

# The Spectral Hermite Transform as a Means of Simulating Brownian Motion

Mike Johnson

NERS 590: Stochastic Differential Equations

April 23<sup>rd</sup>, 2024

# Part I: What is the Sechmet?

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$$\mathbb{I}_N^M \left[ f(\bar{x}) \right] (x) = \mu \sum_{n=0}^{N-1} \sum_{m=0}^M f_n^m \mathcal{H}_\sigma^m \left( \frac{x - n\mu}{\sigma} \right),$$

$$f_n^m = \frac{1}{m!} \int_{-\infty}^{\infty} f(\bar{x}) \mathcal{H}_\sigma^m \left( \frac{\bar{x} - n\mu}{\sigma} \right) d\bar{x}, \text{ and}$$

$$\mathcal{H}_C^m(x) = \left( -\frac{d}{dx} \right)^m \left[ \frac{1}{C\sqrt{2\pi}} e^{-\frac{x^2}{2}} \right].$$

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- ▶ Looks complicated!! Let's try building up slowly...

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- ▶  $m$ -th order Hermite function:

$$\mathcal{H}_C^m(x) = \left(-\frac{d}{dx}\right)^m \left[ \frac{1}{C\sqrt{2\pi}} e^{-\frac{x^2}{2}} \right] = (m^{th} \text{ order polynomial}) \cdot \frac{1}{C\sqrt{2\pi}} e^{-\frac{x^2}{2}}.$$

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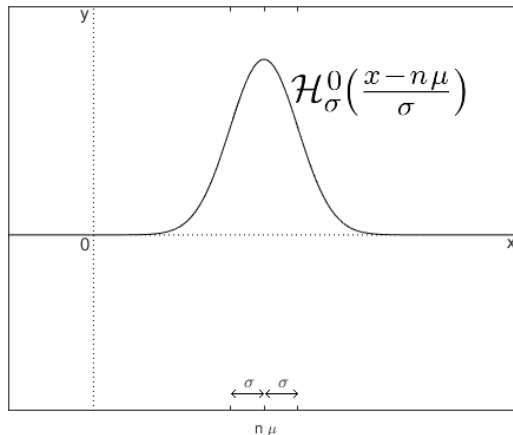
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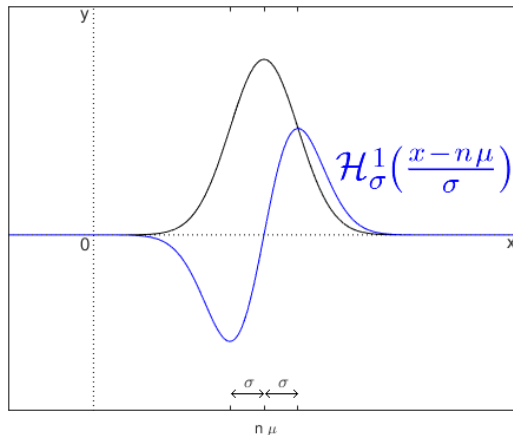
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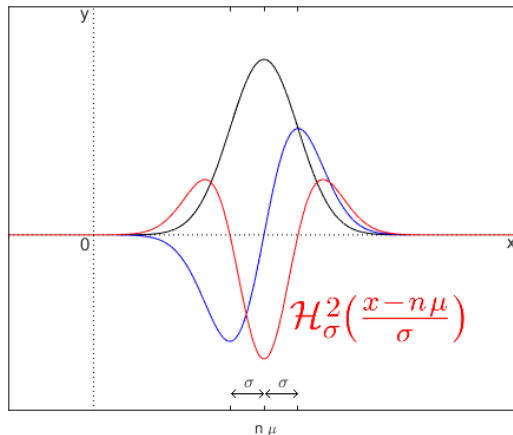
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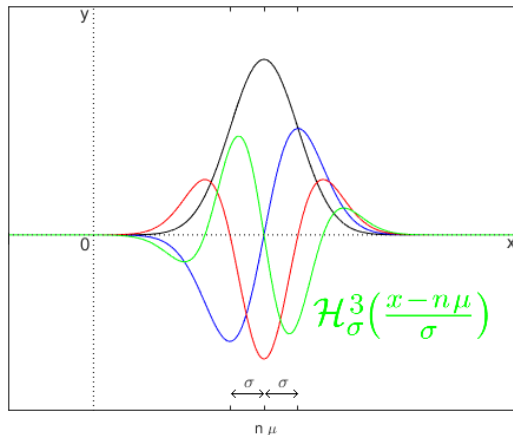
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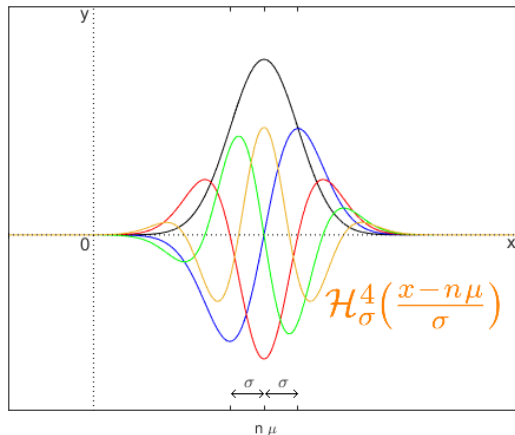
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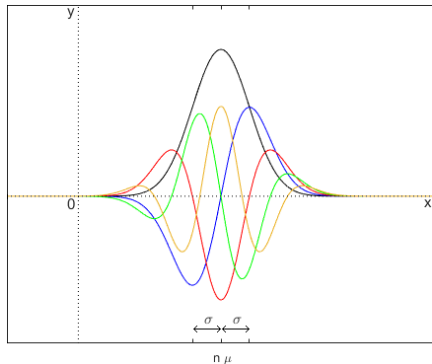
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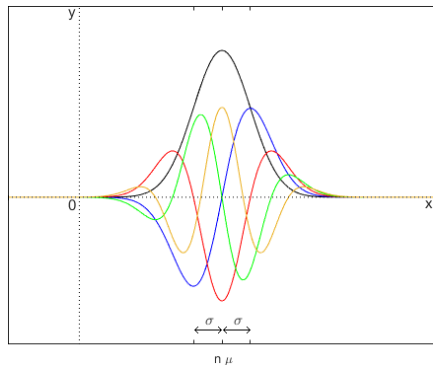
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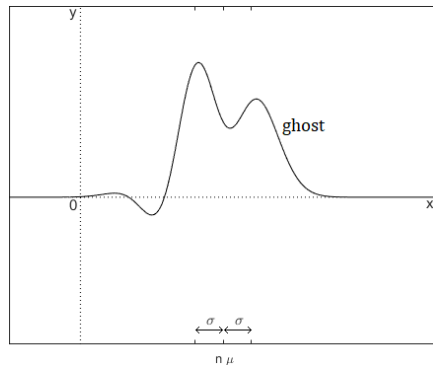


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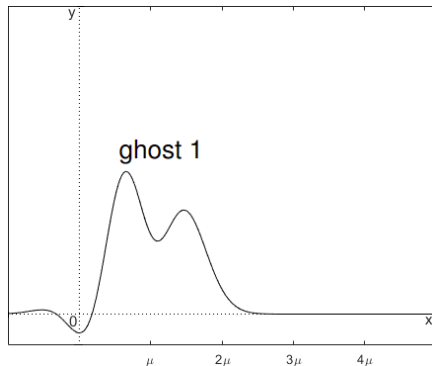
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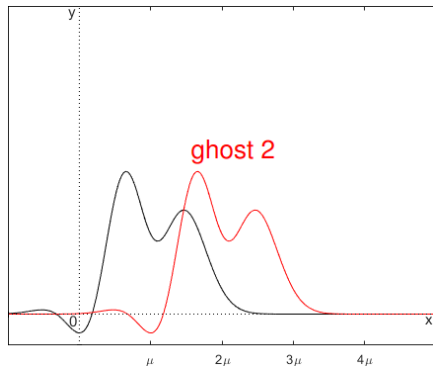
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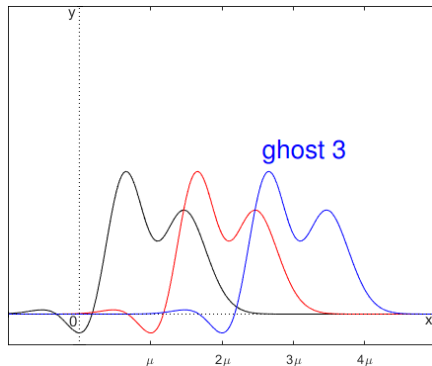
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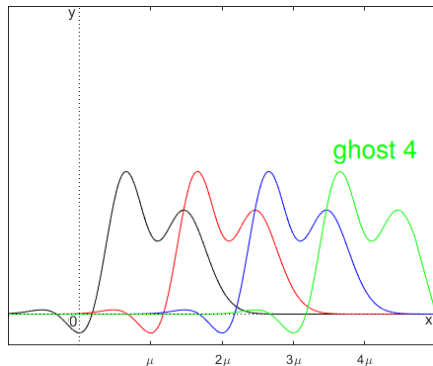
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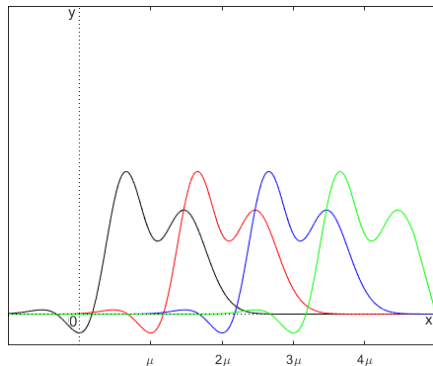
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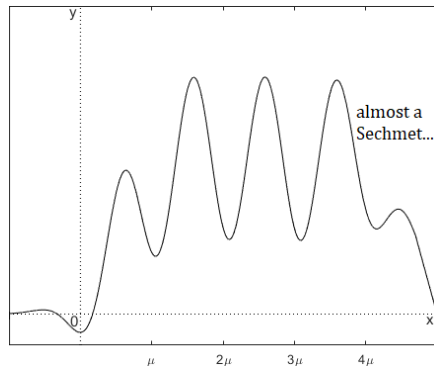


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- ▶ Hopefully this doesn't look like gibberish now:

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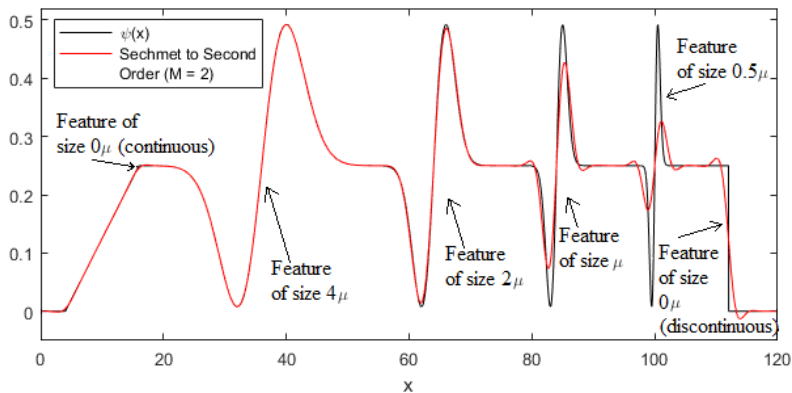
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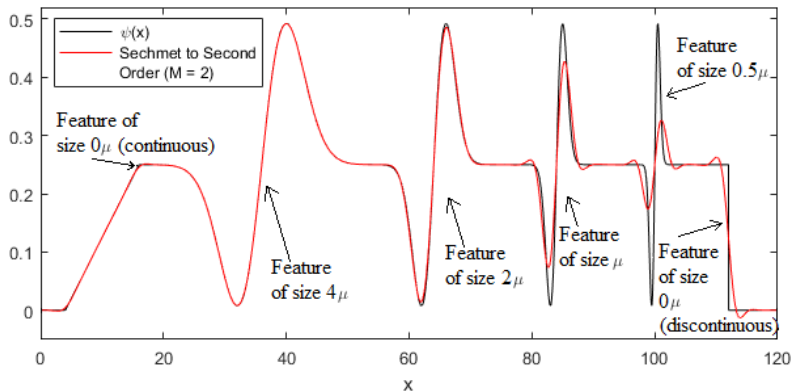
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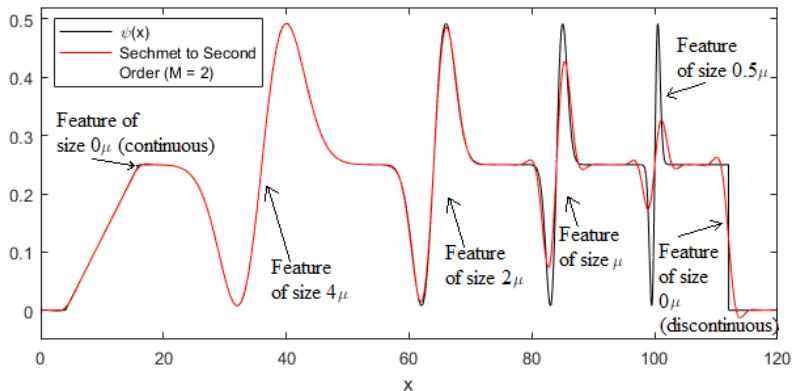
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- Features larger than  $\sim \mu$  are accurate.
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- ▶ Integrals involving Sechmet functions practically always converge.
- ▶ Integrals practically always come out in closed form.
- ▶ Extremely efficient computations via Fourier tricks.
- ▶ Basis functions are localized (unlike e.g. sine or cosine series).

## Part II: Brownian Motion

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- Let's define  $f(x)$  by

$$f'(x) = \mu \sum_{n=0}^{N-1} \sum_{m=0}^M \frac{f_n^m}{\sqrt{m!} \mu} \mathcal{H}_\sigma^m \left( \frac{x - n\mu}{\sigma} \right),$$

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- ▶ Also  $f(x)$  has independent non-overlapping increments when  $N \rightarrow \infty$ . (Not so trivially.)
- ▶ Finally  $f(x) - f(y) \approx f(x - y) \approx \mathcal{N}(0; x - y)$  for large  $N$ . (Very non-trivially!)

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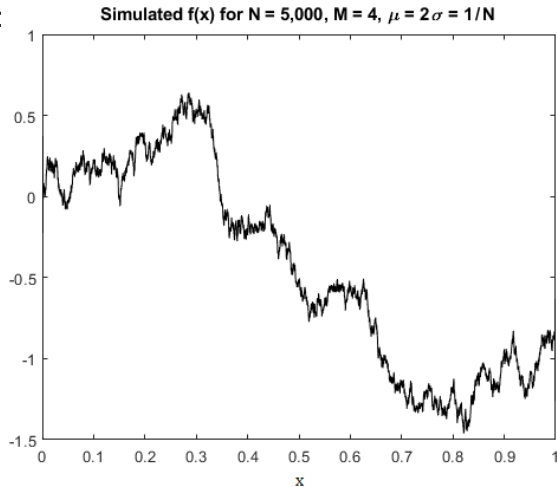
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- Also  $f(x)$  has independent non-overlapping increments when  $N \rightarrow \infty$ . (Not so trivially.)
- Finally  $f(x) - f(y) \approx f(x - y) \approx \mathcal{N}(0; x - y)$  for large  $N$ . (Very non-trivially!)
- In other words,  $f(x) \approx B_x$ !

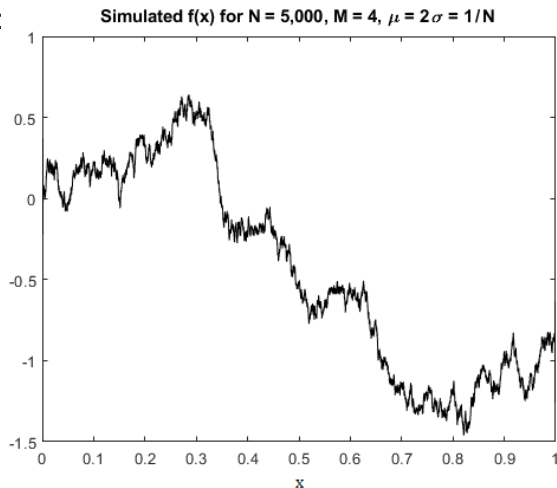
## Part II: Brownian Motion (2/5)

- A picture of  $f(x)$ :



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- ▶ A picture of  $f(x)$ :



- ▶  $f(x)$  looks like Brownian motion, but let's test with some integrals.

## Part II: Brownian Motion (3/5)

- If  $f(x) \approx B_x$ , then  $f(x)$  approximately obeys

$$\int_I \equiv \int_0^x B_s dB_s = \frac{1}{2} (B_x^2 - x) \quad (\text{Itô's Integral}), \text{ and}$$

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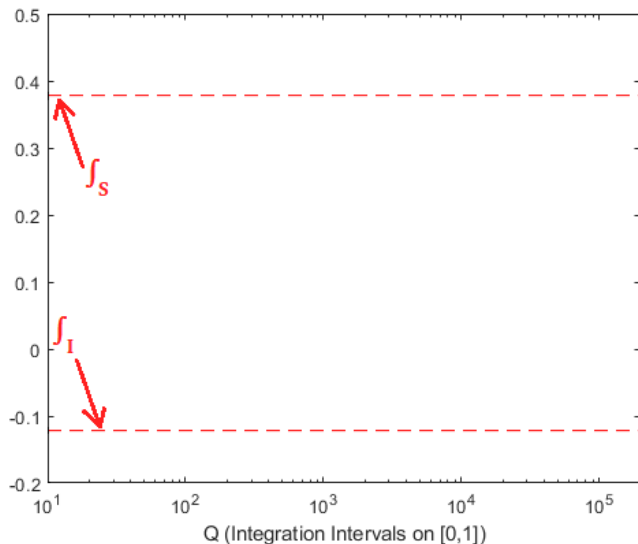
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- ▶ Left-hand Riemann sums over  $f(x) df(x)$  should converge to  $\int_I$ . Let's call these sums  $\Sigma_I$ .
- ▶ Centered Riemann sums over  $f(x) df(x)$  should converge to  $\int_S$ . Let's call these sums  $\Sigma_S$ .

## Part II: Brownian Motion (4/5)

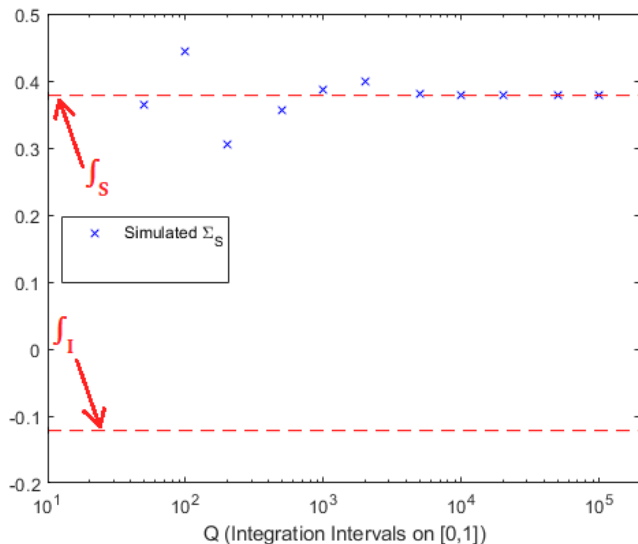
- Numerical results for integrals on  $[0, 1]$  over the  $f(x)$  shown previously:





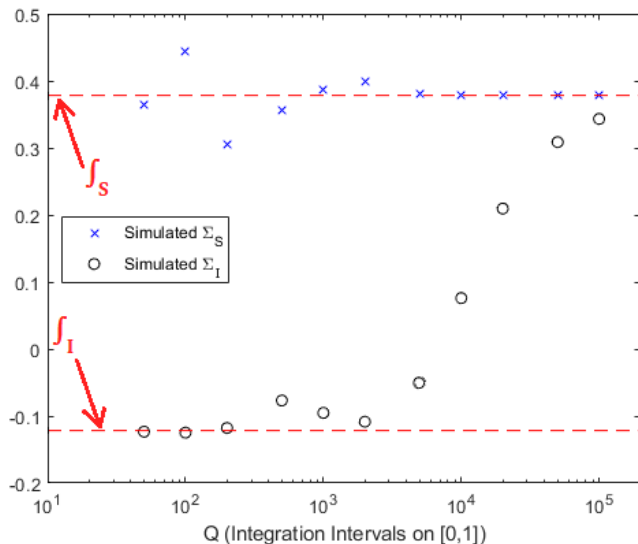
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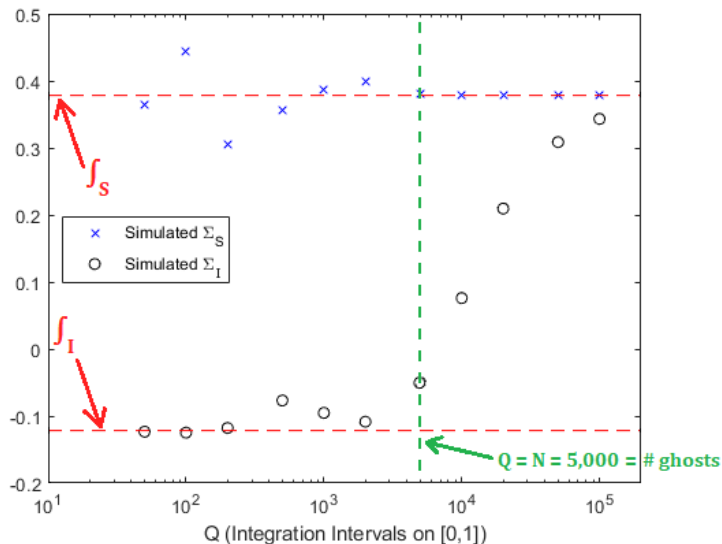
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- ▶ Allows us to cheaply modify particle advection in plasma simulations to include collisional effects via a Langevin-like term. (Future work!)



**Thank you for your attention!**  
**Any questions?**