

The Algebraic Theory
of
Topological Quantum Information

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Abstract

This book aims to give a comprehensive account of the algebraic theory of topological quantum information. It is intended to be accessible both to mathematicians unfamiliar with quantum mechanics and theoretical physicists unfamiliar with category theory. Additionally, this text should make a good reference for working researchers in the field. A primary focus of this text is balancing powerful algebraic generalities with concrete examples, principles, and applications.

For my mentors

0 Preface

This book is a mathematical treatment of topological quantum information, with a focus on formal algebraic aspects and a special eye towards topological quantum computation. This manuscript began as an extended set of notes from a course on topological quantum field theory given by Zhenghan Wang in the winter of 2022 at UC Santa Barbara. Through his courses, his private tutoring, and his recommendations, Zhenghan took me from a state of almost complete ignorance of mathematical physics to being a young researcher in the field. I am greatly emdebtetd to him for this, and it is certain that this book would not have existed without his guidance - he richly deserves of my apple.

This book would not have been possible without the tutelage of my esteemed mentors. Those who most directly contributed are the ones who used their time and energy to teach me topological quantum information - Dave Aasen, Mike Freedman, and Yuri Lensky. There are also those who took a chance on a young mathematician when they had every reason not to - Roald Dejean, Peter Bloomsburgh, Edward Frenkel, and Ken Ribet.

WORK: There's some other people to thank. Andrew Sylvester for letting me try out my arguments on him. Alexei Kitaev and Daniel Ranard now for advising me. Maybe Sam Packman for being a peer mentor?

Great pains have been taken to make this book as pedagogical and accessable as possible. The hope is that it should be readable by both mathematicians unfamiliar with quantum mechanics as well as theoretical physicists unfamiliar with category theory. A primary focus of this text is balancing powerful algebraic generalities with concrete examples, principles, and applications. The prerequisites for this book are a undergraduate-level understanding of topology, linear algebra, and group theory, as well as a popular-science level of familiarity with quantum mechanics.

There are already many great references to learn aspects of the material covered in this book. An excellently written and relatively complete book on topological quantum information from the perspective of a physicist is Steven Simon's text [Sim23]. Simon's book is algebraic, but does *not* include any category theory. The main references for the relevant category theory are Bakalov-Kirillov [BK⁺01] and Etingof-Gelaki-Nikshych-Ostrik [EGNO16]. While both excellent texts, they suffer notable shortcomings for learning topological quantum information. Bakalov-Kirillov was written in 2001, making it outdated. Etingof-Gelaki-Nikshych-Ostrik is modern, but makes no connections to physics and does not use the language of string diagrams. The manuscript most similar to this one is Kong-Zhang's preprint [KZ22]. We distinguish ourself from Kong-Zhang by our rigorous mathematical treatment, our different choice of topics, and our extended scope. Other relevant books and review articles include Wang's monograph [Wan10] and Kauffman-Lomonaco's quantum topology themed review [KL09].

WORK: I will add a section detailing the structure of this book, and how it should be read. I have not written enough for this to be useful yet.

Contents

| | |
|---------------------------------------------------------------------------------|----------|
| 0 Preface | 3 |
| 1 Modular categories | 7 |
| 1.1 Overview | 7 |
| 1.1.1 Introduction | 7 |
| 1.1.2 Using the final product | 8 |
| 1.2 First properties | 9 |
| 1.2.1 Definition | 9 |
| 1.2.2 Anyons in modular categories | 10 |
| 1.2.3 States in modular categories and unitarity | 14 |
| 1.2.4 Topological charge measurement | 17 |
| 1.3 The modular category toolkit | 21 |
| 1.3.1 Duality | 21 |
| 1.3.2 Trace | 22 |
| 1.3.3 Quantum dimension and Frobenius-Perron dimension | 24 |
| 1.3.4 Twist | 27 |
| 1.3.5 Functors, natural transformations, and equivalence | 31 |
| 1.3.6 Deligne tensor product | 31 |
| 1.4 The category of G -graded G -representations | 33 |
| 1.4.1 Overview | 33 |
| 1.4.2 Higher linear algebra | 34 |
| 1.4.3 Spherical fusion structures | 34 |
| 1.4.4 Braiding and modularity | 34 |
| 1.5 The modular representation | 35 |
| 1.5.1 Definition | 35 |
| 1.5.2 Torus perspective | 36 |
| 1.5.3 Bruguieres's modularity theorem and the Verlinde formula | 38 |
| 1.5.4 Verlinde formula | 42 |
| 1.5.5 Proof of modularity | 43 |
| 1.5.6 Vafa's theorem, unitarity of S -matrix, and the Chiral central charge . | 46 |
| 1.6 Skeletonization | 49 |
| 1.6.1 Principle | 49 |
| 1.6.2 F -symbols | 49 |
| 1.6.3 R -symbols | 49 |
| 1.6.4 θ -symbols | 49 |
| 1.6.5 Reconstruction theorem | 49 |
| 1.7 Quantum double modular categories | 49 |
| 1.7.1 The Drinfeld center | 49 |
| 1.7.2 Muger's theorem | 51 |
| 1.7.3 Discrete gauge theory as a quantum double and Morita equivalence . | 51 |
| 1.7.4 Factorizability and time reversal symmetry | 52 |
| 1.7.5 Levin-Wen model | 52 |
| 1.8 Unitarity | 52 |
| 1.8.1 Characterization of unitarizability | 52 |
| 1.8.2 Unitary braidings | 54 |
| 1.8.3 Uniqueness of unitary structure | 54 |

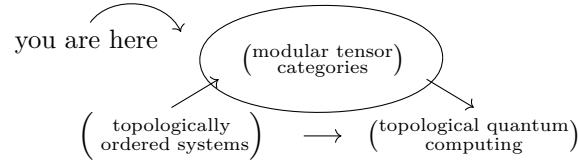
| | | |
|-------|-----------------------------------------------|----|
| 1.8.4 | Skeletonization of unitarity | 54 |
| 1.9 | Number theory in modular categories | 54 |
| 1.9.1 | Techniques and first results | 54 |
| 1.9.2 | Galois conjugation | 54 |
| 1.9.3 | Ocneanu rigidity | 54 |
| 1.9.4 | Rank-finiteness theorem | 54 |
| 1.9.5 | Schauenberg-Ng theorem | 54 |

1 Modular categories

1.1 Overview

1.1.1 Introduction

In this chapter we will be giving a detailed analysis of modular categories, the abstract algebraic structures used to describe anyons in topological order. We recall below how this fits into the general framework of this book:



Describing exactly what an anyon is and how it can transform in terms of states and unitary operators on a Hilbert space can be difficult. However, describing abstractly how these transformations compose with one another can be done relatively simply. Hence we take a composition-first category-theoretic approach to anyons. We will make heavy use of the diagrammatic language of braided monoidal categories established in chapter ???. Roughly, we can think of a modular category as being the category with the following data:

$$\begin{pmatrix} \text{objects: finite collections of anyons} \\ \text{morphisms: motions/behaviors of anyons} \end{pmatrix}$$

We have the following general picture for our algebraic theory:

Physics-math dictionary 1.1. Topological phases of matter are algebraically described by unitary modular categories \mathcal{C} called the *anyon theory* of the phase, and a choice of integer $c_- \in \mathbb{Z}$ called the *chiral central charge* of the phase. [Kit06]

Remark 1.2. Throughout these notes, there are some aspects of the general dictionary 1.1 that we have not emphasized. For instance, we have not emphasized that the modular categories describing topological phases are supposed to be *unitary* modular categories - we will discuss this later, but by analogy we can say that a unitary modular category is related to a modular category just like a Hilbert space is related to a vector space.

Another aspect we have not emphasized is the *chiral central charge*. This is a real invariant of topological phases which is beyond the description of modular categories. In particular, there are topological phases with no non-trivial anyons (and hence correspond to the trivial modular category) yet are non-trivial as topological phases. These are called *invertible* topological phases. Invertible topological phases are parameterized by an integer invariant, their chiral central charge. Not every pair (\mathcal{C}, c_-) describes a valid topological phase. In particular, \mathcal{C} determines c_- modulo 8 but there is no other condition. The way that \mathcal{C} determines c_- modulo 8 is known as the Gauss-Milgram formula and is discussed in subsection 1.5.6.

Remark 1.3. The major structures of a modular tensor category can be motivated by considering abstractly the possible motions and behaviors of anyons. The most basic thing anyons can do is move anyons around each other - this is known as *braid*. If the anyons

touch each other then they can congeal into a single anyon - this is known as *fusion*. Even if there are no anyons in a system, however, there is always something possible. Anyons can be spontaneously created, so long as every anyon which is created comes along with its corresponding antiparticle. This is known as *pair-creation*. These three operations are the fundamental structures which we will build into modular categories:

1. braiding;
2. fusion;
3. pair-creation.

Remark 1.4. Another potentially useful way of thinking about modular categories comes from analogy with classical physics. We saw in chapter ?? that topological classical systems have an algebraic description in terms of finite groups. Namely, quasiparticles in the system of ordered media with order space M is algebraically characterized by the fundamental group $\pi_1(M, m)$ of M relative to some basepoint $m \in M$. Seeing as topological order is a vast quantum generalization of classical ordered media, we can think of modular categories as being a vast quantum generalization of finite groups. Every finite group G induces a modular category $\mathfrak{D}(G)$, by first constructing the Kitaev quantum double model based on that finite group and then describing its anyons. Many modular categories do not come from the group-theoretical construction.

1.1.2 Using the final product

Before developing the theory of modular categories, it is good to get a feel for what using the final product is like. A modular category itself will be a big infinite structure, with infinitely many objects and infinitely many morphisms between those objects. However, all modular categories are in a real sense *finitely generated*. What we mean by this is that plugging in a finite number of objects and morphisms, the rest of the objects and morphisms can be recovered by the abstract rules encoded in the formalism. In this way, the axioms of modular categories are not only necessary by the fact that they restrict which categories can be modular categories, but they are also vital in the fact that they allow us to generate a full description of anyons from a minimal collection of data. For practically-minded readers, this can be viewed as the main motivation for putting so much work into defining modular categories..

Example 1.5. To highlight how the formalism used to define an object impacts its finitely-generated description, we take an example from group theory. Consider the 3-strand braid group B_3 . This group has infinitely many elements and the group operation $B_3 \times B_3 \rightarrow B_3$ naively takes an infinite amount of data to describe. However, the presentation

$$B_3 = \langle \sigma_1, \sigma_2 \mid \sigma_1\sigma_2\sigma_1 = \sigma_2\sigma_1\sigma_2 \rangle$$

gives completely finite description of B_3 . It is important to note, however, that this presentation would *not* have been enough to recover B_3 if we had just been told that B_3 is a monoid. The fact that B_3 is a group implied the existence of elements σ_1^{-1} , σ_2^{-1} , and defined how they interacted with σ_1, σ_2 . We see in this way that the axioms of a group not only serve as a restriction on what mathematical objects are allowed to be groups, but they also serve as a compression technique. They give the rules by which a minimal collection of data can be used to generate the rest.

An important step in going from modular categories to their description in terms of a finite set of data is in coming up with an efficient standard way of describing morphisms in a modular category. This is done using skeletonization, as discussed in section 1.6. A large table of these descriptions are found in appendix ???. We now give a worked example of how this data is used to compute observable quantities.

Example 1.6. WORK: add toric code modular category data

Example 1.7. Or, for a more complicated example, we can consider the data for $G = S_3$:
WORK: add $G = S_3$ modular category data

1.2 First properties

1.2.1 Definition

In this section we define modular categories, which are the main mathematical subject of this book. Seeing as lots of data is involved, we spread out the definition over a series of steps as to not overload the senses. These intermediate definitions are also important in their own right, because they will be used in other places in the algebraic theory of topological phases.

Definition 1.8 (Fusion category). A fusion category is the following data:

1. A category \mathcal{C} ;
2. The structure of a right-rigid monoidal category on \mathcal{C} ;
3. The structure of a \mathbb{C} -linear category on \mathcal{C} .

Such that:

1. The tensor product functor $\otimes : \mathcal{C} \times \mathcal{C} \rightarrow \mathcal{C}$ induces bilinear maps hom-spaces;
2. There is an equivalence $\mathcal{C} \cong \mathbf{Vec}_{\mathbb{C}}^n$ as \mathbb{C} -linear categories;
3. $\text{End}_{\mathcal{C}}(\mathbf{1}) \cong \mathbb{C}$ as \mathbb{C} -vector spaces

Remark 1.9. A fusion category is part of the way towards having all of the requisite structures of a modular category: it has a method for fusion inherited from the tensor product, and it has half of a method for pair-creation coming from right-rigidity. The \mathbb{C} -linearity allows us to think of hom-spaces as vector spaces, which allows us to treat hom-spaces as quantum systems. The condition (1) is a compatibility between the \mathbb{C} -linear structure and the monoidal structure. The conditions (2)-(3) are strong niceness and finiteness conditions - we will explain them in detail later.

Definition 1.10 (Spherical fusion category). A spherical fusion category is the following data:

1. A fusion category \mathcal{C} ;
2. A left-rigid structure on \mathcal{C} .

Such that:

1. The left-rigid and right-rigid structures on \mathcal{C} satisfy the axioms of a pivotal structure on \mathcal{C} ;
2. For every object $A \in \mathcal{C}$ and for every morphism $f : A \rightarrow A$, we have

$$\begin{array}{c} B \\ \text{---} \\ f \\ \text{---} \\ A \end{array} = \begin{array}{c} B \\ \text{---} \\ f \\ \text{---} \\ A \end{array}$$

Remark 1.11. A spherical fusion category now has a structure for fusion, and a full structure for pair-creation. The 2nd axiom in definition 1.10 is known as the *spherical axiom*. We will explain this axiom in more detail later.

Definition 1.12 (Pre-modular category). A pre-modular category is the following data:

1. A spherical fusion category \mathcal{C} ;
2. A braided structure on \mathcal{C} .

No extra compatibility conditions are required.

Remark 1.13. A pre-modular category has all of the structure we require of a modular category: fusion, pair-creation, and braiding. The final axiom it is missing is a non-degeneracy condition. The non-degeneracy condition is subtle in its interpretation, and we will explain it several different ways throughout this chapter.

Definition 1.14 (Modular category). A modular category is a pre-modular category satisfying the following condition. Let $A \in \mathcal{C}$ be an object. If

$$\begin{array}{c} A & B \\ \text{---} & \text{---} \\ \text{---} & \text{---} \\ A & B \end{array} = \begin{array}{c} A & B \\ | & | \\ A & B \end{array}$$

for all $B \in \mathcal{C}$, then $A \cong \mathbf{1}$.

Remark 1.15. Our usage of the term “modular category” is a bit nonstandard. Where we call “modular categories” most authors would call “modular tensor categories”. I find “modular category” more convenient here because it is shorter than “modular tensor category” and I want to be succinct. For many authors, the term “modular category” refers to a potentially non-semisimple modular category, such as in Turaev’s original work on modular categories [Tur92].

1.2.2 Anyons in modular categories

Modular categories are supposed to be theories of anyons in topological order. So, now that we have the definition of modular category, it is natural to ask: what do anyons mathematically correspond to, in modular categories? The answer lies within the condition in a fusion category \mathcal{C} that there is an equivalence $\mathcal{C} \cong \mathbf{Vec}_{\mathbb{C}}^n$ as \mathbb{C} -linear categories. We explore the importance of this condition.

Remark 1.16. Suppose we are given an object $V = (V_1, V_2 \dots V_n) \in \mathbf{Vec}_{\mathbb{C}}^n$. For all $1 \leq i \leq n$, let $\mathbb{C}_i \in \mathbf{Vec}_{\mathbb{C}}^n$ denote the object which has dimension zero in every index $j \neq i$ and is equal to \mathbb{C} in index i . We observe the isomorphism

$$\begin{aligned} V &\cong \bigoplus_{i=1}^n (0 \dots V_i \dots 0) \\ &\cong \bigoplus_{i=1}^n (0 \dots \mathbb{C}^{\dim(V_i)} \dots 0) \\ &\cong \bigoplus_{i=1}^n \dim(V_i) \cdot \mathbb{C}_i \end{aligned}$$

where $\dim(V_i) \cdot (\mathbb{C}_i) = \mathbb{C}_i \oplus \mathbb{C}_i \dots \oplus \mathbb{C}_i$, $\dim(V_i)$ many times. This computation shows that any object in $\mathbf{Vec}_{\mathbb{C}}^n$ can be decomposed into irreducible components \mathbb{C}_i . These objects \mathbb{C}_i are in a real sense the building blocks of $\mathbf{Vec}_{\mathbb{C}}^n$. In the language of definition 1.17, the objects \mathbb{C}_i are the simple objects of $\mathbf{Vec}_{\mathbb{C}}^n$.

Definition 1.17. A *simple object* A in a fusion category \mathcal{C} is an object which has no direct sum decomposition into smaller objects. That is, $A \not\cong B \oplus C$ for any non-zero objects $B, C \in \mathcal{C}$ where \oplus denotes the biproduct in \mathcal{C} .

Physics-math dictionary 1.18. Isomorphism classes of simple objects in \mathcal{C} correspond to anyon types in a topological phase described by the MTC \mathcal{C} .

Definition 1.19. For all fusion categories \mathcal{C} , we define $\mathcal{L}(\mathcal{C})$ (or simple \mathcal{L} when \mathcal{C} is clear from context) to the set of isomorphism classes of simple objects. By entry 1.18 of the physics-math direction, elements of $\mathcal{L}(\mathcal{C})$ correspond bijectively to anyon types in any topological phase described by \mathcal{C} .

Proposition 1.20. Let \mathcal{C} be a fusion category. The biproduct of any two elements in \mathcal{C} exists, and \mathcal{C} has a zero object. The set $\mathcal{L} = \mathcal{L}(\mathcal{C})$ is finite. Choose an object $X \in \mathcal{C}$. There exist unique nonnegative integers $c_{[A]}$, $[A] \in \mathcal{L}$ such that

$$X \cong \bigoplus_{[A] \in \mathcal{L}} N_{[A]} \cdot A.$$

Proof. We first show that \mathcal{C} has biproducts. Let $F : \mathcal{C} \rightarrow \mathbf{Vec}_{\mathbb{C}}^n$, $G : \mathbf{Vec}_{\mathbb{C}}^n \rightarrow \mathcal{C}$ be a pair of functors which induces an equivalence of categories, for some $n \geq 1$. Let $A, B \in \mathcal{C}$ be objects. Since G and F are fully faithful, the universal property of the direct sum $F(A) \oplus F(B)$ guarantees that $G(F(A) \oplus F(B))$ will be a direct sum of $G(F(A))$ and $G(F(B))$. Since $G(F(A)) \cong A$ and $G(F(B)) \cong B$, we conclude that \mathcal{C} has biproducts. The object $G(0)$ is a zero object for \mathcal{C} .

We now prove that \mathcal{C} has finitely many isomorphism classes of simple objects. It is clear that an object $A \in \mathcal{C}$ is simple if and only if $F(A) \in \mathbf{Vec}_{\mathbb{C}}^n$ is simple. Thus, since G serves as an inverse, F establishes a bijection between isomorphism classes of simple objects in \mathcal{C} and isomorphism classes of simple objects in $\mathbf{Vec}_{\mathbb{C}}^n$. Every simple object in $\mathbf{Vec}_{\mathbb{C}}^n$ will be isomorphic to \mathbb{C}_i for some $1 \leq i \leq n$. Hence, there are n simple objects in $\mathbf{Vec}_{\mathbb{C}}^n$. Hence, there are n simple objects in \mathcal{C} , which is finite. The unique direct sum decomposition is clearly true in $\mathbf{Vec}_{\mathbb{C}}^n$. It is immediate that it passes to a unique direct sum decomposition in \mathcal{C} . \square

Proposition 1.21 (Schur's Lemma). *Let \mathcal{C} be a fusion category. An object $A \in \mathcal{C}$ is simple if and only if its endomorphism ring $\text{End}(A)$ is one-dimensional. Additionally, if $A, B \in \mathcal{C}$ are nonisomorphic simple objects then $\text{Hom}(A, B) = 0$.*

Proof. Let $F : \mathcal{C} \rightarrow \mathbf{Vec}_{\mathbb{C}}^n$, $G : \mathbf{Vec}_{\mathbb{C}}^n \rightarrow \mathcal{C}$ be a pair of \mathbb{C} -linear functors which establishes an equivalence between \mathcal{C} and $\mathbf{Vec}_{\mathbb{C}}^n$ as \mathbb{C} -linear categories. The simple objects in $\mathbf{Vec}_{\mathbb{C}}^n$ are all isomorphic to \mathbb{C}_i for some $1 \leq i \leq n$. We compute that

$$\dim \left(\text{Hom}_{\mathbf{Vec}_{\mathbb{C}}^n} \left(\bigoplus_{i=1}^n n_i \mathbb{C}_i, \bigoplus_{i=1}^n m_i \mathbb{C}_i \right) \right) = \sum_{i=1}^n n_i m_i.$$

As a corollary of this formula, we find that if $A = \bigoplus_{i=1}^n n_i \mathbb{C}_i$ then $\dim(\text{End}_{\mathbf{Vec}_{\mathbb{C}}^n}(A)) = \sum_{i=1}^n n_i^2$. Clearly, this dimension is equal to one if $A = \mathbb{C}_i$ for some $1 \leq i \leq n$, and is greater than one otherwise. As a second corollary, we compute that $\text{Hom}(\mathbb{C}_i, \mathbb{C}_j) = 0$ whenever $i \neq j$. The functor G induces a bijection between isomorphism classes of simple objects in $\mathbf{Vec}_{\mathbb{C}}^n$ and isomorphism classes of simple objects in $\mathbf{Vec}_{\mathbb{C}}^n$, and it induces vector space isomorphisms on hom spaces. This means that the results for $\mathbf{Vec}_{\mathbb{C}}^n$ translate to the desired result on \mathcal{C} . \square

Remark 1.22. Schur's lemma is a first verification that simple objects are a good choice of mathematical characterization of anyons. If A, B are distinct anyon types, then there should not be any physical process which goes from one to another. There is no physical mechanism for locally turning one anyon type into another. This is captured by the formula $\text{Hom}(A, B) = 0$. Similarly, given an anyon A , there is no nontrivial action that can be locally performed on A . This comes from the fact that information is topologically protected, and thus cannot be changed by acting on a single particle - topological information processing requires global braiding between multiple particles. This is encoded in the fact that $\text{Hom}(A, A) \cong \mathbb{C}$ is one dimensional and hence consists only of trivial phase gates.

Remark 1.23. As an application of Schur's lemma, we can observe that the monoidal unit $\mathbf{1}$ is a simple object in every fusion category. By our physics-math dictionary, this means that $\mathbf{1}$ corresponds to an anyon type. This anyon $\mathbf{1}$ is the trivial “no-anyon” type, corresponding to a quasiparticle which happens to be the same as the homogenous bulk.

Physics-math dictionary 1.24. The monoidal unit $\mathbf{1}$ corresponds to the trivial anyon type, which is the same as the homogenous bulk (sometimes called the *vacuum* anyon type).

Remark 1.25. Continuing in our expansion of the physics-math dictionary, we can assert that dual objects (defined using the right-rigid structure on a fusion category) correspond to antiparticles (entry 1.26). This is consistent with the fact that anyon types correspond to simple objects, by proposition 1.27.

Physics-math dictionary 1.26. For every simple object A in a modular category \mathcal{C} , the dual object A^* corresponds to the *antiparticle* of A . Expanding our physics-math dictionary, we say that for every anyon A its *antiparticle* is the dual A^* which comes from right-rigidity. This gives a valid anyon type by proposition 1.27.

Proposition 1.27. *Let \mathcal{C} be a fusion category. If $A \in \mathcal{C}$ is a simple object, then so is A^* .*

Proof. By proposition ?? duality induces a bijection on hom-spaces. Since composition is bilinear, duality is thus an isomorphism of vector spaces on all hom-spaces. Hence, for all $A \in \mathcal{C}$ there is an isomorphism $\text{Hom}(A, A) \cong \text{Hom}(A^*, A^*)$. So, $\dim \text{Hom}(A, A) = 1$ if and only if $\dim \text{Hom}(A^*, A^*) = 1$, and thus the result follows from Schur's lemma. \square

Physics-math dictionary 1.28. The tensor product \otimes physically corresponds to joining anyons, forming a composite anyon configuration. That is, the object $A \otimes B$ corresponds to the configuration with one A -type anyon and one B -type anyon.

Remark 1.29. By entry 1.28 of the physics-math dictionary and proposition 1.20, it behooves us to take a moment to interpret the direct sum decomposition

$$A \otimes B \cong \bigoplus_{[C] \in \mathcal{L}} N_C^{A,B} \cdot C, \quad (1.1)$$

where $A, B \in \mathcal{C}$ are objects in a fusion category. We call the non-negative integers $N_C^{A,B}$ the *fusion coefficients* of \mathcal{C} . We observe that

$$\dim \text{Hom}(A \otimes B, C) = N_C^{A,B}. \quad (1.2)$$

Morphisms in \mathcal{C} correspond to physical processes. So, equation 1.2 can be interpreted as saying that whenever $N_C^{A,B} > 0$, there is a physical process $A \otimes B \rightarrow C$. That is, there is a physical process which goes from the composite system of one A -type anyon and one B -type anyon and outputs one C -type anyon. Neccearily, such a process can be interpreted as *fusion* of A and B to C . Moreover, equation 1.2 tells us that the space of possible physical processes which fuse A and B to C (known as *fusion channels* $A \otimes B \rightarrow C$) is $N_C^{A,B}$ -dimensional.

Remark 1.30. We can interpret the modularity axiom in terms of anyons. To begin, we observe the algebraic fact that a pre-modular category is modular if and only if for all *simple* objects $A \in \mathcal{C}$,

$$\begin{array}{ccc} A & B \\ \text{---} & \text{---} \\ \text{---} & \text{---} \end{array} = \begin{array}{ccc} A & B \\ | & | \\ A & B \end{array}$$

for all *simple* objects $B \in \mathcal{C}$ implies $A \cong \mathbf{1}$. This is physically interpreted as saying that if an anyon braids trivially with all other anyons, then it must be the trivial anyon type $\mathbf{1}$. In otherwords, there are no *transparent* anyons. This fact is physically motivated in subsection [ref].

Proposition 1.31. Let \mathcal{C} be a fusion category. For all simple objects $A, B, C, E \in \mathcal{C}$,

$$\sum_E N_E^{A,B} N_D^{E,C} = \sum_G N_G^{B,C} N_D^{A,G}. \quad (1.3)$$

Proof. We observe on one hand that

$$\begin{aligned} (A \otimes B) \otimes C &\cong \bigoplus_E N_E^{A,B} \cdot E \otimes C \\ &\cong \bigoplus_{E,D} N_E^{A,B} N_D^{E,C} \cdot D \end{aligned}$$

and on the other hand

$$\begin{aligned} A \otimes (B \otimes C) &\cong \bigoplus_G N_G^{B,C} \cdot A \otimes G \\ &\cong \bigoplus_{G,D} N_G^{B,C} N_D^{A,G} \cdot D. \end{aligned}$$

Since $(A \otimes B) \otimes C \cong A \otimes (B \otimes C)$, we can compare coefficients to get the result. \square

Proposition 1.32. *Let \mathcal{C} be a braided fusion category. For all simple objects A, B, C ,*

$$N_C^{A,B} = N_C^{B,A}. \quad (1.4)$$

Proof. This follows immediately from the fact that $A \otimes B \cong B \otimes A$ in a braided fusion category. \square

1.2.3 States in modular categories and unitarity

It is now worth reflecting on how to interpret states in a topologically ordered system in terms modular categories. Certainly, objects in modular categories do *not* correspond to quantum systems. They don't have vector space structure. Objects correspond roughly to anyon configurations, which are classical observables. The hom-spaces are what have vector space structure, by \mathbb{C} -linearity. However, they are not yet Hilbert spaces. It is exactly for this reason we need to define *unitary modular categories*. These are modular categories in which hom-spaces have Hilbert space structure.

Definition 1.33 (Unitary fusion category). A unitary fusion category is the following data:

1. An fusion category \mathcal{C} .
2. (Conjugation) A linear map $\dagger : \text{Hom}(A, B) \rightarrow \text{Hom}(B, A)$ for all $A, B \in \mathcal{C}$.

Such that:

1. (Unitarity) Given $f : A \rightarrow A$ an endomorphism of $A \in \mathcal{C}$, define

$$\text{tr}(f) = \text{ev}_A \circ (\text{id}_{A^*} \otimes f) \circ (\text{ev}_A)^\dagger.$$

The map $\langle \cdot | \cdot \rangle : \text{Hom}(A, B) \times \text{Hom}(A, B) \rightarrow \mathbb{C}$ defined by $\langle f | g \rangle = \text{tr}(f^\dagger \circ g)$ is an inner product, endowing $\text{Hom}(A, B)$ with the structure of a Hilbert space.

2. $(f^\dagger)^\dagger = f$ for all $f \in \text{Hom}(A, B)$, $A, B \in \mathcal{C}$.
3. $(f \circ g)^\dagger = g^\dagger \circ f^\dagger$ for all $f \in \text{Hom}(B, C), g \in \text{Hom}(A, B)$, $A, B, C \in \mathcal{C}$.
4. $(f \otimes g)^\dagger = f^\dagger \otimes g^\dagger$ for all $f \in \text{Hom}(A, B), g \in \text{Hom}(C, D)$, $A, B, C, D \in \mathcal{C}$.
5. $(\text{coev}_A)^\dagger \circ (f \otimes \text{id}_{A^*}) \circ \text{coev}_A = \text{tr}(f)$ for all $A \in \mathcal{C}$

Remark 1.34. Unitary fusion categories make for a pleasant object of study because by proposition 1.106 the distinguished maps $(\text{ev}_A)^\dagger : 1 \rightarrow A^* \otimes A$ and $(\text{coev}_A)^\dagger : A \otimes A^* \rightarrow 1$ induce a pivotal structure.

Remark 1.35. In subsection [ref] proposition 1.106, we will show that in every unitary fusion category the maps $\text{ev}_A^L = (\text{coev}_A)^\dagger$ and $\text{coev}_A^L = (\text{ev}_A)^\dagger$ give a left-rigid structure on \mathcal{C} . This left-rigidity endows \mathcal{C} with the structure of a spherical fusion category. Thus, every unitary fusion category is automatically a spherical fusion category.

Definition 1.36 (Unitary pre-modular category). A unitary modular category is the following data:

1. A modular category \mathcal{C} ;
2. (Conjugation) A linear map $\dagger : \text{Hom}(A, B) \rightarrow \text{Hom}(B, A)$ for all $A, B \in \mathcal{C}$.

Such that:

1. Forgetting the left-rigid structure and braiding, (\mathcal{C}, \dagger) forms a unitary fusion category.
2. $(\text{ev}_A^R)^\dagger = \text{coev}_A^L$;
3. $(\text{coev}_A^R)^\dagger = \text{ev}_A^L$;
4. $(\beta_{A,B})^\dagger = \beta_{B,A}^{-1}$.

Definition 1.37 (Unitary modular category). A unitary modular category is a unitary pre-modular category which satisfies the non-degeneracy axiom of a modular category.

Remark 1.38. The compatibility conditions for the twist are chosen so that the definition of trace as a modular category and the definition of trace as a unitary fusion category coincide.

Physics-math dictionary 1.39. For any phase whose anyons are described by a unitary modular category \mathcal{C} .

$$\left(\begin{array}{c} \text{states of topological order } \mathcal{C} \\ \text{on the sphere } S^2 \\ \text{with anyon configuration } A_1, A_2 \dots A_n \end{array} \right) = \left(\begin{array}{c} \text{normalized vectors in the Hilbert space} \\ \text{Hom}_{\mathcal{C}}(\mathbf{1}, A_1 \otimes A_2 \dots \otimes A_n) \end{array} \right) \quad (1.5)$$

where by “anyon configuration $A_1, A_2 \dots A_n$ ” we mean that the state has anyons present in n sites, arranged left to right on a line segment in the sphere, with corresponding anyon type $A_1, A_2 \dots A_n$.

Remark 1.40. It is not immediately clear why the physical space in equation 1.5 should be the sphere. As one source of motivation, we can observe the following corollary of equation 1.5:

$$\dim \left(\begin{array}{c} \text{Hilbert space of topological order } \mathcal{C} \\ \text{on the sphere } S^2 \\ \text{with exactly one anyon of type } A \end{array} \right) = \dim \text{Hom}(\mathbf{1}, A) = \begin{cases} 1 & A = \mathbf{1} \\ 0 & \text{otherwise.} \end{cases}$$

Hence, we conclude that if the sphere has exactly one anyon on it then that anyon type must be trivial. Moreover, there is unique state on the sphere with no anyons. This is true on physical grounds. Suppose there was a state on the sphere with a single anyon of type A . Now choose any other anyon B . We observe that braiding B around A acts by the phase $+1$, because the loop of B can be contracted around the other side of the sphere, as illustrated in figure 1.1. Diagrammatically, this means that for all simple objects $B \in \mathcal{L}$

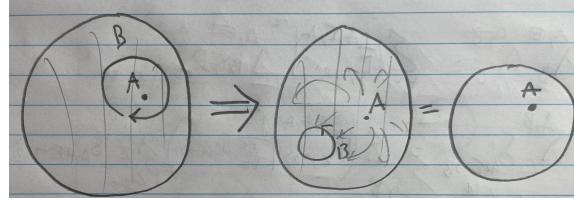


Figure 1.1: Illustration of braiding around a single anyon on the sphere

$$\begin{array}{c} A \quad B \\ \swarrow \quad \searrow \\ A \quad B \end{array} = \begin{array}{c} A \quad B \\ | \quad | \\ A \quad B \end{array}$$

By modularity (c.f. remark 1.30), this implies $A \cong \mathbf{1}$. Thus, the if there is only one anyon on the sphere it must have trivial type. The fact that there is a unique state on the sphere with no anyons is also true on physical grounds, but is more subtle. It is discussed in subsection [ref].

Physics-math dictionary 1.41. For any modular category \mathcal{C} :

$$\left(\begin{array}{l} \text{states of topological order } \mathcal{C} \\ \text{on the infinite flat plane } \mathbb{R}^2 \\ \text{with anyon configuration } A_1, A_2, \dots, A_n \end{array} \right) = \left(\begin{array}{l} \text{normalized vectors in the Hilbert space} \\ \text{Hom}_{\mathcal{C}} \left(\bigoplus_{[B] \in \mathcal{L}} B, A_1 \otimes A_2 \otimes \dots \otimes A_n \right) \end{array} \right) \quad (1.6)$$

Remark 1.42. Replacing $\mathbf{1}$ with $\bigoplus_{[B] \in \mathcal{L}} B$ reflects the differences between the sphere and the infinite flat plane. On a sphere, the only way to have a state with a single isolated anyon is for that anyon to have trivial charge. However, on the plane you can have a state with any single isolated anyon. This is done by creating a particle/antiparticle pair, and then pulling one half of that pair away to infinity, as illustrated in figure 1.2. This is reflected by the fact that equation 1.7 says that for all anyon types A ,

$$\dim \left(\begin{array}{l} \text{Hilbert space of states of topological order } \mathcal{C} \\ \text{on the infinite flat plane } \mathbb{R}^2 \\ \text{with a single anyon of type } A \end{array} \right) = \dim \text{Hom} \left(\bigoplus_{[B] \in \mathcal{L}} B, A \right) = 1. \quad (1.7)$$

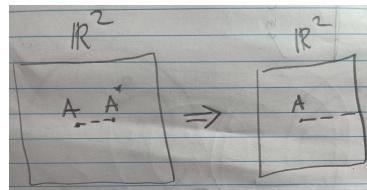


Figure 1.2: Creating a state on the plane with a single isolated anyon

The fact that there is a unique state with a single anyon of type A on the plane \mathbb{R}^2 is a more subtle physical argument, discussed in subsection [ref].

Remark 1.43. Working on higher-genus surfaces or on surfaces with holes is possible, but it is a bit cumbersome in the approach we are taking. We refer to appendix [ref] for a discussion of the topic.

Remark 1.44. The anyon configurations in equations 1.5 and 1.7 are always assumed to be linear. The main reason to do this is because it makes the mathematics much simpler. Keeping track of the positions of each of the anyons in two dimensional space adds unnecessary complexity. Seeing as every anyon configuration can be pushed onto a one-dimensional space, only working with a one-dimensional configuration does not affect the generality of the answers and hence it is preferred.

Remark 1.45. The formula $\text{Hom}_{\mathcal{C}}(\mathbf{1}, A_1 \otimes A_2 \dots \otimes A_n)$ encodes the fact that states can be specified by their history. There are multiple states with anyon configuration $A_1 \otimes A_2 \dots \otimes A_n$. One way of specifying a state $|\psi\rangle$ is to start with the unique state on the sphere with no anyons, and then specify how $|\psi\rangle$ is constructed from the ground state via some unambiguous process. This unambiguous process from the state with no anyons to the state with anyons $A_1 \dots A_n$ is a morphism in $\text{Hom}_{\mathcal{C}}(\mathbf{1}, A_1 \otimes A_2 \dots \otimes A_n)$.

Example 1.46. We can describe two states with four the four anyons A, A^*, A, A^* as follows. The first state is obtained by starting with the unique state with no anyons, and creating two independent (A, A^*) particle/antiparticle pairs adjacent to each other:

$$| \begin{array}{cccc} A & A^* & A & A^* \\ \backslash & / & \backslash & / \end{array} \rangle$$

Another state can be constructed by creating a (A, A^*) particle/antiparticle pair, then moving the halves of the pair far apart from each other, and then creating a (A^*, A) particle/antiparticle pair in the space between the halves of the first pair:

$$| \underbrace{\hspace{1cm}}_{A} \hspace{0.2cm} A^* \hspace{0.2cm} A \hspace{0.2cm} A^* \rangle$$

These two (unnormalized) states have the same anyon configuration. However, they are distinguished by the fact that they were created by different processes. Translating pair-creation and movement into category theoretic language, we see that these two states correspond to different maps in $\text{Hom}_\mathcal{C}(\mathbf{1}, A \otimes A^* \otimes A \otimes A^*)$.

1.2.4 Topological charge measurement

In any reasonable physical theory, one needs to talk about measurements and observables. For the algebraic theory of topological quantum information, the only relevant type of measurement is a *topological charge measurement* [Bon21]. This is the measurement of the overall anyon type (or *topological charge*) of a collection of anyons. When the topological charge of a pair of anyons in a state $|\psi\rangle$ is measured, $|\psi\rangle$ will collapse onto a new state $|\psi'\rangle$ for which that pair of anyons now has a single well-defined charge as shown in equation 1.8. That is, the two anyons will fuse together.

$$\text{measure charge} \quad \begin{array}{c} (A_1 \ A_2) \\ \downarrow \\ B \end{array} \dots A_{n-1} \ A_n \quad (1.8)$$

As a first step towards an algebraic understanding of topological charge measurement, suppose that we are presented with a normalized state on the sphere of the following form:

$$| \psi \rangle = | A_1 | A_2 | \dots | A_n \rangle \quad (1.9)$$

Where $A_1 \dots A_n$, B are simple objects, and $\mu : B \rightarrow A_1 \otimes A_2$ is normalized so that $\langle \mu | \mu \rangle = 1$. If we postcompose $|\psi\rangle$ with the map $\mu^\dagger \otimes \text{id}_{A_3 \otimes \dots \otimes A_n}$, the result will be equal to

$$\mu^\dagger | \psi \rangle = | A_1 | A_2 | \dots | A_n \rangle \quad (1.10)$$

This means that μ gives a physical process whereby $|\psi\rangle$ can transform to a state with an anyon of type B where A_1, A_2 used to be. Moreover, if we postcompose $|\psi\rangle$ with $\eta \otimes \text{id}_{A_3 \otimes \dots \otimes A_n}$ with any map such that $\eta \circ \mu = 0$ so $\eta |\psi\rangle$ will be zero (and thus an invalid quantum state). In particular, since there are no maps $B \rightarrow C$ for non-isomorphism simple objects B, C , we find that there is no physical process on the first two anyons which can be applied to $|\psi\rangle$ which results in a state with charge $C \not\cong B$ where A_1, A_2 used to be. This motivates the following mathematical-physical correspondance:

Physics-math dictionary 1.47. Given a state $|\psi\rangle$ with anyon content $A_1 \dots A_n, B_1 \dots B_m$ we say that the anyons $A_1 \dots A_n$ have overall topological charge C with fusion channel $\mu^\dagger : A_1 \otimes A_2 \dots \otimes A_n \rightarrow C$ if there exists α such that $|\psi\rangle$ can be expressed as follows:

$$| \psi \rangle = | A_1 | A_2 | \dots | A_n | B_1 | B_m \rangle \quad (1.11)$$

If $|\psi\rangle$ is a superposition of states the form 1.11 for a fixed C over different μ , then we say that the anyons $A_1 \dots A_n$ have overall topological charge C but have indefinite fusion channel.

Proposition 1.48. Let $A, B \in \mathcal{C}$ be objects in a modular category. We have that

$$\begin{array}{c} A \ B \\ | \qquad | \\ A \ B \end{array} = \sum_{C, \mu} \begin{array}{c} A \ B \\ | \mu^\dagger | \ C \\ | \mu | \\ A \ B \end{array} \quad (1.12)$$

where $[C]$ runs over isomorphism classes of simple objects in \mathcal{C} , and μ runs over an orthonormal basis of $\text{Hom}(C, A \otimes B)$ in the isotopy normalization.

Proof. WORK: Both maps are the same when precomposed with η' and postcomposed with η^\dagger for any η', η^\dagger in the orthonormal basis. Since it is a basis, we can take linear combinations to conclude that they are the same precomposed and postcomposed with any map, and thus they are the same map. \square

Corollary 1.49. Let \mathcal{C} be a modular category. For all simple objects $A_1, A_2 \dots A_n \in \mathcal{C}$, we have that

WORK: Generalized version of proposition 1.48.

Proof. WORK: This is a straightforward proof by induction. \square

Remark 1.50. Corollary 1.49 tells us every state $|\psi\rangle$ with anyon content $A_1 \dots A_n, B_1 \dots B_m$ can be made into a superposition of states with definite overall charge and fusion channel of $A_1 \dots A_n$. The general principles of quantum mechanics tell us how to measure

Physics-math dictionary 1.51. Let $|\psi\rangle = |\alpha\rangle$ be a normalized state, where $\alpha : \mathbf{1} \rightarrow A_1 \dots \otimes A_n \otimes B_1 \dots \otimes B_m$. For all normalized $\mu^\dagger : A_1 \otimes A_2 \dots \otimes A_n \rightarrow C$, define the state

$$|\psi_{C,\mu}\rangle = |A_1| \boxed{A_2 \dots A_n} |B_1 B_m\rangle$$

\vdots

 $\mathbf{1}$

and

$$|\psi_C\rangle = \sum_{\mu: A_1 \dots \otimes A_n \rightarrow C} |\psi_{C,\mu}\rangle$$

By proposition 1.49, we have that

$$|\psi\rangle = \sum_C |\psi_C\rangle = \sum_{C,\mu} |\psi_{C,\mu}\rangle.$$

A *topological charge measurement* on the anyons $A_1 \dots A_n$ with respect to the orthonormal basis $\{\mu\}$ of $\bigoplus_C \text{Hom}(C, A_1 \otimes \dots \otimes A_n)$ collapses $|\psi\rangle$ onto the state $\frac{1}{\|\psi_{C,\mu}\|} |\psi_{C,\mu}\rangle$ with probability $\|\psi_{C,\mu}\|^2$, and the corresponding observable is (C, μ) . Alternatively, we can think of the state resulting from the measurement as being

$$\frac{\mu^\dagger \circ |\psi_{C,\mu}\rangle}{\|\psi_{C,\mu}\|} = \frac{1}{\|\psi_{C,\mu}\|} |A_1| \boxed{A_2 \dots A_n} |B_1 B_m\rangle$$

\vdots

 $\mathbf{1}$

A topological charge measurement of the anyons $A_1 \dots A_n$ without respect to any basis of $\bigoplus_C \text{Hom}(C, A_1 \otimes \dots \otimes A_n)$ collapses $|\psi\rangle$ onto the state $\frac{1}{\|\psi_C\|} |\psi_C\rangle$ with probability $\|\psi_C\|^2$.

Definition 1.52. Let \mathcal{C} be a modular category. We define the *quantum dimension* of a simple object $A \in \mathcal{C}$ to be $d_A = \text{tr}(\text{id}_A)$.

Remark 1.53. We will discuss quantum dimensions in more detail in subsection [ref].
WORK: is there anything I need to recall now? Probably that they are non-zero, and positive when \mathcal{C} is unitary.

Example 1.54. We give a basic and fundamental example of topological charge measurement. Suppose that we create adjacent particle/antiparticle pairs, $(A^* \otimes A)$ and $(B \otimes B^*)$. What is the probability that we measure the charge C when the middle anyons A, B are fused? When we create the pairs of anyons (A^*, A) , (B, B^*) , the state is represented algebraically as

$$|\psi\rangle = \frac{1}{\sqrt{d_A d_B}} |A^* \ A \ B \ B^*\rangle.$$

The normalization is correct, because

$$\begin{aligned} \langle \psi | \psi \rangle &= \left(\frac{1}{\sqrt{d_A d_B}} \langle A^* \ A \ B \ B^* | \right) \left(\frac{1}{\sqrt{d_A d_B}} |A^* \ A \ B \ B^*\rangle \right) \\ &= \frac{1}{d_A d_B} A^* \bigcirc A \ B \bigcirc B^* = 1. \end{aligned}$$

Now, choosing some an orthonormal basis $\{\mu\}$ of $C \rightarrow A \otimes B$ for all C , we can define

$$|\psi_{C,\mu}\rangle = \frac{1}{\sqrt{d_A d_B}} |A^* \ A \ B \ B^* \begin{array}{c} \mu \\ C \\ \mu^\dagger \end{array} \rangle.$$

We compute

$$\begin{aligned} \langle \psi_{C,\mu} | \psi_{C,\mu} \rangle &= \frac{1}{d_A d_B} A^* \left(\begin{array}{c} \mu \\ C \\ \mu^\dagger \end{array} \right) B^* \\ &= \frac{1}{d_A d_B} A^\dagger \left(\begin{array}{c} \mu^\dagger \\ B \\ \mu \end{array} \right) C^* \\ &= \frac{1}{d_A d_B} C \bigcirc C^* = \frac{d_C}{d_A d_B}. \end{aligned}$$

Thus, when the topological charge measurement is performed with respect to the basis $\{\mu\}$, the probability of the result being (C, μ) is $d_C/(d_A d_B)$. Since for any given C there are $N_C^{A,B}$ different choices of μ , the probability of measuring C when the topological charge measurement is performed without a choice of basis is

$$(\text{probability of measuring } C \text{ when fusing } A, B) = \frac{d_C}{d_A d_B} N_C^{A,B}. \quad (1.13)$$

Remark 1.55. As a sanity check, we can remark that (proposition [ref])

$$\sum_C \frac{d_C}{d_A d_B} N_C^{A,B} = 1.$$

Thus, the probabilities from equation 1.13 really do form a probability distribution.

Remark 1.56. When we measure fusion channel, we need a measurement apparatus with a distinguished choice of fusion basis.

1.3 The modular category toolkit

In this section, we will introduce and prove the basic facts about the most important structures in the theory of modular categories. These facts and structures are the tools used for solving problems about the algebraic theory of anyons.

1.3.1 Duality

Duality is baked into our definition of modular categories as a fundamental part of the structure. In this subsection we will explore the basic properties of duality in modular categories.

Proposition 1.57. *Let \mathcal{C} be a fusion category and let $A, B, C \in \mathcal{C}$ be simple objects. We have the following:*

- (i) (Anti-involution) $N_C^{A,B} = N_{C^*}^{B^*,A^*}$;
- (ii) (Frobenius reciprocity) $N_C^{A,B} = N_B^{A^*,C} = N_A^{C,B^*}$.

Proof. Part (i) follows from the fact that the duality functor is fully faithful and monoidal from proposition ??, so

$$N_C^{A,B} = \dim \text{Hom}(C, A \otimes B) = \dim \text{Hom}(A^* \otimes B^*, C^*) = N_{C^*}^{B^*,A^*}.$$

Part (ii) follows from the following computation. Consider the map

$$i : \text{Hom}(A, B \otimes C) \longrightarrow \text{Hom}(A^* \otimes C, B)$$

$$\begin{array}{ccc} A & B \\ \square & \square \\ \alpha & \\ \downarrow & \\ C & \end{array} \mapsto \begin{array}{ccc} & B \\ & \curvearrowleft \\ & \square \\ & \alpha \\ & \downarrow \\ A^* & C \end{array}$$

Since composition is bilinear, i is a linear map. The map

$$\text{Hom}(A^* \otimes C, B) \longrightarrow \text{Hom}(A, B \otimes C)$$

$$\begin{array}{ccc} B \\ \square \\ \alpha \\ \downarrow \\ A^* C \end{array} \mapsto \begin{array}{ccc} A & B \\ \square & \square \\ \alpha \\ \downarrow \\ C \end{array}$$

serves as an inverse for i by rigidity. Hence, we conclude that

$$N_C^{A,B} = \dim \text{Hom}(C, A \otimes B) = \dim \text{Hom}(A^* \otimes C, B) = N_B^{A^*,C}.$$

The third equality in Frobenius reciprocity follows from an identical argument, and hence we conclude the proof. \square

Corollary 1.58. *Let \mathcal{C} be a fusion category. Let $A, B \in \mathcal{C}$ be simple objects. We find that*

$$N_{\mathbf{1}}^{A,B} = N_{\mathbf{1}}^{B,A} = \begin{cases} 1 & B \cong A^* \\ 0 & \text{otherwise.} \end{cases}$$

Proof. This follows from Frobenius reciprocity and Schur's lemma:

$$N_{\mathbf{1}}^{A,B} = N_A^{A^*, \mathbf{1}} = \dim(\text{Hom}(A, A^*)) = \begin{cases} 1 & B \cong A^* \\ 0 & \text{otherwise.} \end{cases}$$

□

Corollary 1.59. *If \mathcal{C} is a fusion category, then $A \cong A^{**}$ for all $A \in \mathcal{C}$.*

Proof. Since $N_{\mathbf{1}}^{A^*, A^{**}} > 0$, we conclude that $A^{**} \cong A$ by the (iii) \implies (i) implication in proposition 1.57. □

Remark 1.60. Despite corollary 1.59, we *cannot* conclude that every fusion category admits a pivotal structure. The isomorphisms $A \cong A^{**}$ may fail to form a monoidal natural transformation. It is an open problem whether or not every fusion category admits a pivotal structure, and it is furthermore an open problem whether every fusion category admits a spherical structure [ENO05].

1.3.2 Trace

The first structure to define in the theory of modular categories is the *trace*. Let \mathcal{C} be a spherical fusion category. Given any object $A \in \mathcal{C}$ and any endomorphism $f : A \rightarrow A$, we define the *trace of f* by the following formula:

$$\text{tr}(f) = A^* \left(\begin{array}{c} A \\ f \\ A \end{array} \right)$$

Initially, the trace is a morphism, $\text{tr}(f) : \mathbf{1} \rightarrow \mathbf{1}$. However, we will choose to think of the trace of a morphism as a *complex number*, $\text{tr}(f) \in \mathbb{C}$. This can be done because the definition of a fusion category $\text{End}(\mathbf{1}) \cong \mathbb{C}$. This isomorphism can be made canonical by identifying an endomorphism $g \in \text{End}(\mathbf{1})$ with the unique $\lambda \in \mathbb{C}$ such that $g = \lambda \cdot \text{id}_{\mathbf{1}}$.

Remark 1.61. The trace is used mainly as a tool for linearization. Morphisms and objects are hard to describe, but the trace is a complex number.

Proposition 1.62. *Let \mathcal{C} be a spherical fusion category. For all $A, B \in \mathcal{C}$, $f \in \text{End}(A)$ the following claims are all true:*

1. $\text{tr} : \text{End}(A) \rightarrow \mathbb{C}$ is a linear map of vector spaces,
2. $\text{tr}(f^*) = \overline{\text{tr}(f)}$,
3. $\text{tr}(f \oplus g) = \text{tr}(f) + \text{tr}(g)$ for all $g \in \text{End}(B)$,
4. $\text{tr}(f \otimes g) = \text{tr}(f) \cdot \text{tr}(g)$ for all $g \in \text{End}(B)$,
5. $\text{tr}(h \circ g) = \text{tr}(g \circ h)$ for all $g : A \rightarrow B$, $h : B \rightarrow A$.

6. Trace is preserved by functors. That is, let \mathcal{C}, \mathcal{D} be spherical categories with traces $\text{tr}_{\mathcal{C}}, \text{tr}_{\mathcal{D}}$ respectively. Let $F : \mathcal{C} \rightarrow \mathcal{D}$ be a pivotal functor. We have that $\text{tr}_{\mathcal{C}}(f) = \text{tr}_{\mathcal{D}}(F(f))$;

Proof. We prove the claims one by one.

1. This follows immediately from the bilinearity of composition.
2. This is a straightforward computation.
3. WORK: do proof. This uses facts about the direct sum we haven't established yet.
4. We compute as follows:

$$\begin{aligned} \text{tr}(f \otimes g) &= \begin{array}{c} A \otimes B \\ \text{---} \\ f \otimes g \\ \text{---} \\ A \otimes B \end{array} = \begin{array}{c} A \quad B \\ \text{---} \\ f \quad g \\ \text{---} \\ A \quad B \end{array} \\ &= \begin{array}{c} A \quad B \\ \text{---} \\ f \quad g \\ \text{---} \\ A \quad B \end{array} = \text{tr}(f) \cdot \text{tr}(g). \end{aligned}$$

5. We compute as follows:

$$\text{tr}(f \circ g) = \begin{array}{c} A \\ \text{---} \\ g \\ \text{---} \\ f \\ \text{---} \\ A \end{array} = \begin{array}{c} B \\ \text{---} \\ f^* \\ \text{---} \\ g \\ \text{---} \\ B \end{array} = \begin{array}{c} A \\ \text{---} \\ f \\ \text{---} \\ g \\ \text{---} \\ A \end{array} = \text{tr}(g \circ f).$$

6. WORK: do the proof. It's not very hard, but involves a diagram and uses pivotality of the functor.

This completes the proof. \square

Corollary 1.63. Let $f : A \rightarrow A$ be an endomorphism in a fusion category \mathcal{C} . Fix a decomposition $A \cong \bigoplus_{i \in I} A_i$ of A into simple objects A_i . Moreover, we take the decomposition such that if $A_i \cong A_j$ then $A_i = A_j$. We can decompose

$$\text{Hom}(A, A) \cong \text{Hom}\left(\bigoplus_{i \in I} A_i, \bigoplus_{i \in I} A_i\right) = \bigoplus_{i \in I, j \in I} \text{Hom}(A_i, A_j).$$

Let M be the matrix whose columns and rows are labeled by I , and whose (i, j) entry is 0 if $A_i \not\cong A_j$ and $\lambda \cdot d_{A_i}$ if $A_i = A_j$, where $\lambda \in \mathbb{C}$ is the unique value such that the $\text{Hom}(A_i, A_j)$ component of f is $\lambda \cdot \text{id}_{A_i}$. We have that

$$\text{tr}_{\mathcal{C}}(f) = \text{tr}_{\mathbf{Vec}}(M).$$

Proof. Suppose that $A = A_0 \oplus A_1$ is the direct sum of two objects, not necessarily simple. By proposition [ref] we have a canonical decomposition

$$(A_0 \oplus A_1) \otimes (A_0 \oplus A_1)^* \cong (A_0 \otimes A_0^*) \oplus (A_0 \otimes A_1^*) \oplus (A_1 \otimes A_0^*) \oplus (A_1 \otimes A_1^*).$$

Suppose that $h : A \rightarrow A$ is an endomorphism. We can decompose $h = h_{A_0, A_0} + h_{A_0, A_1} + h_{A_1, A_0} + h_{A_1, A_1}$ as a sum of morphisms which restrict to maps $A_i \rightarrow A_j$. We find that $\text{coev}_{A \oplus B}$ restricts to a map whose codomain is $(A_0 \otimes A_0^*) \oplus (A_1 \otimes A_0)^*$ and similarly ev restricts to a map whose domain is $(A_0 \otimes A_0^*) \oplus (A_1 \otimes A_0)^*$, since $\text{coev}_{A \oplus B} = \text{coev}_A \oplus \text{coev}_B$ and $\text{ev}_{A \oplus B} = \text{ev}_A \oplus \text{ev}_B$.

Hence, in the definition of trace, we find that the cross terms $h_{A_0, A_1} + h_{A_1, A_0}$ act by zero since they send the codomain of $\text{coev}_{A \oplus B}$ to elements with no effect on the map $\text{ev}_{A \oplus B}$. Moreover, we compute in this way that $\text{tr}(h) = \text{tr}(h_{A_0, A_0}) + \text{tr}(h_{A_1, A_1})$.

In this way, the trace splits over direct sums and only picks out diagonal elements. Applying this result inductively reduces the proof to the case that A is a simple object. This follows directly from the definition of quantum dimension. \square

1.3.3 Quantum dimension and Frobenius-Perron dimension

WORK: include something about global quantum dimension \mathcal{D} .

Our next tool to discuss is the *quantum dimension*. We have already seen quantum dimensions appear in subsection [ref]. We recall that given any spherical fusion category \mathcal{C} and any simple object $A \in \mathcal{C}$, we the quantum dimension of A by the formula

$$d_A = A \bigcirc A^*$$

As usual, we identify d_A with a complex number via the canonical isomorphism $\text{End}(\mathbf{1}) \cong \mathbb{C}$. The quantum dimension is clearly equal to the trace of the identity map on A , $d_A = \text{tr}(\text{id}_A)$. The first properties of quantum dimension follow from our general analysis of trace:

Proposition 1.64. *For every spherical fusion category \mathcal{C} and any objects $A, B \in \mathcal{C}$, we have the following formulas:*

- (i) *If $A \cong B$, then $d_A = d_B$;*
- (ii) *$d_{A^*} = d_A$;*
- (iii) *$d_A \neq 0$;*
- (iv) *$d_A \in \mathbb{R}$.*

Proof.

- (i) Let $f : A \cong B$ be an isomorphism. Using Proposition [ref] we find

$$d_A = \text{tr}(\text{id}_A) = \text{tr}(f^{-1} \circ f) = \text{tr}(f \circ f^{-1}) = \text{tr}(\text{id}_B) = d_B.$$

- (ii) This follows from proposition [ref].

- (iii) From proposition [ref], we know that $A \otimes A^* \cong \mathbf{1} \oplus X$ for some $X \in \mathcal{C}$ which does not have any factors of $\mathbf{1}$ in its direct sum decomposition. The map $\text{coev}_A^R : \mathbf{1} \rightarrow A \otimes A^*$ is thus a non-zero scalar times the inclusion $\mathbf{1} \hookrightarrow \mathbf{1} \oplus X$, and the map $\text{ev}_A^L : A \otimes A^* \rightarrow \mathbf{1}$ is a non-zero scalar times the projection $\mathbf{1} \oplus X \rightarrow \mathbf{1}$. Since inclusion composed with projection is the identity, we find that $\text{ev}_A^L \circ \text{coev}_A^R$ is a non-zero scalar times the identity, as desired.
- (iv) WORK: There is reference to this fact in the MathOverflow question ‘Modular Tensor Categories: Reasoning behind the axioms’

□

Definition 1.65. For every simple object $A \in \mathcal{C}$ in a fusion category, we define the *fusion matrix* N^A corresponding to A to be the matrix whose rows and columns are labeled by \mathcal{L} , and whose (B, C) entry is $N_C^{A,B}$. As a linear operator, this means that N^A acts as follows:

$$N^A : \mathbb{C}[\mathcal{L}] \rightarrow \mathbb{C}[\mathcal{L}],$$

$$|[B]\rangle \mapsto \sum_{[C] \in \mathcal{L}} N_C^{A,B} |[C]\rangle$$

Proposition 1.66. Let \mathcal{C} be a spherical fusion category.

(i) Let $A, B \in \mathcal{C}$ be simple objects. We have that

$$d_A d_B = \sum_{[C] \in \mathcal{L}} N_C^{A,B} d_C.$$

(ii) Define $\mathbf{d} = \sum_{[B] \in \mathcal{L}} d_B |[B]\rangle \in \mathbb{C}[\mathcal{L}]$. We have that

$$N^A \mathbf{d} = d_A \mathbf{d}$$

for all simple objects A .

Proof. From proposition [ref], we have an isomorphism

$$A \otimes B \cong \bigoplus_{[C] \in \mathcal{L}} N_C^{A,B} \cdot C$$

and thus

$$\text{tr}(\text{id}_{A \otimes B}) = \text{tr}(\text{id}_{\bigoplus_{[C] \in \mathcal{L}} N_C^{A,B} \cdot C}).$$

Expanding using the rules in proposition [ref] gives part (i). Part (ii) follows from expanding the definition of the linear operator and applying part (i). □

Remark 1.67. Proposition 1.66 tells us that d_A is an eigenvalue of N^A . This gives a strong restriction on d_A . A first restriction is that d_A is an algebraic number, since it is the root of the characteristic polynomial of N^A .

Remark 1.68. The question is whether or not the quantum dimensions are *positive* real numbers. We recall that we defined a unitarizable spherical fusion category to be one in which the quantum dimensions are all positive. It is at this point that this becomes relevant. In particular, if \mathcal{C} is unitarizable then its quantum dimensions are eigenvalues of N^A , and their corresponding eigenvector \mathbf{d} has positive real entries. **WORK: Not sure how much of this stuff I will have already said before.**

Theorem 1.69 (Frobenius-Perron theorem, [EGNO16]). *Let B be a square matrix with nonnegative real entries.*

- (i) *B has a non-negative real eigenvalue. The largest non-negative real eigenvalue $\lambda(B)$ of B dominates the absolute values of all other eigenvalues μ of B : $|\mu| \leq \lambda(B)$. Moreover, there is an eigenvector of B with non-negative entries and eigenvalue $\lambda(B)$.*
- (ii) *If B has strictly positive entries then $\lambda(B)$ is a simple positive eigenvalue, and the corresponding eigenvector can be normalized to have strictly positive entries. Moreover, $|\mu| < \lambda(B)$ for any other eigenvalue μ of B .*
- (iii) *If a matrix B with non-negative entries has an eigenvector v with strictly positive entries, then the corresponding eigenvalue is $\lambda(B)$.*

Remark 1.70. We call the largest positive real eigenvalue of a matrix its *Frobenius-Perron eigenvalue*.

WORK: Not sure what direction I should take for statements about unitary fusion categories like corollary 1.71. Should I state them for unitary fusion categories, or for unitarizable fusion categories, which have less baggage (I can define a unitarizable fusion category to be one with positive quantum dimensions, and then prove that that a fusion category admits a unitary structure if and only if it is unitarizable).

Corollary 1.71. *Let \mathcal{C} be a unitarizable spherical fusion category. Let $A \in \mathcal{C}$ be a simple object. The quantum dimension d_A is equal to the Frobenius-Perron eigenvalue of N^A .*

Proof. Since \mathcal{C} is unitarizable, the vector $\mathbf{d} = \sum_{[B] \in \mathcal{L}} d_B |[B]\rangle \in \mathbb{C}[\mathcal{L}]$ has positive entries and has eigenvalue d_A . Hence, d_A is the Frobenius-Perron eigenvalue of N^A as desired. \square

Definition 1.72. Let $A \in \mathcal{C}$ be a simple object in a spherical fusion category. We define the *Frobenius-Perron dimension* of A (denoted $\text{FPdim}(A)$) by the formula

$$\text{FPdim}(A) = (\text{Frobenius-Perron eigenvalue of } N^A).$$

Remark 1.73. In light of corollary 1.71, $\text{FPdim}(A) = d_A$ whenever \mathcal{C} is unitarizable. However, in this chapter we will mostly work with spherical fusion categories with no conditions on unitarizability. Thus, the definition of $\text{FPdim}(A)$ is not redundant and is at times useful. An additional observation is that the definition of quantum dimension strongly uses the spherical structure on \mathcal{C} . However, the Frobenius-Perron dimension only uses the fusion coefficients, and those are well-defined in any fusion category. Hence, the Frobenius-Perron dimension also derives utility from being applicable in a broader set of situations than the quantum dimension.

Proposition 1.74. *Let \mathcal{C} be a fusion category, and let $A \in \mathcal{C}$ be a simple object.*

$$(i) \text{ FPdim}(A) = \lim_{n \rightarrow \infty} \dim(\text{Hom}(A^{\otimes n}, A^{\otimes n}))^{1/(2n)}$$

$$(ii) \text{ FPdim}(A) = \lim_{n \rightarrow \infty} \dim(\text{Hom}(\mathbf{1}, A^{\otimes n}))^{1/n}$$

(iii)

$$\text{FPdim}(A) = \lim_{n \rightarrow \infty} (\# \text{ of simple objects in the direct sum decomposition of } A^{\otimes n})^{1/n}.$$

Proof. WORK: I can do a good part of this when \mathcal{C} is unitarizable, so that its largest eigenvalue is strictly larger than all of the others. When there are multiple large eigenvalues all of the same size then the proofs go wrong. Is there something about the structure of N^A I can exploit? Are these theorems true for fusion categories, or do I need to pass to unitarizable fusion categories? \square

Remark 1.75. Proposition 1.74 gives an alternate interpretation of the Frobenius-Perron dimension in terms of growth in tensor powers. This sort of alternate perspective of dimension applies to several types of objects outside the scope of tensor category theory [COT24].

This proposition can be interpreted as saying that the simple object A has $\text{FPdim}(A)$ internal degrees of freedom “on average”. Elements of the vector space $\text{Hom}(\mathbf{1}, A^{\otimes n})$ correspond to states in the system with n anyons of type A arranged in a line. If the internal configuration space of each anyon was $\text{FPdim}(A)$ -dimensional, then the overall dimension would be $\text{FPdim}(A)^n$. By Proposition [ref], $\text{FPdim}(A)^n$ is approximately $\text{Hom}(\mathbf{1}, A^{\otimes n})$ for large n . Hence, each anyon has approximately $\text{FPdim}(A)$ internal degrees of freedom. Of course, $\text{FPdim}(A)$ has no reason to be an integer! In the Fibonacci theory $\text{FPdim}(\tau) = \phi = 1.61\dots$ Frobenius-Perron dimension just gives an average amount for large values.

WORK: re-do this explanation way better + add diagram for it.

1.3.4 Twist

In this section we will discuss *twists*. The twist is a subtle concept, which we have not explicitly mentioned up to now. The idea is that anyons can *rotate in place*. Since the space of endomorphisms of an anyon is one dimensional, this rotation must act by a phase. This phase is physically relevant, and can be measured in experiment.

For example, consider the Y -type on the toric code. It consists of the fusion of an X -type anyon and a Z -type anyon, as shown below:

WORK: add figure of Y as a thick X and Z together; could be hard to draw these nice

Twisting Y in place will correspond to twisting X and Z around each other. This twisting thus results in a phase of -1 . In general, we can imagine anyons as having some thickness to them. Anyons are not localized at points - they are localized at small regions. Twisting this region all the way around can be viewed visually as

WORK: twisted anyon.

This is the twist. One way of working with the twist is to work with thickened diagrams, where strings are replaced with ribbons. While popular in some parts of the literature, we will continue to work with string diagrams for simplicity. The key observation is that the twist can be constructed using string diagrammatic structures we already have as follows:

WORK: twist as a swirl diagram, compared with ribbon.

Hence, letting \mathcal{C} be a pre-modular fusion category, we *define* the twist θ_A of an object $A \in \mathcal{C}$ to be

$$\begin{array}{c} A \\ | \\ \boxed{\theta_A} \\ | \\ A \end{array} = \begin{array}{c} A \\ | \\ \circlearrowleft \\ | \\ A \end{array}$$

For every simple object $A \in \mathcal{C}$, the map $\theta_A \in \text{End}(A)$ can be identified with the unique complex number λ such that $\theta_A = \lambda \cdot \text{id}_A$. Equivalently, we can identify θ_A with the complex number $\lambda = \text{tr}(\theta_A)/d_A$ which gives the graphical formula

$$\theta_A = \frac{1}{d_A} \bigcirc\!\!\!\bigcirc A$$

We can reinterpret all other twist-like maps in terms of θ :

Lemma 1.76. *Let \mathcal{C} be a pre-modular fusion category. We have that*

$$\begin{array}{ccc} \begin{array}{c} A \\ | \\ \circlearrowleft \\ | \\ A \end{array} & = & \begin{array}{c} A \\ | \\ \boxed{\theta_A} \\ | \\ A \end{array}, \\ \begin{array}{c} A \\ | \\ \circlearrowright \\ | \\ A \end{array} & = & \begin{array}{c} A \\ | \\ \boxed{\theta_A^{-1}} \\ | \\ A \end{array}, \\ \begin{array}{c} A \\ | \\ \circlearrowleft \\ | \\ A \end{array} & = & \begin{array}{c} A \\ | \\ \boxed{\theta_A} \\ | \\ A \end{array}, \\ \begin{array}{c} A \\ | \\ \circlearrowright \\ | \\ A \end{array} & = & \begin{array}{c} A \\ | \\ \boxed{\theta_A^{-1}} \\ | \\ A \end{array} \end{array}$$

Proof. To begin we show that

$$\begin{array}{c} A \\ | \\ \circlearrowleft \\ | \\ A \end{array} = \begin{array}{c} A \\ | \\ \circlearrowright \\ | \\ A \end{array}.$$

When A is simple, this follows from the spherical axiom. Taking the trace of both sides gives the same formula for θ_A as a figure-eight. Additionally, pushing through duals it is clear that both sides in the above proposed equality are natural isomorphisms. Natural isomorphisms are determined by their action on simple objects because they commute with direct sums. Hence, we conclude that the sides are equal for all objects.

To get that the two reversed formulas are equal to θ_A^{-1} , it suffices to compose with θ_A and use string-diagram manipulations to show that it results in the identity. This is a simple exercise and is left as an exercise to the reader. \square

We now characterize the key properties of the twist:

Proposition 1.77. *Let \mathcal{C} be a pre-modular fusion category. The twists θ induce a monoidal natural isomorphism $\text{id}_{\mathcal{C}} \xrightarrow{\sim} \text{id}_{\mathcal{C}}$. Additionally, θ satisfies the identity*

$$\theta_{A \otimes B} = \beta_{B,A} \circ \beta_{A,B} \circ (\theta_A \otimes \theta_B)$$

for all $A, B \in \mathcal{C}$, and $\theta_{A^*} = (\theta_A)^*$.

Proof. Naturality of θ follows from pushing through duals. The formula $\theta_{A \otimes B} = \beta_{B,A} \circ \beta_{A,B} \circ (\theta_A \otimes \theta_B)$ comes from manipulating string diagrams to get the equation

Finally, $\theta_{A^*} = (\theta_A)^*$ comes from the string-diagram manipulation and proposition [ref]:

as desired. \square

The naive reason to care about twists is that they describe a physically relevant quantity and hence should be studied. The more subtle reason to care about twists is that they are the most efficient way of encoding the spherical structure on \mathcal{C} . A spherical structure is first and foremost a pivotal structure, meaning that it has a right and left rigid structure which are compatible. Given a spherical structure one can always obtain twists. Conversely, given a right-rigid structure and twists one can recover the left-rigid structure via the formulas

In this way, giving a spherical structure on a right-rigid monoidal category is the *same* as giving a twist structure. This is codified in the following lemma:

Proposition 1.78 (Deligne's twisting lemma, [Yet92]). *Let \mathcal{C} be a right-rigid braided monoidal category. Every pivotal structure on \mathcal{C} naturally gives a twist natural transformation $\theta : \text{id}_{\mathcal{C}} \rightarrow \text{id}_{\mathcal{C}}$. This assignment induces a canonical bijection between the set of pivotal structures on \mathcal{C} and the set of natural isomorphisms $\theta : \text{id}_{\mathcal{C}} \rightarrow \text{id}_{\mathcal{C}}$ satisfying $\theta_{A \otimes B} = \beta_{B,A} \circ \beta_{A,B} \circ (\theta_A \otimes \theta_B)$ for all $A, B \in \mathcal{C}$.*

Moreover, restricting the assignment to the space of spherical structures on \mathcal{C} induces a canonical bijection between the set of spherical structures on \mathcal{C} and the set of isomorphisms $\theta : \text{id}_{\mathcal{C}} \rightarrow \text{id}_{\mathcal{C}}$ satisfying $\theta_{A \otimes B} = \beta_{B,A} \circ \beta_{A,B} \circ (\theta_A \otimes \theta_B)$ for all $A, B \in \mathcal{C}$ and $\theta_{A^} = (\theta_A)^*$.*

Proof. We already showed in proposition [ref] that every spherical category gives a twist natural transformation satisfying the desired axioms. Restricting the proof to only a possibly non-spherical pivotal category still gives a twist natural transformation satisfying $\theta_{A \otimes B} = \beta_{B,A} \circ \beta_{A,B} \circ (\theta_A \otimes \theta_B)$ for all $A, B \in \mathcal{C}$. The heart of the proof is showing that the formulas [ref] induce pivotal and spherical structures with the twist satisfies the right axioms. The process of inducing a pivotal structure and inducing a twist are inverses to one another because

$$\begin{array}{ccc} A & & A \\ \downarrow \theta_A & = & \text{---} \\ A & & A \end{array}$$

To begin, we assume that $\theta_{A \otimes B} = \beta_{B,A} \circ \beta_{A,B} \circ (\theta_A \otimes \theta_B)$ and we seek to prove that the corresponding ev^L , coev^L maps induce a pivotal structure. We first axiom of pivotality follows from use of the axiom $\theta_{A \otimes B} = \beta_{B,A} \circ \beta_{A,B} \circ (\theta_A \otimes \theta_B)$:

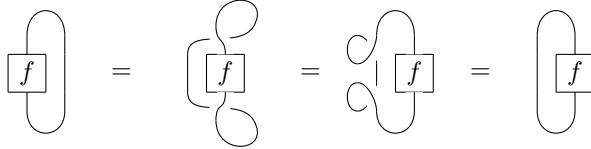
$$\begin{array}{c} B^* \otimes A^* \quad A \otimes B \\ \text{---} \end{array} = \begin{array}{c} B^* A^* \quad A \quad B \\ \text{---} \quad \theta_{A \otimes B} \end{array} = \begin{array}{c} B^* A^* \quad A \quad B \\ \theta_A \quad \theta_B \end{array}$$

$$= \begin{array}{c} B^* A^* \quad A \quad B \\ \theta_A \quad \theta_B \end{array} = \begin{array}{c} B^* A^* \quad A \quad B \\ \theta_A \quad \theta_B \end{array} = \begin{array}{c} B^* \quad A^* \quad A \quad B \\ \text{---} \end{array}$$

The second axiom of pivotality follows from the use of the naturality of θ :

$$\begin{array}{c} A^* \\ \text{---} \quad B \\ f \\ A \end{array} = \begin{array}{c} A^* \\ \text{---} \quad B^* \\ \theta_B^{-1} \\ f \\ \theta_A \\ \text{---} \quad B^* \end{array} = \begin{array}{c} A^* \\ \text{---} \quad B \\ f \\ A \end{array}$$

Finally, we assume that $(\theta_A)^* = \theta_{A^*}$ and we seek to prove the spherical axiom. Taking the dual of theta we can get all of the equalities in Lemma [ref]. Applying them we get that



as desired. \square

1.3.5 Functors, natural transformations, and equivalence

In this section, we will talk about functors, natural transformations, and equivalences between fusion, spherical, pre-modular, and modular categories. Given a topological order, there is *not* a unique modular category describing it. There is a unique modular category *up to equivalence*. Hence, the notion of equivalence of categories is baked into our physics-math correspondance so it is important that we state it explicitly.

Functors which do not induce equivalences of categories are also physically relevant. In certain contexts, a functor $F : \mathcal{C} \rightarrow \mathcal{D}$ is used to model a *phase transition* from \mathcal{C} to \mathcal{D} . We will see a lot more functors and natural transformations between modular categories throughout the book, especially in chapter [ref].

Even though structures in categories require a lot of compatibility conditions, the conditions on the functors do not. This means that we have the following:

- The correct notion of functor between fusion categories is \mathbb{C} -linear monoidal functor. There is no compatibility condition required between the \mathbb{C} -linear structure and the monoidal structure. The correct notion of natural transformation between \mathbb{C} -linear monoidal functors is a monoidal natural transformation.
- The correct notion of functor between spherical fusion categories is \mathbb{C} -linear pivotal monoidal functor. There is no compatibility condition required between the \mathbb{C} -linear structure and the pivotal monoidal structure. The correct notion of natural transformation is monoidal natural transformation.
- The correct notion of functor between pre-modular categories is \mathbb{C} -linear pivotal braided monoidal functor. There is no compatibility condition required between the \mathbb{C} -linear structure, pivotal monoidal structure, or braided monoidal structure. The correct notion of natural transformation is monoidal natural transformation.
- The correct notions of functors/natural transformations for modular categories are the same as for pre-modular categories.

WORK: : this section is very short. I don't have much to say, actually. Should this be moved? Maybe I keep a very short section? I don't know.

1.3.6 Deligne tensor product

In the theory of any class of mathematical object, an important consideration is the ways in which examples can be put together to give new examples. In the case of fusion categories, this basic operation is known as the *Deligne tensor product*. Given any fusion categories \mathcal{C}, \mathcal{D} , their Deligne tensor product $\mathcal{C} \boxtimes \mathcal{D}$ is a new fusion category. The Deligne tensor product of spherical fusion categories will be equipped with the structure of a spherical

fusion category, and the Deligne tensor product of (pre-)modular categories will be equipped with the structure of a (pre-)modular category.

Physically, the Deligne tensor product corresponds to *stacking*. Consider two sheets of material. We choose two modular categories \mathcal{C}, \mathcal{D} . We endow the top sheet with the structure of a topologically ordered quantum system described by \mathcal{C} and we endow the bottom with the structure of a topologically ordered quantum system described by \mathcal{D} . The algebraic description of this bilayer system is $\mathcal{C} \boxtimes \mathcal{D}$. This can be viewed as the physical definition of $\mathcal{C} \boxtimes \mathcal{D}$.

WORK: add bilayer system diagram

We now mathematically define the Deligne tensor product.

Definition 1.79. Let \mathcal{C}, \mathcal{D} be \mathbb{C} -linear categories, isomorphic as \mathbb{C} -linear categories to $\text{Vec}_{\mathbb{C}}^n, \text{Vec}_{\mathbb{C}}^m$ respectively. We define a Deligne tensor product of \mathcal{C} and \mathcal{D} to be the following data:

1. A \mathbb{C} -linear category $\mathcal{C} \boxtimes \mathcal{D}$;
2. A \mathbb{C} -linear functor $\mathcal{C} \times \mathcal{D} \rightarrow \mathcal{C} \boxtimes \mathcal{D}$.

Such that:

1. Every object $X \in \mathcal{C} \boxtimes \mathcal{D}$ has a direct sum decomposition

$$X \cong \bigoplus_{i=1}^n A_i \boxtimes B_i$$

for some $n \geq 1, A_i \in \mathcal{C}, B_i \in \mathcal{D}$.

2. There is an equality of vector spaces

$$\text{Hom}_{\mathcal{C} \boxtimes \mathcal{D}}(A \boxtimes B, A' \boxtimes B') = \text{Hom}_{\mathcal{C}}(A, A') \otimes \text{Hom}_{\mathcal{D}}(B, B').$$

3. Given any $A, A', A'' \in \mathcal{C}, B, B', B'' \in \mathcal{D}, f : A \rightarrow A', f' : A' \rightarrow A'', g : B \rightarrow B', g' : B' \rightarrow B''$, the diagram

$$\begin{array}{ccccc} A \boxtimes B & \xrightarrow{f \boxtimes g} & A' \boxtimes B' & \xrightarrow{f' \boxtimes g'} & A'' \boxtimes B'' \\ & \searrow & & \nearrow & \\ & & (f' \circ f) \boxtimes (g' \circ f) & & \end{array}$$

commutes.

We now state the main existence/uniqueness result about the Deligne tensor product:

Proposition 1.80. Let \mathcal{C}, \mathcal{D} be \mathbb{C} -linear categories isomorphic as \mathbb{C} -linear categories to $\text{Vec}_{\mathbb{C}}^n$ and $\text{Vec}_{\mathbb{C}}^m$ respectively. There exists a Deligne product $\mathcal{C} \boxtimes \mathcal{D}$ for \mathcal{C} and \mathcal{D} . Moreover, given any other deligne tensor product $\mathcal{C} \boxtimes' \mathcal{D}$ of \mathcal{C} and \mathcal{D} there exists a unique functor $F : \mathcal{C} \boxtimes \mathcal{D} \rightarrow \mathcal{C} \boxtimes' \mathcal{D}$ making the diagram

$$\begin{array}{ccc} \mathcal{C} \times \mathcal{D} & \longrightarrow & \mathcal{C} \boxtimes \mathcal{D} \\ & \searrow & \downarrow F \\ & & \mathcal{C} \boxtimes' \mathcal{D} \end{array}$$

commute. This functor is an equivalence of categories.

Proof. It is clear that $\mathbf{Vec}_{\mathbb{C}}^n \boxtimes \mathbf{Vec}_{\mathbb{C}}^m = \mathbf{Vec}_{\mathbb{C}}^{nm}$. Every equivalence of categories $\mathcal{C} \rightarrow \mathcal{C}'$ induces an equivalence of categories $\mathcal{C} \boxtimes \mathcal{D} \rightarrow \mathcal{C}' \boxtimes \mathcal{D}$. Hence, since \mathcal{D} and \mathcal{D}' are equivalent to $\mathbf{Vec}_{\mathbb{C}}^n$ and $\mathbf{Vec}_{\mathbb{C}}^m$ respectively, their Deligne tensor product exists and is equivalent to $\mathbf{Vec}_{\mathbb{C}}^{nm}$.

Any functor making the diagram commute must send $A \boxtimes B$ to $A \boxtimes' B$. The definition of Deligne tensor produce tells us this is enough to conclude that the map is an equivalence of categories, since axiom 3 this map is always a functor, axiom 2 implies it is fully faithful, and axiom 1 implies it is essentially surjective, and hence we can apply proposition [ref]. \square

Now that we have defined the Deligne tensor product of \mathbb{C} -linear categories equivalent to $\mathbf{Vec}_{\mathbb{C}}^n$, we move on to defining the Deligne tensor product of fusion categories, spherical fusion categories, pre-modular categories, and modular categories.

Proposition 1.81. *The following claims are all true.*

- (i) *Let \mathcal{C}, \mathcal{D} be fusion categories. On the level of objects, define a monoidal structure $\mathcal{C} \boxtimes \mathcal{D}$ by the formula*

$$(A \boxtimes B) \otimes (A' \boxtimes B') = (A \otimes A') \boxtimes (B \otimes B').$$

Along with a natural choice of action of the tensor product on morphisms, unit $\mathbf{1}_{\mathcal{C} \boxtimes \mathcal{D}} = \mathbf{1}_{\mathcal{C}} \boxtimes \mathbf{1}_{\mathcal{D}}$, and a natural choice of associator and unitors, this induces the structure of a monoidal category on \mathcal{C} .

Define a right-rigid structure on $\mathcal{C} \boxtimes \mathcal{D}$ as follows. The dual of an object $A \boxtimes B$ is $A^ \boxtimes B^*$. Define $\text{ev}_{A \boxtimes B} = \text{ev}_A \boxtimes \text{ev}_B$, $\text{coev}_{A \boxtimes B} = \text{coev}_A \boxtimes \text{coev}_B$. This induces a well-defined right-rigid structure on $\mathcal{C} \boxtimes \mathcal{D}$.*

The above definitions induce the structure of a fusion category on $\mathcal{C} \boxtimes \mathcal{D}$.

- (ii) *Let \mathcal{C}, \mathcal{D} be spherical fusion categories. The evaluation and coevaluation maps $\text{ev}_{A \boxtimes B}^L = \text{ev}_A^L \boxtimes \text{ev}_B^L$ and $\text{coev}_{A \boxtimes B}^L = \text{coev}_A^L \boxtimes \text{coev}_B^L$ induce a left-rigid structure on $\mathcal{C} \boxtimes \mathcal{D}$. Along with the canonical structure of a fusion category on $\mathcal{C} \boxtimes \mathcal{D}$, this induces the structure of a spherical fusion category on $\mathcal{C} \boxtimes \mathcal{D}$.*
- (iii) *Let \mathcal{C}, \mathcal{D} be pre-modular categories. The braiding map $\beta_{\mathcal{C} \boxtimes \mathcal{D}} = \beta_{\mathcal{C}} \boxtimes \beta_{\mathcal{D}}$ induces the structure of a pre-modular category on $\mathcal{C} \boxtimes \mathcal{D}$. The product $\mathcal{C} \boxtimes \mathcal{D}$ is modular if and only if \mathcal{C}, \mathcal{D} are both modular.*

Proof. Given any of the above structures, all of the axioms on $\mathcal{C} \boxtimes \mathcal{D}$ immediately follow from their respective axioms on \mathcal{C} and \mathcal{D} . Hence, the proof is an exercise in recalling definitions which we omit. \square

1.4 The category of G -graded G -representations

1.4.1 Overview

We've talked about a lot of general theory of modular categories. It's time for us to focus on our main family of *examples*. Namely, the categories $\mathfrak{D}(G)$ of G -graded G -representations. These categories describe discrete gauge theory based on the finite group G .

Before we can prove that $\mathfrak{D}(G)$ is a modular category, we need to endow $\mathfrak{D}(G)$ with the necessary structures. In particular, we will endow $\mathfrak{D}(G)$ with \mathbb{C} -linear, monoidal, braided,

right-rigid, and left-rigid structures. We will need to show that all of these structures are compatible with each other in the necessary ways, and that $\mathfrak{D}(G)$ satisfies the non-degeneracy axiom. We will use this as an opportunity to introduce tools of general use for proving that categories satisfy the axioms of a modular category.

Additionally, we will also study two categories similar to $\mathfrak{D}(G)$ which will serve as extra examples to get our grip on definitions. These categories will also appear later as relevant in and of themselves. The first is \mathbf{Vec}_G , the category of G -graded vector spaces. It is defined as follows:

WORK: define \mathbf{Vec}_G in terms of objects and composition.

Our second structure of interest is $\text{Rep}(G)$, the category of G -representations. It is defined as follows:

WORK: define $\text{Rep}(G)$ in terms of objects and composition.

We will show that both \mathbf{Vec}_G and $\text{Rep}(G)$ can be naturally equipped with the structures of spherical fusion categories. We then show that $\text{Rep}(G)$ admits a braiding which turns it into a pre-modular category. This braiding is symmetric in the sense that $\beta_{B,A} \circ \beta_{A,B} = \text{id}_A \otimes \text{id}_B$ for all $A, B \in \mathcal{C}$, and hence $\text{Rep}(G)$ is not a modular category. The category \mathbf{Vec}_G is shown to not admit a braiding whenever G is non-abelian.

1.4.2 Higher linear algebra

WORK: In this section we define the \mathbb{C} -linear structures on \mathbf{Vec}_G , $\text{Rep}(G)$, and $\mathfrak{D}(G)$. Our goal is to show that they are all equivalent to $\mathbf{Vec}_{\mathbb{C}}^n$ for some $n \geq 1$.

WORK: It seems like the best approach is through higher linear algebra. Namely, we show that if \mathcal{C} is abelian, \mathbb{C} -linear, semisimple, and has finitely many isomorphism classes of simple objects then it must be isomorphic to $\mathbf{Vec}_{\mathbb{C}}^n$. Its a good time to wax philosophical about higher linear algebra and 2-vector spaces. However, its not clear that this approach actually helps at all. It might be easier to immediately note that everybody is the direct sum of irreducibles, prove a Schur's lemma, and call it a day. Of course these approaches are all equivalent but its not clear what's best.

WORK: Here's a lemma I would like to use. Since \times and \boxtimes look too similar, I should use \boxplus to denote the Cartesian product and call it the direct sum. I should then prove this:

Lemma 1.82. *Consider the category $\mathbf{2Vec}$ whose objects are \mathbb{C} -linear categories equivalent to $\mathbf{Vec}_{\mathbb{C}}^n$ for some $n \geq 0$, and whose morphisms are \mathbb{C} -linear functors. The direct sum \boxplus is a biproduct in $\mathbf{2Vec}$*

Proof. WORK: do proof □

1.4.3 Spherical fusion structures

WORK: show that the categories have duals and monoidal structure. This should be pretty easy and painless. Pentagon identity should follow from the pentagon identity on $\mathbf{Vec}_{\mathbb{C}}$.

1.4.4 Braiding and modularity

WORK: Introduce braidings. Show that $\text{Rep}(G)$ is symmetric. Show that \mathbf{Vec}_G does not admit a braiding if G is not abelian and does admit a symmetric braiding if G is abelian. Show that $\mathfrak{D}(G)$ admits a non-degenerate braiding.

1.5 The modular representation

1.5.1 Definition

In this chapter we are going to talk about the *modular representations* of modular categories. Here's the point. Let \mathcal{C} be a modular category. Let \mathcal{L} be the set of isomorphism classes of simple objects of \mathcal{C} . We will define a group homomorphism

$$\rho_{\mathcal{C}} : \mathrm{SL}_2(\mathbb{Z}) \rightarrow \mathrm{Aut}(\mathbb{C}[\mathcal{L}])$$

associated to \mathcal{C} , where $\mathrm{SL}_2(\mathbb{Z})$ is the group of 2-by-2 matrices with integer coefficients and unit determinant. The group $\mathrm{SL}_2(\mathbb{Z})$ is sometimes known as the *modular group*, due to its connection with moduli spaces of elliptic curves. Hence, $\rho_{\mathcal{C}}$ is known as the *modular representation* of \mathcal{C} .

The goal of this chapter is to introduce $\rho_{\mathcal{C}}$, show it is well defined, and then prove a series of theorems related to $\rho_{\mathcal{C}}$.

Before defining $\rho_{\mathcal{C}}$, we recall the basic group theory of $\mathrm{SL}_2(\mathbb{Z})$. It is generated by the matrices

$$s = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}, \quad t = \begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix}.$$

These two matrices satisfy the relations $s^2 = -1$ and $(st)^3 = -1$, where 1 is used to represent the identity matrix. These relations generate $\mathrm{SL}_2(\mathbb{Z})$, in the sense that we have the following presentation:

Proposition 1.83. *The following presentation is valid:*

$$\mathrm{SL}_2(\mathbb{Z}) = \langle s, t \mid s^4 = 1, (st)^3 = s^2 \rangle.$$

Proof. This is a standard fact about $\mathrm{SL}_2(\mathbb{Z})$. See for instance [ref] □

Hence, to define a homomorphism $\rho_{\mathcal{C}} : \mathrm{SL}_2(\mathbb{Z}) \rightarrow \mathrm{Aut}(\mathbb{C}[\mathcal{L}])$ it suffices to choose automorphisms $\rho_{\mathcal{C}}(s)$, $\rho_{\mathcal{C}}(t)$ of $\mathbb{C}[\mathcal{L}]$, and show that they satisfy the relations $\rho_{\mathcal{C}}(s)^4 = 1$ and $(\rho_{\mathcal{C}}(s)\rho_{\mathcal{C}}(t))^3 = \rho_{\mathcal{C}}(s)^2$. Since $\mathbb{C}[\mathcal{L}]$ has a canonical basis, we can think of its automorphisms as being matrices with rows and columns labeled by \mathcal{L} . We define an operator $S : \mathbb{C}[\mathcal{L}] \rightarrow \mathbb{C}[\mathcal{L}]$ via the matrix coefficients

$$S_{A,B} = A^* \bigcircledcirc B^*$$

We next define the matrix $T : \mathbb{C}[\mathcal{L}] \rightarrow \mathbb{C}[\mathcal{L}]$ to be the diagonal matrix with $([A], [A])$ -entry θ_A , for all $[A] \in \mathcal{L}$.

As currently stated, the S and T matrices defined do not satisfy $S^4 = 1$ and $(ST)^3 = S^2$. They only satisfy these formula up to phases in \mathbb{C} . They still need to be normalized before we can define $\rho_{\mathcal{C}}$. The normalization factors come in terms of the *Gauss sums*,

$$p_{\mathcal{C}}^{\pm} = \sum_{[A] \in \mathcal{L}} \theta_A^{\pm 1} d_A^2.$$

We can now state the main theorem of this chapter:

Theorem 1.84. Let \mathcal{C} be a modular category. The values $p_{\mathcal{C}}^+$ and $p_{\mathcal{C}}^-$ are nonzero, and the map

$$\begin{aligned}\rho_{\mathcal{C}} : \mathrm{SL}_2(\mathbb{Z}) &\rightarrow \mathrm{Aut}(\mathbb{C}[\mathcal{L}]) \\ s &\mapsto \frac{1}{D} \cdot S \\ t &\mapsto (p_{\mathcal{C}}^-/p_{\mathcal{C}}^+)^{1/6} \cdot T\end{aligned}$$

is a group homomorphism.

We will prove this theorem and motivate why it should be true over the course of this chapter. We will also prove key facts about the image and kernel of this representation, as well as other formulas of interest relating to twists, S-matrix entries, and Gauss sums.

1.5.2 Torus perspective

It's good to reflect on why MCs have $\mathrm{SL}_2(\mathbb{Z})$ representations associated with them in the first place. Not only does the representation exist, but it is so fundamental to the modular category that it is chosen as the namesake. This begs the question. What's going on?

The answer has to do with topological phases on the torus.

WORK: add torus

Every modular category \mathcal{C} is supposed to describe a topological order. Up to now we have only considered what happens when this topological order is applied to an infinitely large flat sheet. We have not examined what happens when this topological order is put on a space with nontrivial topology. For instance, the torus. Suppose we analyse the system of \mathcal{C} applied to the torus. This amounts to breaking up the torus into some microscopic lattice and applying some Hamiltonian. This Hamiltonian will have group states $V_{\mathrm{g.s.}}^{T^2}$, which are independent of the choice of microscopic realization of \mathcal{C} .

Suppose we start with a torus, cut it across, twist one of its legs, then glue it back together, as shown below:

WORK: add Dehn twist picture.

If the initial torus has some state $|\psi\rangle \in V_{\mathrm{g.s.}}^{T^2}$ on it, then applying this procedure would give back another group state, though possibly a different one. The key phenomenon is that continuous transformations on physical space correspond to linear transformations on state space:

WORK: add schematic.

We can make this more formal as follows. We define the *mapping class group* of a topological space X as follows:

$$\mathrm{MCG}(X) = (\text{homeomorphisms } X \rightarrow X) / (\text{continuous deformations}).$$

If two homeomorphisms can be continuously deformed from one another then they will act the same on the ground states $V_{\mathrm{g.s.}}^{T^2}$. This is because ground states are topologically protected and hence slowly changing the diffeomorphism cannot affect them. Hence, we get a well-defined group homomorphism

$$\rho_{\mathcal{C}}^{T^2} : \mathrm{MCG}(T^2) \rightarrow \mathrm{Aut}\left(V_{\mathrm{g.s.}}^{T^2}\right).$$

This homomorphism connects back to our modular representation as follows:

- **Claim 1:** $\text{MCG}(T^2) \cong \text{SL}_2(\mathbb{Z})$;
- **Claim 2:** $V_{\text{g.s.}}^{T^2} \cong \mathbb{C}[\mathcal{L}]$;
- **Claim 3:** $\rho_{\mathcal{C}}^{T^2} \cong \rho_{\mathcal{C}}$, passing through the identifications in claims 1 and 2.

In general, we see that associated to every modular category \mathcal{C} there should not only be a modular representation, but also a representation of $\text{MCG}(\Sigma)$ for many other choices of topological space Σ . For instance, if $\Sigma = \Sigma_g$ is the g -holed torus then putting \mathcal{C} on Σ_g we get a map

$$\rho_{\mathcal{C}}^{\Sigma_g} : \text{MCG}(\Sigma_g) \rightarrow \text{Aut}(V_{\text{g.s.}}^{\Sigma_g}).$$

WORK: maybe say a few words about these representations. I'm sure they must have an explicit description in terms of generators and relations. A good reference (though a bit early) is this one: [Lyu95]

We now examine and motivate claims 1-3.

Claim 1: $\text{MCG}(T^2) \cong \text{SL}_2(\mathbb{Z})$. This claim is best seen by thinking of the torus as a gluing diagram,

WORK: add gluing diagram

WORK: add diagram with s acting by rotating by 90 degrees. Clearly, $s^4 = 1$.

WORK: add diagram with t as a shift. $(st)^3 = s^2$ can be left as an exercise.

WORK: writing presentation for $\text{MCG}(T^2)$, note that it is the same as $\text{SL}_2(\mathbb{Z})$.

Claim 2: $V_{\text{g.s.}}^{T^2} \cong \mathbb{C}[\mathcal{L}]$.

WORK: explain this. Cut into cylinder, label by charge on boundary

Claim 3: $\rho_{\mathcal{C}}^{T^2} \cong \rho_{\mathcal{C}}$.

WORK: Showing that the Dehn twist acts diagonally by θ_A is obvious. θ_A and Dehn twist are both defined as a 2π twist. For S we need another argument, more subtle but not too hard. I think Simon has it.

WORK: Finish by saying this is something like TQFTs. TQFT = bundled collection of mapping class group representations. Link this to TQFT appendix.

WORK: I wrote a little extra about this for the Kapustin final, might be useful

To demonstrate the TQFT perspective in action, we focus on the case of the genus one surface - the torus. As part of the data of every TQFT $(V_g, \rho_g, Z_{g,g'})$ there is a map

$$\rho_1 : \text{MCG}(\Sigma_1) \rightarrow \text{Aut}(V_1).$$

We seek to understand this map in full detail, and as such gain some level of enlightenment about the behavior of topological order on the torus. On physical groups, the space V_1 is easy to describe. We argued before it should be $\mathbb{C}[\mathcal{L}]$. The mapping class group $\text{MCG}(\Sigma_1) = \text{MCG}(T^2)$ is also straightforward to describe. We state a classical theorem about $\text{MCG}(T^2)$ below:

Theorem 1.85 ([FM11], Theorem 2.5). *There is an isomorphism $\text{MCG}(T^2) \cong \text{SL}_2(\mathbb{Z})$, which sends the element of $\text{MCG}(T^2)$ represented by $\pi/2$ -rotation to s ,*

WORK: add picture

and which sends the element of $\text{MCG}(T^2)$ represented by a Dehn twist around a handle of the torus to t .

Remark 1.86. For simplicity of notation, we will from now on identify $\text{MCG}(T^2)$ with $\text{SL}_2(\mathbb{Z})$, using the letters s, t to represent both the elements of $\text{SL}_2(\mathbb{Z})$ and their corresponding preimage in $\text{MCG}(T^2)$.

A more elementary way of stating this theorem is that every self-homeomorphism of the torus can be deformed to a composition of Dehn twists and rotations, and that the rotation and Dehn twist satisfy the $\text{SL}_2(\mathbb{Z})$ relations $s^4 = 1$ and $(st)^3 = s^2$.

The Dehn twist is described as follows. WORK: describe Dehn twist in terms of theta-symbols θ_A .

The rotation is described as follows. WORK: describe the rotation in terms of the S -matrix.

Now that we have described $\text{MCG}(T^2) = \text{SL}_2(\mathbb{Z})$ and $V_1 = \mathbb{C}[\mathcal{L}]$, we are tasked with understanding $\rho_1 : \text{SL}_2(\mathbb{Z}) \rightarrow \mathbb{C}[\mathcal{L}]$. That is, we must describe the maps $\rho_1(s), \rho_1(t) \in \text{Aut}(\mathbb{C}[\mathcal{L}])$.

]

1.5.3 Bruguières's modularity theorem and the Verlinde formula

In this section we prove Bruguières's modularity theorem. This theorem asserts that, given a pre-modular category \mathcal{C} , the S -matrix S is invertible if and only if \mathcal{C} is modular. Historically, this theorem is backwards. The original definition of modular category included that the S -matrix should be invertible. This was the only definition of modular category, until Bruguières proved in [ref] that the invertability of the S -matrix is equivalent to \mathcal{C} having the non-degenerate braiding property that if

$$\begin{array}{c} A & B \\ \swarrow & \searrow \\ A & B \end{array} = \begin{array}{c} A & B \\ | & | \\ A & B \end{array}$$

for all $B \in \mathcal{C}$ then $A \cong \mathbf{1}$. We are thus stating a historically incorrect definition of modular category, and Bruguières's modularity theorem tells us that this is equivalent to the original definition. The proof of the modularity theorem relies on the *Verlinde algebra* of \mathcal{C} . This algebra will be of use for us in proving other theorems in the future, in particular the Verlinde formula in section [ref].

We define an *algebra* over \mathbb{C} to a vector space V paired with a bilinear map $\cdot : V \times V \rightarrow V$ called multiplication, such that multiplication is associative and has a unit. An algebra is called *commutative* if its multiplication is commutative.

We define the Verlinde algebra $K_{\mathbb{C}}(\mathcal{C})$ of \mathcal{C} as follows:

$$K_{\mathbb{C}}(\mathcal{C}) = \left\{ \mathbb{C}[\mathcal{L}] \text{ with algebra structure } |[A]\rangle \cdot |[B]\rangle = \sum_{[C] \in \mathcal{L}} N_C^{A,B} |[C]\rangle \right\}.$$

We additionally define the function algebra

$$\mathbb{C}[\mathcal{L}]^{\text{func}} = \left\{ \mathbb{C}[\mathcal{L}] \text{ with algebra structure } \left(\sum_{[A] \in \mathcal{L}} c_A |[A]\rangle \right) \cdot \left(\sum_{[A] \in \mathcal{L}} c'_A |[A]\rangle \right) = \sum_{[A] \in \mathcal{L}} c_A c'_A |[A]\rangle \right\}.$$

Lemma 1.87. Both $K_{\mathbb{C}}(\mathcal{C})$ and $\mathbb{C}[\mathcal{L}]^{\text{func}}$ are commutative algebras.

Proof. The fact that $K_{\mathbb{C}}(\mathcal{C})$ is associative follows from the associativity of the tensor product. Its unit is $|[1]\rangle$. It is commutative because \mathcal{C} is braided. The fact that $\mathbb{C}[\mathcal{L}]^{\text{func}}$ is a commutative algebra is a standard exercise in algebra. Its unit is $\sum_{[A] \in \mathcal{L}} |[A]\rangle$. \square

We now state and prove the core theorem which underlies the core properties of the S matrix:

Proposition 1.88. *The map*

$$\begin{aligned}\mathcal{S} : K_{\mathbb{C}}(\mathcal{C}) &\rightarrow \mathbb{C}[\mathcal{L}]^{\text{func}} \\ |[A]\rangle &\mapsto \sum_{[B] \in \mathcal{L}} \frac{1}{d_B} S_{B,A} |[B]\rangle\end{aligned}$$

is a morphism of algebras.

Proof. Since it was defined on a basis, \mathcal{S} is clearly a linear map. We now verify that \mathcal{S} preserves multiplication. In the below computation, we identify endomorphisms of simple objects with the unique scalar they are times the identity. We let A, B, D be simple objects.

$$\begin{aligned}\left(\frac{1}{d_D} S_{D,A}\right) \left(\frac{1}{d_D} S_{D,B}\right) &= \begin{pmatrix} D \\ | \\ A \circlearrowleft \\ | \\ D \end{pmatrix} \cdot \begin{pmatrix} D \\ | \\ B \circlearrowleft \\ | \\ D \end{pmatrix} = \begin{pmatrix} D \\ | \\ B \circlearrowleft \\ | \\ A \circlearrowleft \\ | \\ D \end{pmatrix} \\ &= \begin{pmatrix} D \\ | \\ A \otimes B \circlearrowleft \\ | \\ D \end{pmatrix} = \sum_{[C] \in \mathcal{L}} N_C^{A,B} \cdot \begin{pmatrix} D \\ | \\ C \circlearrowleft \\ | \\ D \end{pmatrix} \\ &= \sum_{[C] \in \mathcal{L}} N_C^{A,B} \left(\frac{1}{d_D} S_{D,C} \right).\end{aligned}$$

Note our key use of the fact that

$$B \oplus C \circlearrowleft = B \circlearrowleft + C \circlearrowleft$$

which follows from the facts that $\text{id}_{B \oplus C}$ can be decomposed as projection onto B plus projection onto C , and composition is bilinear. We now conclude that

$$\begin{aligned}
\mathcal{S}(|[A]\rangle) \cdot \mathcal{S}(|[B]\rangle) &= \sum_{[D] \in \mathcal{L}} \left(\frac{1}{d_D} S_{D,A} \right) \left(\frac{1}{d_D} S_{D,B} \right) |[D]\rangle \\
&= \sum_{[D] \in \mathcal{L}} \left(\sum_{[C] \in \mathcal{L}} N_C^{A,B} \left(\frac{1}{d_D} S_{D,C} \right) \right) |[D]\rangle \\
&= \mathcal{S}(|[A]\rangle \cdot |[B]\rangle)
\end{aligned}$$

as desired. \square

By Proposition [ref], we have constructed a map of algebras $\mathcal{S} : K_{\mathbb{C}}(\mathcal{C}) \rightarrow \mathbb{C}[\mathcal{L}]^{\text{func}}$. As a map of vector spaces, \mathcal{S} is equal to the S -matrix up to a rescaling of rows by nonzero factors. Hence, it is clear that the S -matrix is invertible if and only if the algebra map \mathcal{S} is invertible. We now use special properties of the map \mathcal{S} to prove the main theorem of the section:

Theorem 1.89 (Bruguières's modularity theorem). *Let \mathcal{C} be a pre-modular category. The braiding on \mathcal{C} satisfies the non-degenerate braiding axiom if and only if the S -matrix is invertible.*

Proof. We observe that \mathcal{C} has a degenerate braiding if and only if there exists some $A \not\cong \mathbf{1}$ such that

$$\begin{array}{ccc}
D & & D \\
| & & | \\
A \bigcirc \text{---} & = d_A \cdot & | \\
| & & | \\
D & & D
\end{array}$$

for all $D \in \mathcal{D}$. If such an element A exists, then clearly $\mathcal{S}(|[A]\rangle) = d_A \mathcal{S}(|[1]\rangle)$. Hence, two linearly independent vectors map to linearly dependent vectors and thus \mathcal{S} is not invertible. Thus, the invertibility of the S -matrix implies that the braiding is non-degenerate.

We now prove the converse, and hence we suppose that \mathcal{C} has nondegenerate braiding. The proof is in two main steps. First, we prove that $|[1]\rangle$ is in the image of \mathcal{S} . Then, we use the fact that $|[1]\rangle$ is in the image of \mathcal{S} to construct the rest of the image, which proves that \mathcal{S} is surjective hence invertible.

Part 1: $|[1]\rangle$ is in the image of \mathcal{S} . Since \mathcal{C} has nondegenerate braiding, for all simple objects $A \not\cong \mathbf{1}$ there exists some simple object \tilde{A} such that

$$\begin{array}{ccc}
A & & A \\
| & & | \\
\tilde{A} \bigcirc \text{---} & \neq d_{\tilde{A}} \cdot & | \\
| & & | \\
A & & A
\end{array}$$

Thus, the vector $\mathcal{S}(|[\tilde{A}]\rangle) - \frac{S_{A,\tilde{A}}}{d_A} \mathcal{S}(|[1]\rangle)$ has a coefficient zero of for $|[A]\rangle$ but a non-zero coefficient for $|[1]\rangle$. Thus, using the product structure on $\mathbb{C}[\mathcal{L}]^{\text{func}}$, we find that the vector

$$\prod_{\substack{[A] \in \mathcal{L} \\ A \not\cong \mathbf{1}}} \left(\mathcal{S}(|[\tilde{A}]\rangle) - \frac{S_{A,\tilde{A}}}{d_A} \mathcal{S}(|[1]\rangle) \right)$$

has a coefficient of zero for all $|[A]\rangle$, $A \not\cong \mathbf{1}$, but a non-zero coefficient of $|[\mathbf{1}]\rangle$. Hence, it is a scalar multiple of $|[\mathbf{1}]\rangle$. Since \mathcal{S} is a morphism of algebras, it is in the image of \mathcal{S} . Hence, $|[\mathbf{1}]\rangle$ is in the image of \mathcal{S} .

This completes the first part of the proof. We now use the fact that $|[\mathbf{1}]\rangle$ is in the image of \mathcal{S} to construct the rest of the vectors.

Part 2: \mathcal{S} is surjective. Let $\omega = \sum_{[A] \in \mathcal{L}} \omega_A |[A]\rangle \in K_{\mathbb{C}}(\mathcal{C})$ be a vector such that $\mathcal{S}(\omega) = |[\mathbf{1}]\rangle$, which exists by part 1 of the proof. We now compute the quantity

$$h_{X,Y} = \sum_{[A] \in \mathcal{L}} \omega_A \cdot \text{tr} \left(A \begin{array}{c} X & Y \\ \vdash & \dashv \\ - & \end{array} \right)$$

two ways, for all simple objects $X, Y \in \mathcal{C}$. The first way follows by expanding $X \otimes Y$ as a direct sum and using the fact that $\mathcal{S}(\omega) = |[\mathbf{1}]\rangle$:

$$\begin{aligned} h_{X,Y} &= \sum_{[A] \in \mathcal{L}} \sum_{[B] \in \mathcal{L}} \omega_A N_B^{X,Y} d_B \cdot \begin{array}{c} B \\ \vdash \\ A \circlearrowleft \\ B \end{array} \\ &= \sum_{[B] \in \mathcal{L}} N_B^{X,Y} d_B \sum_{[A] \in \mathcal{L}} \omega_A \cdot \begin{array}{c} B \\ \vdash \\ A \circlearrowleft \\ B \end{array} \\ &= N_{\mathbf{1}}^{X,Y} d_1 = \begin{cases} 1, & X \cong Y^* \\ 0, & \text{otherwise.} \end{cases} \end{aligned}$$

In our second way of computing $h_{X,Y}$, we relate the string diagram trace to S -matrix values:

$$\begin{aligned} \text{tr} \left(A \begin{array}{c} X & Y \\ \vdash & \dashv \\ - & \end{array} \right) &= X^* \begin{array}{c} A \\ \text{---} \\ A \end{array} Y^* \\ &= X^* \begin{array}{c} A \\ \text{---} \\ A^* \end{array} Y^* \\ &= \frac{1}{d_A} S_{X,A} S_{Y,A} \end{aligned}$$

and thus, combining our two computations, we find

$$h_{X,Y} = \sum_{[A] \in \mathcal{L}} \frac{\omega_A}{d_A} S_{X,A} S_{Y,A} = \begin{cases} 1 & X \cong Y^* \\ 0 & \text{otherwise.} \end{cases}$$

We now define the vector

$$\omega^{(X)} = \sum_{[A] \in \mathcal{L}} \frac{\omega_A}{d_A} S_{A,X^*} |[A]\rangle$$

for all simple objects $X \in \mathcal{C}$. We compute

$$\begin{aligned} \mathcal{S}(\omega^{(X)}) &= \sum_{[A] \in \mathcal{L}} \frac{\omega_A}{d_A} S_{A,X^*} \left(\sum_{[Y] \in \mathcal{L}} \frac{1}{d_Y} S_{A,Y} |[Y]\rangle \right) \\ &= \sum_{[Y] \in \mathcal{L}} \frac{1}{d_Y} \left(\sum_{[A] \in \mathcal{L}} \frac{\omega_A}{d_A} S_{A,X^*} S_{A,Y} \right) |[Y]\rangle \\ &= \sum_{[Y] \in \mathcal{L}} \frac{h_{X^*,Y}}{d_Y} |[Y]\rangle = \frac{1}{d_X} |[X]\rangle. \end{aligned}$$

Hence $|[X]\rangle$ is in the image of \mathcal{S} for $[X] \in \mathcal{L}$, as desired. \square

1.5.4 Verlinde formula

In this section, we prove the *Verlinde formula*. This formula was first conjectured by Verlinde [Ver88], and proven the following year by Moore-Seiberg [MS89]. There are now many Verlinde-type formulas. Most importantly, there is one for vertex operator algebras [Hua08] and one in algebraic geometry [Fal94]. With proposition [ref] in hand, the proof is very quick:

Theorem 1.90 (Verlinde formula). *Let \mathcal{C} be a modular category.*

(i) *For all simple objects $A, B, C \in \mathcal{C}$,*

$$N_C^{A,B} = \sum_{[E] \in \mathcal{L}} \frac{S_{A,E} S_{B,E} (S^{-1})_{C,E}}{d_E}$$

where $(S^{-1})_{C,E}$ denotes the (C, E) -coefficient of the inverse of the S matrix.

(ii) *For all simple objects $A \in \mathcal{C}$, the matrix*

$$D^A = S N^A S^{-1}$$

is diagonal with $([B], [B])$ -entry $S_{A,B}/d_B$, where $N^A = (N_C^{A,B})_{([B], [C]) \in \mathcal{L}^2}$ is the fusion matrix of A .

Proof. We begin by proving part (ii). The main observation of the proof is that the operator $N^A : \mathbb{C}[\mathcal{L}] \rightarrow \mathbb{C}[\mathcal{L}]$ is exactly left multiplication by A in $K_{\mathbb{C}}(\mathcal{C})$. Proposition [ref] says that \mathcal{S} is a morphism of algebras, and hence we can commute N^A past \mathcal{S} , and turn it into multiplication by A in $\mathbb{C}[\mathcal{L}]^{\text{func}}$. Hence, using the appropriate multiplication in the appropriate algebra, we find

$$\begin{aligned}
(\mathcal{S}N^A\mathcal{S}^{-1})|[B]\rangle &= \mathcal{S}(|[A]\rangle \cdot_{K_{\mathcal{C}}(\mathcal{C})} \mathcal{S}^{-1}(|[B]\rangle)) \\
&= \mathcal{S}(|A\rangle) \cdot_{\mathbb{C}[\mathcal{L}]^{\text{func}}} |[B]\rangle \\
&= \frac{S_{A,B}}{d_B} |[B]\rangle.
\end{aligned}$$

Thus, $\mathcal{S}N^A\mathcal{S}^{-1}$ is diagonal with $([B], [B])$ entry $S_{A,B}/d_B$. Scaling rows of \mathcal{S} does not change the effect of diagonalization. Hence, we conclude that SN^AS^{-1} is diagonal as well, with the same entries, and thus our proof of (ii) is complete.

We now move on to proving part (i). Expanding the formula $N^A = S^{-1}D^AS$, we find

$$\begin{aligned}
N^A|[B]\rangle &= S^{-1}D^AS|[B]\rangle \\
&= S^{-1} \left(\sum_{[E] \in \mathcal{L}} \frac{S_{A,E}S_{B,E}}{d_E} |[E]\rangle \right) \\
&= \sum_{[E] \in \mathcal{L}} \frac{S_{A,E}S_{B,E}}{d_E} \left(\sum_{[C] \in \mathcal{L}} (S^{-1})_{C,E} |[C]\rangle \right).
\end{aligned}$$

Comparing coefficients with the definition of N^A , we conclude the result. \square

1.5.5 Proof of modularity

In this section we prove that the S -matrix and T -matrix indeed give a representation of the modular group. That is, we will prove Theorem [ref]. At its heart, the fact that the modular representation of modular category is a homomorphisms comes down to proving a series of relations between the coefficients of the S -matrix and the coefficients of the T -matrix. That is, we are proving a series of relations between braiding and twisting. The general method is to take traces of certain diagrams, and then compute those traces in two ways. One way will involve more twists and the other will involve more braiding. This will give some algebraic relation, and choosing the right diagrams we will get enough algebraic relations to deduce Theorem [ref].

We begin with the most fundamental relationship between S -matrix and T -matrix entries:

Lemma 1.91. *Let \mathcal{C} be a pre-modular category. We have that*

$$S_{A,B} = \theta_A^{-1}\theta_B^{-1} \sum_{[C] \in \mathcal{L}} N_C^{A,B} \theta_C d_C$$

and

$$S_{A^*,B} = \theta_A \theta_B \sum_{[C] \in \mathcal{L}} N_C^{A,B} \theta_C^{-1} d_C.$$

Proof. By Proposition [ref], we have $\beta_{B,A} \circ \beta_{A,B} = (\theta_A^{-1} \otimes \theta_B^{-1}) \circ \theta_{A \otimes B}$. Taking the trace of this formula we get

$$S_{A,B} = \theta_A^{-1} \theta_B^{-1} \operatorname{tr}(\theta_{A \otimes B}).$$

Now, since θ is a natural transformation it splits over direct sums. Traces split over direct sums as well by proposition [ref] and hence $\operatorname{tr}(\theta_{A \otimes B}) = \sum_{[C] \in \mathcal{L}} N_C^{A,B} \theta_C d_C$. This concludes the proof of the first formula.

For the second formula, we observe that replacing overcrossings with undercrossings in the definition of $S_{A,B}$ has the effect of taking the dual of one of elements, replacing it with $S_{A^*,B}$. Hence $\operatorname{tr}(\beta_{A,B}^{-1} \circ \beta_{B,A}^{-1}) = S_{A^*,B}$. Thus, taking the trace of the formula $\beta_{A,B}^{-1} \circ \beta_{B,A}^{-1} = \theta_{A \otimes B}^{-1} \circ (\theta_A \otimes \theta_B)$ yields the desired result. \square

Before continuing to our proof of theorem [ref] we observe a key lemma:

Lemma 1.92. *Let \mathcal{C} be a pre-modular category. Let $A \in \mathcal{C}$ be a (possibly non-simple) object. We have that*

$$\sum_{[B] \in \mathcal{L}} d_B \theta_B \cdot \left(\begin{array}{c} A \\ | \\ B \text{ } \bigcirc \\ | \\ A \end{array} \right) = p_C^+ \cdot \boxed{\theta_A^{-1}} \quad \begin{array}{c} A \\ | \\ \boxed{\theta_A^{-1}} \\ | \\ A \end{array}$$

and

$$\sum_{[B] \in \mathcal{L}} d_B \theta_B^{-1} \cdot \left(\begin{array}{c} A \\ | \\ B \text{ } \bigcirc \\ | \\ A \end{array} \right) = p_C^- \cdot \boxed{\theta_A} \quad \begin{array}{c} A \\ | \\ \boxed{\theta_A} \\ | \\ A \end{array}$$

Proof. We only prove the first formula - the second follows by a formally dual argument. We restrict to the case that A is simple. Seeing as both sides are linear with respect to direct sums, the case that A is simple will immediately imply the general case. Since A is simple, it suffices to prove that the traces of both sides are equal. The trace on the left hand side has the effect of replacing the diagram with $S_{A,B}$. Hence, we compute as follows using Lemma [ref] and the fact that $\sum_{[C] \in \mathcal{L}} N_C^{A,B} d_B = d_A d_C$ from proposition [ref]:

$$\begin{aligned} \sum_{[B] \in \mathcal{L}} d_B \theta_B S_{A,B} &= \sum_{[B] \in \mathcal{L}} d_B \theta_B \left(\theta_A^{-1} \theta_B^{-1} \sum_{[C] \in \mathcal{L}} N_C^{A,B} \theta_C d_C \right) \\ &= \theta_A^{-1} \sum_{[C] \in \mathcal{L}} \theta_C d_C \left(\sum_{[B] \in \mathcal{L}} N_C^{A,B} d_B \right) \\ &= \theta_A^{-1} d_A \sum_{[C] \in \mathcal{L}} \theta_C d_C^2 = p_C^+ \theta_A^{-1} d_A. \end{aligned}$$

This result is exact the trace of the right hand sides of the lemma. Hence, the proof is complete. \square

We now give the heart of the proof of theorem [ref]:

Theorem 1.93. *Let \mathcal{C} be a pre-modular category. Define the charge conjugation operator $\check{C} : \mathbb{C}[\mathcal{L}] \rightarrow \mathbb{C}[\mathcal{L}]$ to the matrix with $([A], [A^*])$ coefficient 1 for all $[A] \in \mathcal{L}$, and all other coefficients zero.*

$$(i) \check{C}S = S\check{C}, \check{C}T = T\check{C}, \text{ and } \check{C}^2 = 1;$$

$$(ii) (ST)^3 = p_{\mathcal{C}}^+ S^2;$$

$$(iii) (ST^{-1})^3 = p_{\mathcal{C}}^- S^2 \check{C}.$$

If \mathcal{C} is modular, then

$$(iv) S^2 = p_{\mathcal{C}}^+ p_{\mathcal{C}}^- \check{C}.$$

Proof. Part (i) follows from the fact that $S_{A^*, B^*} = S_{A, B}$, $\theta_{A^*} = \theta_A$, and $A^{**} \cong A$. Parts (ii) and (iii) have formally dual proofs, which arise from replacing θ with θ^{-1} at every opportunity. Part (iv) follows algebraically from combining formulas (ii) and (iii) whenever S is invertible, which is always the case when \mathcal{C} is modular by theorem [ref].

Hence, it suffices to prove part (ii). The proof comes from computing the quantity

$$h_{X,Y} = \sum_{[A] \in \mathcal{L}} d_A \theta_A \cdot \text{tr} \left(A \begin{array}{c} X & Y \\ \diagup & \diagdown \\ - & - \\ \diagdown & \diagup \\ X & Y \end{array} \right)$$

two ways.

In the first way of computing $h_{X,Y}$, we use lemma [ref]. We find the following:

$$h_{X,Y} = p_{\mathcal{C}}^+ \cdot \text{tr} \left(\begin{array}{c} X & Y \\ \diagup & \diagdown \\ \theta_{X \otimes Y}^{-1} \\ \diagdown & \diagup \\ X & Y \end{array} \right) = p_{\mathcal{C}}^+ \theta_X \theta_Y S_{X^*, Y}$$

In our second way of computing $h_{X,Y}$, we use computation of the trace of two lines through a loop in the proof of Theorem [ref]. We find this way that

$$h_{X,Y} = \sum_{[A] \in \mathcal{L}} \theta_A S_{X,A} S_{Y,A}.$$

Thus, we find that $\sum_{[A] \in \mathcal{L}} \theta_A S_{X,A} S_{Y,A} = p_{\mathcal{C}}^+ \theta_X \theta_Y S_{X^*, Y}$. Thinking of these quantities as the $([X], [Y])$ entries in operators $\mathbb{C}[\mathcal{L}] \rightarrow \mathbb{C}[\mathcal{L}]$, we get the equation

$$STS = p_{\mathcal{C}}^+ T S T \check{C}.$$

WORK: This formula is WRONG. It should be

$$STS = p_{\mathcal{C}}^+ T^{-1} S T^{-1}.$$

From this we get

$$(ST)^3 = p_{\mathcal{C}}^+ S^2$$

as desired. I'm not sure where I went wrong, but something is off in here. \square

We now know that the S and T matrices give a modular representation *up to phase!* We still need to work out the details of the phases. A key part is the following computation:

Corollary 1.94. *Let \mathcal{C} be a modular category. The quantities $p_{\mathcal{C}}^+$ and $p_{\mathcal{C}}^-$ are nonzero, and*

$$p_{\mathcal{C}}^+ p_{\mathcal{C}}^- = \mathcal{D}^2.$$

Proof. The values $p_{\mathcal{C}}^+$ and $p_{\mathcal{C}}^-$ must be nonzero because S is invertible and $S^2 = p_{\mathcal{C}}^+ p_{\mathcal{C}}^- \check{C}$. The formula $S^2 = p_{\mathcal{C}}^+ p_{\mathcal{C}}^- \check{C}$, when expanded, says that

$$\sum_{[C] \in \mathcal{L}} S_{C,A} S_{C,B} = \begin{cases} p_{\mathcal{C}}^+ p_{\mathcal{C}}^- & A \cong B^* \\ 0 & \text{otherwise.} \end{cases}$$

Applying this formula to $A = B = \mathbf{1}$, we find

$$\sum_{[C] \in \mathcal{L}} S_{C,\mathbf{1}} S_{C,\mathbf{1}} = \sum_{[C] \in \mathcal{L}} d_C^2 = p_{\mathcal{C}}^+ p_{\mathcal{C}}^-$$

as desired. \square

Now, it is clear that we can normalize the representation appropriately and conclude theorem [ref], as desired.

1.5.6 Vafa's theorem, unitarity of S -matrix, and the Chiral central charge

In this section we discuss some finer points of the structure of the modular representation. In particular, we will prove that modular representation of every modular category is *unitary*. That is, the S and T matrices are both unitary operators on $\mathbb{C}[\mathcal{L}]$ when it is endowed with its canonical inner product coming from its basis

We begin with the matrix T . For a diagonal matrix to be unitary, it is necessary and sufficient for its diagonal entries to have absolute value 1. We will prove something even stronger: that all of the entries are roots of unity! We recall that a number $z \in \mathcal{C}$ is called a root of unity if $z^n = 1$ for some integer $n \geq 1$. We begin with a key topological lemma which will underscore our proof:

Lemma 1.95 (Lantern identity). *Let \mathcal{C} be a pre-modular category. Let $A, B, C \in \mathcal{C}$ be objects. As maps $A \otimes B \otimes C \rightarrow A \otimes B \otimes C$, we have the identity*

$$\theta_{A \otimes B} \circ \theta_{A \otimes C} \circ \theta_{B \otimes C} = \theta_{A \otimes B \otimes C} \circ (\theta_A \otimes \theta_B \otimes \theta_C).$$

Proof. In the language of string diagrams, this formula becomes

WORK: add diagram.

It is a matter of elementary manipulations to convince one's self that these two diagrams are equal. An alternate algebraic approach is to expand the relation both sides using the formula $\theta_{X \otimes Y} = (\beta_{Y,X} \circ \beta_{X,Y}) \circ (\theta_X \otimes \theta_Y)$, cancel $\theta_A^2 \otimes \theta_B^2 \otimes \theta_C^2$, and compare the resulting braids using the hexagon and naturality. \square

We state one more linear-algebraic lemma necessary for the proof:

Lemma 1.96. *Let $n \geq 1$ be an integer and let $M : \mathbb{R}^n \rightarrow \mathbb{R}^n$ be a linear operator whose (a,b) entry is $M_{a,b}$. Suppose that all off diagonal entries of M are positive, and that $\sum_{b=1}^n M_{a,b} < 0$ for all $1 \leq a \leq n$. Then, M is invertible.*

Proof. WORK: This follows from the Gershgorin circle theorem, but that's way to general for our purposes. Find a nice proof which is appropriate for the situation. \square

Theorem 1.97 (Vafa). *Let \mathcal{C} be a pre-modular category. The values θ_A are roots of unity for all simple object $A \in \mathcal{C}$.*

Remark 1.98. This theorem is interesting in part due to its history. WORK: add history.

Proof. To begin, we choose a simple object A . We observe that every endomorphism of $A \otimes A^* \otimes A$ induces a linear map

$$\text{Hom}(A, A \otimes A^* \otimes A) \rightarrow \text{Hom}(A, A \otimes A^* \otimes A)$$

by postcomposition. We now consider the lantern identity (proposition [ref]) on the stands A, A^*, A ,

$$\theta_{A \otimes A^*} \circ \theta_{A \otimes A} \circ \theta_{A^* \otimes A} = \theta_{A \otimes A^* \otimes A} \circ (\theta_A \otimes \theta_{A^*} \otimes \theta_A).$$

viewed as an equality of linear operators on $\text{Hom}(A, A \otimes A^* \otimes A)$. We compute the determinant of both sides. We first compute the determinant of $\theta_{A \otimes A^*}$. The eigenvalues of the operator $\theta_{A \otimes A^*}$ are the twists θ_B . The dimension of the θ_B eigenspace is exactly the dimension of the subspace of $\text{Hom}(A, A \otimes A^* \otimes A)$ in which $A \otimes A^*$ fuse to B . The dimension of this space is $N_B^{A,A^*} N_A^{B,A} = (N_B^{A,A^*})^2$. The determinant is equal to the product of the eigenvalues counted with multiplicity, and hence

$$\det \theta_{A \otimes A^*} = \prod_{[B] \in \mathcal{L}} \theta_B^{(N_B^{A,A^*})^2}.$$

Continuing this way for all the other twists and plugging them into the lantern identity we get

$$\prod_B \theta_B^{2(N_B^{A,A^*})^2 + (N_B^{A,A})^2} = \theta_A^{4 \cdot \dim \text{Hom}(A, A \otimes A^* \otimes A)}.$$

We now define the coefficients

$$M_{A,B} = \begin{cases} 2(N_B^{A,A^*})^2 + (N_B^{A,A})^2 & A \not\cong B \\ 2(N_B^{A,A^*})^2 + (N_B^{A,A})^2 - 4 \cdot \dim \text{Hom}(A, A \otimes A^* \otimes A) & A \cong B. \end{cases}$$

We observe that all of the off-diagonal entries of $M_{A,B}$ are positive and

$$\begin{aligned} \sum_{[B] \in \mathcal{L}} M_{A,B} &= \sum_{[B] \in \mathcal{L}} 2(N_B^{A,A^*})^2 + (N_B^{A,A})^2 - 4 \cdot \dim \text{Hom}(A, A \otimes A^* \otimes A) \\ &= - \dim \text{Hom}(A, A \otimes A^* \otimes A) < 0. \end{aligned}$$

Thus, by lemma [ref], we conclude that the matrix $M = (M_{A,B})_{([A],[B]) \in \mathcal{L}^2} : \mathbb{C}[\mathcal{L}] \rightarrow \mathbb{C}[\mathcal{L}]$ is invertible. Let \tilde{M} be the adjugate of M . That is, an integer valued matrix such that $M \cdot \tilde{M} = \tilde{M} \cdot M = n$ where $n = \det M$. We find for all simple objects A that

$$\theta_A^n = \prod_{[C] \in \mathcal{L}} \left(\prod_{[B] \in \mathcal{L}} \theta_B^{M_{C,B}} \right)^{\tilde{M}_{A,C}} = 1,$$

and hence θ_A is a root of unity as desired. \square

We immediately get several corollaries from this theorem. The first is that braiding enough times gets you back where you started:

Corollary 1.99. *Let \mathcal{C} be a pre-modular category. There exists an integer $n \geq 1$ such that*

$$(\beta_{B,A} \circ \beta_{A,B})^n = \text{id}_{A \otimes B}$$

for all $A, B \in \mathcal{C}$

Proof. Choose $n \geq 1$ so that $\theta_C^n = 1$ for all simple objects $C \in \mathcal{C}$, which exists by Vafa's theorem (theorem [ref]). Seeing that $\beta_{B,A} \circ \beta_{A,B} = \theta_{A \otimes B} \circ (\theta_A^{-1} \otimes \theta_B^{-1})$, in the direct sum decomposition $A \otimes B \cong \bigoplus_{[C] \in \mathcal{L}} N_C^{A,B} |[C]\rangle$ the transformation $\beta_{B,A} \circ \beta_{A,B}$ acts by the scalar $\theta_C / (\theta_A \theta_B)$ on every $|[C]\rangle$ summand. Thus, $(\beta_{B,A} \circ \beta_{A,B})^n$ acts by $(\theta_C / (\theta_A \theta_B))^n = 1$ and thus $(\beta_{B,A} \circ \beta_{A,B})^n$ is the identity as desired. \square

Corollary 1.100. *Let \mathcal{C} be a modular category. The quantity p_C^- / p_C^+ is a root of unity.*

Proof. We take determinants. From $S^2 = p_C^+ p_C^- \check{C}$ we find that $\det(S)^2 = \pm p_C^+ p_C^-$, since $\det \check{C} = \pm 1$. From the formula $(ST)^3 = p_C^+ S^2$ we find that

$$(p_C^+ / \det(S))^2 = \det(T)^6$$

so $p_C^+ / p_C^- = \pm \det(T)^6$. By Vafa's theorem $\det(T)$ is a root of unity. Hence, we conclude that p_C^+ / p_C^- is a root of unity as desired. \square

We take a moment to observe that the quantity p_C^- / p_C^+ is of great interest to physicists. It is related to the *chiral central charge* c_{top} of the theory by the formula

$$p_C^- / p_C^+ = e^{-\frac{2\pi i c_{\text{top}}}{4}}.$$

WORK: probably say a bit more? Maybe there will be a better spot somewhere else.

We now move on to proving that the S matrix is unitary. The main technical ingredient is as follows:

Proposition 1.101. *Let \mathcal{C} be a modular category. For all simple objects $A, B \in \mathcal{C}$ we have that $S_{A^*, B} = \overline{S_{A, B}}$.*

Proof. By the Verlinde formula [ref], we know that for all simple objects A there exists a vector $\mathbf{v}_B \in \mathbb{C}[\mathcal{L}]$ such that

$$N^A \mathbf{v}_B = \frac{S_{A,B}}{d_B} \mathbf{v}_B$$

for all simple objects A . Namely, \mathbf{v}_B is the $[B]$ -column of the S matrix. Let \mathbf{v}_B^* be the row vector which is the Hermitian adjoint to \mathbf{v}_B . We have that

$$\mathbf{v}_B N^A \mathbf{v}_B = \frac{S_{A,B}}{d_B} |\mathbf{v}_B|^2.$$

Now, we observe that by Frobenius reciprocity (proposition [ref]) $(N^A)^\dagger = N^{A^*}$. Hence,

$$\begin{aligned} \mathbf{v}_B N^A \mathbf{v}_B &= \left(N^{A^*} \mathbf{v}_B \right)^* \mathbf{v}_B \\ &= \left(\frac{S_{A^*,B}}{d_B} \mathbf{v}_B \right)^* \mathbf{v}_B \\ &= \frac{\overline{S_{A^*,B}}}{d_B} |\mathbf{v}_B|^2. \end{aligned}$$

Comparing, we get the desired result. \square

We now get the following theorem:

Theorem 1.102 (Etingof-Nikshych-Ostrik). *Every matrix in the image of the modular representation $\mathrm{SL}_2(\mathbb{Z}) \rightarrow \mathrm{Aut}(\mathbb{C}[\mathcal{L}])$ is a unitary operator on $\mathbb{C}[\mathcal{L}]$.*

Proof. We already know by theorem [ref] and corollary [ref] that the normalized T -matrix is unitary. It thus suffices to show that the modular S matrix is unitary. From theorem [ref] we have that $(\frac{1}{D} S)^{-1} = (\frac{1}{D} S) \cdot \check{C}$. By proposition [ref] $(\frac{1}{D} S) \cdot \check{C} = (\frac{1}{D} S)^\dagger$ and thus $\frac{1}{D} S$ is unitary as desired. \square

1.6 Skeletonization

1.6.1 Principle

WORK: lots of big choices need to be made here. Do I call this the skeletonization, or do I call it something else? Do I work with multiplicity-free categories, or do I allow multiplicity? I don't know what the correct statements are or what the proofs look like so this section might be a tough one. A good mathematical reference is [DHW13].

1.6.2 F -symbols

1.6.3 R -symbols

1.6.4 θ -symbols

1.6.5 Reconstruction theorem

1.7 Quantum double modular categories

1.7.1 The Drinfeld center

A quantum double modular category is a special type of modular category. They are particularly important because many of the constructions of topological order only deal with quantum double modular categories. For instance, there are constructions of modular categories/topological order coming from the theory of tensor networks [ref], subfactors [ref], vertex operator algebras [ref], WORK: add more sources. All of these constructions only

give quantum double modular categories. Hence, understanding quantum doubles is key to understanding how topological order work in practice.

At the heart of quantum doubles is a construction known as the *Drinfeld center*. In its most basic form the Drinfeld center induces an assignment

$$\mathcal{Z} : (\text{monoidal categories}) \rightarrow (\text{braided monoidal categories}).$$

In our context, we care about a more structured version of the Drinfeld center. It is a theorem of Muger that the Drinfeld center induces an assignment as follows:

$$\mathcal{Z} : (\text{spherical fusion categories}) \rightarrow (\text{modular categories}).$$

This theorem is fantastic because it allows one to construct modular categories using much less data than would otherwise be necessary. Without needing a braiding, non-degenerate or otherwise, the Drinfeld center allows one to construct a modular category. This makes the Drinfeld center an abundant source of modular categories. We call a modular category \mathcal{C} a quantum double if it is of the form $\mathcal{Z}(\mathcal{C}_0)$ for some spherical fusion category \mathcal{C}_0 . A major goal of this chapter is to set up and prove Muger's theorem.

We now define the Drinfeld center. The Drinfeld center is a somewhat direct categorification of the usual notion of center for finite groups. If G is a finite group, its center is defined as follows:

$$Z(G) = \{g \in G \mid gh = hg \ \forall h \in G\}.$$

The first guess at $Z(\mathcal{C})$ is thus

$$\mathcal{Z}(\mathcal{C}) = \{A \in \mathcal{C} \mid A \otimes B \cong B \otimes A \ \forall B \in \mathcal{C}\}.$$

This is almost correct, but not quite. The issue is that $\mathcal{Z}(\mathcal{C})$ is not quite a braided monoidal category yet. Even though $A \otimes B \cong B \otimes A$ for all $A, B \in \mathcal{Z}(\mathcal{C})$, we don't have a distinguished choice of isomorphism. A braided monoidal category requires a distinguished isomorphism $\beta_{A,B} : A \otimes B \xrightarrow{\sim} B \otimes A$. Moreover, these distinguished isomorphisms are required to satisfy the hexagon equations. Hence, we make a new definition of center which keeps track of the choice of isomorphism and enforces the hexagon equation along the way:

Proposition 1.103. *The Drinfeld center $\mathcal{Z}(\mathcal{C})$ of a monoidal category \mathcal{C} is a braided monoidal category defined as follows:*

- (Objects) Pairs $(A, \beta_{A,-})$, where $A \in \mathcal{C}$, and $\beta_{A,-}$ is a natural isomorphism of monoidal natural isomorphism between the two functors $A \otimes -$ and $- \otimes A$ from \mathcal{C} to \mathcal{C} , satisfying the additional condition that

$$\beta_{A,B \otimes C} = (\text{id}_B \otimes \beta_{A,C}) \circ (\beta_{A,B} \otimes \text{id}_C).$$

- (Morphisms) Given $(A, \beta_{A,-}), (B, \beta_{B,-}) \in \mathcal{Z}(\mathcal{C})$, $\text{Hom}_{\mathcal{Z}(\mathcal{C})}((A, \beta_{A,-}), (B, \beta_{B,-}))$ is the subspace of morphisms $f \in \text{Hom}_{\mathcal{C}}(A, B)$ such that for all $C \in \mathcal{C}$

$$(\text{id}_C \otimes f) \circ \beta_{A,C} = \beta_{B,C} \circ (f \otimes \text{id}_C).$$

- (*Tensor product*) Given $(A, \beta_{A,-}), (B, \beta_{B,-}) \in \mathcal{Z}(\mathcal{C})$, we define

$$(A, \beta_{A,-}) \otimes (B, \beta_{B,-}) = (A \otimes B, (\beta_{A,-} \otimes \text{id}_C) \circ (\text{id}_C \otimes \beta_{B,-})).$$

- (*Unit*) The element $(1, \rho \circ \lambda^{-1})$
- (*Braiding*) We define the braiding between two elements $(A, \beta_{A,-}), (B, \beta_{B,-}) \in \mathcal{Z}(\mathcal{C})$ to be $\beta_{A,B} = \beta_{A,B}$.

Inheriting associativity, unitors, and composition from \mathcal{C} , this gives $\mathcal{Z}(\mathcal{C})$ the structure of a braided monoidal category.

Proof. Since morphisms in $\mathcal{Z}(\mathcal{C})$ are a subspace of morphisms in \mathcal{C} , commutative diagrams don't change when going from \mathcal{C} to $\mathcal{Z}(\mathcal{C})$. Hence, the triangle and pentagon axioms for $\mathcal{Z}(\mathcal{C})$ follow immediately from the triangle and pentagon axioms on \mathcal{C} . One thing to be checked is that evaluation/co-evaluation satisfy the compatibility condition required to a morphism in $\mathcal{Z}(\mathcal{C})$, but this is straightforward. We remark on the hexagon identities. The condition imposed on $\beta_{A,B \otimes C}$ given is technically incorrect. To make the parentheses work in the braiding one has to add associators, and impose the longer condition

$$\beta_{A,B \otimes C} = \alpha_{C,A,B}^{-1} \circ (\text{id}_B \otimes \beta_{A,C}) \circ \alpha_{A,C,B} \circ (\beta_{A,B} \otimes \text{id}_C) \circ \alpha_{A,B,C}^{-1}.$$

This condition makes the second hexagon identity tautological. Similarly, the definition of tensor product given is not strictly correct - one must add the correct associator terms, making the first hexagon identity immediate. Lastly one must verify the half-braidings defined on the tensor unit/tensor product are actually half braidings, i.e., that they satisfy the hexagon condition. These follow from straightforward computations, which we leave as exercises. This completes the proof. \square

1.7.2 Muger's theorem

WORK: through Muger's theorem. The exposition will greatly differ based on what the proof looks like, which I haven't done before. The standard proof people use is completely steeped in the language of module categories. The original proof uses the tube algebra, and is "elementary" in the sense that you don't need tube algebras but it is heinous, and I would very much like to avoid it. I'll figure out what to write hence once I've done the module category section and I've digested the proof. Maybe there's a way to de-module-ify it, but maybe not. Maybe I don't want to do that if it'll lose its essence.

1.7.3 Discrete gauge theory as a quantum double and Morita equivalence

We saw in chapter [ref] that \mathbf{Vec}_G and $\text{Rep}(G)$ are both naturally spherical fusion categories. Thus, Muger's theorem tells us that $\mathcal{Z}(\mathbf{Vec}_G)$ and $\mathcal{Z}(\text{Rep}(G))$ are both modular categories. Hence, given a finite group G we have three different modular categories we can associate to it: $\mathfrak{D}(G)$, $\mathcal{Z}(\mathbf{Vec}_G)$, $\mathcal{Z}(\text{Rep}(G))$. The amazing fact is that these are all the same category:

Proposition 1.104. *Let G be a finite group. There are equivalences of modular categories $\mathfrak{D}(G) \cong \mathcal{Z}(\mathbf{Vec}_G) \cong \mathcal{Z}(\text{Rep}(G))$.*

Proof. WORK: do proof. Shouldn't be too hard. \square

We now make a few comments about this theorem. The first is that it proves that $\mathfrak{D}(G)$ is a quantum double modular category. Secondly, it gives a second proof that $\mathfrak{D}(G)$ has a non-degenerate braiding, using Muger's theorem. Thirdly, it demonstrates the concept of *Morita equivalence*.

WORK: introduce Morita equivalence. I know that there's some important basic facts to tell - I should include those.

1.7.4 Factorizability and time reversal symmetry

Given a modular category \mathcal{C} , we can forget the braiding on \mathcal{C} and only remember its structure as a spherical fusion category. Hence, Muger's theorem tells us that $\mathcal{Z}(\mathcal{C})$ is canonically a modular category. It is a fantastic fact that in this case $\mathcal{Z}(\mathcal{C})$ can be explicitly computed in terms of \mathcal{C} . We describe this computation now.

WORK: Define the time-reversal conjugate $\bar{\mathcal{C}}$. Setup the map $\mathcal{C} \boxtimes \bar{\mathcal{C}} \rightarrow \mathcal{Z}(\mathcal{C})$

Proposition 1.105. *Let \mathcal{C} be a pre-modular category. The canonical map $\mathcal{C} \boxtimes \bar{\mathcal{C}} \rightarrow \mathcal{Z}(\mathcal{C})$ is an equivalence of categories if and only if \mathcal{C} is modular.*

Proof. WORK: proof \square

This theorem is fantastic because it not only computes $\mathcal{Z}(\mathcal{C})$ for every modular category \mathcal{C} , but also it gives an equivalent definition of modularity. This gives us our third definition of modularity. Namely a pre-modular category \mathcal{C} is modular if and only if its braidings are all non-degenerate, or equivalently if its S -matrix is non-degenerate, or equivalently if it is factorizable in the above sense.

1.7.5 Levin-Wen model

WORK: work though the Levin-Wen model. I think that this model is fantastic because it shows how all of the ideas of tensor category theory can manifest themselves extremely concretely on the level of gapped Hamiltonians. Namely, the coherence relations on the category theory side correspond exactly to the formulas needed to make terms in a Hamiltonian commute with one another. It would be nice if I could give a motivation for which the category which describes the Levin-Wen model is the Drinfeld center, though I've never seen that before.

1.8 Unitarity

1.8.1 Characterization of unitarizability

WORK: An early reference about unitary MTCs is [KJ96]. I should read this.

WORK: Include this somewhere:

Proposition 1.106. *Let \mathcal{C} be a unitary fusion category. The maps $\text{ev}_A^L = (\text{coev}_A)^\dagger$ and $\text{coev}_A^L = (\text{ev}_A)^\dagger$ give a left-rigid structure on \mathcal{C} . This left-rigidity endows \mathcal{C} with the structure of a spherical fusion category.*

As we've mentioned before, the algebraic theory of topological order is not modular category theory but *unitary* modular category theory because hom-spaces need Hilbert space

structure to define valid quantum systems. We define unitary fusion categories and unitary modular categories in section [ref].

We've mostly ignored unitary fusion categories for the following reasons: their additional structure is a large semantic burden, and is in a way inessential. We mean this in the following sense:

1. Given a spherical fusion category, \mathcal{C} every unitary structure on \mathcal{C} is equivalent. That is, all the unitary fusion categories obtained by enriching \mathcal{C} with unitary structure are equivalent to one another as unitary fusion categories;
2. All braidings on unitary fusion categories are automatically unitary so there is no subtlety in passing to unitary modular categories;
3. A spherical fusion category admits a unitary structure if and only if its quantum dimensions are all positive.

Still, it is worthwhile to study unitary fusion categories. They are the correct algebraic structure to describe topological order, so understanding *why* points (1-3) above are true gives insight into unitarity in practice. There is also a characterization of unitary fusion categories in the skeleized perspective which is of much utility in practice. Namely, a fusion category admits a unitary structure if and only if its F -matrices and R -matrices can all be made unitary. We discuss this more and offer a proof in section [ref].

We now prove the first main result of the section:

Proposition 1.107. *Let \mathcal{C} be a spherical fusion category. There is a compatible unitary structure on \mathcal{C} if and only if the quantum dimensions d_A are positive for all simple objects $A \in \mathcal{C}$.*

Proof. Suppose first that \mathcal{C} admits a unitary structure. Then, for all simple objects $A \in \mathcal{C}$,

$$d_A = \text{tr}(\text{id}_A) = \langle \text{id}_A | \text{id}_A \rangle > 0$$

because the inner product is positive definite. Conversely, suppose that all of the quantum dimensions of \mathcal{C} are positive. We will suppose for simplicity that \mathcal{C} is skeletal, which is possible by proposition [ref]. Suppose that $f : A \rightarrow B$ is a morphism between objects $A, B \in \mathcal{C}$. We know that there are direct sum decompositions $A = \bigoplus_i A_i$ and $B = \bigoplus_j B_j$. Writing f as an element of $\bigoplus_{i,j} \text{Hom}(A_i, B_j)$, we find that f acts by some scalar $f_{i,j}$ on each homspace $\text{Hom}(A_i, B_j)$, where $f_{i,j} = 0$ if $A \not\cong B$. We define $f^\dagger : B \rightarrow A$ to be the map whose $\text{Hom}(B_j, A_i)$ is $f_{j,i}^\dagger = \overline{f_{i,j}}$.

We now prove that this defines a unitary structure on \mathcal{C} . Given any two morphism $f : A \rightarrow B$, we compute

$$\langle f | f \rangle = \text{tr}(f^\dagger \circ f) = \sum_{i,j} |f_{i,j}|^2 d_{A_i}.$$

Clearly, $\langle f | f \rangle = 0$ if and only if $f = 0$ because all of the quantum dimensions d_{A_i} are positive real numbers. Hence, $\langle \cdot | \cdot \rangle$ is positive definition.

It is clear that $(f^\dagger)^\dagger = f$, $(f \circ g)^\dagger = g^\dagger \circ f^\dagger$, and $(f \otimes g)^\dagger = f^\dagger \otimes g^\dagger$ for appropriate choices of f, g . It remains to show that $(\text{ev}_A^R)^\dagger = \text{coev}_A^L$ and $(\text{coev}_A^R)^\dagger = \text{ev}_A^L$.

WORK: I don't know how to do this part of the proof. It's something to come back to. \square

1.8.2 Unitary braidings

WORK: Show that every braiding is automatically unitary. This is the content of [Gal14].

1.8.3 Uniqueness of unitary structure

WORK: Show that unitary structures are unique. This is the content of [Reu23b].

1.8.4 Skeletonization of unitarity

WORK: Show that a fusion category admits a unitary structure if and only if its F and R symbols can be made unitary in some basis. There are good references for this in the papers cited above.

1.9 Number theory in modular categories

1.9.1 Techniques and first results

WORK: A good general reference about this stuff is [GS19]. To-read, for sure. There is also interesting work in [MS12] which gives some counterexamples. Another reference is [DHW13].

1.9.2 Galois conjugation

WORK: An early reference is [CG93]. Some other papers not to ignore in this area are [PSYZ23, BR22].

1.9.3 Ocneanu rigidity

1.9.4 Rank-finiteness theorem

WORK: Of course there is the original paper on the topic. However, there is also the generalization for G -crossed MTCs and fermionic MTCs - [JMNR21]. Should I bring this up now or later?

1.9.5 Schauenberg-Ng theorem

WORK: Go through Schauenberg-Ng's original paper and understand the proof. It seems on the face of it like it is a hard theorem. Certainly, it uses strongly the theory of the Drinfeld center as well as the modular representation. It is good to push this proof as far down as possible since it will use a lot of machinery. I think it can be boiled down to something elegant, though, if the machinery has been set up.

WORK: I would quite like to show that fusion categories (or at least modular categories) have finitely automorphisms of the identity functor. The proof I know follows from generalities about the universal grading group - [GN08]. I wonder if this proof can be done completely without the use of grading. Then, the fact that the universal grading group can be interpreted in terms of grading can be put as an exercise in the “further structure” section!

WORK: this section is going to host a lot more theorems

History and further reading:

Modular categories were born from conformal field theory in the late 1980s. In a series of papers, Moore and Seiberg analysed deeply the underlying content within conformal field theory to find what essential algebraic data lied within it [MS88, MS89]. They wrote out the axioms of this essential algebraic data in their subsequent notes on conformal field theory [MS90]. They used the name modular category to describe their data, as suggested by Igor Frenkel. This definition was then refined and re-introduced by Turaev [Tur92]. The first major application of modular categories was the Reshetikhin-Turaev construction [RT91, Tur10b]. Prior to this result nobody had succeed in constructing topological quantum field theories. In this way, modular categories and the Reshetikhin-Turaev construction completed Witten's programme of quantizing Chern-Simons theory.

By the early 2000s, the proposal of topological quantum computing was attracting a lot of interest in anyons and their algebraic properties. Seeing as topological order can be described by topological quantum field theory and topological quantum field theory is essentially equivalent to modular categories, it was understood that modular categories could be used to understand topological order. This latent description of anyons in terms of modular categories was made explicit in an appendix in the seminal 2006 paper of Kitaev [Kit06]. This approach to anyons in terms of modular categories was popularized by Wang's early monograph [Wan10]. This has since become the standard approach towards the algebraic theory of topological quantum information.

Exercises:

- 1.1. WORK: apply Verlinde formula to group-theoretical modular categories to recover classical theorem by Burnside
- 1.2. WORK: show that irreducible G -graded G -reps are equivalent to irreducible reps of centralizers of conjugacy classes
- 1.3. WORK: One of the pivotal axioms automatically holds in fusion categories. Namely, the one with the twists going the two different ways. The proof is very standard - reduce to the case A, B are both simple, reduce to case $A = B$, $f = \lambda \cdot \text{id}_A$ and hence commutes with everything. This would be a nice exercise.
- 1.4. WORK: there's a notion of prime factorization of MTCs. If \mathcal{C} is an MTC, then it can be decomposed as a tensor product

$$\mathcal{C} = \mathcal{C}_1 \boxtimes \mathcal{C}_2 \dots \boxtimes \mathcal{C}_n,$$

where $\mathcal{C}_i \neq 0$ all have no proper non-degenerate braided fusion subcategories other than $\text{Vec}_{\mathbb{C}}$. We call such MTCs *prime MTCs*. Suppose that \mathcal{C} has NO abelian anyons. Then, the decomposition is unique. However, if \mathcal{C} has abelian anyons then the decomposition can fail to be unique. In particular, the color code is equivalent to both bilayer toric code and bilayer 3-fermion (c.f. [KPEB18]). The factorization results I asserted come from the (largely ignored?) paper by Muger: [Müg02]. This could make for a very nice exercise. First, prove Muger's double centralizer theorem. Then, prove the factorization

into prime MTCs. Then, if \mathcal{C} has no abelian anyon anyons, prove that the decomposition is unique. Then, prove that bilayer toric code is equivalent to bilayer 3-fermion (or maybe just leave that as a comment)!

[WORK:

List of things to add:

- A big “physics-math dictionary” which allows the translation of everything;
- A table of notation;

]

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