

# **Note**

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## Part I

# Mathematical Fundamentals

# Chapter 1

## Advent of Mathematical Symbols

- Kronecker delta:

$$\delta_{ij} := \begin{cases} 1, & i = j \\ 0, & i \neq j \end{cases} \quad (1.0.1)$$

- Levi-Civita symbol:

$$\varepsilon_{ijk} := \begin{cases} 1, & (i, j, k) = (1, 2, 3) \text{ or } (2, 3, 1) \text{ or } (3, 1, 2) \\ -1, & (i, j, k) = (3, 2, 1) \text{ or } (2, 1, 3) \text{ or } (1, 3, 2) \\ 0, & \text{else.} \end{cases} \quad (1.0.2)$$

**Example 1.0.1.**

$$(a \times b)_i = \sum_{j,k=1}^3 \varepsilon_{ijk} a_j b_k, \quad (1.0.3)$$

where  $a, b$  are three dimensional vectors and " $\times$ " denotes cross product.

- Nabla symbol:

$$\nabla := \begin{pmatrix} \frac{\partial}{\partial x_1} \\ \frac{\partial}{\partial x_2} \\ \frac{\partial}{\partial x_3} \end{pmatrix}. \quad (1.0.4)$$

- Factorial:  $n! := n \cdot (n-1) \cdot (n-2) \cdots 2 \cdot 1$ .

Recursive definition:  $0! := 1$ ,  $n! := n \cdot (n-1)!$ ,  $n \in \mathbb{N}$ .

- Gamma function:

$$\Gamma(z) := \int_0^\infty x^{z-1} \cdot e^{-x} dx, \quad \operatorname{Re}(z) \geq 0. \quad (1.0.5)$$

---

**Property 1.0.1.**

$$\Gamma(n) = (n-1)!, \quad n \in \mathbb{N}; \quad \Gamma(z+1) = z \cdot \Gamma(z). \quad (1.0.6)$$

- Composition:  $(g \circ f)(x) := g(f(x))$ .

- Sum symbol:  $\sum_{k=1}^n a_k := a_1 + a_2 + \cdots + a_n$ .

Recursive defintion:  $\sum_{k=1}^0 a_k := 0$ ,  $\sum_{k=1}^n a_k := \left( \sum_{k=1}^{n-1} a_k \right) + a_n$ .

- Product:  $\prod_{k=1}^n a_k := a_1 \cdot a_2 \cdot \cdots \cdot a_n$ .

Recursive defintion:  $\prod_{k=1}^0 a_k := 1$ ,  $\prod_{k=1}^n a_k := (\prod_{k=1}^{n-1} a_k) \cdot a_n$

- Restriction:  $f|_A : A \rightarrow Y$ . For  $f : X \rightarrow Y$  and  $A \subseteq X$ , we define  $f|_A(x) = f(x)$  for all  $x \in A$ .

- Pauli matrices:

$$\sigma_1 = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}, \quad \sigma_2 = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}, \quad \sigma_3 = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}. \quad (1.0.7)$$

**Property 1.0.2.** We have  $\sigma_k^2 = I$  and  $\sigma_j \sigma_k - \sigma_k \sigma_j = 2i \varepsilon_{jkl} \sigma_l$ .

- Set brackets:  $\{f(x) | x \in A\}$ .

**Example 1.0.2.**

$$\{2x + 1 | x \in \{0, 1, 2, 3\}\} = \{1, 3, 5, 7\}. \quad (1.0.8)$$

- Big  $O$ :  $f(x) = O(g(x))$ ,  $(x \rightarrow a)$ , which means that  $|f(x)| \leq M \cdot |g(x)|$ , i.e.,  $\limsup_{x \rightarrow a} \frac{f(x)}{g(x)} < \infty$ .

$$x^2 + x + 2 = O(x^2), \quad (x \rightarrow \infty) \quad (1.0.9)$$

$$x^2 + x + 2 = O(x^3), \quad (x \rightarrow \infty). \quad (1.0.10)$$

- Binomial coefficient:

$$\binom{n}{k} = \frac{n \cdot (n-1) \cdots (n-k-1)}{k!} \quad (1.0.11)$$

$$= \frac{n!}{k!(n-k)!}. \quad (1.0.12)$$

- Modulo:  $x \bmod n := r \in [0, n)$  with  $x = n \cdot q + r$  where  $q$  is the integer.

**Example 1.0.3.**

$$5 \bmod 3 = 2 \quad (1.0.13)$$

$$6 \bmod 3 = 0 \quad (1.0.14)$$

$$7.1 \bmod 3 = 1.1 \quad (1.0.15)$$

$$9.7 \bmod 2.1 = 1.3. \quad (1.0.16)$$

- Beta function:

$$\beta(x, y) := \int_0^1 t^{x-1} (1-t)^{y-1} dt, \quad (1.0.17)$$

where  $x, y \in \mathbb{C}$ ,  $\operatorname{Re}(x) > 0$  and  $\operatorname{Re}(y) > 0$ .

**Lemma 1.0.1** (Identity between  $\beta$  func. and  $\Gamma$  func.).

$$\beta(x, y) = \frac{\Gamma(x) \cdot \Gamma(y)}{\Gamma(x+y)}, \quad (1.0.18)$$

where  $\Gamma(\cdot)$  is related to factorial and  $\beta(x, y)$  is related to binomial coefficient.

- Map arrows:  $f : X \rightarrow Y$  where  $X$  is the domain and  $Y$  is the codomain. This map can also be denoted as elementwise-mapping as  $x \mapsto f(x)$ .

**Example 1.0.4.**

$$f := \mathbb{R} \rightarrow \mathbb{R} \quad (1.0.19)$$

$$x \mapsto x^2. \quad (1.0.20)$$

- Little  $o$ :  $f(x) = o(g(x))$ ,  $(x \rightarrow a)$ , which means  $\lim_{x \rightarrow a} \left| \frac{f(x)}{g(x)} \right| = 0$ .

**Example 1.0.5.**

$$8 \cdot x^2 \neq o(x^2), \quad (x \rightarrow \infty) \quad (1.0.21)$$

$$8 \cdot x^2 \neq o(x^3), \quad (x \rightarrow \infty). \quad (1.0.22)$$

- Outer product (Kronecker product for vectors):

$$\begin{pmatrix} v_1 \\ v_2 \end{pmatrix} \otimes \begin{pmatrix} w_1 & w_2 & w_3 \end{pmatrix} = \begin{pmatrix} v_1 w_1 & v_1 w_2 & v_1 w_3 \\ v_2 w_1 & v_2 w_2 & v_2 w_3 \end{pmatrix}, \quad (1.0.23)$$

i.e. matrix entries  $(V \otimes W)_{ij} = v_i \cdot w_j$ .

- Euler's phi function:  $\phi : \mathbb{N} \rightarrow \mathbb{N}$  defined as

$$\phi(n) = \text{count numbers } a \in \mathbb{N} \text{ with} \quad (1.0.24)$$

$$(1) a \leq n \quad (1.0.25)$$

$$(2) \gcd(a, n) = 1 (\text{mutually prime}). \quad (1.0.26)$$

**Example 1.0.6.**

$$\phi(4) = 2 \quad (1.0.27)$$

$$\phi(5) = 4 \quad (1.0.28)$$

$$\phi(p) = p - 1 \text{ for } p \text{ prime.} \quad (1.0.29)$$

- Laplace operator (Laplacian):

$$\Delta f(x) = \frac{\partial^2 f}{\partial x_1^2}(x) + \frac{\partial^2 f}{\partial x_2^2}(x) + \frac{\partial^2 f}{\partial x_3^2}(x), \quad (1.0.30)$$

where  $f : \mathbb{R}^3 \rightarrow \mathbb{R}$ .

- Convolution:  $(f * g)(x) := \int_{-\infty}^{\infty} f(\tau) \cdot g(x - \tau) d\tau$ , where  $f : \mathbb{R} \rightarrow \mathbb{R}$ ,  $g : \mathbb{R} \rightarrow \mathbb{R}$  and  $f * g : \mathbb{R} \rightarrow \mathbb{R}$ .
- Heaviside function:

$$H(x) := \begin{cases} 1, & x \geq 0 \\ 0, & x < 0 \end{cases} \quad (1.0.31)$$

**Property 1.0.3.**

$$H' = \delta. \quad (1.0.32)$$

- Quaternions:

$$\mathbb{H} \supseteq \mathbb{C}, \quad (1.0.33)$$

where  $a, b, c, d \in \mathbb{R}$ , the element in  $\mathbb{H}$  is  $a + i \cdot b + j \cdot c + k \cdot d$  with  $i^2 = -1, j^2 = -1, k^2 = -1, ijk = -1$ .  $\mathbb{H}$  is not commutative in multiplication, i.e.,  $i \cdot j = -j \cdot i$ .

- Infinity:  $\infty$ .

**Example 1.0.7.** In measure theory:  $[0, \infty]$ . We have

$$a + \infty = \infty + a = \infty \text{ for } a \in [a, \infty] \quad (1.0.34)$$

$$(1.0.35)$$

- means equivalence relation. For example,  $x \sim y$  means  $x$  is equivalent to  $y$  for some conditions.

## Chapter 2

# Measure Theory

### 2.1 Sigma algebra

**Definition 2.1.1** (Power set). *We define  $\mathcal{P}(X)$  as the power set of set  $X$ . Assume that set  $X = \{a, b\}$ , the power set  $P(X)$  would be  $\{\emptyset, X, \{a\}, \{b\}\}$*

**Definition 2.1.2** (Sigma algebra).  $\mathcal{A} \subseteq \mathcal{P}(X)$  is called a  $\sigma$ -algebra:

$$(a) \emptyset, X \in \mathcal{A} \quad (2.1.1)$$

$$(b) A \in \mathcal{A} \implies A^c := X \setminus A \in \mathcal{A} \quad (2.1.2)$$

$$(c) A_i \in \mathcal{A}, i \in \mathcal{N} \implies \bigcup_{i=1}^{\infty} A_i \in \mathcal{A}. \quad (2.1.3)$$

**Definition 2.1.3** (Measurable sets).  $A \in \mathcal{A}$  is called a  $\mathcal{A}$ -measurable set.

**Example 2.1.1.**

$$(1) \mathcal{A} = \{\emptyset, X\} \quad (2.1.4)$$

$$(2) \mathcal{A} = \{\mathcal{P}(X)\}. \quad (2.1.5)$$

**Lemma 2.1.1.** *Assume  $\mathcal{A}_i$  is  $\sigma$ -algebra on  $X$ ,  $i \in I$  (index set). Then, we have  $\cap_{i \in I} \mathcal{A}_i$  is also a  $\sigma$ -algebra on  $X$ .*

**Definition 2.1.4** (Sigma algebra generated by  $\mathcal{M}$ ). *For  $\mathcal{M} \subseteq \mathcal{P}(X)$ , there is a smallest  $\sigma$ -algebra that contains  $\mathcal{M}$ :*

$$\sigma(\mathcal{M}) := \cap_{\mathcal{A} \supseteq \mathcal{M}, \text{ a } \sigma\text{-algebra}} \mathcal{A}. \quad (2.1.6)$$

**Example 2.1.2.** *We define  $X = \{a, b, c, d\}$  and  $\mathcal{M} = \{\{a\}, \{b\}\}$ . Then we have*

$$\sigma(\mathcal{M}) = \{\emptyset, X, \{a\}, \{b\}, \{a, b\}, \{b, c, d\}, \{a, c, d\}, \{c, d\}\}. \quad (2.1.7)$$

Here is a very nice picture which shows the connections between the topological space, metric space and vector space. It is from this video:

**Definition 2.1.5** (Topological space). *Topological space  $(X, \tau)$ , where  $X$  is a set and  $\tau$  is a collection of subsets of  $X$ . This  $\tau$  is called the topology on  $X$ . The pair  $(X, \tau)$  is called the topological space. In order to be a topological space, the collection of subsets must satisfy three properties:*

- *Empty set  $\emptyset, X \in \tau$ .*
- *Unions must be in  $\tau$ , i.e.,  $\cup_{i=1}^{\infty} \tau_i \in \tau$ .*
- *Intersections must be in  $\tau$ , i.e.,  $\cap_{i=1}^n \tau_i \in \tau$ .*

**Definition 2.1.6** (Indiscrete topology). *Indiscrete topology is defined as  $\tau = \{\emptyset, X\}$ .*

*Proof.* This can be proved by the definition of 2.1.5. □

**Definition 2.1.7** (Discrete topology). *Discrete topology is the power set 2.1.1 of  $X$ .*

**Proposition 2.1.1.** *Any topology  $\tau$  on  $X$  satisfies the following relation:  $\{\emptyset, X\} \subset \tau \subset \mathcal{P}(X)$ , where  $\{\emptyset, X\}$  is the indiscrete topology 2.1.6 and  $\mathcal{P}(X)$  is the discrete topology 2.1.7.*

**Definition 2.1.8** (Borel sigma algebra). *Let  $(X, \mathcal{T})$  be a topological space 2.1.5 (Let  $X$  be a metric space/Let  $X$  be a subset of  $\mathbb{R}^n$ ; We need "open sets"). We then define  $\mathcal{B}(X)$  is the borel  $\sigma$ -algebra on  $X$  as*

$$\mathcal{B}(X) := \sigma(\mathcal{T}), \quad (2.1.8)$$

*which is the  $\sigma$ -algebra generated by the open sets  $\mathcal{T}$ .*

**Definition 2.1.9** (Borel sets). *Any set contained in Borel  $\sigma$ -algebra is called Borel set. If  $A \in \mathcal{B}(X)$ , then  $A$  is borel set.*

**Proposition 2.1.2.** *Let  $\Omega = [0, 1)$  and  $b \in \Omega$ , then the singleton  $\{b\}$  is a Borel set.*

*Proof.*

$$\{b\} = \cap_{n=1}^{\infty} \left[ \left( b - \frac{1}{n}, b + \frac{1}{n} \right) \cap \Omega \right]. \quad (2.1.9)$$

□

[SZQ: 2023.04.08: This is a standard trick to prove.]

**Proposition 2.1.3.** Let  $\Omega = [0, 1)$  and  $b \in \Omega$ , then  $(a, b]$ ,  $[a, b]$  and  $[a, b)$  are Borel sets.

*Proof.* We write

$$(a, b] = \cap_{n=1}^{\infty} (a, b + \frac{1}{n}) \cap \Omega. \quad (2.1.10)$$

Then we can prove  $(a, b]$  is a borel set.

We can also write  $(a, b]$  as the union of singletons and there use 2.1.2 and the fact that the union of borel sets is also a borel set.  $\square$

**Definition 2.1.10** (Borel measure). A borel measure on  $\mathbb{R}$  is a measure on  $(\mathbb{R}, \mathcal{B}(\mathbb{R}))$ .

**Definition 2.1.11** (Cumulative distribution function). A CDF(cumulative distribution function) is a function  $F : \mathbb{R} \rightarrow \mathbb{R}$  such that

- $F$  is nondecreasing ( $x \leq y \implies F(x) \leq F(y)$ )
- $F$  is right continuous ( $\lim_{a \rightarrow a^+} = F(a)$ )
- $\lim_{x \rightarrow \infty} F(x) = 1$
- $\lim_{x \rightarrow -\infty} F(x) = 0$

**Theorem 1.** • If  $F$  is a CDF then there si a unique Borel probability measure on  $\mathbb{R}$  such that  $P((-\infty, x]) = F(x), \forall x \in \mathbb{R}$ .

- If  $P$  is a Borel probability measure on  $\mathbb{R}$  then there is a unique CDF  $F$  such that  $F(x) = P((-\infty, x]), \forall x \in \mathbb{R}$ .

That is, there is an equivalence between CDFs and Borel probability measure.

## 2.2 What is a measure?

**Definition 2.2.1** (Measure).  $(X, \mathcal{A})$  is called a measurable space, where  $X$  is a set and  $\mathcal{A}$  is a  $\sigma$ -algebra on  $X$ . A map  $\mu : \mathcal{A} \rightarrow [0, \infty] := [0, \infty) + \{\infty\}$  is called a measure if it satisfies:

$$(a) \mu(\emptyset) = 0 \quad (2.2.1)$$

$$(b) \mu(\cup_{i=1}^{\infty} \mathcal{A}_i) = \sum_{i=1}^{\infty} \mu(\mathcal{A}_i) \text{ with } \mathcal{A}_i \cap \mathcal{A}_j = \emptyset, i \neq j \text{ for all } \mathcal{A}_i \in \mathcal{A}. (\sigma - \text{additive}) \quad (2.2.2)$$

**Definition 2.2.2.**  $(X, \mathcal{A}, \mu)$  is called a measure space.

**Example 2.2.1.** Given  $X$  and  $\mathcal{A} = \mathcal{P}(X)$ .

- Counting measure ( $A \in \mathcal{A}$ ) is defined as

$$\mu(A) := \begin{cases} \#A, & A \text{ has finitely many elements} \\ \infty & \text{else} \end{cases} \quad (2.2.3)$$

where  $\#A$  means the number of elements in  $A$ .

Calculation rules in  $[0, \infty]$ :

$$x + \infty := \infty \text{ for all } x \in [0, \infty] \quad (2.2.4)$$

$$x \cdot \infty := \infty \text{ for all } x \in (0, \infty] \quad (2.2.5)$$

$$0 \cdot \infty := 0 \text{ (only true in most cases in measure theory!)} \quad (2.2.6)$$

- Dirac measure for  $p \in X$  is defined as

$$\delta_p(A) := \begin{cases} 1, & p \in A \\ 0, & \text{else} \end{cases} \quad (2.2.7)$$

- We search a measure on  $X \in \mathcal{R}^n$  satisfying:

$$(1) \mu([0, 1]^n) = 1 \quad (2.2.8)$$

$$(2) \mu(x + A) = \mu(A) \text{ for all } x \in \mathcal{R}^n, \quad (2.2.9)$$

which is known as Lebesgue measure where the  $\sigma$ -algebra is not equal to power set.

## 2.3 Not everything is lebesgue measurable

**Measure problem:** search measure  $\mu$  on  $\mathcal{P}(\mathbb{R})$  with:

- (1)  $\mu([a, b]) = b - a$ ,  $b > a$ ,
- (2)  $\mu(x + A) = \mu(A)$ ,  $A \in \mathcal{P}(\mathbb{R})$ ,  $x \in \mathbb{R}$ .

$\implies \mu$  does not exist.

**Claim:** Let  $\mu$  be a measure on  $\mathcal{P}(\mathbb{R})$  with  $\mu((0, 1]) < \infty$  and (2).  $\implies \mu = 0$ .

*Proof.* (a) Definitions:  $I \in (0, 1]$  with equivalence relation on  $I$ :  $x \sim y \iff x - y \in \mathbb{Q}$  i.e.,  $[x] := \{x + r | r \in \mathbb{Q}, x + r \in I\}$ . Following this definition, we have a disjoint decomposition of  $I$  into boxes, possibly uncountable many of them! We then pick one element  $a_n$  from each box  $[x_n]$  and form a set  $A \in I$ , i.e.,  $\{a_1, a_2, \dots\} = A$ . We have  $A \in I$  with prperty:

- (1) For each  $[x]$ , there is an  $a \in A$  with  $a \in [x]$ .
- (2) For all  $a, b \in A$ :  $a, b \in [x] \implies a = b$ .

In uncountable case, the existence of  $A \in I$  with the above property is guaranteed by the axiom of choice of set theory.

We define  $A_n := r_n + A$ , where  $(r_n)_{n \in \mathbb{N}}$  enumeration of  $\mathbb{Q}_n(-1, 1)$ .

(b) We then claim that  $A_n \cap A_m = \emptyset \iff n \neq m$ . The proof is as follows:  $x \in A_n \cap A_m \implies x = r_n + a_n, a_n \in A$  and  $x = r_m + a_m, a_m \in A \implies r_n + a_n = r_m + a_m \implies a_n - a_m = r_n - r_m \in \mathbb{Q} \implies a_n \sim a_m \implies a_m, a_n \in [a_m] \implies a_n = a_m \implies r_n = r_m \implies n = m$ .

(c) We claim that  $(0, 1] \subseteq \bigcup_{n \in \mathbb{N}} A_n \subseteq (-1, 2]$ . The proof is as follows:

Assume now:  $\mu$  measure on  $\mathcal{P}(\mathbb{R})$  with  $\mu((0, 1]) < \infty$  and (2).

By (2):  $\mu(1 + A) = \mu(A)$  for all  $n \in \mathbb{N}$ .

By (c): we have

$$\mu((0, 1]) \leq \mu(\bigcup_{n \in \mathbb{N}} A_n) \leq \mu((-1, 2]) \quad (2.3.1)$$

We know:  $\mu((0, 1]) =: C < \infty$ . By using (2) and  $\sigma$ -additivity, we get  $\mu((-1, 2]) = \mu((-1, 0] \cup (0, 1] \cup (1, 2]) = 3C \implies \text{2.3.1, (b)} C \leq \sum_{n=1}^{\infty} \mu(A_n) \leq 3C \implies C \leq \sum_{n=1}^{\infty} \mu(A) \leq 3C \implies \mu(A) = 0 \implies C = 0$  (hence  $\mu((0, 1]) = 0 \implies \mu(\mathbb{R}) = \mu(\bigcup_{n \in \mathbb{Z}} (m, m+1]) = 0 \implies \mu = 0$ ).  $\square$

## 2.4 Measurable maps

**Definition 2.4.1** (Measurable maps).  $(\Omega_1, \mathcal{A}_1)$  and  $(\Omega_2, \mathcal{A}_2)$  are measurable spaces.  $f : \Omega_1 \rightarrow \Omega_2$  is a measurable map w.r.t.  $\mathcal{A}_1$  and  $\mathcal{A}_2$  if  $f^{-1}(A_2) \in \mathcal{A}_1$  for all  $A_2 \in \mathcal{A}_2$ .

**Example 2.4.1.** •  $(\Omega, \mathcal{A})$  and  $(\mathbb{R}, \mathcal{B}(\mathbb{R}))$  are two measurable spaces. We define characteristic function (aksi indicator function) as  $\chi_A : \Omega \rightarrow \mathbb{R}$ , where

$$\chi_A(w) := \begin{cases} 1, & w \in A \\ 0, & w \notin A \end{cases} \quad (2.4.1)$$

For all measurable  $A \in \mathcal{A}$ ,  $\chi_A$  is a measurable map. We have

$$\chi_A^{-1}(\emptyset) = \emptyset \in \mathcal{A}, \quad \chi_A^{-1}(\mathbb{R}) = \Omega \in \mathcal{A} \quad (2.4.2)$$

$$\chi_A^{-1}(\{\textcolor{red}{A}\}) = A, \quad \chi_A^{-1}(\{0\}) = A^c \in \mathcal{A}. \quad (2.4.3)$$

- Composition of measurable maps.

**Lemma 2.4.1.**  $(\Omega_1, \mathcal{A}_1), (\Omega_2, \mathcal{A}_2), (\Omega_3, \mathcal{A}_3)$  are measurable space. We define  $\Omega_1 \xrightarrow{f} \Omega_2 \xrightarrow{g} \Omega_3$ . Then  $f, g$  are measurable implies  $g \circ f$  is measurable.

*Proof.*

$$(g \circ f)^{-1}(A_3) = f^{-1}(g^{-1}(A_3)) \quad (2.4.4)$$

$$\in \mathcal{A}_1 \quad (2.4.5)$$

□

### Important measurable maps

**Lemma 2.4.2.**  $(\Omega, \mathcal{A})$  and  $(\mathbb{R}, \mathcal{B}(\mathbb{R}))$  are measurable spaces.  $f, g : \Omega \rightarrow \mathbb{R}$  are measurable maps indicates that  $f + g, f - g, f \cdot g, |f|$  are measurable maps.

## 2.5 Lebesgue integral

**Example 2.5.1.** Define Characteristic function  $\chi_A : X \rightarrow \mathbb{R}$ ,  $A \in \mathcal{A}$ . We define  $I(A) := \mu(A)$ . Surprisingly,  $I(A)$  is nothing but the integral of  $\chi_A$  over  $A$ .

**Definition 2.5.1** (Simple/Step/Staircsae functions,...). For  $A_1, A_2, \dots, A_n \in \mathcal{A}$ , and  $c_1, c_2, \dots, c_n \in \mathbb{R}$ . We define

$$f(x) := \sum_{i=1}^n c_i \cdot \chi_{A_i}(x). \quad (2.5.1)$$

We then have  $f(x)$  is measurable and the integraal of  $f$  is defined as  $I(f) := \sum_{i=1}^n c_i \mu(A_i)$ .

**Remark 2.5.1.** The problem of the integral  $I(f)$  is that it is undefined when  $\mu(A_i) = \infty$ . The problem can be solved by exclude  $\infty$  by defintion or the following way.

**Definition 2.5.2** (Lebesgue integral). Define  $S^+ := \{f : X \rightarrow \mathbb{R} \mid f \text{ simple function, } f \geq 0\}$ .  $f \in S^+$  and choose representation  $f(x) = \sum_{i=1}^n c_i \chi_{A_i}(x)$ ,  $c_i \geq 0$ . The lebesgue integral of  $f$  w.r.t.  $\mu$  is defined as

$$\int_X f(x) \, d\mu(x) = \int_X f \, d\mu \quad (2.5.2)$$

$$= I(f) \quad (2.5.3)$$

$$= \sum_{i=1}^n c_i \cdot \mu(A_i) \quad (2.5.4)$$

$$= [0, \infty]. \quad (2.5.5)$$

**Property 2.5.1.** •  $I(\alpha f + \beta g) = \alpha I(f) + \beta I(g)$ ,  $\alpha, \beta \geq 0$ .

- $f \leq g \implies I(f) \leq I(g)$  (monotonicity)

**Definition 2.5.3.** Define a measurable map  $f : X \rightarrow [0, \infty)$ .  $h = \sum_{i=1}^n c_i \cdot \chi_{A_i}$ . The lebesgue integral of  $f$  w.r.t.  $\mu$  is defined as

$$\int_X f \, d\mu := \sup \{I(h) | h \in S^+, h \leq f\} \quad (2.5.6)$$

$$\in [0, \infty]. \quad (2.5.7)$$

$f$  is called  $\mu$ -integrable if  $\int_X f \, d\mu < \infty$ .

**Property 2.5.2.** Define measurable maps  $f, g : X \rightarrow [0, \infty)$ , we have

- 1.  $f = g$  for  $\mu$ -almost everywhere(a.e.), which satisfies  $\mu(\{x \in X | f(x) \neq g(x)\}) = 0 \implies \int_X f \, d\mu = \int_X g \, d\mu$ .
- 2.  $f \leq g$  for  $\mu$  a.e.  $\implies \int_X f \, d\mu \leq \int_X g \, d\mu$
- 3.  $f = 0$  for  $\mu$ -a.e.  $\iff \int_X f \, d\mu = 0$ .

*Proof of 2.: monotonicity.* Let  $h := X \rightarrow [0, \infty)$  be a simple function, i.e.,

$$h(x) = \sum_{i=1}^n c_i \chi_{A_i}(x) \quad (2.5.8)$$

$$= \sum_{t \in h(X)} t \cdot \chi_{\{x \in X | h(x) = t\}}. \quad (2.5.9)$$

Let  $X = \tilde{X}^c \cup \tilde{X}$  with  $\mu(\tilde{X}^c) = 0$ ,

$$\tilde{h}(x) := \begin{cases} h(x), & x \in \tilde{X} \\ a, & x \in \tilde{X}^c \end{cases} \quad (2.5.10)$$

$$\tilde{h}(x) = \sum_{t \in h(X)} t \cdot \chi_{\{x \in \tilde{X} | h(x) = t\}} + a \cdot \chi_{\tilde{X}^c} \quad (2.5.11)$$

$$I(\tilde{h}) = \sum_{t \in h(X)} t \cdot \mu(\{x \in \tilde{X} | h(x) = t\}) + a \cdot \mu(\tilde{X}^c) \quad (2.5.12)$$

$$= \sum_{t \in h(X)} t [\mu(\{x \in \tilde{X} | h(x) = t\}) + \mu(\{x \in \tilde{X}^c | h(x) = t\})] \quad (2.5.13)$$

$$= \sum_{t \in h(X)} t [\mu(\{x \in \tilde{X} | h(x) = t\} \cup \{x \in \tilde{X}^c | h(x) = t\})] \quad (2.5.14)$$

$$I(h) = \sum_{t \in h(X) \setminus \{0\}} t \cdot \mu(\{x \in X | h(x) = t\}). \quad (2.5.15)$$

We define

$$\tilde{X} := \{x \in X | f(x) \leq g(x)\}, \quad (2.5.16)$$

$$\mu(\tilde{X}^c) = 0 \quad (2.5.17)$$

$$\int_X f \, d\mu = \sup \{I(h) | h \in S^+, h \leq f\} \quad (2.5.18)$$

$$= \sup \{I(\tilde{h}) | \tilde{h} \in S^+, \tilde{h} \leq f \text{ on } \tilde{X}\} \quad (2.5.19)$$

$$\leq \sup \{I(\tilde{h}) | \tilde{h} \in S^+, h \leq g \text{ on } \tilde{X}\} \quad (2.5.20)$$

$$= \sup \{I(h) | h \in S^+, h \leq g\} \quad (2.5.21)$$

$$= \int_X g \, d\mu. \quad (2.5.22)$$

□

**Theorem 2** (Monotone convergence theorem). *( $X, \mathcal{A}, \mu$ ) measurable spaces,  $f_n : X \rightarrow [0, \infty]$ , ( $f : X \rightarrow [0, \infty]$ ) measurable for all  $n \in \mathbb{N}$  with*

$$f_1 \leq f_2 \leq f_3 \leq \dots \quad \mu - a.e. \quad (2.5.23)$$

$$\left( \lim_{n \rightarrow \infty} \int_X f_n \, d\mu = \int_X f \, d\mu \quad \mu - a.e. (x \in X) \right) \quad (2.5.24)$$

This implies that

$$\lim_{n \rightarrow \infty} \int_X f_n \, d\mu = \int_X \lim_{n \rightarrow \infty} f_n \, d\mu. \quad (2.5.25)$$

*Proof.*  $\int_X f_1 \, d\mu \leq \int_X f_2 \, d\mu \leq \dots$  and  $\int_X f_n \, d\mu \leq \int_X f \, d\mu$  for  $n \in \mathbb{N}$ . Then we have

$$\lim_{n \rightarrow \infty} \int_X f_n \, d\mu \leq \int_X f \, d\mu, \quad (2.5.26)$$

which is the first part of 2.5.25.

Let  $h$  be a simple function  $0 \leq h \leq f$  and  $\varepsilon > 0$ . We define

$$X_n := \{x \in X | f_n(x) \geq (1 - \varepsilon)h(x)\} \quad (2.5.27)$$

with  $\cup_{n=1}^{\infty} X_n = \tilde{X}$ , and  $\mu(\tilde{X}^c) = 0$ . We have

$$\int_X f_n \, d\mu \geq \int_{X_n} f_n \, d\mu \geq \int_{X_n} (1 - \varepsilon)h \, d\mu \quad (2.5.28)$$

$$\lim_{n \rightarrow \infty} \int_X f_n \, d\mu \geq \lim_{n \rightarrow \infty} \int_{X_n} (1 - \varepsilon)h \, d\mu \quad (2.5.29)$$

$$= \int_{\tilde{X}} (1 - \varepsilon)h \, d\mu \quad (2.5.30)$$

$$= \int_X (1 - \varepsilon)h \, d\mu. \quad (2.5.31)$$

This implies

$$\lim_{n \rightarrow \infty} \int_X f_n \, d\mu \geq \int_X h \, d\mu, \quad (2.5.32)$$

since  $\varepsilon > 0$  arbitrarily. Then we have

$$\lim_{n \rightarrow \infty} \int_X f_n \, d\mu \geq \int_X f \, d\mu, \quad (2.5.33)$$

since  $h$  is arbitrary and  $h \leq f$ , which is second part of 2.5.25.  $\square$

**Applications** Given a series  $(g_n)_{n \in \mathbb{N}}$ ,  $g_n : X \rightarrow [0, \infty]$  measurable for all  $n$ . Then we have  $\sum_{n=1}^{\infty} g_n : X \rightarrow [0, \infty]$  measurable and

$$\int_X \sum_{n=1}^{\infty} g_n \, d\mu = \sum_{n=1}^{\infty} \int_X g_n \, d\mu, \quad (2.5.34)$$

which means the integral and sum can exchange.

## 2.6 Fatou' lemma

**Lemma 2.6.1** (Fatou' lemma). *Given  $(X, \mathcal{A}, \mu)$  measurable space,  $f_n : X \rightarrow [0, \infty]$  measurable for all  $n \in \mathbb{N}$ . Then we have*

$$\int_X \liminf_{n \rightarrow \infty} f_n \, d\mu \leq \liminf_{n \rightarrow \infty} \int_X f_n \, d\mu. \quad (2.6.1)$$

**Remark 2.6.1.**  $\liminf_{n \rightarrow \infty} f_n : X \rightarrow [0, \infty]$  is a function. This is

$$g(x) := \left( \liminf_{n \rightarrow \infty} f_n \right) (x) \quad (2.6.2)$$

$$:= \lim_{n \rightarrow \infty} \left( \inf_{k \geq n} f_k(x) \right) \quad (2.6.3)$$

$$\in [0, \infty] \quad (2.6.4)$$

$$g_n(x) := \inf_{k \geq n} f_k(x). \quad (2.6.5)$$

We have

$$g_1 \leq g_2 \leq g_3 \leq \cdots, \quad (2.6.6)$$

which is monotonically increasing. All these functions are measurable.

*Proof.*

Since (2),

$$\int_X \lim_{n \rightarrow \infty} g_n \, d\mu = \lim_{n \rightarrow \infty} \int_X g_n \, d\mu \quad (2.6.7)$$

$$= \liminf_{n \rightarrow \infty} \int_X g_n \, d\mu. \quad (2.6.8)$$

We know that  $g_n \leq f_n$  for all  $n \in \mathbb{N}$ . By (2.5.2), we have

$$\int_X g_n \, d\mu \leq \int_X f_n \, d\mu, \quad (2.6.9)$$

for all  $n \in \mathbb{N}$ . Then we have

$$\int_X \liminf_{n \rightarrow \infty} f_n \, d\mu = \liminf_{n \rightarrow \infty} \int_X g_n \, d\mu \quad (2.6.10)$$

$$\leq \liminf_{n \rightarrow \infty} \int_X f_n \, d\mu. \quad (2.6.11)$$

□

## 2.7 Lebesgue's dominated convergence theorem

$(X, \mathcal{A}, \mu)$ ,  $\mathcal{L}^1 := \{f : X \rightarrow \mathbb{R} \text{ measurable} \mid \int_X |f|^1 \, d\mu < \infty\}$ . For  $f \in \mathcal{L}^1(\mu)$ , write  $f = f^+ - f^-$ , where  $f^+, f^- \geq 0$ . Define  $\int_X f \, d\mu := \int_X f^+ \, d\mu - \int_X f^- \, d\mu$ .

**Theorem 3** (Lebesgue's dominated convergence theorem).  $f_n : X \rightarrow \mathbb{R}$  measurable for all  $n \in \mathbb{N}$ .  $f : X \rightarrow \mathbb{R}$  with  $f(x) \text{ for } x \in X$  ( $\mu$ -a.e.) and  $|f_n| \leq g$  with  $g \in \mathcal{L}^1(\mu)$  for all  $n \in \mathbb{N}$ , where  $g$  is called integral majorant. Then: we have  $f_1, f_2, \dots \in \mathcal{L}^1(\mu)$ ,  $f \in \mathcal{L}^1(\mu)$  and

$$\lim_{n \rightarrow \infty} \int_X f_n \, d\mu = \int_X f \, d\mu. \quad (2.7.1)$$

*Proof.*

$$|f_n| \leq g \xrightarrow{\text{monotonicity}} \int_X g \, d\mu < \infty \quad (2.7.2)$$

$$\implies f_1, f_2, \dots \in \mathcal{L}^1(\mu) \quad (2.7.3)$$

$$|f| \leq g \text{ for } \mu\text{-a.e.} \implies f \in \mathcal{L}^1(\mu) \quad (2.7.4)$$

We will show  $\int_X |f_n - f| \, d\mu \xrightarrow{n \rightarrow \infty} 0$ .

$$|f_n - f| \leq |f_n| + |f| \leq 2g \quad (2.7.5)$$

$$\implies h_n := 2g - |f_n - f| \geq 0 \quad (2.7.6)$$

Hence:  $h_n : X \rightarrow [0, \infty]$  measurable for all  $n \in \mathbb{N}$ . Then by (2.6.1),

$$\implies \int_X \liminf_{n \rightarrow \infty} h_n \, d\mu \leq \liminf_{n \rightarrow \infty} \int_X h_n \, d\mu \quad (2.7.7)$$

$$\implies \int_X 2g \, d\mu \leq \int_X 2g \, d\mu - \limsup_{n \rightarrow \infty} \int_X |f_n - f| \, d\mu \quad (2.7.8)$$

$$\implies 0 \leq \liminf_{n \rightarrow \infty} \int_X |f_n - f| \, d\mu \leq \limsup_{n \rightarrow \infty} \int_X |f_n - f| \, d\mu \leq 0 \quad (2.7.9)$$

$$\implies \quad (2.7.10)$$

Limits exists and  $\lim_{n \rightarrow \infty} |f_n - f| \, d\mu = 0$ . We conclude that

$$0 \leq \left| \int_X f_n \, d\mu - \int_X f \, d\mu \right| = \left| \int_X (f_n - f) \, d\mu \right| \leq \int_X |f_n - f| \, d\mu \xrightarrow{n \rightarrow \infty} 0, \quad (2.7.11)$$

$$(2.7.12)$$

where the third inequality is due to the integral's triangle inequality.

$$\implies \lim_{n \rightarrow \infty} \int_X f_n \, d\mu = \int_X f \, d\mu. \quad (2.7.13)$$

□

## 2.8 Caratheodory's extension theorem

**Theorem 4** (Caratheodory's extension theorem). *X set,  $\mathcal{A} \in \mathcal{P}(X)$  semiring of sets. A map  $\mu : \mathcal{A} \rightarrow [0, \infty]$ . Note that  $\mu$  is not a measure, it is called A pre-measure.*

- Then  $\mu$  has an extension  $\tilde{\mu} : \sigma(\mathcal{A}) \rightarrow [0, \infty]$ , where  $\tilde{\mu}$  is a measure and  $\sigma(\mathcal{A})$  is a  $\sigma$ -algebra generated by  $\mathcal{A}$ , i.e.,  $\mu(A) = \tilde{\mu}(A)$ .
- If there is sequence  $(S_j)$  with  $S_j \in \mathcal{A}$ ,  $\cup_{j=1}^{\infty} S_j = X$ , then the extension  $\tilde{\mu}$  from (a) is unique. ( $\tilde{\mu}$  is also  $\sigma$ -finite)

**Definition 2.8.1** (Semiring set). *Semiring of sets  $\mathcal{A} \subseteq \mathcal{P}(X)$ :*

- $\emptyset \in \mathcal{A}$  (as for  $\sigma$ -algebra)
- $A, B \in \mathcal{A} \implies A \cap B \in \mathcal{A}$
- For  $A, B \in \mathcal{A}$ , there are pairwise disjoint sets  $S_1, S_2, \dots, S_n \in \mathcal{A}$ :  $\cup_{j=1}^n S_j = A \setminus B$

**Example 2.8.1.**  $A := \{[a, b] | a, b \in \mathbb{R}, a \leq b\}$  not a  $\sigma$ -algebra because  $\mathbb{R} \notin \mathcal{A}$ . But  $\sigma(\mathcal{A}) = \mathcal{B}(\mathbb{R})$  (Borel  $\sigma$ -algebra). Check that  $\mathcal{A}$  is semiring set:

- $\emptyset \in \mathcal{A}$

•

$$[a, b) \cap [c, d) = \begin{cases} \emptyset, & b \leq c, d \leq a \\ [c, b), & c \in [a, b), d \notin [a, b) \\ \dots \end{cases} \quad (2.8.1)$$

•

$$[a, b) \setminus [c, d) = \begin{cases} [a, b), & d \leq a, b \leq c \\ [a, c), & c \in [a, b), d \notin [a, b) \\ [a, c) \cup [d, b), & c > a, d < b \\ \dots \end{cases} \quad (2.8.2)$$

**Definition 2.8.2** (Pre-measure).  $\mu : \mathcal{A} \rightarrow [0, \infty]$  with  $\mathcal{A}$  semiring os sets:

- $\mu(\emptyset) = 0$
- $\mu(\cup_{j=1}^{\infty} A_j) = \sum_{j=1}^{\infty} \mu(A_j)$ , for  $A_j \in \mathcal{A}$ ,  $A_i \cap A_j = \emptyset$  for  $i \neq j$  and  $\cup_{j=1}^{\infty} A_j \in \mathcal{A}$ .

**Application:**  $A := \{[a, b) | a, b \in \mathbb{R}, a \leq b\}$ ,  $\mu : \mathcal{A} \rightarrow [0, \infty]$ ,  $\mu([a, b)) = b - a$  is a pre-measure (We can check by the definition of pre-measure). Then by (4), there is a unique extension to  $\mathcal{B}(\mathbb{R}) \rightarrow \text{lebesgue measure}$ .

## 2.9 Lebesgue-Stieltjes measures

$F : \mathbb{R} \rightarrow \mathbb{R}$  monotonically increasing (non-decreasing).  $[a, b)$  is the length of the interval. Now we consider new kinds of intervals:

$$F(b^-) - F(a^-) =: \mu_F([a, b)), \quad (2.9.1)$$

where  $F(a^-) := \lim_{\varepsilon \rightarrow 0^+} F(a - \varepsilon)$ . Alternatively, we also have

$$F(b^+) - F(a^+) =: \mu_F((a, b]), \quad (2.9.2)$$

where  $F(a^+) := \lim_{\varepsilon \rightarrow 0^+} F(a + \varepsilon)$ . We consider the previous one hereafter.

**Definition 2.9.1.**  $\mathcal{A} := \{[a, b) : a, b \in \mathbb{R}, a \leq b\}$  semiring of sets. Then by Caratheodory' theorem, we have that there exists exactly one measure

$$\mu_F : \mathcal{B}(\mathbb{R}) \rightarrow [0, \infty] \quad (2.9.3)$$

with  $\mu_F([a, b))$

$$\cdot \quad (2.9.4)$$

**Example 2.9.1.** •  $F(x) = x$ ,  $\mu_F([a, b]) = b - a \rightarrow$  Lebesgue measure.

- $F(x) = 1$ ,  $\mu_F([a, b]) = 0 \rightarrow$  zero measure.

•

$$F(x) = \begin{cases} 0, & x < 0 \\ 1, & x \geq 0 \end{cases} \quad (2.9.5)$$

$\mu_F([-ε, ε]) = 1 \rightarrow$  Dirac measure  $δ_0$ .

- $F : \mathbb{R} \rightarrow \mathbb{R}$  monotonically increasing + continuously differentiable. Then we have

$$F' : \mathbb{R} \rightarrow [0, \infty) \quad (2.9.6)$$

and

$$\mu_F([a, b]) = F(b) - F(a) \quad (2.9.7)$$

$$= \int_a^b F'(x) \, dx, \quad (2.9.8)$$

which implies

$$\mu_F : A \longmapsto \int_A F'(x) \, dx, \quad (2.9.9)$$

where  $F'(x)$  is called the density function.

## 2.10 Radon-Nikodym theorem and Lebesgue's decomposition theorem

$(X, \mathcal{A}, \lambda)$  measure space. Special case:  $X = \mathbb{R}$ ,  $\mathcal{A} = \mathcal{B}(\mathbb{R})$ , and  $\lambda$  is lebesgue measure. Recall that  $\lambda([a, b]) = b - a$ . Another measure  $\mu : \mathcal{B}(\mathbb{R}) \rightarrow [0, \infty]$ . We will look how  $\mu$  acts w.r.t. the given reference measure: lebesgue measure.

**Definition 2.10.1.** •  $\mu$  is called absolutely continuous (w.r.t.  $\lambda$ ) if  $\lambda(A) = 0 \implies \mu(A) = 0$  for all  $A \in \mathcal{B}(\mathbb{R})$ . One writes:  $\mu \ll \lambda$ .

- $\mu$  is called singular (w.r.t.  $\lambda$ ) if there is  $N \in \mathcal{B}(\mathbb{R})$  with  $\lambda(N) = 0$  and  $\mu(N^c) = 0$ . One writes:  $\mu \perp \lambda$ .

**Example 2.10.1.**  $\delta_0$  Dirac measure ( $\delta_0(\{0\}) = 1 \implies \delta_0 \perp \lambda$  (Choose  $N = \{0\}$ )).

**Theorem 5** (Lebesgue's decomposition theorem).  $\mu : \mathcal{B}(\mathbb{R}) \rightarrow [0, \infty]$  ( $\sigma$ -finite)

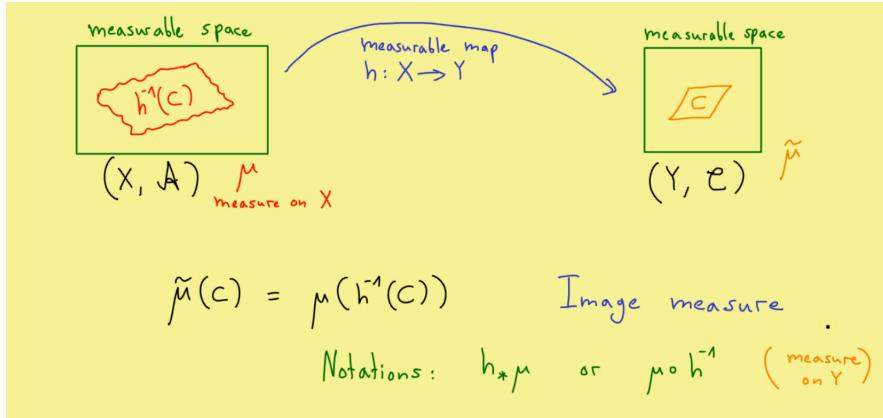
- There are measures (uniquely determined)  $\mu_{ac}, \mu_s : \mathcal{B}(\mathbb{R}) \rightarrow [0, \infty]$  with  $\mu = \mu_{ac} + \mu_s$ ,  $\mu_{ac} \ll \lambda$ ,  $\mu_s \perp \lambda$ .

**Theorem 6** (Radon-Nikodym theorem).  $\mu : \mathcal{B}(\mathbb{R}) \rightarrow [0, \infty]$  ( $\sigma$ -finite)

- There is a measurable map  $h : \mathbb{R} \rightarrow [0, \infty)$  with  $\mu_{ac} = \int_A d \lambda$  for all  $A \in \mathcal{B}(\mathbb{R})$ , where  $h$  is called the density function.

## 2.11 Image measure and substitution formula

Image measure is also called pushforward measure. Substitution formula is also called change of variable.



**Definition 2.11.1** (Image measure). Measure space  $(X, \mathcal{A})$ ,  $\mu$  is a measure on  $X$ . Measure space  $(Y, \mathcal{E})$ ,  $\tilde{\mu}$  is a measure on  $Y$ . Define a measure map  $h : X \rightarrow Y$ . See the above figure. We then define the image measure as

$$\tilde{\mu}(c) = \mu(h^{-1}(c)). \quad (2.11.1)$$

The notations:  $h_*\mu$  or  $\mu \circ h^{-1}$ .  $h_*\mu$  means pushforward and  $\mu \circ h^{-1}$  is readable. Remember that  $\tilde{\mu}$  is a measure on  $Y$ .

**Lemma 2.11.1** (Substitution formula). A integrable function  $g : Y \rightarrow \mathbb{R}$ . We have

$$\int_Y g \, d(h_*\mu) = \int_X g \circ h \, d\mu, \quad (2.11.2)$$

which can also be written as

$$\int_Y g(y) \, d(\mu \circ h^{-1})(y) = \int_X g(h(x)) \, d\mu(x), \quad (2.11.3)$$

which is called the change of variables:  $y = h(x)$ .

**Example 2.11.1.**  $F$  is a strictly monotonically increasing and continuously differentiable and surjective function from  $(\mathbb{R}, \mathcal{B}(\mathbb{R}))$  with  $\mu_F$  as  $\mu_F(A) = \int_A F'(x) \, dx$  to

$(\mathbb{R}, \mathcal{B}(\mathbb{R}))$ . We have

$$(F * \mu_F)([a, b]) = \mu_F(F^{-1}([a, b])) \quad (2.11.4)$$

$$= \mu_F([F^{-1}(a), F^{-1}(b)]) \quad (2.11.5)$$

$$= \int_{F^{-1}(a)}^{F^{-1}(b)} F'(x) \, dx \quad (2.11.6)$$

$$= \int_a^b dy \quad (2.11.7)$$

$$= \lambda([a, b]) \quad (2.11.8)$$

$$\implies F_x \mu_F = \lambda, \quad (2.11.9)$$

Substitution formula:

$$\int_Y g \, d(F * \mu_F) = \int_X g \circ F \, d\mu_F \quad (2.11.10)$$

$$\implies \int_{\mathbb{R}} g(y) \, dy = \int_{\mathbb{R}} g(F(x)) F'(x) \, dx. \quad (2.11.11)$$

*Proof.* (1) Let  $g = \chi_c$  with  $C \subseteq Y$  measurable. For the left hand side, we have

$$\int_Y \chi_c \, d(h * \mu) = (h * \mu)(c) \quad (2.11.12)$$

$$= \mu(h^{-1}(c)). \quad (2.11.13)$$

For the right hand side, we have

$$\int_X \chi_c \circ h \, d\mu = \int_X \chi_c \circ h \, d\mu \quad (2.11.14)$$

$$= \int_X \chi_c(h(x)) \, d\mu(x) \quad (2.11.15)$$

$$= \int_X \chi_{h^{-1}(c)} \, d\mu \quad (2.11.16)$$

$$= \mu(h^{-1}(c)), \quad (2.11.17)$$

where

$$\chi_c(h(x)) = \begin{cases} 1, & x \in h^{-1}(c) \\ 0, & x \notin h^{-1}(c) \end{cases} \quad (2.11.18)$$

(2) Let  $g$  be a simple function, i.e.,  $g = \sum_{i=1}^n \lambda_i \chi_{c_i}$ . We then obtain

$$\int_Y \sum_{i=1}^n \lambda_i \chi_{c_i} d(h * \mu) = \sum_{i=1}^n \lambda_i \int_Y \chi_{c_i} d(h * \mu) \quad (2.11.19)$$

By (1)

$$= \sum_{i=1}^n \lambda_i \int_X \chi_{c_i}(h(x)) d\mu(x) \quad (2.11.20)$$

$$= \int_X \left( \sum_{i=1}^n \lambda_i \chi_{c_i} \right)(h(x)) d\mu(x) \quad (2.11.21)$$

$$= \int_X g \circ h d\mu. \quad (2.11.22)$$

(3) Let  $g : Y \rightarrow [0, \infty)$  measurable. We have

$$\int_Y g d(h * \mu) = \sup \left\{ \int_Y \tilde{s} d(h * \mu) \mid \tilde{s} : Y \rightarrow [0, \infty) \text{ simple}, \tilde{s} \leq g \right\}. \quad (2.11.23)$$

We have the following equivalence relation:

$$\forall y \in h(x) : \tilde{s}(y) \leq g(y) \quad (2.11.24)$$

$$\iff \forall x \in X : \tilde{s}(h(x)) \leq g(h(x)) \quad (2.11.25)$$

$$[\text{i.e., } \tilde{s} \circ h \leq (g \circ h)(x)]. \quad (2.11.26)$$

Then we have

$$\int_Y g d(h * \mu) = \sup \left\{ \int_X \tilde{s} \circ d\mu \mid \tilde{s} : Y \rightarrow [0, \infty) \text{ simple}, \tilde{s} \circ h \leq g \circ h \right\} \quad (2.11.27)$$

Left as exercise

$$= \sup \left\{ \int_X s \circ d\mu \mid s : X \rightarrow [0, \infty) \text{ simple}, s \circ h \leq g \circ h \right\} \quad (2.11.28)$$

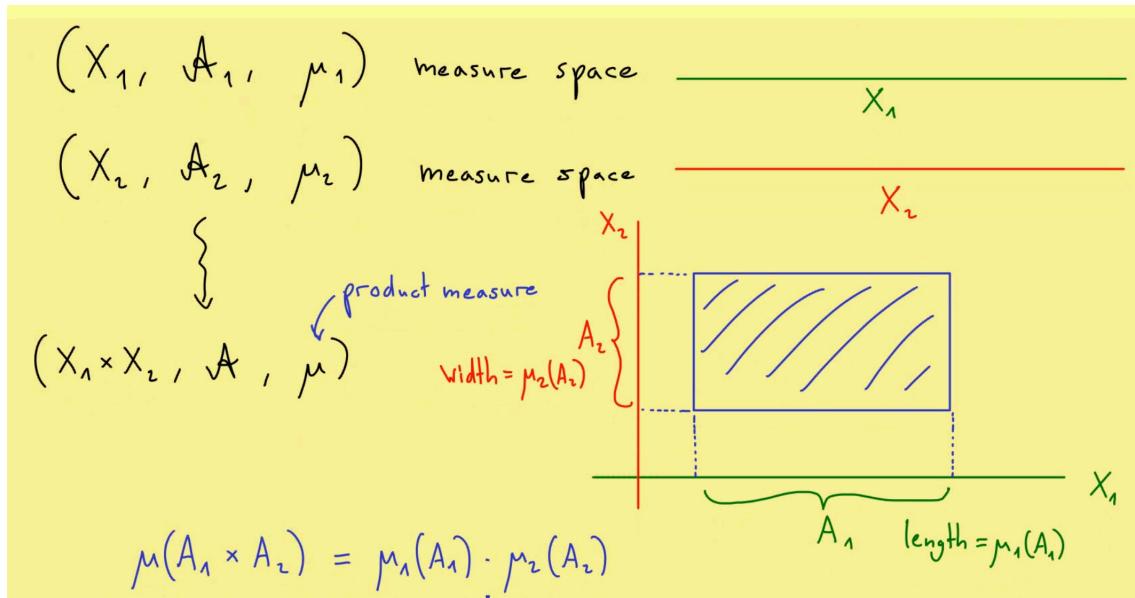
$$= \int_X g \circ h d\mu. \quad (2.11.29)$$

□

## 2.12 Product measure and Cavalieri's principle

$(X_1, \mathcal{A}_1, \mu_1)$  measure space and  $(X_2, \mathcal{A}_2, \mu_2)$  measure space,

$$\implies (X_1 \times X_2, \mathcal{A}, \mu), \text{ where } \mu \text{ is the product measure.} \quad (2.12.1)$$



We have

$$\mu(A_1 \times A_2) = \mu_1(A_1) \cdot \mu_2(A_2). \quad (2.12.2)$$

**Definition 2.12.1** (Product  $\sigma$ -algebra).

$$\mathcal{A} = \sigma(\mathcal{A}_1 \times \mathcal{A}_2). \quad (2.12.3)$$

**Remark 2.12.1.** Set of rectangles ( $= \mathcal{A}_1 \times \mathcal{A}_2$ ) are not a  $\sigma$ -algebra (but a semiring of sets)

**Definition 2.12.2.** Define product measure  $\mu$  as  $\mu(A_1 \times A_2) = \mu_1(A_1) \times \mu_2(A_2)$  for all  $A_1 \in \mathcal{A}_1$  and  $A_2 \in \mathcal{A}_2$ , and use (4).

**Remark 2.12.2.** Product measure in general not unique.

Proposition: If  $\mu_1, \mu_2$  are  $\sigma$ -finite, then there is exactly one measure  $\mu$  with  $\mu(A_1 \times A_2) = \mu_1(A_1) \cdot \mu_2(A_2)$ .

It satisfies:

$$\begin{aligned} \mu(M) &= \int_{X_2} \mu_1(M_y) d\mu_2(y) \\ &= \int_{X_1} \mu_2(M_x) d\mu_1(x) \end{aligned}$$

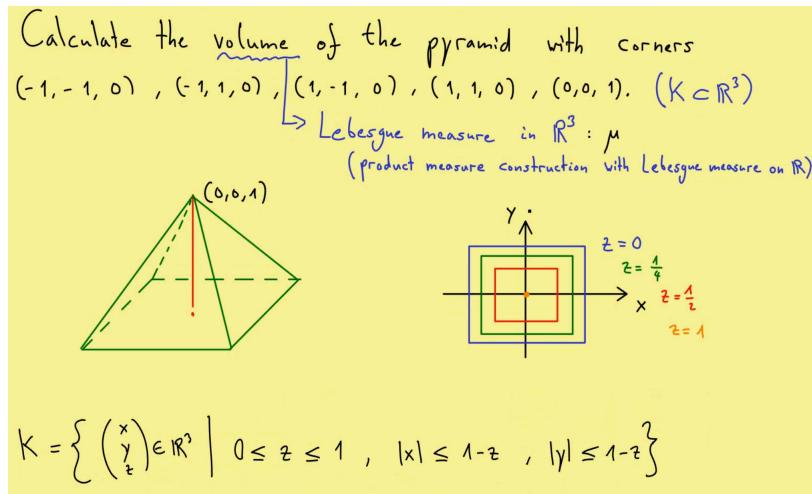
[Cavalieri's principle]

**Proposition 2.12.1** (Cavalieri's principle). *If  $\mu_1, \mu_2$  are  $\sigma$ -finite, then there is exactly one measure  $\mu$  with  $\mu(A_1 \times A_2) = \mu_1(A_1) \cdot \mu_2(A_2)$ . It satisfies:*

$$\mu(M) = \int_{X_2} \mu_1(M_y) d\mu_2(y) \quad (2.12.4)$$

$$= \int_{X_1} \mu_2(M_x) d\mu_1(x). \quad (2.12.5)$$

**Example 2.12.1** (An example for Cavalieri's principle). Calculate the volume of the pyramid with corners  $(-1, -1, 0), (-1, 1, 0), (1, -1, 0), (1, 1, 0), (0, 0, 1)$ ,  $K \subset \mathbb{R}^3$ , where the volume if the lebesgue measure in  $\mathbb{R}^3$ :  $\mu$  (Recall product measure construction with lebesgue measure on  $\mathbb{R}$ ).



*Proof.* Set

$$K = \{(x, y, z)^T \in \mathbb{R}^3 \mid 0 \leq z \leq 1, |x| \leq 1-z, |y| \leq 1-z\}. \quad (2.12.6)$$

Define  $\mu$  as a product measure of  $\mu_1$  and  $\mu_2$ , where  $\mu_1$  is the lebesgue measure in  $\mathbb{R}(z\text{-coordinate})$  and  $\mu_2$  is the lebesgue measure on  $\mathbb{R}^2(x\text{- and } y\text{-coordinate})$ . Following the definition of product measure, we have the volume of  $K$  as

$$\mu(k) = \int_{\mathbb{R}} \mu_2(M_{z_0}) d\mu_1(z_0) \quad (2.12.7)$$

$$= \int_{[0,1]} 4 \cdot (1-z_0)^2 d\mu_1(z_0) \quad (2.12.8)$$

$$= \frac{4}{3}, \quad (2.12.9)$$

where

$$M_{z_0} := \{(x, y)^T \in \mathbb{R}^2 \mid |x| \leq 1-z_0, |y| \leq 1-z_0\}, \quad (2.12.10)$$

and  $\mu_2(M_{z_0})$  is the area of the square only for  $z_0 \in [0, 1]$ .  $\square$

## 2.13 Fubini's theorem

**Theorem 7** (Fubini's theorem). *Let  $\mu_1$  and  $\mu_2$  be  $\sigma$ -finite,  $\mu$  be the product measure and*

$$f : X_1 \times X_2 \rightarrow [0, \infty] \text{ measurable [or } f \in \mathcal{L}^1(\mu)], \quad (2.13.1)$$

*then:*

$$\int_{X_1 \times X_2} f \, d\mu = \int_{X_2} \left( \int_{X_1} f(x, y) \, d\mu_1(x) \right) \, d\mu_2(y) \quad (2.13.2)$$

$$= \int_{X_1} \left( \int_{X_2} f(x, y) \, d\mu_2(y) \right) \, d\mu_1(x). \quad (2.13.3)$$

**Example 2.13.1.**  $\mu$  lebesgue measure for  $\mathbb{R}^2$ . Calculate  $\int_A f \, d\mu = ?$ , where

$$A = \{(x, y) \in [0, 1] \times [0, 1] \mid x \geq y \geq x^2\}, \quad (2.13.4)$$

$$f(x, y) = 2xy. \quad (2.13.5)$$

We have

$$\int_A f \, d\mu = \int_{\mathbb{R}^2} f \cdot \chi_A \, d\mu \quad (2.13.6)$$

$$= \int_{\mathbb{R}} \left( \int_{\mathbb{R}} f(x, y) \chi_A(x, y) \, dy \right) \, dx \quad (2.13.7)$$

$$= \int_0^1 \left( \int_x^{x^2} 2xy \, dy \right) \, dx \quad (2.13.8)$$

$$= \frac{1}{12}. \quad (2.13.9)$$

## 2.14 Outer measure

- tools for the proof of (4)
- "outer measure" is a new notion. "Outer measure" is not an attribute for "measure"! "Outer mesure" do not have to be measures!

**Definition 2.14.1** (Outer measure). *A map  $\phi : \mathcal{P}(X) \rightarrow [0, \infty]$  is called an outer measure if:*

- (a)  $\phi(\emptyset) = 0$
- (b)  $A \subseteq B \implies \phi(A) \leq \phi(B)$ . (monotonicity)
- (c)  $A_1, A_2, \dots, \in \mathcal{P}(X) \implies \phi(\cup_{n=1}^{\infty} A_n) \leq \sum_{n=1}^{\infty} \phi(A_n)$ . ( $\sigma$ -subadditivity)

**Question:**  $\phi : \mathcal{P}(X) \rightarrow [0, \infty]$  outer measure  $\xrightarrow{?} \mu$  measure?

**Definition 2.14.2** ( $\phi$ -measurable). Let  $\phi$  be an outer measure.  $A \in \mathcal{P}(X)$  is called  $\phi$ -measurable if for all  $Q \in \mathcal{P}(X)$  we have:

$$\phi(Q) \geq \phi(Q \cap A) + \phi(Q \cap A^c). \quad (2.14.1)$$

**Proposition 2.14.1.** If  $\phi : \mathcal{P}(X) \rightarrow [0, \infty]$  is an outer measure, then:

- $\mathcal{A}_\phi := \{A \subseteq X | A \text{ } \phi \text{-measurable}\}$  is a  $\sigma$ -algebra.
- $\mu : \mathcal{A}_\phi \rightarrow [0, \infty]$ ,  $\mu(A) := \phi(A)$ , is a measure.

union with two sets:  $A_1, A_2 \in \mathcal{A}_\phi$

$$\begin{aligned} \varphi(Q) &= \varphi(Q \cap A_1) + \varphi(Q \cap A_1^c) = \varphi(Q \cap A_1) + \varphi(\tilde{Q} \cap A_1) + \varphi(\tilde{Q} \cap A_1^c) \\ &\geq \varphi(Q \cap A_1) + \varphi(\tilde{Q} \cap A_2) + \varphi(\tilde{Q} \cap A_2^c) \end{aligned}$$

*Proof.* •  $\phi \in \mathcal{A}_\phi$ ? Is  $\emptyset$   $\phi$ -measurable?

$$\phi(Q) = \phi(Q \cap \emptyset) + \phi(Q \cup \emptyset^c) \quad (2.14.2)$$

$$= 0 + \phi(Q) \quad (2.14.3)$$

- $X \in \mathcal{A}_\phi$ ? Is  $X$   $\phi$ -measurable?

$$\phi(Q) = \phi(Q \cap X) + \phi(Q \cap X^c) \quad (2.14.4)$$

$$= \phi(Q) + \phi(\emptyset). \quad (2.14.5)$$

- $A \in \mathcal{A}_\phi \implies$

$$\phi(Q) = \phi(Q \cap A) + \phi(Q \cap A^c) \quad (2.14.6)$$

$$= \phi(Q \cap A^c) + \phi(Q \cap (A^c)^c) \quad (2.14.7)$$

$$\implies A^c \in \mathcal{A}_\phi. \quad (2.14.8)$$

- union with two sets:  $A_1, A_2 \in \mathcal{A}$

$$\phi(Q) = \phi(Q \cap A_1) + \phi(Q \cap A_1^c) \quad (2.14.9)$$

Define  $\tilde{Q} := Q \cap A_1^c$

$$= \phi(Q \cap A_1) + \phi(\tilde{Q} \cap A_2) + \phi(\tilde{Q} \cap A_2^c) \quad (2.14.10)$$

$$\geq \phi((Q \cap A_1) \cup (\tilde{Q} \cap A_2)) + \phi(\tilde{Q} \cap A_2^c) \quad (2.14.11)$$

$$= \phi(Q \cap (A_1 \cup A_2)) + \phi(Q \cap (A_1 \cup A_2)^c), \quad (2.14.12)$$

$$\implies \phi(Q) \geq \phi(Q \cap (A_1 \cup A_2)) + \phi(Q \cap (A_1 \cup A_2)^c) \quad (2.14.13)$$

$$\implies A_1 \cup A_2 \in \mathcal{A}_\phi, \quad (2.14.14)$$

where the fourth equation is obtain by the above figure.

- countable union:  $A_1, A_2, \dots \in \mathcal{A}_\phi$ ,  $A := \bigcup_{j=1}^{\infty} A_j \in \mathcal{A}_\phi$ ?

$$\phi(Q) = \phi(Q \cap A_1) + \phi(Q \cap A_1^c) \quad (2.14.15)$$

$$\begin{aligned} & \text{Set } Q = \hat{Q} \cap (A_1 \cup A_2) \\ & = \phi(\hat{Q} \cap A_1) + \phi(\hat{Q} \cap A_2). \end{aligned} \quad (2.14.16)$$

Induction:  $\phi(\hat{Q} \cap \bigcup_{j=1}^n A_j) = \sum_{j=1}^n \phi(\hat{Q} \cap A_j)$ . We have:

$$\phi(\hat{Q}) = \phi(\hat{Q} \cap \bigcup_{j=1}^n A_j) + \phi(\hat{Q} \cap (\bigcup_{j=1}^n A_j)^c) \quad (2.14.17)$$

$$\geq \sum_{j=1}^n \phi(\hat{Q} \cap A_j) + \phi(\hat{Q} \cap A^c) \quad (2.14.18)$$

$$\implies \phi(\hat{Q}) \geq \sum_{j=1}^n \phi(\hat{Q} \cap A_j) + \phi(\hat{Q} \cap A^c) \quad (2.14.19)$$

$$\geq \phi(\hat{Q} \cap A) + \phi(\hat{Q} \cap A^c) \quad (2.14.20)$$

$$\geq \phi(\hat{Q}) \quad (2.14.21)$$

$$\implies A \in \mathcal{A}_\phi. \quad (2.14.22)$$

□

**Example 2.14.1.** (1)  $\phi : \mathcal{P}(\mathbb{R}) \rightarrow [0, \infty]$ ,

$$\phi(A) = \begin{cases} 0, & A = \emptyset \\ 1, & A \neq \emptyset. \end{cases} \implies \text{outer measure but not a measure!} \quad (2.14.23)$$

**Example 2.14.2.**  $\phi : \mathcal{P}(\mathbb{N}) \rightarrow [0, \infty]$ ,

$$\phi(A) = \begin{cases} |A|, & A \text{ finite} \\ \infty, & A \text{ not finite.} \end{cases} \quad (2.14.24)$$

$$\implies \text{outer measure but a measure! (counting measure)} \quad (2.14.25)$$

$$\begin{aligned} (3) \quad \mathcal{X} &= \left\{ [a, b] \mid a, b \in \mathbb{R}, a \leq b \right\}, \quad \mu([a, b]) = b - a \quad ("length") \\ \text{Define} \quad \varphi : \mathcal{P}(\mathbb{R}) &\longrightarrow [0, \infty] \quad \text{by:} \quad \bigcup_{i=1}^{\infty} I_i \cup A \quad \text{A} \subseteq \bigcup_{j=1}^{\infty} I_j \\ \varphi(A) &:= \inf \left\{ \sum_{j=1}^{\infty} \mu(I_j) \mid I_j \in \mathcal{X}, A \subseteq \bigcup_{j=1}^{\infty} I_j \right\} \\ \rightsquigarrow \varphi &\text{ is an outer measure!} \end{aligned}$$

**Example 2.14.3.**  $\mathcal{I} = \{[a, b] | a, b \in \mathbb{R}, a \leq b\}$ ,  $\mu([a, b]) = b - a$  ("length").

Define  $\phi : \mathcal{P}(\mathbb{R}) \rightarrow [0, \infty)$  by:

$$\phi(A) := \inf \left\{ \sum_{j=1}^{\infty} \mu(I_j) \mid I_j \in \mathcal{I}, A \subseteq \bigcup_{j=1}^{\infty} I_j \right\} \quad (2.14.26)$$

$$\implies \phi \text{ is an outer measure!} \quad (2.14.27)$$

*Proof.* check (a) of (2.14.1):  $\phi(\emptyset) = 0$ .

check (b) of (2.14.1): monotonicity,

$$A \subseteq B \implies \phi(B) \quad (2.14.28)$$

$$= \inf \left\{ \sum_{j=1}^{\infty} \mu(I_j) \mid I_j \in \mathcal{I}, B \subseteq \bigcup_{j=1}^{\infty} I_j \right\} \quad (2.14.29)$$

$$\geq \inf \left\{ \sum_{j=1}^{\infty} \mu(I_j) \mid I_j \in \mathcal{I}, A \subseteq \bigcup_{j=1}^{\infty} I_j \right\}, \quad (2.14.30)$$

since  $A \subseteq B$ .

check (c) of (2.14.1): show that  $\phi(\bigcup_{n \in \mathbb{N}} A_n) \leq \sum_{n \in \mathbb{N}} \phi(A_n)$ . Let  $\varepsilon > 0$ . Choose  $\varepsilon_n > 0$  with  $\sum_{n \in \mathbb{N}} \varepsilon_n = \varepsilon$ . Then there are intervals  $I_{j,n}$  with:

$$\phi(A_n) \geq \sum_{j=1}^{\infty} \mu(I_{j,n}) - \varepsilon_n, \quad (2.14.31)$$

and

$$A_n \subseteq \bigcup_{j=1}^{\infty} I_{j,n}. \quad (2.14.32)$$

Then:  $\bigcup_{n \in \mathbb{N}} \subseteq \bigcup_{n \in \mathbb{N}} \bigcup_{j \in \mathbb{N}} I_{j,n} = \bigcup_{j \in \mathbb{N}} I_{j,n}$ .

$$\implies \phi(\bigcup_{n \in \mathbb{N}}) \stackrel{(b)}{\leq} \phi(\bigcup_{j \in \mathbb{N}} I_{j,n}) \quad (2.14.33)$$

$$\leq \sum_{j \in \mathbb{N}} \mu(I_{j,n}) \quad (2.14.34)$$

$$= \sum_{n \in \mathbb{N}} \left\{ \sum_{j \in \mathbb{N}} \mu(I_{j,n}) \right\} \quad (2.14.35)$$

$$\leq \sum_{n \in \mathbb{N}} (\phi(A_n) + \varepsilon_n) \quad (2.14.36)$$

$$= \sum_{n \in \mathbb{N}} \phi(A_n) + \varepsilon. \quad (2.14.37)$$

□

# Chapter 3

## Linear algebra

### 3.1 Exercises

- 3.C.6: Are the construced  $v_i$  form a basis of vector space  $V$ ?
-

## Chapter 4

# Functional Analysis

### 4.1 Metric Space

**Definition 4.1.1** (Metrix Spaces). Define set  $X$ . Define a metric:  $d : X \times X \rightarrow [0, \infty)$

- $d(x, y) = 0 \iff x = y$
- $d(x, y) = d(y, x)$
- $d(x, y) \leq d(x, z) + d(z, y)$  (triangle inequality)

Summary:  $X$  set +  $d : X \times X \rightarrow [0, \infty)$  metric = metric space  $(X, d)$ .

**Example 4.1.1.** •  $X = \mathbb{C}, d(x, y) = |x - y|$

- $X = \mathbb{R}^n, d(x, y) = \sqrt{(x_1 - y_1)^2 + (x_2 - y_2)^2 + \dots + (x_n - y_n)^2}$  (Euclidean metric)
- $X = \mathbb{R}^n, d(x, y) = \max \{|x_1 - y_1|, \dots, |x_n - y_n|\}$ . It is possible that  $d(x, y) = d(x, z)$ .
- $X$  any set ( $\neq \emptyset$ ), we define

$$d(x, y) = \begin{cases} 0, & x = y \\ 1, & x \neq y \end{cases} \quad (4.1.1)$$

Now, let's check whether this is a metrix space. Conditions (1) and (2) are verified easily. We focus on the verification of (3): choose  $x, y, z \in X$ .

For the first case:  $x = y$ :  $d(x, y) = 0 \leq d(x, z) + d(z, y)$ .

For the second case:  $x \neq y$ :  $d(x, y) = 1 \leq \{d(x, z) \text{ or } d(z, y)\} = d(x, z) + d(z, y)$ .

This is called **discrete metric**.

## 4.2 Open and Closed Sets

**Definition 4.2.1** (Open Ball).  $(X, d)$  a metrix space. Define

$$B_\epsilon(x) := \{y \in X \mid d(x, y) < \epsilon\}, \quad (4.2.1)$$

whihc is called the open ball of radieue  $\epsilon > 0$  centered at  $x$ .

**Definition 4.2.2** (Open Sets).  $A \subseteq X$  is called open if for each  $x \in A$  there is an open ball with  $B_\epsilon(x) \subseteq A$ .

We have nice pic for this.

**Definition 4.2.3** (Boundary Points).  $A \subseteq X$ ,  $x \in X$  is called a boundary point for  $A$  if for all  $\epsilon > 0$ :  $B_\epsilon(x) \cap A = \emptyset$  and  $B_\epsilon(x) \cap A^c = \emptyset$ .  $[A^c := X \setminus A]$

Notation:  $\partial A : \{x \in X \mid x \text{ is boundary point for } A\}$

We have nice pic for this.

Remember:  $A$  open  $\iff A \cap \partial A = \emptyset$ .

**Definition 4.2.4** (Closed Sets).  $A \subseteq X$  is called closed if  $A^c := X \setminus A$  is open.

**Definition 4.2.5** (Closure).  $\overline{A} := A \cup \partial A$  (always closed!)

**Example 4.2.1.**  $X := (1, 3] \cup (4, \infty)$ ,  $d(x, y) := |x - y|$ ,  $(X, d)$  is a metrix space.

- $A := (1, 3] \subseteq X$  open?

For  $x \in A$ ,  $x \neq 3$ , define  $\epsilon := \frac{1}{2} \min(|1 - x|, |3 - x|)$ . Then  $B_\epsilon(x) \subseteq A$ .

For  $x = 3$ :  $B_1(x) = \{y \in X \mid d(x, y) < 1\} = (2, 3] \subseteq A$ .

- $A$  is also closed!

- $C := (1, 2]$ ,  $\partial C = \{2\}$ ,  $\overline{C} = C$ .

## 4.3 Sequence, Limits and Closed Sets

**Definition 4.3.1** (Sequence). Sequence in  $X$ :  $x_1, x_2, \dots$  or  $(x_n)_{n \in \mathbb{N}}$  or map  $x : \mathbb{N} \rightarrow X$  /  $n \mapsto x_n$ .

**Definition 4.3.2** (Convergence). A sequence  $(x_n)_{n \in \mathbb{N}}$  in a metrix space  $(X, d)$  is called convergence if there is  $\tilde{x} \in X$  with  $\forall \varepsilon \geq 0, \exists N \in \mathbb{N}, \forall n \geq N : d(x_n, \tilde{x}) < \varepsilon$ . We write:  $X_n \xrightarrow{n \rightarrow \infty} \tilde{x}$  or  $\lim_{n \rightarrow \infty} x_n = \tilde{x}$ .

**Proposition 4.3.1** (Closed Sets).  $A \subseteq X$  is closed  $\iff$  For every convergent sequence  $(a_n)_{n \in \mathbb{N}} \subseteq A$ , one has  $\lim_{n \rightarrow \infty} a_n \in A$ .

*Proof.* ( $\Leftarrow$ ): Show it by contraposition! Assume  $A$  is not closed.

$$\implies A^c := X \setminus A \text{ is not open.} \quad (4.3.1)$$

$$\implies \text{There is an } \tilde{x} \in A^c \text{ with } B_\varepsilon(\tilde{x}) \cap A \neq \emptyset, \forall \varepsilon > 0. \quad (4.3.2)$$

$$\implies \text{There is a sequence } (a_n)_{n \in \mathbb{N}} \text{ with } a_n \in B_{1/n}(\tilde{x}) \cap A \quad (4.3.3)$$

$$\implies \lim_{n \rightarrow \infty} a_n = \tilde{x} \neq A. \quad (4.3.4)$$

( $\Rightarrow$ ): Show it by contraposition! Assume there is  $(a_n)_{n \in \mathbb{N}} \subseteq A$  with  $\tilde{x} := \lim_{n \rightarrow \infty} a_n \notin A$ .

$$\implies B_\varepsilon(\tilde{x}) \neq \emptyset, \forall \varepsilon > 0 \quad (4.3.5)$$

$$\implies A^c \text{ is not open} \quad (4.3.6)$$

$$\implies A \text{ is not closed.} \quad (4.3.7)$$

□

[SZQ: This proof uses many contrapositions. It is nice to get familiar with how to use contrapositions to prove. It also needs to know the inverse of  $\forall$ .]

## 4.4 Cauchy Sequence and complete metric spaces

**Example 4.4.1.**  $X = (0, 3)$  with  $d(x, y) = |x - y|$ .

$(0, 3)$  is closed:

- complement  $\emptyset$  is open
- each convergent sequence  $(x_n)_{n \in \mathbb{N}} \subseteq (0, 3)$  (with limit  $\tilde{x} \in X$ ) satisfies  $\tilde{x} \in (0, 3)$

What is about the sequence  $(\frac{1}{n})_{n \in \mathbb{N}}$ ?

- sequence in  $X$
- $d(x_n, x_m) \xrightarrow{n, m \rightarrow \infty} 0$
- it does not converge  $\implies (X, d)$  is not complete

**Definition 4.4.1** (Cauchy Sequence and Complete). Let  $(X, d)$  be a metrix space. A sequence  $(x_n)_{n \in \mathbb{N}} \subseteq X$  is called Cauchy sequence if  $\forall \varepsilon > 0, \exists N \in \mathbb{N}, \forall n, m \geq N : d(x_n, x_m) < \varepsilon$ .  $(X, d)$  is called complete if all Cauchy sequences converge.

**Example 4.4.2.** •  $X = [0, 3]$  with  $d(x, y) = |x - y|$  is complete.

- $X = (0, 3)$  with

$$\begin{cases} 1, & x \neq y \\ 0, & x = y \end{cases} \quad (4.4.1)$$

is complete.

*Proof.* Let  $(x_n)_{n \in \mathbb{N}} \subseteq X$  be a Cauchy sequence. Take  $\varepsilon = 1/2$ . Then there is an  $N \in \mathbb{N}$  such that for all  $n, m \geq N$ , we have  $d(x_n, x_m) < \frac{1}{2}$ . By definition,  $d(x_n, x_m) = 0$ . Hence,  $x_n = x_m$ .  $\square$

## 4.5 Norms and Banach Spaces

**Definition 4.5.1** (Norm).  $\mathbb{F} \in \{\mathbb{R}, \mathbb{C}\}$ . Let  $X$  be a  $\mathbb{F}$ -vector space. A map  $\|\cdot\| : X \rightarrow [0, \infty)$  is called norm if

- $\|x\| = 0 \iff x = 0$  (positive definite)
- $\|\lambda \cdot x\| = |\lambda| \|x\|$  for all  $\lambda \in \mathbb{F}$ ,  $x \in X$  (absolutely homogeneous)
- $\|x\| \leq \|x\| + \|y\|$  for all  $x, y \in X$  (triangle inequality)

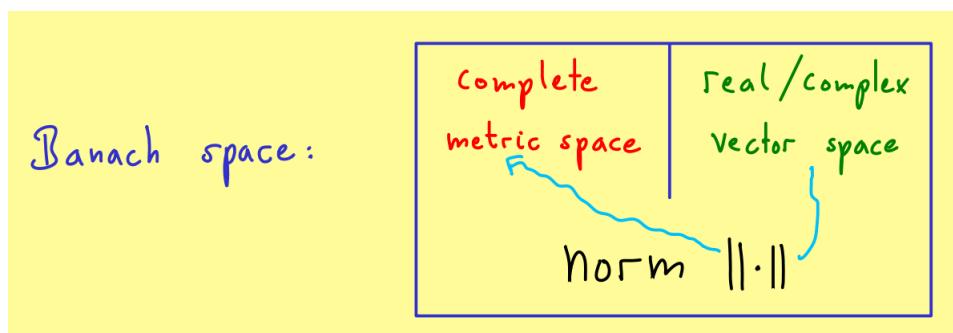
**Definition 4.5.2** (Normed Space).  $(X, \|\cdot\|)$  is then called a normed space.

Normed space is a special case of metric space!

**Lemma 4.5.1** (Relationship between normed space and metric space). If  $\|\cdot\|$  is a norm for the  $\mathbb{F}$ -vector space  $X$ , then  $d_{\|\cdot\|}(x, y) := \|x - y\|$  defines a metric for the set  $X$ .

*Proof.* This can be proved by the definition of norm.  $\square$

**Definition 4.5.3** (Banach space). If  $(X, d_{\|\cdot\|})$  is a complete metric space, then the normed space  $(X, \|\cdot\|)$  is called a Banach space.



The figure is very impressive and informative!

## 4.6 Examples of Banach Spaces

### 4.7 Part 12: Continuity

**Definition 4.7.1** (Continuity for metric spaces).  $(X, d_X), (Y, d_Y)$  are two metric spaces. A map  $f : X \rightarrow Y$  is called:

- continuous if  $f^{-1}[B]$  is open in  $X$  for all open sets  $B \subseteq Y$ .
- sequentially continuous if for all  $\tilde{x} \in X$  and  $(x_n)_{n \in \mathbb{N}} \subseteq X$  with  $x_n \xrightarrow{n \rightarrow \infty} \tilde{x}$  holds  $f(x_n) \xrightarrow{n \rightarrow \infty} f(\tilde{x})$ .

These are shown by figs.

**Lemma 4.7.1.** For metric spaces, continuous and sequentially continuous are equivalent. But for topological spaces, they are different.

**Example 4.7.1.** •  $(X, d_X)$  discrete metric space,  $(Y, d_Y)$  any metric space  $\Rightarrow$  all  $f : X \rightarrow Y$  are continuous.

- $(X, d_X), (Y, d_Y)$  metric spaces,  $y_0 \in Y$  fixed.  $\Rightarrow f : X \rightarrow Y, x \mapsto y_0$  is always continuous.
- $(X, \|\cdot\|)$  normed space,  $Y = \mathbb{R}$  with standard metric.  $\Rightarrow f : X \rightarrow \mathbb{R}, x \mapsto \|x\|$  is continuous.

*Proof.* Let  $(x_n)_{n \in \mathbb{N}} \subseteq X$  sequence with limit  $\tilde{x} \in X$ . Then:

$$f(x_n) = \|x_n\| \tag{4.7.1}$$

$$= \|x_n - \tilde{x} + \tilde{x}\| \tag{4.7.2}$$

By triangle inequality 4.5.1

$$\leq \|x_n - \tilde{x}\| + \|\tilde{x}\| \tag{4.7.3}$$

$$= d(x_n, \tilde{x}) + f(\tilde{x}) \tag{4.7.4}$$

$$\Rightarrow \lim_{n \rightarrow \infty} f(x_n) \leq f(\tilde{x}). \tag{4.7.5}$$

We also hold:

$$f(\tilde{x}) = \|\tilde{x}\| \quad (4.7.6)$$

$$= \|\tilde{x} - x_n + x_n\| \quad (4.7.7)$$

By triangle inequality 4.5.1

$$\|\tilde{x} - x_n\| + \|x_n\| \quad (4.7.8)$$

$$= d(\tilde{x}, x_n) + f(x_n) \quad (4.7.9)$$

$$\implies f(\tilde{x}) \leq \lim_{n \rightarrow \infty} f(x_n). \quad (4.7.10)$$

□

- $(X, \langle \cdot, \cdot \rangle)$  inner product space,  $Y \in \mathbb{C}$  with the standard metric,  $x_0 \in X$  fixed.  $\implies f : X \rightarrow \mathbb{C}$ ,  $x \mapsto \langle x_0, x \rangle$  is continuous.

*Proof.* Let  $(x_n)_{n \in \mathbb{N}} \subseteq X$  sequence with limit  $\tilde{x} \in X$ . Then:

$$|f(x_n) - f(\tilde{x})| = |\langle x_0, x_n \rangle - \langle x_0, \tilde{x} \rangle| \quad (4.7.11)$$

$$= |\langle x_0, x_n - \tilde{x} \rangle| \quad (4.7.12)$$

By Cauchy Schiwz inequality

$$\leq \|x_0\| \cdot \|x_n - \tilde{x}\| \xrightarrow{n \rightarrow \infty} 0. \quad (4.7.13)$$

Analogously,  $g : X \rightarrow \mathbb{C}$ ,  $x \mapsto \langle x, x_0 \rangle$  is continuous. □

**Lemma 4.7.2** (Orthogonal complement is closed).  $(X, \langle \cdot, \cdot \rangle)$  inner product space,  $U \subseteq X$ . Then  $U^\perp$  is closed.

*Proof.* Let  $(x_n)_{n \in \mathbb{N}} \subseteq U^\perp$  with limit  $\tilde{x} \in X$ .

$$\implies \langle x_n, u \rangle = 0, \forall u \in U \quad (4.7.14)$$

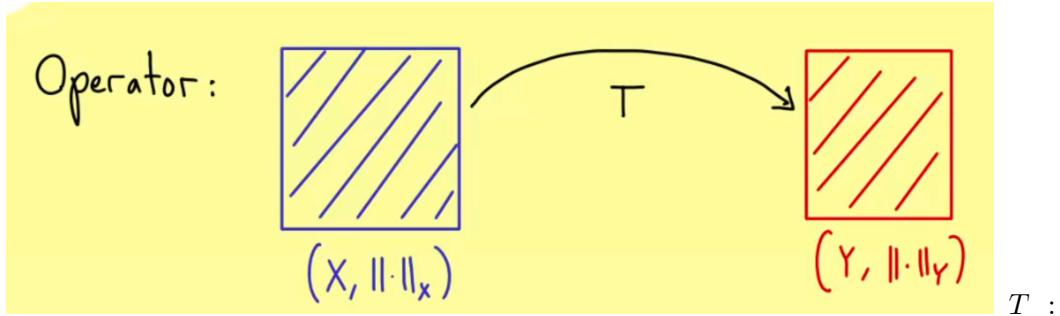
$$\implies \lim_{n \rightarrow \infty} \langle x_n, u \rangle = 0, \forall u \in U \quad (4.7.15)$$

$$\implies \langle \tilde{x}, u \rangle = 0, \forall u \in U. \quad (4.7.16)$$

$$\implies \tilde{x} \in U^\perp. \quad (4.7.17)$$

□

## 4.8 Part 13: Bounded Operators



$X \rightarrow Y$ , which satisfies

- linear (conserves the algebraic structure)
- continuous (bounded) (conserves the topological structure)

**Definition 4.8.1** (Operator norm and bounded).  $(X, \|\cdot\|_X)$ ,  $(Y, \|\cdot\|)$  two normed spaces,  $T : X \rightarrow Y$  linear, which means

$$T(x + \tilde{x}) = Tx + T\tilde{x} \quad (4.8.1)$$

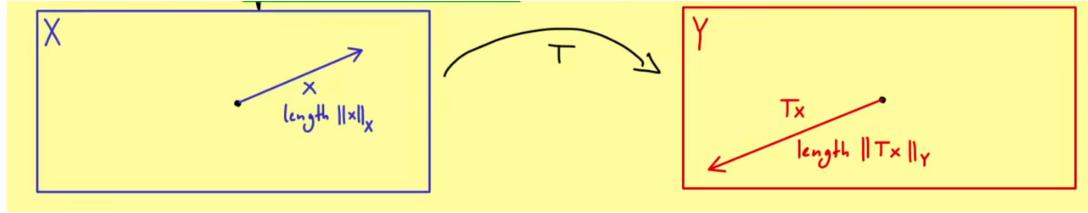
$$T(\lambda x) = \lambda Tx \quad (4.8.2)$$

for all  $x, \tilde{x} \in X$ ,  $\lambda \in \mathbb{F}$ . Then

$$\|T\| = \|T\|_{X \rightarrow Y} \quad (4.8.3)$$

$$:= \sup \left\{ \frac{\|Tx\|_Y}{\|x\|_X} \mid x \in X, x \neq 0 \right\} \quad (4.8.4)$$

is called the operator norm of  $T$ . If  $\|T\| < \infty$ ,  $T$  is called bounded.



**Proposition 4.8.1** (Continuous equivalent to bounded). Let  $(X, \|\cdot\|_X)$ ,  $(Y, \|\cdot\|_Y)$  two normed spaces,  $T : X \rightarrow Y$  linear. Then the following claims are equivalent:

- $T$  is continuous.
- $T$  is continuous at  $x = 0$ .
- $T$  is bounded.

*Proof.* (a)  $\implies$  (b) is easily seen.

(b)  $\implies$  (c): proposition (\*): For all sequence  $(x_n)_{n \in \mathbb{N}} \subseteq X$  with  $x_n \xrightarrow{n \rightarrow \infty} 0$ , we have  $Tx_n \xrightarrow{n \rightarrow \infty} 0$  due to the properties of linear map.

Claim: proposition (\*)  $\implies$  proposition (+): there is a  $\delta > 0$  such that  $\|Tx\|_Y < 1$  for all  $x \in X$  with  $\|x\|_X < \delta$ .

proof of this claim: we prove its contraposition such that  $\neg(*) \implies$  For all  $n \in \mathbb{N}$ , we find  $x_n \in X$  with  $\|x_n\|_X < \frac{1}{n}$  and  $\|Tx_n\|_Y \geq 1 \implies \neg(*)$ .

$$\frac{\|Tx\|_Y}{\|x\|_X} = \frac{\|Tx\|_Y \cdot \frac{\delta}{2} \cdot \frac{1}{\|x\|_X}}{\|x\|_X \cdot \frac{\delta}{2} \cdot \frac{1}{\|x\|_X}} \quad (4.8.5)$$

$$= \frac{\|T(\frac{\delta}{2} \frac{x}{\|x\|_X})\|_Y}{\|\frac{\delta}{2} \frac{x}{\|x\|_X}\|_X} \quad (4.8.6)$$

$$\leq \frac{2}{\delta} \quad (4.8.7)$$

$$\implies \|T\| = \sup \left\{ \frac{\|Tx\|_Y}{\|x\|_X} \mid x \in X, x \neq 0 \right\} \leq \frac{2}{\delta} < \infty. \quad (4.8.8)$$

(c)  $\implies$  (a): Let  $(x_n)_{n \in \mathbb{N}} \subseteq X$  be convergent with limit  $\tilde{x} \in X$ . Then  $\|Tx_n - T\tilde{x}\|_Y = \|T(x_n - \tilde{x})\|_Y \leq \|T\| \cdot \|x_n - \tilde{x}\|_X \xrightarrow{n \rightarrow \infty} 0$ .  $\square$

[SZQ: 2023.04.02: This proof needs to be understood later.]

[SZQ: 2023.04.02: It should be proved that  $\|T\|$  defined above is indeed a norm.]

## 4.9 Part 14: Example Operator Norm

**Example 4.9.1.**  $X = (C[0, 1], \mathbb{F}, \|\cdot\|_\infty)$ ,  $Y = (\mathbb{F}, |\cdot|)$ . For  $g \in X$  with  $g(t) \neq 0$  for all  $t \in [0, 1]$ , define  $T_g : X \rightarrow Y$  by  $T_g(f) := \int_0^1 g(t) \cdot f(t) dt$ . So what is  $\|T_g\|$ ?

*Proof.* Recall that

$$\|F_g\| = \sup \left\{ \frac{|T_g(f)|}{\|f\|_\infty} \mid f \in X, f \neq 0 \right\} \quad (4.9.1)$$

This trick has been used before.

$$= \sup \left\{ \frac{|T_g(f)|}{\|f\|_\infty} \mid f \in X, f \neq 0 \right\} \quad (4.9.2)$$

$$= \sup \{ |T_g(f)| \mid f \in X, \|f\|_\infty = 1 \} \quad (4.9.3)$$

$$= \sup \left\{ \left| \int_0^1 g(t) \cdot f(t) dt \right| \mid f \in X, \|f\|_\infty = 1 \right\} \quad (4.9.4)$$

Since  $\left| \int_0^1 g(t) \cdot f(t) dt \right| \leq \int_0^1 |g(t)| \cdot |f(t)| dt$  and  $|f(t)| \leq \|f\|_\infty = 1$

$$\leq \int_0^1 |g(t)| dt \quad (4.9.5)$$

$$< \infty. \quad (4.9.6)$$

Check the other inequality:  $h(t) := \frac{\overline{g(t)}}{|g(t)|}$  with  $\|h\|_\infty = 1$ . We then have

$$\|T_g\| \geq |T_g(h)| \quad (4.9.7)$$

$$= \left| \int_0^1 g(t) \frac{\overline{g(t)}}{|g(t)|} dt \right| \quad (4.9.8)$$

$$= \int_0^1 \frac{|g(t)|^2}{|g(t)|} dt \quad (4.9.9)$$

$$= \int_0^1 |g(t)| dt. \quad (4.9.10)$$

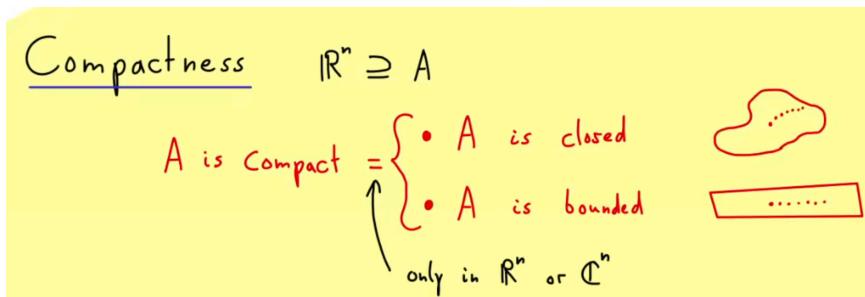
□

## 4.10 Part 16: Compact Sets

**Example 4.10.1.** Compactness  $A \subseteq \mathbb{R}^n$ .  $A$  is compact such that

- $A$  is closed.
- $A$  is bounded.

This is only true in  $\mathbb{R}^n$  or  $\mathbb{C}^n$ .



**Definition 4.10.1** (Sequentially compact). Let  $(X, d)$  be a metric space.  $A \subseteq X$  is called (sequentially) compact if for each sequence  $(x_n)_{n \in \mathbb{N}} \subseteq A$ , one finds a convergent subsequence  $(x_{n_k})_{k \in \mathbb{N}}$  with  $\tilde{x} := \lim_{k \rightarrow \infty} x_{n_k} \in A$ .

Definition: Let  $(X, d)$  be a metric space.  $A \subseteq X$  is called (sequentially) compact if for each sequence  $(x_n)_{n \in \mathbb{N}} \subseteq A$  one finds a convergent subsequence  $(x_{n_k})_{k \in \mathbb{N}}$  with  $\tilde{x} := \lim_{k \rightarrow \infty} x_{n_k} \in A$

**Example 4.10.2.** •  $(\mathbb{R}, d_{eucl.}), A = [0, 1]$  compact by Bolzano-Weierstrass theorem.

- $(\mathbb{R}, d_{discr}), A = [0, 1]$  not compact because: The sequence  $(x_n)_{n \in \mathbb{N}} \subseteq A$  with  $x_n = \frac{1}{n}$  satisfies  $d_{discr.}(x_n, x_m) = 1$  for all  $n, m \in \mathbb{N}$  with  $n \neq m$ .  $\Rightarrow$  no convergent subsequence.

Proposition: Let  $(X, d)$  be a metric space and  $A \subseteq X$  compact. Then  $A$  is closed and bounded. There is an  $x \in X$  and an  $\varepsilon > 0$  such that  $B_\varepsilon(x) \supseteq A$

**Definition 4.10.2** (Bounded). Let  $(X, d)$  be a metrix space and  $A \subseteq X$  compact.  $A$  is bounded means that there is an  $x \in X$  and an  $\varepsilon > 0$  such that  $A \subseteq B_\varepsilon(x)$ .

**Proposition 4.10.1** (Compact implies closed and bounded). Let  $(X, d)$  be a metrix space and  $A \subseteq X$  compact. Then  $A$  is closed and bounded.

*Proof.* Let  $A \subseteq X$  be compact.

(1) Let  $(x_n)_{n \in \mathbb{N}} \subseteq A$  be convergent with limit  $\tilde{x} \in X$ .

$$\Rightarrow \text{There is a convergent subsequence } (x_{n_k})_{k \in \mathbb{N}} \text{ with limit } \tilde{x} \in A \quad (4.10.1)$$

$$\Rightarrow \tilde{x} = \tilde{x} \in A \quad (4.10.2)$$

$$\Rightarrow A \text{ is closed!} \quad (4.10.3)$$

(2) contraposition:  $A$  is not bounded.

$\Rightarrow$  For given  $a \in A$ , there are  $x_n \in A$  with  $d(a, x_n) > n$ .

$\Rightarrow$  For any subsequence  $(x_{n_k})_{k \in \mathbb{N}}$  and any point  $b \in A$ :

$$n_K < d(a, x_{n_k}) \quad (4.10.4)$$

$$\leq d(a, b) + d(b, x_{n_k}) \quad (4.10.5)$$

$$\begin{aligned} &\implies n_k - d(a, b) \leq d(b, x_{n_k}). \\ &\implies d(b, x_{n_k}) \xrightarrow{k \rightarrow \infty} 0 \text{ for all } b \in A \\ &\implies A \text{ not compact!} \end{aligned}$$

□

## 4.11 Part 18: Compact operators

**Definition 4.11.1** (Compact operators). *Let  $(X, \|\cdot\|_X)$ ,  $(Y, \|\cdot\|_Y)$  be two normed spaces. A bounded linear operator  $T : X \rightarrow Y$  is called compact if  $\overline{T[B_1(0)]} \subseteq Y$  is a compact set.*

[SZQ: 2023.04.02: the definition of  $\overline{T[B_1(0)]} \subseteq Y$  should be learned first.]

## 4.12 Part 19: Holder's inequality

**Lemma 4.12.1** (Young's inequality). *For all  $a, b > 0$ , we have  $a, b \leq \frac{a^p}{p} + \frac{b^{p'}}{p'}$ .*

*Proof.* Note that function  $f : x \mapsto e^x$  is convex, we have

$$f(\lambda x + (1 - \lambda)y) \leq \lambda f(x) + (1 - \lambda)f(y). \quad (4.12.1)$$

Let  $\lambda = \frac{1}{p}$ ,  $x = \ln a^p$ ,  $1 - \lambda = \frac{1}{p'}$ , and  $y = \ln b^{p'}$ . We have

$$a \cdot b = f\left(\frac{1}{p} \ln a^p + \frac{1}{p'} \ln b^{p'}\right) = f(\lambda x + (1 - \lambda)y) \leq \lambda f(x) + (1 - \lambda)f(y) \quad (4.12.2)$$

$$= \frac{1}{p} f(\ln a^p) + \frac{1}{p'} f(\ln b^{p'}) \quad (4.12.3)$$

$$= \frac{a^p}{p} + \frac{b^{p'}}{p'}. \quad (4.12.4)$$

□

**Theorem 8** (Holder's inequality). *For all  $x, y \in \mathbb{F}^n$ , we have*

$$\|xy\|_1 \leq \|x\|_p \cdot \|y\|_{p'}, \quad (4.12.5)$$

where  $x \in \mathbb{F}^n$  and the  $p$ -norm of  $x$  is

$$\|x\|_q := \left( \sum_{j=1}^n |x_j|^q \right)^{\frac{1}{q}}, \quad (4.12.6)$$

$q \in [1, \infty)$ ,  $\frac{1}{p} + \frac{1}{p'} = 1$ , and for  $x, y \in \mathbb{F}^n$ , we write

$$xy := \begin{pmatrix} x_1 y_1 \\ x_2 y_2 \\ \dots \\ x_n y_n \end{pmatrix}. \quad (4.12.7)$$

*Proof.* Case 1:  $x = 0$  or  $y = 0$ .

Case 2: We have

$$\frac{1}{\|x\|_p \|y\|_{p'}} \|xy\|_1 = \frac{1}{\|x\|_p \|y\|_{p'}} \sum_{j=1}^n |x_j y_j| \quad (4.12.8)$$

$$= \sum_{j=1}^n \frac{|x_j|}{\|x\|_p} \frac{|y_j|}{\|y\|_{p'}} \quad (4.12.9)$$

By Young's lemma (4.12.1)

$$\leq \sum_{j=1}^n \frac{1}{p} \cdot \frac{|x_j|^p}{\|x\|_p^p} + \sum_{j=1}^n \frac{1}{p'} \cdot \frac{|y_j|^{p'}}{\|y\|_{p'}^{p'}} \quad (4.12.10)$$

$$= \frac{1}{p} + \frac{1}{p'} \quad (4.12.11)$$

$$= 1. \quad (4.12.12)$$

□

## 4.13 Part 21: Isomorphism?

**Definition 4.13.1** (Homomorphism). *What is Homomorphism: map that preserves structures.*

[SZQ: When we talk about homomorphism, we must know the underlying structures!]

**Example 4.13.1.** • Let  $X, Y$  be vector spaces and  $f : X \rightarrow Y$  be a map. We want

$$f(\lambda \cdot x) = \lambda \cdot f(x) \quad (4.13.1)$$

$$f(x + x') = f(x) + f(x'), \quad (4.13.2)$$

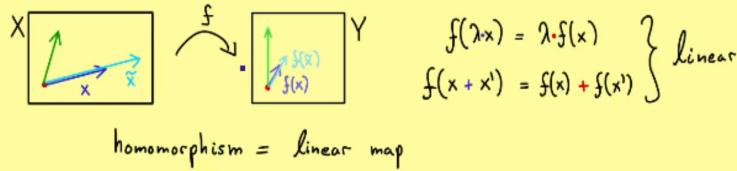
which is called linear. Thus homomorphism = linear map!

• Let  $(X, d_X), (Y, d_Y)$  be two metric spaces and  $f : X \rightarrow Y$  be a map. We want

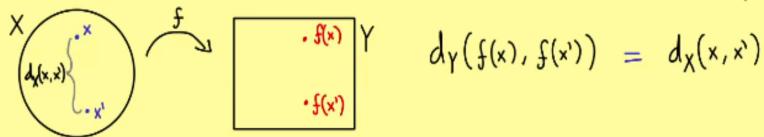
$$d_Y(f(x), f(x')) \leq d_X(x, x'). \quad (4.13.3)$$

homomorphism = map that satisfies (4.13.3)

Example: (a) Let  $X, Y$  be vector spaces and  $f: X \rightarrow Y$  be a map.



(b) Let  $(X, d_X), (Y, d_Y)$  be two metric spaces and  $f: X \rightarrow Y$  be a map.



**Definition 4.13.2** (Isomorphism). *Isomorphism = homomorphism + bijective + inverse map is also homomorphism*

**Definition 4.13.3** (Isomorphism for Banach spaces). *Isomorphism for banach spaces  $X, Y$ :  $f : X \rightarrow Y$  with: linear + bijective +  $\|f(x)\|_Y = \|x\|_X$  (oftern called isometric isomorphism).*

**Example 4.13.2.** •  $S_R : l^p(\mathbb{N}) \rightarrow l^p(\mathbb{N})$ ,  $(x_1, x_2, x_3, \dots) \mapsto (0, x_1, x_2, \dots)$

$$\implies \text{linear}, \|S_R x\|_p = \|x\|_p \text{ not surjective} \quad (4.13.4)$$

$$\implies \text{not an isomorphism} \quad (4.13.5)$$

•  $S : l^p(\mathbb{Z}) \rightarrow l^p(\mathbb{Z})$ ,  $(\dots, x_{-1}, x_0, x_1, \dots) \mapsto (\dots, x_{-2}, x_{-1}, x_0, \dots)$

$$\implies \text{linear}, \|Sx\|_p = \|x\|_p \text{ and bijective} \quad (4.13.6)$$

$$\implies \text{isomorphism} \quad (4.13.7)$$

## 4.14 Part 22: Dual spaces

**Proposition 4.14.1.** *Let  $X$  be a normed space. Then  $(X', \|\cdot\|_{X \rightarrow \mathbb{F}})$  is a Banach space.*

*Proof.*

□

[SZQ: 2023.04.02: I need to learn Riesz representation theorem!]

## 4.15 Part 28: Spectrum for bounded linear operators

Recall:  $A \in \mathbb{C}^{n \times n}$  matrix with  $n$  rows and  $n$  columns.  $\lambda \in \mathbb{C}$  is called an eigenvalue of  $A$  if:

$$\exists x \in \mathbb{C}^n \setminus \{0\} : Ax = \lambda x \quad (4.15.1)$$

$$\iff \exists x \in \mathbb{C}^n \setminus \{0\} : (A - \lambda I)x = 0 \quad (4.15.2)$$

$$\iff \text{Ker}(A - \lambda I) \neq \{0\} \quad (4.15.3)$$

$$\iff \text{map } x \mapsto (A - \lambda I)x \text{ not injective}. \quad (4.15.4)$$

**Theorem 9** (Rank-nullity theorem). *For any matrix  $M \in \mathbb{C}^{m \times n}$ :*

$$\dim(\text{Ran}(M)) + \dim(\text{Ker}(M)) = n. \quad (4.15.5)$$

**Definition 4.15.1** (Spectrum and resolvent). *Let  $X$  be a complex Banach space and  $T : X \rightarrow X$  be a bounded linear operator. Then the spectrum of  $T$  is defined by:*

$$\sigma(T) := \{\lambda \in \mathbb{C} | (T - \lambda I) \text{not bijective}\}. \quad (4.15.6)$$

*The resolvent of  $T$  is defined by:*

$$\rho(T) := \{\lambda \in \mathbb{C} | (T - \lambda I) \text{ bijective and } (T - \lambda I)^{-1} \text{ bounded}\}. \quad (4.15.7)$$

**Corollary 4.15.1.** *By bounded inverse theorem, we have*

$$\sigma(T) = \mathbb{C} \setminus \rho(T). \quad (4.15.8)$$

**Definition 4.15.2** (Point/continuous/residual spectrum). *We have the disjoint union:  $\sigma(T) = \sigma_p(T) \cup \sigma_c(T) \cup \sigma_r(T)$ . We have*

$$\text{Point spectrum :} \sigma_p := \{\lambda \in \mathbb{C} | (T - \lambda I) \text{ not injective}\}, \quad (4.15.9)$$

$$\text{Continuous spectrum :} \sigma_c := \{\lambda \in \mathbb{C} | (T - \lambda I) \text{ injective but not surjective with } \overline{\text{Ran}(T - \lambda I)} = X\}, \quad (4.15.10)$$

$$(T - \lambda I) \text{ injective but not surjective with } \overline{\text{Ran}(T - \lambda I)} = X, \quad (4.15.11)$$

$$\text{Residual spectrum :} \sigma_r := \{\lambda \in \mathbb{C} | (T - \lambda I) \text{ injective but not surjective with } \overline{\text{Ran}(T - \lambda I)} \neq X\}. \quad (4.15.12)$$

$$(T - \lambda I) \text{ injective but not surjective with } \overline{\text{Ran}(T - \lambda I)} \neq X, \quad (4.15.13)$$

$$(4.15.14)$$

[SZQ: 2023.04.02: I need to learn injective, surjective, bijective first!]

## 4.16 Part 31: Spectral Radius

**Definition 4.16.1** (Spectral radius). *X complex Banach space.  $T : X \rightarrow X$  bounded linear operator. We define the spectral radius as*

$$r(T) := \sup \{|\lambda| \}, \quad (4.16.1)$$

where  $\lambda \in \sigma(T)$ .

Here, we have a fig to show this def.

**Theorem 10.** *X complex Banach space,  $T : X \rightarrow X$  bounded linear operator. Then*

- $\sigma(T) \subseteq \mathbb{C}$  is compact
- $X \neq \{0\} \implies \sigma(T) \neq \emptyset$
- $r(T) := \sup |\lambda| = \lim_{k \rightarrow \infty} \|T^k\|^{\frac{1}{k}} = \inf_{k \rightarrow \infty} \|T^k\|^{\frac{1}{k}} \leq \|T\| < \infty,$

where  $\lambda \in \sigma(T)$

*Proof.*

□

[SZQ: 2023.04.02: I need to learn the properties of sepctrum, dual spaces, Vom-Neuman series, Liouville's theorem, Hahn-Banach Theorem!]

## 4.17 Part 32: Normal and Self-Adjoint Operators

**Definition 4.17.1** (Adjoint operator). *Let X be a Hilbert space and  $T : X \rightarrow X$  a bounded lineaar operator. The bounded linear operator  $T^* : X \rightarrow X$  defined by*

$$\langle y, Tx \rangle = \langle T^*y, x \rangle, \forall x, y \in X \quad (4.17.1)$$

is called the adjoint operator of T.

**Definition 4.17.2** (Self Adjoint operator). *Let X be a Hilbert space and  $T : X \rightarrow X$  a bounded lineaar operator. T is called*

- self-adjoint if  $T^* = T$
- skew-adjoint if  $T^* = -T$
- normal if  $T^*T = TT^*$

**Proposition 4.17.1.** *T is normal  $\implies r(T) = \|T\|$ .*

## **Part II**

# **Physics Fundamentals**

## Chapter 5

# Theory of Electrodynamics

### 5.1 零、数学与物理常量基础公式

1. 矢量叉乘公式

$$a \times (b \times c) = b(a \cdot c) - c(a \cdot b) \quad (5.1.1)$$

[\* 请自行结合  $\nabla$  的微分性与线性叠加, 得到  $\nabla \times (a \times b), \nabla \times (\nabla \times a)$  等]

2. 关于  $\nabla$  的相关公式 (推导用):

$$\nabla \cdot \left( \frac{R}{R^3} \right) = 4\pi\delta(R) \quad \nabla f(r) = \frac{\partial f}{\partial r} \nabla r = \frac{\partial f}{\partial r} \hat{r} \quad (5.1.2)$$

3. 关于  $\nabla$  积分的相关公式:

$$\int_V \nabla \cdot A d\tau = \oint_S A \cdot dS \quad \int_V \nabla \psi d\tau = \oint_S \psi dS \quad (5.1.3)$$

$$\int_V \nabla \times A d\tau = \oint_S dS \times A \quad (5.1.4)$$

$$\int_S \nabla \times A \cdot dS = \oint_C A \cdot dl \quad \int_S dS \times \nabla \psi = \oint_C \psi dl \quad (5.1.5)$$

4. 物理常量基本公式

$$\mu_0 \epsilon_0 = \frac{1}{c^2} \quad (5.1.6)$$

### 5.2 麦克斯韦方程组

1. 电场及标量势, 及 2 者关系

$$E(r) = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(r')}{R^3} R d\tau' = -\nabla \varphi \quad \varphi(r) = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(r')}{R} d\tau' \quad (5.2.1)$$

2. 电偶极子 ( $p = ql$ ) 的电场及电势

$$\varphi = \frac{1}{4\pi\epsilon_0} \frac{p \cdot r}{r^3} \quad E = -\frac{1}{4\pi\epsilon_0} \left[ \frac{p - 3(p \cdot \hat{r})\hat{r}}{r^3} \right] \quad (5.2.2)$$

3. 电荷守恒/连续性方程:

$$\frac{\partial \rho}{\partial t} + \nabla \cdot j = 0 \quad j = \rho v \quad (5.2.3)$$

4. 磁场及矢势, 及 2 者关系

$$B = \frac{\mu_0}{4\pi} \int \frac{j(r') d\tau' \times R}{R^3} = \nabla \times A \quad A = \frac{\mu_0}{4\pi} \int \frac{j(r')}{R} d\tau' \quad (5.2.4)$$

5. 点电荷在电磁场中的受力

$$F = q(E + v \times B) \quad (5.2.5)$$

6. 磁偶极子 ( $m = IS$ ) 的矢势及磁场

$$A = \frac{\mu_0}{4\pi} \frac{m \times r}{r^3} \quad B = -\frac{\mu_0}{4\pi} \left[ \frac{m - 3(m \cdot \hat{r})\hat{r}}{r^3} \right] \quad (5.2.6)$$

7. 真空及介质中的 Maxwell 方程组

$$\begin{cases} \nabla \cdot E = \rho/\epsilon_0 \\ \nabla \times E = -\frac{\partial}{\partial t} B \\ \nabla \cdot B = 0 \\ \nabla \times B = \mu_0 j + \mu_0 \epsilon_0 \frac{\partial}{\partial t} E \end{cases} \quad \begin{cases} \nabla \cdot D = \rho_f \\ \nabla \times E = -\frac{\partial}{\partial t} B \\ \nabla \cdot B = 0 \\ \nabla \times H = j_f + \frac{\partial}{\partial t} D \end{cases} \quad (5.2.7)$$

8. 各项同性材料中的本构关系

$$D = \epsilon E, H = B/\mu \text{ (无色散)} \quad (5.2.8)$$

$$D_\omega(r) = \epsilon(\omega)E_\omega(r), B_\omega(r) = \mu(\omega)H_\omega(r) \text{ (色散介质)} \quad (5.2.9)$$

$$(j_\omega = \sigma(\omega)E_\omega) \quad (5.2.10)$$

9. 材料的线性响应

$$P = \epsilon_0 \chi_e E, \quad \epsilon = \epsilon_r \epsilon_0 = (1 + \chi_e) \epsilon_0 \quad (5.2.11)$$

$$M = \frac{1}{\mu_0} \frac{\chi_m}{1 + \chi_m} B, \quad \mu = \mu_r \mu_0 = (1 + \chi_m) \mu_0, \quad (5.2.12)$$

10. 介质中的电荷和电流 (自由、极化、磁化)

$$\rho = \rho_f + \rho_p, \rho_p = -\nabla \cdot P \quad (5.2.13)$$

$$j = j_f + j_m + j_\rho, j_m = \nabla \times M, j_p = \frac{\partial P}{\partial t} \quad (5.2.14)$$

11. 麦克斯韦方程组的边界条件 ( $\nabla \rightarrow n$ )

$$\left\{ \begin{array}{l} n \cdot (D_1 - D_2) = \sigma_f \text{ 自由电荷面密度} \\ n \times (E_1 - E_2) = 0 \\ n \cdot (B_1 - B_2) = 0 \\ n \times (H_1 - H_2) = \alpha_f \text{ 面电流密度} \end{array} \right. \quad (5.2.15)$$

## 5.3 电磁场的守恒定律和对称性

1. 电(磁)场对电荷做功:

$$dR = E \cdot j d\tau dt \quad (5.3.1)$$

2. 真空中电磁场的能量守恒定律及各项的含义

$$\frac{d}{dt} \left[ W_m + \int u \, d\tau \right] = - \oint S_p \cdot dS \quad (5.3.2)$$

$$u(r, t) = \frac{1}{2} \left( \epsilon_0 E^2 + \frac{1}{\mu_0} B^2 \right), \quad S_p(r, t) = \frac{1}{\mu_0} E \times B \quad (5.3.3)$$

其物理意义为: 在一个闭合空间内物理量 ( $W_m + \int u \, d\tau$ ) 的增加等于从边界流入闭合空间的  $S_p$  的大小。其中  $u(r, t)$  是  $r$  点处  $t$  时刻电磁场的能量密度,  $S_p$  即为相应的能流密度, 也称做坡印廷矢量。

3. 真空中电磁场的动量守恒定律及各项的含义

$$\frac{dG_m}{dt} = - \oint dS \cdot \vec{T} - \frac{d}{dt} \int g d\tau \quad (5.3.4)$$

$$\text{电磁场的动量密度 } g = \epsilon_0 (E \times B) = \frac{1}{c^2} S_p \quad (5.3.5)$$

$$\text{动量流密度 } \vec{T} = \frac{1}{2} \left( \epsilon_0 E^2 + \frac{1}{\mu_0} B^2 \right) - \epsilon_0 E E - \frac{1}{\mu_0} B B \quad (5.3.6)$$

$$\text{受力 } F_m = \frac{dG_m}{dt} \quad (5.3.7)$$

4. 带电的运动粒子在外磁场中的总动量

$$p = mv + qA \quad (5.3.8)$$

5. 介质中的电磁能量守恒定律

$$u(r, t) = \frac{1}{2} (\epsilon E^2 + \mu H^2), \quad S_p(r, t) = E \times H \quad (5.3.9)$$

6. 介质中电磁场的动量守恒定律

$$g' = D \times B, \quad \vec{T}' = \frac{1}{2} (E \cdot D + B \cdot H) \vec{I}' - DE - BH \quad (5.3.10)$$

## 5.4 导体静电学

1. 导体静电问题中电势满足的方程及边界条件

$$\begin{cases} \nabla^2 \varphi = -\rho/\epsilon \\ \varphi \text{ 在边界有限} \\ \left. \frac{\partial \varphi}{\partial n} \right|_{\text{边界}} = -\frac{\sigma}{\epsilon}, \quad Q = -\epsilon \oint \frac{\partial \varphi}{\partial n} dS \end{cases} \quad (5.4.1)$$

2. 格林互易定理给定一个有  $m$  个导体组成的体系，假设当导体上的电荷为  $q_1, q_2, \dots$  时，它们的电势等于  $\phi_i$ ，而对应另外一种电荷分布  $q'_i$ ，导体的电势分布为  $\phi'_i$ ，那么有关系式  $\sum_i q_i \phi'_i = \sum_i q'_i \phi_i$ .

3. 导体系中的总能和相互作用能

$$\text{总能 } W = \frac{1}{2} \int \varphi \rho d\tau = \frac{1}{2} \sum_i \phi_i Q_i \quad (5.4.2)$$

$$\text{相互作用能 } W = \frac{1}{2} \sum_i q_i \phi'_i, \quad \phi'_i = \sum_{j \neq i} \phi_j \text{ 为其余电荷在 } q_i \text{ 处电势} \quad (5.4.3)$$

4. 电容的定义

$$q_i = \sum_j C_{ij} \phi_j, \quad \phi_i = \sum_j C_{ij}^{-1} q_j \rightarrow W = \frac{1}{2} \sum_{i,j} C_{ij} \phi_i \phi_j \quad (5.4.4)$$

5. 静电体系的稳定性—汤姆逊定理和恩肖定理

汤姆逊定理：若导体系中每个导体的位置固定不变，电荷可再分布，则体系基态对应，电荷的分布使所有导体均为等势体。

恩肖定理：静电体系的平衡条件是体系中任一导体所处的位置的电场均为 0，因此无约束下静电体系没有平衡态。

6. 导体表面所受静电力

$$F_s = \frac{\epsilon_0}{2} E^2 \hat{e}_n = \frac{1}{2\epsilon_0} \sigma^2 \hat{e}_n \quad (5.4.5)$$

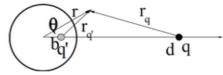
## 5.5 电介质静电学

1. 电介质的边界条件

$$\begin{cases} \varphi_1 = \varphi_2 \\ \epsilon_1 \cdot \frac{\partial \varphi_1}{\partial n} - \epsilon_2 \cdot \frac{\partial \varphi_2}{\partial n} = -\sigma_f (\vec{n} : 2 \rightarrow 1) \end{cases} \quad (5.5.1)$$

2. 唯一性定理对于确定的静电(磁)问题(已知电荷分布和电介质分布  $\rho(r), \epsilon(r)$  或电流和磁介质分布), 边条唯一决定解(前提: 介质中 D 与 E(H 与 A) 之间的本构关系为单调单值)

3. 镜像法(会做)



$$q' = \frac{R}{d}q, \quad b = \frac{R^2}{d} \quad (5.5.2)$$

4. 本征函数展开法

外场为均匀电场时, 对于球/柱体系, 解只具有  $l = 1$  的项。

$$\text{轴对称球坐标系 } \varphi = \left( Ar + \frac{B}{r^2} \right) \cos \theta \quad (5.5.3)$$

$$\text{与 } z \text{ 无关的柱对称 } \varphi = B_0 \ln \rho + (A_1 \rho + B_1 \rho^{-1}) \cos \phi + (C_1 \rho + D_1 \rho^{-1}) \sin \phi \quad (5.5.4)$$

5. 退极场

$$E_d = E_0 - E_{\text{in}} = -L \cdot P / \epsilon_0 \quad (5.5.5)$$

$L$  为退极因子, 取决于物体形状, 对于球  $L = 1/3$ .

6. 多极矩展开

$$\varphi = \frac{Q}{4\pi\epsilon_0 r} - \frac{P}{4\pi\epsilon_0} \nabla \frac{1}{r} + \frac{1}{4\pi\epsilon_0} \frac{1}{6} \vec{D} : \nabla \nabla \frac{1}{r} \quad (5.5.6)$$

$$Q = \int \rho(r') d\tau', P = \int \rho(r') r' d\tau', \vec{D} = 3 \int \rho(r') r' r' d\tau' \quad (5.5.7)$$

7. 多极矩在外场中的作用

$$W = Q\varphi - p \cdot E + \frac{1}{6} \vec{D} : \nabla E \quad (5.5.8)$$

8. 电偶极矩在外场中的受力与力矩

$$F_e = p \cdot \nabla E, \quad M = p \times E \quad (5.5.9)$$

## 5.6 静磁场

1. 磁场的矢势方程及边界条件

$$\nabla^2 A = -\mu_0 j \quad (5.6.1)$$

$$\hat{e}_n \times (A_1 - A_2) = 0 \quad (5.6.2)$$

$$\hat{e}_n \times \left[ \frac{1}{\mu_1} (\nabla \times A_1) - \frac{1}{\mu_2} (\nabla \times A_2) \right] = \alpha_f \quad (5.6.3)$$

2. 静磁场总能

$$U_m = \frac{1}{2} \int B \cdot H d\tau = \frac{1}{2} \int A \cdot j d\tau \quad (5.6.4)$$

3. 磁场的标量势解法中磁标势的定义，引入磁标势的条件及其意义

$$H = -\nabla \varphi_m \quad (5.6.5)$$

1) 无传导电流；2) 引入磁壳使得空间为单连通；

从而保证了  $H$  是个保守场，且保证了磁标势单值性。

4. 磁介质中的边界条件

$$\begin{cases} \varphi_1 = \varphi_2 \\ \mu_1 \cdot \frac{\partial \varphi_1}{\partial n} = \mu_2 \cdot \frac{\partial \varphi_2}{\partial n} \end{cases} \quad (5.6.6)$$

5. 铁磁介质中的边界条件 (饱和磁化为  $M_0^i$ )

$$\begin{cases} \varphi_1 = \varphi_2 \\ \frac{\partial \varphi_1}{\partial n} - \frac{\partial \varphi_2}{\partial n} = e_n \cdot (M_0^1 - M_0^2) \end{cases} \quad (5.6.7)$$

6. 磁多极展开

$$A = A^{(1)} = \frac{\mu_0}{4\pi} \frac{m \times r}{r^3}, m = \frac{1}{2} \int r' \times j \, d\tau' \quad (5.6.8)$$

7. 磁偶极子的能量、受力和力矩

$$U = -m \cdot B; F = \nabla(m \cdot B); \tau = m \times B \quad (5.6.9)$$

## 5.7 似稳场 (准静)

1. 似稳条件

(a) 电磁场变化频率远小于金属特征频率  $\omega \ll \omega_\sigma = \sigma_c/\epsilon$ , 其中  $j = \sigma_c E$ .

(b)  $R \ll \lambda/2\pi$  从而忽略位移电流与辐射

## 2. 似稳场方程

$$\frac{\partial}{\partial t} \begin{pmatrix} H \\ E \end{pmatrix} = \frac{1}{\mu\sigma_c} \nabla^2 \begin{pmatrix} H \\ E \end{pmatrix} \quad (5.7.1)$$

## 3. 趋肤效应与趋肤深度

$$E = \hat{x}A \exp[pz - i\omega t] = \hat{x}E_0 e^{-\alpha z} \cos(\omega t - \alpha z) \quad (5.7.2)$$

$$\delta = \frac{1}{\alpha} = \sqrt{\frac{2}{\mu\omega\sigma_c}} \quad (5.7.3)$$

## 5.8 电磁波的传播

### 1. 电磁波的传播方程

$$\left( \nabla^2 - \epsilon\mu \frac{\partial}{\partial t^2} \right) \begin{pmatrix} E \\ B \end{pmatrix} = 0 \quad (5.8.1)$$

[\* 请自行从 Maxwell 方程推导得到该式.]

### 2. 电磁波的解

$$E(r, t) = E_0 \cos(k \cdot r - \omega t + \phi) \quad (5.8.2)$$

$$B(r, t) = B_0 \cos(k \cdot r - \omega t + \phi) \quad (5.8.3)$$

### 3. 电磁波的色散关系、波速、波长和折射率

$$k^2 = \epsilon\mu\omega^2 \quad (5.8.4)$$

$$v^2 = \frac{1}{\epsilon\mu}, v_p = \omega/k \quad (5.8.5)$$

$$k = 2\pi/\lambda \quad (5.8.6)$$

$$n = \sqrt{\epsilon_r\mu_r} = k\omega/c \quad (5.8.7)$$

### 4. 波矢与电场、磁场的关系

$$k \cdot E_0 = 0, \quad k \cdot B_0 = 0 \quad (5.8.8)$$

$$k \times E_0 = \omega B_0, \quad k \times B_0 = -\epsilon\mu\omega E_0 \quad (5.8.9)$$

### 5. 阻抗的定义

$$Z = \sqrt{\mu/\epsilon}, |E_0| = Z |H_0| \quad (5.8.10)$$

## 6. 电磁波的能流

$$\langle S_p \rangle = \frac{1}{2} \operatorname{Re} (E \times H^*) = \frac{1}{2Z} E_0^2 \hat{k} = \langle u \rangle \cdot v \quad (5.8.11)$$

$$\langle u \rangle = \frac{\epsilon}{2} E_0^2 \quad (5.8.12)$$

7.  $E_0 = \hat{x}E_{x0}e^{i\phi_x} + \hat{y}E_{y0}e^{i\phi_y}$  的线偏振和圆偏振条件

线偏振:  $\phi_x = \phi_y$ ;

圆偏振  $\phi_x - \phi_y = \pm\pi/2, |E_{x0} = E_{y0}|$ , 正负对应右旋、左旋偏振  $\hat{e}_{\text{right}} = (\hat{x} - i\hat{y})/\sqrt{2}, \hat{e}_{\text{left}} = (\hat{x} + i\hat{y})/\sqrt{2}$

## 8. 金属的有效电导率——Drude 模型及其推导

利用散射力模型 (平均  $\tau$  时间受到一次异种粒子的散射而丢失所有的动量)

$$\sigma(\omega) = \frac{n_e e^2}{m(-i\omega + 1/\tau)} \quad (5.8.13)$$

## 9. 金属的有效介电函数及其不同频段的行为

$$\epsilon(\omega) = \epsilon_0 \epsilon_v + i \frac{\sigma(\omega)}{\omega}, \epsilon_r = \epsilon_v - \frac{\omega_p^2}{\omega(\omega + i/\tau)}, \quad \epsilon_v \approx 1 \quad (5.8.14)$$

$$\epsilon_r(\omega) \approx \begin{cases} i \frac{\sigma_c}{\epsilon_0 \omega}, & \leq \text{GHz} \\ 1 - \left(\frac{\omega_p}{\omega}\right)^2, & \text{可见光波段} \end{cases} \quad (5.8.15)$$

在可见光波段, 电子高频下碰撞时间远大于周期, 几乎不表现出损耗。在 GHz 波段, 电子平均碰撞时间远小于电场周期, 表现为欧姆损耗直流良导体特征。

## 10. 等离共振频率及其物理含义

$$\omega_p = \sqrt{\frac{n_e e^2}{\epsilon_0 m}} \quad (5.8.16)$$

表示自由电子气在外场的驱动下集体震荡。

## 11. 导电介质的色散关系

$$k^2 = \left(\frac{\omega}{c}\right)^2 \epsilon_r(\omega) \quad (5.8.17)$$

## 12. 良导体在不同电磁波频率的行为

- (a) 在 GHz 以下,  $1/\omega \gg \tau$  既震荡又衰减, 电磁场有  $\pi/4$  相位差,  $\alpha = \sqrt{\sigma_c \mu_0 \omega / 2}$ , 趋肤深度  $\delta = \alpha^{-1}$ .

- (b) 在光波段  $k^2 = (\omega^2 - \omega_p^2)/c^2$ , 电磁场  $\pi/2$  相位差. 当  $\omega < \omega_p$ , 不传播, 造成反射;  
当  $\omega = \omega_p$  隧穿至极值;  
当  $\omega > \omega_p$ , 紫外透明, 有一定反射.

### 13. 非良导体在不同电磁波频率的行为

非良导体  $\omega \ll 1/\tau$ , 电磁场几乎无相位差, 电磁波传播伴随很小的衰减。

### 14. 各向异性介质/旋光介质中的本构关系

$$D_\omega = \epsilon_0 \vec{\epsilon}_r(\omega) E. \quad (5.8.18)$$

$$\vec{\epsilon}_r(\omega) = \begin{pmatrix} \epsilon_1 & i\epsilon_2 & 0 \\ -i\epsilon_2 & \epsilon_1 & 0 \\ 0 & 0 & \epsilon_3 \end{pmatrix} \quad (5.8.19)$$

$$\epsilon_1 = 1 - \frac{\omega_p^2}{\omega^2 - \omega_B^2}, \epsilon_2 = \frac{\omega_p^2 \omega_B}{(\omega^2 - \omega_B^2) \omega}, \epsilon_3 = 1 - \frac{\omega_p^2}{\omega^2}, \omega_B = -\frac{eB_0}{m} \quad (5.8.20)$$

### 15. 左 ( $k_-$ ) 右 ( $k_+$ ) 旋光的色散关系

$$k_\pm = \frac{\omega}{c} \sqrt{\epsilon_1 \pm \epsilon_2} \quad (5.8.21)$$

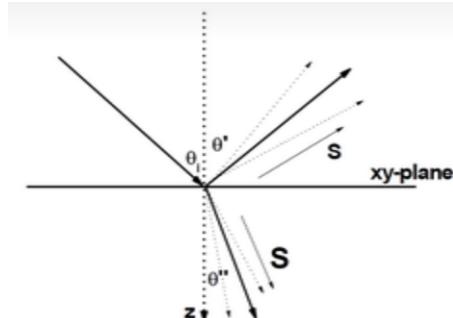
### 16. 法拉第效应

与入射光比较偏振方向旋转了如下的角度  $\Delta\phi = (k_+ - k_-) \cdot d/2$ .

### 17. 电磁波在介质界面反射、折射的基本规律及其物理本质

- (a) 反射波、折射波与入射波  $\omega$  相等——时间平移不变性
- (b) 入射线、反射线和折射线在同一平面内,  $k_{\parallel}$  相同——空间平移不变性
- (c) 在正常介质中, 反射波  $k'_z$  应取负根, 折射波  $k''_z$  应取正根——因果律
- (d) 入射角等于反射角——空间平移不变性
- (e) Snell law 空间平移不变性

$$\frac{\sin \theta}{\sin \theta''} = \frac{k''}{k} = \frac{n_2}{n_1} \quad (5.8.22)$$



## 18. 电磁波在介质界面反射、折射的振幅关系

(a) 横电波 S/TE , 电场垂直于入射面, 有效阻抗  $Z_{eff} = Z / \cos \theta$ 

$$E'_0 = \frac{Z_{eff,2} - Z_{eff,1}}{Z_{eff,2} + Z_{eff,1}} E_0 = r_s \cdot E_0, E''_0 = \frac{2Z_{eff,2}}{Z_{eff,2} + Z_{eff,1}} E_0 = t_s \cdot E_0 \quad (5.8.23)$$

(b) 横磁波 P/TM, 磁场垂直于入射面,  $Z_{eff} = Z \cos \theta$ 

$$H'_0 = \frac{Z_{eff,1} - Z_{eff,2}}{Z_{eff,2} + Z_{eff,1}} H_0 = r_p \cdot H_0, H''_0 = \frac{2Z_{eff,1}}{Z_{eff,2} + Z_{eff,1}} H_0 = t_p \cdot H_0 \quad (5.8.24)$$

19. 电磁波的反射率  $R$  与透射率  $T$ 

$$R = |r|^2, T = \begin{cases} |t_s|^2 \frac{Z_{eff,1}}{Z_{eff,2}} \\ |t_p|^2 \frac{Z_{eff,2}}{Z_{eff,1}} \end{cases} \quad (5.8.25)$$

## 20. 布鲁斯特 Brewster 角及其意义

当反射波与折射波相互垂直时,  $P$  偏振的电磁波完全不被反射, 入射角

$$\theta_B = \arctan(n_2/n_1) \quad (5.8.26)$$

## 21. 全反射临界角

$$\theta_C = \arcsin(n_2/n_1) \quad (5.8.27)$$

**5.9 波导和谐振腔**

## 1. 波导管 (PEC) 的边界条件

$$n \cdot B = 0, n \times E = 0 \quad (5.9.1)$$

$$n \cdot D = \sigma, n \times H = j \quad (5.9.2)$$

## 2. 波导管中波的模式 (偏振)

(a) 横电波 TE:  $B_{0z} \neq 0, E_{0z} = 0$ (b) 横电波 TM:  $B_{0z} = 0, E_{0z} \neq 0$ 

## 3. TE 波的解及基模

$$B_{0z} = B_0 \cos\left(\frac{m\pi}{a}x\right) \cos\left(\frac{n\pi}{b}y\right); B_z = B_{0z} \exp[ik_z z - \omega t] \quad (5.9.3)$$

$$k_c^2 = \pi^2 \left( \frac{m^2}{a^2} + \frac{n^2}{b^2} \right) = k_0^2 - k_z^2 \quad (5.9.4)$$

截止频率, 电磁波能够传播的最低频率:  $\omega_c = ck_c$  (5.9.5)

$$\text{色散关系 } k_z = \frac{\omega}{c} \sqrt{1 - \left( \frac{\omega_c}{\omega} \right)^2} \quad (5.9.6)$$

TE 模的最低阶模式 (基模) 为  $TE_{01}$  或  $TE_{10}$

#### 4. TM 波的解及基模

$$E_{0z} = E_0 \sin \left( \frac{m\pi}{a} x \right) \sin \left( \frac{n\pi}{b} y \right); E_z = E_{0z} \exp [ik_z z - \omega t] \quad (5.9.7)$$

$$k_c^2 = \pi^2 \left( \frac{m^2}{a^2} + \frac{n^2}{b^2} \right) \quad (5.9.8)$$

TM 模的最低阶模式 (基模) 为  $TM_{11}$ .

#### 5. 谐振腔的频率及基模

$$\omega_{\{mnp\}} = c\pi \sqrt{\frac{m^2}{a^2} + \frac{n^2}{b^2} + \frac{p^2}{d^2}} \quad (5.9.9)$$

谐振腔中  $m, n, p$  三个指标中只能有一个不是 0 , 故谐振腔的最低阶模式为 (110), (101) 或 (011).

## 5.10 电磁波的辐射

### 1. 用电势和矢势求得电场

$$E = -\frac{\partial}{\partial t} A - \nabla \varphi \quad (5.10.1)$$

[\* 请自行从 Maxwell 方程推导得到]

### 2. 库仑规范条件与洛伦兹规范条件

$$\nabla \cdot A = 0 \quad (5.10.2)$$

$$\nabla \cdot A + \frac{1}{c^2} \frac{\partial \varphi}{\partial t} = 0 \quad (5.10.3)$$

## 3. 洛伦兹规范条件下势所满足的方程

$$\left( \nabla^2 - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} \right) \begin{pmatrix} \varphi \\ A \end{pmatrix} = \begin{pmatrix} -\rho/\epsilon_0 \\ -\mu_0 j \end{pmatrix} \quad (5.10.4)$$

## 4. 推迟势

$$\varphi(r, t) = \frac{1}{4\pi\epsilon_0} \int \frac{[\rho]}{R} d\tau', \quad [\rho] = \rho(r', t' = t - R/c), R = r - r' \quad (5.10.5)$$

$$A = \frac{\mu_0}{4\pi} \int \frac{[j]}{R} d\tau' \quad (5.10.6)$$

5. 电偶极辐射 ( $[p] = p_0 e^{-i\omega t} e^{ikr}$ )

$$\varphi_1 = -\nabla \cdot \frac{[p]}{4\pi\epsilon_0 r}, [p] = \left[ \int r' \rho d\tau' \right] \quad (5.10.7)$$

$$A_0 = \frac{\mu_0}{4\pi r} [\dot{p}] = \frac{-i\omega\mu_0}{4\pi} \frac{[p]}{r} \quad (5.10.8)$$

远场下的电磁场为 ( $\nabla \rightarrow ik\hat{r}$ )

$$B = \nabla \times A = \frac{\omega^2 \mu_0}{4\pi c r} \hat{r} \times [p] \quad (5.10.9)$$

$$E = -\hat{r} \times (cB) = -\frac{\omega^2}{4\pi\epsilon_0 c^2 r} \hat{r} \times (\hat{r} \times [p]) \quad (5.10.10)$$

## 6. 辐射能流的角分布

$$\langle f(\theta, \phi) \rangle = \langle S_p \rangle r^2 \quad (5.10.11)$$

## 7. 远场下磁偶极辐射

$$A = \frac{i\mu_0\omega}{4\pi r^2 c} r \times [m] \quad (5.10.12)$$

$$B = \nabla \times A = -\frac{\mu_0}{4\pi} \frac{k^2}{r} \hat{r} \times (\hat{r} \times [m]) \quad (5.10.13)$$

$$E = -\hat{r} \times (cB) = -\frac{\mu_0 k^2 c}{4\pi} \hat{r} \times [m] \quad (5.10.14)$$

## 8. 天线电流及矢势

$$I(z', t') = I_0 e^{-i\omega T} \sin \left( k \left( \frac{l}{2} - |z'| \right) \right) \quad (5.10.15)$$

$$A_z = \frac{\mu_0 I_0}{2\pi k r} e^{-i\omega(t-r/c)} \left( \frac{\cos(kl/2) \cos\theta - \cos(kl)}{\sin^2\theta} \right) \quad (5.10.16)$$

## 9. 天线阵的辐射角分布

利用相邻路程差  $l \cos\theta$ ,  $E_i = C(\theta) \exp[ikR_i]/R_i$ , 求和得到, 总的辐射角分布

$$f_{\text{total}}(\theta, \phi) = f_{\text{single}}(\theta, \phi) \frac{\sin^2(m\alpha/2)}{\sin^2(\alpha/2)}, \alpha = kl \cos\theta \quad (5.10.17)$$

## 5.11 相对论电动力学

### 1. 狹义相对论的两条基本假设

- (a) 相对性原理: 自然规律在不同惯性系中的表达式相同
- (b) 光速不变原理, 选择 Maxwell 方程在一切惯性系中形式不变

### 2. 洛伦兹变换

$$x' = (x - vt)/\sqrt{1 - v^2/c^2}, t' = \left(t - vx/c^2\right)/\sqrt{1 - v^2/c^2} \quad (5.11.1)$$

$$x_\mu = (x, y, z, ict), \beta = v/c, \gamma = 1/\sqrt{1 - \beta^2} \quad (5.11.2)$$

$$x_\mu = \alpha_{\nu\mu} x'_\nu, \alpha = \begin{pmatrix} \gamma & & i\beta\gamma & \\ & 1 & & \\ & & 1 & \\ -i\beta\gamma & & & \gamma \end{pmatrix} \quad (5.11.3)$$

$$(5.11.4)$$

### 3. 电荷守恒定律的协变形式

$$J_\mu = (j, ic\rho), \quad \partial_\mu J_\mu = 0, \partial_\mu = \left(\nabla, -i\frac{\partial_t}{c}\right) \quad (5.11.5)$$

### 4. 协变形式的麦克斯韦方程组

$$\partial_\mu F_{\mu\nu} = -\mu_0 J_\nu, \partial_\mu F_{\nu\alpha} + \partial_\alpha F_{\mu\nu} + \partial_\nu F_{\alpha\mu} = 0 (\mu \neq \nu \neq \alpha) \quad (5.11.6)$$

$$A_\mu = (A, i\varphi/c), F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu \quad (5.11.7)$$

## Part III

# Machine Learning Fundamentals

## 5.12 Computational learning theory

# **Part IV**

# **Quantum Fundamentals**

# Chapter 6

## Prerequisites

### 6.1 Hilbert Spaces and Linear Operators

Throughout this course,  $\mathcal{H}$  denotes a finite-dimensional Hilbert space (complex vector space with an associated inner product). Using Dirac's "bra-ket" notation we denote elements of the Hilbert space (called kets) as

$$|\psi\rangle \in \mathcal{H}. \quad (6.1.1)$$

The elements of the dual Hilbert space are called bras and are denoted

$$\langle\psi| \in \mathcal{H}^*, \quad (6.1.2)$$

where  $\langle\psi| = (|\psi\rangle)^\dagger$ . Here,  $X^\dagger := \bar{X}^T$  denotes the Hermitian adjoint (also called the conjugate transpose). We denote

$$B(\mathcal{H}_1, \mathcal{H}_2) := \{\text{linear maps from } \mathcal{H}_1 \text{ to } \mathcal{H}_2\} \quad (6.1.3)$$

and the set of all linear maps to and from the same space will be denoted  $B(\mathcal{H}) = B(\mathcal{H}, \mathcal{H})$ . An operator  $X \in B(\mathcal{H})$  is *normal* if  $XX^T = X^TX$ . Every normal operator has a *spectral decomposition*. That is, there exists a unitary  $U$  and a diagonal matrix  $D$  whose entries are the eigenvalues  $\lambda_1, \dots, \lambda_d \in \mathbb{C}$  of  $X$  such that

$$X = UDU^\dagger. \quad (6.1.4)$$

In other words,

$$X = \sum_{i=1}^d \lambda_i |\psi_i\rangle \langle\psi_i| \quad (6.1.5)$$

where  $X |\psi_i\rangle = \lambda_i |\psi_i\rangle$  and  $U = (|\psi_1\rangle, \dots, |\psi_d\rangle)$ . If  $X$  is Hermitian,  $X = X^\dagger$ , then  $\lambda_i \in \mathbb{R}$ . An operator  $X$  is positive semi-definite (PSD) if

$$\langle \varphi | X | \varphi \rangle \geq 0 \quad \forall |\varphi\rangle \in \mathcal{H}. \quad (6.1.6)$$

As a consequence,  $X \geq 0$  and  $\lambda_i \geq 0$ . It holds that PSD  $\implies$  Hermitian  $\implies$  normal. Unless otherwise stated, we will always assume we are working in an orthonormal basis.

## 6.2 Quantum States

A quantum state  $\rho$  in a Hilbert space  $\mathcal{H}$  is a PSD linear operator with

$$\rho \in B(\mathcal{H}), \quad \rho \geq 0, \quad \text{tr}\rho = 1. \quad (6.2.1)$$

This means that the state has eigenvalues  $\{\lambda_i\}_{i=1}^d$  satisfying  $\lambda_i \geq 0$  and  $\sum_{i=1}^d \lambda_i = 1$ . Thus,  $\{\lambda_i\}_{i=1}^d$  forms a probability distribution.

A *pure quantum state*  $\psi$  is a quantum state with rank 1. We can find  $|\psi\rangle \in \mathcal{H}$  such that  $\psi = |\psi\rangle \langle \psi|$ . In this case,  $\psi$  is called a *projector*. A *mixed state* is a quantum state with rank  $> 1$ . Mixed states are convex combinations of pure states. That is, for every quantum state  $\rho$  with  $r = \text{rank}(\rho)$  there are pure states  $|\psi_i\rangle_{i=1}^k$  ( $k \geq r$ ) and a probability distribution  $\{p_i\}_{i=1}^k$  such that

$$\rho = \sum_{i=1}^k p_i |\psi_i\rangle \langle \psi_i|. \quad (6.2.2)$$

The spectral decomposition of  $\rho$  is a special case of this property.

## 6.3 Composite systems, partial trace, entanglement

Let  $A$  and  $B$  be two quantum systems with Hilbert spaces  $\mathcal{H}_A$  and  $\mathcal{H}_B$ . The *joint system*  $AB$  is described by the Hilbert space  $\mathcal{H}_{AB} := \mathcal{H}_A \otimes \mathcal{H}_B$ . We denote quantum states of the joint system as  $\rho_{AB} \in \mathcal{H}_{AB}$ . The marginal of the bipartite state, denoted  $\rho_A$ , is uniquely defined as the operator satisfying

$$\rho_A := \text{tr}_B \rho_{AB}, \quad (6.3.1)$$

which is defined via  $\text{tr}(\rho_{AB}(X_A \otimes \mathbb{I}_B)) = \text{tr}\rho_A X_A \quad \forall X_A \in B(\mathcal{H}_A)$ . For a Hilbert space with  $|B| := \dim \mathcal{H}_B$ , the explicit form of the partial trace is

$$\text{tr}_B \rho_{AB} = \sum_{i=1}^{|B|} (\mathbb{I}_A \otimes \langle i|_B) \rho_{AB} (\mathbb{I}_A \otimes |i\rangle_B), \quad (6.3.2)$$

for some basis  $\{|i\rangle_B\}_{i=1}^{|B|}$  of  $\mathcal{H}_B$ .

A *product state* on  $AB$  is a state of the form  $\rho_A \otimes \sigma_B$ . The state is called *separable* if it lies in the convex hull of product states:

$$\rho_{AB} = \sum_i p_i \rho_A^i \otimes \sigma_B^i \quad (6.3.3)$$

for some states  $\{\rho_A^i\}_i$  and  $\{\sigma_B^i\}_i$  and probability distribution  $\{p_i\}_i$ . A state is called *entangled*, if it is not separable. An entangled state of particular interest is the *maximally entangled state*. Let  $d = \dim \mathcal{H}$ ,  $\{|i\rangle\}_{i=1}^d$  be a basis for  $\mathcal{H}$ . A maximally entangled state is expressed as

$$|\phi^+\rangle = \frac{1}{\sqrt{d}} \sum_{i=1}^d |i\rangle \otimes |i\rangle \in \mathcal{H} \otimes \mathcal{H} \quad (6.3.4)$$

## 6.4 Measurements

The most general measurement is given by a *positive operator-valued measure* (POVM)  $E = \{E_i\}_i$  where  $E_i \geq 0 \quad \forall i$  and  $\sum_i E_i = \mathbb{I}$ . Then, for a quantum system  $\mathcal{H}$  in state  $\rho$ , the probability of obtaining measurement outcome  $i$  is given by  $p_i = \text{tr}[\rho E_i]$ . So, we have

$$\sum_i p_i = \sum_i \text{tr}[\rho E_i] = \text{tr} \left[ \rho \sum_i E_i \right] = \text{tr}[\rho \mathbb{I}] = \text{tr} \rho = 1, \quad (6.4.1)$$

for all normalized quantum states. A *projective measurement*  $\Pi = \{\Pi_i\}$  is a POVM with the added property of orthogonality, which for projectors means

$$\Pi_i \Pi_j = \delta_{ij} \Pi_i. \quad (6.4.2)$$

Any basis  $\{|e_i\rangle\}_{i=1}^{\dim \mathcal{H}}$  gives rise to a projective measurement  $\Pi = \{|e_i\rangle \langle e_i|\}_{i=1}^{\dim \mathcal{H}}$ .

## 6.5 Entropies

The *Shannon entropy*  $H(p)$  of a probability distribution  $p = \{p_1, \dots, p_d\}$  is defined as  $H(p) = -\sum_{i=1}^d p_i \log p_i$ , where the logarithm is base 2 unless otherwise specified. Note that when the logarithm is base 2, the entropy has units of *bits*. The *von Neumann entropy*  $S(\rho)$  of a quantum state  $\rho$  is defined as

$$S(\rho) = -\text{tr} [\rho \log \rho] = H(\{\lambda_1, \dots, \lambda_d\}), \quad (6.5.1)$$

where  $\rho = \sum_i \lambda_i |\psi_i\rangle \langle \psi_i|$  is a spectral decomposition of  $\rho$  and where the logarithm of an operator is obtained by first diagonalizing the matrix representing the operator and then taking the logarithm of the diagonal elements. That is,

$$\log \rho = \sum_{i:\lambda_i > 0} \log(\lambda_i) |\psi_i\rangle \langle \psi_i|. \quad (6.5.2)$$

# Chapter 7

## Prerequisites

### 7.1 Hilbert Spaces and Linear Operators

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where  $\langle\psi| = (|\psi\rangle)^\dagger$ . Here,  $X^\dagger := \bar{X}^T$  denotes the Hermitian adjoint (also called the conjugate transpose). We denote

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In other words,

$$X = \sum_{i=1}^d \lambda_i |\psi_i\rangle \langle\psi_i| \quad (7.1.5)$$

where  $X |\psi_i\rangle = \lambda_i |\psi_i\rangle$  and  $U = (|\psi_1\rangle, \dots, |\psi_d\rangle)$ . If  $X$  is Hermitian,  $X = X^\dagger$ , then  $\lambda_i \in \mathbb{R}$ . An operator  $X$  is positive semi-definite (PSD) if

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This means that the state has eigenvalues  $\{\lambda_i\}_{i=1}^d$  satisfying  $\lambda_i \geq 0$  and  $\sum_{i=1}^d \lambda_i = 1$ . Thus,  $\{\lambda_i\}_{i=1}^d$  forms a probability distribution.

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$$\rho = \sum_{i=1}^k p_i |\psi_i\rangle \langle \psi_i|. \quad (7.2.2)$$

The spectral decomposition of  $\rho$  is a special case of this property.

## 7.3 Composite systems, partial trace, entanglement

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$$\rho_A := \text{tr}_B \rho_{AB}, \quad (7.3.1)$$

which is defined via  $\text{tr}(\rho_{AB}(X_A \otimes \mathbb{I}_B)) = \text{tr}\rho_A X_A \quad \forall X_A \in B(\mathcal{H}_A)$ . For a Hilbert space with  $|B| := \dim \mathcal{H}_B$ , the explicit form of the partial trace is

$$\text{tr}_B \rho_{AB} = \sum_{i=1}^{|B|} (\mathbb{I}_A \otimes \langle i|_B) \rho_{AB} (\mathbb{I}_A \otimes |i\rangle_B), \quad (7.3.2)$$

for some basis  $\{|i\rangle_B\}_{i=1}^{|B|}$  of  $\mathcal{H}_B$ .

A *product state* on  $AB$  is a state of the form  $\rho_A \otimes \sigma_B$ . The state is called *separable* if it lies in the convex hull of product states:

$$\rho_{AB} = \sum_i p_i \rho_A^i \otimes \sigma_B^i \quad (7.3.3)$$

for some states  $\{\rho_A^i\}_i$  and  $\{\sigma_B^i\}_i$  and probability distribution  $\{p_i\}_i$ . A state is called *entangled*, if it is not separable. An entangled state of particular interest is the *maximally entangled state*. Let  $d = \dim \mathcal{H}$ ,  $\{|i\rangle\}_{i=1}^d$  be a basis for  $\mathcal{H}$ . A maximally entangled state is expressed as

$$|\phi^+\rangle = \frac{1}{\sqrt{d}} \sum_{i=1}^d |i\rangle \otimes |i\rangle \in \mathcal{H} \otimes \mathcal{H} \quad (7.3.4)$$

## 7.4 Measurements

The most general measurement is given by a *positive operator-valued measure* (POVM)  $E = \{E_i\}_i$  where  $E_i \geq 0 \quad \forall i$  and  $\sum_i E_i = \mathbb{I}$ . Then, for a quantum system  $\mathcal{H}$  in state  $\rho$ , the probability of obtaining measurement outcome  $i$  is given by  $p_i = \text{tr}[\rho E_i]$ . So, we have

$$\sum_i p_i = \sum_i \text{tr}[\rho E_i] = \text{tr} \left[ \rho \sum_i E_i \right] = \text{tr}[\rho \mathbb{I}] = \text{tr} \rho = 1, \quad (7.4.1)$$

for all normalized quantum states. A *projective measurement*  $\Pi = \{\Pi_i\}$  is a POVM with the added property of orthogonality, which for projectors means

$$\Pi_i \Pi_j = \delta_{ij} \Pi_i. \quad (7.4.2)$$

Any basis  $\{|e_i\rangle\}_{i=1}^{\dim \mathcal{H}}$  gives rise to a projective measurement  $\Pi = \{|e_i\rangle \langle e_i|\}_{i=1}^{\dim \mathcal{H}}$ .

## 7.5 Entropies

The *Shannon entropy*  $H(p)$  of a probability distribution  $p = \{p_1, \dots, p_d\}$  is defined as  $H(p) = -\sum_{i=1}^d p_i \log p_i$ , where the logarithm is base 2 unless otherwise specified. Note that when the logarithm is base 2, the entropy has units of *bits*. The *von Neumann entropy*  $S(\rho)$  of a quantum state  $\rho$  is defined as

$$S(\rho) = -\text{tr} [\rho \log \rho] = H(\{\lambda_1, \dots, \lambda_d\}), \quad (7.5.1)$$

where  $\rho = \sum_i \lambda_i |\psi_i\rangle \langle \psi_i|$  is a spectral decomposition of  $\rho$  and where the logarithm of an operator is obtained by first diagonalizing the matrix representing the operator and then taking the logarithm of the diagonal elements. That is,

$$\log \rho = \sum_{i:\lambda_i > 0} \log(\lambda_i) |\psi_i\rangle \langle \psi_i|. \quad (7.5.2)$$

# Chapter 8

## Classes of Quantum Channels

### 8.1 Qubit channels

Note: please see Nielsen and Chuang for figures depicting these channels' actions on the Bloch sphere. I may make nice ones at some point but for now, I am too lazy.

1. **Bit flip channel.** First, we look at the classical bit flip channel, the properties of which were established by Claude Shannon in his seminal 1948 paper. The quantum version is simply given by the Pauli  $X$  operator,

$$X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad (8.1.1)$$

which acts as

$$X |0\rangle = |1\rangle, \quad (8.1.2)$$

$$X |1\rangle = |0\rangle. \quad (8.1.3)$$

It's eigenbasis is  $|\pm\rangle := \frac{1}{\sqrt{2}}(|0\rangle \pm |1\rangle)$ . So,  $X |+\rangle |+\rangle$  and  $X |-\rangle = -|-\rangle$ . The channel acts as

$$\mathcal{F}_p^X : \rho \mapsto (1-p)\rho + pX\rho X \quad (8.1.4)$$

So, the Kraus operators are  $\sqrt{1-p}\mathbb{I}, \sqrt{p}X$

## 2. Phase-flip/Z-dephasing channel

$$Z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \quad (8.1.5)$$

$$Z |0\rangle = |0\rangle \quad (8.1.6)$$

$$Z |1\rangle = -|1\rangle \quad (8.1.7)$$

$$Z |+\rangle = |-\rangle \quad (8.1.8)$$

$$Z |-\rangle = |+\rangle \quad (8.1.9)$$

The channel is given as

$$\mathcal{F}_p^Z : \rho \mapsto (1-p)\rho + pZ\rho Z \quad (8.1.10)$$

## 3. Bit-phase flip/Y-dephasing channel

$$Y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \quad (8.1.11)$$

$$|\pm i\rangle = \frac{1}{\sqrt{2}}(|0\rangle \pm i|1\rangle) \quad (8.1.12)$$

$$Y |i\rangle = |i\rangle \quad (8.1.13)$$

$$Y |-i\rangle = -|-i\rangle \quad (8.1.14)$$

$$Y |0\rangle = i|1\rangle \quad (8.1.15)$$

$$Y |1\rangle = -i|0\rangle \quad (8.1.16)$$

The channel is then

$$\mathcal{F}_p^Y : \rho \mapsto (1-p)\rho + pY\rho Y \quad (8.1.17)$$

$$= (1-p)\rho + pXZ\rho ZX \quad (8.1.18)$$

**4. Depolarizing channel.** Corresponds to an error model in which each Pauli error occurs with equal probability. The channel is

$$\mathcal{D}_p : \rho \mapsto (1-p)\rho + \frac{p}{3}(X\rho X + Y\rho Y + Z\rho Z) \quad (8.1.19)$$

with equal probability  $p/3$ . The Kraus operators are then  $\sqrt{1-p}\mathbb{I}$ ,  $\sqrt{p/3}X$ ,  $\sqrt{p/3}Y$ ,  $\sqrt{p/3}Z$ . Thus, the Kraus rank is 4 for this channel when  $p \in (0, 1)$ . An alternative representation is given by

$$\rho \mapsto (1-q)\rho + q\text{tr}(\rho)\frac{\mathbb{I}}{2} \quad (8.1.20)$$

This represents replacing the input state with the maximally mixed state with “probability”  $q$  ( $q$  can be greater than 1).

What about the relation  $p \leftrightarrow q$ ? To deduce this, we can use the identity

$$\frac{1}{2}\mathbb{I}_2 = \frac{1}{4}(\rho + X\rho X + Y\rho Y + Z\rho Z) \quad \forall \rho \in B(\mathbb{C}^2) \quad (8.1.21)$$

We can then derive the relationship between  $p$  and  $q$ . We have

$$\frac{1}{2}\mathbb{I}_2 = \frac{1}{4}(\rho + X\rho X + Y\rho Y + Z\rho Z) \quad \forall \rho \in B(\mathbb{C}^2) \quad (8.1.22)$$

$$(1-p)\rho + p\frac{\mathbb{I}}{2} = (1-q)\rho + \frac{q}{4}(\rho + X\rho X + Y\rho Y + Z\rho Z) \quad (8.1.23)$$

$$\Rightarrow \boxed{q = \frac{4}{3}p} \quad (8.1.24)$$

**5. Generalized Pauli channel.** Let  $\vec{p} = (p_0, p_1, p_2, p_3)$  be a probability distribution.

The generalized Pauli channel is

$$\mathcal{N}_{\vec{p}}(\rho) = p_0\rho + p_1X\rho X + p_2Y\rho Y + p_3Z\rho Z \quad (8.1.25)$$

where we recover the depolarizing channel by setting  $p_0 = 1 - p$  and  $p_i = p/3$ . Note that for any  $\vec{p}$ , this channel is unital.

Pauli channels are interesting from an information theoretic standpoint. Classically, they are very easily understood. Quantum mechanically, they very much are not understood (except in the case of flip or dephasing channels).

In what sense are these flip channels de-phasing? Let us look at the example of the phase flip channel

$$\mathcal{F}_p^Z : \rho \mapsto (1-p)\rho + pZ\rho Z. \quad (8.1.26)$$

We can understand the justification of the term de-phasing by looking at the action on the density matrix

$$\begin{pmatrix} \rho_{11} & \rho_{12} \\ \rho_{21} & \rho_{22} \end{pmatrix} \mapsto \begin{pmatrix} \rho_{11} & (1-2p)\rho_{12} \\ (1-2p)\rho_{21} & \rho_{22} \end{pmatrix} \quad (8.1.27)$$

where  $\rho \geq 0, \text{tr}(\rho) = 1 \implies \rho_{11} + \rho_{22} = 1$ . We can make the following observations

1. If  $\rho = x|0\rangle\langle 0| + (1-x)|1\rangle\langle 1|$ , then

$$\mathcal{F}_p^Z(\rho) = \rho \quad \forall p \in [0, 1] \quad (8.1.28)$$

2.  $p = \frac{1}{2}$ : the channel is diagonal in the z-basis for all input states

We can also think about sending classical information through this channel. Let us encode 0 as  $|0\rangle\langle 0|$  and 1 as  $|1\rangle\langle 1|$ . That is, we are encoding one bit in one qubit. We can do this reliably with this channel because

$$\mathcal{F}_p^Z(|0\rangle\langle 0|) = |0\rangle\langle 0|, \quad (8.1.29)$$

$$\mathcal{F}_p^Z(|1\rangle\langle 1|) = |1\rangle\langle 1|. \quad (8.1.30)$$

This means that we can send one bit through the channel. The key idea is: this channel preserves classical information because it preserves the orthogonality of the basis states. The same would hold for the  $X$  channel but you would have to encode info in the  $|\pm\rangle$  basis.

- 6. Amplitude damping channel.** Physical model: 2-level system (e.g. an atom with a ground state  $|0\rangle$  and excited state  $|1\rangle$ ). If the system is in an excited state  $|1\rangle$ , it decays with a certain probability,  $\gamma$ , emitting a photon to the environment. How can we capture this decay process mathematically? Consider the isometry

$$|0\rangle_A \mapsto |0\rangle_B |0\rangle_E, \quad (8.1.31)$$

$$|1\rangle \mapsto \sqrt{1-\gamma} |1\rangle_B |0\rangle_E + \sqrt{\gamma} |0\rangle_B |1\rangle_E. \quad (8.1.32)$$

Compactly we can write

The Kraus operators for this channel are

$$K_0 = \langle 0|_E V = |0\rangle\langle 0|_B + \sqrt{1-\gamma} |1\rangle\langle 1|_B = \begin{pmatrix} 1 & 0 \\ 0 & \sqrt{1-\gamma} \end{pmatrix}, \quad (8.1.33)$$

$$K_1 = \langle 1|_E V = \sqrt{\gamma} |0\rangle_B \langle 1|_A = \begin{pmatrix} 0 & \gamma \\ 0 & 0 \end{pmatrix}. \quad (8.1.34)$$

In the Kraus representation, then, we can write the amplitude damping channel

$$\mathcal{A}_\gamma : \rho \mapsto K_0 \rho K_0^\dagger + K_1 \rho K_1^\dagger. \quad (8.1.35)$$

We note that the amplitude damping channel is not unital. Therefore, it is not a Pauli (or even a mixed-unitary) channel. Interestingly, it is completely understood how to send *quantum information* through the amplitude damping channel; however, it is not known how much classical information can be sent through this channel.

- 7. Erasure channel.** Classical model:

Quantum version:  $\mathcal{E}_p : B(\mathcal{H}) \rightarrow B(\mathcal{H} \otimes \mathbb{C})$ . The channel acts as

$$\rho \mapsto (1-p)\rho + p \text{ptr}(\rho) |e\rangle\langle e|. \quad (8.1.36)$$

This channel has Kraus operators given by  $K_0\sqrt{1-p}\mathbb{I}$ ,  $K_1\sqrt{p}|0\rangle\langle 0|$ ,  $K_2\sqrt{p}|e\rangle\langle 1|$ . Note that  $|e\rangle$  must be orthogonal to all  $\rho \in B(\mathcal{H})$ . That is,  $\langle |e\rangle\langle e|, \rho \rangle = 0 \quad \forall \rho \in B(\mathcal{H})$ . Bob can always tell whether erasure happened by performing a measurement! This is a very well-understood quantum channel.

## 8.2 Generalized dephasing channels

A generalized dephasing channel leaves a *fixed* orthonormal basis invariant and it dephases off-diagonal elements with respect to that fixed basis. That is, we lose quantum coherences under the action of a generalized dephasing channel.

Construction:  $\mathcal{H} = \mathbb{C}^d$  with orthonormal basis  $\{|i\rangle\}_{i=1}^d$ . Choose an environment  $\mathcal{H}_E$  with  $\dim \mathcal{H}_E := |E| \geq 2$ , and let  $\{\varphi_i\}_{i=1}^{|E|}$  be *some* set of pure states on  $E$ . This set is normalized but not orthogonal. Then, define the isometry

$$V : |i\rangle_A \mapsto |i\rangle_B \otimes |\varphi_i\rangle_E. \quad (8.2.1)$$

The Stinespring extension is

$$\mathcal{N}(\rho_A) = \text{tr}_E V \rho_A V^\dagger, \quad (8.2.2)$$

$$= \sum_{i,j} |i\rangle \rho \langle i| |i\rangle \langle j|_B \text{tr}(|\varphi_i\rangle \langle \varphi_j|), \quad (8.2.3)$$

$$= \sum_{i,j} \langle \varphi_i | \varphi_j \rangle \langle i | \rho | j \rangle |i\rangle \langle i|_B. \quad (8.2.4)$$

Let us confirm this channel acts as expect.

$$[\mathcal{N}(\rho)]_{kk} = \langle \varphi_k | \varphi_k \rangle \langle k | \rho | k \rangle = \rho_{kk}, \quad (8.2.5)$$

$$[\mathcal{N}(\rho)]_{jk} = \langle \varphi_j | \varphi_k \rangle \langle j | \rho | k \rangle, \quad (8.2.6)$$

but  $0 < \langle \varphi_j | \varphi_k \rangle \leq 1$ . Thus, we see that the diagonals are preserved but the off-diagonals are dephased.

Higher-dimensional example: Let  $d \geq 2$ , and define two unitaries

$$X |i\rangle = |i + 1 \bmod d\rangle \quad (\text{shift operator}) \quad (8.2.7)$$

$$Z |j\rangle = \omega^j |j\rangle, \text{ where } \omega = \exp\left(\frac{2\pi i}{d}\right) \quad (\text{clock operator}) \quad (8.2.8)$$

This generalizes the Pauli operators:  $X^d = Z^d = \mathbb{I}$ . These are the generators of the Heisenberg-Weyl group:

$$\{\omega^j Z^k X^l : j, k, l \in [d]\}. \quad (8.2.9)$$

So, for  $\rho \in B(\mathbb{C}^d)$  we maps like

$$\rho \mapsto (1-p)\rho + pX\rho X^\dagger, \quad (8.2.10)$$

$$\rho \mapsto (1-p)\rho + \frac{p}{3}Z\rho Z^\dagger + \frac{2p}{3}Z^2\rho(Z^\dagger)^2. \quad (8.2.11)$$

The generalized de-phasing channels have full classical capacity

$$C(\mathcal{N}) = \log d \quad (\text{in general, } C(\mathcal{N}) \leq \log d) \quad (8.2.12)$$

Sketch of a proof: we know that for a generalized dephasing channel, there exists an orthonormal basis  $\{|i\rangle\}_i$  such that

$$\mathcal{N}(|i\rangle\langle i|) = |i\rangle\langle i| \quad (8.2.13)$$

of classical signals/messages  $x_i, \dots, x_d$ . We can use the encoding

$$x_i \mapsto |i\rangle\langle i| \quad (8.2.14)$$

where  $\langle i|j\rangle = \delta_{ij}$ . Thus,  $d$  messages can be sent perfectly. So  $\log d$  bits of classical info can be sent through  $N$ . Bob can measure the output to retrieve the classical message.

### 8.2.1 Detour: Holevo information ( $\mathcal{X}$ -quantity)

Recall: the von Neumann entropy is defined as

$$S(\rho) = -\text{tr}\rho \log \rho. \quad (8.2.15)$$

When one has the spectral decomposition of  $\rho = \sum_i \lambda_i |\psi_i\rangle\langle\psi_i|$ , the von Neumann entropy is just the Shannon entropy for the eigenvalue distribution

$$S(\rho) = -\sum_i \lambda_i \log \lambda_i. \quad (8.2.16)$$

If we then have an ensemble of quantum states  $E = \{p_x, \rho_x\}$ , we have

$$\mathcal{X}(E, \mathcal{N}) = S\left(\sum_x p_x \mathcal{N}(\rho_x)\right) - \sum_x p_x S(\mathcal{N}(\rho_x)) \quad (8.2.17)$$

The Holevo information is then defined as

$\boxed{\text{Holevo Information, } \mathcal{X}(\mathcal{N}) = \max_E \mathcal{X}(E, \mathcal{N})}$

(8.2.18)

A fundamental result due to Holevo, Schumacher, and Westmoreland is

$\boxed{C(\mathcal{N}) \geq \mathcal{X}(\mathcal{N})}$

(8.2.19)

If  $\mathcal{N}$  is a generalized dephasing channel, there exists an orthonormal basis  $\{|i\rangle\}$  such that  $\mathcal{N}(|i\rangle\langle i|) = |i\rangle\langle i|$  for all  $i$ . Let our ensemble be

$$\rho_i = |i\rangle\langle i| \quad (8.2.20)$$

$$p_i = \frac{1}{d}, \quad (8.2.21)$$

which implies

$$\bar{\rho} = \sum_i p_i \rho_i = \frac{1}{d} \mathbb{I} \quad (8.2.22)$$

Then,  $E_n = \{\frac{1}{d}, \rho_i\}$ , so the Holevo quantity is

$$\mathcal{X}(E_n, \mathcal{N}) = S\left(\sum_i p_i \mathcal{N}(\rho_i)\right) - \sum_i p_i S(\mathcal{N}(\rho_i)) \quad (8.2.23)$$

$$= S\left(\sum_i p_i \rho_i\right) - \sum_i p_i S(\rho_i) \quad (8.2.24)$$

$$= \log d - 0 \quad (8.2.25)$$

$$= \log d \quad (8.2.26)$$

Thus, we have the following chain of inequalities

$$\log d \leq \mathcal{X}(\mathcal{N}) \leq C(\mathcal{N}) \leq \log d \implies C(\mathcal{N}) = \log d \quad (8.2.27)$$

as desired.

### 8.2.2 Some comments on last lecture

- The rank of the Kraus operators is not unitarily invariant. We can see this by examining the 50-50 dephasing channel.

$$\rho \mapsto \frac{1}{2}\rho + \frac{1}{2}Z\rho Z \implies K_0 = \frac{1}{\sqrt{2}}\mathbb{I}, \quad K_1 = \frac{1}{\sqrt{2}}Z \quad (8.2.28)$$

We can apply a Hadamard transform to the operators to obtain new Kraus operators

$$L_i = \sum_j H_{ij} K_j, \quad (8.2.29)$$

$$\implies L_0 = |0\rangle\langle 0|, L_1 = |1\rangle\langle 1|, \quad (8.2.30)$$

which are rank-1 operators now when the previous Kraus operators were full rank.

- Every unital qubit channel is unitarily equivalent to a Pauli channel. That is, if  $\mathcal{N} : B(\mathbb{C}^2) \rightarrow B(\mathbb{C}^2)$  is a unital channel, then there are unitaries  $U, V$  such that

$$M(\rho) = U\mathcal{N}(V\rho V^\dagger)U^\dagger \quad (8.2.31)$$

is Pauli.

- mixed unitary channels

$$\rho \mapsto \sum_i p_i U_i \rho U_i^\dagger, \quad U_i \text{ unitary} \quad (8.2.32)$$

$$\mathbb{I} \mapsto \sum_i p_i U_i U_i^\dagger = \mathbb{I} \quad \text{unital!} \quad (8.2.33)$$

Every unital qubit channel is mixed unitary. Watrous' book provides a nice example of a unital channel that is *not* mixed unitary:

$$X \in B(\mathbb{C}^3), \quad \rho \mapsto \frac{1}{2} \text{tr}(X) \mathbb{I} - \frac{1}{2} X^T \quad (8.2.34)$$

### 8.3 Entanglement-breaking channels

Reminder: A bipartite  $\rho_{AB}$  is called separable if it lies in the convex hull of product states. That is,

$$\rho_{AB} \in \text{conv}\{\omega_A \otimes \sigma_B : \omega_{A(B)} \text{ state on } \mathcal{H}_{A(B)}\} \quad (8.3.1)$$

Explicitly,

$$\rho_{AB} = \sum_i p_i \omega_A^i \otimes \sigma_B^i \quad (8.3.2)$$

**Definition 8.3.1.** A channel  $\mathcal{N} : A \rightarrow B$  is entanglement-breaking if

$$(\mathbb{I}_R \otimes \mathcal{N})(\rho_{RA}) \quad (8.3.3)$$

is separable for any  $\rho_{RA}$ .

**Proposition 8.3.1.** The following are all equivalent

- $\mathcal{N} : A \rightarrow B$  is entanglement breaking
- $\tau_{AB}^{\mathcal{N}} = (\mathbb{I}_A \otimes \mathcal{N})(\gamma_{AA'})$  is separable
- $\mathcal{N}$  has a Kraus representation with rank-1 Kraus operators
- $\mathcal{N}$  is a measure-and-prepare channel: there exists POVM  $E = \{E_i\}_i$  and states  $\{\sigma_i\}_i$  such that

$$\mathcal{N}(\rho) = \sum_i \text{tr}(\rho E_i) \sigma_i \quad (8.3.4)$$

where we recall that a POVM must satisfy  $E_i \geq 0, \sum_i E_i = \mathbb{I}$ .

*Proof.* (a  $\implies$  b):  $(\mathbb{I}_R \otimes \mathcal{N})(\rho_{RA})$  is separable for all  $\rho_{RA}$ . In particular, for  $\gamma_{RA}$ ,  $(|\gamma\rangle_{RA}) = \sum_i |i\rangle_R |i\rangle_A$ .

(b  $\implies$  c):  $\tau_{AB}^{\mathcal{N}}$  is separable: there exist pure states  $\psi_i = |\psi_i\rangle \langle \psi_i|_A$  and  $\varphi_i = |\varphi_i\rangle \langle \varphi_i|_B$  such that  $\frac{1}{d}\tau_{AB}^{\mathcal{N}} = \sum_i p_i \psi_i \otimes \varphi_i$ , where  $d = |A|$ . Then, set  $K_i = \sqrt{p_i d} |\varphi_i\rangle_B \langle \bar{\psi}_i|_A$ . We can then check the action of these operators

$$\frac{1}{d} \sum_i (\mathbb{I} \otimes K_i) (\gamma_{AA'})(\mathbb{I}_A \otimes K_i)^\dagger = \frac{1}{d} \sum_{i,j,k} |j\rangle \langle k|_A \otimes K_i |j\rangle \langle k|_{A'} K_i^\dagger \quad (8.3.5)$$

$$= \sum_{i,j,k} dp_i |j\rangle \langle k|_A \otimes \langle \bar{\psi}_i| |j\rangle \langle k| |\bar{\psi}_i\rangle |\varphi_i\rangle \langle \varphi_i|_B \quad (8.3.6)$$

$$= \sum_{i,j,k} p_i |j\rangle \langle k|_A \otimes \langle \bar{\psi}_i| |j\rangle \langle k| |\bar{\psi}_i\rangle |\varphi_i\rangle \langle \varphi_i|_B \quad (8.3.7)$$

$$= \sum_i p_i \left( \sum_{j,k} \langle k| \bar{\psi}_i |j\rangle |j\rangle \langle k| \right) \otimes |\varphi_i\rangle \langle \varphi_i|_B \quad (8.3.8)$$

$$= \sum_i p_i (\psi_i) \otimes \varphi_i \quad (8.3.9)$$

$$= \frac{1}{d} \tau_{AB}^{\mathcal{N}} \quad (8.3.10)$$

$$\sum_i K_i^\dagger K_i = \mathbb{I} \quad (8.3.11)$$

$$= d \sum_i p_i |\bar{\psi}_i\rangle \langle \varphi_i| \langle \bar{\psi}_i| \quad (8.3.12)$$

$$= d \sum_i p_i \bar{\psi}_i \quad (8.3.13)$$

$$= d \left( \sum_i p_i \bar{\psi}_i \right) \quad (8.3.14)$$

$$= d \frac{1}{d} \mathbb{I}_A \quad (8.3.15)$$

$$= \mathbb{I}_A \quad (8.3.16)$$

(c  $\implies$  d):  $\text{rank}(K_i) = 1$ :  $K_i = |\mathcal{X}_i\rangle_B \langle \omega_i|_A$  for some vectors  $|\mathcal{X}_i\rangle_B, |\omega_i\rangle_A, \langle \mathcal{X}_i| \mathcal{X}_i\rangle = 1$ .

$$\mathcal{N}(\rho) = \sum_i K_i \rho K_i^\dagger \quad (8.3.17)$$

$$= \sum_i \langle \omega_i | \rho | \omega_i \rangle |\mathcal{X}_i\rangle \langle \mathcal{X}_i|_B \quad (8.3.18)$$

POVM:  $\omega = \{\omega_i\}$ , states:  $\mathcal{X}_i$

$$\mathbb{I} = \sum_i K_i^\dagger K_i \quad (8.3.19)$$

$$= \sum_i |\omega_i\rangle \langle \mathcal{X}_i | \mathcal{X}_i \rangle \langle \omega_i| \quad (8.3.20)$$

$$= \sum_i |\omega_i\rangle \langle \omega_i| \quad (8.3.21)$$

(d  $\implies$  a): Let  $\rho_{RA}$  be arbitrary:

$$(\mathbb{I}_R \otimes \mathcal{N})(\rho_{RA}) = \sum_i \text{tr}_B [(\mathbb{I}_A \otimes E_i)\rho_{RA}] \otimes \sigma_i \quad (8.3.22)$$

$$= \sum_i \text{tr}_B [(\mathbb{I}_R \otimes \sqrt{E_i})\rho_{RA}(\mathbb{I}_R \otimes \sqrt{E_i})] \otimes \sigma_i \quad (8.3.23)$$

$$= \sum_i p_i \omega_i \otimes \sigma_i \quad (8.3.24)$$

where  $p_i = \text{tr}(E_i \rho_{RA})$ ,  $\omega_i = \frac{1}{p_i} \text{tr}_B [(\mathbb{I}_R \otimes \sqrt{E_i})\rho_{RA}(\mathbb{I}_R \otimes \sqrt{E_i})] \geq 0$ .  $\square$

We note that there are *some* channel capacities of entanglement-breaking channels that are understood.

- Quantum information transmission is equivalent to generating entanglement. Because entanglement-breaking channels cannot do that, their quantum capacity is zero.
- $C(\mathcal{N}) \geq \mathcal{N}$ ,  $C(\mathcal{N}) = \sup_{n \in \mathbb{N}} \frac{1}{n} \mathcal{X}(\mathcal{N}^{\otimes n})$ . Entanglement-breaking channels destroy entanglement between different inputs. This implies that

$$C(\mathcal{N}) = \mathcal{X}(\mathcal{N}) \quad (8.3.25)$$

BUT:  $\mathcal{X}(\mathcal{N})$  is NP-hard to compute, so this relationship doesn't actually get us much.

## 8.4 PPT-channels

Checking separability is NP-hard. So, is there some easier criterion? The standard relaxed criterion is following.

**Definition 8.4.1.** *Peres-Horodecki criterion: if  $\rho_{AB}$  is separable if  $\rho_{AB}$  has a positive partial transpose with respect to either party.*

Let us see why this is a reasonable criterion for separability. If  $\rho_{AB}$  is separable,

$$\rho_{AB} = \sum_i p_i \omega_A^i \otimes \sigma_B^i \quad (8.4.1)$$

$$\implies \rho_{AB}^{T_B} = \sum_i p_i \omega_A^i \otimes (\sigma_B^i)^T \geq 0 \quad (8.4.2)$$

- if  $\rho_{AB}$  is NPT  $\implies \rho_{AB} \notin \text{SEP}$ .
- if  $|A| \cdot |B| \leq 6$ , then  $\rho_{AB} \in \text{SEP} \Leftrightarrow \rho_{AB} \in \text{PPT}$ .

**Definition 8.4.2.** A channel  $\mathcal{N} : A \rightarrow B$  is called PPT if  $(\mathbb{I}_R \otimes \mathcal{N})(\rho_{RA})$  is PPT for all  $\rho_{RA}$ .

**Proposition 8.4.1.** The following are all equivalent

- $\mathcal{N} : A \rightarrow B$  is PPT
- $\tau_{AB}^{\mathcal{N}}$  is PPT
- $\vartheta \circ \mathcal{N}$  is CP (Recall that when  $\vartheta : X \mapsto X^T : (\mathbb{I}_R \otimes \vartheta)(\gamma) = \mathbb{F} \not\geq 0$ )

*Proof.* (a  $\Rightarrow$  b): True by definition. We know

$$(\mathbb{I}_A \otimes \mathcal{N})(\rho_{AA'}) \quad (8.4.3)$$

is PPT, in particular for  $\rho = \gamma_{AA'}$ .

(b  $\Rightarrow$  c):

$$(\mathbb{I}_A \otimes \vartheta \circ \mathcal{N})(\gamma) = (\mathbb{I} \otimes \vartheta)(\mathbb{I} \otimes \mathcal{N})(\gamma_{AA'}) \quad (8.4.4)$$

$$= (\mathbb{I}_A \otimes \vartheta)(\tau_{AB}^{\mathcal{N}}) \quad (8.4.5)$$

$$= (\tau_{AB}^{\mathcal{N}})^{T_B} \geq 0 \quad (8.4.6)$$

(c  $\Rightarrow$  a)  $\vartheta \circ \mathcal{N}$  is CP:

$$(\mathbb{I}_R \otimes \vartheta \circ \mathcal{N})(\rho_{RA}) \geq 0 \quad \forall \rho_{RA} \geq 0 \quad (8.4.7)$$

This implies  $\mathcal{N}$  is PPT.  $\square$

**Remarks:** In general,  $\vartheta \circ \mathcal{N}$  is *not* CP. However, for every CP map  $\mathcal{N}$ , the map  $\vartheta \circ \mathcal{N} \circ \vartheta$  is CP (see exercises).

What about the capacities of PPT channels?

- Horodeckis: PPT states are undistillable. Entanglement distillation: given iid copies of a state  $\rho_{AB}$ , the goal is to convert these copies into a smaller number

of maximally entangled states. If there is an entanglement distillation protocol (LOCC) such that

1. error of the protocol tends to 0 as  $n \Rightarrow \infty$
2. rate  $c := \frac{1}{n} \log m_n \rightarrow c > 0$  as  $n$  goes to infinity

then  $\rho_{AB}$  is *distillable*. We conclude that the quantum capacity for PPT channels is zero. This result even holds if two-way classical communication is allowed!

We note that the set of protocols allowing only one-way communication is a strict subset of the set of protocols allowing two-way communication. Even still, the Horodecki's showed that the quantum capacity of all PPT channels is zero in both cases. Also note that the classical capacity is generally unknown.

## 8.5 Anti-degradable channels

Let  $\mathcal{N}$  be a quantum channel from  $A$  to  $B$  with an isometry  $V : \mathcal{H}_A \rightarrow \mathcal{H}_B \otimes \mathcal{H}_E$  such that  $\mathcal{N}(\rho) = \text{tr}_E V \rho V^\dagger$ . Recall that the complementary channel is  $\mathcal{N}^c(\rho) = \text{tr}_B V \rho V^\dagger$ .

**Definition 8.5.1.**  $\mathcal{N}$  as above is called antidegradable if there exists a channel  $\mathcal{A} : E \rightarrow B$  such that

$$\mathcal{N} = \mathcal{A} \circ \mathcal{N}^c \quad (8.5.1)$$

Intuition: Eve (environment) can locally obtain Bob's output via the channel  $\mathcal{A}$ . Antidegradable channels cannot transmit quantum information and thus have zero quantum capacity,  $Q(\mathcal{N}) = 0$ .

A “proof” of this is as follows. Assume that this channel has non-zero quantum capacity. This means that Alice can faithfully send qubits to Bob at a positive rate. But there is a protocol based on the channel  $\mathcal{A}$  that lets Eve implement the same protocol that Alice and Bob use. This violates no-cloning and is thus prohibited.

**Examples of anti-degradable channels:**

1. erasure channel  $\mathcal{E}_p : \rho \mapsto (1 - p)\rho + \text{ptr}(\rho)|e\rangle\langle e|$  for  $p \geq \frac{1}{2}$ .

Let  $\mathcal{H}_1 = \mathbb{C}^2$  be the input space. Let the  $\mathcal{H}_2 = \mathbb{C}$  be an erasure flag. Let  $\mathcal{E}_p : B(\mathcal{H}_1) \rightarrow B(\mathcal{H}_1 \oplus \mathcal{H}_2)$ . That is, an input state maps as

$$\rho \mapsto (1 - p)\tilde{\rho} + \text{ptr}(\rho)|e\rangle\langle e| \quad (8.5.2)$$

So, how do we embed our input state in the larger space? We use a pretty trivial embedding that is represented as:

$$\hat{\rho} = \begin{pmatrix} \rho_{00} & \rho_{01} & 0 \\ \rho_{10} & \rho_{11} & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad (8.5.3)$$

and the erasure flag is simply

$$|e\rangle\langle e| = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix}. \quad (8.5.4)$$

Note that the complementary channel of the erasure channel has the same form but with the probabilities flipped:

$$\mathcal{E}_p^c : B(\mathcal{H}_1) \rightarrow B(\mathcal{H}_1 \oplus \mathcal{H}_2), \rho \mapsto p\tilde{\rho} + (1-p)\text{tr}(\rho)|e\rangle\langle e|. \quad (8.5.5)$$

Now consider  $p \geq \frac{1}{2}$  and define  $q = \frac{2p-1}{p}$ . the idea is that we erase  $\tilde{\rho}$  with probability  $q$  and do nothing with  $|e\rangle\langle e|$ .

$$\mathcal{E}_q(p\rho) = p(1-q)\tilde{\rho} + pq\text{tr}(\rho)|e\rangle\langle e| \quad (8.5.6)$$

$$= (1-p)\tilde{\rho} + (2p-1)|e\rangle\langle e| \quad (8.5.7)$$

$$= (1-p)|e\rangle\langle e| \quad (8.5.8)$$

**almost certainly messed up p's and q's here.**

We need to extend the action of  $\mathcal{E}_q$  to  $B(\mathcal{H}_1 \otimes \mathcal{H}_2)$ . The solution is to define  $\mathcal{A}$  in the Kraus representation. The Kraus operators are

$$K_0 = \sqrt{1-q} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad (8.5.9)$$

$$K_1 = \sqrt{q} \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{pmatrix} \quad (8.5.10)$$

$$K_2 = \sqrt{q} \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 1 & 0 \end{pmatrix} \quad (8.5.11)$$

$$K_3 = \sqrt{q} \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix} \quad (8.5.12)$$

where the first three operators correspond to erasure on  $B(\mathcal{H}_1 \oplus 0)$  and the last to doing nothing on  $B(0 \oplus \mathcal{H}_2)$ . This definition ensures that  $\mathcal{E}_p = \mathcal{A} \circ \mathcal{E}_p^c$  for  $p \geq \frac{1}{2}$ . We have  $B(\mathcal{H}_1 \oplus 0), B(0 \oplus \mathcal{H}_2) : \mathcal{A}_p = \mathcal{E}_p \oplus \mathbb{I}$ .

For  $\rho \in B(\mathcal{H}_1)$ ,  $\sigma = (1 - \lambda)\tilde{\rho} + \lambda |e\rangle\langle e|$ . This has the block matrix:

$$\begin{pmatrix} (1 - \lambda)\rho & 0 \\ 0 & \lambda \end{pmatrix} \quad (8.5.13)$$

So

$$(\mathcal{E}_p \oplus \mathbb{I})(\sigma) = \begin{pmatrix} \mathcal{E}_p((1 - \lambda)\rho) & 0 \\ 0 & \mathbb{I}(\lambda) \end{pmatrix} \quad (8.5.14)$$

Then for  $\omega \in B(\mathcal{H}_1 \oplus \mathcal{H}_2)$  we have

$$\omega = \begin{pmatrix} \omega_{\mathcal{H}_1} & * \\ * & \omega_{\mathcal{H}_2} \end{pmatrix} \quad (8.5.15)$$

which has off-diagonal elements. We can get rid of these by measuring with respect to  $\mathcal{H}_1 \oplus \mathcal{H}_2 = \mathcal{H}$ . Let  $P_2 = |e\rangle\langle e|, P_1 = \mathbb{I} - |e\rangle\langle e|$ . Then  $P_1\omega P_1 + P_2\omega P_2$  is diagonal. Finally,  $\mathcal{A}_M = P_1 \cdot P_1 + P_2 \cdot P_2$  which implies  $\mathcal{A} = \mathcal{A}_p \circ \mathcal{A}_M$  where  $\mathcal{A}_p = \mathcal{E}_p \oplus \mathbb{I}$ .

2. amplitude damping channel  $\mathcal{A}_\gamma$  for  $\gamma \geq \frac{1}{2}$ .
3. depolarizing channel  $\mathcal{D}_p$  for  $p \geq \frac{1}{4}$

**Definition 8.5.2.** Let  $\mathcal{N} : A \rightarrow B$  be a quantum channel and let  $G$  be a group with unitary representations  $U_g$  on  $\mathcal{H}_A$  and  $V_g$  on  $\mathcal{H}_B$ . Then  $\mathcal{N}$  is called covariant with respect to  $(G, U_g, V_g)$  if

$$V_g \mathcal{N}(\cdot) V_g^\dagger = \mathcal{N}(U_g \cdot U_g^\dagger) \quad (8.5.16)$$

for all  $g \in G$ .

### Representation theory basics:

- $(\varphi, V)$  is a representation of  $G$ :  $\varphi : G \rightarrow GL(V), \varphi(gh) = \varphi(g)\varphi(h)$ . Subspace  $W \leq V$  is called  $G$ -invariant, if  $\varphi(g)w \in W \quad \forall w \in W, \forall g \in G$ .
- $\{0\}, V$  are always  $G$ -invariant
- $(\varphi, V)$  is called irreducible if  $\{0\}, V$  are the only  $G$ -inv subspaces.
- Schur's lemma:  $(\varphi, V), (\psi, W)$  representations of  $G$ .  $G$ -linear map  $f : f \circ \varphi(g) = \psi(g) \circ f$ . If  $\varphi, \psi$  are irreducible representations, then either  $V \not\cong W$  and  $f = 0$  or  $V \cong W$  and  $f = \lambda \mathbb{I}_{V \rightarrow W}$  for some  $\lambda \in \mathbb{C}$ .

**Some examples:**

- Pauli channels:  $\rho \mapsto p_0\rho + p_1X\rho X + p_2Y\rho Y + p_3Z\rho Z$

Covariance group: Pauli group,  $P = \{\pm 1, \pm i\} \cup \{\mathbb{I}, X, Y, Z\}$

$$\text{For } \theta_1, \theta_2 \in P: \vartheta_1 \theta_2 \cdot \theta_2^\dagger \vartheta_1^\dagger = \vartheta_2 \theta_1 \cdot \theta_1^\dagger \vartheta_2^\dagger$$

- depolarizing channel:  $\rho \mapsto (1-p)\rho + \frac{p}{3}(X\rho X + Y\rho Y + Z\rho Z)$ . The covariance group is  $\mathcal{U}(2)$ . This is easy to see in the "q-representation":  $\rho \mapsto (1-q)\rho + q\text{tr}\rho \frac{1}{2}\mathbb{I}$

$$\mathcal{D}_q(U\rho U^\dagger) = (1-q)U\rho U^\dagger + q\text{tr}(U\rho U^\dagger)\frac{1}{2}\mathbb{I} = U\mathcal{D}_p(\rho)U^\dagger \quad (8.5.17)$$

In  $d \geq$  dimensions:  $\rho \mapsto (1-q)\rho + q\text{tr}\rho \frac{1}{d}\mathbb{I}$ . The covariance group is  $\mathcal{U}(d)$ .

- Amplitude damping channel  $A_\gamma = \{K_0, K_1\}$  as normal. Covariance group:  $\{\mathbb{I}, Z\} \cong \mathbb{Z}_2$ .
- Erasure channel  $\mathcal{E}_p(\rho) = (1-p)\rho + p\text{tr}\rho |e\rangle\langle e|$ . The covariance group is again  $\mathcal{U}(2)$ . Similarly, in  $d \geq 2$  dimensions, the  $d$ -dimensional erasure channel with covariance group  $\mathcal{U}(d)$ .

**Proposition 8.5.1.** A channel  $\mathcal{N} : A \rightarrow B$  is  $(G, U_g, V_g)$ -covariant iff  $\tau^{\mathcal{N}_{AB}} = (\bar{U}_g \otimes V_g)\tau_{AB}^{\mathcal{N}}(\bar{U}_g \otimes V_g)^\dagger$  for all  $g \in G$ .

*Proof.* ( $\Rightarrow$ ):  $V_g\mathcal{N}(\cdot)V_g^\dagger = \mathcal{N}(U_g \cdot U_g^\dagger)$  for all  $g \in G$ . This is equivalent to  $\mathcal{N}(\cdot) = V_g^\dagger \mathcal{N}(U_g \cdot U_g^\dagger) V_g$  for all  $g \in G$ .

$$\tau_{AB}^{\mathcal{N}} = (\mathbb{I} \otimes \mathcal{N})(\gamma) \quad (8.5.18)$$

$$= (\mathbb{I}_1 \otimes V_g^\dagger)(\mathbb{I} \otimes \mathcal{N}) \left( (\mathbb{I}_1 \otimes U_g)\gamma(\mathbb{I} \otimes U_g^\dagger) \right) (\mathbb{I}_1 \otimes V_g) \quad (8.5.19)$$

$$= (\mathbb{I} \otimes V_g)(\mathbb{I} \otimes \mathcal{N}) \left( (U_g^T \otimes \mathbb{I}_2)\gamma(U_g^T \otimes \mathbb{I}_2)^\dagger \right) (\mathbb{I} \otimes V_g) \quad (8.5.20)$$

$$= (U_g^T \otimes V_g^\dagger)(\mathbb{I} \otimes \mathcal{N})(\gamma)(\bar{U}_g \otimes V_g) \quad (8.5.21)$$

$$\implies \tau_{AB}^{\mathcal{N}} = (U_g^T \otimes V_g^\dagger)\tau_{AB}^{\mathcal{N}}(\bar{U}_g \otimes V_g) \quad \text{for all } g \in G \quad (8.5.22)$$

( $\Leftarrow$ ):  $\tau_{AB}^{\mathcal{N}} = (\bar{U}_g \otimes V_g)\tau_{AB}^{\mathcal{N}}(\bar{U}_g \otimes V_g)^\dagger$  for all  $g \in G$ . This implies  $(U_g^T \otimes \mathbb{I})\tau_{AB}^{\mathcal{N}}(U_g^T \otimes \mathbb{I})^\dagger = (\mathbb{I} \otimes V_g)\tau_{AB}^{\mathcal{N}}(\mathbb{I} \otimes V_g)^\dagger$ .

Recall the choi isomorphism:  $\mathcal{N}(X) = \text{tr}_1\tau_{AB}^{\mathcal{N}}(X^T \otimes \mathbb{I})$ .

$$\mathcal{N}(U_g X U_g^\dagger) = \text{tr}_1 \left[ \tau_{AB}^{\mathcal{N}} ((U_g X U_g^\dagger)^T \otimes \mathbb{I}) \right] \quad (8.5.23)$$

$$= \text{tr}_1 \left[ (U_g^T \otimes \mathbb{I}) \tau_{AB} (U_g^T \otimes \mathbb{I})^\dagger (X^T \otimes \mathbb{I}) \right] \quad (8.5.24)$$

$$= \text{tr}_1 \left[ (\mathbb{I} \otimes V_g) \tau_{AB} (\mathbb{I} \otimes V_g)^\dagger (X^T \otimes \mathbb{I}) \right] \quad (8.5.25)$$

$$= V_g \text{tr}_1 [\tau_{AB}^{\mathcal{N}} (X^T \otimes \mathbb{I})] V_g^\dagger \quad (8.5.26)$$

$$= V_g \mathcal{N}(X) V_g^\dagger \quad (8.5.27)$$

□

**Problem:** with the above proposition is that it is basis-dependent (via the Choi operator).

**Solution:** is that there is also a basis independent version based on the Jamialkowski operator  $J_{AB}^{\mathcal{N}} = (\mathbb{I} \otimes \mathcal{N})(\mathbb{F}_{AA'})$  where  $\mathbb{F}_{AB} = |\gamma\rangle\langle\gamma|^{T_A}$ . In terms of the Jamialkowski operator,  $(G, U_g, V_g)$ -covariance of  $\mathcal{N}$  is equivalent to  $(U_g \otimes V_g) J_{AB}^{\mathcal{N}} (U_g \otimes V_g)^\dagger = J_{AB}^{\mathcal{N}}$  for all  $g \in G$ .

We note that a  $d$ -dimensional depolarizing channel;  $\rho \mapsto (1-q)\rho + q\text{tr}\rho \frac{1}{d}\mathbb{I}_d$  has  $\mathcal{U}(d)$  as its covariance group.

**Proposition 8.5.2.** *Let  $\mathcal{N} : A \rightarrow B$  be a channel with input and output spaces of equal dimension  $d$ . If  $U\mathcal{N}(\cdot)U^\dagger = \mathcal{N}(U \cdot U^\dagger)$  for all  $U \in \mathcal{U}(d)$ , then*

$$\mathcal{N} = (1 - q) \cdot + q\text{tr}(\cdot) \frac{1}{d}\mathbb{I}_d \quad (8.5.28)$$

with  $q = (1 - f)/(1 - d)$  where  $f = \langle\gamma|\tau^{\mathcal{N}}|\gamma\rangle$ .

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# Chapter 9

# Introduction to Quantum Optics

## 9.1 Introduction

- classicaal: classical atom and light
- semiclassical: quantized atom and classical light
- quantum mechanical: quantized atom and light

### Light-Atom Interaction Hamiltonian

- classical dipole in eletric field: dipole moment  $\vec{d} = q\vec{r}$ ,  $U_I = -\vec{d} \cdot \vec{E}$ . We have

$$\hat{H}_I = -\hat{d} \cdot \vec{E}(\vec{v}_0, t), \quad (9.1.1)$$

where  $\hat{d} = q\hat{v}$  is the dipole operator.

- induced atomic dipole

## 9.2 Light Atom Quantum Evolution

**Time Evolution** We have the Schrodinger equation (both sides) as

$$i\hbar \frac{\partial}{\partial t} |\Psi(t)\rangle = (\hat{H}_0 + \hat{H}_I(t)) |\Psi(t)\rangle, \quad (9.2.1)$$

where the general ansatz (assumption) is

$$|\Psi(t)\rangle = \sum_n c_n(t) e^{-iE_n t/\hbar} |n\rangle, \quad (9.2.2)$$

and

$$\hat{H}_0|n\rangle = E_n|n\rangle \quad (9.2.3)$$

is the atomic eigenstates. Inserting  $|\Psi(t)\rangle$  and  $\hat{H}_0|n\rangle$  into Schrodinger equation, we get

$$i\hbar \sum_n \left\{ \dot{c}_n e^{-iE_nt/\hbar} |n\rangle - \frac{iE_n}{\hbar} c_n e^{-iE_nt/\hbar} |n\rangle \right\} = \sum_n \left\{ c_n e^{-iE_nt/\hbar} |n\rangle + c_n e^{-iE_nt/\hbar} \hat{H}_I |n\rangle \right\} \quad (9.2.4)$$

$$\implies i\hbar \sum_n \dot{c}_n e^{-iE_nt/\hbar} |n\rangle = \sum_n c_n e^{-iE_nt/\hbar} \hat{H}_I |n\rangle \quad (9.2.5)$$

$$\implies i\hbar \dot{c}_n e^{-iE_k t/\hbar} |k\rangle = \sum_n c_n(t) e^{-iE_nt/\hbar} \langle k | \hat{H}_I(t) | n \rangle \quad (9.2.6)$$

$$\implies i\hbar \dot{c}_k = \sum_n c_n(t) e^{-iE_{n,k}t/\hbar} \langle k | \hat{H}_I(t) | n \rangle, \quad (9.2.7)$$

where we use

$$\langle k | n \rangle = \delta_{kn}, \quad (9.2.8)$$

$$E_{n,k} = E_n - E_k, \quad (9.2.9)$$

$$\omega_{nk} = (E_n - E_k)/\hbar. \quad (9.2.10)$$

and  $\langle k | \hat{H}_I(t) | n \rangle$  is the matrix element.

### 9.3 Time Dependent Perturbation Theory

Recall the time evolution:

$$i\hbar \dot{c}_k = \sum_n c_n(t) e^{-i\omega_{nk}t} \langle k | \hat{H}_I(t) | n \rangle, \quad (9.3.1)$$

and

$$\omega_{nk} = (E_n - E_k)/\hbar. \quad (9.3.2)$$

Consider the Simplification (Perturbation Theory)

- System only in state  $|1\rangle$  at  $t = 0 \implies c_1|0\rangle = 1$  (only the ground state  $|1\rangle$ ),
- Perturbative treatment of interaction term: weak perturbation  $\forall |c_k(t)|^2 \ll 1$ .

We then have

$$i\hbar \dot{c}_k = e^{i\omega_{1k}t} \langle k | \hat{H}_I(t) | 1 \rangle, \quad (9.3.3)$$

with  $c_k(0) = 0$ , we obtain:

$$c_k(t) = \frac{1}{i\hbar} \int_0^t e^{-i\omega_{1k}t'} \langle k | \hat{H}_I(t') | 1 \rangle dt'. \quad (9.3.4)$$

**Example 9.3.1** (Sinusoidal perturbation). Define

$$\hat{H}(t) = \hat{H}_I e^{-i\omega t}. \quad (9.3.5)$$

Given the figure in the video, we have

$$c_k(T) = \frac{1}{i\hbar} \int_0^T e^{i\Delta\omega t} \langle k | \hat{H}_I | 1 \rangle dt \quad (9.3.6)$$

$$\implies \text{Transition probability } P_{k1}(T) = |c_k(T)|^2 = \frac{1}{\hbar^2} |\langle k | \hat{H}_I | 1 \rangle|^2 Y(\Delta\omega, T), \quad (9.3.7)$$

with

$$Y(\Delta\omega, T) = \frac{\sin^2(\Delta\omega T/2)}{(\Delta\omega/2)^2} \quad (9.3.8)$$

$$\sim \text{sinc}^2 x, \quad (9.3.9)$$

where  $\Delta\omega = \omega - \omega_{1k}$  is the detuning.

Let's take a look at the sinc function  $Y(\Delta\omega, T) = \text{sinc}^2 x$ . Transition for  $\Delta\omega \leq \frac{2\pi}{T}$ , we have  $\Delta\omega \cdot T \leq 2\pi$ , which implies

$$\Delta E \cdot T \leq \hbar, \quad (9.3.10)$$

which is the time-frequency uncertainty. (The expression in the video seems wrong, so I make corrections above.) We have the following case

$$\frac{1}{2\pi T} Y(\Delta\omega, T) \xrightarrow{T \rightarrow \infty} \delta(\Delta\omega), \quad (9.3.11)$$

then we have

$$P_{k1}(T \rightarrow \infty) = \frac{2\pi}{\hbar^2} |\langle k | \hat{H}_I | i \rangle|^2 \delta(\Delta\omega) T. \quad (9.3.12)$$

**Fermi's Golden Rule**  $|k\rangle$  Quasi continuum of final states. We have the transition probability

$$P_{k1} = \Gamma_{k1} T, \quad (9.3.13)$$

where

$$\Gamma_{k1} = \frac{2\pi}{\hbar} |\langle k | \hat{H}_I | 1 \rangle|^2 \rho(E_k = E_1 + \hbar\omega) \quad (9.3.14)$$

is called the Femi's Golden Rule,

$$|\langle k | \hat{H}_I | 1 \rangle|^2 \quad (9.3.15)$$

is the coupling strength  $\propto E_0^2$  and  $\propto I$ ,

$$\rho(E_k = E_1 + \hbar\omega) \quad (9.3.16)$$

is the density states which is number of available final states to the system,

$$\Gamma_{k1} \hat{=} Transition\ Rate = \frac{dP_{k1}}{dT}, \quad (9.3.17)$$

and density states

$$\rho(E) = \frac{dN}{dE}, \quad (9.3.18)$$

where  $\Delta N$  is the number of states in an energy interval  $\Delta E$  around energy  $E_k$  and we let  $\Delta E$  approaches 0.

## 9.4 Two Level Atom (TLA)

Given by the figure, in state  $|1\rangle$ , we have  $E_1 = \hbar\omega_1$  and in state  $|2\rangle$ , we have  $E_2 = \hbar\omega_2$  and  $E_2 - E_1 = \hbar(\omega_2 - \omega_1) = \omega_{21}$ . We have the Hamiltonian

$$\hat{H} = \hat{H}_0 - \hat{d} \cdot E(t), \quad (9.4.1)$$

where

$$E(t) = \varepsilon E_0 \cos(\omega t), \quad (9.4.2)$$

where  $\varepsilon$  is the polarization vector,  $E_0$  is the field amplitude, and  $\omega$  is the frequency of the light field.

**Ansatz for Solving TLA** We have

$$|\Psi(t)\rangle = c_1(t)e^{-i\omega_1 t}|1\rangle + c_2(t)e^{-i\omega_2 t}|2\rangle. \quad (9.4.3)$$

**Time Evolution Amplitude** We have

$$\dot{c}_1(t) = i \frac{d_{12}^\varepsilon E_0}{\hbar} e^{-\omega_{21}} \cos(\omega t) c_2(t) \quad (9.4.4)$$

$$\dot{c}_2(t) = i \frac{d_{12}^\varepsilon E_0}{\hbar} e^{+\omega_{21}} \cos(\omega t) c_1(t), \quad (9.4.5)$$

where

$$d_{12}^\varepsilon = \langle 1 | \hat{d} \cdot \varepsilon | 2 \rangle \quad (9.4.6)$$

$$= \langle 1 | \hat{d} | 2 \rangle \cdot \varepsilon \quad (9.4.7)$$

$$= \langle 1 | \hat{d}_x | 2 \rangle \cdot \varepsilon_x + \langle 1 | \hat{d}_y | 2 \rangle \cdot \varepsilon_y + \langle 1 | \hat{d}_z | 2 \rangle \cdot \varepsilon_z. \quad (9.4.8)$$

is the Dipole Matrix Element, which is the atomic property and we assume it's real. We also define

$$\Omega_0 = \frac{d_{12}^\varepsilon E_0}{\hbar} \quad (9.4.9)$$

as the Rabi frequency.

**Time Evolution** Using Euler' form, we have

$$\dot{c}_1(t) = i \frac{\Omega_0}{2} e^{-\omega_{21}} (e^{i\omega t} + e^{-i\omega t}) c_2(t) \quad (9.4.10)$$

$$\dot{c}_2(t) = i \frac{\Omega_0}{2} e^{+\omega_{21}} (e^{i\omega t} + e^{-i\omega t}) c_1(t) \quad (9.4.11)$$

by

$$\cos \alpha = \frac{1}{2} (e^{i\alpha} + e^{-i\alpha}) \quad (9.4.12)$$

and

$$e^{i\alpha} = \cos \alpha + i \sin \alpha. \quad (9.4.13)$$

**Rotating Wave Approximation** We have

$$\dot{c}_1(t) = i \frac{\Omega_0}{2} (e^{+i(\omega-\omega_{21})t} + e^{-i(\omega+\omega_{21})t}) c_2(t) \quad (9.4.14)$$

$$\dot{c}_2(t) = i \frac{\Omega_0}{2} (e^{-i(\omega-\omega_{21})t} + e^{+i(\omega+\omega_{21})t}) c_1(t), \quad (9.4.15)$$

and we ignore the sum frequency term and get

$$\dot{c}_1(t) = i \frac{\Omega_0}{2} e^{+i(\omega-\omega_{21})t} c_2(t) \quad (9.4.16)$$

$$\dot{c}_2(t) = i \frac{\Omega_0}{2} e^{-i(\omega-\omega_{21})t} c_1(t), \quad (9.4.17)$$

which is a good approcimation for detwining  $\delta = \omega - \omega_{21} \approx 0$ . We introduce

$$\tilde{c}_1(t) = c_1(t) e^{-i\frac{\delta}{2}t} \quad (9.4.18)$$

$$\tilde{c}_2(t) = c_2(t) e^{+i\frac{\delta}{2}t}. \quad (9.4.19)$$

$$(9.4.20)$$

**Ansatz Wavefunctions for TLA** Whole time evolution in state amplitudes

$$|\Psi(t)\rangle = c'_1(t)|1\rangle + c'_2(t)|2\rangle. \quad (9.4.21)$$

Time evolution when field is off

$$|\Psi(t)\rangle = c'_1(0)e^{-i\omega_1 t}|1\rangle + c'_2(0)e^{-i\omega_2 t}|2\rangle. \quad (9.4.22)$$

However, this is boring. We chose different ansatz as

$$|\Psi(t)\rangle = c_1(t)e^{-i\omega_1 t}|1\rangle + c_2(t)e^{-i\omega_2 t}|2\rangle \quad (9.4.23)$$

$$\iff |\Psi(t)\rangle = c_1(t)|1\rangle + c_2(t)e^{-i\omega_{21} t}|2\rangle, \quad (9.4.24)$$

where  $c_1(t)$  and  $c_2(t)$  capture time evolution on top of eigenstate evolution! We now have

$$|\Psi(t)\rangle = c_1(t)|1\rangle + c_2(t)e^{-i\omega_{21} t}|2\rangle, \quad (9.4.25)$$

which is called the rotating frame of atom. We also have Rotating frame of light field as

$$|\Psi(t)\rangle = \tilde{c}_1(t)|1\rangle + \tilde{c}_2(t)e^{-i\omega t}|2\rangle, \quad (9.4.26)$$

where  $\omega$  is the light frequency,  $\tilde{c}_1$  and  $\tilde{c}_2$  describe time evolution on top of fast light field oscillation.

**Solving the TLA Dynamics** We have the following equations:

$$\frac{d}{dt} \begin{pmatrix} \tilde{c}_1(t) \\ \tilde{c}_2(t) \end{pmatrix} = \frac{i}{2} \begin{pmatrix} -\delta & \Omega_0 \\ \Omega_0 & +\delta \end{pmatrix} \begin{pmatrix} \tilde{c}_1(t) \\ \tilde{c}_2(t) \end{pmatrix}. \quad (9.4.27)$$

Considering the simplest case  $\delta = 0$

$$\frac{d}{dt} \tilde{c}_1(t) = \frac{i}{2} \Omega_0 \tilde{c}_2(t) \quad (9.4.28)$$

$$\frac{d}{dt} \tilde{c}_2(t) = \frac{i}{2} \Omega_0 \tilde{c}_1(t). \quad (9.4.29)$$

Take time derivative of the first equation, then we have

$$\ddot{\tilde{c}}_1(t) = -\frac{\Omega_0^2}{4} \tilde{c}_1(t), \quad (9.4.30)$$

the solutions of which are

$$\tilde{c}_1(t) = \cos(\Omega_0 t/2) \quad (9.4.31)$$

$$\tilde{c}_2(t) = i \sin(\Omega_0 t/2) \quad (9.4.32)$$

for  $\tilde{c}_1(0) = 1$  and  $\tilde{c}_2(0) = 0$ . Also we can obtain the excited state probability as

$$P_2(t) = |c_2(t)|^2 \quad (9.4.33)$$

$$= |\tilde{c}_2(t)|^2. \quad (9.4.34)$$

**Rabi Oscillations (Resonant Case)** Nonlinear Response can be seen from the figure.

**General Rabi Oscillations (with detuning)** Given the figurem.

$$|\tilde{c}_2(t)|^2 = \frac{\Omega_0^2}{\Omega} \sin^2 \left( \frac{1}{2} \Omega t \right) \quad (9.4.35)$$

$$= \frac{\Omega_0^2}{2\Omega^2} \{1 - \cos(\Omega t)\}, \quad (9.4.36)$$

where  $\Omega = \sqrt{\Omega_0^2 + \delta^2}$  is the effective Rabi frequency.

**Interesting Special Cases** a) Pi-Puls  $\Omega_0\tau = \pi$ : swap population

$$|1\rangle \rightarrow i|2\rangle \quad (9.4.37)$$

$$|2\rangle \rightarrow i|1\rangle. \quad (9.4.38)$$

b) 2Pi-Puls  $\Omega_0\tau = 2\pi$ : flip the sign

c) Pi/2-Puls  $\Omega_0\tau = \pi/2$ : superposition state

## 9.5 Oscillating Dipoles

### Atomic Eigenstates

$$|\Psi_{nlm}(t)\rangle = e^{-iE_{nlm}t/\hbar}|\Psi_{nlm}(0)\rangle, \quad (9.5.1)$$

$$\hat{H}_0|\Psi_{nlm}(0)\rangle = E_{nlm}|\Psi_{nlm}\rangle, \quad (9.5.2)$$

and the electron density is

$$\rho(r, \theta, \phi) = |\Psi(r, \theta, \phi, t=0)|^2. \quad (9.5.3)$$

**Atomic Dipole** Calculate (Oscillating) Dipole Moment for Atomic Eigenstate. We denote  $|1\rangle = |\Psi_{nlm}\rangle$ . We have

$$d(t) = \langle 1(t) | \hat{d} | 1(t) \rangle \quad (9.5.4)$$

$$= \langle \hat{d} | 1 \rangle \quad (9.5.5)$$

$$= -e \langle 1 | \hat{r} | 1 \rangle. \quad (9.5.6)$$

Then,

$$-e \langle 1 | \hat{r} | 1 \rangle = -e \langle 1 | \hat{P} \hat{P}^{-1} \hat{r} \hat{P} \hat{P}^{-1} | 1 \rangle \quad (9.5.7)$$

$$= +e \langle 1 | \hat{r} | 1 \rangle, \quad (9.5.8)$$

which implies

$$\langle 1 | \hat{r} | 1 \rangle = 0. \quad (9.5.9)$$

**Atomic Dipole - Superposition States** Calculate (Oscillating) Dipole Moment for Atomic Superposition State

$$|\Psi(0)\rangle = \frac{1}{\sqrt{2}}(|1\rangle + i|2\rangle). \quad (9.5.10)$$

Evolution

$$|\Psi(t)\rangle = \frac{1}{\sqrt{2}}(|1\rangle + ie^{-i\omega_{21}t}|2\rangle). \quad (9.5.11)$$

We have

$$d(t) = \langle\Psi(t)|\hat{d}|\Psi(t)\rangle \quad (9.5.12)$$

$$= \frac{1}{2} \left\{ \langle 1|\hat{d}|1\rangle + \langle 2|\hat{d}|2\rangle + ie^{-i\omega_{21}t}\langle 1|\hat{d}|2\rangle - ie^{-i\omega_{21}t}\langle 2|\hat{d}|1\rangle \right\} \quad (9.5.13)$$

$$= d_{12}i\frac{1}{2} \left\{ e^{-i\omega_{21}t} - e^{i\omega_{21}t} \right\} \quad (9.5.14)$$

$$= d_{12} \sin(\omega_{21}t), \quad (9.5.15)$$

where  $d_{12}$  is the dipole moment amplitude,  $\omega_{21}$  is the natural resonance frequency.

**Electron Density - Superposition States** Calculate Electron Probability Density for Superposition State. The superposition state is

$$\Psi(r, t) = \frac{1}{\sqrt{2}} (\Psi_1(r) + ie^{-i\omega_{21}t}\Psi_2(r)). \quad (9.5.16)$$

The Electron Probability Density is

$$\rho(r, t) = |\Psi(r, t)|^2 \quad (9.5.17)$$

$$= \Psi^*\Psi \quad (9.5.18)$$

$$= \frac{1}{2} \left\{ |\Psi_1(r)|^2 + |\Psi_2(r)|^2 + 2\text{Re}(ie^{-i\omega_{21}t}\Psi_1^*(r)\Psi_2(r)) \right\}, \quad (9.5.19)$$

where  $2\text{Re}(ie^{-i\omega_{21}t}\Psi_1^*(r)\Psi_2(r))$  is the interference term.

**Examples** This is shown by animation and figure in the video.

## 9.6 The Bloch Sphere

### General Two-Level State

- General State Description

$$|\Psi\rangle = c'_1|1\rangle + c'_2|2\rangle \quad (9.6.1)$$

$$\text{Up to a global phase} \quad (9.6.2)$$

$$= |c'_1||1\rangle + e^{i\phi}|c'_2||2\rangle \quad (9.6.3)$$

satisfying  $|c'_1|^2 + |c'_2|^2 = 1$ .

- Alternative way

$$|\Psi\rangle = \cos(\theta/2)|1\rangle + e^{i\phi} \sin(\theta/2)|2\rangle, \quad (9.6.4)$$

since  $\cos^2(\theta/2) + \sin^2(\theta/2) = 1$ .

**Geometric Description - Bloch Sphere** We then have

$$|\Psi\rangle = \cos(\theta/2)|1\rangle + e^{i\phi} \sin(\theta/2)|2\rangle \quad (9.6.5)$$

with  $0 \leq \theta \leq \pi$  as the latitude and  $0 \leq \phi \leq 2\pi$  as the longitude. This is the Bloch Sphere representation. The definition of  $\theta$  and  $\phi$  and their ranges are different from my familiar coordinate system.

### Special States on Bloch Sphere

**Analogy to Spin -1/2 States** Is is shown in the figure.

## 9.7 Density Operator and Density Matrix

**The Problem** How do we describe "imperfect state preparation" in an experiment? For example, 50% $|1\rangle$  and 50% $|2\rangle$ . We may think of

$$|\Psi\rangle = \frac{1}{\sqrt{2}}(|1\rangle + |2\rangle).??? \quad (9.7.1)$$

This is 100% $|\Psi\rangle$  pure state. We need stable relative phase between the two states!

**Optical Analogy - Controlled Phase** The double slit problem is shown in the video.

Intensity on Detection Screen:

$$I \propto |E|^2 = |E_1 + e^{i\phi} E_2|^2 \quad (9.7.2)$$

$$= |E_1|^2 + |E_2|^2 + 2\text{Re}(E_1 E_2 e^{i\phi}). \quad (9.7.3)$$

As  $\phi$  varies, Interference pattern "washed out"!

We need new formalism to describe mixed states!(imperfect state preparation, spontaneous emission,...)

**Density Operator and Matrix** The description of mixed states can be handled by the density operator (matrix) formalism!

- Density operator (hermitian)

$$\hat{\rho} = \sum p_i |\Psi_i\rangle\langle\Psi_i| \quad (9.7.4)$$

$$\hat{\rho} = I \hat{\rho} I \quad (9.7.5)$$

$$= \sum_{i,j} |i\rangle\langle i| \hat{\rho} |j\rangle\langle j| \quad (9.7.6)$$

$$= \rho_{11}|1\rangle\langle 1| + \rho_{12}|1\rangle\langle 2| + \rho_{21}|2\rangle\langle 1| + \rho_{22}|2\rangle\langle 2|, \quad (9.7.7)$$

where  $I = \sum_i |i\rangle\langle i|$ .

- Density matrix

$$\rho = \begin{bmatrix} \rho_{11} & \rho_{12} \\ \rho_{21} & \rho_{22} \end{bmatrix}, \quad (9.7.8)$$

where  $\rho_{11}$  and  $\rho_{22}$  are the populations,  $\rho_{12}$  and  $\rho_{21}$  are the coherence. Since  $\rho$  is hermitian, we have

$$\rho_{12} = \rho_{21}^*. \quad (9.7.9)$$

**Example 9.7.1** (Example: Density Matrix of Pure State). *We have*

$$|\Psi\rangle = |c_1||1\rangle + e^{i\phi}|c_2||2\rangle. \quad (9.7.10)$$

The corresponding density operator of the **pure state** is  $\hat{\rho} = |\Psi\rangle\langle\Psi|$ . Then the corresponding density matrix is

$$\rho = \begin{bmatrix} |c_1|^2 & |c_1||c_2|e^{-i\phi} \\ |c_1||c_2|e^{i\phi} & |c_2|^2 \end{bmatrix}, \quad (9.7.11)$$

where  $|c_1||c_2|e^{-i\phi}$  and  $|c_1||c_2|e^{i\phi}$  are relative phase between states  $|1\rangle$  and  $|2\rangle$ .

specific example:

$$|\Psi\rangle = \frac{1}{\sqrt{2}}(|1\rangle + |2\rangle), \quad (9.7.12)$$

so

$$\rho = \begin{bmatrix} 1/2 & 1/2 \\ 1/2 & 1/2 \end{bmatrix}. \quad (9.7.13)$$

**Example 9.7.2** (Example: Fully Incoherent Mixture).

$$\hat{\rho} = \frac{1}{2}|1\rangle\langle 1| + \frac{1}{2}|2\rangle\langle 2| \quad (9.7.14)$$

with

$$\rho = \begin{bmatrix} 1/2 & 0 \\ 0 & 1/2 \end{bmatrix}, \quad (9.7.15)$$

where vanishingly coherence and the phase varies from 0 to  $2\pi$ . It means that we did not control phase.

### Useful Facts

- Expectation values:  $\langle \hat{A} \rangle = \text{Tr}(\hat{\rho}\hat{A}) = \text{Tr}(\rho A)$
- Time evolution (von Neumann equation)

$$i\hbar \frac{\partial \hat{\rho}}{\partial t} = [\hat{H}, \hat{\rho}] \quad (9.7.16)$$

- Pure state:  $\text{Tr}(\rho^2) = 1$
- Mixed states:  $\text{Tr}(\rho^2) < 1$

## 9.8 Optical Bloch Equations

**Time Evolution of Density Matrix** How to calculate time evolution of density matrix?

$$i\hbar \frac{\partial \hat{\rho}}{\partial t} = [\hat{H}, \hat{\rho}]. \quad (9.8.1)$$

Assume pure state

$$\frac{d}{dt}\rho_{11} = \frac{d}{dt}(c_1 c_1^*) \quad (9.8.2)$$

$$= \dot{c}_1 c_1^* + c_1 \dot{c}_1^* \quad (9.8.3)$$

$$= i \frac{\Omega_0}{2} \left( e^{i\delta t} \rho_{21} - e^{-i\delta t} \rho_{12} \right) \quad (9.8.4)$$

Transformation to rotatory frame of light

$$= i \frac{\Omega_0}{2} (\tilde{\rho}_{21} - \tilde{\rho}_{12}), \quad (9.8.5)$$

where

$$\dot{c}_1(t) = i \frac{\Omega_0}{2} e^{+i\delta t} c_2(t) \quad (9.8.6)$$

$$\dot{c}_2(t) = i \frac{\Omega_0}{2} e^{-i\delta t} c_1(t) \quad (9.8.7)$$

$$\tilde{\rho}_{12} = e^{-i\delta t} \rho_{12} \quad (9.8.8)$$

$$\tilde{\rho}_{21} = e^{+i\delta t} \rho_{21}. \quad (9.8.9)$$

Other elements obtained in analogy!

$$\frac{d}{dt} \rho_{11} = i \frac{\Omega_0}{2} (\tilde{\rho}_{21} - \tilde{\rho}_{12}) \quad (9.8.10)$$

$$\frac{d}{dt} \rho_{22} = i \frac{\Omega_0}{2} (\tilde{\rho}_{12} - \tilde{\rho}_{21}) \quad (9.8.11)$$

$$\frac{d}{dt} \tilde{\rho}_{12} = -i\delta \tilde{\rho}_{12} + i \frac{\Omega_0}{2} (\rho_{22} - \rho_{11}) \quad (9.8.12)$$

$$\frac{d}{dt} \tilde{\rho}_{21} = +i\delta \tilde{\rho}_{21} + i \frac{\Omega_0}{2} (\rho_{11} - \rho_{22}). \quad (9.8.13)$$

Noting that  $\tilde{\rho}_{12} = \tilde{\rho}_{21}$  due to hermitian matrix, the third and the forth equations are the same. So we have

$$\frac{d}{dt} \rho_{11} = i \frac{\Omega_0}{2} (\tilde{\rho}_{21} - \tilde{\rho}_{12}) \quad (9.8.14)$$

$$\frac{d}{dt} \rho_{22} = i \frac{\Omega_0}{2} (\tilde{\rho}_{12} - \tilde{\rho}_{21}) \quad (9.8.15)$$

$$\frac{d}{dt} \tilde{\rho}_{12} = -i\delta \tilde{\rho}_{12} + i \frac{\Omega_0}{2} (\rho_{22} - \rho_{11}). \quad (9.8.16)$$

**Optical Bloch Equations with Damping** Phenomenological damping and spontaneous emission in the figure. Combine the decay, we have

$$\frac{d}{dt} \rho_{11} = i \frac{\Omega_0}{2} (\tilde{\rho}_{21} - \tilde{\rho}_{12}) + \gamma \rho_{22} \quad (9.8.17)$$

$$\frac{d}{dt} \rho_{22} = i \frac{\Omega_0}{2} (\tilde{\rho}_{12} - \tilde{\rho}_{21}) - \gamma \rho_{22} \quad (9.8.18)$$

$$\frac{d}{dt} \tilde{\rho}_{12} = -i\delta \tilde{\rho}_{12} + i \frac{\Omega_0}{2} (\rho_{22} - \rho_{11}) - (\gamma/2) \tilde{\rho}_{12}. \quad (9.8.19)$$

We now define the inversion  $w = \rho_{22} - \rho_{11}$ . We have Optical Bloch Equations with Damping

$$\frac{d}{dt} \tilde{\rho}_{21} = -(\gamma/2 - i\delta) \tilde{\rho}_{21} - \frac{i\omega\Omega_0}{2} \quad (9.8.20)$$

$$\frac{d}{dt} \omega = -\gamma(\omega + 1) - i\Omega_0 (\tilde{\rho}_{21} - \tilde{\rho}_{12}) \quad (9.8.21)$$

in the Density Matrix Form.

## 9.9 Optical Bloch Equations - Dynamics and Steady State

**Dynamical Evolution of System** Shown in the figure in the picture.

**Steady State Solution** Conditions:  $\frac{d}{dt}\tilde{\rho}_{21} = 0$  and  $\frac{d}{dt}\omega = 0$ . Then we have the solutions

$$\omega = -\frac{1}{1+S} \quad (9.9.1)$$

$$\tilde{\rho}_{21} = \frac{2\Omega_0}{2(\gamma/2 - \delta)(1+S)} \quad (9.9.2)$$

$$S = \frac{\Omega_0^2/2}{\delta^2 + \gamma^2/4} = \frac{S_0}{1 + 4\delta^2/\gamma^2} \quad (9.9.3)$$

$$S_0 = \frac{2\Omega_0^2}{\gamma^2} = \frac{I}{O_{sat}}, \quad (9.9.4)$$

where  $S$  is called the saturation parameter,  $S_0$  is called resonant saturation parameter.

Limiting Cases:

- $S \leq 1$ :  $w \rightarrow -1$  where  $w = \rho_{22} - \rho_{11}$ . Atom is mainly in ground state.
- $S \geq 1$ :  $S \rightarrow \infty$ ,  $w \rightarrow 0$ .
- Excited State Population:

$$\rho_{22} \quad (9.9.5)$$

$$\begin{aligned} & \text{Combine with } \rho_{22} + \rho_{11} = 1 \\ &= \frac{1}{2}(1+w) \end{aligned} \quad (9.9.6)$$

$$= \frac{S}{2(1+S)} \quad (9.9.7)$$

$$= \frac{S_0/2}{1 + S_0 + 4\delta^2/\gamma^2} \quad (9.9.8)$$

$$\xrightarrow[S_0 \rightarrow \infty; \delta=0]{} \frac{1}{2}. \quad (9.9.9)$$

- Photon Scattering Rate:  $\Gamma_{ph} = \gamma\rho_{22} = \frac{\gamma}{2} \frac{S_0}{1 + S_0 + 4\delta^2/\gamma^2}$ .  $\Gamma_{ph} \rightarrow \gamma/2$  for  $S_0 \rightarrow \infty$  and  $\delta = 0$ . We rewrite it as

$$\Gamma_{ph} = \left( \frac{S_0}{1 + S_0} \right) \left( \frac{\gamma/2}{1 + 4\delta^2/\gamma^2} \right) \quad (9.9.10)$$

$$\gamma' = \gamma\sqrt{1 + S_0}. \quad (9.9.11)$$

It has a figure in the video. The saturation broadening is shown in the figure.

## 9.10 Lambert-Beer Law

**Attenuation of Light** It is shown in the figure.

**Scattered Light from Slab of Atoms** scattered light power by slab of length  $dz$

$$dP_{sc} = \Gamma_{ph} \times nAdz \times \hbar\omega, \quad (9.10.1)$$

where  $\Gamma_{ph}$  is the single atom photon scattering rate,  $\hbar\omega$  is the energy of single atom,  $nAdz$  is the number of atoms. Then we have

$$\frac{dP_{sc}}{dz} = \Gamma_{ph} \times nA \times \hbar\omega. \quad (9.10.2)$$

**Scattered Light from Slab of Atoms** Energy conservation requires

$$\frac{dP}{dz} = -\frac{dP_{sc}}{dz} \quad (9.10.3)$$

$$\frac{dP}{dz} = \frac{dI}{dI} A. \quad (9.10.4)$$

Put every thing together:

$$\frac{dI}{dz} = -\Gamma n \hbar \omega. \quad (9.10.5)$$

We have

$$\frac{dI(z)}{dz} = -n\sigma I(z), \quad (9.10.6)$$

where  $\sigma$  is the atomic scattering cross section.

**Lambert-Beer Law (no saturation)** We compute the solutions

$$I(z) = I(0)e^{-n\sigma z}, \quad (9.10.7)$$

which is the Lambert-Beer Law of Absorption.

**Laser induced Fluorescence** Shown in a video.

## 9.11 Bloch Vector

**Density Matrix Revisited** Density Matrix of TLA

$$\rho = \begin{pmatrix} \rho_{11} & \rho_{12} \\ \rho_{21} & \rho_{22} \end{pmatrix} \quad (9.11.1)$$

Density Matrix hermitian

$$\rho = \rho^\dagger = (\rho^T)^*, \quad (9.11.2)$$

so we have

$$\rho = \begin{pmatrix} \rho_{11} & \text{Re}\rho_{12} + i\text{Im}\rho_{12} \\ \text{Re}\rho_{12} - i\text{Im}\rho_{12} & \rho_{22} \end{pmatrix}. \quad (9.11.3)$$

Pauli matrices are

$$\sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}. \quad (9.11.4)$$

The decomposition of Density matrix into Pauli matrices

$$\rho = \frac{1}{2} (I + b_x \sigma_x + b_y \sigma_y + b_z \sigma_z), \quad (9.11.5)$$

where  $b_x, b_y, b_z \in \mathbb{R}$ .

**Bloch Vector** We have the density matrix in rotating frame of light

$$\tilde{\rho} = \begin{pmatrix} \rho_{11} & \tilde{\rho}_{12} \\ \tilde{\rho}_{21} & \rho_{22} \end{pmatrix}, \quad (9.11.6)$$

where  $\tilde{\rho}_{12} = \rho_{12}e^{-i\omega t}$ . We use following sign convention and have

$$\tilde{\rho} = \frac{1}{2} (I + u\sigma_x - v\sigma_y - w\sigma_z), \quad (9.11.7)$$

and the bloch vector is defined as

$$\begin{pmatrix} u \\ v \\ w \end{pmatrix}. \quad (9.11.8)$$

It can be easily shown that

$$u = 2\text{Re}(\tilde{\rho}_{12}) = \tilde{\rho}_{12} + \tilde{\rho}_{12}^* \quad (9.11.9)$$

$$v = 2\text{Im}(\tilde{\rho}_{12}) = i(\tilde{\rho}_{12}^* - \tilde{\rho}_{12}) \quad (9.11.10)$$

$$w = \rho_{22} - \rho_{11}, \quad (9.11.11)$$

$$(9.11.12)$$

where  $u$  is the dispersive component,  $v$  is the absorption component and  $w$  is the inversion.

Bloch vector can be used to describe any state of TLA density matrix!

Properties of Bloch Vector

- Mixed State:  $u^2 + v^2 + w^2 < 1$
- Pure State:  $u^2 + v^2 + w^2 = 1$

## 9.12 Understanding Bloch Vector

What physical behaviour do the components stand for?

- $w = -1$  atom in ground state.  $w = +1$  atom in excited state.
- What about  $u, v$ ?

$$\langle \hat{d}_i(t) \rangle = \text{Tr}(\hat{\rho}\hat{d}) \quad (9.12.1)$$

$$= \text{Tr} \left[ \begin{pmatrix} \rho_{11} & \rho_{12} \\ \rho_{12}^* & \rho_{22} \end{pmatrix} \begin{pmatrix} 0 & d_{12}^i \\ d_{12}^i & 0 \end{pmatrix} \right], \quad (9.12.2)$$

where  $d_{12}^x = \langle 1| -q\hat{x}|2\rangle$ .

Written in the vector form, we have

$$\langle \hat{d} \rangle(t) = d_{12} (\rho_{12} + \rho_{12}^*) \quad (9.12.3)$$

$$= d_{12} (\tilde{\rho}_{12} e^{i\omega t} + \tilde{\rho}_{12}^* e^{-i\omega t}) \quad (9.12.4)$$

$$= d_{12} [u \cos(\omega t) - v \sin(\omega t)], \quad (9.12.5)$$

where we use  $\rho_{12} = \tilde{\rho}_{12} e^{i\omega t}$ ,  $u$  denotes in phase and  $v$  denotes  $90^\circ$  out of phase component.

Reminder:  $E(t) = \epsilon E_0 \cos(\omega t)$ .

- Which component responsible for absorption/emission? We have a figure in the video to show the classical picture.

Average absorbed power per atom (classical ensemble average)

$$\langle \frac{dW}{dt} \rangle = \epsilon E_0 \cos(\omega t) \langle -q \frac{dr}{dt} \rangle \quad (9.12.6)$$

$$= \epsilon E_0 \cos(\omega t) \langle \dot{d} \rangle. \quad (9.12.7)$$

Quantum mechanical analogue (Ehrenfest)

$$\langle \frac{dW}{dt} \rangle = \epsilon E_0 \cos(\omega t) \langle \dot{d} \rangle \quad (9.12.8)$$

$$\langle \hat{d} \rangle(t) = d_{12} [u \cos(\omega t) - v \sin(\omega t)]. \quad (9.12.9)$$

$$\langle \frac{dW}{dt} \rangle = -d_{12} \cdot \epsilon E_0 \omega (u \cos(\omega t) \sin(\omega t) + v \sin(\omega t)^2) \quad (9.12.10)$$

$$\langle \frac{dW}{dt} \rangle = \frac{1}{T} \int dt \langle \frac{dW}{dt} \rangle \quad (9.12.11)$$

$$= -\frac{d_{12} \cdot \epsilon E_0 \omega v}{2} \quad (9.12.12)$$

$$= -\hbar \frac{d_{12} \epsilon E_0}{\hbar} \omega \frac{v}{2} \quad (9.12.13)$$

$$= -\hbar \Omega_0 \omega \frac{v}{2}, \quad (9.12.14)$$

which is the absorption.

## 9.13 Optical Bloch Equations using Bloch Vector

## 9.14 Interlude: The Mach-Zehnder Interferometer

## 9.15 Ramsey Interferometer

## 9.16 Review: QM of the Harmonic Oscillator

[SZQ: 2023.04.20: I have understand the content in this video.]

## 9.17 Wave equation and energy density of classical radiation field

This section is also known as the review of Maxwell equations vector potentials.

**Fundamentals** Maxwell equations in free space

$$\nabla \cdot E = 0, \nabla \times E = -\frac{\partial B}{\partial t} \quad (9.17.1)$$

$$\nabla \cdot B = 0, \nabla \times B = \frac{1}{c^2} \frac{\partial E}{\partial t}. \quad (9.17.2)$$

**Lemma 9.17.1** (Coulomb Gauge). *Considering Coulomb Gauge, we have*

$$\nabla \cdot A = 0. \quad (9.17.3)$$

*Then we can express the electric field and magnetic field in terms of the vector potential*

$$B(r, t) = \nabla \times A(r, t) \quad (9.17.4)$$

$$E(r, t) = -\frac{\partial A(r, t)}{\partial t}. \quad (9.17.5)$$

**Lemma 9.17.2** (Wave equation). *Considering Coulomb Gauge, the wave equation is*

$$\nabla^2 A - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} A = 0. \quad (9.17.6)$$

*Proof.* Using 9.17.1 and the forth equation in 9.17.1, we have

$$\nabla \times B = \nabla \times (\nabla \times A(r, t)), \quad (9.17.7)$$

and

$$\nabla \times B = \frac{1}{c^2} \frac{\partial E}{\partial t} \quad (9.17.8)$$

$$= -\frac{1}{c^2} \frac{\partial^2}{\partial t^2} A(r, t). \quad (9.17.9)$$

So we have

$$\nabla \times (\nabla \times A) = -\frac{1}{c^2} \frac{\partial^2}{\partial t^2}. \quad (9.17.10)$$

Then use the rule in vector Calculus

$$\nabla \times (\nabla \times A) = \nabla(\nabla \cdot A) - \Delta A. \quad (9.17.11)$$

Use lemma 9.17.1, we then have

$$-\Delta A = -\frac{1}{c^2} \frac{\partial^2}{\partial t^2} A \quad (9.17.12)$$

$$= -\nabla^2 A. \quad (9.17.13)$$

So we have

$$\nabla^2 A - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} A = 0. \quad (9.17.14)$$

□

### Solutions of Wave Equation

**Lemma 9.17.3** (Solutions of Wave Equation). *Plane waves:*

$$A_{k,\alpha} = \epsilon_{k,\alpha} A_{k,\alpha} \exp[i(\mathbf{k}\mathbf{r} - \omega_k t)], \quad (9.17.15)$$

where  $\epsilon_{k,\alpha}$  is polarization,  $A_{k,\alpha}$  is complex amplitude,  $|k| = \frac{2\pi}{\lambda}$  is wavenumber i.e., the magnitude of the wave vector,  $\mathbf{k}$  is the wave vector,  $\omega_k = ck$ .

Which wave vectors are possible? (a). in finite space,  $\mathbf{k}$  distributed continuous; (b). finite box of length  $L$ ,  $\mathbf{k}$  distributed discretely (periodic boundary conditions)

$$k_x = \frac{2\pi}{L} n_x, k_y = \frac{2\pi}{L} n_y, k_z = \frac{2\pi}{L} n_z \quad (9.17.16)$$

$$\mathbf{A}(\mathbf{r}, t) = \sum_{\mathbf{k}, \alpha} \epsilon_{\mathbf{k}, \alpha} (A_{\mathbf{k}, \alpha} \exp[i(\mathbf{k}\mathbf{r} - \omega_k t)] + A_{\mathbf{k}, \alpha}^* \exp[-i(\mathbf{k}\mathbf{r} - \omega_k t)]) \quad (9.17.17)$$

$$\mathbf{E}(\mathbf{r}, t) = -\frac{\partial}{\partial t} \mathbf{A}(r, t) = \sum_{\mathbf{k}, \alpha} \epsilon_{\mathbf{k}, \alpha} i\omega_k [A_{\mathbf{k}, \alpha} \exp[i(\mathbf{k}\mathbf{r} - \omega_k t)] - A_{\mathbf{k}, \alpha}^* \exp[-i(\mathbf{k}\mathbf{r} - \omega_k t)]] \quad (9.17.18)$$

$$\mathbf{B}(\mathbf{r}, t) = \sum_{\mathbf{k}, \alpha} i(\mathbf{k} \times \epsilon_{\mathbf{k}, \alpha}) [A_{\mathbf{k}, \alpha} \exp[i(\mathbf{k}\mathbf{r} - \omega_k t)] - A_{\mathbf{k}, \alpha}^* \exp[-i(\mathbf{k}\mathbf{r} - \omega_k t)]]. \quad (9.17.19)$$

[SZQ: 2023.04.20: The complex conjugate term is used to eliminate the imaginary part.]

**Total Energy of Radiation Field** Total energy of radiation field in volume  $V = L^3$ .

[SZQ: 2023.04.20: The total energy is the integration of the electric density and magenetic density over the volume.]

The electric density is

$$\frac{1}{2}\varepsilon_0 E(r, t)^2. \quad (9.17.20)$$

The magenetic density is

$$\frac{1}{2\mu_0} B(r, t)^2. \quad (9.17.21)$$

Then the total energy of radiation field in volume  $V = L^3$  is

$$H = \frac{1}{2} \int_V dV \left[ \varepsilon_0 E(r, t)^2 + \frac{1}{\mu_0} B(r, t)^2 \right] \quad (9.17.22)$$

$$= \sum_{\mathbf{k}, \alpha} \varepsilon_0 V \omega_k^2 [A_{\mathbf{k}, \alpha} A_{\mathbf{k}, \alpha}^* + A_{\mathbf{k}, \alpha}^* A_{\mathbf{k}, \alpha}] \quad (9.17.23)$$

$$= \sum_{\mathbf{k}, \alpha} E_{\mathbf{k}, \alpha}, \quad (9.17.24)$$

where

$$E_{\mathbf{k}, \alpha} = \varepsilon_0 V \omega_k^2 [A_{\mathbf{k}, \alpha} A_{\mathbf{k}, \alpha}^* + A_{\mathbf{k}, \alpha}^* A_{\mathbf{k}, \alpha}]. \quad (9.17.25)$$

[SZQ: 2023.04.20: This expression is similar to the quantum harmonic oscillators.] [SZQ: 2023.04.20: I ignore the bold. So you should understand where you should use the bold.]

## 9.18 Quantization of the e.m. field

**Fundamental Idea** RadiationMode ( $k, \alpha$ )

- **To every radiation mode, we associate a harmonic oscillator!** Creation and annihilation operators can change the degree of excitation of mode (occupation with photons)
- **A photon is an excitation quantum of the harmonic oscillator associated with a mode!**

**Creation and Annihilation Operators**  $\hat{a}_k |n_k\rangle = \sqrt{n_k} |n_k - 1\rangle$ : decrease photon number by one photon.

$\hat{a}_k^\dagger |n_k\rangle = \sqrt{n_k + 1} |n_k + 1\rangle$ : increase photon number by one photon.

**Number operator:**  $\hat{n}_k |n_k\rangle = n_k |n_k\rangle$ .  $\hat{a}_k^\dagger \hat{a}_k = \hat{n}_k$ .

**Fock state:**  $|n_k\rangle$ . Fock state is the eigenstate of quantum harmonic oscillator.

**Hamiltonian of Radiation Field** The Hamiltonian of Radiation Field is the sum of the hamitonian of harmonic oscillator of each mode as

$$\hat{H}_R = \sum_k \hat{H}_k, \quad (9.18.1)$$

where

$$\hat{H}_k = \frac{1}{2} \hbar \omega_k (\hat{a}_k \hat{a}_k^\dagger + \hat{a}_k^\dagger \hat{a}_k). \quad (9.18.2)$$

We can compare it with classall expression

$$E_{k,\alpha} = \epsilon_0 V \omega_k^2 (A_{k,\alpha} A_{k,\alpha}^* + A_{k,\alpha}^* A_{k,\alpha}). \quad (9.18.3)$$

If we replace  $A_k$  with

$$A_k = \sqrt{\frac{\hbar}{2\epsilon_0 V \omega_k}} \hat{a}_k, \quad (9.18.4)$$

and replace  $A_k^*$  with

$$A_k^* = \sqrt{\frac{\hbar}{2\epsilon_0 V \omega_k}} \hat{a}_k^\dagger. \quad (9.18.5)$$

We will arrive at  $\hat{H}_k$ . Also we can obtain the quantum version of vector potential operator. The classical vector potential operator is

$$A_k(r, t) = \epsilon_k [A_k \exp[i(kr - \omega_k t)] + A_k^* \exp[-i(kr - \omega_k t)]]. \quad (9.18.6)$$

The quantum version will be

$$\hat{A}_k(r, t) = \epsilon_k \sqrt{\frac{\hbar}{2\epsilon_0 V \omega_k}} [\hat{a}_k \exp[i(kr - \omega_k t)] + \hat{a}_k^\dagger \exp[-i(kr - \omega_k t)]]. \quad (9.18.7)$$

Use the quantum vector potential, we can derive the quantum electric field operator as

$$\hat{E}_k(r, t) = -\frac{\partial}{\partial t} \hat{A}_k(r, t) \quad (9.18.8)$$

$$= -\epsilon_k \sqrt{\frac{\hbar}{2\epsilon_0 \omega_k V}} (-\omega_k) [i \hat{a}_k \exp[i(kr - \omega_k t)] - i \hat{a}_k^\dagger \exp[-i(kr - \omega_k t)]]. \quad (9.18.9)$$

Recall that  $i = \exp[i\pi/2]$  and define

$$\chi_k(r, t) = -kr + \omega_k t - \pi/2. \quad (9.18.10)$$

We then have the compact form

$$\hat{E}(r, t) = \sum_k \epsilon_k \sqrt{\frac{\hbar \omega_k}{2\epsilon_0 V}} [\hat{a}_k \exp[-i\chi_k(r, t)] + \hat{a}_k^\dagger \exp[i\chi_k(r, t)]] \quad (9.18.11)$$

$$= \sum_k \hat{E}_k(r, t) \quad (9.18.12)$$

$$:= \hat{E}^+(r, t) + \hat{E}^-(r, t). \quad (9.18.13)$$

**Hamiltonian of Radiation Field** The Hamiltonian of Radiation Field is

$$\hat{H}_R = \frac{1}{2} \int_V dV \left[ \epsilon_0 \hat{E} \cdot \hat{E} + \frac{1}{\mu_0} \hat{B} \cdot \hat{B} \right] \quad (9.18.14)$$

$\hat{B}, \hat{E}$  are the quantum operator of  $B, E$

$$= \sum_k \frac{\hbar\omega_k}{2} \left[ \hat{a}_k^\dagger \hat{a}_k + \hat{a}_k \hat{a}_k^\dagger \right] \quad (9.18.15)$$

Use the commutation relation

$$= \sum_k \left( \hat{a}_k^\dagger \hat{a}_k + 1/2 \right). \quad (9.18.16)$$

Use this hamiltonian, we derive the energy of multi-mode Fock states as

$$\hat{H}_R |n_{k_1}, n_{k_2}, \dots\rangle = \sum_k \hbar\omega_k \left( n_k + \frac{1}{2} \right) |n_{k_1}, n_{k_2}, \dots\rangle \quad (9.18.17)$$

using the fact that  $\hat{a}_k^\dagger \hat{a}_k$  is the number operator  $\hat{n}_k$ .

Also the vacuum state energy will be

$$E_0 = \sum_k \frac{1}{2} \hbar\omega_k \quad (9.18.18)$$

corresponds to

$$|0\rangle = |0\rangle \otimes \dots \otimes |0\rangle. \quad (9.18.19)$$

This is divergent, but do not worry. When we calculate the difference, this term will be canceled.

## 9.19 Field state of single radiation field mode: Fock States

We focus discussion on a **single mode of the radiation field (wave vector  $k$ )**

We define the phase factor

$$\chi = \chi_k(r, t) = \omega_k t - \mathbf{k}\mathbf{r} - \pi/2. \quad (9.19.1)$$

Then we have

$$\hat{E}(\chi) = \hat{E}^+(\chi) + \hat{E}^-(-\chi) \quad (9.19.2)$$

$$= \left( \frac{\hbar\omega}{2\varepsilon_0 V} \right)^{1/2} (\hat{a} \exp[-i\chi] + \hat{a}^\dagger \exp[i\chi]). \quad (9.19.3)$$

We write the field operator in natural units  $2 \left( \frac{\hbar\omega}{2\varepsilon_0 V} \right)^{1/2}$ , which is also called vacuum field strength. We then have

$$\hat{E}(\chi) = \frac{1}{2} (\hat{a} \exp[-i\chi] - \hat{a}^\dagger \exp[i\chi]). \quad (9.19.4)$$

**Fock states:**  $|n\rangle$  means  $n$  photons in radiation mode, also means eigenstate of number operator  $\hat{n}$ .

$$\hat{H}|n\rangle = \hbar\omega(\hat{a}^\dagger\hat{a} + \frac{1}{2})|n\rangle \quad (9.19.5)$$

$$= \hbar\omega(\hat{n} + \frac{1}{2})|n\rangle \quad (9.19.6)$$

$$= \hbar\omega(n + \frac{1}{2})|n\rangle. \quad (9.19.7)$$

**Lemma 9.19.1.** *Fluctuations in  $n$  is*

$$(\Delta n)^2 = 0. \quad (9.19.8)$$

*Proof.*

$$(\Delta n)^2 = \langle n|(\hat{n} - \langle \hat{n} \rangle)^2|n\rangle \quad (9.19.9)$$

$$= \langle n|(\hat{n}^2 - 2\langle \hat{n} \rangle \hat{n} + \langle \hat{n} \rangle^2)|n\rangle \quad (9.19.10)$$

$$= \langle n|\hat{n}^2|n\rangle - 2\langle \hat{n} \rangle^2 + \langle \hat{n} \rangle^2 \quad (9.19.11)$$

$$= \langle n|\hat{n}^2|n\rangle - \langle n|\hat{n}|n\rangle^2 \quad (9.19.12)$$

$$= \langle n|n^2|n\rangle - \langle n|n|n\rangle^2 \quad (9.19.13)$$

$$= n^2 - n^2 \quad (9.19.14)$$

$$= 0, \quad (9.19.15)$$

where we use

$$\langle n|\hat{n}|n\rangle = \langle \hat{n} \rangle. \quad (9.19.16)$$

□

**Lemma 9.19.2.** *The expectation value of the field is*

$$E = \langle n|\hat{E}(\chi)|n\rangle = 0. \quad (9.19.17)$$

*Proof.*

$$E = \langle n|\hat{E}(\chi)|n\rangle \quad (9.19.18)$$

$$= \frac{1}{2}\langle n|\hat{a} \exp[-i\chi] + \hat{a}^\dagger \exp[i\chi]|n\rangle \quad (9.19.19)$$

$$= \frac{1}{2}\langle n|\hat{a} \exp[-i\chi]|n\rangle + \frac{1}{2}\langle n|\hat{a}^\dagger \exp[i\chi]|n\rangle \quad (9.19.20)$$

$$= 0 + 0 \quad (9.19.21)$$

$$= 0. \quad (9.19.22)$$

□

**Lemma 9.19.3.** *Field fluctuations is*

$$(\Delta E(\chi))^2 = \frac{1}{2}(n + \frac{1}{2}). \quad (9.19.23)$$

*Proof.*

$$(\Delta E(\chi))^2 = \langle n | \hat{E}(\chi)^2 | n \rangle - \langle n | \hat{E}(\chi) | n \rangle^2 \quad (9.19.24)$$

$$= \langle n | \hat{E}(\chi)^2 | n \rangle - 0 \quad (9.19.25)$$

$$= \frac{1}{4} \langle n | (\hat{a} \exp[-i\chi] + \hat{a}^\dagger \exp[i\chi])^2 | n \rangle \quad (9.19.26)$$

$$= \frac{1}{4} \langle n | \hat{a}^2 \exp[-2i\chi] + \hat{a}^{\dagger 2} \exp[2i\chi] + \hat{a}\hat{a}^\dagger + \hat{a}^\dagger\hat{a} | n \rangle \quad (9.19.27)$$

$$= \frac{1}{4} \langle n | 2\hat{n} + 2 | n \rangle \quad (9.19.28)$$

$$= \frac{1}{4} \langle n | 2n + 1 | n \rangle \quad (9.19.29)$$

$$= \frac{1}{2}(n + \frac{1}{2}). \quad (9.19.30)$$

□

When  $n = 0$ , which is vacuum state, we have the standard deviation as  $1/2$ , which is half of the unit, i.e., vacuum field strength.

## 9.20 Field state of single radiation field mode: Coherent States

**How to reproduce classical motion** Superposition of Fock states reproduces oscillating wavepacket motion!

$$|\alpha\rangle \propto \exp[-i\frac{1}{2}\omega t]|0\rangle + \exp[-i\frac{3}{2}\omega t]\alpha|1\rangle + \exp[-i\frac{5}{2}\omega t]\frac{\alpha^2}{\sqrt{2}}|2\rangle + \dots \quad (9.20.1)$$

$$= \sum_{n=0}^{\infty} \frac{\alpha^n}{\sqrt{n!}} \exp[-iE_n t/\hbar]. \quad (9.20.2)$$

**What we want:** states of light, whose expectation value corresponds to classical e.m. waves!

**Solotion:** Coherent States

**Definition 9.20.1.** *Coherent States:*

$$|\alpha\rangle = \exp\left[-\frac{1}{2}|\alpha|^2\right] \sum_{n=0}^{\infty} \frac{\alpha^n}{\sqrt{n!}} |n\rangle, \quad (9.20.3)$$

where  $\alpha$  is complex number, amplitude of coherent state.  $\alpha = |\alpha| \exp[i\theta]$ ,  $|n\rangle$  is the fock state.

**Lemma 9.20.1.** *Coherent state is normalized:  $\langle \alpha | \alpha \rangle = 1$ .*

*Proof.*

$$\langle \alpha | \alpha \rangle = \exp[-|\alpha|^2] \sum_{n,n'} \frac{(\alpha^*)^n \alpha^{n'}}{\sqrt{n!} \sqrt{n'!}} \langle n | n' \rangle \quad (9.20.4)$$

$$= \exp[-|\alpha|^2] \sum_{n=0}^{\infty} \frac{|\alpha|^{2n}}{n!} \quad (9.20.5)$$

$$= \exp[-|\alpha|^2] \exp[+|\alpha|^2] \quad (9.20.6)$$

$$= 1. \quad (9.20.7)$$

□

**Lemma 9.20.2.** *Coherent state is quasi orthogonal:  $|\alpha - \beta| >> 1 \rightarrow |\langle \alpha | \alpha \rangle| \rightarrow 0$ .*

*Proof.*

$$\langle \alpha | \beta \rangle = \exp \left[ -\frac{1}{2} |\alpha|^2 - \frac{1}{2} |\beta|^2 + \alpha^* \beta \right]. \quad (9.20.8)$$

$$|\langle \alpha | \beta \rangle|^2 = \langle \alpha | \beta \rangle^* \langle \alpha | \beta \rangle \quad (9.20.9)$$

$$= \exp \left[ -|\alpha|^2 - |\beta|^2 + \alpha^* \beta + \beta^* \alpha \right] \quad (9.20.10)$$

$$= \exp[-|\alpha - \beta|^2]. \quad (9.20.11)$$

We have  $|\alpha - \beta| >> 1 \rightarrow |\langle \alpha | \alpha \rangle| \rightarrow 0$ , which is called quasi-orthogonal. □

**Lemma 9.20.3.** *Coherent states are eigenstates of destruction operator  $\hat{a} | \alpha \rangle = \alpha | \alpha \rangle$ .*

*Proof.*

$$\hat{a} | \alpha \rangle = \exp \left[ -\frac{1}{2} |\alpha|^2 \right] \sum_{n=0}^{\infty} \frac{\alpha^n}{\sqrt{n!}} |n\rangle \quad (9.20.12)$$

$$= \exp \left[ -\frac{1}{2} |\alpha|^2 \right] \sum_{n=0}^{\infty} \frac{\alpha^n}{\sqrt{n!}} \hat{a} |n\rangle \quad (9.20.13)$$

$$= \exp \left[ -\frac{1}{2} |\alpha|^2 \right] \sum_{n=1}^{\infty} \frac{\alpha^n}{\sqrt{n!}} \sqrt{n} |n-1\rangle \quad (9.20.14)$$

$$= \alpha \exp \left[ -\frac{1}{2} |\alpha|^2 \right] \sum_{n=1}^{\infty} \frac{\alpha^{n-1}}{\sqrt{(n-1)!}} |n-1\rangle \quad (9.20.15)$$

k := n-1

$$= \alpha \exp \left[ -\frac{1}{2} |\alpha|^2 \right] \sum_{k=0}^{\infty} \frac{\alpha^k}{\sqrt{k!}} |k\rangle \quad (9.20.16)$$

$$= \alpha | \alpha \rangle. \quad (9.20.17)$$

□

**Lemma 9.20.4.** *The average photon number of coherent states:  $\bar{n} = |\alpha|^2$ .*

*Proof.*

$$\bar{n} = \langle \alpha | \hat{n} | \alpha \rangle \quad (9.20.18)$$

$$= \langle \alpha | \hat{a}^\dagger \hat{a} | \alpha \rangle \quad (9.20.19)$$

Use 9.20.3

$$= \langle \alpha^* \alpha | \alpha \rangle \quad (9.20.20)$$

$$= |\alpha|^2 \langle \alpha | \alpha \rangle \quad (9.20.21)$$

$$= |\alpha|^2. \quad (9.20.22)$$

□

[SZQ: 2023.04.20: coherent state is robust!]

**Lemma 9.20.5.** *Photon number variance of coherent states is*

$$(\Delta n)^2 = \bar{n}. \quad (9.20.23)$$

*Proof.*

$$(\Delta n)^2 = \langle \alpha | \hat{n}^2 | \alpha \rangle - \langle \alpha | \hat{n} | \alpha \rangle^2 \quad (9.20.24)$$

$$= \langle \hat{a}^\dagger \hat{a} \hat{a}^\dagger \hat{a} | \alpha \rangle - |\alpha|^4 \quad (9.20.25)$$

$$\hat{a} \hat{a}^\dagger = \hat{a}^\dagger \hat{a} + 1$$

$$= \langle \alpha | \hat{a}^\dagger \hat{a}^\dagger \hat{a} \hat{a} | \alpha \rangle + \langle \hat{a}^\dagger \hat{a} | \alpha \rangle - |\alpha|^4 \quad (9.20.26)$$

$$= |\alpha|^2 \quad (9.20.27)$$

$$= \bar{n}. \quad (9.20.28)$$

□

**Lemma 9.20.6.** *The photon number distribution is*

$$P(n) = |\langle n | \alpha \rangle|^2 \quad (9.20.29)$$

$$= |\langle n | \exp \left[ -\frac{|\alpha|^2}{2} \right] \sum_{n'=0}^{\infty} \frac{\alpha^{n'}}{\sqrt{n!}} |n' \rangle|^2 \quad (9.20.30)$$

$$= |\exp \left[ -\frac{|\alpha|^2}{2} \right] \sum_{n'=0}^{\infty} \frac{\alpha^{n'}}{\sqrt{n!}} \langle n | n' \rangle|^2 \quad (9.20.31)$$

$$= |\exp \left[ -\frac{|\alpha|^2}{2} \right] \frac{\alpha^n}{\sqrt{n!}}|^2 \quad (9.20.32)$$

$$= \exp[-|\alpha|^2] \frac{|\alpha|^{2n}}{n!} \quad (9.20.33)$$

*Use lemma 9.20.4*

$$= \exp[-\bar{n}] \frac{\bar{n}^n}{n!}. \quad (9.20.34)$$

This is known as **Poisson distribution**. The standard deviation is

$$\Delta n = \sqrt{\bar{n}}. \quad (9.20.35)$$

The standard deviation relative to the mean is

$$\frac{\Delta n}{\bar{n}} = \frac{1}{\sqrt{\bar{n}}}, \quad (9.20.36)$$

which shows that the fluctuations is smaller and smaller when the mean becomes larger and larger.

For large  $\bar{n}$ , we have

$$P(n) \simeq \frac{1}{\sqrt{2\pi\bar{n}}} \exp \left[ -\frac{1}{2} \frac{(n - \bar{n})^2}{\bar{n}} \right]. \quad (9.20.37)$$

[SZQ:  $\simeq$  means asymptotically equal to.]

**Lemma 9.20.7.** *The expectation value of the field operator*

$$\langle \alpha | \hat{E}(\chi) | \alpha \rangle = \frac{1}{2} (\langle \alpha | \hat{a} \exp[-i\chi] + \hat{a}^\dagger \exp[i\chi] | \alpha \rangle) \quad (9.20.38)$$

$$= \frac{1}{2} \left( \langle \alpha | \hat{a} \exp[-i\chi] | \alpha \rangle + \langle \alpha | \hat{a}^\dagger \exp[i\chi] | \alpha \rangle \right) \quad (9.20.39)$$

$$= \frac{1}{2} (\alpha \exp[-i\chi] + \alpha^* \exp[i\chi]) \quad (9.20.40)$$

$$= |\alpha| \cos(\chi - \theta), \quad (9.20.41)$$

where  $\alpha$  is the complex amplitude and  $\alpha = |\alpha| \exp[i\theta]$ .

We can plot  $\langle \alpha | \hat{E}(\chi) | \alpha \rangle$  in Phasor Diagram in terms of  $|\alpha|$  and  $\theta$ .

**Lemma 9.20.8.** *The fluctuations (variance) of the E-field*

$$(\Delta \hat{E}(\chi))^2 = \langle \alpha | \hat{E}(\chi)^2 | \alpha \rangle - \langle \alpha | \hat{E}(\chi) | \alpha \rangle^2 = \frac{1}{4}. \quad (9.20.42)$$

*Proof.*

$$(\Delta \hat{E}(\chi))^2 = \langle \alpha | \hat{E}(\chi)^2 | \alpha \rangle - \langle \alpha | \hat{E}(\chi) | \alpha \rangle^2 \quad (9.20.43)$$

$$= \frac{1}{4} \langle \alpha | \left[ \hat{a} \exp[-i\chi] + \hat{a}^\dagger \exp[i\chi] \right]^2 | \alpha \rangle - |\alpha|^2 \cos^2(\chi - \theta) \quad (9.20.44)$$

$$= \frac{1}{4} \langle \alpha | \hat{a}^2 \exp[-2i\chi] + \hat{a}^{*2} \exp[2i\chi] + \hat{a}\hat{a}^\dagger + \hat{a}^\dagger\hat{a} | \alpha \rangle - |\alpha|^2 \cos^2(\chi - \theta) \quad (9.20.45)$$

$$= \frac{1}{4} (\alpha^2 \exp[-2i\chi] + \alpha^{*2} \exp[2i\chi] + 2|\alpha|^2 + 1) - |\alpha|^2 \cos^2(\chi - \theta) \quad (9.20.46)$$

$$= \frac{1}{4} (|\alpha \exp[-i\chi] + \alpha^* \exp[i\chi]|^2 + 1) - |\alpha|^2 \cos^2(\chi - \theta) \quad (9.20.47)$$

$$= \frac{1}{4} (4|\alpha|^2 \cos^2(\chi - \theta) + 1) - |\alpha|^2 \cos^2(\chi - \theta) \quad (9.20.48)$$

$$= |\alpha|^2 \cos^2(\chi - \theta) + \frac{1}{4} - |\alpha|^2 \cos^2(\chi - \theta) \quad (9.20.49)$$

$$= \frac{1}{4}. \quad (9.20.50)$$

□

**Lemma 9.20.9.** *The expectation value of the Energy of Coherent States*

$$\langle \alpha | \hat{H} | \alpha \rangle = \hbar\omega(\bar{n} + 1/2). \quad (9.20.51)$$

*Proof.*

$$\langle \alpha | \hat{H} | \alpha \rangle = \hbar\omega \langle \alpha | \hat{n} + \frac{1}{2} | \alpha \rangle \quad (9.20.52)$$

$$= \hbar\omega (|\alpha|^2 + 1/2) \quad (9.20.53)$$

$$= \hbar\omega \left( \bar{n} + \frac{1}{2} \right). \quad (9.20.54)$$

□

**Lemma 9.20.10.** *The fluctuations of the energy of coherent States*

$$\Delta H = \hbar\omega|\alpha|. \quad (9.20.55)$$

*Proof.*

$$(\Delta H)^2 = \langle \alpha | \hat{H}^2 | \alpha \rangle - \langle \alpha | \hat{H} | \alpha \rangle^2, \quad (9.20.56)$$

where

$$\langle \alpha | \hat{H}^2 | \alpha \rangle = \langle \alpha | (\hbar\omega(\hat{n} + \frac{1}{2}))^2 | \alpha \rangle \quad (9.20.57)$$

$$= \hbar^2\omega^2 \langle \alpha | \hat{n}^2 + \hat{n} + \frac{1}{4} | \alpha \rangle \quad (9.20.58)$$

$$= \hbar^2\omega^2 \left( |\alpha|^4 + |\alpha|^2 + \frac{1}{4} \right). \quad (9.20.59)$$

Then

$$(\Delta H)^2 = \hbar^2\omega^2 \left( |\alpha|^4 + |\alpha|^2 + \frac{1}{4} \right) - \hbar^2\omega^2 \left( |\alpha|^2 + \frac{1}{2} \right)^2 \quad (9.20.60)$$

$$= 0 \quad (9.20.61)$$

□

[SZQ: 2023.04.21: I derive wrong answer. Where is the mistake?]

## 9.21 Quadrature Operators and Phase Space of Field States

**Definition 9.21.1.** *Classical electromagnetic field is*

$$E(t) = E_0 \cos(\omega t + \theta) \quad (9.21.1)$$

$$= E_0 \cos \theta \cos \omega t - E_0 \sin \theta \sin \omega t \quad (9.21.2)$$

$$= X_1 \cos \omega t + X_2 \sin \omega t, \quad (9.21.3)$$

where  $X_1, X_2$  are quadrature variables defined as

$$X_1 = E_0 \cos \theta \quad (9.21.4)$$

$$X_2 = -E_0 \sin \theta. \quad (9.21.5)$$

**Definition 9.21.2.** *Phasor representation of field*

$$a(t) = E_0 \exp[-i\theta] \exp[-i\omega t] = a \exp[-i\omega t], \quad (9.21.6)$$

where  $a$  is defined as  $E_0 \exp[-i\theta]$ .  $a(t)$  is called the phaor.

**Lemma 9.21.1** (Relations between the Phasor representation and quadrature variables).

$$a = X_1 + iX_2. \quad (9.21.7)$$

*Proof.* By definition 9.21.2,

$$a = E_0 \exp [-i\theta] \quad (9.21.8)$$

$$= E_0 (\cos \theta - i \sin \theta) \quad (9.21.9)$$

$$= E_0 \cos \theta - i E_0 \sin \theta \quad (9.21.10)$$

By definition 9.21.1

$$= X_1 + iX_2. \quad (9.21.11)$$

□

[SZQ: 2023.04.21:  $a$  is determined by  $E_0$  and  $\theta$ .]

**Corollary 9.21.1.**

$$X_1 = \operatorname{Re}(a) = \frac{1}{2} (a + a^*) \quad (9.21.12)$$

$$X_2 = \operatorname{Im}(a) = \frac{1}{2i} (a - a^*). \quad (9.21.13)$$

*Proof.* By lemma 9.21.1, we can prove this corollary. □

### Quantum-Quadrature Operators

**Definition 9.21.3.** *Quadrature operators:*

$$\hat{x} \hat{=} \hat{X}_1 = \frac{1}{2} (\hat{a} + \hat{a}^\dagger) \quad (9.21.14)$$

$$\hat{p} \hat{=} \hat{X}_2 = \frac{1}{2i} (\hat{a} - \hat{a}^\dagger). \quad (9.21.15)$$

[SZQ:  $\hat{=}$  means "define"] [SZQ: Quadrature operators are **Hermitian operators**, i.e., observables.]

**Lemma 9.21.2.** *The commutation and uncertainty relations:*

$$[\hat{X}_1, \hat{X}_2] = \frac{i}{2} \quad (9.21.16)$$

$$\Delta X_1 \Delta X_2 \geq \frac{1}{4}. \quad (9.21.17)$$

*Proof.*

□

**Definition 9.21.4.** *Generalized quadrature operators have the same commutation relations with quadrature operators:*

$$\hat{X}_\phi = \frac{1}{2} (\hat{a} \exp [-i\phi] + \hat{a}^\dagger \exp [i\phi]) \quad (9.21.18)$$

$$\hat{X}_{\phi+\frac{\pi}{2}} = \frac{1}{2i} (\hat{a} \exp [-i\phi] - \hat{a}^\dagger \exp [i\phi]). \quad (9.21.19)$$

### Phase space distribution of field states

**Lemma 9.21.3.** *In vacuum state  $|0\rangle$ , we have*

$$P^{(0)}(X_1) = |\langle X_1 | 0 \rangle|^2 = \sqrt{\frac{2}{\pi}} \exp [-2X_1^2] \quad (9.21.20)$$

$$P^{(0)}(X_2) = |\langle X_2 | 0 \rangle|^2 = \sqrt{\frac{2}{\pi}} \exp [-2X_2^2]. \quad (9.21.21)$$

$$\text{Fluctuations: } \Delta X_1 = \sqrt{\langle 0 | \hat{X}_1^2 | 0 \rangle - \langle 0 | \hat{X}_1 | 0 \rangle^2} = \frac{1}{2}. \quad (9.21.22)$$

**Lemma 9.21.4.** *In fock state  $|n\rangle$ , we have*

$$P^{(n)}(X_1) = |\langle X_1 | n \rangle|^2 \quad (9.21.23)$$

$$= \sqrt{\frac{2}{\pi}} \frac{1}{2^n n!} \exp [-2X_1^2] \left( H_n(\sqrt{2}X_1) \right)^2 \quad (9.21.24)$$

$$\Delta X_1 = \frac{1}{2} \sqrt{2n+1}. \quad (9.21.25)$$

### Coherent state - A displaced vacuum

**Definition 9.21.5.** *Displacement operator (shifts any coherent state by  $\alpha$ )*

$$\hat{D}(\alpha) = \exp \left[ \alpha \hat{a}^\dagger - \alpha^* \hat{a} \right]. \quad (9.21.26)$$

**Corollary 9.21.2.** *Coherent state from vacuum state*

$$|\alpha\rangle = \hat{D}(\alpha)|0\rangle. \quad (9.21.27)$$

**Squeezed states of light** We have nice pictures in the video.

We have phase squeezed state and amplitude squeezed state.

## 9.22 Thermal Radiation States and Planck's Black Body Radiation Formula

**Definition 9.22.1** (Incoherent superposition).

**Thermal radiation** Thermal radiation states of a single field mode is

$$\hat{\rho} = \sum_n P_n |n\rangle \langle n|, \quad (9.22.1)$$

where  $P_n = e^{-E_n/(k_B T)}/Z$ ,  $Z$  is the normalization factor,  $E_n = \hbar\omega n$  and  $k_B$  is Boltzmann factor. We should note that  $\hat{\rho}$  is incoherent superposition!

**Average thermal mode occupation**

$$\bar{n}_\omega = \sum_n P_n = \sum n \left( e^{-\hbar\omega/(k_B T)} \right)^n / Z \quad (9.22.2)$$

$$= \sum n q^n / Z \quad (9.22.3)$$

$$= \sum q \left( \frac{\partial}{\partial q} q^n \right) / Z \quad (9.22.4)$$

$$= q \frac{\partial}{\partial q} \left( \sum q^n \right) / Z \quad (9.22.5)$$

$$= \frac{1}{q} \frac{1}{1 - q} \quad (9.22.6)$$

$$= \frac{1}{e^{\hbar\omega/(k_B T)} - 1}. \quad (9.22.7)$$

where  $q = e^{-\hbar\omega/(k_B T)}$  and  $Z = \sum_n e^{-n\hbar\omega/(k_B T)} = \sum_{n=0}^{\infty} q^n = 1/(1 - q)$  since  $q$  is real and  $q < 1$ . So photons,

$$\bar{n}_\omega = \frac{1}{e^{\hbar\omega/(k_B T)} - 1} \quad (9.22.8)$$

is **Bose-Einstein distribution**.

**Average energy in mode** The thermal energy in single radiation mode is

$$\overline{E_\omega} = \bar{n}_\omega \hbar\omega, \quad (9.22.9)$$

where  $\bar{n}_\omega$  is the average number of photons and  $\hbar\omega$  is the single photon energy.

**Plank radiation formula**

$$u(\omega)d\omega = \bar{n}_\omega \hbar\omega \frac{dN}{d\omega} d\omega \frac{1}{V} \quad (9.22.10)$$

is the energy density in frequency interval  $[\omega, \omega + d\omega]$ , where  $\bar{n}_\omega \hbar\omega$  is average energy stored in single mode,  $\frac{dN}{d\omega} d\omega$  is the number of modes in  $[\omega, \omega + d\omega]$  and  $V$  is the volume of box.

**Density of states in a box** Quantized (discrete) radiation modes is

$$k_x = \frac{2\pi}{L} n_x, k_y = \frac{2\pi}{L} n_y, k_z = \frac{2\pi}{L} n_z, \quad (9.22.11)$$

where  $n_x, n_y, n_z$  are intergers and  $\frac{dN}{d\omega} = \frac{dN}{dk}$  due to  $\omega = ck$ . Now we can calculate

$$N(k) = \frac{4}{3}\pi k^3 / (2\pi/L)^3 = \frac{1}{6} \frac{1}{\pi^2} k^3 L^3. \quad (9.22.12)$$

Then we have

$$\frac{dN}{dk} = \frac{1}{2} \frac{k^2}{\pi^2} V. \quad (9.22.13)$$

Using that  $\omega = ck$  so  $k = \omega/c$ , we have

$$\frac{dN}{d\omega} = \frac{dN}{cdk} = \frac{1}{2} \cdot 2 \frac{\omega^2}{\pi^2 c^3} V \quad (9.22.14)$$

$$= \frac{\omega^2}{\pi^2 c^3} V. \quad (9.22.15)$$

Then we have

$$\frac{dN}{d\omega} d\omega = \frac{\omega^2}{\pi^2 c^3} V d\omega. \quad (9.22.16)$$

We finally arrive at the Plank Radiation Formula as

$$u(\omega) = \frac{1}{e^{\hbar\omega/(k_B T)} - 1} \hbar\omega \frac{\omega^2}{\pi^2 c^3} = \frac{\hbar\omega^3}{\pi^2 c^3} \frac{1}{e^{\hbar\omega/(k_B T)} - 1}. \quad (9.22.17)$$

## 9.23 The Classical Beamsplitter

Assume all beams have same polarization and frequency. We input  $E_1, E_2$  and output  $E_3, E_4$ . There is a nice picture to illustrate it. Then we have

$$E_3 = RE_1 + TE_2 \quad (9.23.1)$$

$$E_4 = T'E_1 + R'E_2, \quad (9.23.2)$$

which is

$$\begin{pmatrix} E_3 \\ E_4 \end{pmatrix} = \begin{pmatrix} R & T \\ T' & R' \end{pmatrix} \begin{pmatrix} E_1 \\ E_2 \end{pmatrix}, \quad (9.23.3)$$

where  $R, R'$  denote reflection coefficients, which are complex coefficients,  $T, T'$  denote transmission coefficients, which are complex coefficients.

The simplified case: symmetric Beamsplitter

$$\begin{pmatrix} E_3 \\ E_4 \end{pmatrix} = \begin{pmatrix} R & T \\ T & R \end{pmatrix} \begin{pmatrix} E_1 \\ E_2 \end{pmatrix}. \quad (9.23.4)$$

**Lemma 9.23.1.** *The energy conservation*

$$|E_1|^2 + |E_2|^2 = |E_3|^2 + |E_4|^2 \quad (9.23.5)$$

leads to

$$|R|^2 + |T|^2 = 1 \quad (9.23.6)$$

$$RT^* + TR^* = 0. \quad (9.23.7)$$

*Proof.*

$$|E_3|^2 + |E_4|^2 = (RE_1 + TE_2)(RE_1 + TE_2)^* + (TE_1 + RE_2)(TE_1 + RE_2)^* \quad (9.23.8)$$

$$= |R|^2|E_1|^2 + |T|^2|E_2|^2 + RT^*E_1E_2^* + TR^*E_2E_1^* \quad (9.23.9)$$

$$+ |T|^2|E_1|^2 + |R|^2|E_2|^2 + TR^*E_1E_2^* + RT^*E_2E_1^* \quad (9.23.10)$$

$$= (|R|^2 + |T|^2)|E_1|^2 + (|R|^2 + |T|^2)|E_2|^2 \quad (9.23.11)$$

$$+ (RT^* + TR^*)E_1E_2^* + (RT^* + TR^*)E_2E_1^* \quad (9.23.12)$$

$$= |E_1|^2 + |E_2|^2. \quad (9.23.13)$$

Therefore,

$$|R|^2 + |T|^2 = 1 \quad (9.23.14)$$

$$RT^* + TR^* = 0. \quad (9.23.15)$$

□

**Lemma 9.23.2.** Define  $R = |R|e^{i\phi_R}$  and  $T = |T|e^{i\phi_T}$ . We have  $\phi_R - \phi_T = \frac{\pi}{2}$ .

*Proof.* By the second requirement in 9.23.1, we have

$$|R||T|e^{i(\phi_R - \phi_T)} + |R||T|e^{-i(\phi_R - \phi_T)} = 0. \quad (9.23.16)$$

So

$$2 \cos(\phi_R - \phi_T) = 0, \quad (9.23.17)$$

leads to

$$\phi_R - \phi_T = \frac{\pi}{2}. \quad (9.23.18)$$

□

Set  $\phi_T = 0$ , then  $\phi_R = \frac{\pi}{2}$ . We then have

$$E_3 = i|R|E_1 + |T|E_2 \quad (9.23.19)$$

$$E_4 = |T|E_1 + i|R|E_2. \quad (9.23.20)$$

**50/50 beamsplitter input-output** Symmetrized beamsplitter input-output relations

$$\begin{pmatrix} E_3 \\ E_4 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} i & 1 \\ 1 & i \end{pmatrix} \begin{pmatrix} E_1 \\ E_2 \end{pmatrix}. \quad (9.23.21)$$

In this case, light is always equally split in port 3 and port 4.

## 9.24 The Quantum Beamsplitter

In the quantum beamsplitter,  $E_1, E_2, E_3, E_4$  are replaced by  $\hat{a}_1, \hat{a}_2, \hat{a}_3, \hat{a}_4$ . There is a nice picture. We have

$$\hat{a}_3 = R\hat{a}_1 + T\hat{a}_2 \quad (9.24.1)$$

$$\hat{a}_4 = T\hat{a}_1 + R\hat{a}_2, \quad (9.24.2)$$

which is

$$\begin{pmatrix} \hat{a}_3 \\ \hat{a}_4 \end{pmatrix} = \begin{pmatrix} R & T \\ T & R \end{pmatrix} \begin{pmatrix} \hat{a}_1 \\ \hat{a}_2 \end{pmatrix}, \quad (9.24.3)$$

where  $R$  denotes reflection coefficient, which is complex coefficient,  $T$  denotes transmission coefficient, which is complex coefficient.

**What about  $\hat{a}_1$  and  $\hat{a}_2$ ?**

**Lemma 9.24.1.**

$$\hat{a}_1 = R^*\hat{a}_3 + T^*\hat{a}_4 \quad (9.24.4)$$

$$\hat{a}_2 = T^*\hat{a}_3 + R^*\hat{a}_4. \quad (9.24.5)$$

*Proof.* Recall that

$$\hat{a}_3 = R\hat{a}_1 + T\hat{a}_2 \quad (9.24.6)$$

$$\hat{a}_4 = T\hat{a}_1 + R\hat{a}_2. \quad (9.24.7)$$

We have

$$R^*\hat{a}_3 = |R|^2\hat{a}_1 + R^*T\hat{a}_2 \quad (9.24.8)$$

$$T^*\hat{a}_4 = |T|^2\hat{a}_1 + T^*R\hat{a}_2. \quad (9.24.9)$$

Sum two equations, we have

$$(|R|^2 + |T|^2)\hat{a}_1 + R^*T\hat{a}_2 + T^*R\hat{a}_2 = R^*\hat{a}_3 + T^*\hat{a}_4. \quad (9.24.10)$$

Use lemma 9.23.1, we have

$$\hat{a}_1 = R^*\hat{a}_3 + T^*\hat{a}_4. \quad (9.24.11)$$

Similarly, we can get

$$\hat{a}_2 = T^*\hat{a}_3 + R^*\hat{a}_4. \quad (9.24.12)$$

□

**Lemma 9.24.2.** *We have the creation operators*

$$\hat{a}_1^\dagger = R\hat{a}_3^\dagger + T\hat{a}_4^\dagger \quad (9.24.13)$$

$$\hat{a}_2^\dagger = T\hat{a}_3^\dagger + R\hat{a}_4^\dagger \quad (9.24.14)$$

$$\hat{a}_3^\dagger = R^*\hat{a}_1^\dagger + T^*\hat{a}_2^\dagger \quad (9.24.15)$$

$$\hat{a}_4^\dagger = T^*\hat{a}_1^\dagger + R^*\hat{a}_2^\dagger. \quad (9.24.16)$$

**Single photon on BS** **Input state** We have a nice picture in the video.

$$|1\rangle_1|0\rangle_2 \quad (9.24.17)$$

$$|1\rangle_1|0\rangle_2 = \hat{a}_1^\dagger|0\rangle_1|0\rangle_2 \quad (9.24.18)$$

**Output state**

$$\hat{a}_1^\dagger = R\hat{a}_3^\dagger + T\hat{a}_4^\dagger. \quad (9.24.19)$$

This shows how  $\hat{a}_1^\dagger$  is splitted in port 3 and port 4. We have the output state as

$$(R\hat{a}_3^\dagger + T\hat{a}_4^\dagger)|0\rangle_3|0\rangle_4 = R|1\rangle_3|0\rangle_4 + T|0\rangle_3|1\rangle_4, \quad (9.24.20)$$

which is an entangled state of photon between field modes.

**Single photon on 50/50 BS** We have

$$\frac{1}{\sqrt{2}}(i|1\rangle_3|0\rangle_4 + |0\rangle_3|1\rangle_4). \quad (9.24.21)$$

**Lemma 9.24.3.** *Average output photon number at port 3*

$$\langle \hat{n}_3 \rangle = \frac{1}{2}, \quad (9.24.22)$$

and at port 4

$$\langle \hat{n}_4 \rangle = \frac{1}{2}, \quad (9.24.23)$$

*Proof.*

$$\langle \hat{n}_3 \rangle = \langle \hat{a}_3^\dagger \hat{a}_3 \rangle \quad (9.24.24)$$

$$= \langle 0|_1 \langle 1| \hat{a}_3^\dagger \hat{a}_3 |1\rangle_1 |0\rangle_2 \quad (9.24.25)$$

$$= \langle 0|_1 \langle 1| (R^*\hat{a}_1^\dagger + T^*\hat{a}_2^\dagger)(R\hat{a}_1 + T\hat{a}_2) |1\rangle_1 |0\rangle_2 \quad (9.24.26)$$

number operator  $\hat{n} = \hat{a}^\dagger \hat{a}$

$$= |R|^2 \langle 0|_1 \langle 1| 1 |1\rangle_1 |0\rangle_2 \quad (9.24.27)$$

$$= \frac{1}{2}. \quad (9.24.28)$$

Similarly, we have

$$\langle \hat{n}_4 \rangle = \frac{1}{2}. \quad (9.24.29)$$

□

**Lemma 9.24.4.** *Correlations*

$$\langle \hat{n}_3 \hat{n}_4 \rangle =_2 \langle 0|_1 \langle 1 | \hat{n}_3^\dagger \hat{n}_4 | 1 \rangle_1 | 0 \rangle_2 = 0. \quad (9.24.30)$$

**Remark 9.24.1.** • *non-classical Correlations.*

- due to entangled single photon state between field modes.

**What about coherent states?**

## 9.25 The Quantum Mach-Zehnder Interferometer

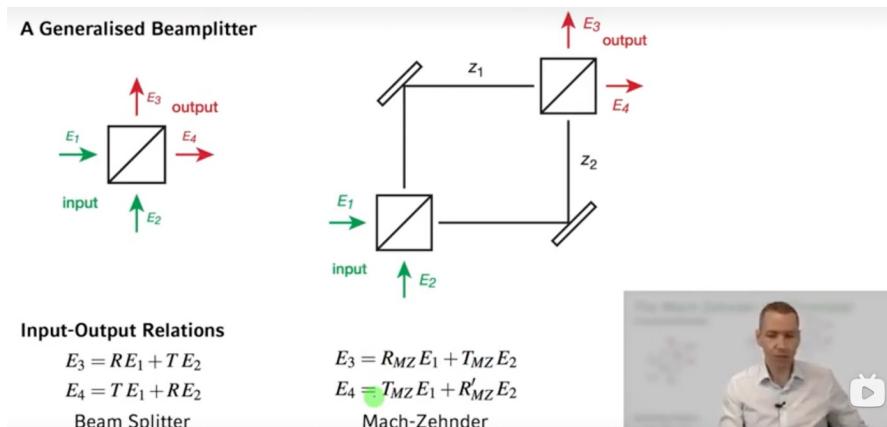


Figure 9.1: Mach-Zehnder Interferometer

The generalised transmission-reflection coefficients are

$$R_{MZ} = R^2 e^{ikz_1} + T^2 e^{ikz_2} \quad (9.25.1)$$

$$T_{MZ} = RT(e^{ikz_1} + e^{ikz_2}) \quad (9.25.2)$$

$$R'_{MZ} = T^2 e^{ikz_1} + R^2 e^{ikz_2}. \quad (9.25.3)$$

Assuming that there is no input at  $E_2$ , then the output light intensities are

$$|E_3|^2 = |R_{MZ}|^2 |E_1|^2 \quad (9.25.4)$$

$$|E_4|^2 = |T_{MZ}|^2 |E_1|^2. \quad (9.25.5)$$

Plug the transmission-reflection coefficients into the  $|E_4|^2$ , we have

$$|E_4|^2 = |R|^2 |T|^2 |e^{ikz_1} + e^{ikz_2}|^2 |E_1|^2 = 4|R|^2 |T|^2 \cos^2(k(z_1 - z_2)/2) |E_1|^2. \quad (9.25.6)$$

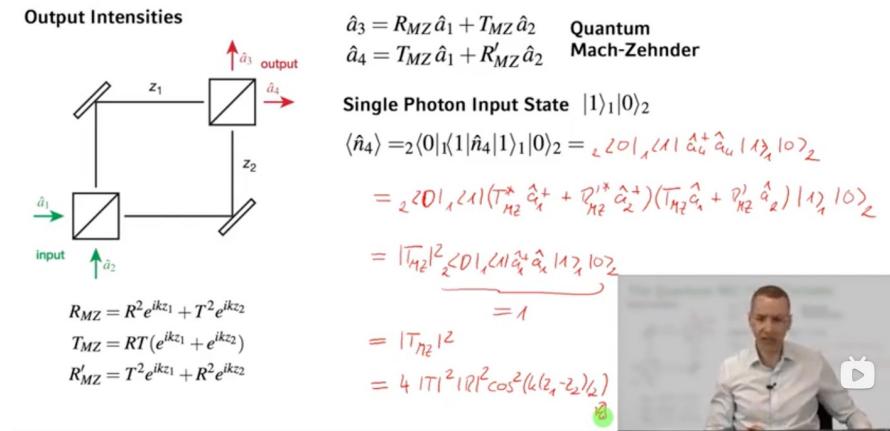


Figure 9.2: Quantum Mach-Zehnder Interferometer

**The quantum MZ interferometer** For the quantum Mach-Zehnder Interferometer, the output creation operators are

$$\hat{a}_3 = R_{MZ}\hat{a}_1 + T_{MZ}\hat{a}_2 \quad (9.25.7)$$

$$\hat{a}_4 = T_{MZ}\hat{a}_1 + R'_{MZ}\hat{a}_2. \quad (9.25.8)$$

Then for the single photon input state  $|1\rangle_1|0\rangle_2$ , the expectation value of the  $\hat{n}_4$  is

$$\langle \hat{n}_4 \rangle =_2 \langle 0|_1 \langle 1|\hat{n}_4|1\rangle_1|0\rangle_2 =_2 \langle 0|_1 \langle 1|\hat{a}_4^\dagger \hat{a}_4|1\rangle_1|0\rangle_2 \quad (9.25.9)$$

$$=_2 \langle 0|_1 \langle 1|(T_{MZ}^* \hat{a}_1^\dagger + R'_{MZ}^* \hat{a}_2^\dagger)(T_{MZ}\hat{a}_1 + R'_{MZ}\hat{a}_2)|1\rangle_1|0\rangle_2 \quad (9.25.10)$$

$$= |T_{MZ}|_2^2 \langle 0|_1 \langle 0|\hat{a}_1^\dagger \hat{a}_1|1\rangle_1|0\rangle_2 \quad (9.25.11)$$

$$= |T_{MZ}|^2 \quad (9.25.12)$$

$$= 4|T|^2|R|^2 \cos^2(k(z_1 - z_2)/2). \quad (9.25.13)$$

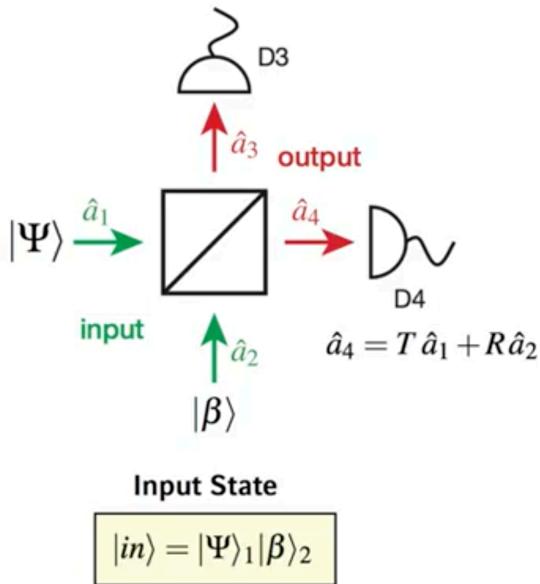
We see that the results are the same as the classical one.

## 9.26 Balanced Homodyne Detection

Input states:  $|in\rangle = |\Psi\rangle_1|\beta\rangle_2$ , where  $|\beta\rangle$  is the coherent state.

**Lemma 9.26.1.** For a balanced 50/50 beamsplitter, difference photocurrent

$$i_{34} \propto -2|\beta|\langle \hat{X}(\phi + \pi/2) |\Psi\rangle. \quad (9.26.1)$$



*Proof.* For a balanced 50/50 beamsplitter,  $T = \frac{1}{\sqrt{2}}$ ,  $R = \frac{1}{\sqrt{2}}i$ . Recall that 9.24.2,  $\hat{a}_4 = T\hat{a}_1 + R\hat{a}_2$ . We then have

$$\hat{n}_4 = \hat{a}_4^\dagger \hat{a}_4 \quad (9.26.2)$$

$$= (T^* \hat{a}_1^\dagger + R^* \hat{a}_2^\dagger)(T\hat{a}_1 + R\hat{a}_2) \quad (9.26.3)$$

$$= \left(\frac{1}{\sqrt{2}}\hat{a}_1^\dagger - \frac{1}{\sqrt{2}}i\hat{a}_2^\dagger\right)\left(\frac{1}{\sqrt{2}}\hat{a}_1 + \frac{1}{\sqrt{2}}i\hat{a}_2\right) \quad (9.26.4)$$

$$= \frac{1}{2}(\hat{a}_1^\dagger - i\hat{a}_2^\dagger)(\hat{a}_1 + i\hat{a}_2) \quad (9.26.5)$$

$$= \frac{1}{2}(\hat{a}_1^\dagger \hat{a}_1 - i\hat{a}_2^\dagger \hat{a}_1 + i\hat{a}_1^\dagger \hat{a}_2 + \hat{a}_2^\dagger \hat{a}_2). \quad (9.26.6)$$

Similarly, we have

$$\hat{n}_3 = \hat{a}_3^\dagger \hat{a}_3 \quad (9.26.7)$$

$$= \frac{1}{2}(\hat{a}_1^\dagger \hat{a}_1 + i\hat{a}_2^\dagger \hat{a}_1 - i\hat{a}_1^\dagger \hat{a}_2 + \hat{a}_2^\dagger \hat{a}_2). \quad (9.26.8)$$

We have  $\beta = |\beta| \exp[i\phi]$ . Then,

$$i_{34} = i_3 - i_4 \propto \langle in | \hat{n}_3 - \hat{n}_4 | in \rangle = -2|\beta|_2 \langle \Psi | \frac{1}{2i} (\hat{a}_2^\dagger \hat{a}_1 - \hat{a}_1^\dagger \hat{a}_2) | \beta \rangle_2 | \Psi \rangle_1 \quad (9.26.9)$$

To derive the following equation, we need to know that  $\hat{a}_1^\dagger, \hat{a}_1$  are applied to  $|\Psi\rangle_1$

according to its subindex. Similar for subindex 2

$$= -2|\beta|_1 \langle \Psi | \frac{1}{2i} (\beta^* \hat{a}_1 - \hat{a}_1^\dagger \beta) | \Psi \rangle_1 \quad (9.26.10)$$

$$= -2|\beta|_1 \langle \Psi | \frac{1}{2i} (\hat{a}_1 \exp[-i\phi] - \hat{a}_1^\dagger \exp[i\phi]) | \Psi \rangle_1 \quad (9.26.11)$$

By 9.21.4

$$= -2|\beta|_1 \langle \Psi | \hat{X}_1(\phi + \frac{\pi}{2}) | \Psi \rangle_1, \quad (9.26.12)$$

where  $|\beta|$  is the amplification of quadrature signal.

[SZQ:  $\hat{a}_1^\dagger \hat{a}_2$  must mean  $\hat{a}_1^\dagger \otimes \hat{a}_2$ . In addition, the order of  $\hat{a}_1^\dagger$  and  $\hat{a}_2$  needs to be changed according to subindex. These correspondences can be easily understood by the above picture showing the physical setting, where we know that the subindex must coincide.]  $\square$

## 9.27 Quantized Light-Atom Interaction

This section deals with quantized fields and quantized atom.

We have

$$\hat{H} = \hat{H}_A + \hat{H}_R + \hat{H}_I. \quad (9.27.1)$$

There is a nice picture to show the quantized light field and quantized atom.

Atomic Hamiltonian is

$$\hat{H}_A = \hbar\omega_{21} |2\rangle\langle 2| \quad (9.27.2)$$

$$= \hbar\omega_{21} \hat{\sigma}^+ \hat{\sigma}^-, \quad (9.27.3)$$

where  $\hat{\sigma}^+ = |2\rangle\langle 1|$  and  $\hat{\sigma}^- = |1\rangle\langle 2|$ .

Light Hamiltonian is

$$\hat{H}_R = \sum_k \hbar\omega_k (\hat{a}_k^\dagger \hat{a}_k + \frac{1}{2}). \quad (9.27.4)$$

Light-Atom Interaction Hamiltonian is

$$\hat{H}_I = -\hat{d} \cdot \hat{E}(t). \quad (9.27.5)$$

Rewrite in atomic ladder operators

$$\hat{d} = e\hat{r} = i\hat{d}\hat{d}^\dagger = \sum_{i,j} |i\rangle\langle i|\hat{d}|j\rangle j|. \quad (9.27.6)$$

Using parity argument and for the two level system, we have

$$\hat{d} = d_{12}(\hat{\sigma}^+ + \hat{\sigma}^-) := d_{12}(\hat{\sigma}^\dagger + \hat{\sigma}). \quad (9.27.7)$$

E-field operator in Schrodinger picture is

$$\hat{E}(r) = \sum_k \epsilon_k \sqrt{\frac{\hbar\omega_k}{2\varepsilon_0 V}} [\hat{a}_k e^{ikr} + \hat{a}_k^\dagger e^{-ikr}]. \quad (9.27.8)$$

This is a time-independent operator. Then the Interaction Hamiltonian is

$$\hat{H}_I = \sum_k \hbar g_k (\hat{a}_k e^{ikr} + \hat{a}_k^\dagger e^{-ikr}) (\hat{\sigma}^\dagger + \hat{\sigma}), \quad (9.27.9)$$

where  $g_k := \sqrt{\frac{\omega_k}{2\varepsilon_0 \hbar V}} d_{12} \cdot \epsilon_k$  is the light-atom coupling constant (mode-dependent).

Simplified  $r = 0$ , we have

$$\hat{H}_I = \sum_k \hbar g_k (\hat{a}_k + \hat{a}_k^\dagger) (\hat{\sigma}^\dagger + \hat{\sigma}). \quad (9.27.10)$$

We have four terms:  $\hat{a}_k^\dagger \hat{\sigma}^\dagger$  create atomic excitation and create photon;  $\hat{a}_k \hat{\sigma}^\dagger$  create atomic excitation and absorpt photon;  $\hat{a}_k^\dagger \hat{\sigma}$  destroy atomic excitation and emmits photon;  $\hat{a}_k \hat{\sigma}$  destroy photon and destroy atomic excitation.  $\hat{a}_k^\dagger \hat{\sigma}^\dagger$  and  $\hat{a}_k \hat{\sigma}$  do not obey energy conservation law, thus are dropped. So we have

$$\hat{H}_I = \sum_k \hbar g_k (\underbrace{\hat{a}_k^\dagger \hat{\sigma}}_{\text{emission}} + \underbrace{\hat{a}_k \hat{\sigma}^\dagger}_{\text{absorption}}). \quad (9.27.11)$$

Absorption and emission matrix elements. Light field in mode  $k$  and atom are described by states

$$|n_k, i\rangle = |n_k\rangle \otimes |i\rangle, \quad (9.27.12)$$

where  $|n_k\rangle$  is the photon state and  $|i\rangle$  is the atomic state. Now, we consider two-level atom. Then we have the absorption of photon from mode  $k$  as

$$\langle n_k - 1, 2 | \hat{H}_I | n_k, 1 \rangle = \hbar g_k \langle n_k - 1, 2 | \hat{a}_k^\dagger \hat{\sigma} + \hat{a}_k \hat{\sigma}^\dagger | n_k, 1 \rangle \quad (9.27.13)$$

$$= \hbar g_k \sqrt{n_k}. \quad (9.27.14)$$

We have the emission of photon into mode  $k$  as

$$\langle n_k + 1, 1 | \hat{H}_I | n_k, 2 \rangle = \hbar g_k \langle n_k + 1, 1 | \hat{a}_k^\dagger \hat{\sigma} + \hat{a}_k \hat{\sigma}^\dagger | n_k, 2 \rangle \quad (9.27.15)$$

$$= \hbar g_k \sqrt{n_k + 1}, \quad (9.27.16)$$

where the  $+1$  is surprisingly the **spontaneous emission**.

## 9.28 Jaynes-Cummings Model

Interaction of Time-independent light-atom interaction with single mode of radiation field. Since it is single mode, we drop  $k$ . The Jaynes-Cummings hamiltonian is

$$\hat{H}_{JC} = \hbar\omega_{21}\hat{\sigma}^\dagger\hat{\sigma} + \hbar\omega\hat{a}^\dagger\hat{a} + \hbar g(\hat{a}^\dagger\hat{\sigma} + \hat{a}\hat{\sigma}^\dagger). \quad (9.28.1)$$

Switch to interaction picture. Interaction hamiltonian in interaction picture is

$$\hat{H}_I^{in}(t) = e^{\frac{i}{\hbar}\hat{H}_0 t}, \quad (9.28.2)$$

$$|\Psi^{in}(t)\rangle = e^{\frac{i}{\hbar}\hat{H}_0 t}|\Psi(t)\rangle. \quad (9.28.3)$$

Interaction Hamiltonian (in interaction picture) is

$$\hat{H}_I^{in}(t) = \hbar g \left( \hat{a}^\dagger\hat{\sigma}e^{i\Delta t} + \hat{a}\hat{\sigma}^\dagger e^{-i\Delta t} \right), \quad (9.28.4)$$

where  $\Delta = \omega - \omega_{21}$ .

Time evolution is

$$i\hbar\frac{\partial}{\partial t}|\Psi^{in}(t)\rangle = \hat{H}_I^{in}(t)|\Psi^{in}(t)\rangle. \quad (9.28.5)$$

Assume Ansatz wavefunction

$$|\Psi^{in}(t)\rangle = \sum_n c_{1,n}|n, 1\rangle + c_{2,n}|n, 2\rangle. \quad (9.28.6)$$

[SZQ: Ansatz means "an assumption about the form of an unknown function which is made in order to facilitate solution of an equation or other problem."]

**Time evolution in JC model** Which states are coupled through  $\hat{H}_I^{in}(t)$ ?

$$|1, n+1\rangle \rightleftharpoons_{emission}^{absorption} |2, n\rangle. \quad (9.28.7)$$

Plug the ansatz function into the time evolution equation and have

$$\dot{c}_{1,n+1}(t) = -ig\sqrt{n+1}e^{+i\Delta t}c_{2,n}(t), \quad (9.28.8)$$

$$\dot{c}_{2,n}(t) = -ig\sqrt{n+1}e^{-i\Delta t}c_{1,n+1}(t). \quad (9.28.9)$$

**Time evolution in JC model-solutions** Initial conditions ( $n+1$  photons, atom in ground state) is assumed as

$$c_{1,n+1}(0) = 1. \quad (9.28.10)$$

The resonant interaction  $\Delta = 0$ . The solution is

$$c_{1,n+1}(t) = \cos(g\sqrt{n+1}t) \quad (9.28.11)$$

$$c_{2,n}(t) = -i \sin(g\sqrt{n+1}t). \quad (9.28.12)$$

The corresponding probability is

$$P_{1,n+1}(t) = |c_{1,n+1}(t)|^2 = \frac{1}{2} [1 + \cos(\Omega_n t)], \quad (9.28.13)$$

where  $\Omega_n = 2g\sqrt{n+1}$  is the quantized rabi frequency. Recall that in the semiclassical case,  $\Omega = dE/\hbar$  is continuous.

Initial conditions (0 photons, atom in excited state) with  $n = 0$ ,  $c_{2,0}(0) = 1$  and  $\Delta = 0$ . The probability is

$$P_{2,0}(t) = |c_{1,n+1}(t)|^2 = \frac{1}{2} [1 + \cos(\Omega_0 t)], \quad (9.28.14)$$

where  $\Omega_0 = 2g$  with vacuum state  $n = 0$ . There is nice picture showing that state  $|2,0\rangle$  and  $|1,1\rangle$  swaps continuously. This is called the Vacuum Rabi-Oscillation! This do not happen in classical case!

## 9.29 Interaction of TLA with a Coherent State

This section is also called the Collapse and Revival of Rabi Oscillations.

**Time-independent light-atom interaction and coherent state** The light field is

$$|\alpha\rangle = e^{-\frac{1}{2}|\alpha|^2} \sum_n \frac{\alpha^n}{\sqrt{n!}} |n\rangle. \quad (9.29.1)$$

Assume that atom is initially in excited state  $|2\rangle$ . Then for the combined atom-field state amplitudes are

$$c_{1,n}(0) = 0 \quad (9.29.2)$$

$$|c_{2,n}(0)|^2 = e^{\bar{n}} \frac{\bar{n}^n}{n!}. \quad (9.29.3)$$

Recall from the last lecture, we have

$$c_{2,n}(t) = \cos(g\sqrt{n+1}t) |c_{2,n}(0)|, \quad (9.29.4)$$

$$c_{1,n}(t) = -i \sin(g\sqrt{n}t) |c_{2,n-1}(0)|, \forall n \geq 1. \quad (9.29.5)$$

Then inversion is

$$w(t) = P_2(t) - P_1(t) \quad (9.29.6)$$

$$= \sum_{n=1}^{\infty} |c_{2,n}(t)|^2 - \sum_{n=0}^{\infty} |c_{1,n}(t)|^2 \quad (9.29.7)$$

$$= \sum_{n=0}^{\infty} |c_{2,n}(0)|^2 \cos(2g\sqrt{n+1}t) \quad (9.29.8)$$

$$= e^{\bar{n}} \sum_{n=0}^{\infty} \frac{\bar{n}^n}{n!} \cos(2g\sqrt{n+1}t). \quad (9.29.9)$$

There is a nice picture to plot this inversion for  $\bar{n} = 20$ . This shows a Collapse and Revival of Rabi-Oscillations.

**Estimate collapse time** First, what is collapse time? The collapse time  $t_c$  should satisfy

$$\Omega_{\bar{n}+\Delta n} t_c - \Omega_{\bar{n}-\Delta n} t_c = \pi. \quad (9.29.10)$$

For large  $\bar{n}$ , we have  $\Delta n = \sqrt{\bar{n}}$ . We then get

$$(2g\sqrt{\bar{n} + \sqrt{\bar{n}}} - 2g\sqrt{\bar{n} - \sqrt{\bar{n}}})t_c \sim \pi \quad (9.29.11)$$

$$2g \left( \sqrt{\bar{n}} \sqrt{1 + \frac{1}{\sqrt{\bar{n}}}} - \sqrt{\bar{n}} \sqrt{1 - \frac{1}{\sqrt{\bar{n}}}} \right) \quad (9.29.12)$$

$$2g \left( \sqrt{\bar{n}} 1 + \frac{1}{2} \frac{1}{\sqrt{\bar{n}}} - \sqrt{\bar{n}} 1 - \frac{1}{2} \frac{1}{\sqrt{\bar{n}}} \right) \sim \pi. \quad (9.29.13)$$

We have

$$2gt_c \sim \pi, \quad (9.29.14)$$

so

$$t_c \simeq \frac{\pi}{2g} = \frac{\pi}{\Omega_0}. \quad (9.29.15)$$

[SZQ: " $\sim$ " means asymptotically equal.]

**Estimate revival time** What is revival time?

$$\Omega_{\bar{n}+1-\Omega_{\bar{n}}} = 2\pi. \quad (9.29.16)$$

Revival when 'neighbouring' Rabi Oscillations become in phase! Plug the Rabi frequency, we have

$$(2g\sqrt{\bar{n}+1} - 2g\sqrt{\bar{n}})t_R \sim 2\pi \quad (9.29.17)$$

$$(2g\sqrt{\bar{n}}\sqrt{1 + \frac{1}{\bar{n}}} - 2g\sqrt{\bar{n}})t_R \sim 2\pi \quad (9.29.18)$$

Using taylor approximation

$$(2g\sqrt{\bar{n}}(1 + \frac{1}{2}\frac{1}{\bar{n}}) - 2g\sqrt{\bar{n}})t_R \sim 2\pi \quad (9.29.20)$$

$$(2g\sqrt{\bar{n}}\frac{1}{2\sqrt{\bar{n}}})t_R \sim 2\pi \quad (9.29.21)$$

$$t_R \sim \frac{2\pi\sqrt{\bar{n}}}{g} = \frac{4\pi\sqrt{\bar{n}}}{\Omega_0}. \quad (9.29.22)$$

**Classical limit** Quantized box  $V \rightarrow \infty$ , then  $g \rightarrow 0$ .

Keep Rabi frequency constant by  $g \rightarrow 0, \bar{n} \rightarrow \infty$ , which means

$$2g\sqrt{\bar{n}} = \Omega = \text{constant} \quad (9.29.23)$$

$$t_c \sim \frac{\pi}{2g} \rightarrow \infty, \text{ for } g \rightarrow 0. \quad (9.29.24)$$

This means collapse becomes absent.

There is a nice experiment of the test of quantized Rabi Oscillations.

## 9.30 Primer on Cavity Quantum Electrodynamics

This section provides the experimental setup and results for the quantized light-atom interaction. There are many pictures reffering to the video.

How can we experimentally observe quantized light-atom interaction?

$$\hat{H}_I = \hbar g(\hat{a}^\dagger \hat{\sigma} + \hat{a} \hat{\sigma}^\dagger), \quad (9.30.1)$$

$$g = \sqrt{\frac{\hbar\omega}{2\varepsilon_0 V}} d_{12} \cdot \epsilon. \quad (9.30.2)$$

$V$  small! Use cavity!!

**Light field cavity** finite reflectivity.

$k$ : transmission losses of cavity (finite reflectivity mirrors).

$\gamma$ : spontaneous emission of atom into free space.

$g$ : coherent coupling.

$k, \gamma$ : incoherent dynamics.

$g$ : coherent dynamics.

Define  $C := \frac{g^2}{k\gamma}$  cooperativity.

**Experimentally used CQED Setups** 1) Superconducting mirrors for photons in microwave regime (Haroche, Walter)

3) On-chip mitrowaave resonators with superconducting qubits (Schoelkopf, Martinis, Wallraff)

Details ref to the video.

**Rydberg atoms**  $d_{12}$  large.

Energy spectrum

$$E_{n,l} = -\frac{R}{(n - \delta_l)^2} = \simeq -\frac{R}{n^2}. \quad (9.30.3)$$

Huge dipole moment

$$\langle \phi_{n,l'} | q\hat{r} | \phi_{n,l} \rangle \simeq qa_0 n^2. \quad (9.30.4)$$

## 9.31 Seeing a Single Photon without Destroying it

**Detecting a photon with photodetector** Usually through absorption: single photon  $\rightarrow$  photodetector.

Can we detect a photon without destroying it ???

Yes! interferometrically. Experimental setup: ref to video.

Step 1:  $\pi/2$  pulse in Ramsey Zone 1 (on i-g transition)

Step 2: a) if 1 photon in cavity  $|g, 1\rangle \rightarrow e^{i\pi}|g, 1\rangle = -|g, 1\rangle$ .  $2gt_{int} = 2\pi$ . b) if 0 photons in cavity  $|g, 0\rangle \rightarrow |g, 0\rangle$ .

Step 3:  $\pi/2$  pulse in Ramsey Zone 2 (on i-g transition)

There are nice pictures about the experimental results.

## 9.32 Dressed States

**Atom-light field without interaction** Atom:  $\hat{H}_A = \hbar\omega_{21}\hat{\sigma}^\dagger\hat{\sigma}$ .

Fight field:  $\hat{H}_L = \hbar\omega\hat{a}^\dagger\hat{a}$ .

**Atom-light field** Group the eigenstates:  $|1, n+2\rangle, |2, n+1\rangle, |1, n+1\rangle, |2, n\rangle, |1, n\rangle, |2, n-1\rangle$ .

**Dressed states**  $\hat{H}_A + \hat{H}_R$  eigenstates of uncoupled hamiltonian  $|1, n+1\rangle, |2, n\rangle \rightleftarrows |1(n)\rangle, |2(n)\rangle$   $\hat{H}_A + \hat{H}_R + \hat{H}_I$  eigenstates of coupled hamiltonian.

Dressed state eigenstates

$$|1(n)\rangle = \sin(\theta/2)|1, n+1\rangle + \cos(\theta/2)|2, n\rangle \quad (9.32.1)$$

$$|2(n)\rangle = \cos(\theta/2)|1, n+1\rangle - \sin(\theta/2)|2, n\rangle, \quad (9.32.2)$$

where  $\tan \theta = -\Omega_n/\Delta$ ,  $\sin(\theta/2) = \sqrt{\frac{\Omega_n^\Delta - \Delta}{2\Omega_n^\Delta}}$ ,  $\cos(\theta/2) = \sqrt{\frac{\Omega_n^\Delta + \Delta}{2\Omega_n^\Delta}}$  and the generalized Rabi frequency is  $\Omega_n^\Delta = \sqrt{\Omega_n^2 + \Delta^2}$ ,  $\Omega_n = 2g\sqrt{n+1}$ .

There is a nice picture.

**Avoided corssing of energy levels** There is a nice picture.

Dressed state for  $\Delta = 0 \rightarrow \theta = \pi/2$ , we have

$$|1(n)\rangle = \frac{1}{\sqrt{2}}(|1, n+1\rangle + |2, n\rangle) \quad (9.32.3)$$

$$|2(n)\rangle = \frac{1}{\sqrt{2}}(|1, n+1\rangle - |2, n\rangle). \quad (9.32.4)$$

**Re-interpreting Rabi Oscillations** Initial state

$$|\Psi(0)\rangle = |1, n+1\rangle = \frac{1}{\sqrt{2}}(|1(n)\rangle + |2(n)\rangle) \quad (9.32.5)$$

$$|\psi(t)\rangle = \frac{1}{\sqrt{2}}(e^{i\Omega_n t}|1(n)\rangle + |2(n)\rangle). \quad (9.32.6)$$

$$P_1(t) = |\langle 1, n+1 | \psi(t) \rangle|^2 \quad (9.32.7)$$

$$= \frac{1}{2} \left| \frac{1}{\sqrt{2}} e^{-i\Omega_n t} + \frac{1}{\sqrt{2}} \right|^2 \quad (9.32.8)$$

$$= \frac{1}{4} |e^{-i\Omega_n t} + 1|^2 \quad (9.32.9)$$

$$= \frac{1}{4} |e^{-i\Omega_n t/2} + e^{i\Omega_n t/2}|^2 \quad (9.32.10)$$

$$= \cos^2(\Omega_n t/2). \quad (9.32.11)$$

# Chapter 10

## A Nutshell of Quantum Mechanics

### 一、量子论基础

#### 1. Plank 黑体辐射假说

能量量子化  $E_\nu = h\nu$

#### 2. 光电效应的实验现象 (为什么不能从波动的观点进行解释)

光强大小正比于光子数目；截止频率  $\nu = A/h$ ；光子吸收、发射时间很短

#### 3. Einstein 的光量子假说

频率  $\nu$  的电磁波辐射场由宏观多个光量子组成， $E_\nu = h\nu = \hbar\omega$ ，逸出功  $A = h\nu - \frac{1}{2}mv^2$ .

#### 4. Bohr 的氢原子理论

定态假设；量子化条件即角动量量子化  $L = n\hbar$ ；跃迁条件  $\Delta E = h(\nu_1 - \nu_2)$ ；给出了能级公式  $E_n \propto -1/n^2$ .

#### 5. 德布罗意物质波

$$\lambda = \frac{h}{p}, \quad p = \hbar k$$

### 二、量子力学的基本概念

#### 1. 量子力学的基本假设 (sy 的 3 条版)

1. 量子孤立系统由量子态描述，量子态是希尔伯特空间中是态矢  $|\psi\rangle$ . 其随时间的演化满足薛定方程

$$i\hbar \frac{\partial}{\partial t} \Psi = \hat{H} \Psi$$

其中， $\hat{H}$  是哈密顿量

II. 每一个可观测量  $A$ , 与希尔伯特空间中的一个厄米算符  $\hat{A}$  相关联。测量结果是它的本征值之一, 概率幅是原量子态在该本征值相应的本征态上的分量 (即展开系数)

$$\hat{A}|i\rangle = a_i|i\rangle; |\psi(t)\rangle = \sum_i C_i(t)|i\rangle; p_i(t) = |\langle i | \psi(t)\rangle|^2 = |C_i(t)|^2$$

III. 对于由态矢  $|\psi\rangle$  描述的系统, 如测量可观测量  $A$ , 得到结果  $a_n$ , 那么测量刚结束时, 系统的状态是

$$\frac{P_n|\psi\rangle}{\sqrt{\langle\psi|P_n|\psi\rangle}}$$

$P_n$  为投影到相应于  $a_n$  对应的  $A$  的本征矢量张成的子空间。

2. 定态薛定谔方程的得到

$$i\hbar\frac{\partial}{\partial t}\Psi = \hat{H}\Psi = \left(-\frac{\hbar^2}{2m}\nabla^2 + V\right)\Psi$$

当  $H$  不含时间时, 可分离变量  $\Psi(\mathbf{r}, t) = \psi(\mathbf{r})\phi(t)$

$$\begin{aligned} i\hbar\frac{1}{\phi}\frac{\partial\phi}{\partial t} &= -\frac{\hbar^2}{2m}\frac{1}{\psi}\nabla^2\psi + V = E \\ \left(-\frac{\hbar^2}{2m}\nabla^2 + V\right)\psi &= E\psi, \phi(t) = \exp[-iEt/\hbar] \end{aligned}$$

3. 波函数的统计诠释

$\int_a^b |\Psi(x, t)|^2 dx$  表示  $t$  时刻发现粒子处于  $a$  和  $b$  之间的几率.

4. 波函数的归一化条件及其意义

$$\int_{-\infty}^{\infty} |\Psi(x, t)|^2 dx = 1$$

全空间的总几率为 1。

5. 量子力学中概率流密度的连续性方程

$$\begin{aligned} \frac{\partial}{\partial t}\rho + \nabla \cdot j(\mathbf{r}, t) &= 0 \\ \rho = |\Psi|^2, j &= \frac{i\hbar}{2m} [\Psi\nabla\Psi^* - \Psi^*\nabla\Psi] \end{aligned}$$

6. 可观测量  $A$  的期待值

$$\langle \hat{A} \rangle = \int \Psi^*(\mathbf{r}, t)\hat{A}\Psi(\mathbf{r}, t)d^3r$$

7. 为什么可观测量对应厄米算符

$$\langle \hat{Q} \rangle = \langle \hat{Q} \rangle^* \rightarrow \langle \psi | \hat{Q}\psi \rangle = \langle \hat{Q}\psi | \psi \rangle \rightarrow \hat{Q} = \hat{Q}^\dagger$$

8. 证明不同本征值的本征函数正交

$$\hat{Q}\psi_i = q_i\psi_i; \quad q_2 \langle \psi_1 | \psi_2 \rangle = \langle \psi_1 | \hat{Q}\psi_2 \rangle = \langle \hat{Q}\psi_1 | \psi_2 \rangle = q_1 \langle \psi_1 | \psi_2 \rangle$$

9. 同时对角化定理 ( $\hat{A}, \hat{B}$  具有共同本征态的条件)

$$[\hat{A}, \hat{B}] = 0$$

10. Gram-Schmidt 正交化法则

$$\psi'_i = \psi_i - \sum_{k=1}^{i-1} \frac{\langle \psi'_k | \psi_i \rangle}{\langle \psi'_k | \psi'_k \rangle} \psi'_k$$

11. 不确定原理

$$\sigma_x = \sqrt{\langle x^2 \rangle - \langle x \rangle^2}; \quad \sigma_x \sigma_p \geq \frac{\hbar}{2}; \quad \Delta E \Delta t \geq \frac{\hbar}{2}$$

$$\sigma_A^2 \sigma_B^2 \geq \left( \frac{1}{2i} \langle [\hat{A}, \hat{B}] \rangle \right)^2$$

[\* 请利用 Schwarz 不等式自行推导]

12. Heisenberg 动力学方程

$$\frac{d}{dt} \langle \hat{Q} \rangle = \frac{i}{\hbar} \langle [\hat{H}, \hat{Q}] \rangle + \left\langle \frac{\partial \hat{Q}}{\partial t} \right\rangle$$

13. 动量空间的本征函数及其正交与完备条件

$$\hat{p}\psi_p(x) = p\psi_p(x), \quad \psi_p(x) = \frac{1}{\sqrt{2\pi\hbar}} \exp[ipx/\hbar]$$

$$\int \psi_{p'}(x)\psi_p(x)dx = \delta(p - p')$$

$$\int \psi_p^*(x)\psi_p(y)dp = \delta(x - y) \text{ 或 } \int |p\rangle\langle p| = 1$$

14. Dirac 符号下能量本征函数的正交与完备条件

$$\langle m | n \rangle = \delta_{mn}, \sum_n |n\rangle\langle n| = 1$$

15. 位置算符与动量算符的对易关系

$$[\hat{x}, \hat{p}_x] = i\hbar, \quad [\hat{x}, \hat{p}_y] = 0$$

### 三、定态薛定谔方程的解

1. 无限深势阱的解及函数图像

$$\psi_n = \sqrt{\frac{2}{a}} \sin\left(\frac{n\pi x}{a}\right), \quad E_n = \frac{n^2\pi^2\hbar^2}{2ma^2}$$

2. 谐振子的阶梯算符, 对偶关系及其对本征波函数的作用效果

$$\hat{a}_{\pm} = \frac{1}{\sqrt{2\hbar m\omega}} (\mp i\hat{p} + m\omega x), \quad [\hat{a}_-, \hat{a}_+] = 1$$

$$\hat{a}_+ |n\rangle = \sqrt{n+1} |n+1\rangle, \quad \hat{a}_- |n\rangle = \sqrt{n} |n-1\rangle$$

3. 谐振子的本征函数图像及总能

$$\hat{H} = \left( \hat{a}_+ \hat{a}_- + \frac{1}{2} \right) \hbar\omega, \quad E_n = \left( n + \frac{1}{2} \right) \hbar\omega$$

#### 4. 束缚态与散射态的定义

束缚态: 能量  $E < V(-\infty)$  and  $V(\infty)$ , 波函数可归一化

散射态: 能量  $E > V(-\infty)$  or  $V(\infty)$ , 波函数不可归一化

#### 5. $\delta$ 势阱中波函数满足的条件

波函数连续  $\psi(0^-) = \psi(0^+)$

$$\begin{aligned} \text{波函数的一阶导连续 } & \int_{0^-}^{0^+} \left[ -\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} - \alpha\delta(x)\psi \right] dx = E \int_{0^-}^{0^+} \psi dx \\ & \rightarrow \psi'(0^+) - \psi'(0^-) = -\frac{2\alpha m}{\hbar^2} \psi(0) \end{aligned}$$

#### 6. $V = -\alpha\delta$ 势阱中束缚态的解及函数图像

$$\psi(x) = \frac{\sqrt{m\alpha}}{\hbar} \exp \left[ -m\alpha|x|/\hbar^2 \right], \quad E = -\frac{m\alpha^2}{2\hbar^2}$$

7.  $V = -\alpha\delta$  势阱中反射率  $R$  与透射率  $T$ , 其中左侧波函数为  $e^{ikx} + re^{-ikx}$ , 右侧波函数为  $te^{ikx}$

$$\begin{aligned} \beta &= \frac{\alpha m}{\hbar^2 k}, \quad r = \frac{i\beta}{1 - i\beta}, \quad t = \frac{1}{1 - i\beta} \\ R &= |r|^2 = \frac{\beta^2}{1 + \beta^2}, T = |t|^2 = \frac{1}{1 + \beta^2} \end{aligned}$$

#### 8. 有限深势阱的束缚态的奇函数解

$$\begin{aligned} \psi(x) &= \begin{cases} Be^{-\kappa x} & x > a \\ D \cos(lx), & 0 < x < a \\ \psi(-x) & x < 0 \end{cases} \\ \kappa &= \frac{\sqrt{-2mE}}{\hbar}, l = \frac{\sqrt{2m(E + V_0)}}{\hbar}, \kappa a = la \tan(la) \end{aligned}$$

#### 9. 自由粒子解

$$\begin{aligned} \psi(x, t) &= \sqrt{\frac{1}{L}} e^{ik_n x} e^{-iE_n t/\hbar}, k_n = \frac{2n\pi}{L} = \pm \frac{\sqrt{2mE_n}}{\hbar} \\ \Psi(x, t) &= \frac{1}{\sqrt{2\pi}} \int \phi_k \exp \left[ ikx - i\frac{\hbar k^2}{2m} t \right] dk \end{aligned}$$

### 四、三维空间中的量子力学

#### 1. 氢原子薛定谔方程中的哈密顿量

$$\hat{H} = -\frac{\hbar^2}{2m} \nabla^2 - \frac{e^2}{4\pi\epsilon_0} \frac{1}{r}$$

#### 2. 氢原子薛定谔方程的解

$$\begin{aligned} \hat{H}\psi_{nlm} &= E_n \psi_{nlm}, \quad \psi_{nlm} = R_{nl}(r) Y_l^m(\theta, \phi) \\ n &= 1, 2, 3, \dots, \quad l = 0, 1, \dots, n-1 \quad m = -l, -l+1, \dots, l-1, l \\ E_n &= \frac{E_1}{n^2}, E_1 = \frac{m}{2\hbar^2} \left( \frac{e^2}{4\pi\epsilon_0} \right)^2 = -13.6 \text{eV} \end{aligned}$$

3. (自旋/轨道) 角动量的对易关系

$$[\hat{L}_x, \hat{L}_y] = i\hbar \hat{L}_z, \quad [\hat{L}_y, \hat{L}_z] = i\hbar \hat{L}_x, \quad [\hat{L}_z, \hat{L}_x] = i\hbar \hat{L}_y, \quad [\hat{L}^2, \hat{L}] = 0$$

$\hat{S}$  同理

[\* 请自行从  $L = r \times p$  推导得到]

4. 轨道角动量与自旋角动量的本征方程

$$\begin{aligned}\hat{L}^2|lm\rangle &= \hbar^2 l(l+1)|lm\rangle, & \hat{L}_z|lm\rangle &= \hbar m|lm\rangle \\ l &= 0, 1, 2, \dots & m &= -l, -l+1, \dots, l-1, l \\ \hat{S}^2|sm\rangle &= \hbar^2 s(s+1)|sm\rangle, & \hat{S}_z|sm\rangle &= \hbar m|sm\rangle \\ s &= 0, \frac{1}{2}, 1, \frac{3}{2}, 2, \dots & m &= -s, -s+1, \dots, s-1, s\end{aligned}$$

5. 角动量的上升下降算符，及其对本征波函数的作用

$$\hat{L}_{\pm} = \hat{L}_x \pm i\hat{L}_y, \quad \hat{L}_{\pm}|lm\rangle = \hbar\sqrt{l(l+1)-m(m\pm 1)}|l, m\pm 1\rangle$$

$\hat{S}$  同理

6.  $1/2$  自旋的角动量在本征波函数空间的矩阵表示

$$\begin{aligned}\hat{S} &= \frac{\hbar}{2}\sigma \\ \sigma_x &= \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}\end{aligned}$$

7. 自旋在磁场中受到的相互作用及拉莫尔进动频率

$$\boldsymbol{\mu} = \gamma \boldsymbol{S}, \quad \hat{H} = -\gamma \boldsymbol{B} \cdot \hat{\boldsymbol{S}}, \quad \omega = \gamma B_0$$

8. 角动量的叠加， $\hat{S} = \hat{S}_1 + \hat{S}_2$ , 则  $\hat{S}^2, \hat{S}_z$  ?

$$\hat{S}^2 = \hat{S}_1^2 + \hat{S}_2^2 + 2\hat{S}_1 \cdot \hat{S}_2, \quad \hat{S}_z = \hat{S}_{1z} + \hat{S}_{2z}$$

9. 2 个自旋  $1/2$  的相加

$$\text{三重态}(s=1) : \begin{cases} |1, 1\rangle = |\uparrow\uparrow\rangle \\ |1, 0\rangle = \frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle) \\ |1, -1\rangle = |\downarrow\downarrow\rangle \end{cases}$$

$$\text{单态 } (s=0) : |0, 0\rangle = \frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

10. 角动量叠加后，本征态用 Clebsch-Gordan (CG) 系数和原本的态表示

$$|sm\rangle = \sum_{m_1+m_2=m} C_{s_1 m_1; s_2 m_2}^{sm} |s_1 m_1; s_2 m_2\rangle$$

## 五、全同粒子

### 1. 玻色子和费米子的区别

玻色子: 自旋为整数,  $\psi_+(r_1, r_2) = \psi_+(r_2, r_1)$

费米子: 自旋为半整数,  $\psi_-(r_1, r_2) = -\psi_-(r_2, r_1)$

### 2. 由两个全同粒子组成的波函数为 (考虑自旋与否)

$$\psi_{\pm}(\mathbf{r}_1, \mathbf{r}_2) = A [\psi_a(\mathbf{r}_1) \psi_b(\mathbf{r}_2) \pm \psi_b(\mathbf{r}_1) \psi_a(\mathbf{r}_2)]$$

考虑自旋的玻色子:  $\psi_{\pm}(r_1, r_2) [\chi_a(s_1) \chi_b(s_2) \pm \chi_b(s_1) \chi_a(s_2)]$

考虑自旋的费米子:  $\psi_{\pm}(\mathbf{r}_1, \mathbf{r}_2) [\chi_a(s_1) \chi_b(s_2) \mp \chi_b(s_1) \chi_a(s_2)]$

### 3. 证明费米子满足泡利不相容原理

若  $\psi_a = \psi_b$ , 则  $\psi_- = 0$

### 4. 交换相互作用的表现

在不考虑自旋情况下, 空间对称波函数 (玻色子) 受到牵引力, 空间反对称波函数 (费米子) 受到排斥力

## 六、定态微扰理论

### 1. 非简并微扰 (一阶波函数, 二阶能量, 微扰项为 $\hat{H}'$ )

$$E_n = E_n^0 + \left\langle \psi_n^0 \left| \hat{H}' \right| \psi_n^0 \right\rangle + \sum_{m \neq n} \frac{|\langle \psi_m^0 | \hat{H}' | \psi_n^0 \rangle|^2}{E_n^0 - E_m^0}$$

$$\psi_n = \psi_n^0 + \sum_{m \neq n} \frac{\langle \psi_m^0 | \hat{H}' | \psi_n^0 \rangle}{E_n^0 - E_m^0} \psi_m^0$$

[\* 请自行推导]

### 2. 非简并微扰成立条件

能级非简并, 且  $|\langle \psi_m^0 | \hat{H}' | \psi_n^0 \rangle| \ll |E_n^0 - E_m^0|$

### 3. 近简并微扰

$$\begin{pmatrix} H'_{11} & H'_{12} & \cdot & H'_{1g} \\ \vdots & \vdots & \ddots & \vdots \\ H'_{g1} & H'_{g2} & \cdot & H'_{gg} \end{pmatrix} \begin{pmatrix} a_1 \\ \vdots \\ a_g \end{pmatrix} = E^1 \begin{pmatrix} a_1 \\ \vdots \\ a_g \end{pmatrix}$$

$$H'_{ij} = \langle \psi_{n,i}^0 | \hat{H}' | \psi_{n,j}^0 \rangle, \quad \phi_n = \sum_i a_i \psi_{n,i}$$

### 4. 氢原子哈密顿量的相对论效应修正

$$T = \sqrt{p^2 c^2 + m^2 c^4} - mc^2 \approx \frac{p^2}{2m} - \frac{p^4}{8m^3 c^2} \left( \text{泰勒展开} \right) \rightarrow H' = -\frac{p^4}{8m^3 c^2}$$

5. 氢原子的自旋轨道耦合修正的哈密顿量及对应的好量子数

$$\begin{aligned} B &= \frac{\mu_0 I}{2r} = \frac{1}{4\pi\epsilon_0} \frac{e}{mc^2 r^3} \mathbf{L}, \boldsymbol{\mu}_e = -\frac{e}{m} \mathbf{S} \\ \hat{H}' &= \frac{1}{2} \frac{e^2}{4\pi\epsilon_0} \frac{1}{m^2 c^2 r^3} \hat{\mathbf{S}} \cdot \hat{\mathbf{L}} \\ \left\{ \hat{H}_0, \hat{L}^2, \hat{S}^2, \hat{J}^2, \hat{J}_z \right\}, \hat{\mathbf{J}} &= \hat{\mathbf{L}} + \hat{\mathbf{S}} \end{aligned}$$

6. 外加电场下，氢原子 Stark 效应的微扰哈密顿量

$$\hat{H}' = -qEz = eE\hat{z}$$

7. 氢原子的精细结构能级修正

$$E_n^{(1)} = \frac{E_n^2}{2mc^2} \left( 3 - \frac{4n}{j+1/2} \right)$$

8. 外加磁场下，氢原子 Zeeman 效应的微扰哈密顿量

$$\hat{H}' = -(\boldsymbol{\mu}_L + \boldsymbol{\mu}_S) \cdot \mathbf{B} = \frac{e}{2m} (\hat{\mathbf{L}} + 2\hat{\mathbf{S}}) \cdot \mathbf{B} = \frac{eB}{2m} (\hat{L}_z + 2\hat{S}_z)$$

[\* 请自行推导，在强弱磁场下该如何选择好量子数及对应的能量修正]

七、变分原理与 WKB 近似

1. 变分原理

$$E_{gs} \leq \langle \psi | \hat{H} | \psi \rangle$$

取到等号时即对应基态能量与波函数。

2. WKB 近似公式

$$\text{经典区域 } (E > V) : \psi(x) = \frac{C}{\sqrt{p(x)}} \exp \left[ \pm \frac{i}{\hbar} \int p(x) dx \right], \quad p(x) = \sqrt{2m(E - V(x))}$$

$$\text{隧穿区域 } (E < V) : \psi(x) = \frac{C}{\sqrt{|p(x)|}} \exp \left[ \pm \frac{1}{\hbar} \int p(x) dx \right]$$

八、含时微扰理论

1. 含时微扰理论的系数满足的方程

$$i\hbar \frac{d}{dt} c_m(t) = \sum_n c_n(t) \left\langle m \left| \hat{H}'(t) \right| n \right\rangle e^{i\omega_{mn} t}, \quad \omega_{mn} = \frac{E_m - E_n}{\hbar}$$

2. 2 能级系统含时微扰理论的一级修正 ( $c_m(0) = 0, c_n(0) = 1$ )

$$\begin{aligned} c_m(t) &= \frac{1}{i\hbar} \int_0^t \left\langle m \left| \hat{H}'(t') \right| n \right\rangle e^{i\omega_{mn} t'} dt', \quad c_n(t) = 1 \\ \omega_{mn} &= \frac{E_m - E_n}{\hbar}, \quad P_{n \rightarrow m} = |c_m(t)|^2 \end{aligned}$$

3. 正弦微扰  $H'(\mathbf{r}, t) = V(r) \cos(\omega t)$  ( $\omega_0 + \omega \gg |\omega_0 - \omega|$ ) 下的跃迁几率

$$P_{a \rightarrow b}(t) = \frac{|V_{ab}|^2}{\hbar^2} \frac{\sin^2 [(\omega_0 - \omega) t / 2]}{(\omega_0 - \omega)^2}$$

#### 4. 光与原子相互作用的三种类型及辐射的物理原因

##### 吸收、受激辐射与自发辐射

受激发射：一个光子进入，两个光子出来，包含引起跃迁的原始光子和另一个来自原子的光子、可以通过计算跃迁几率得到。是激光的理论基础。

$$P_{2 \rightarrow 1}(t) = \left( \frac{|e \langle \psi_1 | z | \psi_2 \rangle| E_0}{\hbar} \right)^2 \frac{\sin^2 [(\omega_0 - \omega) t / 2]}{(\omega_0 - \omega)^2}, \quad E_2 - E_1 = \hbar \omega_0$$

自发辐射其实并不是真正的“自发”，而是受到热涨落或电磁波的量子零点涨落（真空涨落）的刺激。

#### 5. 跃迁的选择定则及物理原因

$$\mathbf{P} = e \langle \psi_2 | \mathbf{r} | \psi_1 \rangle \neq 0 \rightarrow \Delta m = 0, \pm 1 \text{ 且 } \Delta l = \pm 1$$

因为光子的自旋为 1，角动量 ( $z$  分量) 守恒要求原子失去的等于光子获得的角动量，角动量的叠加规律只允许  $l' = l + 1, l' = l, l' = l - 1$ ，但对于电偶极辐射， $l' = l$  不会发生。

### 九、散射理论

#### 1. 考虑平面波入射的散射情形下的波函数

$$\psi(r, \theta, \phi) = A \left[ e^{ikx} + f(\theta, \phi) \frac{e^{ikr}}{r} \right]$$

#### 2. 微分散射截面与散射截面

$$\frac{d\sigma}{d\Omega} = |f(\theta, \phi)|^2, \sigma = \int \frac{d\sigma}{d\Omega} d\Omega$$

#### 3. $l$ 分波具有 $\delta_l$ 的相移时，对应的散射截面

$$\sigma = \frac{4\pi}{k^2} \sum_{l=0}^{\infty} (2l+1) \sin^2 (\delta_l)$$

#### 4. 散射理论中的 Born 近似及球对称下的表述

$$\begin{aligned} f(\theta, \phi) &= -\frac{m}{2\pi\hbar^2} \int V(\mathbf{r}_0) \exp [i(\mathbf{k}' - \mathbf{k}) \cdot \mathbf{r}_0] d^3 r_0 \\ f(\theta, \phi) &= -\frac{2m}{\hbar^2 \kappa} \int r V(r) \sin(\kappa r) dr, \quad \kappa = 2k \sin \frac{\theta}{2} \end{aligned}$$

$\theta$  为  $k, k'$  之间的夹角

# **Part V**

## **My toolboxes**

# Chapter 11

## Introduction

I collect all the tools here. I have two approaches to collect these tools, i.e., the bottom-top and top-bottom approach.

The bottom-top one means that I collect it from basic materials such as textbooks or courses. For example, I learn the linear programming in the optimization textbook. Then in the research side, I may use linear programming to solve research problems. In a word, I first have a knife and then use it to cut things. This is different from the textbook learning. In textbook learning, I learn general knowledge and do not know what applications they have. Only after they have applications, I will call them tools. So generally, I put the notes of textbook or course learning in other parts. In those parts, my aim is to build a whole framework of those theories. But in this part, I aim to learn deeper for a particular technique.

The top-bottom one means that I extract it directly from the papers. This means I learn a protocol from a paper to solve a particular problem. Then I extract the fundamental technique from the protocol. This technique will be in my toolbox and I will use it to solve other problems. This approach is more efficient especially for beginners.

## Chapter 12

# Importance Sampling

This chapter collects all the information of importance sampling. This mathematical tool is used to study all the fields in quantum information.

Importance sampling is a very useful topic in statistics. It has numerous applications in science and engineering.

## **Part VI**

# **Quantum Information**

# Chapter 13

## Estimation

### 13.1 DFE

# Chapter 14

## Sensing

### 14.1 Quantum state discrimination with Bosonic channels and Gaussian states

This PhD thesis is written by Sihui Tan, who is the one that first propose Quantum Illumination. So this thesis is a very valuable material to learn QI.

#### 14.1.1 Mathematical preliminaries

**Definition 14.1.1** (Direct sum of matrix).

#### 14.1.2 Physics preliminaries

**Lemma 14.1.1** (Eq. 2.23 in [?]). *A system made of multiple harmonic oscillators is described by their annihilation operators,*

$$\hat{a}_k, \quad k = 1, \dots, n. \quad (14.1.1)$$

*They satisfy the canonical commutation relation*

$$[\hat{a}_k, \hat{a}_l^\dagger] = \delta_{kl}. \quad (14.1.2)$$

**Lemma 14.1.2** (Eq. 2.27 in [?]). *The corresponding quadrature operators of a system made of multiple harmonic oscillators satisfy the canonical commutation relation*

$$[\hat{q}_k, \hat{p}_l] = \frac{i}{2} \delta_{kl}. \quad (14.1.3)$$

**Lemma 14.1.3.** *We describe the phase space of the single bosonic modes with the vector*

$$\mathcal{R} = (\hat{q}, \hat{p})^T, \quad (14.1.4)$$

so  $\mathcal{R}_1 = \hat{q}$  and  $\mathcal{R}_2 = \hat{p}$ . Then we have

$$[\mathcal{R}_k, \mathcal{R}_l] = \frac{i}{2} \omega_{kl}, \quad (14.1.5)$$

where

$$\omega = \begin{bmatrix} 0 & 1 \\ -1 & 0 \end{bmatrix} \quad (14.1.6)$$

is the symplectic form of the phase space of single bosonic mode.

**Lemma 14.1.4.** We describe the phase space of the  $n$  bosonic modes with the vector

$$\mathcal{R} = (\hat{q}^1, \hat{p}^1, \dots, \hat{q}^n, \hat{p}^n)^T. \quad (14.1.7)$$

We denote  $\mathcal{R}_1^i = \hat{q}^i$ ,  $\mathcal{R}_2^i = \hat{p}^i$ , where  $i$  denotes  $i$ th mode. Then we have

$$[\mathcal{R}_k^m, \mathcal{R}_l^n] = \frac{i}{2} \omega_{kl} \delta_{mn}, \quad (14.1.8)$$

where

$$\omega = \begin{bmatrix} 0 & 1 \\ -1 & 0 \end{bmatrix} \quad (14.1.9)$$

is the symplectic form of the phase space of single bosonic mode. Then the symplectic of  $n$  bosonic modes is

$$\Omega = \bigoplus_{k=1}^n \omega, \quad (14.1.10)$$

where  $\bigoplus$  denotes the direct sum.

### 14.1.3 Introduction

#### 14.1.4 Problem formulation

#### 14.1.5 Protocol

#### 14.1.6 Protocol Details

#### 14.1.7 Performance evaluation

## 14.2 Joint Measurement of TFE via SFG [?]

### 14.2.1 Mathematical preliminaries

**Proposition 14.2.1** (The Fourier Transform of 1).

### 14.2.2 Physics preliminaries

**Definition 14.2.1** (Sum Frequency Generation (SFG)). *The SFG process in a  $\chi^{(2)}$  nonlinear medium could be modeled as the following evolution operator:*

$$V = I + \varepsilon \left( \int d\omega_p d\omega_s d\omega_i a_p^\dagger(\omega_p) a_s(\omega_s) a_i(\omega_i) \delta(\omega_p - \omega_s - \omega_i) - H.C. \right), \quad (14.2.1)$$

where photons in the signal mode  $a_s(\omega_s)$  and ideler mode  $a_i(\omega_i)$  are annihilated to generate photons in the pump mode  $a_p(\omega_p)$  and  $\varepsilon$  characterizes the interaction strength. The SPDC process is the time-reversal process of SFG, which can also be described by the same revolution operator.

**Definition 14.2.2** (SPDC).

**Definition 14.2.3** (Evolution operator).

**Theorem 11** (Schmidt' decomposition [1]).

**Definition 14.2.4** (Schmidt's number [1]).

**Definition 14.2.5** (Heisenberg picture).

**Proposition 14.2.2** (The relationship between annihilation operator and frequency).

**Definition 14.2.6** (Time Correlation).

**Definition 14.2.7** (Frequency/Spectral/Energy Correlation).

### 14.2.3 Introduction

SPDC is widely adopted to generate TFE photons. SFG can be used to perform TFE joint-measurement.

[SZQ: 2023.06.03: TFE joint-measurement is just entangled measurement. SPDC is similart to unitary evolution for state preparation, in particular, the time-frequency entangled states. Thus, the reversed process SPDC, i.e., SFG process, is the entangled measurement. This is analogy to bell state preparation and bell measurements.]

### 14.2.4 Problem formulation

The frequency sum and time difference of two photons could be simultaneously measured through the sum-frequency generation process.

### 14.2.5 Protocol

Given the close connection between the spontaneous parametric down-conversion (SPDC) process and time-frequency entanglement (TFE), it's natural to utilize the time-reversal of the SPDC process, i.e., sum frequency generation (SFG) to obtain a TFE joint measurement based protocol.

**Definition 14.2.8** (Frequency sum (FS) operator). *The frequency sum operator  $P_{\delta_\omega}(\omega)$  that selects states with the frequency sum  $\omega_s + \omega_i$  of the signal and idler photon being around  $\omega$  within uncertainty  $\delta_\omega$  is defined as:*

$$P_{\delta_\omega}(\omega) = \int \int d\omega_s d\omega_i a_s^\dagger(\omega_s) a_i^\dagger(\omega_i) a_s(\omega_s) a_i(\omega_i) \text{Gate}\left(\frac{\omega - \omega_s - \omega_i}{\delta_\omega}\right), \quad (14.2.2)$$

where  $\text{Gate}(x) = 1$  for  $|x| \leq 1/2$  and  $\text{Gate}(x) = 0$  otherwise.

**Lemma 14.2.1** (Frequency sum operator is a projection operator). *The frequency sum operator  $P_{\delta_\omega}(\omega)$  is a projection operator satisfying*

$$P_{\delta_\omega}(\omega)^2 = P_{\delta_\omega}(\omega). \quad (14.2.3)$$

**Definition 14.2.9** (Time difference (TD) operator). *The time difference operator  $P_{\delta_t}(t)$  that selects states with the time difference  $t_s - t_i$  of the signal and idler photon being around  $t$  within uncertainty  $\delta_t$  is defined as:*

$$P_{\delta_t}(t) = \int \int dt_s dt_i \tilde{a}_s^\dagger(t_s) \tilde{a}_i^\dagger(t_i) \tilde{a}_s(t_s) \tilde{a}_i(t_i) \text{Gate}\left(\frac{t_s - t_i - t}{\delta_t}\right), \quad (14.2.4)$$

where  $\text{Gate}(x) = 1$  for  $|x| \leq 1/2$  and  $\text{Gate}(x) = 0$  otherwise and

$$\tilde{a}_x = \frac{1}{\sqrt{2\pi}} \int d\omega \exp(-i\omega t) a_x(\omega) \quad (14.2.5)$$

**Lemma 14.2.2** (Time difference operator is a projection operator). *The time difference operator  $P_{\delta_t}(t)$  is a projection operator satisfying*

$$P_{\delta_t}(t)^2 = P_{\delta_t}(t). \quad (14.2.6)$$

**Definition 14.2.10** (Joint projection operator). *The joint projection operator of the time difference and frequency sum is defined as*

$$P_{\delta_\omega, \delta_t}(\omega, t) = P_{\delta_\omega}(\omega) P_{\delta_t}(t), \quad (14.2.7)$$

which means selecting states of which the time difference between the signal and idler photon  $t_s - t_i$  is around  $t$  within uncertainty  $\delta_t$  and frequency sum of the signal and idler photon being around  $\omega$  within uncertainty  $\delta_\omega$ , simultaneously.

**Lemma 14.2.3** (Commutation relationship between frequency-time operators). *We have the commutation relationship*

$$[P_{\delta_\omega}(\omega), P_{\delta_t}(t)] = 0. \quad (14.2.8)$$

**Definition 14.2.11** (TD an FS probability density operator (PDF)). *We define*

$$P(\omega, t) = \lim_{\delta_t \rightarrow 0, \delta_\omega \rightarrow 0} \frac{1}{\delta_\omega \delta_t} P_{\delta_\omega}(\omega) P_{\delta_t}(t). \quad (14.2.9)$$

**Lemma 14.2.4.** *We have*

$$P(\omega, t) = \frac{1}{2\pi} B_p^\dagger B_p, \quad (14.2.10)$$

where

$$B_p = \int \int d\omega_s d\omega_i \delta(\omega_s + \omega_i - \omega) \exp[i\omega_i t] a_s(\omega_s) a_i(\omega_i). \quad (14.2.11)$$

**Lemma 14.2.5** (Connection between TD and FS PDF and SFG Process). *We have*

$$B^\dagger B = 2\pi P(\omega, 0). \quad (14.2.12)$$

This lemma did not show the SFG process

**Lemma 14.2.6** (Discrete sum of evolution operator of SFG process). *The discrete sum of the evolution operator of SFG process can be obtained by a two-step Schmidt decomposition as:*

$$v = I + \varepsilon \sum_m \left( \sqrt{\lambda_m^{(1)}} A_m^\dagger B_m - H.C. \right), \quad (14.2.13)$$

where

$$B_m = \sum_n \sqrt{\lambda_{m,n}^{(2)}} F_{m,n} G_{m,n}, \quad (14.2.14)$$

$$A_m = \int d\omega \psi_{A,m}(\omega) a_p(\omega), \quad (14.2.15)$$

$$B_m = \int d\omega_s d\omega_i \psi_{B,m}(\omega_s, \omega_i) a_s(\omega_s) a_i(\omega_i), \quad (14.2.16)$$

$$F_{m,n} = \int d\omega \psi_{F,m,n}(\omega) a_s(\omega), \quad (14.2.17)$$

$$G_{m,n} = \int d\omega \psi_{G,m,n}(\omega) a_i(\omega). \quad (14.2.18)$$

**Lemma 14.2.7** (Non-uniqueness of the first step Schmidt Decomposition). *If the function  $f_0$  can be written in the following form:*

$$\delta(\omega_p - \omega_s - \omega_i) f_0(\omega_p - \omega_s - \omega_i) = \delta(\omega_p - \omega_s - \Omega_i) f\left(\frac{\omega_s - \Omega_i}{\sqrt{2}}\right), \quad (14.2.19)$$

then the first step Schmidt decomposition in the main text is not unique.

**Lemma 14.2.8** (An useful commutation relationship). *An useful commutation relationship:*

$$[B_{m'}, B_{m''}^\dagger] = \delta_{m'm''} + \int d\omega'_s d\omega_i d\omega''_s \psi_{B,m''}^*(\omega'_s, \omega_i) a_s^\dagger(\omega''_s) a_s(\omega'_s) \quad (14.2.20)$$

$$+ \int d\omega_s d\omega'_i d\omega''_i \psi_{B,m''}^*(\omega_s, \omega''_i) \psi_{B,m'} a_i^\dagger(\omega''_i) a_i(\omega'_i). \quad (14.2.21)$$

**Corollary 14.2.1.** *By (14.2.8), we have*

$$[B_{m'}, B_{m''}^\dagger] = \delta_{m'm''} |0\rangle. \quad (14.2.22)$$

**Lemma 14.2.9.** *We have the commutation relation between time difference projection operator and frequency sum projection operator:*

$$[P_{\delta_\omega}(\omega), P_{\delta_t}(t)] = 0. \quad (14.2.23)$$

**Lemma 14.2.10.** *The frequency spectrum  $S(\omega)$  of the generated pump photon is given by the expectation value of the spectral density operator  $a_p^\dagger(\omega_p) a_p(\omega_p)$*

$$S(\omega_p) = \frac{\epsilon^2 \exp \left[ \frac{1}{8} \left( -4\Delta t^2 \sigma_-^2 - \frac{\Delta \omega^2}{\sigma_-^2} - \frac{4(\Delta \omega + \omega_0 - \Omega_p)^2}{\sigma_+^2} \right) \right]}{2\sqrt{\pi} \sigma_+}. \quad (14.2.24)$$

**Definition 14.2.12.** *The discrete mode operator  $F_{m,n}^{(b)}$  for the noise photons is defined as*

$$F_{m,n}^{(b)} = \int d\omega \psi_{F_{m,n}}(\omega) a_s^{(b)}(\omega). \quad (14.2.25)$$

**Definition 14.2.13.** *The virtual beam-splitter is modeled as the following unitary transform:*

$$U_{loss} = \Pi_n \exp \left[ i \arccos(\eta) (F_{0,n}^\dagger F_{0,n}^{(b)} + H.C.) \right]. \quad (14.2.26)$$

**Definition 14.2.14.** *We use a density matrix  $\rho_b$  that satisfies the following conditions to describe the noise photons:*

$$\text{Tr}[F_{0,n''}^{(b)\dagger} F_{0,n'}^{(b)} \rho_b] = \delta_{n',n''} \mu_b, \quad (14.2.27)$$

$$\text{Tr}[F_{0,n'}^{(b)} \rho_b] = 0 \quad (14.2.28)$$

**Definition 14.2.15.** *The signal and idler photon pair source is described by the biphoton state  $|pair\rangle$ :*

$$|pair\rangle = B_0^\dagger |0\rangle \quad (14.2.29)$$

$$= \sum_n \sqrt{\lambda_{0,n}^{(2)}} F_{0,n}^\dagger G_{0,n} |0\rangle. \quad (14.2.30)$$

**Definition 14.2.16.** *The unitary transform of the SFG process is given by:*

$$V = I + \epsilon \sum_m \left[ \sqrt{\lambda_m^{(1)}} A_m^\dagger B_m - H.C. \right]. \quad (14.2.31)$$

**Definition 14.2.17.** *In the Heisenberg picture, the photon number operator of the generated pump photon in each pump mode  $A_m$  after the beam-splitter transform and the SFG process is given by:*

$$U_{loss}^\dagger V^\dagger A_m^\dagger V U_{loss}. \quad (14.2.32)$$

**Proposition 14.2.3.** *When the transmission of the signal photon is perfect ( $\eta = 1$ ), the pump photon can only generate in mode  $A_0 (m = 0)$ .*

**Lemma 14.2.11.** *We have*

$$\langle U_{loss}^\dagger V^\dagger A_0^\dagger A_0 V | U_{loss} \rangle = \epsilon^2 \lambda_0^{(1)} (\eta + \mu_b \sum_n \lambda_{0,n}^2). \quad (14.2.33)$$

**Definition 14.2.18.** *The generated SPDC state is given by:*

$$V = |0\rangle - \epsilon \sqrt{\lambda_0^{(1)}} \alpha B_0^\dagger |0\rangle. \quad (14.2.34)$$

**Lemma 14.2.12.** *The joint density operator of the noise-idler state  $\rho_j$  is given by the tensor product of  $\rho_i$  and  $\rho_b$ :*

$$\rho_j = \rho_i \otimes \rho_b \quad (14.2.35)$$

$$= \mu_b \int \int d\omega_s d\omega'_s \int \int d\omega'_i d\omega''_i \phi_0^*(\omega'_s, \omega'_i) \phi_0(\omega'_s, \omega''_i) a_i^\dagger(\omega'_i) a_s^\dagger(\omega_s) |0\rangle \langle 0| a_i(\omega''_i) a_s(\omega_s). \quad (14.2.36)$$

**Lemma 14.2.13.** *The spectral density  $S(\omega)$  of the upconverted photons is*

$$S(\omega) = \frac{\epsilon^2 \mu_b \exp \left[ -\frac{(\omega - \omega_0)^2}{8\sigma_-^2 - 2\sigma_+^2} \right]}{\sqrt{\pi} \sqrt{4\sigma_-^2 + \sigma_+^2}}. \quad (14.2.37)$$

**Theorem 12.** *The error exponent of the classical Chernoff bound of the TFE QI protocol is given by  $C_{QI}$  [2]:*

$$C_{QI} = -\log \min_{s \in [0,1]} \left\{ \sum_{b \in 0,1} p_0(b)^s p_1(b)^{(1-s)} \right\}. \quad (14.2.38)$$

## 14.2.6 Performance evaluation

**Lemma 14.2.14.**

$$(14.2.39)$$

### 14.2.7 Ideas

## 14.3 Optimum Mixed-State Discrimination for Noisy Entanglement-Enhanced Sensing [?]

### 14.3.1 Mathematical preliminaries

### 14.3.2 Physics preliminaries

### 14.3.3 Introduction

### 14.3.4 Problem formulation

### 14.3.5 Protocol

### 14.3.6 Protocol Details

### 14.3.7 Performance evaluation

**Lemma 14.3.1.**

(14.3.1)

### 14.3.8 Ideas

## 14.4 Quantum Estimation Methods for Quantum Illumination [?]

### 14.4.1 Mathematical preliminaries

**Definition 14.4.1** (Quantum Fisher Information (QFI)). *Define a quantum state  $\rho_\eta$  parameterized by  $\eta$ . Then the QFI for  $\rho_\eta$  is*

$$H := 2 \sum_{mn} \frac{|\langle \phi_m | (\partial_\eta \rho_\eta|_{\eta=0}) | \phi_n \rangle|^2}{\lambda_m + \lambda_n}, \quad (14.4.1)$$

where  $\lambda_n$  is the eigenvalue of  $\rho_{\eta=0}$  corresponding to the eigenstate  $|\phi_n\rangle$ , and the derivative is evaluated as  $\eta = 0$ .

**Definition 14.4.2** (Cramér-Rao Bound (CROB) in Eq. (2) [?]). *The CROB asserts that the limits on the achievable precision of an unbiased estimator  $\tilde{\eta}$  is*

$$\Delta \tilde{\eta}^2 \geq \frac{1}{MH}, \quad (14.4.2)$$

where  $H$  is the QFI (??) of  $\rho_\eta$  and  $M$  is the number of copies.

**Theorem 13** (Cramér-Chernoff theorem [?]).

#### 14.4.2 Physics preliminaries

#### 14.4.3 Introduction

#### 14.4.4 Problem formulation

The quantum illumination problem is modeled as a **Reflectivity Estimation Problem** as the following.

Let us consider a general bipartite pure state representation of signal-ideler system written in the Schmidt decomposition form

$$|\psi\rangle_{SI} = \sum_{\alpha} \sqrt{p_{\alpha}} |w_{\alpha}\rangle_S |v_{\alpha}\rangle_I. \quad (14.4.3)$$

In the Quantum Illumination (QI) protocol, the signal modes of  $M$  copies of  $|\psi\rangle_{SI}$  are sent to the target region embedded in a bright thermal noise, in which there could possibly be an object. We receive  $M$  copies of the state

$$\rho_{\eta} = \text{Tr} \left[ U_{\eta} \psi \bigotimes \rho_B U_{\eta}^{\dagger} \right], \quad (14.4.4)$$

where

$$U_{\eta} := \exp[\sin^{-1}(\eta) - s^{\dagger}b - sb^{\dagger}] \simeq \exp[\eta(s^{\dagger}b - sb^{\dagger})] \quad (14.4.5)$$

is the signal-object interaction, modeled as a beamsplitter with amplitude reflectivity  $\eta \ll 1$ , and

$$\rho_B := \sum_n \frac{N_B^n}{(1 + N_B)^{1+n}} |n\rangle \langle n| \quad (14.4.6)$$

is the thermal state with mean photon number  $N_B$ .

In this framework, at the receiver side,  $\eta = 0$  corresponds to the absence of the object in the target region. We note  $\eta$  is unknown parameter. So we measurement  $\rho_{\eta}$  many times and get the estimation  $\tilde{\eta}$ . We then use  $\hat{\eta}$  to determine whether is target is absent or present. This lie in the Estimation problem and state discrimination problem!

#### 14.4.5 Protocol

The quantum illumination protocol has been presented the last subsection. Here, we focus on the estimation protocol.

The main idea of the estimation protocol is to use QFI (??) and CROB (??) for estimating  $\eta$ .

**Lemma 14.4.1** (The QFI for the returned state  $\rho_\eta$  in Eq. (3) [?]). *The QFI for  $\rho_\eta$  is*

$$H := \frac{4}{1 + N_B} \sum_{\alpha\alpha'} \frac{p_\alpha p_{\alpha'}}{p_{\alpha'} + p_\alpha \frac{N_B}{1+N_B}} |\langle w_{\alpha'} | s | w_\alpha \rangle|^2. \quad (14.4.7)$$

*Proof.* □

We choose the observable that optimize the QFI of the returned state (??) and measure  $\rho_\eta$ , the outcome is  $o_i$ . Then by the strong law of large numbers, we perform the sample mean as

$$\tilde{\eta} = \frac{1}{M} \sum_{i=1}^M o_i. \quad (14.4.8)$$

---

**Algorithm 1** QI estimation protocol

**Input:**  $|\psi_{SI}\rangle$ : the signal-idler state,

$\eta$ : the reflectivity parameter,

$\rho_B$ : the thermal state

$M$ : the number of copies of  $|\psi_{SI}\rangle$

**Output:** Estimation of the reflectivity  $\tilde{\eta}$ .

- 1: **for**  $i = 1, \dots, M$  **do**
  - 2:   Sent the signal mode of  $\psi_{SI}^{(1)}$  to the target region,
  - 3:   Rceive the state  $\rho_\eta^{(1)}$ ,
  - 4:   **Choose the observable that optimize the QFI**,
  - 5:   Measure the observable and get the outcome  $o_i$ .
  - 6: **end for**
  - 7: Calculate  $\tilde{\eta}$ (14.4.8).
  - 8: Output the estimated reflectivity parameter  $\tilde{\eta}$ .
- 

After that, we use the estimation  $\tilde{\eta}$  for state discrimination.

**Algorithm 2** QI discrimination protocol

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**Input:**  $\tilde{\eta}$ : the estimation reflectivity,

$\xi$ : the scaling parameter.

**Output:** Absent or present.

- 1: **if**  $\tilde{\eta} > \xi\eta$  **then**
  - 2:   Output present.
  - 3: **else if**  $\tilde{\eta} \leq \xi\eta$  **then**
  - 4:   Output absent.
  - 5: **else**
  - 6:   Output 'Wrong input!'
  - 7: **end if**
- 

#### 14.4.6 Protocol Details

#### 14.4.7 Performance evaluation

**Theorem 14** (Type I-II error probabilities Theorem in [?]). Let  $|\psi\rangle_{SI} = \sqrt{p_\alpha}|n\rangle_S|v_n\rangle_I$  be the Schmidt decomposition of the signal-idler state, and denote  $\rho_S$  the state of the signal. Then  $P_{I,II} \sim \exp\left(-\frac{\eta_{I,II}^2 HM}{2}\right)$  provided that  $\exists C > 0$  s.t.  $\langle s^k s^{\dagger k} \rangle_{\rho_S} \leq k!C^k, \forall k \in \mathcal{N}$ .

*Proof.* □

#### 14.4.8 Examples

**Example 14.4.1** (Gaussian states).

**Example 14.4.2** (Schrödinger's cat state).

#### 14.4.9 Main ideas summarized

This subsection summarizes the main ideas in this paper.

Choose the observable that optimize the QFI.

#### 14.4.10 Ideas

How many copies of  $\psi_{SI}$  is needed to achieve a given confidence interval? This is a problem about the quantum resources?

# Bibliography

- [1] Michael A. Nielsen and Isaac L. Chuang. *Quantum computation and quantum information*. Cambridge University Press, 2019.
- [2] Michael Nussbaum and Arleta Szkoła. The chernoff lower bound for symmetric quantum hypothesis testing. 2009.

# **Chapter 15**

## **Imaging**

**15.1 Quantum and non-local effects offer over 40 dB noise resilience advantage towards quantum lidar**