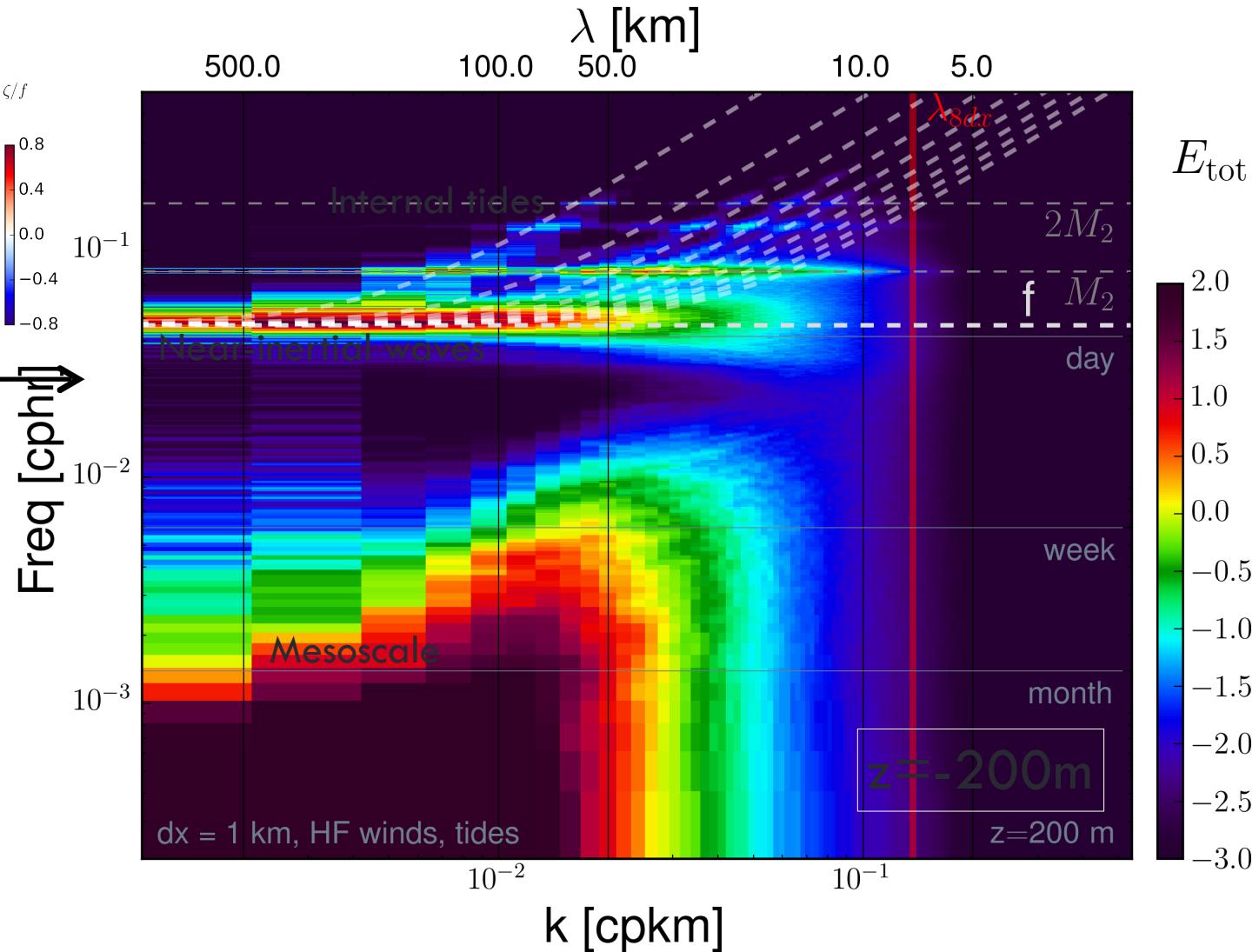
Qualitatively similar results for Northwestern Pacific [Rocha et al., 2016a; Qiu et al., 2017] or the Drake Passage [Rocha et al., 2016]. See global census in [Qiu et al., 2018].



# Can we separate waves from submesoscales?

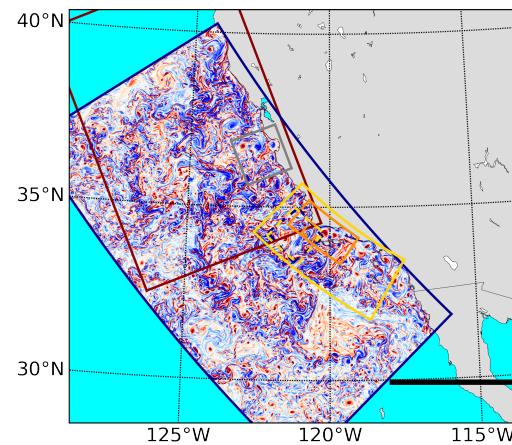


Below the mixed-layer,  
separation between  
internal waves and  
balanced dynamics is  
easy.

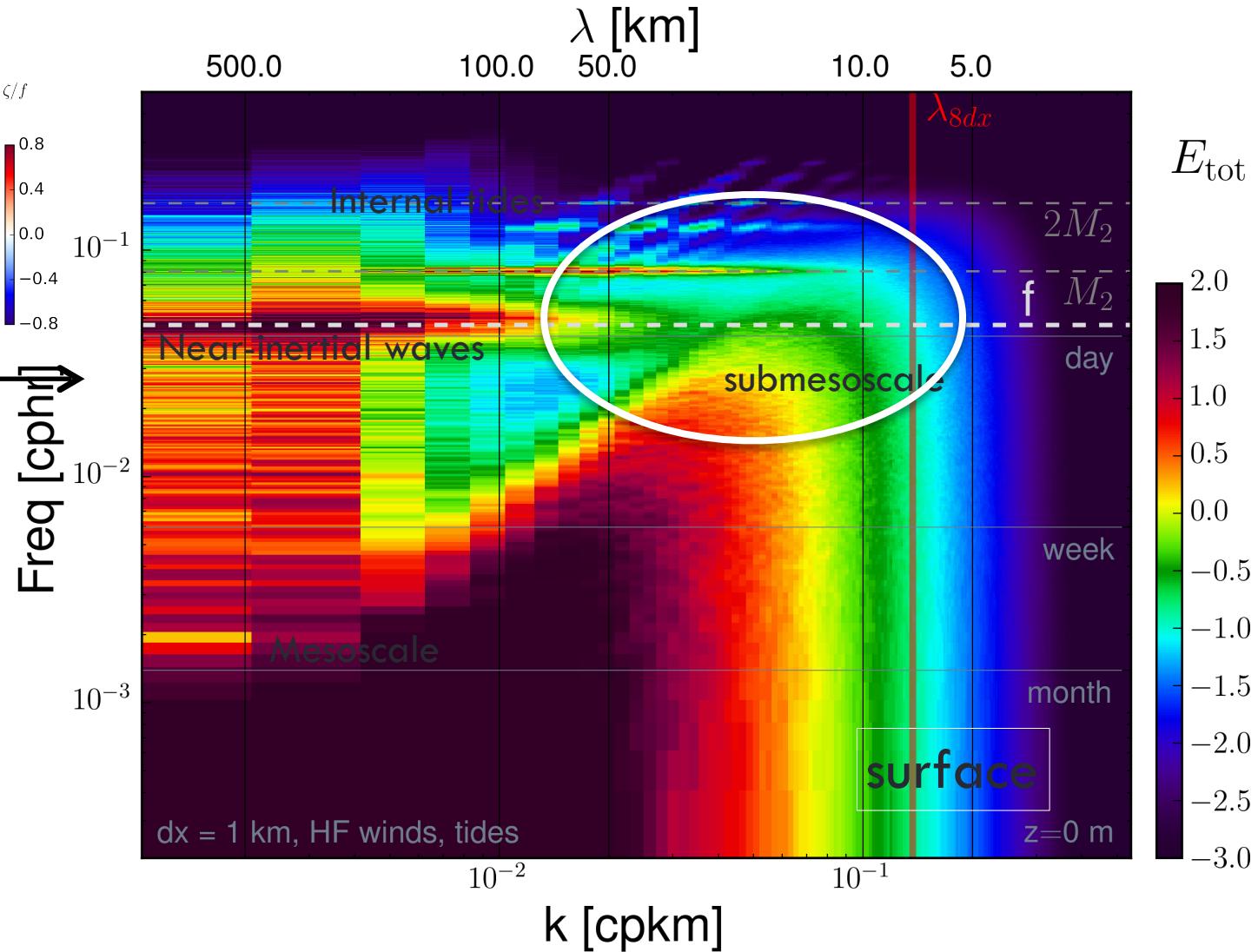


Azimuthally-averaged 2D frequency-wavenumber kinetic energy spectra  
in the California Current.

# Can we separate waves from submesoscales?



In the mixed-layer it is much more complicated



Azimuthally-averaged 2D frequency-wavenumber kinetic energy spectra in the California Current.

## Activity: Numerical simulation of Internal waves

- Download and unzip Fluid2d: [jqula.fr/Ondes/fluid2d.tar.gz](http://jqula.fr/Ondes/fluid2d.tar.gz)
- Install (see also <http://mespages.univ-brest.fr/~roullet/fluid2d/docs/howto.html#download-and-install>)

To have the code run you need to have the following python packages installed

numpy

matplotlib

netCDF4

mpi4py >= 2.0.0

```
tar -xf fluid2d.tar.gz
```

```
cd fluid2d
```

```
make
```

```
source activate.sh
```

## Activity: Numerical simulation of Internal waves

- Download and unzip Fluid2d: [jqula.fr/Ondes/fluid2d.tar.gz](http://jqula.fr/Ondes/fluid2d.tar.gz)
- Install (see also <http://mespages.univ-brest.fr/~roullet/fluid2d/docs/howto.html#download-and-install>)

To have the code run you need to have the following python packages installed

numpy

matplotlib

netCDF4

mpi4py >= 2.0.0

```
tar -xf fluid2d.tar.gz
```

```
cd fluid2d
```

```
make
```

```
source activate.sh
```

## Activity: Numerical simulation of Internal waves

Run the Interfacial wave case:

- cd nhom/Interfacial
- python interfacialwave.py

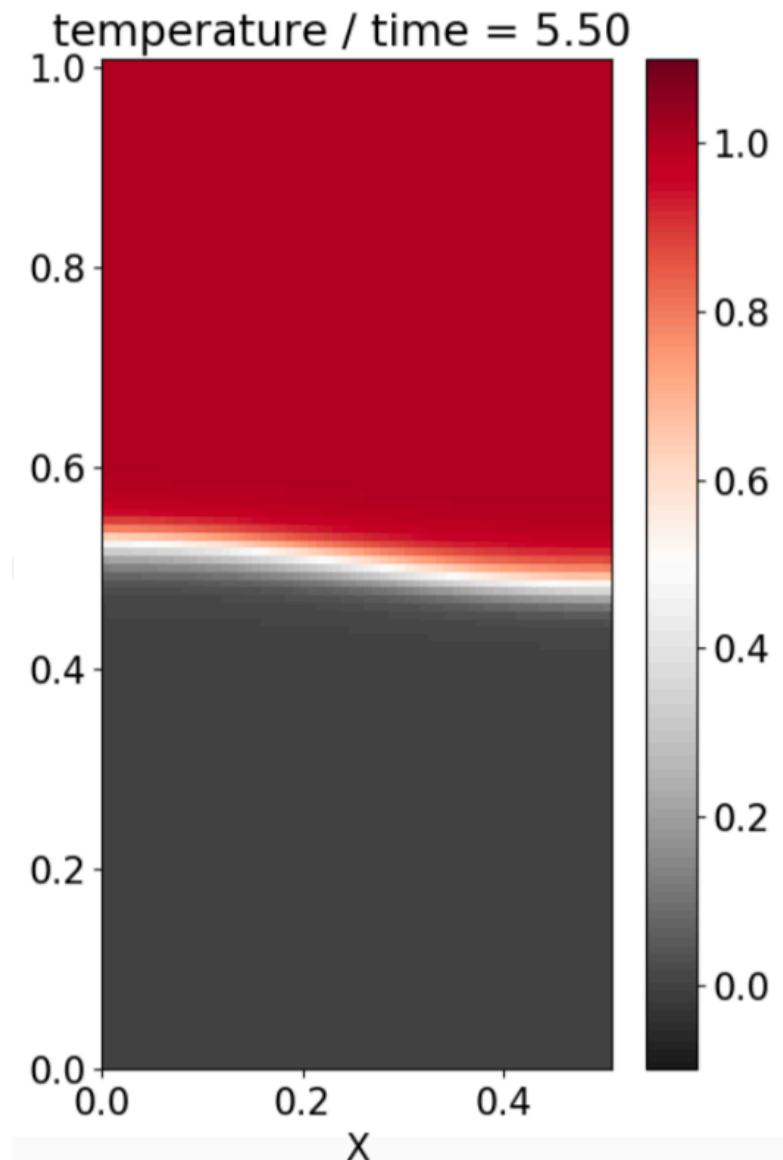
Either

- 1) a sine wave in a tank (stationary wave) or
- 2) a localized perturbation on the left  
(propagating wave)

In 2) look at how the Stokes drift deforms the tracer field

In 1) play with hydroepsilon (not smaller than 0.2),  
look at the structure of the velocity field

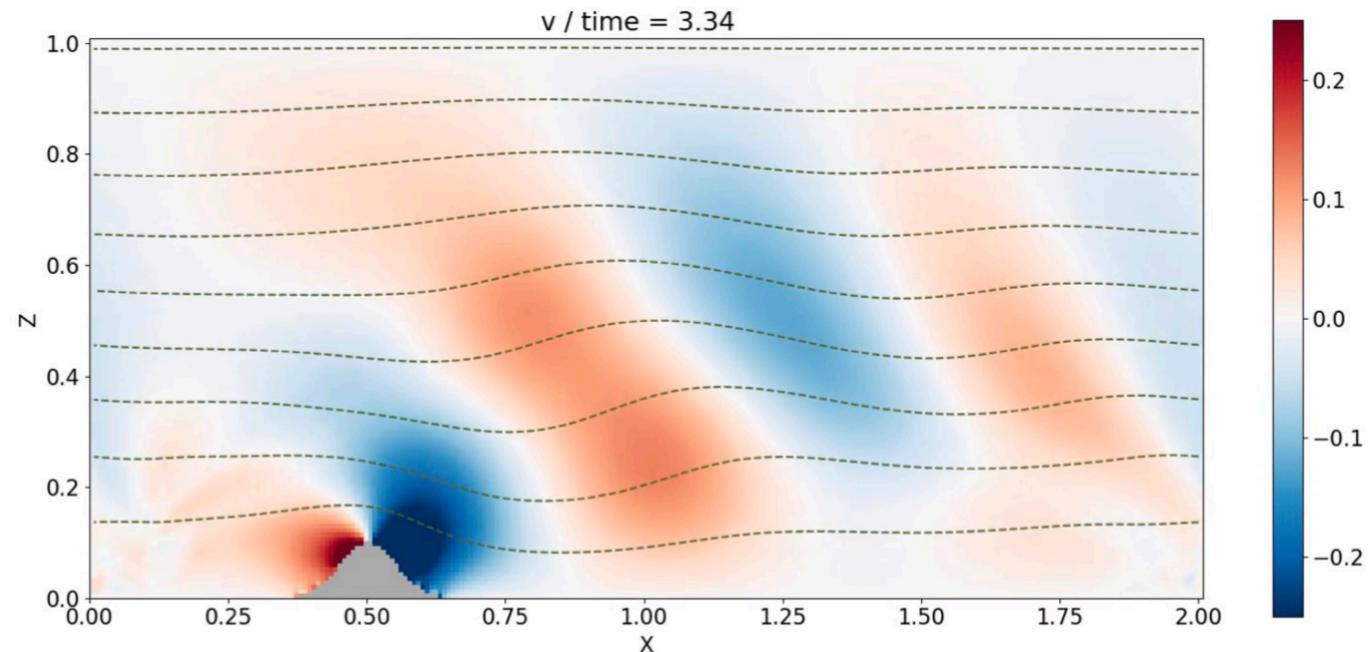
In 1) increase the amplitude until you trigger KHI  
and wave breaking



## Activity: Numerical simulation of Internal waves

Run the Lee wave case:

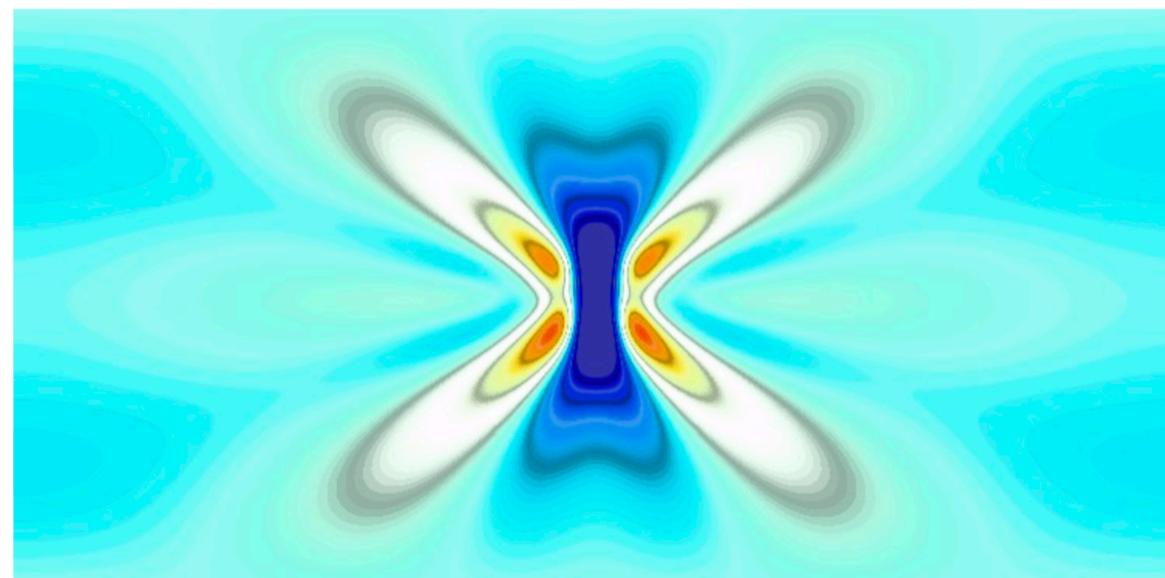
- cd nhom/Leewave
  - python leewave.py
- 
- A flow past a seamount
  - See the wave setting up
  - Increase N and observe mixing downstream



## Activity: Numerical simulation of Internal waves

Run the Internal wave case:

- cd nhom/Internal\_IVP
  - python internalwave.py
- 
- This is an initial value problem: gaussian perturbation of  $b'$
  - Play with the wave amplitude
  - Play with the hydroepsilon (0.2 is the minimum)
  - Play with the size of the initial perturbation



## Activity: Numerical simulation of Internal waves

Run the Internal wave case:

- cd nhom/Internal\_forced
  - Python forcedinternal.py
- 
- This is a localized periodic forcing
  - Play with the forcing frequency (in forcedinternal.py)
  - Generate evanescent waves

