

Multi-particle azimuthal correlation with subevent method

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for the ATLAS collaboration

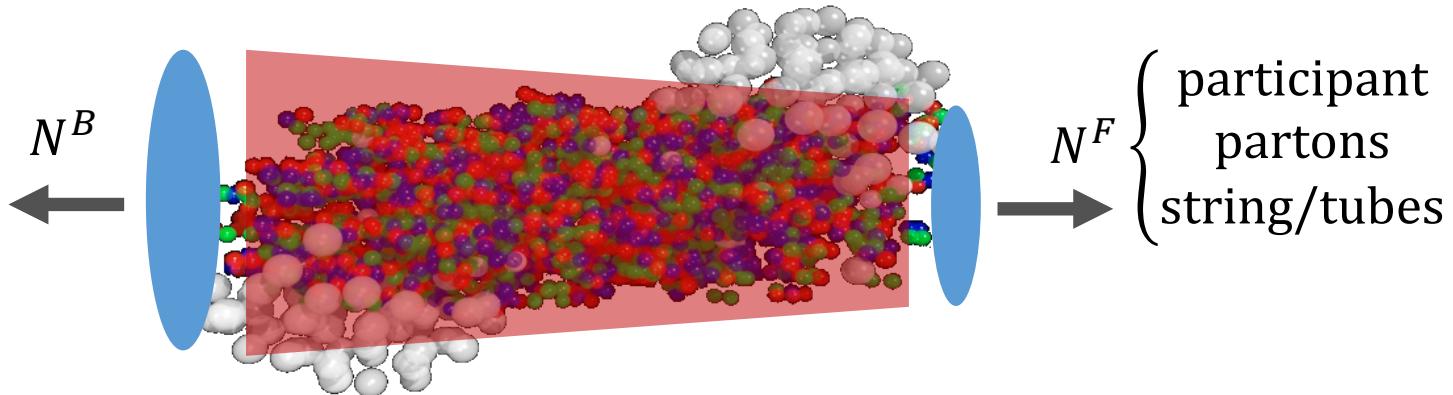
Quark Matter 2017, Feb. 5-11, Chicago

arXiv: 1606.08170

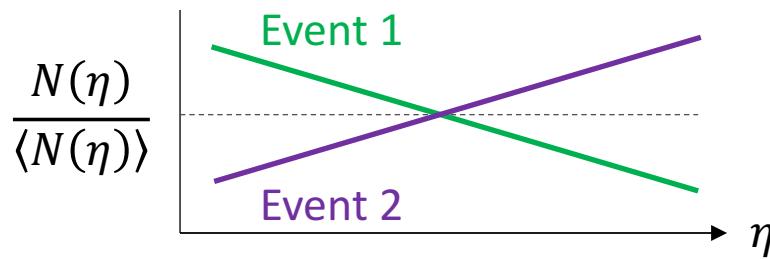
ATLAS-CONF-2017-002



Nature of sources seeding the long-range collective behavior?



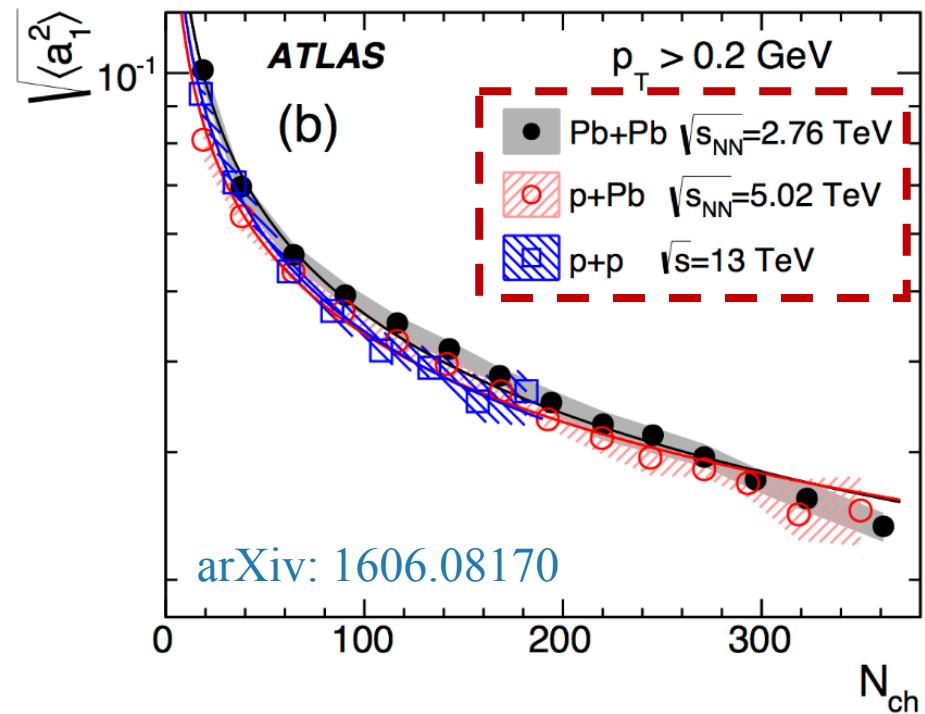
- $dN/d\eta$ shape reflects asymmetry in the number of forward-backward sources;



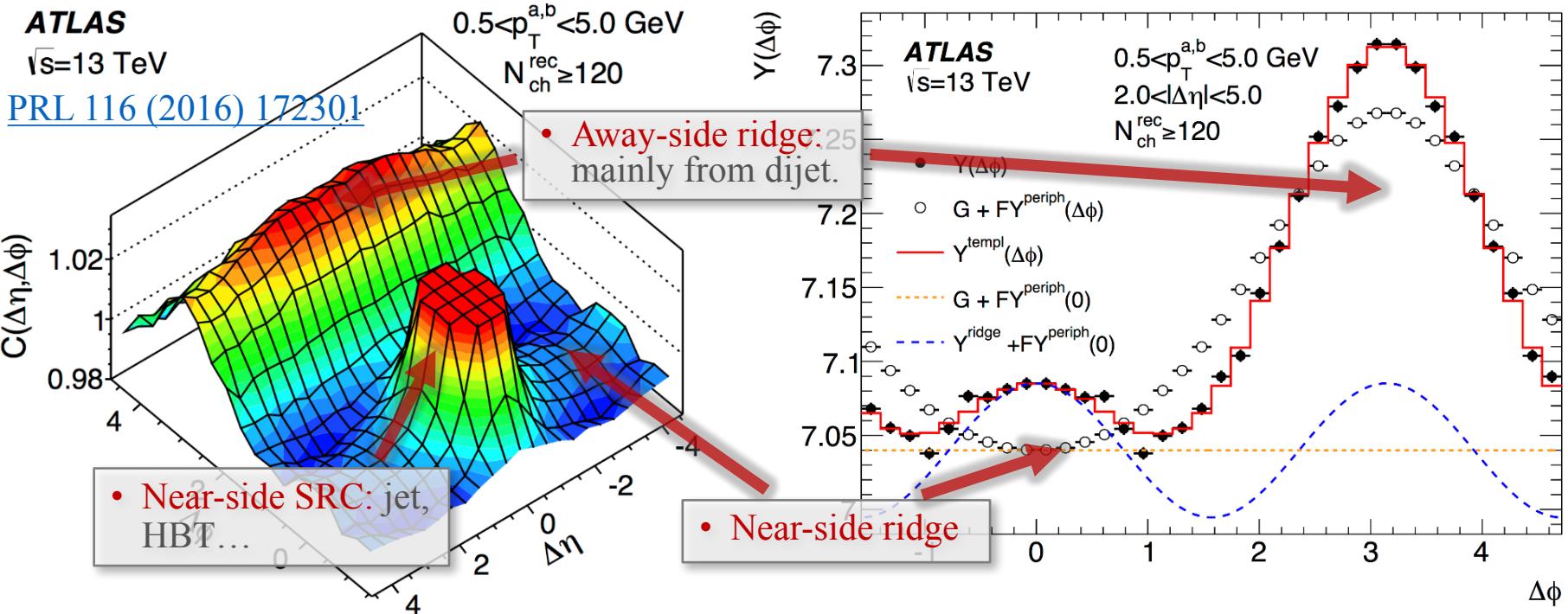
$$\frac{N(\eta)}{\langle N(\eta) \rangle} \approx 1 + a_1 \eta$$

[PRC 87, 024906 \(2013\)](#)
[PRC 93, 044905 \(2016\)](#)

- Dominated by linear fluctuation!
- Independent of collision system!



Long-range azimuthal correlation



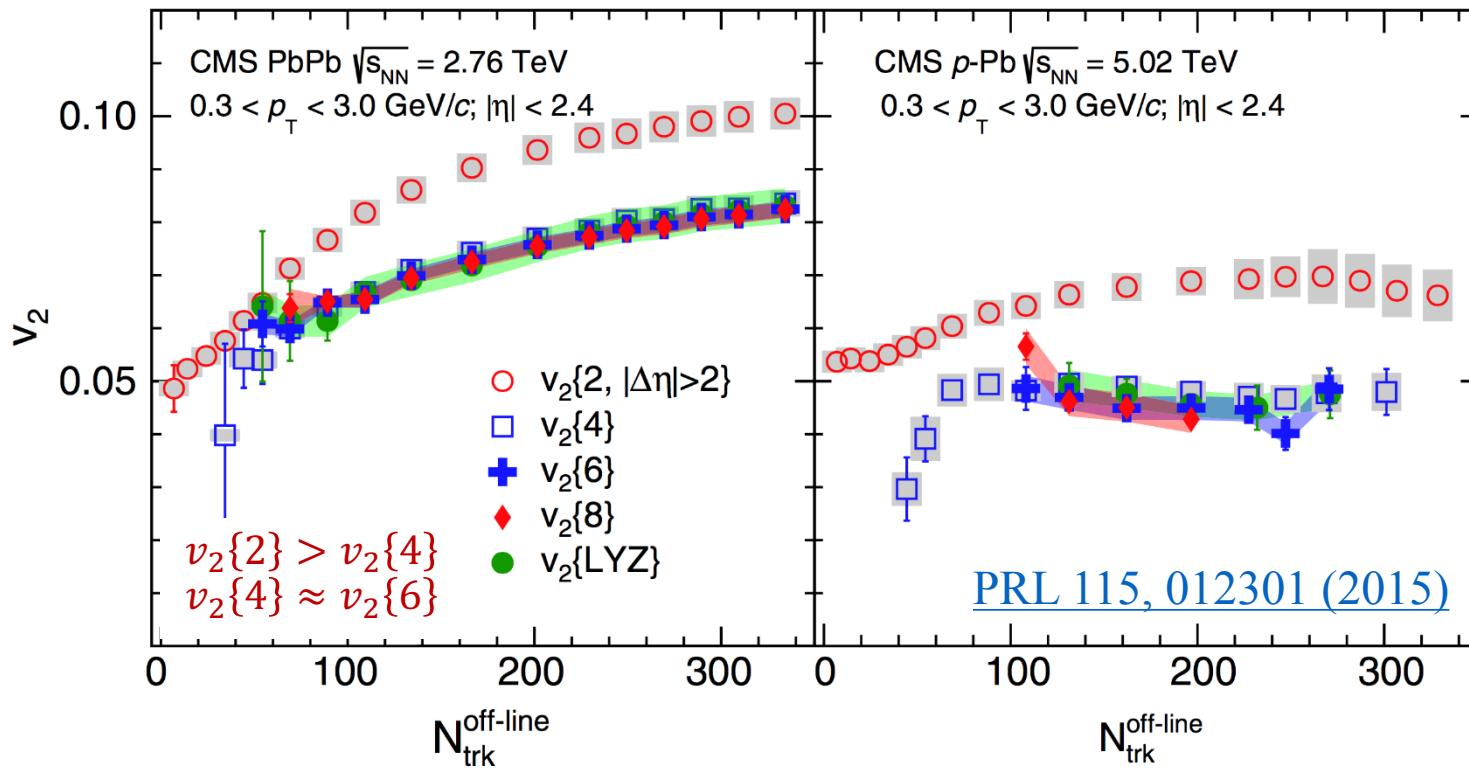
- Ridge structure also observed in small systems! Competing theories:
 - Initial stage interaction (CGC, ...)? [PRD 87 \(2013\) 094034](#)
 - Initial stage fluctuation + final stage interaction? [PRC 88 \(2013\) 014903](#)
- Fluctuation in proton can be constrained in pp and $p+\text{Pb}$ collisions.
- However, ridge hard to extract in data: larger non-flow in small system.

From the experimental perspective, central question about the ridge is whether it involves all particles (collectivity?).

Standard cumulant

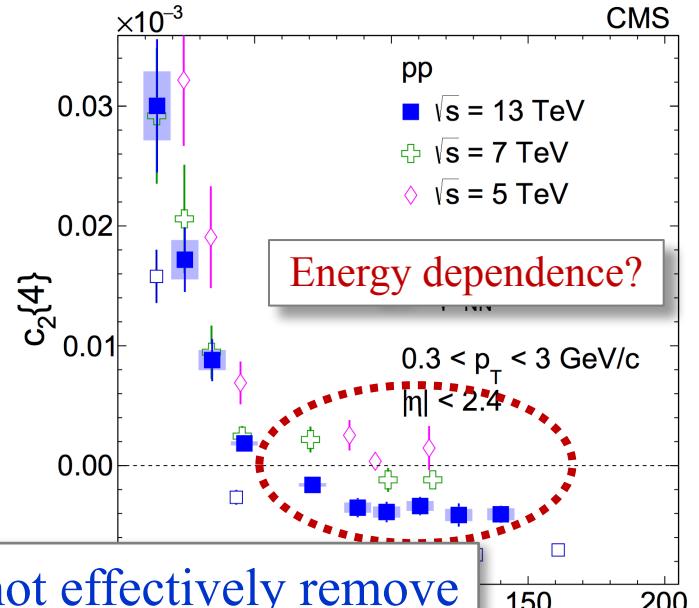
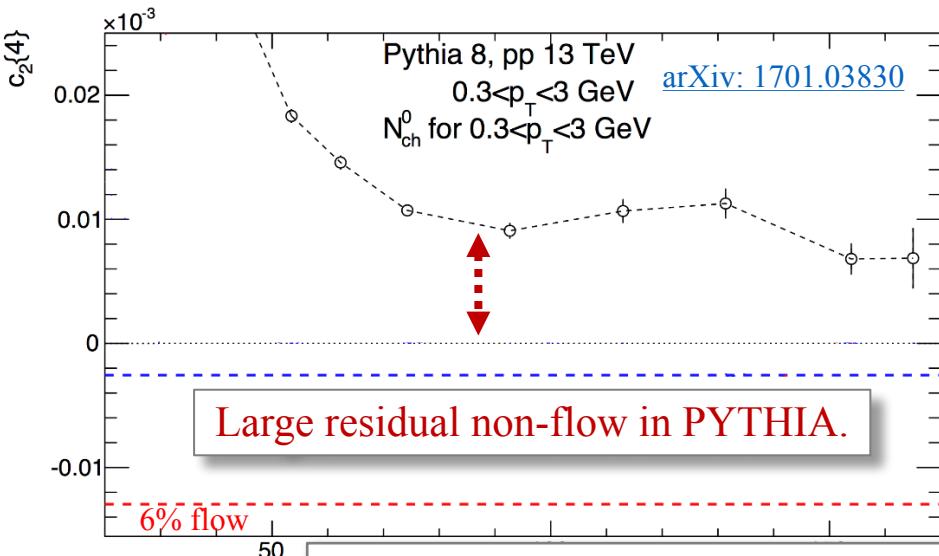
- Cumulant method: natural to probe collectivity:

$$C_2\{4\} \equiv \langle 4 \rangle - 2\langle 2 \rangle^2$$

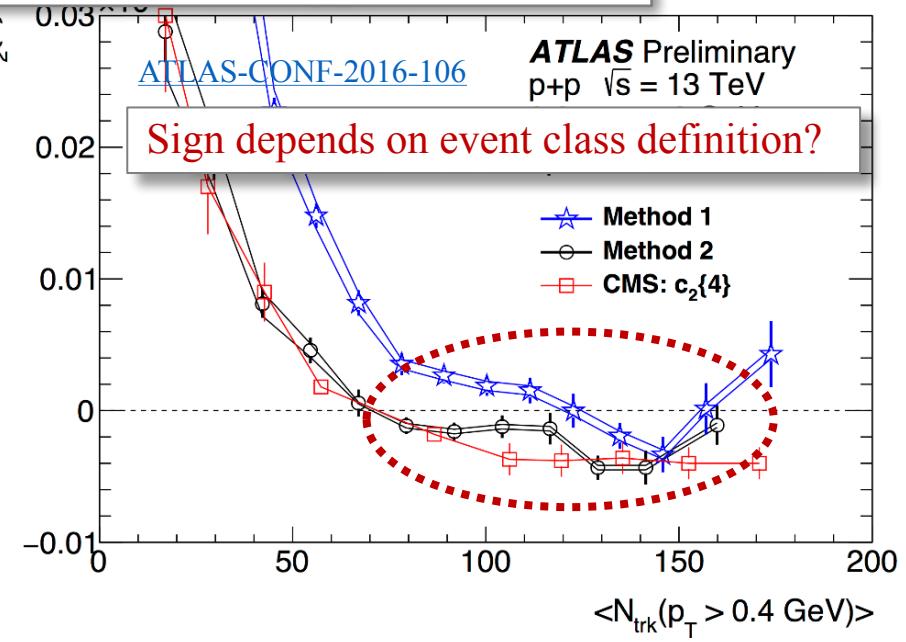
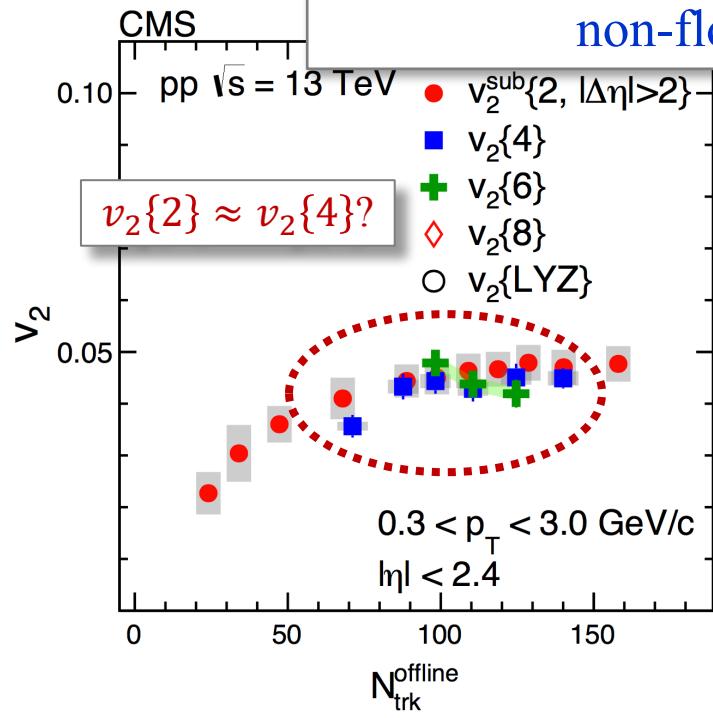


However, does cumulant work in pp ?
 N_{ch} is smaller, non-flow contribution is larger.

Limitations of standard cumulant

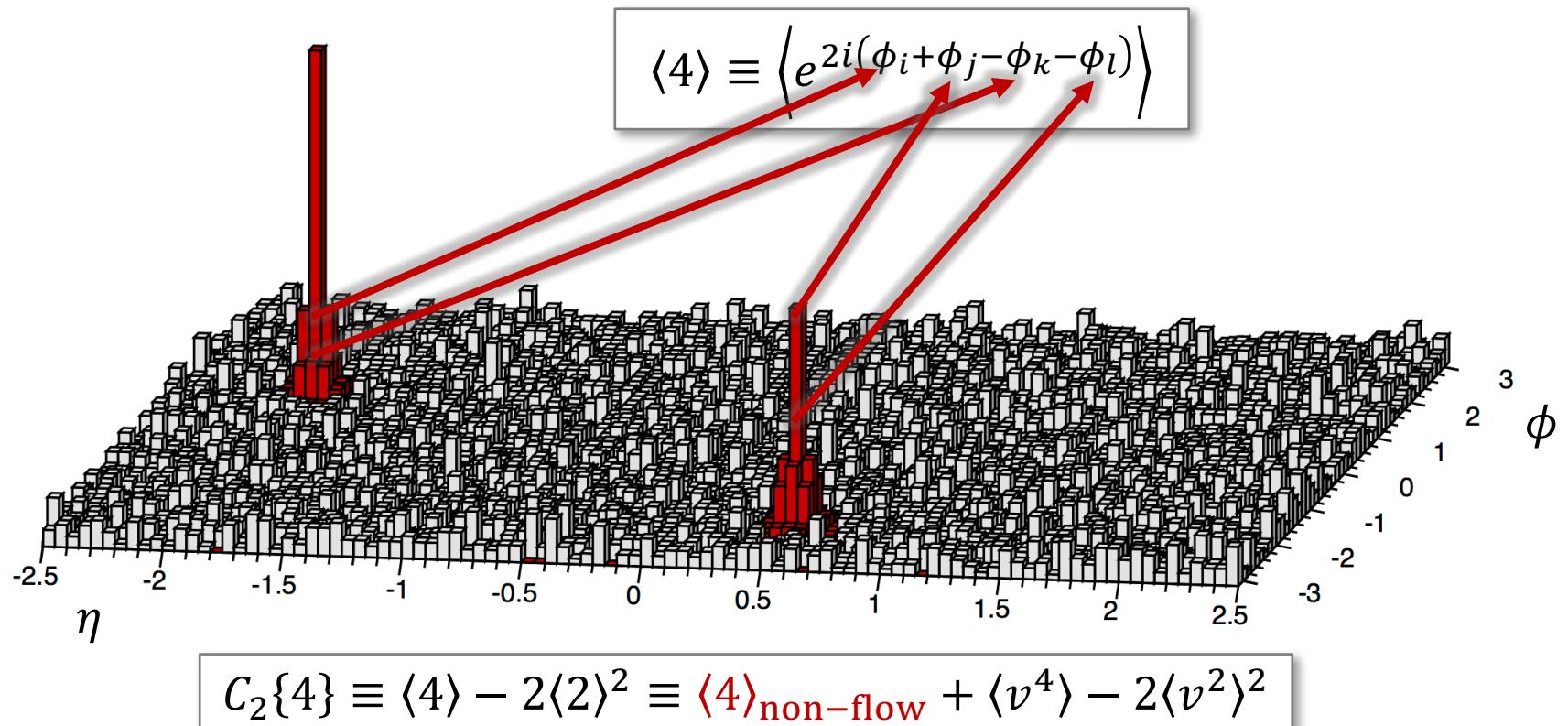


Standard cumulant $c_2\{4\}$ cannot effectively remove non-flow contribution in pp !



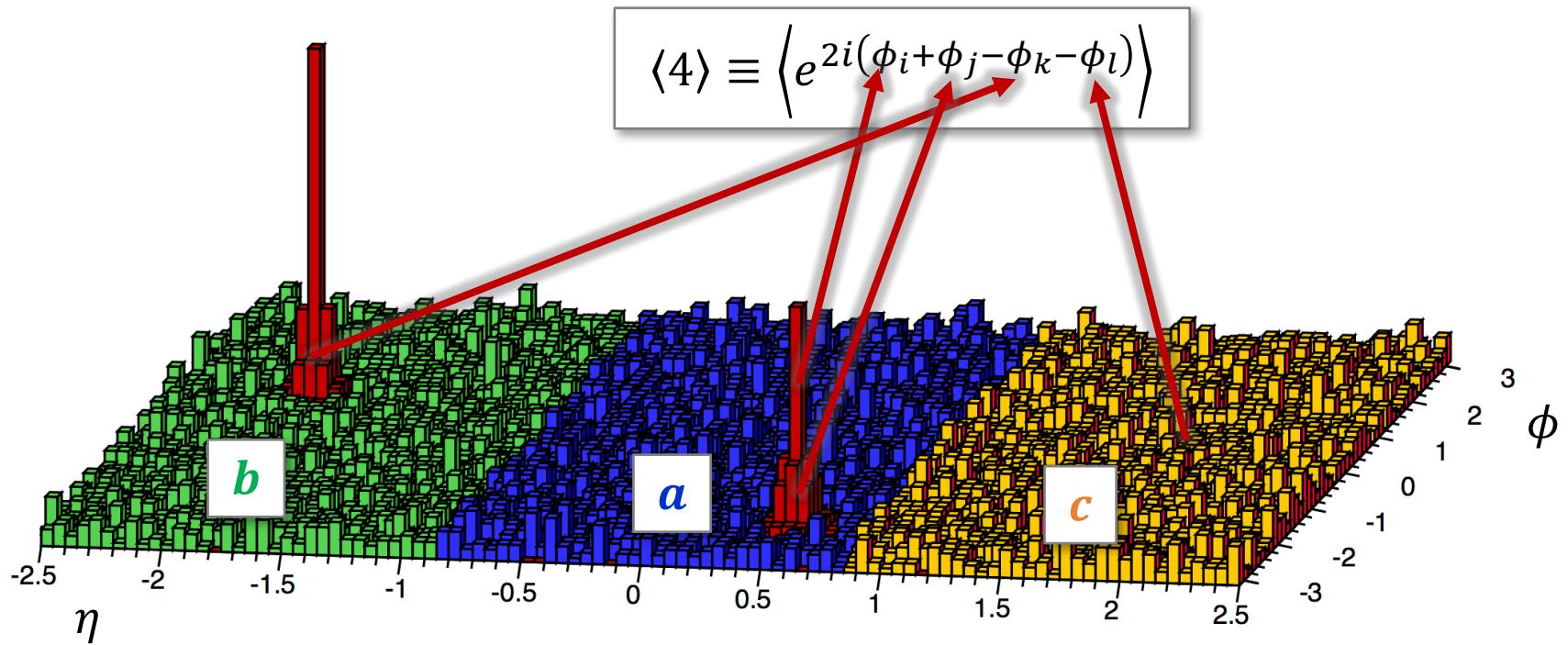
Subevent cumulant

- In standard cumulant, non-flow sources contribute to four-particle correlation $\langle 4 \rangle$;



Subevent cumulant

- In standard cumulant, non-flow sources contribute to four-particle correlation $\langle 4 \rangle$;
- In the subevent method, particles are correlated across all subevents (long-range).
 - 3 subevent cumulant can further suppress away-side jet contribution;
- New method validated in PYTHIA, but is there a data-driven way to check the residual non-flow in data?



$$C_2^{2a|b,c} \{4\} \equiv \langle 4 \rangle_{2a|b,c} - 2\langle 2 \rangle_{a|b}\langle 2 \rangle_{a|c}$$

Test of residual non-flow

$$C_2\{4\} \equiv$$

\langle nonflow + flow \rangle_{evt}



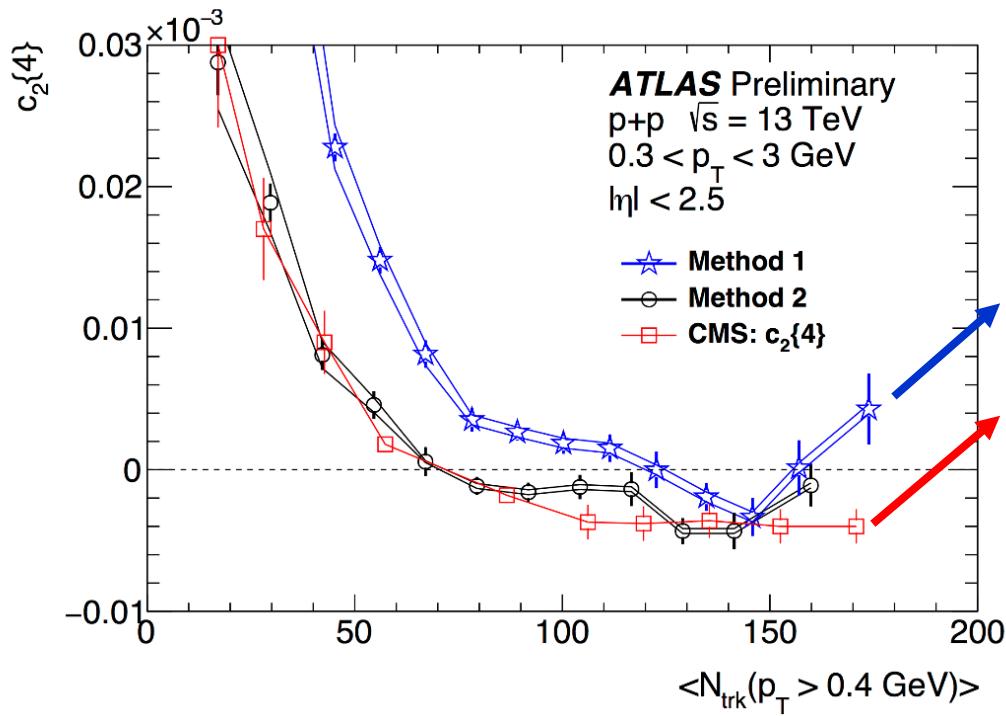
\langle nonflow + flow \rangle_{evt}



Non-flow changes greatly EbyE

Flow changes little EbyE

non-flow fluc. ← multiplicity fluc. ← how $\langle \dots \rangle_{evt}$ is defined: N_{ch}^{Sel}



- N_{ch}^{Sel} defined with different p_T : very different non-flow fluctuation.

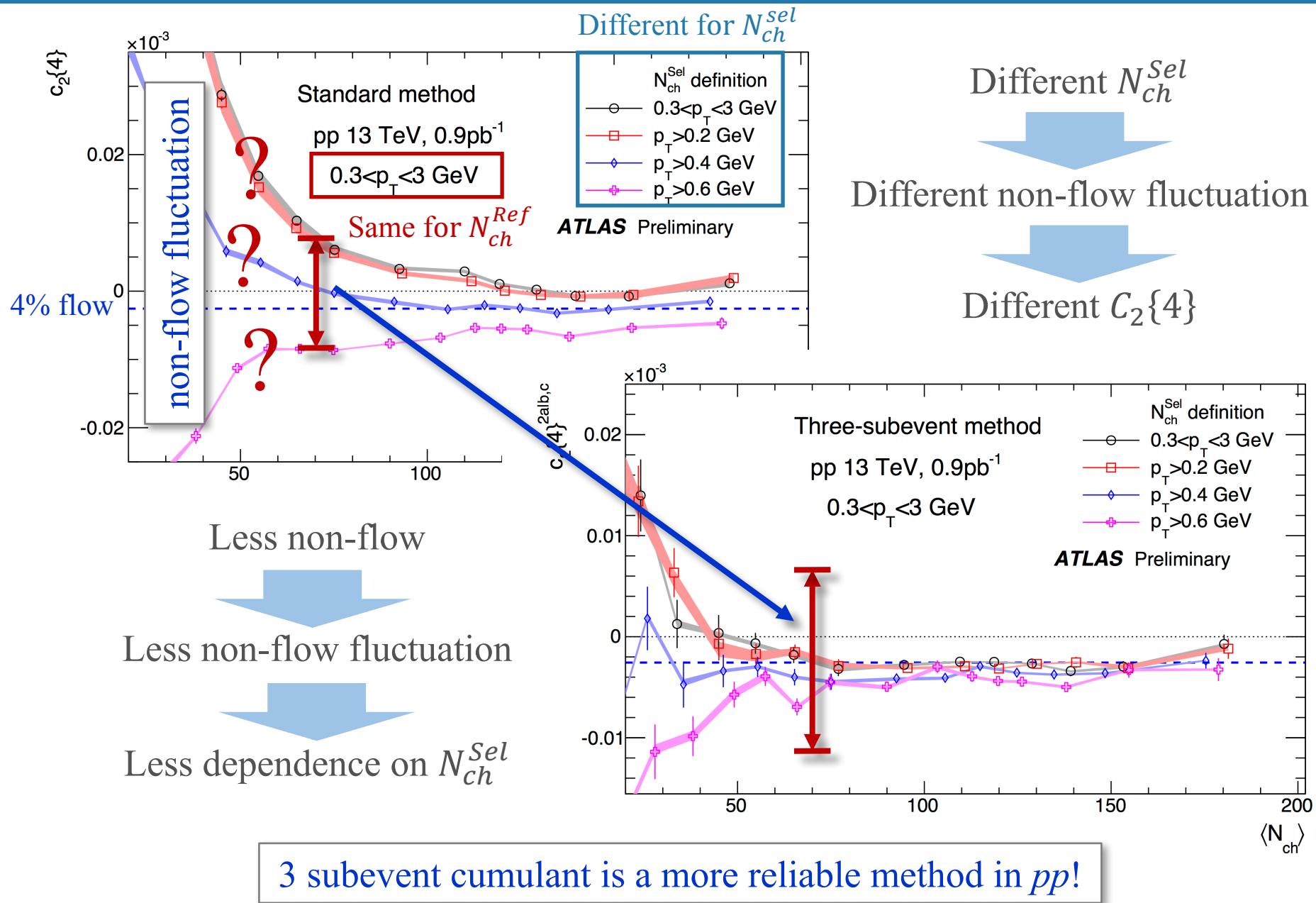
N_{ch}^{Sel} defined with $0.3 < p_T < 3.0 \text{ GeV}$

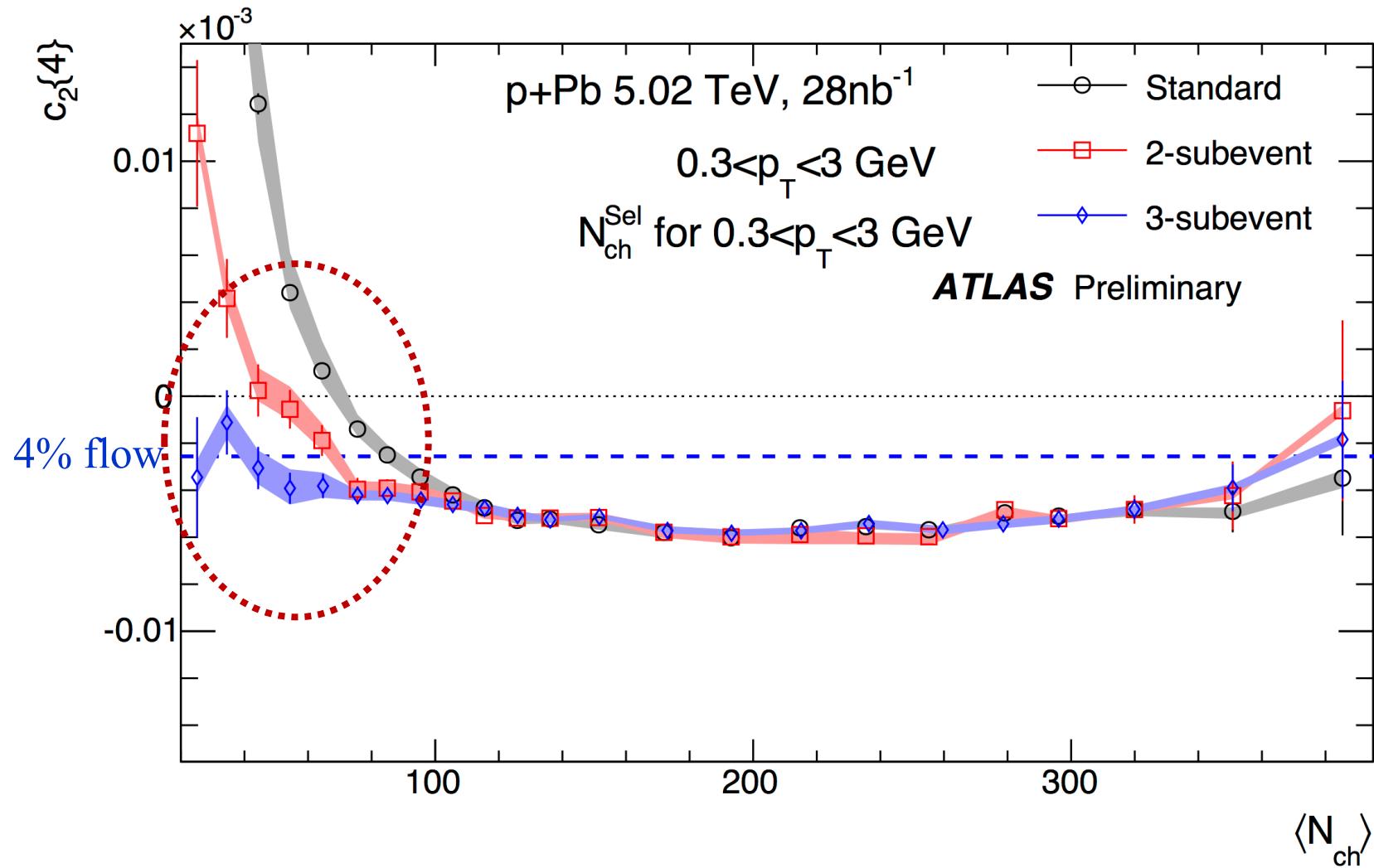
N_{ch}^{Sel} defined with $p_T > 0.4 \text{ GeV}$

- Non-flow fluctuation might mimic the flow signal (negative $C_2\{4\}$)!

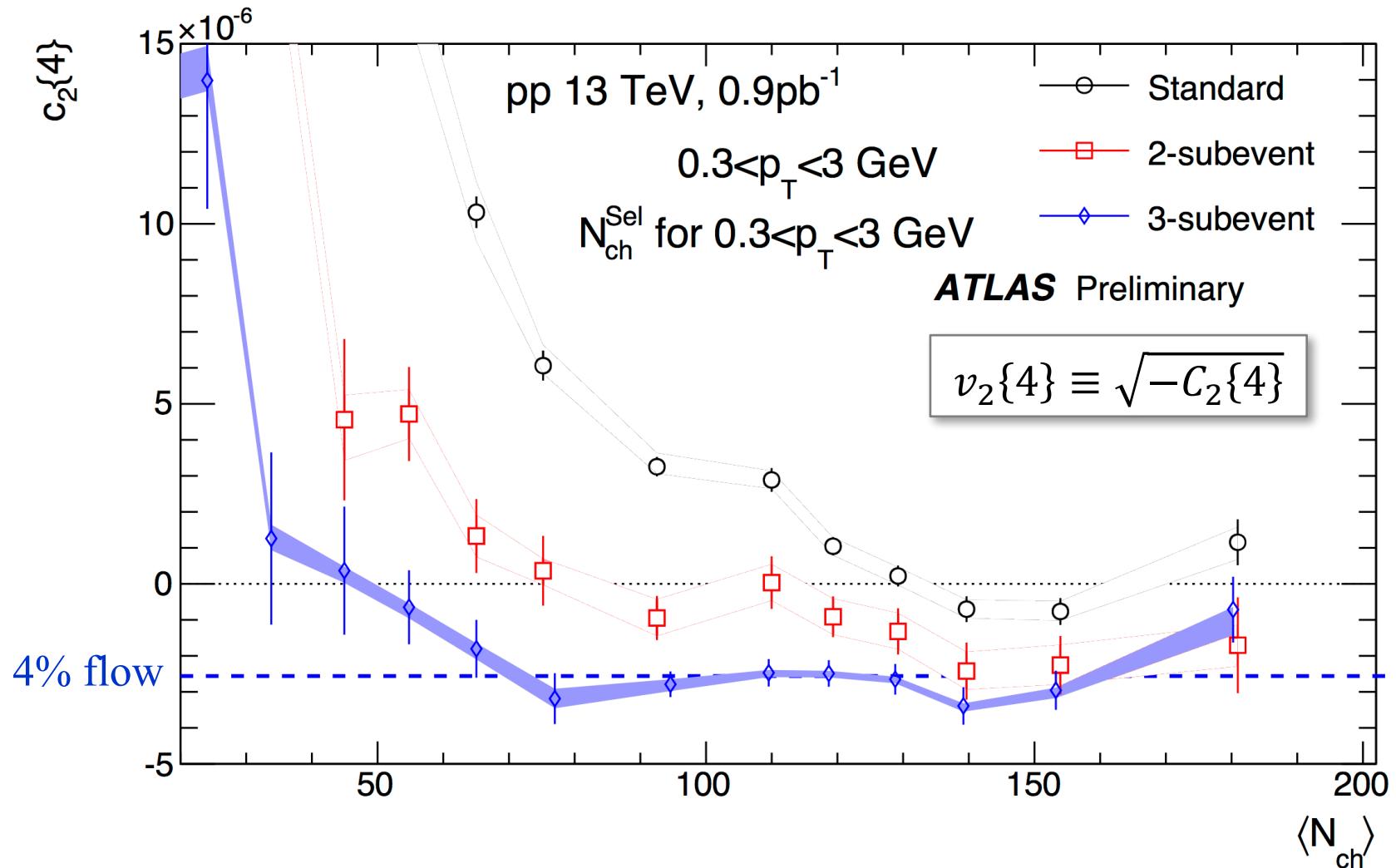
Standard cumulant v.s. Subevent cumulant

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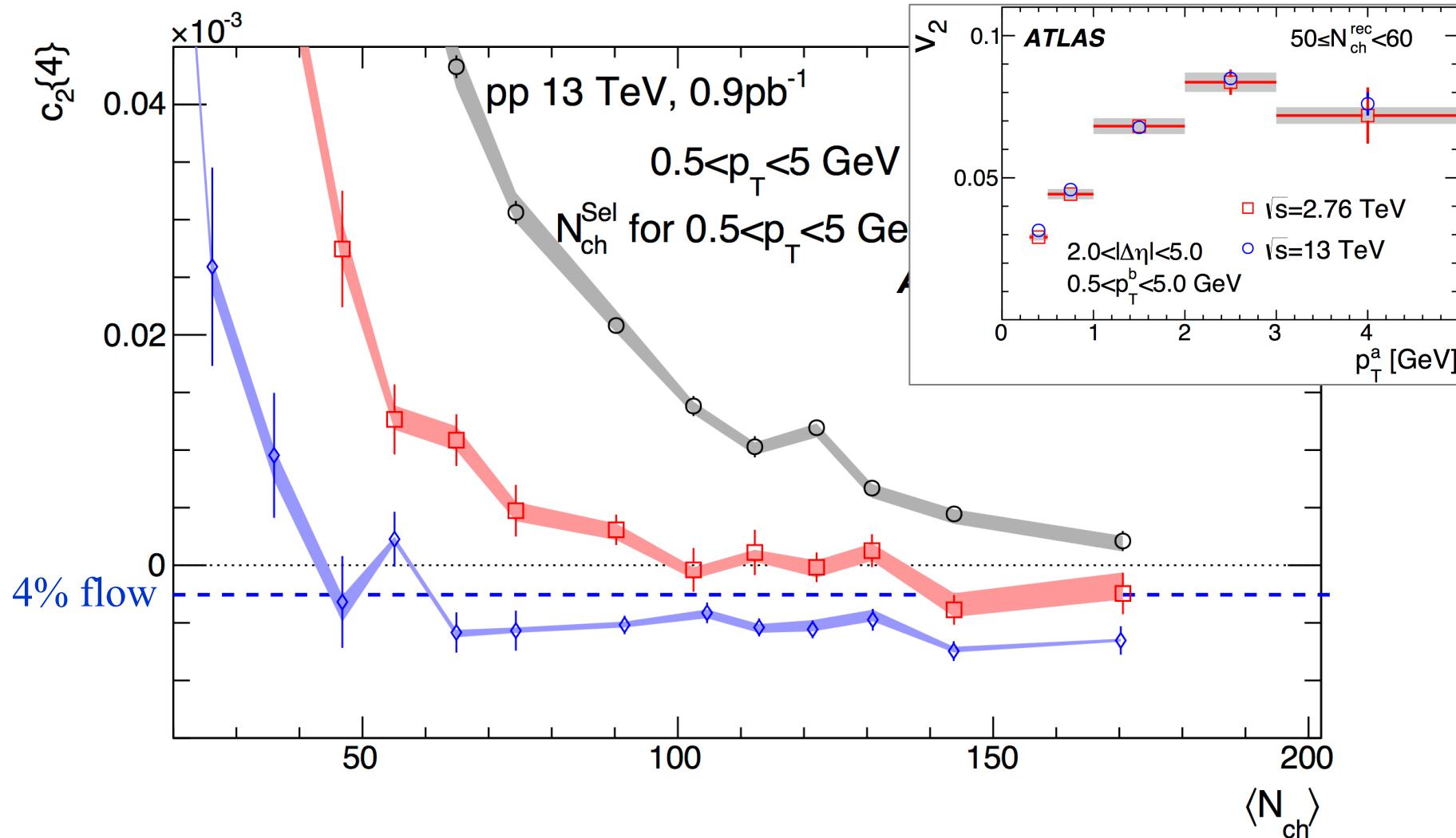
- Consistent at large N_{ch} : non-flow is smaller;
- Split observed at low N_{ch} : suppression of non-flow;



- Standard cumulant has positive $C_2\{4\}$: large residual non-flow;
- 2 subevent cumulant already suppresses non-flow;
- 3 subevent cumulant measures 4% flow down to 70 tracks.

13 TeV pp : higher p_T region

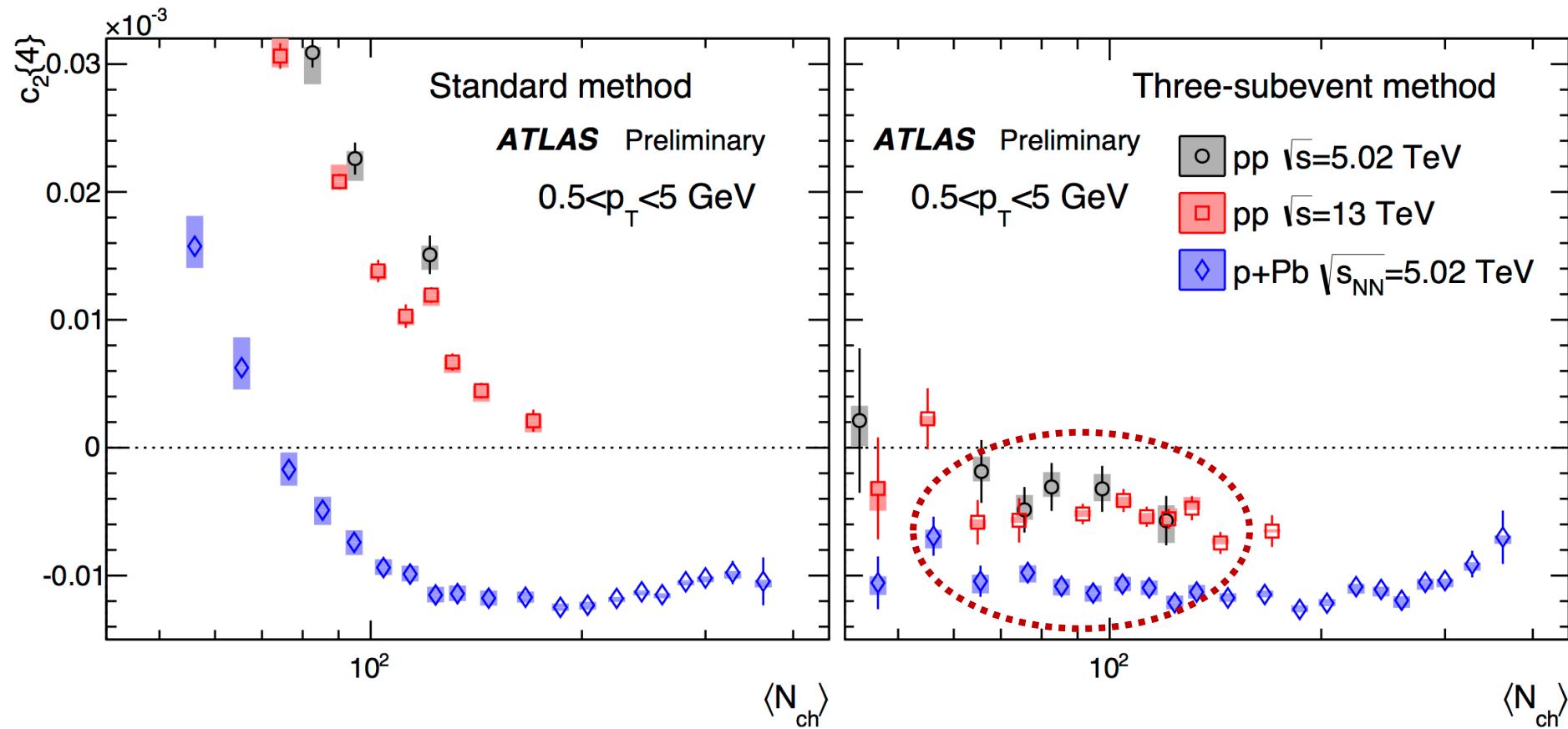
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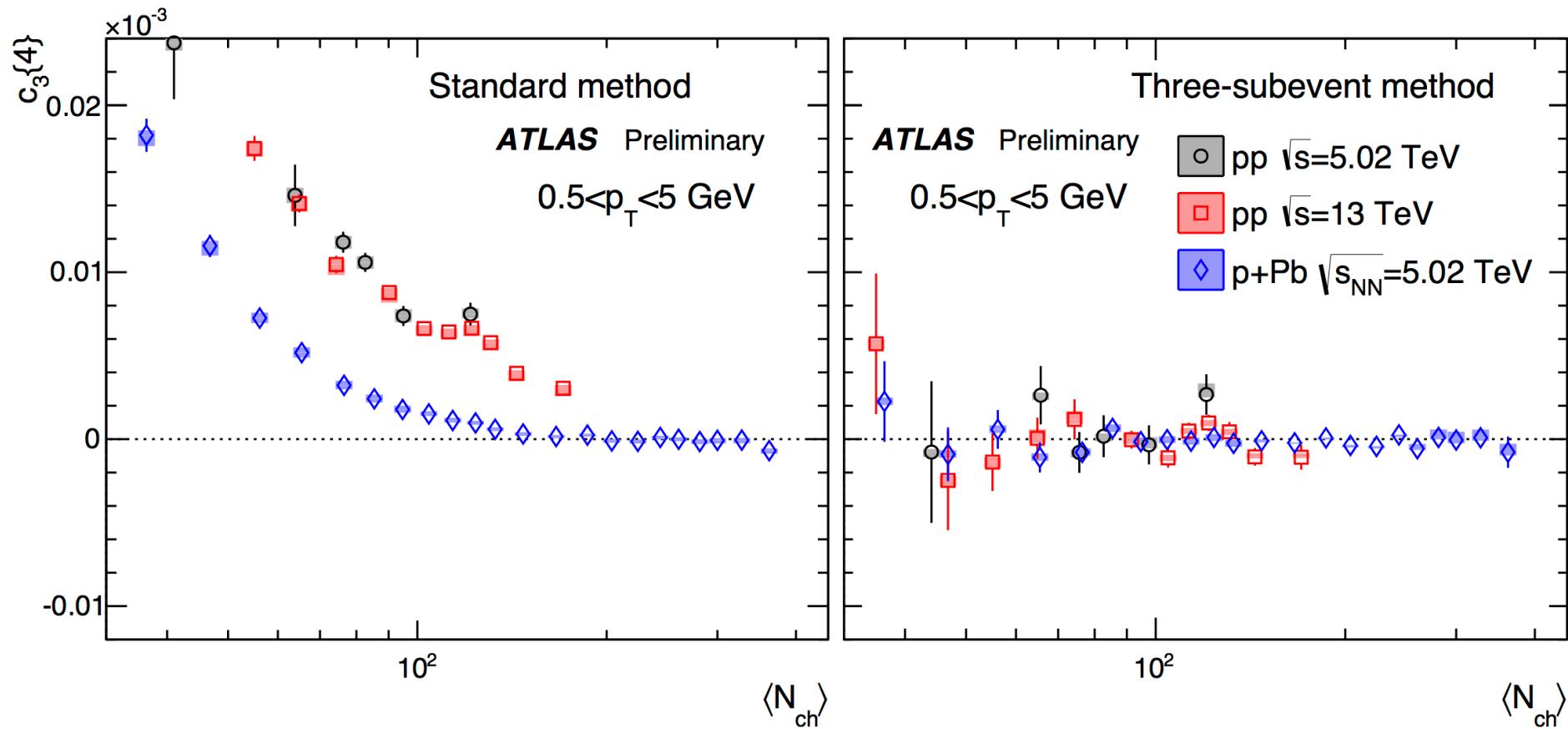
- Higher p_T range: higher fraction of non-flow;
- ONLY 3 subevent method gives larger flow signal;

Comparison among three collision systems

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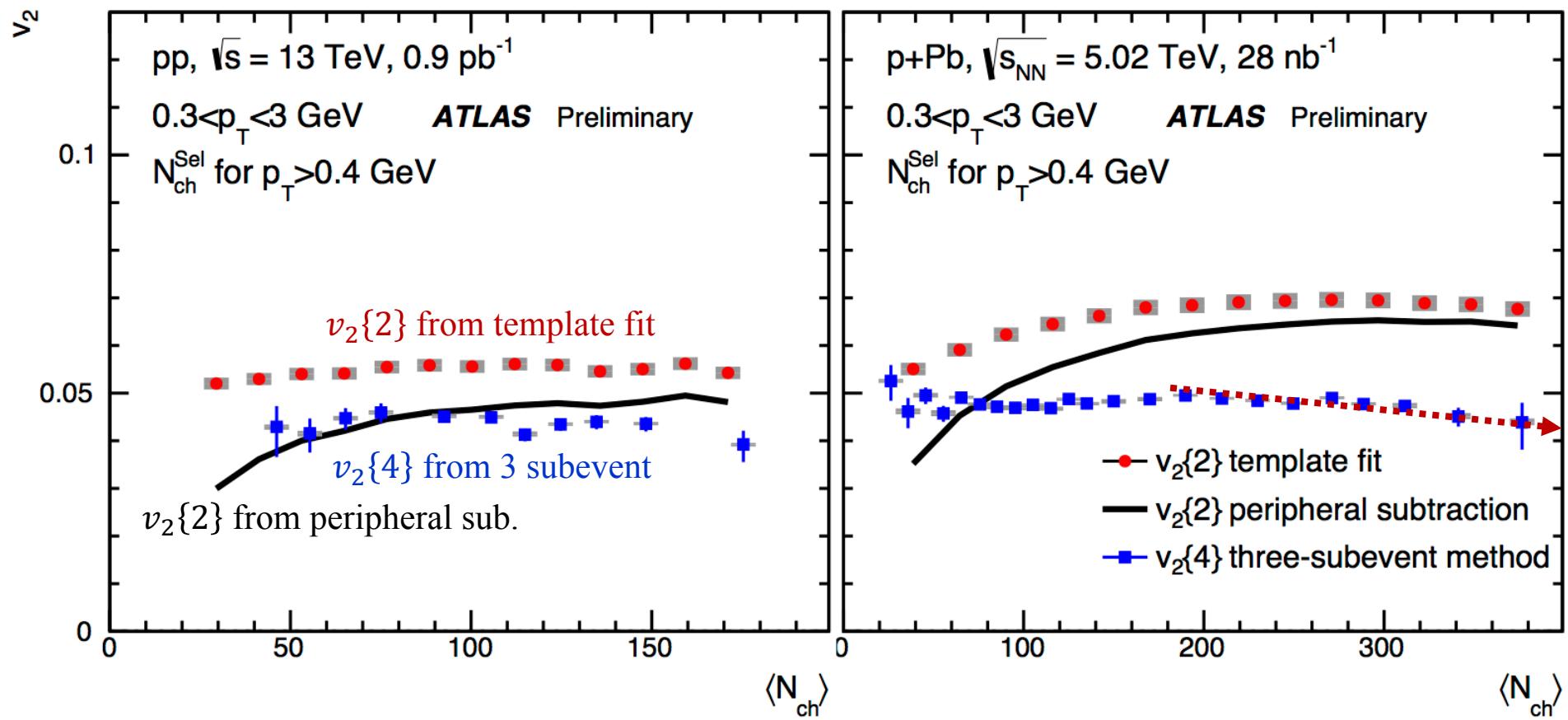
- pp results from standard cumulant unable to compare with $p+\text{Pb}$;
- With 3 subevent, negative $C_2\{4\}$ observed in 5.02 TeV pp ;
- Weak energy dependence in pp : more data for 5.02 TeV pp in 2017?
- $p+\text{Pb}$ has larger flow than pp in the comparable N_{ch} region;



- $C_3\{4\}$ still has residual non-flow in standard method;
- $C_3\{4\}$ from 3 subevent is consistent with 0:
 - $\bar{\nu}_3 \ll \bar{\nu}_2$
 - Fluctuation kills $C_3\{4\}$?

Comparison with 2-particle correlation method

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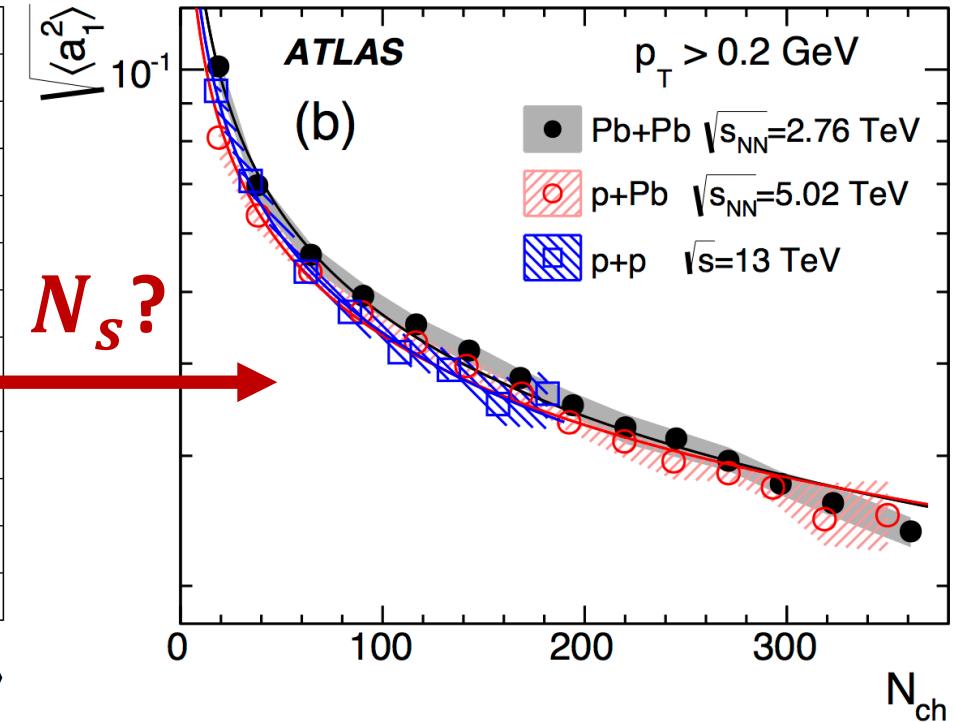
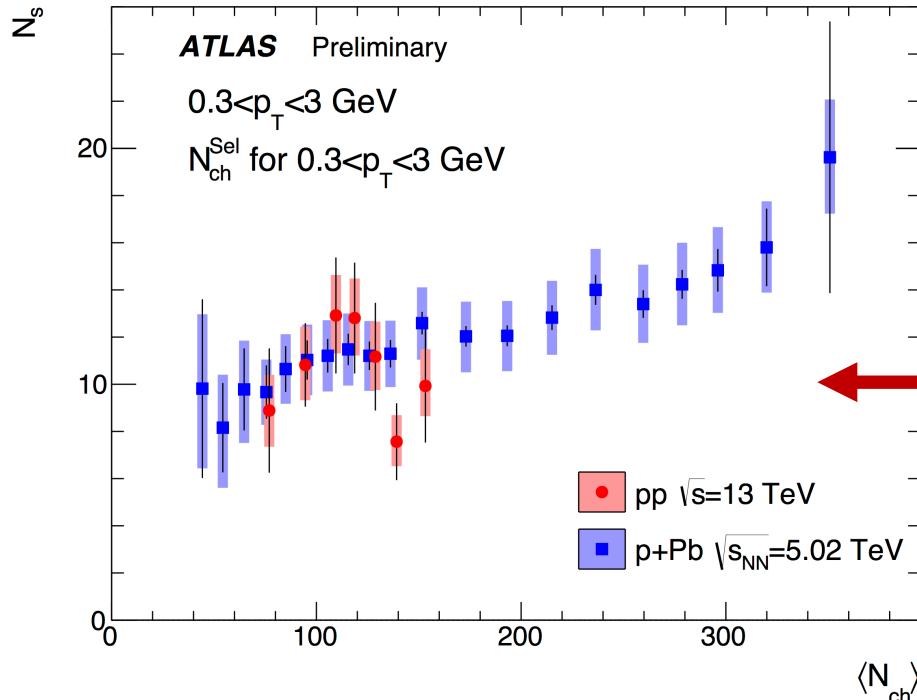


- $v_2\{4\} < v_2\{2\}$ (template fit): flow fluctuation;
- $v_2\{4\} \approx v_2\{2\}$ (peripheral subtraction): underestimation of $v_2\{2\}$;
- Decreasing trend of $C_2\{4\}$ for central $p+\text{Pb}$:
 - Many sources in central $p+\text{Pb}$?

- $\nu_2\{2\} \neq \nu_2\{4\}$: EbyE flow fluctuations associated with fluctuating initial conditions. [PRL 112, 082301 \(2014\)](#)
- Fluctuation can be quantified to the number of sources N_s in the initial stage:

$$\frac{\nu_2\{4\}}{\nu_2\{2\}} = \left(\frac{4}{3 + N_s} \right)^{1/4}$$

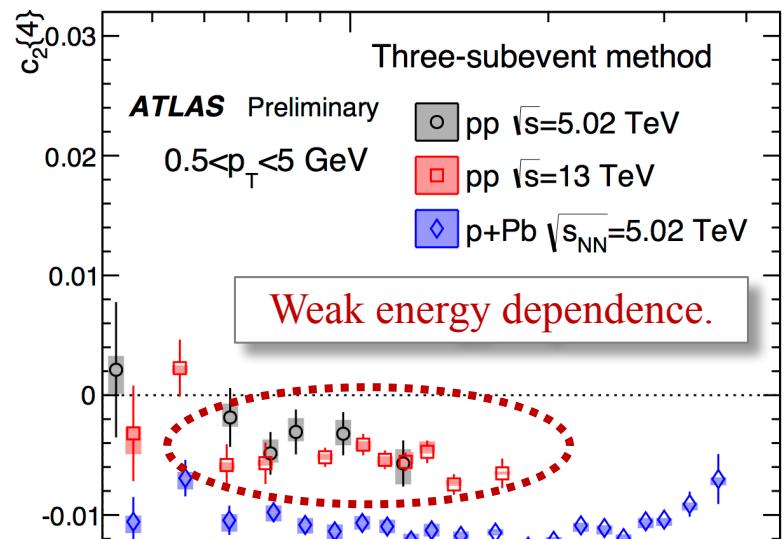
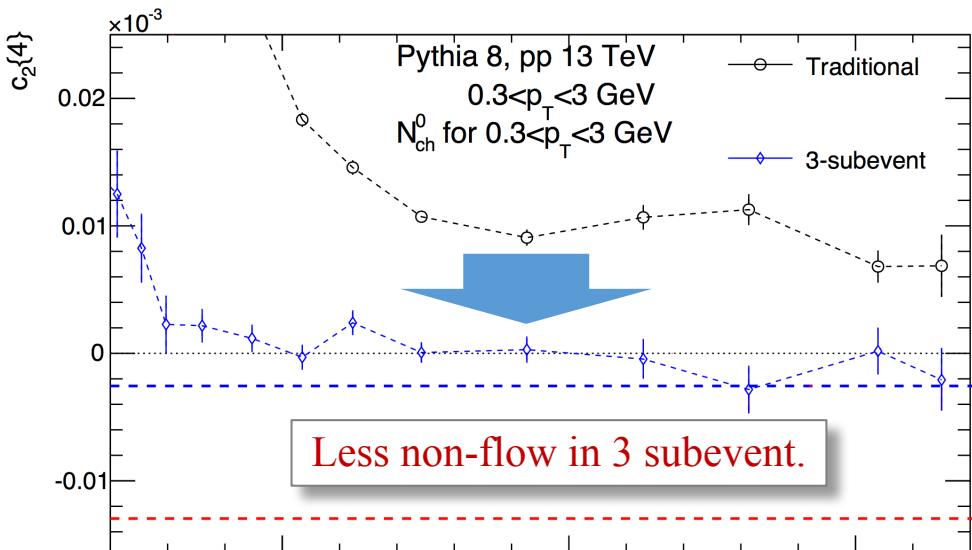
$$\frac{N(\eta)}{\langle N(\eta) \rangle} \approx 1 + a_1 \eta, \quad a_1 \propto \frac{1}{\sqrt{N_s}}$$



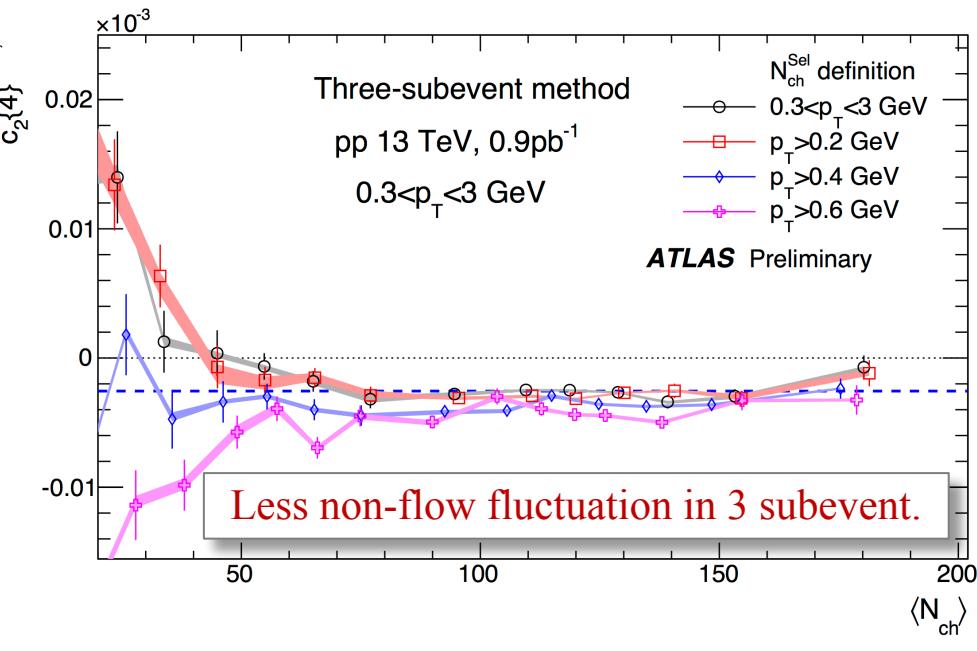
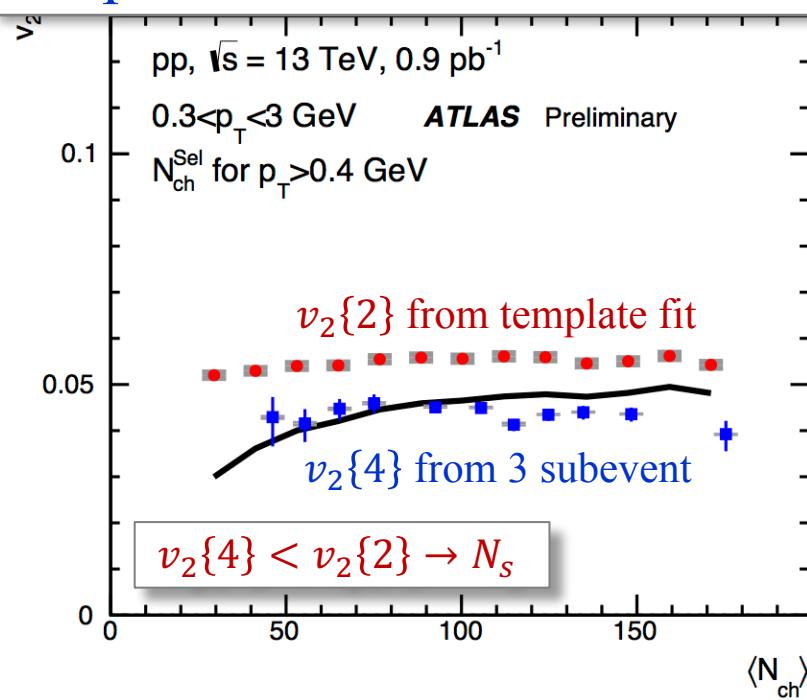
- N_s for $p+Pb$ goes up to 20 at high multiplicity;
- N_s for pp approximately consistent with $p+Pb$ at comparable $\langle N_{ch} \rangle$ value.

Summary

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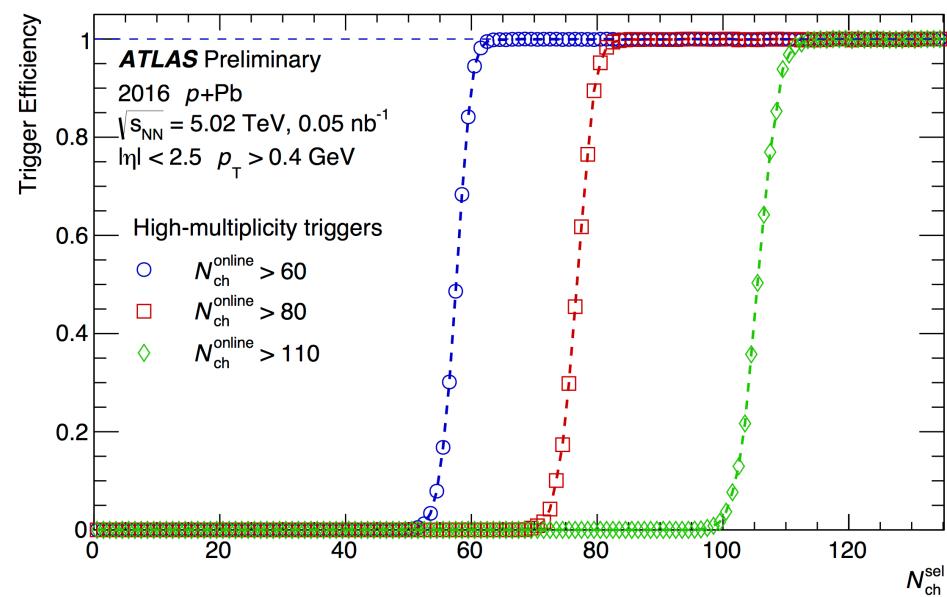
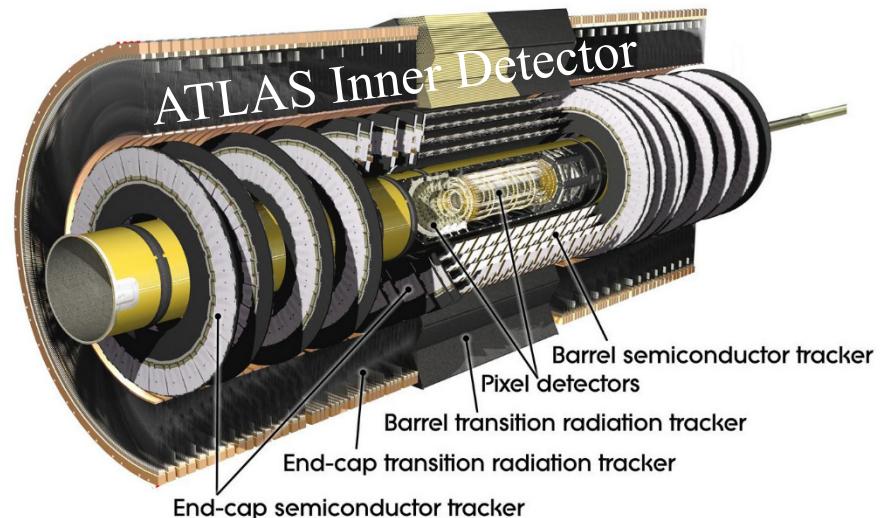
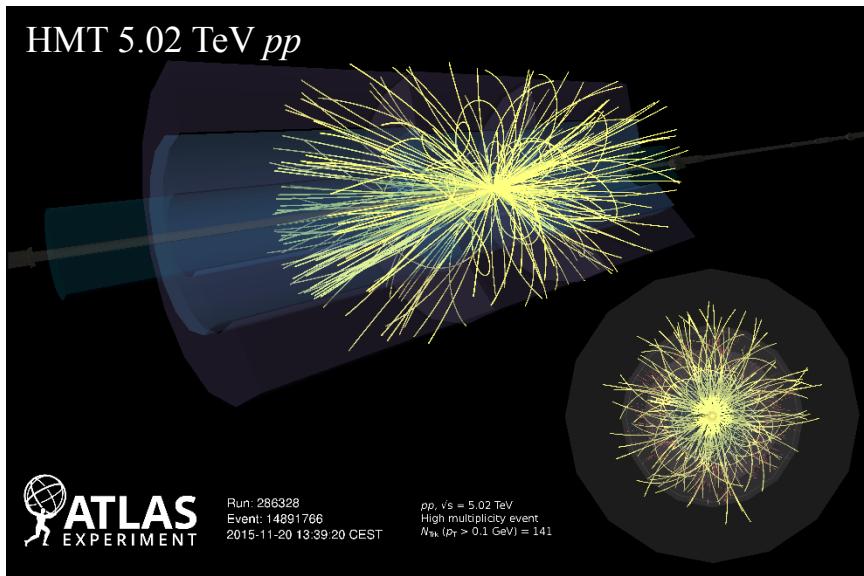
Hope the new results can better constrain models in small systems.



Back Up

Data sets

- Studied in this analysis:
 - Energy dependence of $C_2\{4\}$ in pp ;
 - Comparison between pp and $p+\text{Pb}$;
- 3 collision systems included:
 - 13 TeV pp (2015+2016)
 - 5.02 TeV pp (2015)
 - 5.02 TeV $p+\text{Pb}$ (2013+2016)
- Two p_T ranges: $0.3 < p_T < 3.0 \text{ GeV}$ and $0.5 < p_T < 5.0 \text{ GeV}$;
- High-multiplicity track triggers developed to enhance statistics at large N_{ch} :



Formulas without particle weight

$$C_n\{4\} \equiv \langle 4 \rangle - 2 \langle 2 \rangle^2$$

$$\langle 2 \rangle \equiv \frac{Q_n^2 - M}{M(M-1)}$$

$$\langle 4 \rangle \equiv \frac{Q_n^4 + Q_{2n}^2 - 2Q_{2n}Q_n^*Q_n^* - 4(M-2)Q_n^2 + 2M(M-3)}{M(M-1)(M-2)(M-3)}$$

$$C_n^{a,a|b,b}\{4\} \equiv \langle 4 \rangle_{a,a|b,b} - 2 \langle 2 \rangle_{a|b}^2$$

$$\langle 2 \rangle_{a|b} \equiv \frac{Q_{n,a}Q_{n,b}}{M_aM_b}$$

$$\langle 4 \rangle_{a,a|b,b} \equiv \frac{(Q_{n,a}^2 - Q_{2n,a})(Q_{n,b}^2 - Q_{2n,b})^*}{M_a(M_a-1)M_b(M_b-1)}$$

$$C_n^{a,a|b,c}\{4\} \equiv \langle 4 \rangle_{a,a|b,c} - 2 \langle 2 \rangle_{a|b} \langle 2 \rangle_{a|c}$$

$$\langle 2 \rangle_{a|b} \equiv \frac{Q_{n,a}Q_{n,b}}{M_aM_b}$$

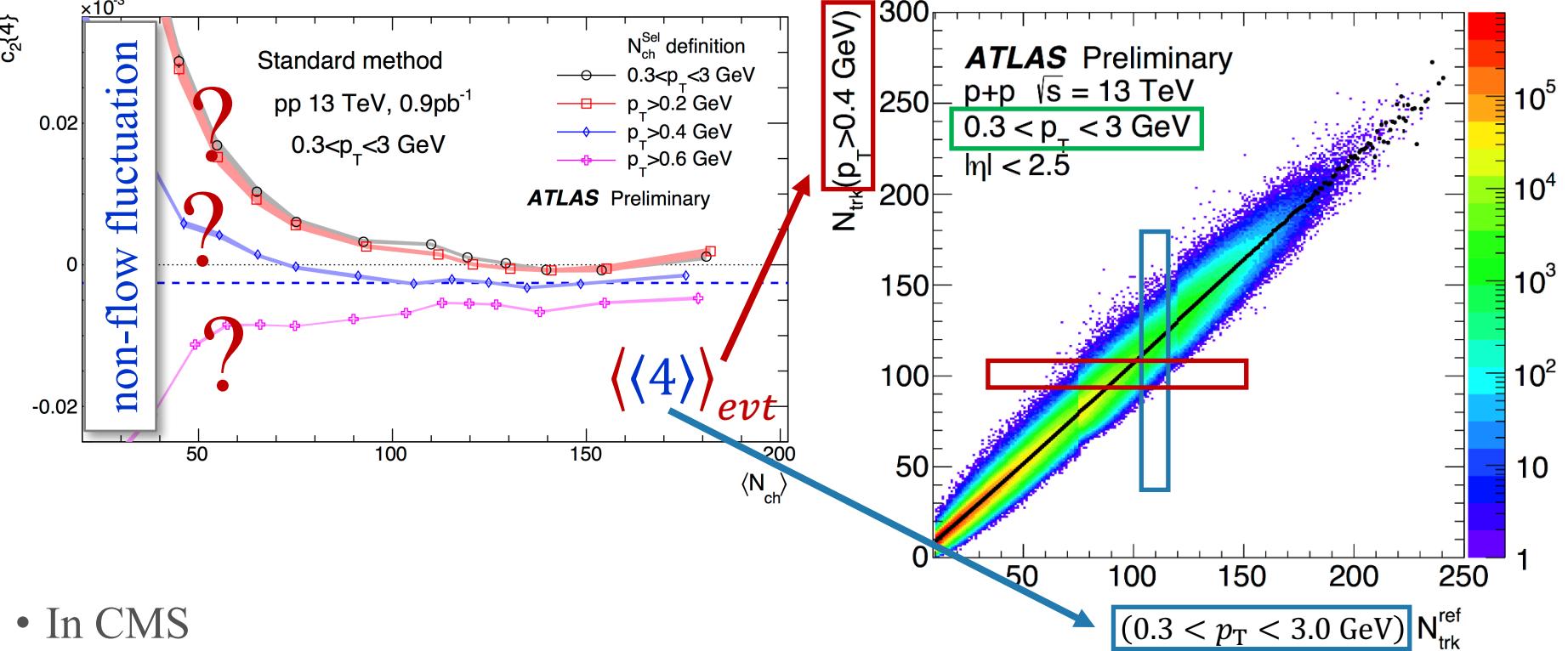
$$\langle 2 \rangle_{a|c} \equiv \frac{Q_{n,a}Q_{n,c}}{M_aM_c}$$

$$\langle 4 \rangle_{a,a|b,c} \equiv \frac{(Q_{n,a}^2 - Q_{2n,a})Q_{n,b}^*Q_{n,c}^*}{M_a(M_a-1)M_bM_c}$$



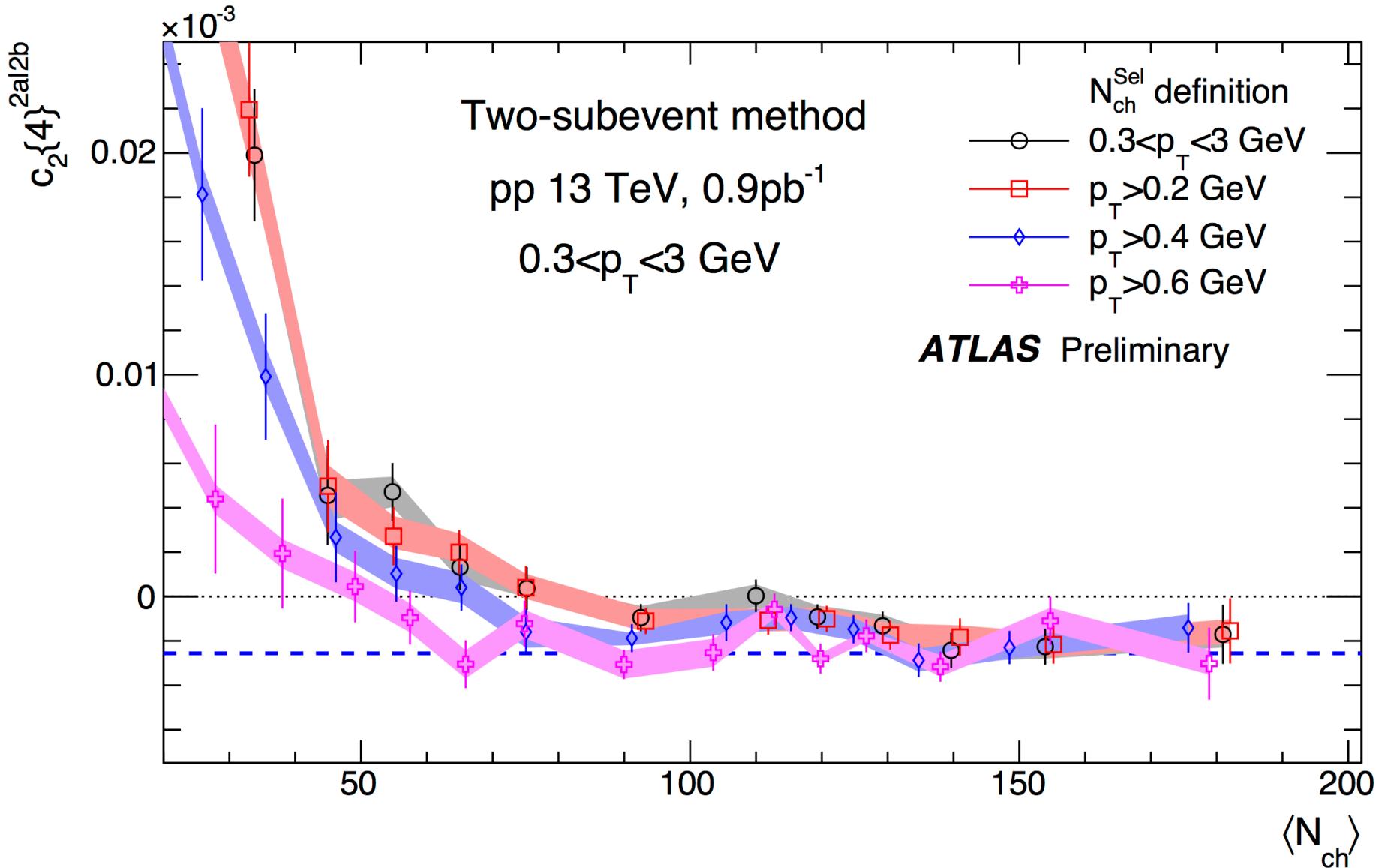
$$= \left\langle e^{in(\phi_1^a + \phi_2^b - \phi_3^a - \phi_4^c)} \right\rangle_c + \langle \mathbf{v}_{n,a}^2 \mathbf{v}_{n,b}^* \mathbf{v}_{n,c}^* \rangle - \langle v_{n,a}^2 \rangle \langle \mathbf{v}_{n,b} \mathbf{v}_{n,c}^* \rangle - \langle \mathbf{v}_{n,a} \mathbf{v}_{n,b}^* \rangle \langle \mathbf{v}_{n,a} \mathbf{v}_{n,c}^* \rangle$$

Multiplicity fluctuation

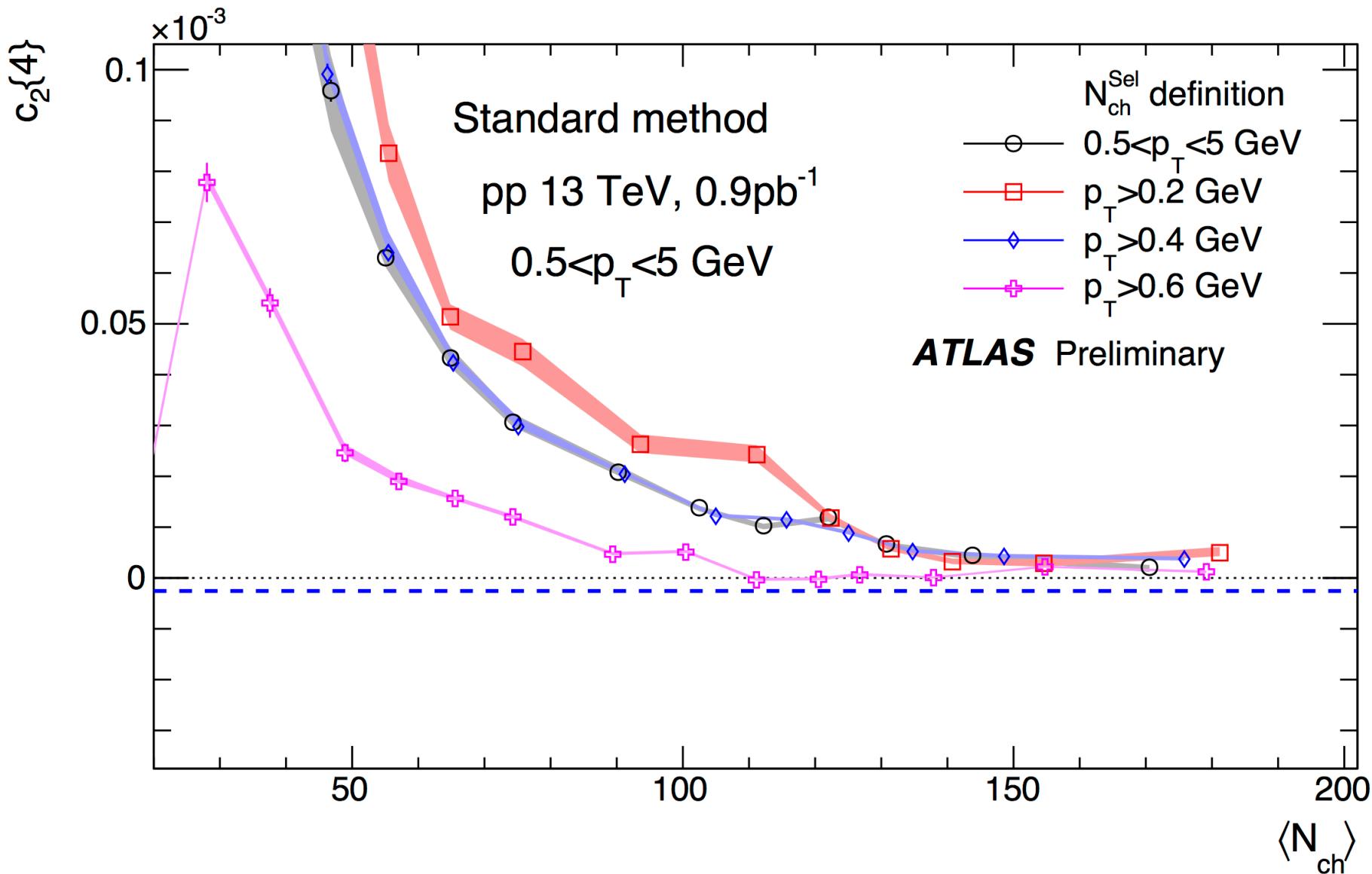


- In CMS
 - $\langle 4 \rangle$ is calculated in $0.3 < p_T < 3.0 \text{ GeV}$;
 - $\langle \dots \rangle_{evt}$ is binned with N_{ch}^{Sel} of $p_T > 0.4 \text{ GeV}$;
- In ATLAS
 - $\langle 4 \rangle$ is calculated in $0.3 < p_T < 3.0 \text{ GeV}$;
 - $\langle \dots \rangle_{evt}$ is binned with N_{ch}^{Sel} of $0.3 < p_T < 3.0 \text{ GeV}$;
 - X-axis of $C_2\{4\}$ is projected to $\langle N_{ch} \rangle$ with $p_T > 0.4 \text{ GeV}$;

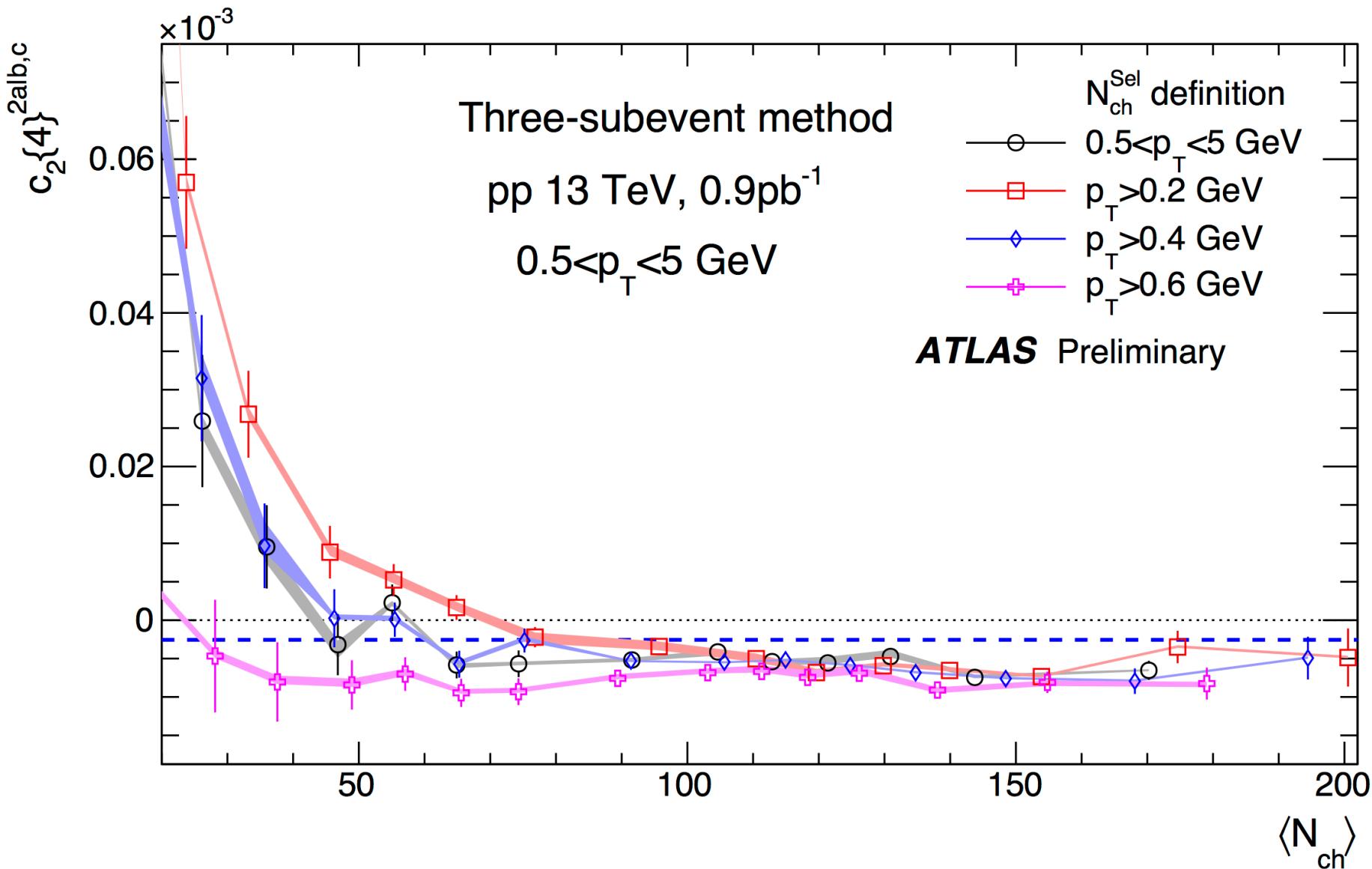
Residual non-flow check in 2 subevent



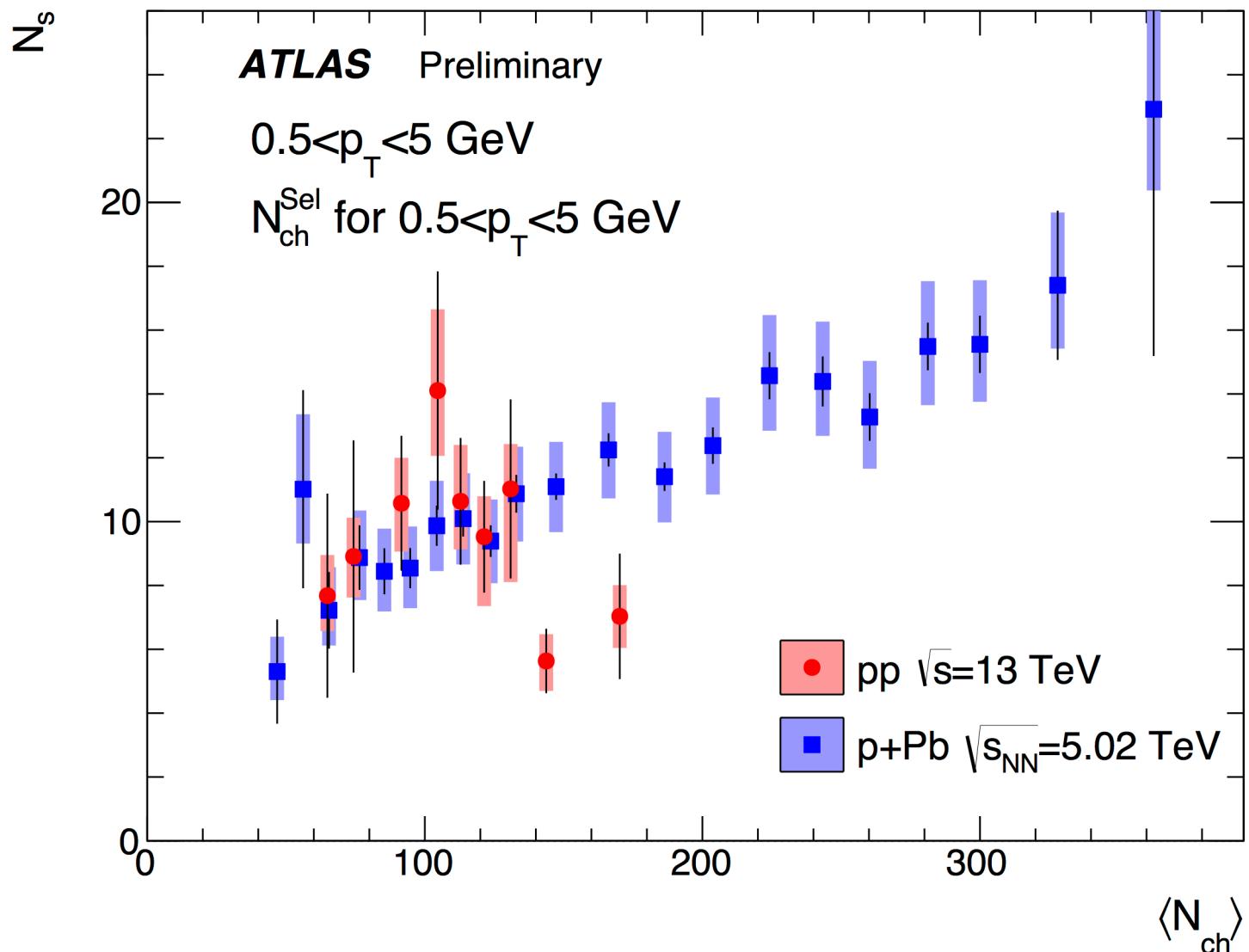
Residual non-flow check in higher p_T



Residual non-flow check in higher p_T

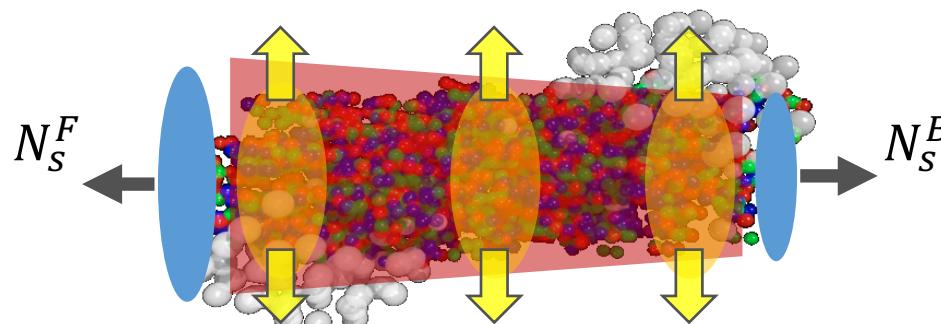
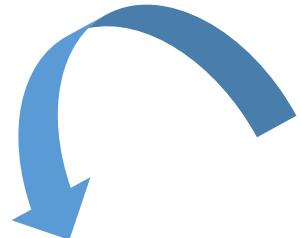


N_s in $0.5 < p_T < 5.0$ GeV

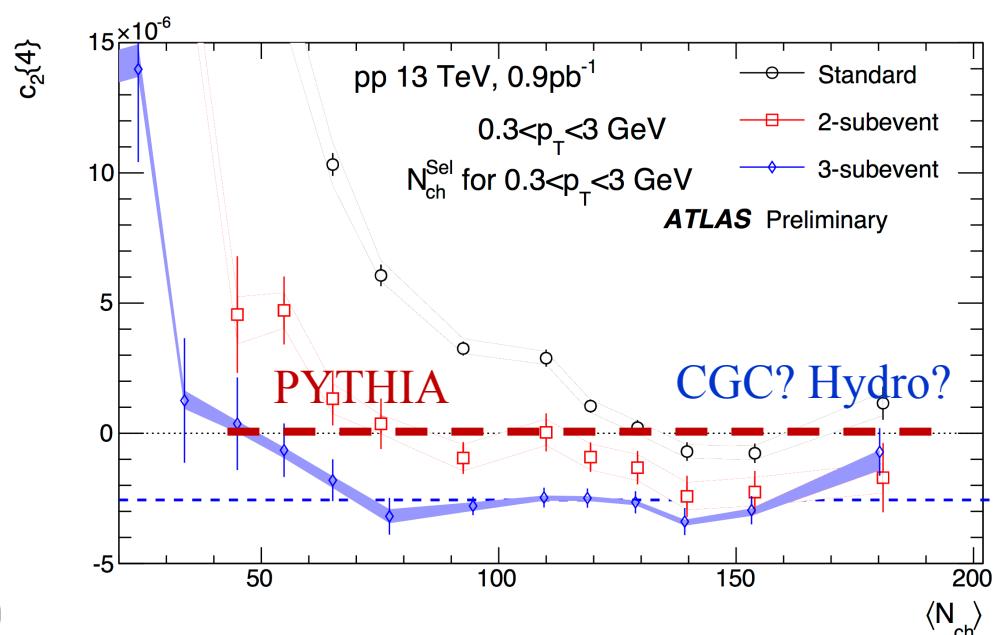
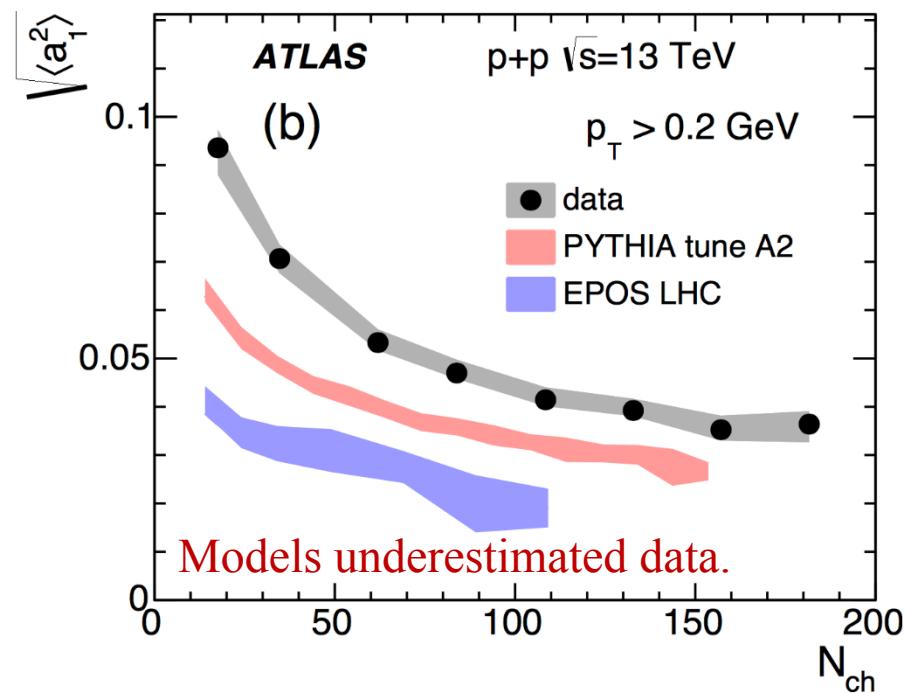
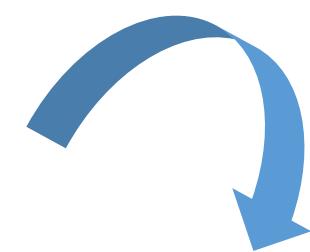


Nature of sources seeding the long-range collective behavior?

Longitudinal correlation



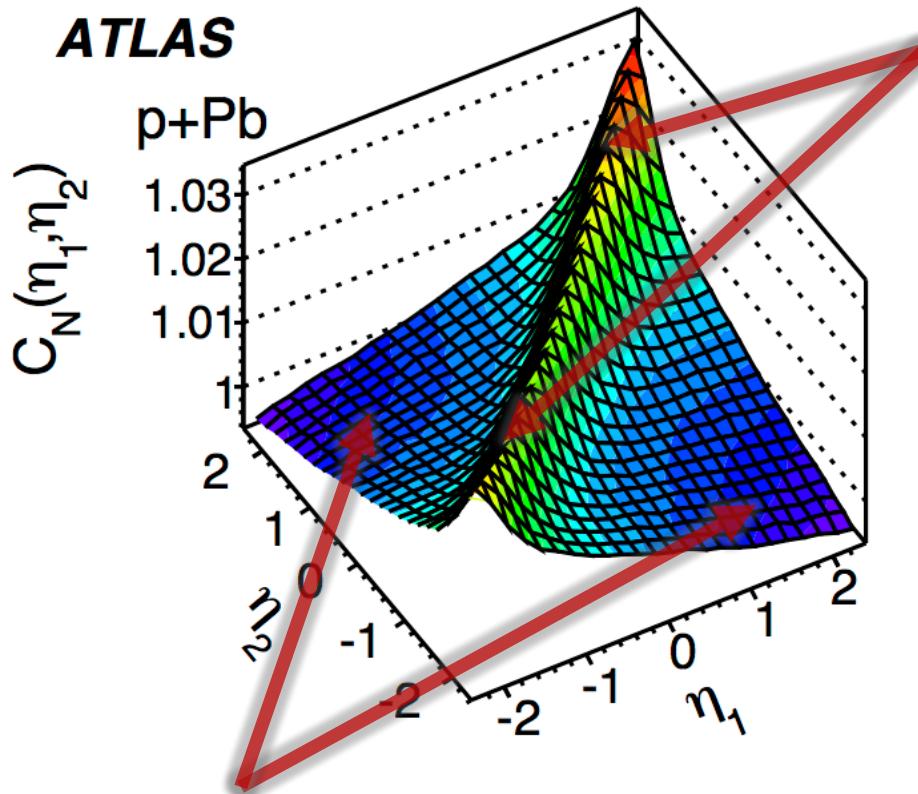
Azimuthal correlation



Physics signal and background

- 2-particle pseudo-rapidity correlation:

$$C(\eta_1, \eta_2) \equiv \frac{\langle N(\eta_1)N(\eta_2) \rangle}{\langle N(\eta_1) \rangle \langle N(\eta_2) \rangle}$$



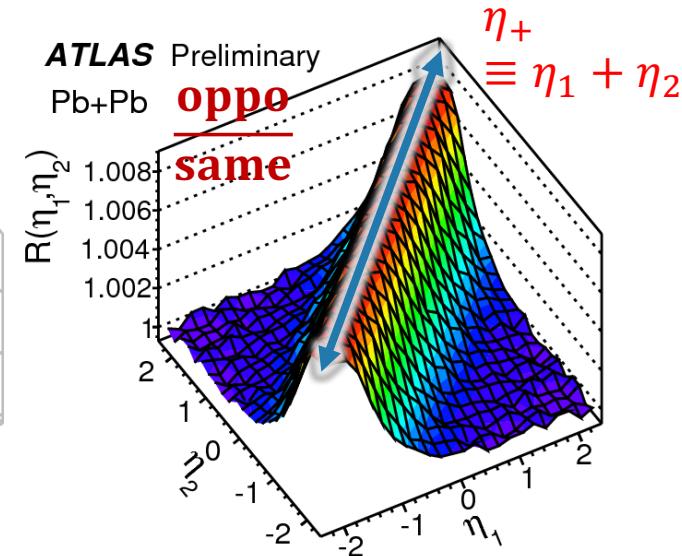
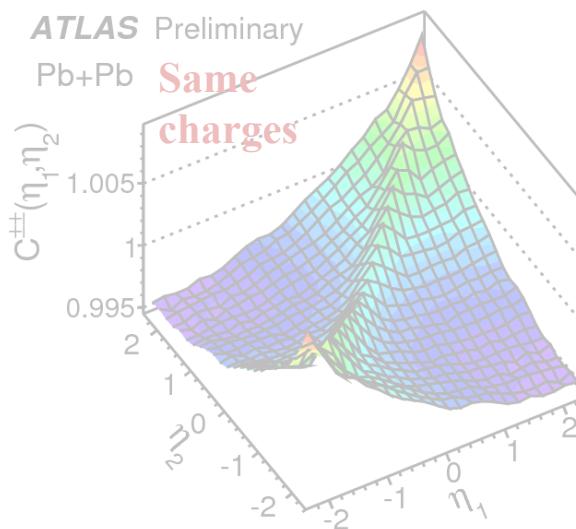
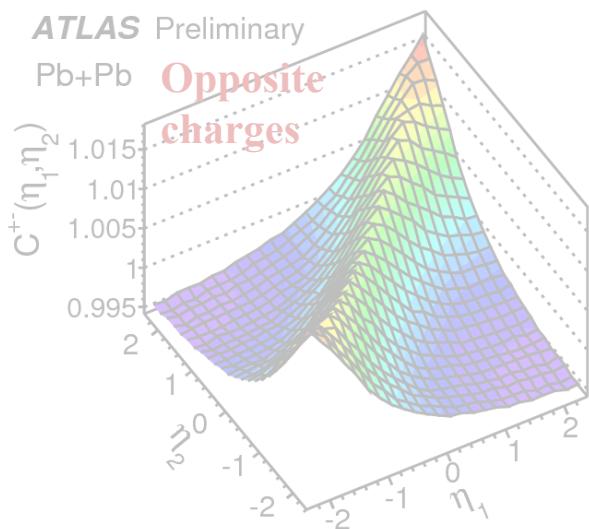
- Long-range correlation (LRC)

- Short-range correlation (SRC) reflects correlation in the same source: jet fragmentation, resonance decay...
- Since SRC has a strong charge dependence, a data-driven way was developed to remove the SRC.
- An alternative way is through multi-particle correlation (cumulant) in pseudo-rapidity. [PRC.93.024903](#)

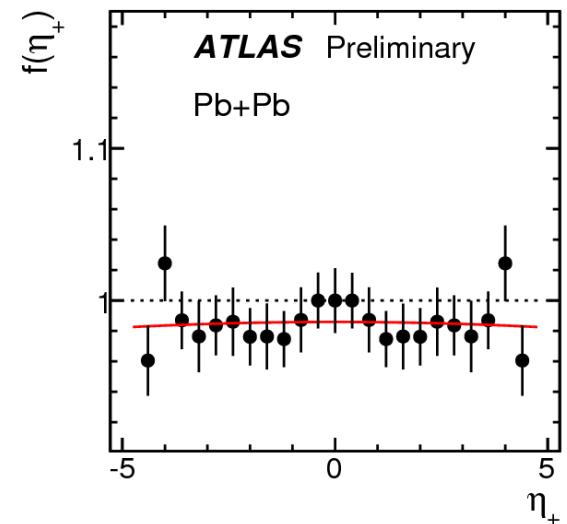
From the experimental perspective, the key is to extract the physics signal from the background.

Charge dependence of $C(\eta_1, \eta_2)$

- Particles from the same source (SRC) have strong charge dependence.



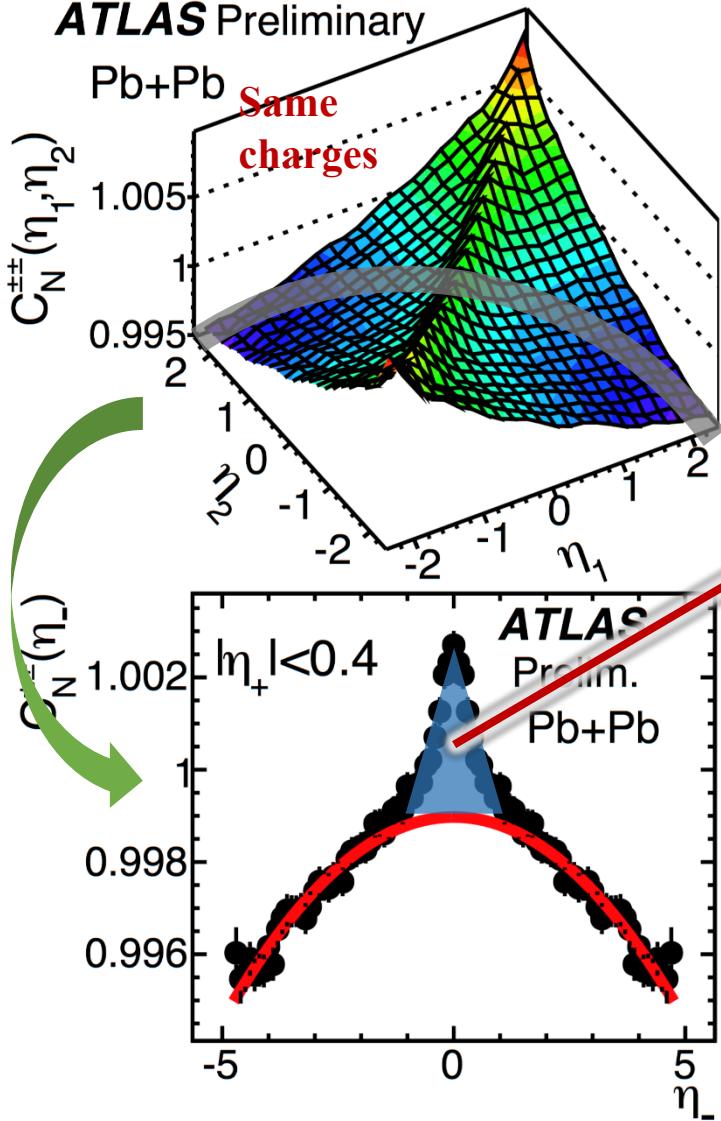
- Ratio of opposite to same charges $R(\eta_1, \eta_2)$
 - Very strong Gaussian-like SRC;
 - Very weak LRC: charge-independent;
- Amplitude of $R(\eta_1, \eta_2)$ along η_+ : $f(\eta_+)$, reflects the strength of SRC in the longitudinal direction;
- Assumption: strength of SRC along η_+ is same for same charge and opposite charge.



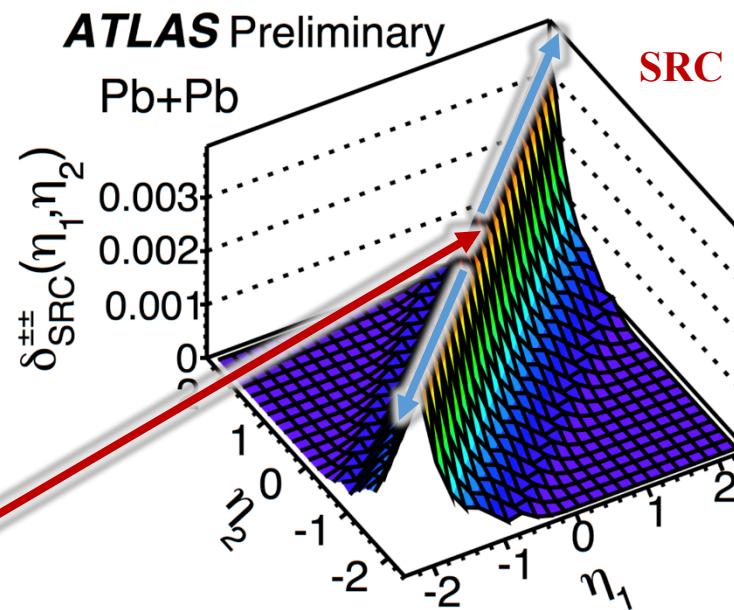
Estimation of short-range correlation

- To estimate SRC, LRC pedestal is estimated first.

ATLAS Preliminary



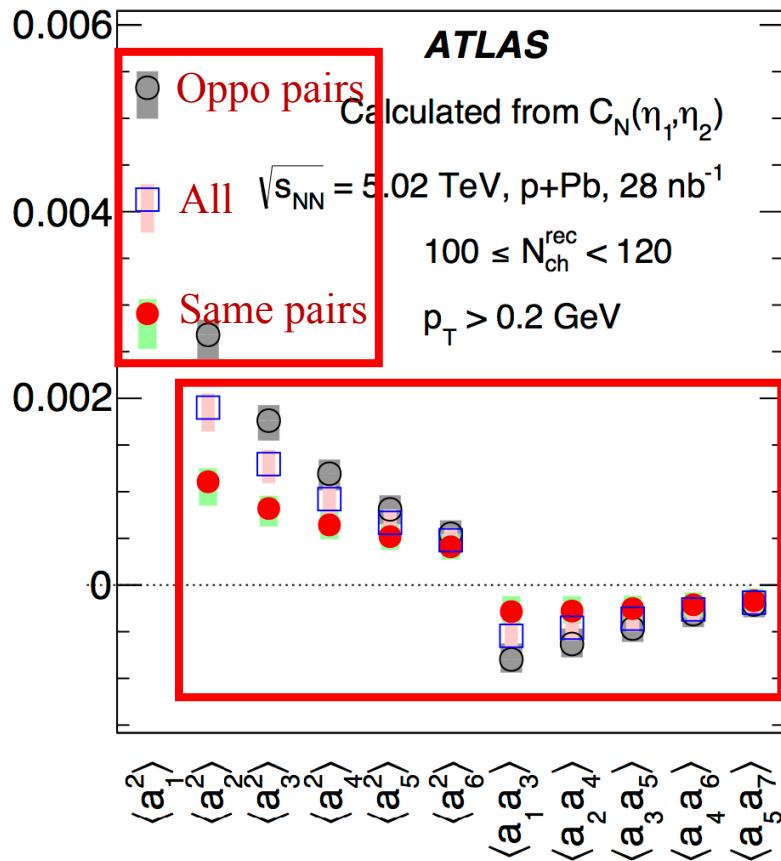
ATLAS Preliminary



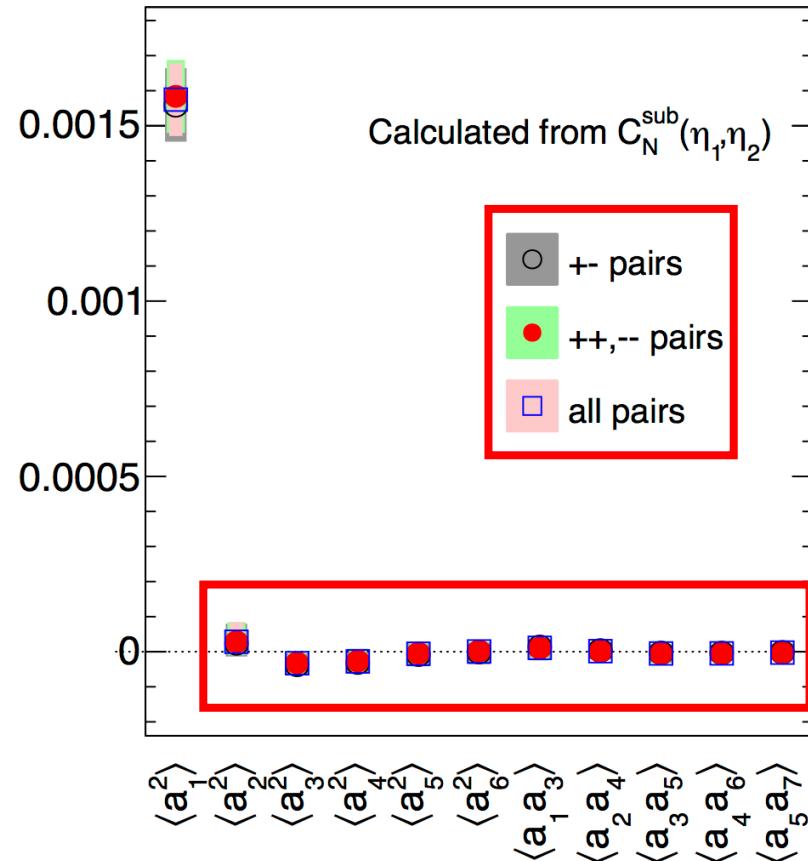
- $C(\eta_1, \eta_2)$ from same charge used to estimate LRC pedestal because of small SRC;
- LRC pedestal is fitted with quadratic function;
- The additional structure upon LRC pedestal determines the shape of SRC;
- The full $\delta_{SRC}(\eta_1, \eta_2)$ is then populated using $f(\eta_+)$ scaling.

Why it is important to remove SRC?

Before SRC removal



After SRC removal

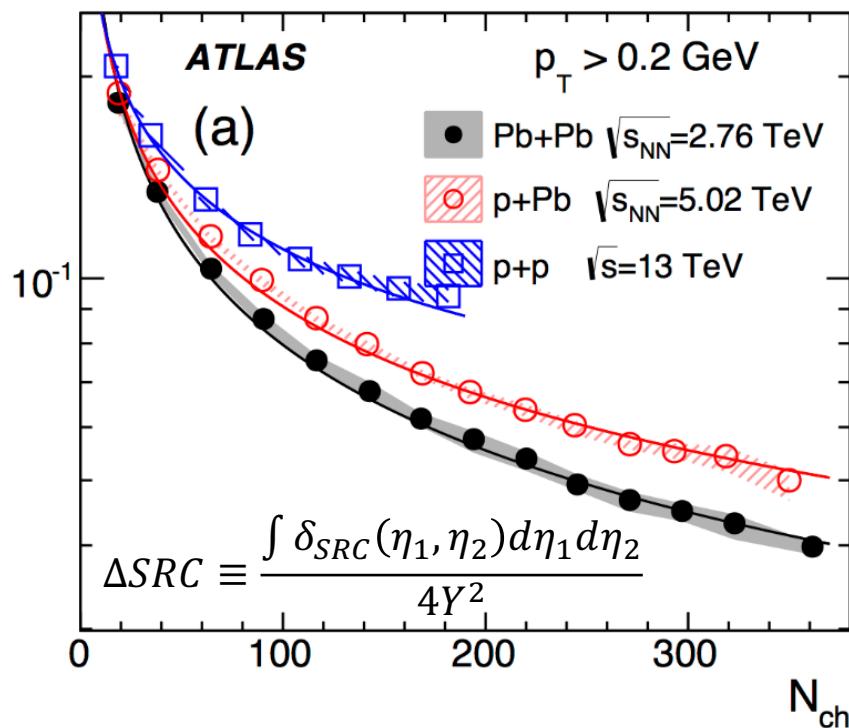


- Higher order coefficients observed;
- Coefficients have charge dependent;
- Results hard to interpret: due to SRC!

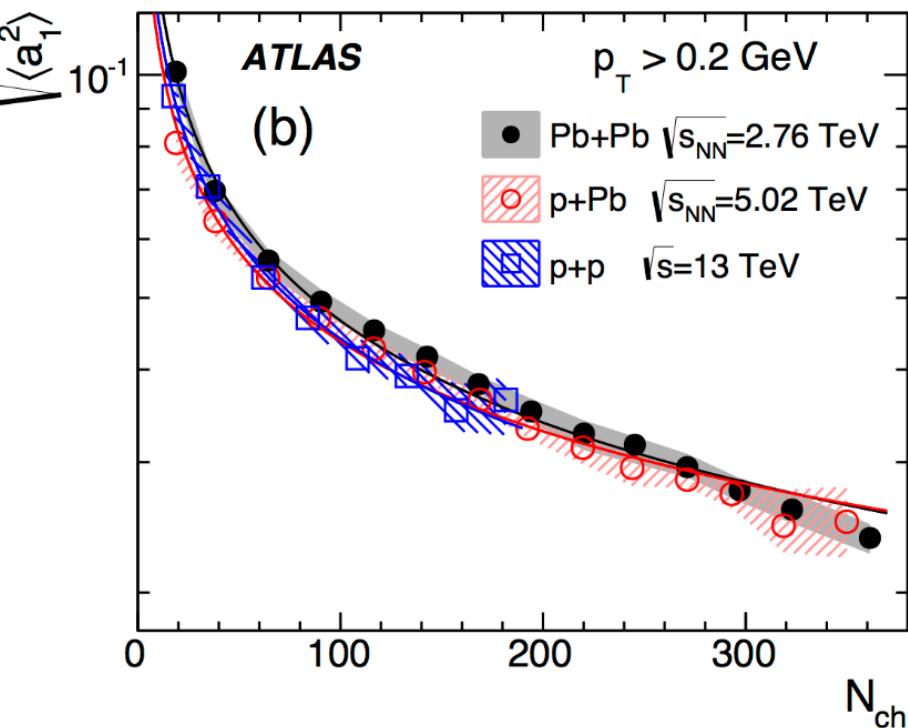
- Simpler picture after SRC removal!
- LRC dominated by linear fluctuation;
- LRC is charge independent.

Why it is important to remove SRC?

Short-range correlation



Long-range correlation



- SRC increases towards peripheral;
- SRC is stronger in small systems;
- Follow-up model studies
 - Viscous hydro: [PLB 2015.11.063](#)
 - Length of sources fluctuation: [PRC 93, 064910 \(2016\)](#)
 - Saturation scale fluctuation: [PRC 94, 044918 \(2016\)](#)

- Dominated by linear fluctuation!
- Independent of collision system!