NSFD Schemes for the Heat Equation

Miles Smith, 2023

July 27, 2022

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Structure

- Introduce the Heat Equation
- Analytic Solution
- Numerical Solutions
- NSFD Approach
- Ourrent Progress and Next Steps

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The Heat Equation In \mathbb{R}^1 :

$$\frac{\partial u}{\partial t} = \alpha \frac{\partial^2 u}{\partial x^2}$$

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The Heat Equation In \mathbb{R}^1 :

$$\frac{\partial u}{\partial t} = \alpha \frac{\partial^2 u}{\partial x^2}$$

What? A *Partial* Differential Equation (PDE), that describes heat, u(x,t), in terms of time, t, and space, x α represents thermal diffusivity

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The Heat Equation In \mathbb{R}^1 :

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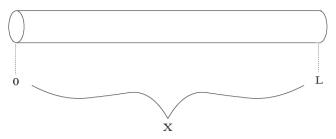
Solution

Heat Flow: How heat spreads out over time



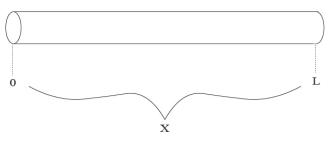
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What 1-D Means



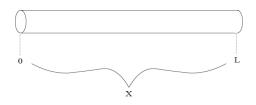
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What 1-D Means



Question: If this rod starts out with some initial heat distribution, how does heat flow out of it over time? (What's the rod's heat in 1 second?)

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Question: how does heat flow over time?

To Answer...

Initial heat distribution (IC): u(x,0) = f(x)

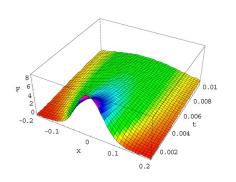
Heat at end-points (BC): u(0, t) = ? u(L, t) = ?

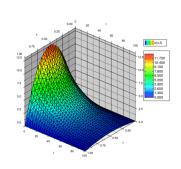


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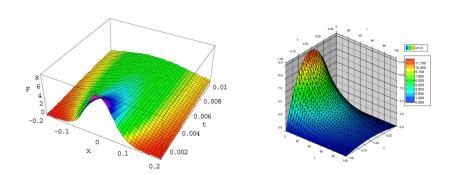
Heat Equation Visualization





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Heat Equation Visualization



Here, u(x,0) is the initial heat distribution and u(0,t) = u(L,t) = 0

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Methods

Fourier Series (traditional)

Green's Function

Dirac-delta function for fundamental solution

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Methods

Fourier Series (traditional)

Green's Function

Dirac-delta function for fundamental solution

Let
$$u(0, t) = u(L, t) = 0$$

and $u(x, 0) = f(x)$

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Methods

Fourier Series (traditional)

Green's Function

Dirac-delta function for fundamental solution

Let
$$u(0, t) = u(L, t) = 0$$

and $u(x, 0) = f(x)$

The Answer...

$$u(x,t) = \sum_{n=1}^{\infty} b_n \cdot \sin\left(\frac{n\pi x}{L}\right) \cdot e^{-\left(\frac{\alpha n\pi}{L}\right)^2 t}$$
$$b_n = \frac{2}{L} \int_0^L f(x) \cdot \sin\left(\frac{n\pi x}{L}\right) dx$$

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The Answer...

$$u(x,t) = \sum_{n=1}^{\infty} b_n \cdot \sin\left(\frac{n\pi x}{L}\right) \cdot e^{-\left(\frac{\alpha n\pi}{L}\right)^2 t}$$
$$b_n = \frac{2}{L} \int_0^L f(x) \cdot \sin\left(\frac{n\pi x}{L}\right) dx$$

Problem: This is complicated... and what if f(x) is complicated and makes it difficult to integrate?



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General Problem

Let
$$L=1$$
, $T=1$



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General Problem

Let
$$L = 1$$
, $T = 1$
BC: $u(0, t) = u(L, t) = 0$
IC: $u(x, 0) = sin(\pi x)$



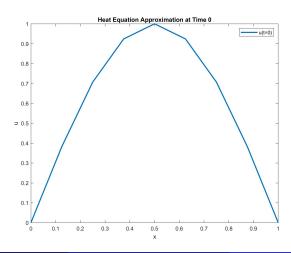
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General Problem

Let L = 1, T = 1

BC: u(0,t) = u(L,t) = 0

IC: $u(x,0) = sin(\pi x)$



Numerical Solutions to the Heat Equations

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Numerical Solutions to the Heat Equations

Solution Methods:

- Explicit Euler
- Implicit Euler
- Crank-Nicolson
- Thomas Algorithm

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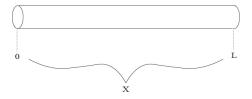
Recall:

$$\frac{df(x)}{dx} \approx \frac{f(x+h) - f(x)}{h}$$

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Recall:

$$\frac{df(x)}{dx} \approx \frac{f(x+h) - f(x)}{h}$$

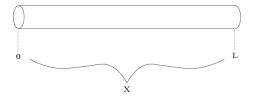


If L=1, and we want to find the change in heat around x=.5 at t=0 with respect to *time*.

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Recall:

$$\frac{df(x)}{dx} \approx \frac{f(x+h) - f(x)}{h}$$



If L=1, and we want to find the change in heat around x=.5 at t=0 with respect to $\it time.$ Then,

$$u_t(.5,0) \approx \frac{u(.5,0+\Delta t) - u(.5,0)}{\Delta t}$$

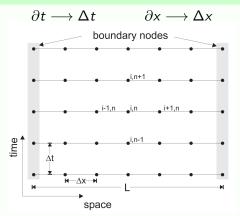
M. Smith Heat Equation, NSFD

If both \times and t go from 0 to 1, and we want to break them into discrete chunks...

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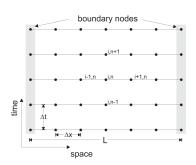
If both x and t go from 0 to 1, and we want to break them into discrete chunks...

Finite Difference Mesh



This mesh includes approximations of heat, Δx and Δt apart, for all values of x and t.

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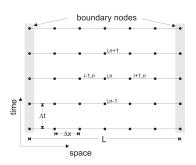


Bottom line:

$$u^{j=1} = \begin{bmatrix} u_{i=1}^{j=1} & u_{i=2}^{j=1} & \dots & u_{i=7}^{j=1} \end{bmatrix}$$

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Bottom line:

$$u^{j=1} = \begin{bmatrix} u_{i=1}^{j=1} & u_{i=2}^{j=1} & \dots & u_{i=7}^{j=1} \end{bmatrix}$$
$$u_{i=5}^{j=2} \approx u(x = .5, t = .2)$$

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$$u_t(.5,0) \approx \frac{u(.5,0+\Delta t) - u(.5,0)}{\Delta t}$$



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$$u_t(.5,0) \approx \frac{u(.5,0+\Delta t) - u(.5,0)}{\Delta t}$$

Discretize Derivatives Using Finite Difference

$$\frac{\partial u}{\partial t} \approx \frac{u_i^{j+1} - u_i^j}{\Delta t}$$

and

$$\frac{\partial^2 u}{\partial x^2} \approx \frac{u_{i+1}^j - 2u_i^j + u_{i-1}^j}{\Delta x^2}$$

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Numerical Set-Up: (3) Discrete Difference Equation

Bringing pieces together...

$$\frac{u_i^{j+1} - u_i^j}{\Delta t} = \alpha \frac{u_{i+1}^j - 2u_i^j + u_{i-1}^j}{\Delta x^2}$$

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Numerical Set-Up: (3) Discrete Difference Equation

Bringing pieces together...

$$\frac{u_{i}^{j+1} - u_{i}^{j}}{\Delta t} = \alpha \frac{u_{i+1}^{j} - 2u_{i}^{j} + u_{i-1}^{j}}{\Delta x^{2}}$$

Rearranged:

Discrete Heat Equation

$$u_i^{j+1} = (1 - 2R)u_i^j + R\left(u_{i+1}^j + u_{i-1}^j\right)$$

$$R = \frac{\Delta t \cdot \alpha}{\Delta x^2}$$

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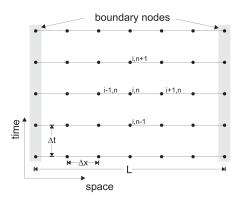
Numerical Set-Up: (4) Solution Methods

- Explicit Euler
- 2 Implicit Euler



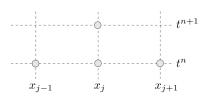
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$$u_i^{j+1} = (1-2R)u_i^j + R\left(u_{i+1}^j + u_{i-1}^j\right)$$



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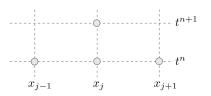
$$u_i^{j+1} = (1 - 2R)u_i^j + R\left(u_{i+1}^j + u_{i-1}^j\right)$$





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$$u_i^{j+1} = (1-2R)u_i^j + R\left(u_{i+1}^j + u_{i-1}^j\right)$$



$$u_3^4 = (1 - 2R)u_3^3 + R(u_4^3 + u_2^3) \approx u(x = .5, t = .05)$$

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How do you create each time-step vector?

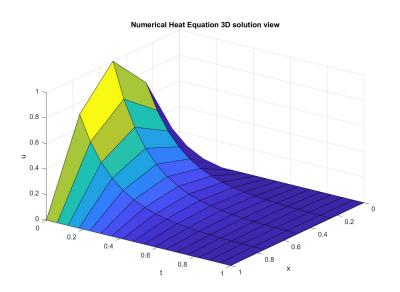
Explicit For-Loop

```
43 -
      a = 0(t) 0;
                                %boundary condition x=0
44 -
    b = @(t) 0;
                                %boundary condition x=L
45-
    g = Q(x) \sin(pi.*x); %initial condition, t=0
46
47 -
     u=zeros(M+1, K+1);
48
49-
     u(1,:)=a(t'); %input BCs
50 -
      u(end,:)=b(t');
51 -
     u(:,1)=g(x); %input ICs
52
53
      % Forward Difference Scheme
54 -
    \exists for i=1:K
55 -
    for i=2:M
56-
              u(i, j+1) = (1-2*R)*u(i,j) + R*(u(i+1,j) + u(i-1,j));
57 -
          end
58 -
     --end
```

K and M are the number of steps for time and space, respectively

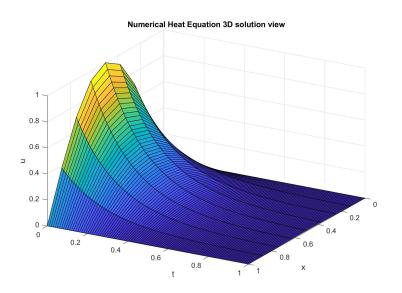
4 D F 4 B F 4 B F

Explicit Euler Example: Simple Mesh



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Explicit Euler Example: Refined Mesh



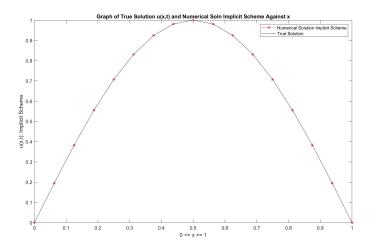
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Approximate True Solution

How do we know if our approximation is good?

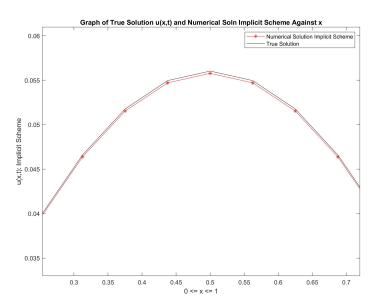
True Solution Baseline

True versus Explicit Euler: First Time Step



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True versus Explicit Euler: 300th Time Step



Order of Convergence:

$$\mathcal{O}(\Delta x)^2 + \mathcal{O}(\Delta t)$$

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Order of Convergence:

$$\mathcal{O}(\Delta x)^2 + \mathcal{O}(\Delta t)$$

$$\mathcal{O}(\Delta t) \approx 1$$
, and $\mathcal{O}(\Delta x)^2 \approx 2$

The total order is dominated by Δt , so it should be closer to 1



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Conditionally Stable: Under the CFL condition $R = \frac{\Delta t \cdot \alpha}{\Delta x^2} < 1/2$



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Conditionally Stable: Under the CFL condition $R = \frac{\Delta t \cdot \alpha}{\Delta x^2} < 1/2$

Means: Δt has to be *much* smaller than Δx



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Numerical Set-Up: (4) Solution Methods

- Explicit Euler
- Implicit Euler



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$$u_i^{j+1} = (1-2R)u_i^j + R\left(u_{i+1}^j + u_{i-1}^j\right)$$

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$$u_i^{j+1} = (1-2R)u_i^j + R\left(u_{i+1}^j + u_{i-1}^j\right)$$

$$A \cdot \vec{x} = \vec{b}$$

Coefficient Matrix, A: A matrix representing the coefficients in the system of equations

Heat Equation, NSFD

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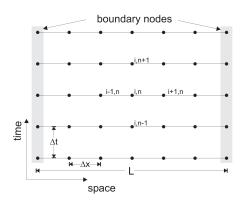
$$u_i^{j+1} = (1-2R)u_i^j + R\left(u_{i+1}^j + u_{i-1}^j\right)$$

$$A \cdot \vec{x} = \vec{b}$$

Coefficient Matrix, A: A matrix representing the coefficients in the system of equations

$$\vec{x} = A^{-1}\vec{b}$$

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$$u_{i}^{j+1} = (1-2R)u_{i}^{j} + R\left(u_{i+1}^{j} + u_{i-1}^{j}\right)$$

$$\begin{bmatrix} 1-2R & R & 0 & 0 & 0 & \dots & 0 \\ R & 1-2R & R & 0 & 0 & \dots & 0 \\ 0 & R & 1-2R & R & 0 & \dots & 0 \\ \vdots & & & & & \vdots \\ 0 & 0 & 0 & 0 & 0 & R & 1-2R \end{bmatrix} \cdot \begin{bmatrix} u_{i+1}^{j+1} \\ u_{j+1}^{j+1} \\ \vdots \\ u_{M}^{j+1} \end{bmatrix} = \begin{bmatrix} u_{2}^{j} \\ u_{3}^{j} \\ u_{4}^{j} \\ \vdots \\ u_{M}^{j} \end{bmatrix} + R \begin{bmatrix} u_{1}^{j} \\ 0 \\ \vdots \\ 0 \\ u_{M+1}^{j} \end{bmatrix}$$

Note: A is a Tridiagonal Matrix

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Since out boundary condition is always 0,

$$\begin{bmatrix} u_1^{j+1} \\ u_2^{j+1} \\ u_3^{j+1} \\ \vdots \\ u_M^{j+1} \end{bmatrix} = A^{-1} \cdot \begin{bmatrix} u_1^j \\ u_2^j \\ u_3^j \\ \vdots \\ u_M^j \end{bmatrix}$$

$$A = \begin{bmatrix} 1 - 2R & R & 0 & 0 & 0 & \dots & 0 \\ R & 1 - 2R & R & 0 & 0 & \dots & 0 \\ 0 & R & 1 - 2R & R & 0 & \dots & 0 \\ \vdots & & & & & \vdots \\ 0 & 0 & 0 & 0 & 0 & R & 1 - 2R \end{bmatrix}$$

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Implicit Code

```
8 Data structures for the linear tridiagonal system
57 - A = zeros(N x-1, N x-1);
58
59 - \bigcirc for i=1:N_x-1
A(i,i)=1-2*lambda; % diagonal
61 -
        end
62
63 - = for i=2:N x-1
64 -
   A(i,i-1)=lambda; %subdiagonal
65 - end
66
67 - \boxed{\text{for } i=1:N x-2}
68 - A(i,i+1)=lambda; %superdiagonal
69- end
```

Implicit Code

```
71- v=zeros(N_x-1,1); %BC

72- for j=1:N_t-1

73  % Solve triagonal system

74- v(1,1)=0; %BC at start

75- v(N_x-1,1)=0; %BC at end

76

77- u(:,j+1)=A\(u(:,j)+lambda*v); %x = A^-1 * b

end
```

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Implicit Code

```
71  v=zeros(N_x-1,1); %BC

72  for j=1:N_t-1

73  %Solve triagonal system

74  v(1,1)=0; %BC at start

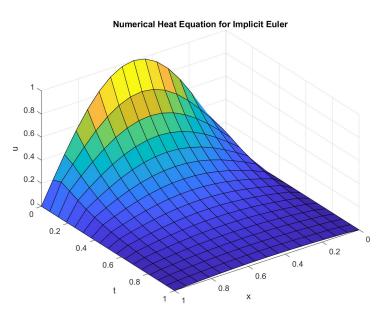
75  %BC at end

76  u(:,j+1)=A\(u(:,j)+lambda*v); %x = A^-1 * b

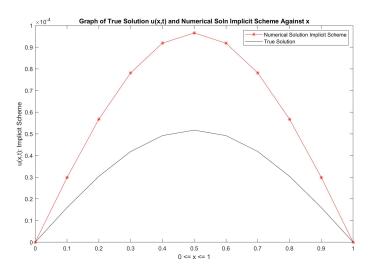
end
```

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Implicit Graph



True versus Implicit Euler



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Unconditionally Stable: always works

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Unconditionally Stable: always works

 L^1,L^2,L^∞ Error for Implicit Scheme for the Heat Equation: For fixed Δx and when Δt is cut in half.

Δχ	Δt	L ¹ Error	L ² Error	L^{∞} Error
0.01000	0.01000	2.42598905 <i>e</i> - 03	2.69481799 <i>e</i> — 04	3.81104814 <i>e</i> - 05
0.01000	0.00500	1.08459172 <i>e</i> - 03	1.20477760 <i>e</i> — 04	1.70381282 <i>e</i> — 05
0.01000	0.00250	5.13610852 <i>e</i> - 04	5.70525148 <i>e</i> — 05	8.06844401e - 06
0.01000	0.00125	2.50928217 <i>e</i> - 04	2.78734099 <i>e</i> - 05	3.94189542 <i>e</i> - 06
0.01000	0.000625	1.25033217e - 04	1.38888410 <i>e</i> — 05	1.96417873 <i>e</i> — 06

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Implicit Method Analysis: Order

$$\mathcal{O}(\Delta x)^2 + \mathcal{O}(\Delta t) \approx 1$$

Δχ	Δt	Δt Ratio	L [∞] Error	Ratio of Errors	Order
0.01000	0.01000	2	3.81104814e - 05	~	~
0.01000	0.00500	0.50000	1.70381282e - 05	1.70381282e - 05/3.81104814e - 05 = 0.44707	1.1614
0.01000	0.00250	0.50000	8.06844401e - 06	8.06844401e - 06/1.70381282e - 05 = 0.47355	1.0784
0.01000	0.00125	0.50000	3.94189542e - 06	3.94189542e - 06/8.06844401e - 06 = 0.48855	1.0334
0.01000	0.000625	0.50000	1.96417873e - 06	1.96417873e - 06/3.94189542e - 06 = 0.49828	1.0049

$$log_2(0.44707)/log_2(0.5) = 1.1614$$

$$log_2(0.47355)/log_2(0.5) = 1.0784$$

$$log_2(0.48855)/log_2(0.5) = 1.0334$$

$$log_2(0.49828)/log_2(0.5) = 1.0049$$



NSFD Approach



Change Finite Difference Scheme: Then, see how explicit and implicit change as a result

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Unity Approximations: Averaging

$$u_i^{j+1} = (1-2R)u_i^j \times 1_0 + R(u_{i+1}^j + u_{i-1}^j) \times 1_1$$



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Unity Approximations: Averaging

$$u_i^{j+1} = (1-2R)u_i^j \times 1_0 + R(u_{i+1}^j + u_{i-1}^j) \times 1_1$$

2-point Average:

$$\frac{u_A^j + u_B^j}{2u_C^j}$$
, $A, B, C \in \{i, i+1, i-1\}$



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Unity Approximations: Averaging

$$u_i^{j+1} = (1-2R)u_i^j \times 1_0 + R(u_{i+1}^j + u_{i-1}^j) \times 1_1$$

2-point Average:

$$\frac{u_A^{j} + u_B^{j}}{2u_C^{j}}, \quad A, B, C \in \{i, i+1, i-1\}$$

3-point Average:

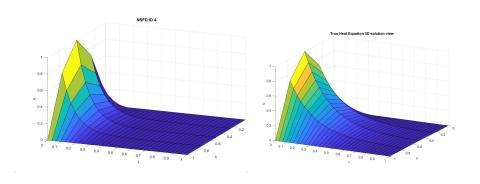
$$\frac{u_A^j + u_B^j + u_C^j}{3u_D^j}, \quad A, B, C, D \in \{i, i+1, i-1\}$$

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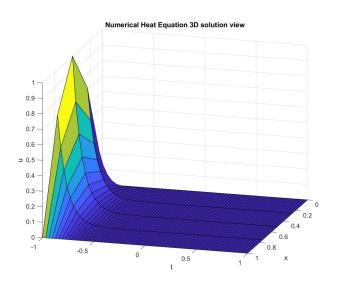
2-pt Avg Technique at 1_0

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NSFD ID 4: Explicit Euler

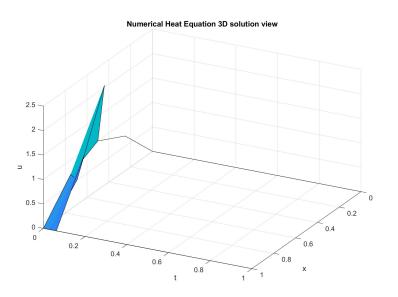


NSFD ID 16 (3pt Avg): Explicit Euler

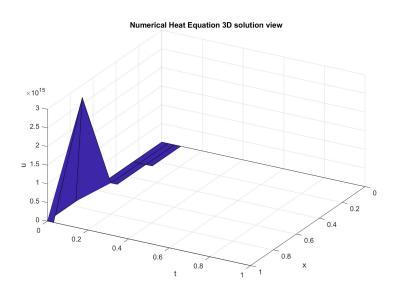


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NSFD ID 17 (3pt Avg): Explicit Euler



NSFD ID 25 (3pt Avg): Explicit Euler



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Implicit NSFD Schemes

Goal: Create the A matrix

2-point Average:

$$\frac{u_A^X + u_B^Y}{2u_C^Z}, \quad A, B, C \in \{i, i+1, i-1\} \quad X, Y, Z \in \{j, j+1, j-1\}$$

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Implicit NSFD Schemes

Goal: Create the A matrix

2-point Average:

$$\frac{u_A^X + u_B^Y}{2u_C^Z}, \quad A, B, C \in \{i, i+1, i-1\} \quad X, Y, Z \in \{j, j+1, j-1\}$$

When j = 1, $u_i^1 = \text{function in terms of } u_i^0$

When j = 2, $u_i^2 = \text{function in terms of } u_i^1$

Then can find the matrix A such that

$$A \cdot u^{j+1} = u^j$$

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Next Steps

- More Averages
- Build out Implicit NSFDs
- Find orders and compare against Implicit SFDs



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Thank you!

Questions?



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