

Lectures on Vertex Operator Algebras and Conformal Blocks

BIN GUI

February 21, 2022

Contents

0	Notations	2
1	Segal's picture of 2d CFT; motivations of VOAs and conformal blocks	3
	Index	12
	References	13

0 Notations

- $\mathbb{N} = \{0, 1, 2, \dots\}$, $\mathbb{Z}_+ = \{1, 2, \dots\}$.
- $\mathbf{i} = \sqrt{-1}$, \mathbb{S}^1 =unit circle, $\mathbb{C}^\times = \mathbb{C} \setminus \{0\}$.
- $\mathbb{D}_r = \{z \in \mathbb{C} : |z| < r\}$, $\mathbb{D}_r^\times = \{z \in \mathbb{C} : 0 < |z| < r\}$, $\mathbb{D}_r^{\text{cl}} = \{z \in \mathbb{C} : |z| \leq r\}$
- $\mathcal{O}(X)$ is the space of holomorphic functions on a complex manifold X .
- Configuration space $\text{Conf}^n(X) = \{(x_1, \dots, x_n) \in X^n : x_i \neq x_j \text{ if } i \neq j\}$.
- z and ζ could mean either points, or the standard coordinate of \mathbb{C} , or formal variables. We will give their meanings when the context is unclear.
- All vector spaces are over \mathbb{C} , unless otherwise stated. If W is a vector space equipped with a Hermitian form $\langle \cdot | \cdot \rangle$, we let $|\cdot\rangle$ be the linear variable and $\langle \cdot |$ be the antilinear (i.e. conjugate linear) one, following physicists' convention.
- If W, W' are vector spaces, then $\text{Hom}(W, W')$ denote the space of linear operators from W to W' . We let $\text{End}(W) = \text{Hom}(W, W)$.
- We use symbols $\langle \cdot, \cdot \rangle$ or (\cdot, \cdot) to denote bilinear forms (i.e., linear on both variables).
- Given a vector space W and a formal variable z ,

$$W[z] = \{\text{polynomials of } z \text{ whose coefficients are elements of } W\}$$

$$W[[z]] = \left\{ \sum_{n \in \mathbb{N}} w_n z^n : w_n \in W \right\}$$

$$W((z)) = \left\{ \sum_{n \in \mathbb{Z}} w_n z^n : w_n \in W, \text{ and } w_n = 0 \text{ when } n \text{ is sufficiently negative} \right\}$$

$$W[[z^{\pm 1}]] = \left\{ \sum_{n \in \mathbb{Z}} w_n z^n : w_n \in W \right\}.$$

Each line is a subspace of the subsequent line. In case there are several formal variables, the spaces are defined in a similar way, expect $W((\dots))$. For instance,

$$W[[z, \zeta^{\pm 1}]] := W[[z]][[\zeta^{\pm 1}]] = W[[\zeta^{\pm 1}]][[z]]$$

consists of $\sum_{m \in \mathbb{N}, n \in \mathbb{Z}} w_{m,n} z^m \zeta^n$ where each $w_{m,n} \in W$. However, note that $W((z))((\zeta))$ and $W((\zeta))((z))$ are not equal. (For instance, $\sum_{m \geq -n} \sum_{n \geq -1} z^m \zeta^n$ belongs to $\mathbb{C}((z))((\zeta))$ but not $\mathbb{C}((\zeta))((z))$.)

- We set

$$\text{Res}_{z=0} \sum_{n \in \mathbb{Z}} w_n z^n dz = w_{-1}. \quad (0.2)$$

This is in line with the complex analytic residue.

- A vector of $W_1 \otimes \cdots \otimes W_N$ written as w_\bullet means that it is of the form $w_1 \otimes \cdots \otimes w_N$ where each $w_i \in W_i$. Depending on the context, w_\bullet will also mean a tuple (w_1, \dots, w_N) .
- Unless otherwise stated, by a manifold, we mean one *without* boundaries. Also, "with boundaries" means "possibly with boundaries".

1 Segal's picture of 2d CFT; motivations of VOAs and conformal blocks

1.1

Vertex operator algebras (VOAs) are mathematical objects defined to understand and construct 2-dimensional conformal field theory (CFT for short). A CFT describes propagations and interactions of strings. There are two types of strings: closed strings $\simeq \mathbb{S}^1$ and open strings $\simeq [0, 1]$. In this course, we will mainly focus on closed strings.

Let me explain how mathematicians understand CFT. Just like any quantum field theory (QFT), in CFT we must have a Hilbert space \mathcal{H} . The vectors in \mathcal{H} are called "states", but unlike ordinary QFT, a vector $\xi \in \mathcal{H}$ is not a state of a particle, but a state of a closed string \mathbb{S}^1 .

The most important and non-trivial part in CFT is to define/understand string interactions. According to Segal's picture [Seg88], an interaction is uniquely determined by a compact Riemann surface Σ with boundaries $\partial\Sigma$, where $\partial\Sigma$ is a disjoint union of some circles (strings). Each string is called either an incoming string or an outgoing one. Suppose $\partial\Sigma$ has N incoming strings and M outgoing ones, then this picture describes an interaction where N strings are going inside, and M strings are going outside.

Moreover, the boundary $\partial\Sigma$ must be **parametrized**. This means that to each connected component $\partial\Sigma_i$ a diffeomorphism $\eta_i : \partial\Sigma_i \xrightarrow{\sim} \mathbb{S}^1$ is associated. The orientation on $\partial\Sigma_i$ defined by pulling back the one of \mathbb{S}^1 along η_i is assumed to be the opposite of the one defined in Stokes' theorem, shown as follows



1.2

Unless otherwise stated, we assume that the boundary parametrization is also **analytic**. Roughly speaking, this means that Σ can be obtained by removing some open discs from a compact Riemann surface C (without boundary) such that the parametrizations of $\partial\Sigma$ are given by local holomorphic functions of C .

Here is a more rigorous explanation. By a **local coordinate** η of C at $x \in C$, we mean η is a holomorphic injective function on a neighborhood U of x such that $\eta(x) = 0$. So

η is a biholomorphism between U and a neighborhood $\eta(U)$ of 0. Now, suppose we have local coordinates η_1, \dots, η_N at distinct points $x_1, \dots, x_N \in C$. The data

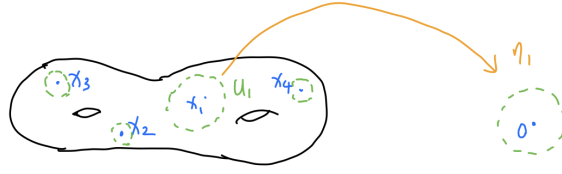
$$\mathfrak{X} := (C; x_\bullet; \eta_\bullet) = (C; x_1, \dots, x_N; \eta_1, \dots, \eta_N) \quad (1.2)$$

is called an **N -pointed compact Riemann surface with local coordinates**.

Let each η_i be defined on a neighborhood $U_i \ni x_i$. We assume moreover the following

Assumption 1.1. $U_i \cap U_j = \emptyset$ if $i \neq j$, and $\eta_i(U_i) \supset \mathbb{D}_1^{\text{cl}}$ for each i . Here \mathbb{D}_1^{cl} is the closed unit disc.

By removing all $\eta_i^{-1}(\mathbb{D}_1)$, we get Σ with boundary strings $\eta_i^{-1}(\partial\mathbb{D}_1^{\text{cl}}) = \eta_i^{-1}(\mathbb{S}^1)$ whose parametrizations are η_i .



1.3

Any Σ as above determines uniquely an interaction of strings. Suppose it has N incoming strings and M outgoing ones. Then mathematically, such an interaction is described by a bounded linear map $T = T_\Sigma : \mathcal{H}^{\otimes N} \rightarrow \mathcal{H}^{\otimes M}$. (The boundedness is automatic thanks to the uniform boundedness principle. But this is not an important point in this course.) Given $\xi_\bullet = \xi_1 \otimes \dots \otimes \xi_N \in \mathcal{H}^{\otimes N}$ and $\eta_\bullet = \eta_1 \otimes \dots \otimes \eta_M \in \mathcal{H}^{\otimes M}$, the value

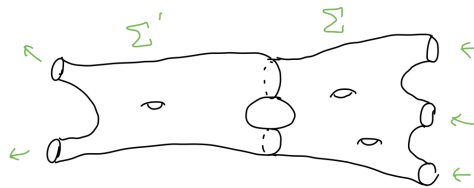
$$\langle \eta_\bullet | T \xi_\bullet \rangle \quad (1.3)$$

describes the probability amplitude that the N incoming closed strings with states ξ_1, \dots, ξ_N become η_1, \dots, η_M after interaction.

The word “conformal” in conformal field theory reflects the fact that T depends only on the complex structure of Σ and its parametrization, but not on the metric for instance. Thus, a CFT is more rigid than a topological quantum field theory (TQFT): in the latter case, T depends only on the topological structures of the manifolds.

1.4

Suppose we have another interaction $S : \mathcal{H}^{\otimes M} \rightarrow \mathcal{H}^{\otimes L}$ corresponding to the parametrized surface Σ' , then the composition of them $S \circ T : \mathcal{H}^{\otimes N} \rightarrow \mathcal{H}^{\otimes L}$ corresponds to the **sewing** $\Sigma \# \Sigma'$ of Σ and Σ' , where the j -th outgoing string $\partial_+ \Sigma_j$ of Σ is sewn with the j -th incoming one $\partial_- \Sigma'_j$ of Σ' .



It is important to specify how $\partial_+ \Sigma_j$ (with parametrization η_j) is identified with $\partial_- \Sigma'_j$ (with parametrization η'_j). Pick $x \in \partial_+ \Sigma_j$ and $y \in \partial_- \Sigma'_j$. Then

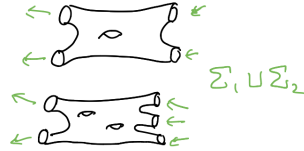
$$x = y \iff \eta_j(x) \eta'_j(y) = 1. \quad (1.4)$$

It is clear from the picture that the orientations of $\partial_+ \Sigma_j$ and $\partial_- \Sigma_j$ are opposite to each other. This is related to the fact that our rule for sewing is $\eta_j(x) = 1/\eta'_j(y)$ but not (say) $\eta_j(x) = \eta'_j(y)$.

Recall we assume that the parametrizations are analytic. We leave it to the readers to check that the sewing of Σ and Σ' , a priori only a topological surface, has a natural complex analytic structure.

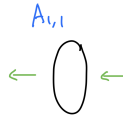
1.5

Suppose $T_1 : \mathcal{H}^{\otimes N_1} \rightarrow \mathcal{H}^{\otimes M_1}$ corresponds to Σ_1 and $T_2 : \mathcal{H}^{\otimes N_2} \rightarrow \mathcal{H}^{\otimes M_2}$ to Σ_2 , then $T_1 \otimes T_2 : \mathcal{H}^{\otimes (N_1+N_2)} \rightarrow \mathcal{H}^{\otimes (M_1+M_2)}$ corresponds to the disjoint union $\Sigma_1 \sqcup \Sigma_2$.



1.6

Consider an annulus $A_{r,R} = \{z \in \mathbb{C} : r < |z| < R\}$ obtained by removing two open discs from the compact Riemann sphere \mathbb{P}^1 via the local coordinate $\eta_1(z) = z/r$ at 0 and $\eta_2(z) = R/z$ at ∞ . We call such $A_{r,R}$ (with the given boundary parametrization) a **standard annulus**. Let $r \nearrow 1, R \searrow 1$. The limit of this annulus is a “degenerate” Riemann surface with 1 incoming boundary circle and 1 outgoing one. Both circles are \mathbb{S}^1 . The incoming one has parametrization $z \mapsto z$ and the outgoing one $z \mapsto z^{-1}$. We call this annulus the **standard thin annulus** and denote it by $A_{1,1}$. The map $T : \mathcal{H} \rightarrow \mathcal{H}$ associated to $A_{1,1}$ is the identity map. This reflects the fact that sewing any Σ with a disjoint union of $A_{1,1}$ gives Σ .



1.7

We give a fancy way to summarize what we have so far: Let \mathcal{C} be the monoidal category of compact 1-dimensional smooth manifolds such that a morphism from an object S_1 to another S_2 is a compact Riemann surface with incoming parametrized boundary $\simeq S_1$ and outgoing one $\simeq S_2$, that the identity morphism for a union of N circles is a disjoint union of N pieces of $A_{1,1}$, that the unit object is the empty set, and that the tensor product of objects and morphisms are respectively the disjoint unions of strings and Riemann surfaces. Then a CFT is a monoidal functor from \mathcal{C} to the

monoidal category of Hilbert spaces. So, roughly speaking, a CFT is a representation of \mathcal{C} .

Since we choose Hilbert spaces as our underlying spaces, we should expect that the representation of \mathcal{C} is unitary. Technically, the functor mentioned above should be a $*$ -functor: this means that for each morphism Σ from N strings to M strings, we should define its adjoint morphism Σ^* from M strings to N ones whose corresponding map is the adjoint $T^* : \mathcal{H}^{\otimes M} \rightarrow \mathcal{H}^{\otimes N}$ of T . Σ^* is defined simply to be the **complex conjugate** $\bar{\Sigma}$ of Σ :

Definition 1.2. $\bar{\Sigma}$ consists of points \bar{x} where $x \in \Sigma$; the local holomorphic functions on $\bar{\Sigma}$ are η^* where η is a locally defined holomorphic function on Σ and

$$\eta^*(\bar{x}) = \overline{\eta(x)} \quad (1.5)$$

whenever η is defined on $x \in \Sigma$; similarly, boundary parametrizations are given by η_j^* . Note that if Σ is obtained by removing open discs from an N pointed $\mathfrak{X} = (C; x_\bullet; \eta_\bullet)$, then $\bar{\Sigma}$ is obtained by removing discs from

$$\bar{\mathfrak{X}} := (\bar{C}; \bar{x}_1, \dots, \bar{x}_N; \eta_1^*, \dots, \eta_N^*) \quad (1.6)$$

η^* should not be confused with $\bar{\eta}$ defined on Σ by

$$\bar{\eta}(x) = \overline{\eta(x)}.$$

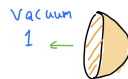
In the present context, we should assume that an incoming (resp. outgoing) string of Σ becomes an outgoing (resp. incoming) one of $\bar{\Sigma}$ via the conjugate map $\mathbb{C} : x \in \Sigma \mapsto \bar{x} \in \bar{\Sigma}$. In the future, we will often consider all strings as incoming ones if necessary (cf. 1.9). In that case, we shall also assume all the boundary strings of $\bar{\Sigma}$ as incoming.

We should point out that although unitarity is a very important condition, there are important non-unitary CFTs, for instance, the logarithmic CFTs. (In such cases, \mathcal{H} is a vector space without inner products.) Also, many VOA results and techniques do not rely on the unitarity. Nevertheless, assuming unitarity will often reasonably simplify discussions or give motivations.

Example 1.3. Let $\mathfrak{X} = (\mathbb{P}^1; 0; \lambda\zeta)$ where ζ is the standard coordinate of \mathbb{C} and $\lambda \in \mathbb{C}^\times$. We can identify the conjugate of \mathbb{P}^1 with \mathbb{P}^1 by letting $x \in \mathbb{P}^1 \mapsto \bar{x}$ be the standard conjugate of \mathbb{C} : $z \mapsto \bar{z}$. Then $(\lambda\zeta)^*(\bar{z}) = \overline{\lambda\zeta(z)} = \bar{\lambda} \cdot \bar{z} = \bar{\lambda}\zeta(\bar{z})$. So the conjugate of \mathfrak{X} is isomorphic to $\bar{\mathfrak{X}} = (\mathbb{P}^1; 0; \bar{\lambda}\zeta)$.

1.8

An interaction process could have no incoming or outgoing strings. *The Hilbert space for the empty string \emptyset is \mathbb{C} .* The most elementary and important example with no incoming boundary is the closed unit disc \mathbb{D}_1^{cl} with 1 outgoing boundary parametrized by $z \mapsto z^{-1}$. The corresponding map $\mathbb{C} \rightarrow \mathcal{H}$ can be identified with its value at 1. This element in \mathcal{H} is denoted by **1** and called the **vacuum vector**.



$$(1.7)$$

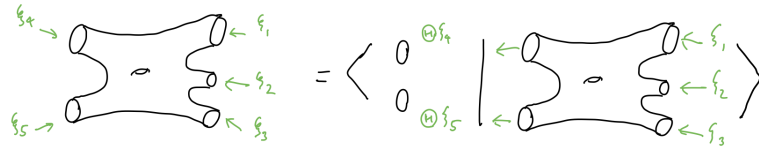
Assume as before that our theory is unitary. Then conjugate of the above disk is the same disk and boundary parametrization, but the original outgoing string is now the incoming one. The corresponding map $\mathcal{H} \rightarrow \mathbb{C}$ is, according to 1.7, the linear functional $\langle \Omega | \cdot \rangle$.

1.9

In general, one may wonder what the interaction $T : \mathcal{H}^{\otimes N} \rightarrow \mathbb{C}$ means physically for a surface Σ with N incoming strings but no outgoing ones. Choose $0 < M < N$, and make M of the N strings of $\partial\Sigma$ be outgoing strings. Then the corresponding interaction is a map $\tilde{T} : \mathcal{H}^{\otimes(N-M)} \rightarrow \mathcal{H}^{\otimes M}$. In unitary CFT, T can be related to \tilde{T} by a anti-unitary (i.e. conjugate-unitary) map Θ on \mathcal{H} , called the **CPT operator**, such that for $\xi_1, \dots, \xi_N \in \mathcal{H}$ (where the last M vectors are associated to the outgoing strings), we have

$$T(\xi_1 \otimes \dots \otimes \xi_N) = \langle \Theta \xi_{N-M+1} \otimes \dots \otimes \Theta \xi_N | \tilde{T}(\xi_1 \otimes \dots \otimes \xi_{N-M}) \rangle, \quad (1.8)$$

interpreted pictorially as



The operator Θ is an involution, i.e., $\Theta^2 = 1_{\mathcal{H}}$.

Such a linear functional T corresponding to an interaction with no outgoing strings is called a **correlation function** (or an **N -point function**). These functions are the central objects in CFT (and indeed, in any quantum field theory). Relation (1.8) teaches us that: (1) correlation functions can be interpreted as probability amplitudes in string interactions with the help of Θ , and (2) to study arbitrary interactions, it suffices to study those with no outgoing strings.

Let me close this subsection by mentioning an important fact: suppose the complex structure of Σ and the (assumed analytic) boundary parametrizations are parametrized holomorphically by some complex variables $\tau_{\bullet} = (\tau_1, \dots, \tau_k)$, then the value of $T(\xi_{\bullet})$ is now a *real analytic function* of τ_{\bullet} , i.e., it is locally a power series of τ_1, \dots, τ_k and their conjugates. Actually, the word “function” in “correlation function” means a function of τ_{\bullet} , but not of ξ_{\bullet} .

1.10

You must be curious what CPT means. Indeed, Θ is responsible for the simultaneous symmetry of charge conjugation (C), parity transformation (P), and time reversal (T). P+T together means an *anti-biholomorphism* $\Sigma \rightarrow \Sigma'$. Now we have arrived at a point that we missed previously: since anti-holomorphic maps are also conformal maps, should we expect that the interaction maps (or the correlation functions) for anti-biholomorphic surfaces are equal? The answer is no. (Namely, P+T are not preserved.) Indeed, if we let Σ has N incomes and no outcomes, let $\bar{\Sigma}$ be its complex

conjugate (cf. 1.7) but still with N incomes, and let $T_\Sigma, T_{\bar{\Sigma}}$ be the correlation functions associated to them, then from 1.7 and relation (1.8), one can in fact check that

$$T_\Sigma(\xi_1 \otimes \cdots \otimes \xi_N) = \overline{T_{\bar{\Sigma}}(\Theta \xi_1 \otimes \cdots \otimes \Theta \xi_N)}. \quad (1.9)$$

This relation explains CPT symmetry: the symmetries of charge (taking complex conjugate of the values of correlation functions) and parity+time (the conjugate biholomorphism $\mathbb{C} : \Sigma \rightarrow \bar{\Sigma}$) are preserved, and the operator realizing this simultaneous symmetry is Θ .

Note that mathematically, charge conjugation C is related to taking complex conjugate of numbers (but not of Σ). Physically, it means making a string into its “anti-string”, or (in general QFT) making a particle (e.g. an electron with negative charge) to its anti-particle (e.g. an antielectron with positive charge).

1.11

The CFT we have described so far is actually very special: it has no conformal anomaly. There are indeed no nontrivial CFTs which are both unitary and without anomaly. In this course, we will be mainly interested in CFTs with conformal anomaly. Technically, the conformal anomaly is determined by a complex number c (positive for unitary CFT), called **central charge**. To describe such CFT, we modify the previous descriptions as follows: The map (or the correlation function) T_Σ for Σ is only up to a positive scalar multiplication depending on Σ . $T_{\Sigma_1} \circ T_{\Sigma_2} = \lambda T_{\Sigma_1 \# \Sigma_2}$ where $\lambda > 0$. (The constants are not necessarily positive in non-unitary CFT.) If Σ is parametrized holomorphically by some complex variables τ_\bullet , then by shrinking the domain of τ_\bullet , we can choose T_Σ depending real analytically on τ_\bullet .

There are many important cases where a real analytic (or even a holomorphic) T_Σ can be chosen globally for τ_\bullet . This will be studied later in details.

Unless otherwise stated, a CFT always means one with (possible) conformal anomaly. Using the fancy language of 1.7, one can say that a unitary CFT is a *projective* monoidal $*$ -functor from the category \mathcal{C} in 1.7 to the category of Hilbert spaces. Namely, it is a projective unitary representation of \mathcal{C} .

1.12

To study the representations of a topological group G , one must first understand very well the topological and the algebraic structures of G . Similarly, the study of CFTs relies heavily on the geometric and analytic structures of compact Riemann surfaces. However, from what we have discussed, there is a huge obstacle for studying CFTs: the correlation functions are real analytic, but not complex analytic (i.e. holomorphic) functions of the parameters τ_\bullet . Thus, in order to study CFTs using the powerful tools of complex analysis (residue theorem, for instance), we make the following Ansatz: A correlation function T is a sum : $T_\Sigma = \sum_j \Phi_\Sigma^j \cdot \Psi_{\bar{\Sigma}}^j$, where each Φ^j and Ψ^j relies holomorphically on Σ and $\bar{\Sigma}$ respectively (so $\Psi_{\bar{\Sigma}}^j$ relies anti-holomorphically on Σ).

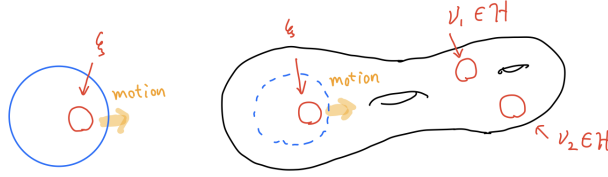
This Ansatz is very vague. Let me explain it in more details. Consider the annulus $A_{r,R}$ with boundary parametrization as in 1.6. We move the inside circle to another one

centered at z (where $z \in A_{r,R}$ is reasonably small), still with radius r . The new eccentric annulus $A_{z,r,R}$ has larger outgoing string parametrized by $\zeta \mapsto R/\zeta$ and the smaller incoming one parametrized by $\zeta \mapsto (\zeta - z)/r$. Let $T_z : \mathcal{H} \rightarrow \mathcal{H}$ be the corresponding map. As we have said, for general vectors $\xi, \eta \in \mathcal{H}$, the expression $\langle \eta | T_z \xi \rangle = \langle \Theta \eta, T_z \xi \rangle$ can be chosen to be real analytic with respect to z . We now let

$$\begin{aligned} \mathbb{V} = \{ \xi \in \mathcal{H} : & \text{For all } r, R, \text{ the map } T \text{ can be chosen such that} \\ & z \mapsto \langle \nu | T_z \xi \rangle \text{ is holomorphic for all } \nu \in \mathcal{H}, \text{ and} \\ & \xi \text{ has "finite energy"} \} \end{aligned} \quad (1.10)$$

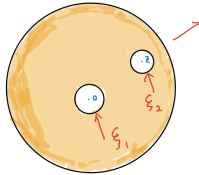
“Finite energy” is a minor condition to be explained later. (See ??.)

We can sew $A_{z,r,R}$ with any Σ , and the motion of the smaller string inside the annulus becomes, after sewing, the motion of a boundary string of Σ :



Therefore, if a vector $\xi \in \mathbb{V}$ is assigned to an incoming string of Σ with (analytic) boundary parametrization η_i , then, when translating this parametrized string with respect to η_i , the correlation function $T_\Sigma(\xi \otimes \dots)$ should be holomorphic with respect to the motion, whatever states we assign to the other strings. We can therefore study \mathbb{V} with the help of complex analysis. \mathbb{V} is called a **vertex operator algebra** (VOA).

We have only described \mathbb{V} as a vector space. But in which sense is \mathbb{V} an algebra? An obvious candidate is as follows: consider \mathbb{P}^1 with three marked points $0, z, \infty$ and usual coordinates, e.g. $\eta_0(\zeta) = \zeta/r_1, \eta_z(\zeta) = (\zeta - z)/r_2, \eta_\infty(\zeta) = R/\zeta$ at $0, z, \infty$ where $r_1, r_2 > 0$ are small and $R > 0$ is large. We assume the strings around 0 and z are ingoing and that around ∞ outgoing. If we assign $\xi_1, \xi_2 \in \mathbb{V}$ to the incoming strings, then the outcome can be viewed as a product of ξ_1 and ξ_2 .



Although this product does not have finite energy, it does satisfy the statement before the last line in (1.10). Thus, this product is almost a vector in \mathbb{V} . By modifying this product suitably, we can ensure that the products of vectors in \mathbb{V} are always in \mathbb{V} . Details will be given in later sections.

Similarly to (1.10), we define $\hat{\mathbb{V}} \subset \mathcal{H}$ to be the set of finite energy vectors ξ such that $\langle \nu | T_z \xi \rangle$ is anti-holomorphic over z . The vacuum vector 1 belongs to $\mathbb{V} \cap \hat{\mathbb{V}}$: The result of gluing the unit disc into the inside of $A_{z,r,R}$ is just the disc with radius R and parametrization $\zeta \mapsto R/\zeta$, which is independent of z . So $T_z 1$ and hence $\langle \nu | T_z 1 \rangle$ are constant over z , and hence both holomorphic and anti-holomorphic over z .

1.13

Now we can give a more detailed presentation of our Ansatz. We let \mathcal{H}^{fin} be the (indeed dense) subspace of vectors in \mathcal{H} with “finite energy”, which is acted on by $\mathbb{V} \otimes \widehat{\mathbb{V}}$. Ansatz:

1. \mathcal{H}^{fin} as a $\mathbb{V} \otimes \widehat{\mathbb{V}}$ -module has decomposition

$$\mathcal{H}^{\text{fin}} = \bigoplus_{i \in \mathcal{I}} \mathbb{W}_i \otimes \widehat{\mathbb{W}}_i \quad \supset \mathbb{V} \otimes \widehat{\mathbb{V}} \quad (1.11)$$

where each $\mathbb{W}_i, \widehat{\mathbb{W}}_i$ are respectively irreducible \mathbb{V} -modules and $\widehat{\mathbb{V}}$ -modules. \mathbb{V} and $\widehat{\mathbb{V}}$ are (according to their definition cf. (1.10)) subspaces of \mathcal{H}^{fin} by identifying them with $\mathbb{V} \otimes 1$ and $1 \otimes \widehat{\mathbb{V}}$ respectively. The vacuum vector 1 of \mathcal{H} is identified with $1 \otimes 1$ (which belongs to $\mathbb{V} \otimes \widehat{\mathbb{V}}$).

2. For some Σ without outgoing boundaries, let $T_\Sigma : \mathcal{H}^{\otimes N} \rightarrow \mathbb{C}$ be the corresponding map. Then, corresponding to the above direct sum decomposition, we have

$$T_\Sigma \Big|_{(\mathcal{H}^{\text{fin}})^{\otimes N}} = \sum_{i_1, \dots, i_N \in \mathcal{I}} \Phi_{\Sigma, i_\bullet} \otimes \Psi_{\bar{\Sigma}, i_\bullet} \quad (1.12)$$

where

$$\begin{aligned} \Phi_{\Sigma, i_\bullet} &: \mathbb{W}_{i_1} \otimes \dots \otimes \mathbb{W}_{i_N} \rightarrow \mathbb{C}, \\ \Psi_{\bar{\Sigma}, i_\bullet} &: \widehat{\mathbb{W}}_{i_1} \otimes \dots \otimes \widehat{\mathbb{W}}_{i_N} \rightarrow \mathbb{C} \end{aligned}$$

are linear. Moreover, when the complex structure and boundary parametrization are parametrized analytically by complex variables τ_\bullet , then locally (with respect to the domain of τ_\bullet), $T_\Sigma, \Phi_{\Sigma, i_\bullet}, \Psi_{\bar{\Sigma}, i_\bullet}$ can be chosen such that Φ_{Σ, i_\bullet} is holomorphic over τ_\bullet (for all input vectors), and $\Psi_{\bar{\Sigma}, i_\bullet}$ holomorphic over $\bar{\tau}_\bullet$. Φ_{Σ, i_\bullet} and $\Psi_{\bar{\Sigma}, i_\bullet}$ are called **conformal blocks** associated to Σ (resp. $\bar{\Sigma}$) and \mathbb{V} (resp. $\widehat{\mathbb{V}}$).

In part one, \bigoplus could be finite (our main focus in this course), infinite but discrete, or continuous.

The second part can be summarized by saying that the CFT is separated into the **chiral halves** (those Φ or \mathbb{W}_i) and the **anti-chiral halves** (those Ψ or $\widehat{\mathbb{W}}_i$). Here, “chiral”=“holomorphic”.

When physicists say a CFT is **rational**, they usually mean that the above direct sum is finite, and each $\mathbb{W}_{i,k}$ is semi-simple (hence, by further decomposition, can be irreducible). So far, the mathematical theory of conformal blocks is complete almost only for rational CFTs. These will be the main examples of this course. For non-rational logarithmic CFTs, even the above Ansatz needs to be modified. (So far, it is not even clear how to do it.)

Physicists more or less consider the above description as the definition of conformal blocks. We mathematicians should do the opposite: define conformal blocks in a different way, and use them to *construct* CFTs following the above Ansatz.

1.14

You may notice that to make this Ansatz compatible with 1.4 and 1.5, it is necessarily to assume that

1. The tensor product of conformal blocks $\Phi_{\Sigma_1}, \Phi_{\Sigma_2}$ associated to Σ_1, Σ_2 respectively should be a conformal block associated to $\Sigma_1 \sqcup \Sigma_2$.
2. The composition of $\Phi_{\Sigma_1}, \Phi_{\Sigma_2}$ (or more precisely, their contractions) should be conformal blocks associated to the sewings of Σ_1 and Σ_2 , where the pair of \mathbb{V} -modules to be contracted must be dual to each other.

$$T_{\Sigma_1}(\xi_1 \otimes \xi_2 \otimes \xi_3) \otimes T_{\Sigma_2}(\nu_1 \otimes \nu_2 \otimes \nu_3) = T_{\Sigma_1 \sqcup \Sigma_2}(\xi_1 \otimes \dots \otimes \nu_3)$$

(A side note on linear algebra: If V^\vee is the dual space (or a suitable dense subspace of the dual space) of a vector space V , we choose a basis $\{v_\alpha\}_{\alpha \in \mathfrak{A}}$ labeled by elements of \mathfrak{A} , and choose a dual basis $\{v_\alpha^\vee\}_{\alpha \in \mathfrak{A}}$ of V^\vee (i.e. the one determined by $\langle v_\alpha, v_\beta \rangle = \delta_{\alpha, \beta}$), then taking contraction means substituting $\sum_{\alpha \in \mathfrak{A}} v_\alpha \otimes v_\alpha^\vee$ inside the linear functional on a tensor product of vector spaces such that V, V^\vee are tensor components.)

After we define conformal blocks rigorously, we will see that the first point is obvious, while the second one is a non-trivial theorem.

We briefly explain the meaning of “dual”: For instance, in the above picture, the unitary \mathbb{V} -module containing ξ_2 is dual to the one containing η_1 . As vector spaces, they are (“graded”) dual spaces of each other. (It is a dense subspace of the full dual space. We will talk about this in future sections.) In unitary CFTs, all \mathbb{V} and $\widehat{\mathbb{V}}$ modules are unitary, and the CPT operator Θ maps each $\mathbb{W}_i \otimes \widehat{\mathbb{W}}_i$ to some $\mathbb{W}'_i \otimes \widehat{\mathbb{W}}'_i$ in \mathcal{H}^{fin} where \mathbb{W}'_i is a \mathbb{V} -module dual to \mathbb{W}_i , and $\widehat{\mathbb{W}}'_i$ a $\widehat{\mathbb{V}}$ -module dual to $\widehat{\mathbb{W}}_i$. (\mathbb{V} is the dual of itself, and $\widehat{\mathbb{V}}$ similarly. So Θ restricts to involutions on \mathbb{V} and $\widehat{\mathbb{V}}$ respectively. Such \mathbb{V} and $\widehat{\mathbb{V}}$ are called **unitary VOAs**.) The formal name for dual module is **contragredient module**, to be defined rigorously in later sections.

1.15

Note that a conformal block with $M + N$ inputs $\Phi_\Sigma : \mathbb{W}_{i_1} \otimes \dots \otimes \mathbb{W}_{i_N} \otimes \mathbb{W}_{j_1} \otimes \dots \otimes \mathbb{W}_{j_M} \rightarrow \mathbb{C}$ can be regarded as one with N inputs and M outputs $\Phi_\Sigma : \mathbb{W}_{j_1} \otimes \dots \otimes \mathbb{W}_{j_M} \rightarrow \mathcal{H}'_{i_1} \otimes \dots \otimes \mathcal{H}'_{i_M}$ where \mathcal{H}'_{i_k} is the Hilbert space completion of \mathbb{W}'_{i_k} and \mathbb{W}'_{i_k} is the contragredient \mathbb{V} -module of \mathbb{W}_{i_k} .

Index

1 , the vacuum vector, [6](#)

$A_{r,R}$ and the standard thin annulus $A_{1,1}$, [5](#)

$\overline{C}, \overline{\Sigma}$, the complex conjugate of C and Σ , [6](#)

f^*, η^*, \dots and $\overline{f}, \overline{f}, \dots$, where $f^*(\overline{x}) = \overline{f(x)}$ and $\overline{f}(x) = \overline{f(x)}$, [6](#)

$T_\Sigma, T_{\mathfrak{X}}$: The interaction map/correlation function, [4](#)

Θ , the CPT operator, [7](#)

References

- [Seg88] Segal, G.B., 1988. The definition of conformal field theory. In Differential geometrical methods in theoretical physics (pp. 165-171). Springer, Dordrecht.
- [Was10] Wassermann, A., 2010. Kac-Moody and Virasoro algebras. arXiv preprint arXiv:1004.1287.

YAU MATHEMATICAL SCIENCES CENTER, TSINGHUA UNIVERSITY, BEIJING, CHINA.
E-mail: binguimath@gmail.com bingui@tsinghua.edu.cn