

Presentación 2

1. Si $f = f(r)$ con $r = \sqrt{x^2 + y^2 + z^2}$, demuestra que

$$\nabla f(r) = \hat{r} \frac{df(r)}{dr}$$

Por la ec. (6), y usando coordenadas esféricas

$$\nabla f(r) = \sum_{i=1}^3 \frac{\hat{e}_i}{h_i} \frac{\partial f(r)}{\partial \theta} = \frac{\hat{e}_r}{h_r} \frac{\partial f(r)}{\partial r} + \frac{\hat{e}_\theta}{h_\theta} \frac{\partial f(r)}{\partial \theta} + \frac{\hat{e}_\phi}{h_\phi} \frac{\partial f(r)}{\partial \phi}$$

y como f solo depende de r , entonces

$$\nabla f(r) = \frac{\hat{e}_r}{h_r} \frac{\partial f(r)}{\partial r}$$

ahora, como $\mathbf{r} = \hat{i}r \cos\theta \cos\phi + \hat{j}r \cos\theta \sin\phi + \hat{k}r \sin\theta$

$$\begin{aligned} h_r = \left\| \frac{\partial \mathbf{r}}{\partial r} \right\| &= \sqrt{\sin^2\theta \cos^2\phi + \sin^2\theta \sin^2\phi + \cos^2\theta} = \sqrt{\sin^2\theta [\cos^2\phi + \sin^2\phi] + \cos^2\theta} \\ &= \sqrt{\sin^2\theta + \cos^2\theta} = 1 \end{aligned}$$

así

$$\nabla f(r) = \hat{e}_r \frac{\partial f}{\partial r} = \hat{r} \frac{\partial f}{\partial r}$$

2. Demuestra que el campo eléctrico de una carga puntual

$$E = \frac{q\hat{r}}{4\pi\epsilon_0 r^2}$$

cumple $\nabla \cdot E = 0$, para $r \neq 0$

Por la ec. (8), y usando coordenadas esféricas

$$\text{Div} E = \frac{1}{h} \sum_{i=1}^3 \frac{\partial}{\partial u_i} \left(\frac{E_i h}{h_i} \right) = \frac{1}{h} \left[\frac{\partial}{\partial u_r} \left(\frac{E_r h}{h_r} \right) + \frac{\partial}{\partial u_\theta} \left(\frac{E_\theta h}{h_\theta} \right) + \frac{\partial}{\partial u_\phi} \left(\frac{E_\phi h}{h_\phi} \right) \right]$$

y como E solo tiene componente E_r

$$\nabla \cdot E = \frac{1}{h} \frac{\partial}{\partial r} \left(\frac{E_r h}{h_r} \right)$$

y como ya tenemos h_r , solo faltan h_θ y h_ϕ

$$h_\theta = \left\| \frac{\partial \mathbf{r}}{\partial \theta} \right\| = \sqrt{r^2 \cos^2\theta \cos^2\phi + r^2 \cos^2\theta \sin^2\phi + r^2 \sin^2\theta} = \sqrt{r^2 [\cos^2\theta \cos^2\phi + \cos^2\theta \sin^2\phi + \sin^2\theta]}$$

$$= r \sqrt{\cos^2 \theta (\cancel{\cos^2 \phi + \sin^2 \phi}) + \sin^2 \theta} = r \sqrt{\cancel{\cos^2 \theta + \sin^2 \theta}} = r$$

$$h_\phi = \left\| \frac{\partial \mathbf{r}}{\partial \theta} \right\| = \sqrt{r^2 \sin^2 \theta \sin^2 \phi + r^2 \sin^2 \theta \cos^2 \phi} = \sqrt{r^2 [\sin^2 \theta \sin^2 \phi + \sin^2 \theta \cos^2 \phi]}$$

$$= r \sqrt{\sin^2 \theta (\cancel{\cos^2 \phi + \sin^2 \phi})} = r \sin \theta$$

así

$$\nabla \cdot E = \frac{1}{h_r h_\theta h_\phi} \frac{\partial}{\partial r} \left(\frac{q \cancel{h_r} h_\theta h_\phi}{4\pi \epsilon_o r^2 \cancel{h_r}} \right) = \frac{\cancel{\sin \theta}}{r^2 \cancel{\sin \theta}} \frac{\partial}{\partial r} \left(\frac{q \cancel{r^2}}{4\pi \epsilon_o \cancel{r^2}} \right) = 0$$

para $r \neq 0$ (así no se indetermina)

3. La ley de Gauss para el campo eléctrico tiene la forma :

$$\oint E \cdot dS = \frac{q}{\epsilon_o}$$

donde $q = \int \rho dV$ es la carga encerrada en la superficie y ρ su densidad volumétrica. Demuestra la ley de Gauss en su forma diferencial

$$\nabla \cdot E = \frac{\rho}{\epsilon_o}$$

Por el teorema de la divergencia,

$$\oint E \cdot dS = \int \nabla \cdot E dV = \frac{q}{\epsilon_o} = \frac{1}{\epsilon_o} \int \rho dV$$

como la expresión anterior debe ser cierta para cualquier volumen (incluso uno infinitesimal) se tiene que se pueden igualar los integrandos

$$\nabla \cdot E = \frac{\rho}{\epsilon_o}$$

4. Demuestra que : $\nabla \times (\phi A) = \phi \nabla \times A + \nabla \phi \times A$

Usando la ec. (10) en su forma matricial:

$$\begin{aligned} \nabla \times (\phi \mathbf{A}) &= \frac{1}{h} \begin{vmatrix} h_1 \hat{e}_1 & h_2 \hat{e}_2 & h_3 \hat{e}_3 \\ \frac{\partial}{\partial u_1} & \frac{\partial}{\partial u_2} & \frac{\partial}{\partial u_3} \\ \phi A_1 h_1 & \phi A_2 h_2 & \phi A_3 h_3 \end{vmatrix} \\ &= \frac{1}{h} \left[h_1 \left(\frac{\partial}{\partial u_2} (\phi A_3 h_3) - \frac{\partial}{\partial u_3} (\phi A_2 h_2) \right) \hat{e}_1 - h_2 \left(\frac{\partial}{\partial u_1} (\phi A_3 h_3) - \frac{\partial}{\partial u_3} (\phi A_1 h_1) \right) \hat{e}_2 \right. \end{aligned}$$

$$+h_3 \left(\frac{\partial}{\partial u_1}(\phi A_2 h_2) - \frac{\partial}{\partial u_2}(\phi A_1 h_1) \right) \hat{e}_3]$$

o más compacto, usando notación de índices:

$$\nabla \times (\phi \mathbf{A})_i = \frac{h_i}{h} \epsilon_{ijk} \frac{\partial}{\partial u_j} (\phi A_k h_k)$$

ahora metamos el diferencial al paréntesis y reordenemos

$$\nabla \times (\phi \mathbf{A})_i = \frac{h_i}{h} \epsilon_{ijk} \left[\frac{\partial}{\partial u_j} (\phi) A_k h_k + \phi \frac{\partial}{\partial u_j} (A_k h_k) \right]$$

cambiamos a la notación presentada en el material adicional de índices

$$= \frac{h_i}{h} \epsilon_{ijk} [\phi \partial_j (A_k h_k) + \partial_j (\phi) A_k h_k]$$

$$\frac{h_i}{h} \phi \epsilon_{ijk} \partial_j (A_k h_k) + \frac{h_i}{h} \epsilon_{ijk} \partial_j (\phi) A_k h_k$$

$$= \phi \left(\frac{h_i}{h} \epsilon_{ijk} \partial_j A_k h_k \right) + \left(\frac{h_i}{h} \epsilon_{ijk} \partial_j (\phi) A_k h_k \right)$$

$$= \phi (\nabla \times \mathbf{A})_i + (\nabla \phi \times \mathbf{A})_i$$

$$\therefore \nabla \times (\phi \mathbf{A}) = \phi \nabla \times \mathbf{A} + \nabla \phi \times \mathbf{A}$$

5. El campo electrostático de un dipolo eléctrico $p = p_o \hat{e}_z$ es

$$E = \frac{p_o (2\hat{e}_r \cos\theta + \hat{e}_\theta \sin\theta)}{r^3}$$

Demuestra que:

$$a) \nabla \times E = 0$$

$$b) \text{ para } r \neq 0, \text{ se tiene } \nabla \cdot E = 0$$

Suponiendo que E está en esféricas, entonces:(y usando los factores de escala calculados en los problemas 1 y 2)

$$\begin{aligned} \nabla \times E &= \frac{1}{h} \begin{vmatrix} h_1 \hat{e}_r & h_2 \hat{e}_\theta & h_3 \hat{e}_\phi \\ \frac{\partial}{\partial u_r} & \frac{\partial}{\partial u_\theta} & \frac{\partial}{\partial u_\phi} \\ E_r h_r & E_\theta h_\theta & E_\phi h_\phi \end{vmatrix} \\ &= \frac{1}{r^2 \sin\theta} \left[\left(\frac{\partial}{\partial \theta} (E_\phi r \sin\theta) - \frac{\partial}{\partial \phi} (E_\theta r) \right) \hat{e}_r + r \left(\frac{\partial}{\partial \phi} (E_r) - \frac{\partial}{\partial r} (E_\phi r \sin\theta) \right) \hat{e}_\theta + \right. \end{aligned}$$

$$\begin{aligned}
& r \operatorname{sen} \theta \left(\frac{\partial}{\partial r} (E_\theta r) - \frac{\partial}{\partial \theta} (E_r) \right) \hat{e}_\phi \\
&= \frac{1}{r^2 \operatorname{sen} \theta} \left[-\frac{\partial}{\partial \phi} \left(\frac{\rho_0 \operatorname{sen} \theta r}{r^3} \right) \hat{e}_r + r \frac{\partial}{\partial \phi} \left(\frac{2\rho_0 \cos \theta}{r^3} \right) \hat{e}_\theta + r \operatorname{sen} \theta \left(\frac{\partial}{\partial r} \left(\frac{\rho_0 \operatorname{sen} \theta r}{r^2} \right) - \frac{\partial}{\partial \theta} \left(\frac{2\rho_0 \cos \theta}{r^3} \right) \right) \hat{e}_\phi \right] \\
&= \frac{1}{r} \left((0) \hat{e}_r + (0) \hat{e}_\theta + \left(\frac{-2\rho_0 \operatorname{sen} \theta}{r^3} - \left(-\frac{2\rho_0 \operatorname{sen} \theta}{r^3} \right) \right) \hat{e}_\phi \right) \\
&= (0, 0, 0) \forall r \neq 0
\end{aligned}$$

$$\begin{aligned}
\nabla \cdot E &= \frac{1}{h} \left[\frac{\partial}{\partial r} \left(\frac{E_r h}{h_r} \right) + \frac{\partial}{\partial \theta} \left(\frac{E_\theta h}{h_\theta} \right) + \frac{\partial}{\partial \phi} \left(\frac{E_\phi h}{h_\phi} \right) \right] \\
&= \frac{1}{r^2 \operatorname{sen} \theta} \left[\frac{\partial}{\partial r} (E_r r^2 \operatorname{sen} \theta) + \frac{\partial}{\partial \theta} \left(\frac{E_\theta r^2 \operatorname{sen} \theta}{r} \right) + \frac{\partial}{\partial \phi} \left(\frac{E_\phi r^2 \operatorname{sen} \theta}{r \operatorname{sen} \theta} \right) \right] \\
&= \frac{1}{r^2 \operatorname{sen} \theta} \left[\frac{\partial}{\partial r} \left(\frac{2\rho_0 \cos \theta r^2 \operatorname{sen} \theta}{r^2} \right) + \frac{\partial}{\partial \theta} \left(\frac{\rho_0 \operatorname{sen}^2 \theta r}{r^2} \right) \right] \\
&= \frac{1}{r^2 \operatorname{sen} \theta} \left[\frac{-2\rho_0 \cos \theta \operatorname{sen} \theta}{r^2} + \frac{2\rho_0 \operatorname{sen} \theta \cos \theta}{r^2} \right] = \frac{1}{r^2 \operatorname{sen} \theta} [0] = 0
\end{aligned}$$

6. Demuestra que el Laplaciano en coordenadas cilíndricas y esféricas es el que se presenta, para ello tendrás que calcular los respectivos factores de escala.

de la ecuación (12) se sigue que:

$$\nabla^2 f = \frac{1}{h} \sum_{i=1}^3 \frac{\partial}{\partial u_i} \left(\frac{h}{h_i^2} \frac{\partial f}{\partial u_i} \right)$$

con coordenadas esféricas, del ejercicio 1 y 2 se tiene:

$$h_r = 1; h_\theta = r; h_\psi = r \operatorname{sen} \theta$$

$$\begin{aligned}
\therefore \nabla^2 f &= \frac{1}{r^2 \operatorname{sen} \theta} \left[\frac{\partial}{\partial r} (r^2 \operatorname{sen} \theta \frac{\partial f}{\partial r}) + \frac{\partial}{\partial \theta} \left(\frac{r^2 \operatorname{sen} \theta}{r^2} \frac{\partial f}{\partial \theta} \right) + \frac{\partial}{\partial \psi} \left(\frac{r^2 \operatorname{sen} \theta}{r^2 \operatorname{sen} \theta^2} \frac{\partial f}{\partial \psi} \right) \right] \\
&= \frac{1}{r^2 \operatorname{sen} \theta} \left[\operatorname{sen} \theta \frac{\partial}{\partial r} (r^2 \frac{\partial f}{\partial r}) + \frac{\partial}{\partial \theta} (\operatorname{sen} \theta \frac{\partial f}{\partial \theta}) + \frac{1}{\operatorname{sen} \theta} \frac{\partial^2 f}{\partial \psi^2} \right] \\
&= \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 \frac{\partial f}{\partial r}) + \frac{1}{r^2 \operatorname{sen} \theta} \frac{\partial}{\partial \theta} (\operatorname{sen} \theta \frac{\partial f}{\partial \theta}) + \frac{1}{r^2 \operatorname{sen} \theta} \frac{\partial^2 f}{\partial \psi^2}
\end{aligned}$$

en coordenadas cilíndricas, dado que:

$$\mathbf{r} = \hat{i}\rho \cos \psi + \hat{j}\rho \sin \psi + \hat{k}z$$

entonces:

$$\frac{\partial \mathbf{r}}{\partial \rho} = \hat{i} \cos \psi + \hat{j} \sin \psi$$

$$\frac{\partial \mathbf{r}}{\partial \psi} = -\hat{i}\rho \sin \psi + \hat{j}\rho \cos \psi$$

$$\frac{\partial \mathbf{r}}{\partial z} = \hat{k}$$

así:

$$h_\rho = \left\| \frac{\partial \mathbf{r}}{\partial \rho} \right\| = \sqrt{\cos^2 \psi + \sin^2 \psi} = 1$$

$$h_\psi = \left\| \frac{\partial \mathbf{r}}{\partial \psi} \right\| = \sqrt{\rho^2 \sin^2 \psi + \rho^2 \cos^2 \psi} = \rho \sqrt{\sin^2 \psi + \cos^2 \psi} = \rho$$

$$h_z = \left\| \frac{\partial \mathbf{r}}{\partial z} \right\| = \sqrt{1} = 1$$

$$\begin{aligned} \therefore \nabla^2 f &= \frac{1}{\rho} \left[\frac{\partial}{\partial \rho} \left(\rho \frac{\partial f}{\partial \rho} \right) + \frac{\partial}{\partial \psi} \left(\frac{\rho}{\rho^2 \frac{\partial f}{\partial \psi}} + \frac{\partial}{\partial z} \left(\rho \frac{\partial f}{\partial z} \right) \right) \right] \\ &= \frac{1}{\rho} \left[\frac{\partial}{\partial \rho} \left(\rho \frac{\partial f}{\partial \rho} \right) + \frac{1}{\rho} \frac{\partial^2 f}{\partial \psi^2} + \rho \frac{\partial^2 f}{\partial z^2} \right] \\ &= \frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial f}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2 f}{\partial \psi^2} + \frac{\partial^2 f}{\partial z^2} \end{aligned}$$

Presentación 3

1. Para el sistema de coordenadas esferoidales prolatas (ξ, η, ϕ) , cuyas reglas de transformación son:

$$x = a \sinh \xi \sin \eta \cos \phi$$

$$y = a \sinh \xi \sin \eta \sin \phi$$

$$z = a \cosh \xi \cos \eta$$

a) **Describe las superficies coordenadas del sistema**

Siguiendo el procedimiento de las notas, elevemos al cuadrado x, y, z:

$$x^2 = a^2 \sinh^2 \xi \sin^2 \eta \cos^2 \phi$$

$$y^2 = a^2 \sinh^2 \xi \sin^2 \eta \sin^2 \phi$$

$$z^2 = a^2 \cosh^2 \xi \cos^2 \eta$$

Ahora desarrollando:

$$\begin{aligned} \frac{x^2 + y^2}{a^2 \sinh^2 \xi} + \frac{z^2}{a^2 \cosh^2 \xi} &= \frac{\cancel{a^2 \sinh^2 \xi} \sin^2 \eta (\cos^2 \phi + \sin^2 \phi)}{\cancel{a^2 \sinh^2 \xi}} + \frac{\cancel{a^2 \cosh^2 \xi} \cos^2 \eta}{\cancel{a^2 \cosh^2 \xi}} \\ &= \cancel{\sin^2 \eta} + \cos^2 \eta = 1 \end{aligned}$$

$$\begin{aligned} \frac{x^2 + y^2}{a^2 \sinh^2 \eta} - \frac{z^2}{a^2 \cosh^2 \eta} &= \frac{\cancel{a^2 \sinh^2 \eta} \sinh^2 \xi (\cos^2 \phi + \sin^2 \phi)}{\cancel{a^2 \sinh^2 \eta}} - \frac{\cancel{a^2 \cosh^2 \eta} \cosh^2 \xi}{\cancel{a^2 \cosh^2 \eta}} \\ &= \sinh^2 \xi - \cosh^2 \xi = -(\cosh^2 \xi - \sinh^2 \xi) = 1 \end{aligned}$$

y así,

$$\frac{x^2 + y^2}{a^2 \sinh^2 \xi} + \frac{z^2}{a^2 \cosh^2 \xi} = 1 \quad (1)$$

$$\frac{z^2}{a^2 \cos^2 \eta} - \frac{x^2 + y^2}{a^2 \sin^2 \eta} = 1 \quad (2)$$

Si $\xi = cte$, la ec. (1) tiene la forma de una elipsoide degenerada (esferoide) $(\frac{x^2}{a^2} + \frac{y^2}{a^2} + \frac{z^2}{b^2} = 1)$ con ejes, $a \sinh \xi$ y $a \cosh \xi$

Si $\eta = cte$, la ec. (2) tiene la forma de un hiperboloide (hiperboloide de revolución) $(-\frac{x^2}{a^2} - \frac{y^2}{b^2} + \frac{z^2}{c^2} = 1)$

Si

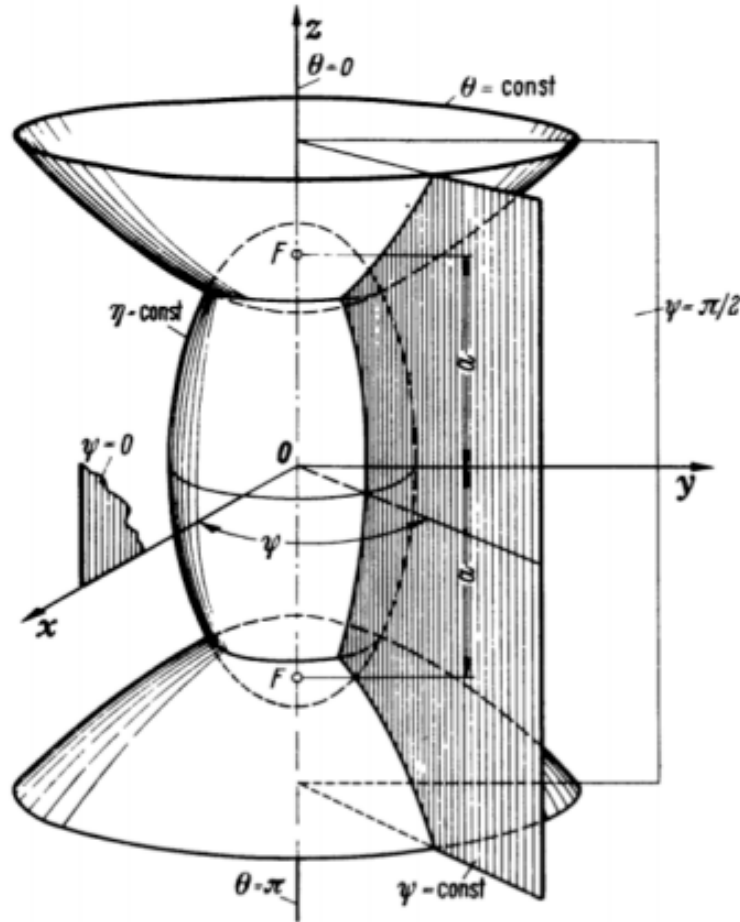
$$\phi = cte \Rightarrow r^2 = x^2 + y^2 + (z \pm a)^2 = x^2 + y^2 + z^2 \pm 2az + a^2$$

$$= a^2 \sinh^2 \xi \sin^2 \eta \cos^2 \phi + a^2 \sinh^2 \xi \sin^2 \eta \sin^2 \phi + a^2 \cosh^2 \xi \cos^2 \eta \pm 2a^2 \cos \xi \cos \eta + a^2$$

$$= a^2 [\sinh^2 \xi \sin^2 \eta (\cancel{\sin^2 \phi + \cos^2 \phi}) + \cosh^2 \xi \cos^2 \eta \pm 2 \cos \xi \cos \eta + 1]$$

$$\begin{aligned}
&= a^2[(\cosh^2 \xi - 1) \sin^2 \eta + \cosh^2 \xi \cos^2 \eta \pm 2 \cos \xi \cos \eta + 1] \\
&= a^2[\cosh^2 \xi (\cancel{\sin^2 \eta + \cos^2 \eta}) - \sin^2 \eta \pm 2 \cos \xi \cos \eta + 1] \\
&= a^2[\cosh^2 \xi + \cos^2 \eta - 1 \pm 2 \cos \xi \cos \eta + 1] \\
&= a^2(\cos \xi \pm \cos \eta)^2
\end{aligned}$$

Lo que describe un plano (o la mitad de uno)



b) **Calcula de manera explícita los factores de escala** (h_ξ, h_η, h_ϕ)

Como $\mathbf{r} = a \sinh \xi \sin \eta \cos \phi \hat{i} + a \sinh \xi \sin \eta \sin \phi \hat{j} + a \cosh \xi \cos \eta \hat{k}$
así,

$$\begin{aligned}
h_\xi = \left\| \frac{\partial \mathbf{r}}{\partial \xi} \right\| &= \sqrt{a^2 \sin^2 \eta \cos^2 \phi \cosh^2 \phi + a^2 \sin^2 \eta + \sin^2 \phi \cosh^2 \xi + a^2 \cos^2 \eta \sinh^2 \xi} \\
&= a \sqrt{\sin^2 \eta \cosh^2 \xi (\cancel{\cos^2 \phi + \sin^2 \phi}) + \cos^2 \eta \sinh^2 \xi}
\end{aligned}$$

$$\begin{aligned}
&= a \sqrt{\sin^2 \eta (\sinh^2 \xi + 1) + \cos^2 \eta \sinh^2 \xi} \\
&= a \sqrt{\sinh^2 \xi (\sin^2 \eta + \cos^2 \eta) + \sin^2 \eta} = a \sqrt{\sinh^2 \xi + \sin^2 \eta} \\
h_\eta &= \left\| \frac{\partial \mathbf{r}}{\partial \eta} \right\| = \sqrt{a^2 \sinh^2 \xi \cos^2 \eta \cos^2 \phi + a^2 \sinh^2 \xi \cos^2 \eta \sin^2 \phi + a^2 \cosh^2 \xi \sin^2 \eta} \\
&= a \sqrt{\sinh^2 \xi \cos^2 \eta (\cos^2 \phi + \sin^2 \phi) + \cosh^2 \xi \sin^2 \eta} \\
&= a \sqrt{\sinh^2 \xi (1 - \sin^2 \eta) + \cosh^2 \xi \sin^2 \eta} \\
&= a \sqrt{\sinh^2 \xi + \sin^2 \eta (\cosh^2 \xi - \sinh^2 \xi)} = h_\xi \\
h_\phi &= \left\| \frac{\partial \mathbf{r}}{\partial \phi} \right\| = \sqrt{a^2 \sinh^2 \xi \sin^2 \eta \sin^2 \phi + a^2 \sinh^2 \xi \sin^2 \eta \cos^2 \phi} \\
&= a \sqrt{\sinh^2 \xi \sin^2 \eta (\sin^2 \phi + \cos^2 \phi)} = a \sinh \xi \sin \eta
\end{aligned}$$

2. Ocupando el mismo sistema de coordenadas esferoidales prolatas (ξ, η, ψ) del ejercicio anterior:

a) Calcula los operadores diferenciables $\nabla \phi$, $\nabla \cdot B$, $\nabla \times B$ y $\nabla^2 \phi$

Usaremos los factores de escala calculados anteriormente:

$$\begin{aligned}
\nabla \phi &= \sum \frac{\hat{e}_i}{h_i} \frac{\partial \phi}{\partial u_i} = \frac{\hat{e}_\xi}{h_\xi} \frac{\partial \phi}{\partial \xi} + \frac{\hat{e}_\eta}{h_\eta} \frac{\partial \phi}{\partial \eta} + \frac{\hat{e}_\psi}{h_\psi} \frac{\partial \phi}{\partial \psi} \\
&= \frac{\hat{e}_\xi}{a \sqrt{\sinh^2 \xi + \sin^2 \eta}} \frac{\partial \phi}{\partial \xi} + \frac{\hat{e}_\eta}{a \sqrt{\sinh^2 \xi + \sin^2 \eta}} \frac{\partial \phi}{\partial \eta} + \frac{\hat{e}_\psi}{a \sinh \xi \sin \eta} \frac{\partial \phi}{\partial \psi} \\
&= \frac{1}{a} \left[\frac{1}{\sqrt{\sinh^2 \xi + \sin^2 \eta}} \left(\hat{e}_\xi \frac{\partial \phi}{\partial \xi} + \hat{e}_\eta \frac{\partial \phi}{\partial \eta} \right) + \frac{\hat{e}_\psi}{\sinh \xi \sin \eta} \frac{\partial \phi}{\partial \psi} \right]
\end{aligned}$$

$$\begin{aligned}
\nabla \cdot B &= \frac{1}{h} \sum \frac{\partial}{\partial u_i} \left(\frac{B_i h}{h_i} \right) = \frac{1}{h} \left[\frac{\partial}{\partial \xi} \left(\frac{B_\xi \mathcal{K}}{\cancel{h_\xi}} \right) + \frac{\partial}{\partial \eta} \left(\frac{B_\eta \mathcal{K}}{\cancel{h_\eta}} \right) + \frac{\partial}{\partial \psi} \left(\frac{B_\psi \mathcal{K}}{\cancel{h_\psi}} \right) \right] \\
&= \frac{1}{h_\xi h_\eta h_\psi} \left[\frac{\partial}{\partial \xi} (B_\xi h_\eta h_\psi) + \frac{\partial}{\partial \eta} (B_\eta h_\xi h_\psi) + \frac{\partial}{\partial \psi} (B_\psi h_\xi h_\eta) \right]
\end{aligned}$$

$$\begin{aligned}
&= \frac{\cancel{a}^2}{a^2(\sinh^2 \xi + \sinh^2 \eta) \sinh \xi \sinh \eta} \left[\frac{\partial}{\partial \xi} (B_\xi \sqrt{\sinh^2 \xi + \sinh^2 \eta} \sinh \xi \sinh \eta) + \right. \\
&\quad \left. \frac{\partial}{\partial \eta} (B_\eta \sqrt{\sinh^2 \xi + \sinh^2 \eta} \sinh \xi) + \frac{\partial}{\partial \psi} (B_\psi (\sinh^2 \xi + \sinh^2 \eta)) \right] \\
&\quad \nabla \times B = \frac{1}{h} \begin{vmatrix} h_\xi \hat{e}_\xi & h_\eta \hat{e}_\eta & h_\psi \hat{e}_\psi \\ \frac{\partial}{\partial \xi} & \frac{\partial}{\partial \eta} & \frac{\partial}{\partial \psi} \\ B_\xi h_\xi & B_\eta h_\eta & B_\psi h_\psi \end{vmatrix} \\
&= \frac{1}{h} \left[h_\xi \left(\frac{\partial}{\partial \eta} (B_\psi h_\psi) - \frac{\partial}{\partial \psi} (B_\eta h_\eta) \right) \hat{e}_\xi + h_\eta \left(\frac{\partial}{\partial \psi} (B_\xi h_\xi) - \frac{\partial}{\partial \xi} (B_\psi h_\psi) \right) \hat{e}_\eta + \right. \\
&\quad \left. h_\psi \left(\frac{\partial}{\partial \xi} (B_\eta h_\eta) - \frac{\partial}{\partial \eta} (B_\xi h_\xi) \right) \hat{e}_\psi \right] \\
&= \frac{1}{h_\xi h_\psi} \left(\frac{\partial}{\partial \eta} (B_\psi h_\psi) - \frac{\partial}{\partial \psi} (B_\eta h_\eta) \right) \hat{e}_\xi + \frac{1}{h_\xi h_\psi} \left(\frac{\partial}{\partial \psi} (B_\xi h_\xi) - \frac{\partial}{\partial \xi} (B_\psi h_\psi) \right) \hat{e}_\eta + \\
&\quad \frac{1}{h_\xi h_\eta} \left(\frac{\partial}{\partial \xi} (B_\eta h_\eta) - \frac{\partial}{\partial \eta} (B_\xi h_\xi) \right) \hat{e}_\psi \\
&= \frac{1}{a^2 \sinh \xi \sinh \eta \sqrt{\sinh^2 \xi + \sinh^2 \eta}} \left(\frac{\partial}{\partial \eta} (B_\psi a \sinh \xi \sinh \eta) - \frac{\partial}{\partial \psi} (B_\eta a \sqrt{\sinh^2 \xi + \sinh^2 \eta}) \right) \hat{e}_\xi \\
&\quad + \frac{1}{a^2 \sinh \xi \sinh \eta \sqrt{\sinh^2 \xi + \sinh^2 \eta}} \left(\frac{\partial}{\partial \psi} (B_\xi a \sinh \eta \sinh \xi) - \frac{\partial}{\partial \xi} (B_\psi a \sinh \eta \sinh \xi) \right) \hat{e}_\eta \\
&\quad + \frac{1}{a^2 (\sinh^2 \xi + \sinh^2 \eta)} \left(\frac{\partial}{\partial \xi} (B_\eta a \sqrt{\sinh^2 \xi + \sinh^2 \eta}) - \frac{\partial}{\partial \eta} (B_\xi a \sqrt{\sinh^2 \xi + \sinh^2 \eta}) \right) \hat{e}_\psi \\
&= \frac{1}{a \sqrt{\sinh^2 \xi + \sinh^2 \eta}} \left(\frac{\frac{\partial}{\partial \eta} (B_\eta \sinh \eta)}{\sinh \eta} - \frac{\sqrt{\sinh^2 \xi + \sinh^2 \eta} \frac{\partial}{\partial \xi} B_\eta}{\sinh \xi \sinh \eta} + \frac{\partial}{\partial \psi} B_\xi - \frac{\frac{\partial}{\partial \xi} (B_\psi \sinh \xi)}{\sinh \xi} \right. \\
&\quad \left. + \frac{\frac{\partial}{\partial \xi} (B_\eta \sqrt{\sinh^2 \xi + \sinh^2 \eta})}{\sqrt{\sinh^2 \xi + \sinh^2 \eta}} - \frac{\frac{\partial}{\partial \eta} (B_\xi \sqrt{\sinh^2 \xi + \sinh^2 \eta})}{\sqrt{\sinh^2 \xi + \sinh^2 \eta}} \right)
\end{aligned}$$

$$\begin{aligned}
\nabla^2 \phi &= \frac{1}{h} \sum \frac{\partial}{\partial u_i} \left(\frac{h}{h_i^2} \frac{\partial \phi}{\partial u_i} \right) \\
&= \frac{1}{h} \left[\frac{\partial}{\partial \xi} \left(\frac{h}{h_\xi^2} \frac{\partial \phi}{\partial \xi} \right) + \frac{\partial}{\partial \eta} \left(\frac{h}{h_\eta^2} \frac{\partial \phi}{\partial \eta} \right) + \frac{\partial}{\partial \psi} \left(\frac{h}{h_\psi^2} \frac{\partial \phi}{\partial \psi} \right) \right] \\
&= \frac{1}{h} \left[\frac{\partial}{\partial \xi} \left(h_\psi \frac{\partial \phi}{\partial \xi} \right) + \frac{\partial}{\partial \eta} \left(h_\psi \frac{\partial \phi}{\partial \eta} \right) + \frac{\partial}{\partial \psi} \left(\frac{a(\sinh^2 \xi + \sinh^2 \eta)}{\sinh \xi \sinh \eta} \frac{\partial \phi}{\partial \psi} \right) \right] \\
&= \frac{1}{a^2(\sinh^2 \xi + \sinh^2 \eta) \sinh \xi \sinh \eta} \left[\frac{\partial}{\partial \xi} (\alpha \sinh \xi \sinh \eta \frac{\partial \phi}{\partial \xi}) + \frac{\partial}{\partial \eta} (\alpha \sinh \xi \sinh \eta \frac{\partial \phi}{\partial \eta}) + \right. \\
&\quad \left. \frac{\alpha(\sinh^2 \xi + \sinh^2 \eta)}{\sinh \xi \sinh \eta} \frac{\partial^2 \phi}{\partial \psi^2} \right] \\
&= \frac{1}{a^2(\sinh^2 \xi + \sinh^2 \eta) \sinh \xi \sinh \eta} \left[\sinh \eta \frac{\partial}{\partial \xi} \left(\sinh \xi \frac{\partial \phi}{\partial \xi} \right) + \sinh \xi \frac{\partial}{\partial \eta} \left(\sinh \eta \frac{\partial \phi}{\partial \eta} \right) + \right. \\
&\quad \left. \frac{\sinh^2 \xi + \sinh^2 \eta}{\sinh \xi \sinh \eta} \frac{\partial^2 \phi}{\partial \psi^2} \right]
\end{aligned}$$