

Ριζαίοι variational principle

$$\hat{H} |\phi_n\rangle = E_n |\phi_n\rangle \longrightarrow E_n = \frac{\langle \phi_n | \hat{H} | \phi_n \rangle}{\langle \phi_n | \phi_n \rangle}$$

Pro libovolnou funkcií $|\psi\rangle$ platí

$$E_{\text{akt}} = \frac{\langle \psi | \hat{H} | \psi \rangle}{\langle \psi | \psi \rangle} \geq E_0$$

energie
základního
stavu

Dokážeme RVP

$$\begin{aligned} E_{\text{akt}} - E_0 &= \frac{\langle \psi | \hat{H} - E_0 | \psi \rangle}{\langle \psi | \psi \rangle} = \frac{\langle \psi | \hat{H} | \psi \rangle - E_0 \langle \psi | \psi \rangle}{\langle \psi | \psi \rangle} \\ &= \frac{\sum_n \langle \psi | \hat{H} | n \rangle \langle n | \psi \rangle - E_0 \langle \psi | \psi \rangle}{\langle \psi | \psi \rangle} \\ &= \frac{\sum_n (E_n - E_0) |\langle \psi | n \rangle|^2}{\langle \psi | \psi \rangle} \geq 0 \end{aligned}$$

$1 = \sum_n |\langle n | \psi \rangle|^2$
 $\hat{H} |n\rangle = E_n |n\rangle$

E_n

≥ 0

≥ 0

$> 0 = 1$

Pr:

$H = \begin{pmatrix} \varepsilon & J \\ J & \varepsilon \end{pmatrix} \rightarrow$ příkladu star libovolný variacíci principle

$$|\psi\rangle = \begin{pmatrix} a \\ b \end{pmatrix} \quad E_{\text{akt}} = \langle \psi | \hat{H} | \psi \rangle = \frac{1}{a \sqrt{1-a^2}} \begin{pmatrix} \varepsilon & J \\ J & \varepsilon \end{pmatrix} \begin{pmatrix} a \\ \sqrt{1-a^2} \end{pmatrix}$$

$$= \overline{a \sqrt{\dots}} \begin{pmatrix} \varepsilon a + J \sqrt{1-a^2} \\ J a + \varepsilon \sqrt{1-a^2} \end{pmatrix} = \varepsilon a^2 + J a \sqrt{1-a^2} + J a \sqrt{1-a^2} + \varepsilon (1-a^2)$$

$$= \varepsilon + 2 J a \sqrt{1-a^2}$$

$$\frac{\partial}{\partial a} E_{\text{var}} = 0 \Rightarrow 2 J \sqrt{1-a^2} + 2 J a \frac{1}{2} \frac{-2a}{\sqrt{1-a^2}} = 0 \quad | \cdot \sqrt{1-a^2}$$

$$J(1-a^2) - J a^2 = 0$$

$$J - 2 J a^2 = 0 \Rightarrow a^2 = \frac{1}{2}$$

$$\boxed{a = \pm \frac{1}{\sqrt{2}}}$$

$$\boxed{b = \frac{1}{\sqrt{2}}}$$

$$E_{\pm} = \varepsilon + 2 J \left(\pm \frac{1}{\sqrt{2}} \right) \sqrt{1 - \frac{1}{2}} = \varepsilon \pm 2 J \frac{1}{2}$$

$$\boxed{E_{\pm} = \varepsilon \pm J}$$

P_h

$$H = \begin{pmatrix} \varepsilon_1 & J \\ J & \varepsilon_2 \end{pmatrix}$$

$$|\psi\rangle = \cos\varphi |1\rangle + \sin\varphi |2\rangle$$

$$= \begin{pmatrix} \cos\varphi \\ \sin\varphi \end{pmatrix}$$

$$E_{\text{var}} = (\cos\varphi \sin\varphi) \begin{pmatrix} \varepsilon_1 & J \\ J & \varepsilon_2 \end{pmatrix} \begin{pmatrix} \cos\varphi \\ \sin\varphi \end{pmatrix} =$$

$$= (\cos\varphi \sin\varphi) \begin{pmatrix} \varepsilon_1 \cos\varphi + J \sin\varphi \\ J \cos\varphi + \varepsilon_2 \sin\varphi \end{pmatrix}$$

$$\begin{aligned}
 &= \varepsilon_1 \cos^2 \varphi + J \cos \varphi \sin \varphi + J \sin \varphi \cos \varphi + \varepsilon_2 \sin^2 \varphi \\
 &= \varepsilon_1 \cos^2 \varphi + \varepsilon_2 \sin^2 \varphi + 2J \cos \varphi \sin \varphi \quad \frac{\partial}{\partial \varphi} E_{\text{av}} = 0
 \end{aligned}$$

$$-2\varepsilon_1 \sin \varphi \cos \varphi + 2\varepsilon_2 \cos \varphi \sin \varphi + 2J \cos^2 \varphi - 2J \sin^2 \varphi = 0$$

$$2(\varepsilon_2 - \varepsilon_1) \cos \varphi \sin \varphi + 2J(\cos^2 \varphi - \sin^2 \varphi) = 0$$

$$(\varepsilon_2 - \varepsilon_1) \sin 2\varphi + 2J \cos 2\varphi = 0$$

$$\tan 2\varphi = \frac{2J}{\varepsilon_2 - \varepsilon_1}$$

παράχρησιν να γίνει

$$\varphi = \frac{1}{2} \arctan \left(\frac{2J}{\varepsilon_2 - \varepsilon_1} \right)$$

Προβλεπόμενη ενέργεια εκάστου απόσπαστου ατόμου υδριδίου

Προβλεπόμενη κυματική συνάρτηση: γαμμαίαν $\Psi(r) = N e^{-\alpha r^2}$

Hamiltonián ατόμου υδριδίου

$$\hat{H} = -\frac{1}{2} \nabla^2 - \frac{1}{r}$$

$$\frac{\langle \Psi | \hat{H} | \Psi \rangle}{\langle \Psi | \Psi \rangle} = E_{\text{av}}$$

$$\begin{aligned}
 \int_0^\infty dr r^{2m} e^{-\beta r^2} &= \frac{(2m)! \sqrt{\pi}}{2^{2m+1} m! \beta^{m+1/2}} \\
 \int_0^\infty dr r^{2m+1} e^{-\beta r^2} &= \frac{m!}{2 \beta^{m+1}}
 \end{aligned}$$

$$\langle \psi | \psi \rangle = \int_0^{\infty} dr r^2 |\psi(r)|^2$$

$$\begin{aligned} \nabla^2 e^{-\alpha r^2} &= \left(\frac{\partial^2}{\partial r^2} + \frac{2}{r} \frac{\partial}{\partial r} \right) e^{-\alpha r^2} = \frac{2}{r} (-2\alpha r) e^{-\alpha r^2} + (-2\alpha) e^{-\alpha r^2} \\ &\quad + (-2\alpha r)(-2\alpha r) e^{-\alpha r^2} \\ &= 4\alpha \left(\alpha r^2 e^{-\alpha r^2} - e^{-\alpha r^2} - \frac{1}{2} e^{-\alpha r^2} \right) \\ &= 4\alpha \left(\alpha r^2 - \frac{3}{2} \right) e^{-\alpha r^2} \end{aligned}$$

$$\frac{1}{2} N e^{-\alpha r^2} \nabla^2 N e^{-\alpha r^2} = N^2 2\alpha \left(\alpha r^2 - \frac{3}{2} \right) e^{-2\alpha r^2}$$

$$\begin{aligned} \langle \psi | H | \psi \rangle &= \int_0^{\infty} dr r^2 \left[-N^2 2\alpha \left(\alpha r^2 - \frac{3}{2} \right) e^{-2\alpha r^2} - N \frac{1}{r} e^{-2\alpha r^2} \right] \\ &= -N^2 2\alpha^2 \int_0^{\infty} dr r^4 e^{-2\alpha r^2} + N^2 3\alpha \int_0^{\infty} dr r^2 e^{-2\alpha r^2} - N^2 \int_0^{\infty} dr r e^{-2\alpha r^2} \\ &= -N^2 2\alpha^2 \frac{4! \sqrt{\pi}}{2^5 2! (2\alpha)^{5/2}} + N^2 3\alpha \frac{2! \sqrt{\pi}}{2^3 (2\alpha)^{3/2}} - \frac{N^2}{4\alpha} \end{aligned}$$

$$\langle \psi | \psi \rangle = N^2 \int_0^{\infty} dr r^2 e^{-2\alpha r^2} = N^2 \frac{2! \sqrt{\pi}}{2^3 (2\alpha)^{3/2}}$$

$$E = - \frac{\cancel{N^2} 2\alpha^2 \cancel{4} \cdot \cancel{3} \cdot \cancel{2} \sqrt{\pi}}{\cancel{2^5} \cancel{2} (2\alpha)^{5/2}} \cdot \frac{\cancel{2^3} (2\alpha)^{3/2}}{\cancel{N^2} \cancel{2} \sqrt{\pi}} + 3\alpha - \frac{1}{\cancel{4\alpha}} \frac{\cancel{2^3} (2\alpha)^{3/2}}{\cancel{2} \sqrt{\pi}}$$

$$= - \frac{3\alpha^2}{2\alpha} + 3\alpha - \frac{2\sqrt{2}\alpha^{1/2}}{\sqrt{\pi}} = -\frac{3}{2}\alpha + \frac{6}{2}\alpha - 2\sqrt{\frac{2\alpha}{\pi}} \Rightarrow$$

$$\Rightarrow E = \frac{1}{2} \alpha - 2 \sqrt{\frac{2\alpha}{\pi}}$$

$$\frac{\partial}{\partial \alpha} E = 0 = \frac{1}{2} - 2 \frac{1}{2} \frac{\frac{1}{\alpha}}{\sqrt{\frac{2\alpha}{\pi}}} \Rightarrow \sqrt{\frac{2}{\pi}} \frac{1}{\sqrt{\alpha}} = \frac{1}{2} \quad |^2$$

$$\frac{2}{\pi \alpha} = \frac{1}{4} \Rightarrow$$

$$\alpha = \frac{8}{9\pi}$$

$$E_g = \frac{1}{2} \frac{8}{9\pi} - 2 \sqrt{\frac{2 \cdot 8}{9 \cdot 4^2 \pi}} = \frac{4}{9\pi} - 2 \frac{4}{9\pi} = \frac{4}{9\pi} - \frac{8}{9\pi} = -\frac{4}{9\pi} = -\frac{4}{3\pi}$$

$$E_g = -0,4244$$

Exaktní výsledek je $E_{gH} = -\frac{1}{2}$

← v atomových
jednotkách