# Class Test 1

## KDSMIL001 MAM2046W 2BP

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## 1. Heat Distribution

Given the following:

$$\begin{split} \nabla^2 T &= \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial T}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 T}{\partial \theta^2} = 0, \ 0 < r < 30, \ 0 < \theta < \frac{\pi}{2} \\ T(30,\theta) &= 50, \ 0 < \theta < \frac{\pi}{2} \\ |T(0,\theta)| &< \infty, \ 0 < \theta < \frac{\pi}{2} \\ T(r,0) &= 0, \ 0 < r < 30 \\ T(r,\frac{\pi}{2}) &= 0, \ 0 < r < 30 \end{split}$$

(a) We separate  $T(r, \theta)$  into two functions of only r and  $\theta$ :

$$T(r,\theta) = R(r)\Theta(\theta)$$

Our ICs and BCs become

$$R(30) = 50$$
$$|R(0)| < \infty$$
$$\Theta(0) = 0$$
$$\Theta(\frac{\pi}{2}) = 0$$

and our initial PDE then becomes

$$0 = \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial R\Theta}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 R\Theta}{\partial \theta^2}$$

$$\implies \frac{\Theta}{r} \frac{d}{dr} \left( r \frac{dR}{dr} \right) = -\frac{R}{r^2} \frac{d^2 \Theta}{d\theta^2}$$

$$\implies \frac{1}{R} \left[ r^2 \frac{d^2 R}{dr^2} + r \frac{dR}{dr} \right] = -\frac{1}{\Theta} \frac{d^2 \Theta}{d\theta^2}$$

We then set each ODE to the same constant  $\lambda$ , since they are equal, giving us

$$\frac{d^2\Theta}{d\theta^2} = -\lambda\Theta$$
$$r^2 \frac{d^2R}{dr^2} + r\frac{dR}{dr} - \lambda R = 0$$

Starting with  $\Theta$ , we know the solutions to this type of equation. For  $\lambda > 0$  we have

$$\Theta(\theta) = A\sin(\sqrt{\lambda}\theta) + B\cos(\sqrt{\lambda}\theta)$$

Applying the BCs we get

$$\Theta(0) = A * 0 + B * 1 = 0$$

$$\Longrightarrow B = 0$$

$$\Theta(\frac{\pi}{2}) = A \sin(\sqrt{\lambda} \frac{\pi}{2}) = 0$$

$$\Longrightarrow \sqrt{\lambda} \frac{\pi}{2} = n\pi, \ n \in \mathbb{N}$$

$$\Longrightarrow \lambda = 4\pi^2, \ n \in \mathbb{N}$$

And so we have a solution to the  $\Theta$  ODE for  $\lambda > 0$ :

$$\Theta(\theta) = A\sin(2n\theta), \ n \in \mathbb{N}$$

We know that this equation has no solution for  $\lambda < 0$ , and the solution for  $\lambda = 0$  is  $\Theta = C$ ,  $C \in \mathbb{R}$ .

Moving on to the R ODE, we have

$$r^2 \frac{d^2 R}{dr^2} + r \frac{dR}{dr} - \lambda R = 0$$

which is an Euler ODE, the standard ansatz for which is  $R = r^s$ , r > 0, so it's always valid in our case. Using this ansatz we find

$$r^{2}[s(s-1)r^{s-2}] + r[sr^{s-1}] - 4n^{2}r^{s} = 0$$

$$\implies s^{2} - 4n^{2} = 0$$

$$\implies s = \pm 2n$$

For an Euler ODE the solutions usually depend on n, but in our case n > 0 so we only examine that case. For  $n \neq 0$ , Euler ODEs have the solution:

$$R(r) = Dr^{2n} + Er^{-2n}$$

Here is where our boundedness condition comes in. If we examine the solution as  $r \to 0$ , we can see that it explodes to  $\infty$ , but since we know that  $|R(0)| < \infty$ , we must have that E = 0 to stop this blowing up. So for  $n \neq 0$ 

$$R(r) = Dr^{2n}$$

Recall we had  $T = R\Theta$ , so our solution is

$$T_n(r,\theta) = A\sin(2n\theta)Dr^{2n}$$
$$= C\sin(2n\theta)r^{2n}$$

And we have the infinite series representation:

$$T(r,\theta) = \sum_{n=1}^{\infty} C_n \sin(2n\theta) r^{2n}$$

(b) Applying the final BC, we have

$$T(30,\theta) = 50 = \sum_{n=1}^{\infty} C_n \sin(2n\theta) 30^{2n}$$

$$\implies 50 \sin(2m\theta) = \sum_{n=1}^{\infty} C_n 30^{2n} \sin(2n\theta) \sin(2m\theta)$$

$$\implies \int_0^{\frac{\pi}{2}} 50 \sin(2m\theta) d\theta = \sum_{n=1}^{\infty} C_n 30^{2n} \int_0^{\frac{\pi}{2}} \sin(2n\theta) \sin(2m\theta) d\theta$$

By orthogonality relations, the  $\int_0^{\frac{\pi}{2}} \sin(2n\theta) \sin(2m\theta) d\theta$  term goes to 0 for all  $n \neq m$ , which simplifies it down to

$$\int_0^{\frac{\pi}{2}} 50\sin(2m\theta)d\theta = C_m 30^{2m} \int_0^{\frac{\pi}{2}} \sin^2(2m\theta)d\theta$$
$$= C_m 30^{2m} \frac{\pi}{4}$$
$$\implies C_m = \frac{4}{30^{2m} \pi} \int_0^{\frac{\pi}{2}} 50\sin(2m\theta)d\theta$$

#### 2. Fourier Series

Given the function:

$$f(x) = \begin{cases} 0 & -\pi \le x < 0 \\ x^2 & 0 \le x \le \pi \end{cases}$$

(a) The Fourier Series of f(x) will look like the following

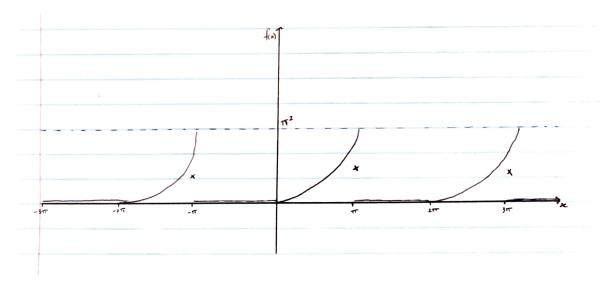


Figure 1: Fourier Series plot of f(x)

(b) The Fourier Coefficients are given by

$$A_0 = \frac{1}{2L} \int_{-L}^{L} f(x) dx$$

$$A_m = \frac{1}{L} \int_{-L}^{L} f(x) \cos\left(\frac{m\pi x}{L}\right) dx$$

$$B_m = \frac{1}{L} \int_{-L}^{L} f(x) \sin\left(\frac{m\pi x}{L}\right) dx$$

Starting with  $A_0$ :

$$A_{0} = \frac{1}{4\pi} \left[ \int_{-\pi}^{0} 0 dx + \int_{0}^{\pi} x^{2} dx \right]$$

$$= \frac{1}{4\pi} \left[ \frac{x^{3}}{3} \right]_{0}^{\pi}$$

$$= \frac{1}{4\pi} \frac{\pi^{3}}{3}$$

$$A_{0} = \frac{\pi^{2}}{12}$$

Now

$$A_{m} = \frac{1}{2\pi} \left[ \int_{-\pi}^{0} 0 dx + \int_{0}^{\pi} x^{2} \cos\left(\frac{mx}{2}\right) dx \right]$$
$$= \frac{1}{2\pi} \int_{0}^{\pi} x^{2} \cos\left(\frac{mx}{2}\right) dx$$
$$= \frac{(\pi^{2}m^{2} - 8) \sin\left(\frac{\pi m}{2}\right) + 4\pi m \cos\left(\frac{\pi m}{2}\right)}{m^{3}\pi}$$

And

$$B_{m} = \frac{1}{2\pi} \left[ \int_{-\pi}^{0} 0 dx + \int_{0}^{\pi} x^{2} \sin\left(\frac{mx}{2}\right) dx \right]$$

$$= \frac{1}{2\pi} \int_{0}^{\pi} x^{2} \sin\left(\frac{mx}{2}\right) dx$$

$$= -\frac{\left((\pi^{2}m^{2} - 8)\cos\left(\frac{\pi m}{2}\right) - 4\pi m \sin\left(\frac{\pi m}{2}\right) + 8\right)}{m^{3}\pi}$$

Since the original function is piecewise-smooth but has jump discontinuities, we know that the Fourier Series will converge to f(x) within the bounds of the domain, and will converge to  $\frac{pi^2}{2}$  at the discontinuities.

(c) The Fourier Series can not be differentiated term by term as

$$f(-\pi) = 0 \neq f(\pi) = \pi^2$$