

1. Calcula las transformadas seno y coseno de Fourier de la función definida por :

$$f(x) = \begin{cases} \sin x & , \quad 0 < x < a \\ 0 & , \quad x > a \end{cases}$$

**SOLUCIÓN:**

Sabemos que las funciones  $y = \sin x$  y  $y = 0$  son continuas en todo su dominio, así que podemos decir que  $f(x)$  es continua a trozos en el intervalo  $(0, \infty)$ , o bien  $f(x) \in P_1(0, \infty)$ , ahora veamos que  $f(x) \in A_1(R^+)$

$$\int_0^{\infty} |f(x)| dx$$

y por propiedades de la integral

$$\begin{aligned} \int_0^{\infty} |f(x)| dx &= \int_0^a |f(x)| dx + \int_a^{\infty} |f(x)| dx \\ &= \int_0^a |\sin x| dx + \int_a^{\infty} 0 dx \end{aligned}$$

pero como  $|\sin x| \leq 1$ , entonces

$$\begin{aligned} &= \int_0^a |\sin x| dx \leq \int_0^a 1 dx = a \\ &\rightarrow \int_0^{\infty} |f(x)| dx \text{ es finito} \quad \therefore f(x) \in A_1(R^+) \end{aligned}$$

por lo anterior, podemos aplicar la transformada seno y coseno de Fourier  
la transformada seno de Fourier es por definición

$$F_s[f; x \rightarrow \xi] = \sqrt{\frac{2}{\pi}} \int_0^{\infty} f(x) \sin \xi x dx$$

y por propiedades de la integral

$$\begin{aligned} F_s[f; x \rightarrow \xi] &= \sqrt{\frac{2}{\pi}} \left[ \int_0^a f(x) \sin \xi x dx + \int_a^{\infty} f(x) \sin \xi x dx \right] \\ &= \sqrt{\frac{2}{\pi}} \left[ \int_0^a \sin x \sin \xi x dx + \int_a^{\infty} (0) \sin \xi x dx \right] \\ &= \sqrt{\frac{2}{\pi}} \int_0^a \sin x \sin \xi x dx \end{aligned}$$

ahora, usemos la identidad trigonométrica  $\sin \alpha \sin \beta = \frac{1}{2}(\cos(\alpha - \beta) - \cos(\alpha + \beta))$

$$\begin{aligned} F_s[f; x \rightarrow \xi] &= \sqrt{\frac{2}{\pi}} \int_0^a \frac{1}{2}(\cos(x - \xi x) - \cos(x + \xi x))dx \\ &= \sqrt{\frac{2}{\pi}} \frac{1}{2} \left[ \int_0^a \cos(x(1 - \xi))dx - \int_0^a \cos(x(1 + \xi))dx \right] \\ &= \frac{1}{\sqrt{2\pi}} \left[ \int_0^a \frac{1 - \xi}{1 - \xi} \cos(x(1 - \xi))dx - \int_0^a \frac{1 + \xi}{1 + \xi} \cos(x(1 + \xi))dx \right] \end{aligned}$$

entonces, sea  $u = x(1 - \xi)$  y  $v = x(1 + \xi)$ ,  $du = (1 - \xi)dx$ ;  $dv = (1 + \xi)dx$ (omitiré los límites de integración por un momento solo por comodidad)

$$\begin{aligned} F_s[f; x \rightarrow \xi] &= \frac{1}{\sqrt{2\pi}} \left[ \frac{1}{u} \int \cos u du - \frac{1}{v} \int \cos v dv \right] \\ &= \frac{1}{\sqrt{2\pi}} \left[ \frac{1}{1 - \xi} \sin u \Big| - \frac{1}{1 + \xi} \sin v \Big| \right] \\ &= \frac{1}{\sqrt{2\pi}} \left[ \frac{1}{1 - \xi} \sin x(1 - \xi) \Big|_0^a - \frac{1}{1 + \xi} \sin x(1 + \xi) \Big|_0^a \right] \\ &= \frac{1}{\sqrt{2\pi}} \left[ \frac{1}{1 - \xi} (\sin a(1 - \xi) - \cancel{\sin 0(1 - \xi)}) - \frac{1}{1 + \xi} (\sin a(1 + \xi) - \cancel{\sin 0(1 + \xi)}) \right] \\ &= \frac{(1 + \xi) \sin a(1 - \xi) - (1 - \xi) \sin a(1 + \xi)}{(1 - \xi^2)\sqrt{2\pi}} = F_s(\xi) \end{aligned}$$

Ahora calculemos la transformada coseno de Fourier, que por definición

$$F_c[f; x \rightarrow \xi] = \sqrt{\frac{2}{\pi}} \int_0^\infty f(x) \cos \xi x dx$$

y por propiedades de la integral

$$\begin{aligned} &= \sqrt{\frac{2}{\pi}} \left[ \int_0^a f(x) \cos \xi x dx + \int_a^\infty f(x) \cos \xi x dx \right] \\ &= \sqrt{\frac{2}{\pi}} \left[ \int_0^a \sin x \cos \xi x dx + \int_a^\infty \cancel{(0) \cos \xi x dx} \right] \\ &= \sqrt{\frac{2}{\pi}} \left[ \int_0^a \sin x \cos \xi x dx \right] \end{aligned}$$

ahora, usemos la identidad trigonométrica  $\sin \alpha \cos \beta = \frac{1}{2}(\sin(\alpha - \beta) + \sin(\alpha + \beta))$

$$\begin{aligned}
F_c[f; x \rightarrow \xi] &= \sqrt{\frac{2}{\pi}} \frac{1}{2} \left[ \int_0^a \sin(x - \xi x) dx + \int_0^a \sin(x + \xi x) dx \right] \\
&= \frac{1}{\sqrt{2\pi}} \left[ \int_0^a \sin x (1 - \xi) dx + \int_0^a \sin x (1 + \xi) dx \right] \\
&= \frac{1}{\sqrt{2\pi}} \left[ \int_0^a \frac{1 - \xi}{1 - \xi} \sin x (1 - \xi) dx + \int_0^a \frac{1 + \xi}{1 + \xi} \sin x (1 + \xi) dx \right]
\end{aligned}$$

entonces, sea  $u = x(1 - \xi)$  y  $v = x(1 + \xi)$ ,  $du = (1 - \xi)dx$ ;  $dv = (1 + \xi)dx$  (omitiré los límites de integración por un momento solo por comodidad)

$$\begin{aligned}
F_c[f; x \rightarrow \xi] &= \frac{1}{\sqrt{2\pi}} \left[ \frac{1}{1 - \xi} \int \sin u du + \frac{1}{1 + \xi} \int \sin v dv \right] \\
&= \frac{1}{\sqrt{2\pi}} \left[ \frac{1}{1 - \xi} (-\cos u) + \frac{1}{1 + \xi} (-\cos v) \right] \\
&= \frac{1}{\sqrt{2\pi}} \left[ \frac{1}{1 - \xi} (-\cos x(1 - \xi)) \Big|_0^a + \frac{1}{1 + \xi} (-\cos x(1 + \xi)) \Big|_0^a \right] \\
&= \frac{1}{\sqrt{2\pi}} \left[ \frac{1}{1 - \xi} (\cos 0(1 - \xi) - \cos a(1 - \xi)) + \frac{1}{1 + \xi} (\cos 0(1 + \xi) - \cos a(1 + \xi)) \right] \\
&= \frac{(1 + \xi)(1 - \cos a(1 - \xi)) + (1 - \xi)(1 - \cos a(1 + \xi))}{(1 - \xi^2)\sqrt{2\pi}} = F_c(\xi)
\end{aligned}$$

## 2. Calcula la transformada de Fourier de la función:

$$f(x) = \begin{cases} e^{-ax} & , \quad x > 0, a > 0 \\ -e^{ax} & , \quad x < 0, a > 0 \end{cases}$$

### SOLUCIÓN:

Las funciones  $y = e^{-ax}$  y  $y = -e^{ax}$  son continuas en todo su dominio, así que podemos decir que  $f(x)$  es continua a trozos en el intervalo  $(-\infty, \infty)$ , o bien  $f(x) \in P_1(-\infty, \infty)$ , ahora veamos que  $f(x) \in A_1(R)$

$$\int_{-\infty}^{\infty} |f(x)| dx$$

y por propiedades de la integral

$$\int_{-\infty}^{\infty} |f(x)| dx = \int_{-\infty}^0 |f(x)| dx + \int_0^{\infty} |f(x)| dx$$

$$\begin{aligned}
&= \int_{-\infty}^0 |-e^{ax}|dx + \int_0^{\infty} |e^{-ax}|dx \\
&= \int_{-\infty}^0 e^{ax}dx + \int_0^{\infty} e^{-ax}dx = \frac{1}{a}(e^0 - e^{-\infty}) - \frac{1}{a}(e^{-\infty} - e^0) = 0 \\
&\rightarrow \int_{-\infty}^{\infty} |f(x)|dx \text{ es finito} \quad (x) \in A_1(R)
\end{aligned}$$

por lo anterior, podemos aplicar la transformada de Fourier  
por definición

$$F[f; x \rightarrow \xi] = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x)e^{i\xi x} dx$$

y por propiedades de la integral

$$\begin{aligned}
F[f; x \rightarrow \xi] &= \frac{1}{\sqrt{2\pi}} \left[ \int_{-\infty}^0 f(x)e^{i\xi x} dx + \int_0^{\infty} f(x)e^{i\xi x} dx \right] \\
&= \frac{1}{\sqrt{2\pi}} \left[ - \int_{-\infty}^0 e^{ax} e^{i\xi x} dx + \int_0^{\infty} e^{-ax} e^{i\xi x} dx \right]
\end{aligned}$$

y por propiedades de la exponencial

$$\begin{aligned}
F[f; x \rightarrow \xi] &= \frac{1}{\sqrt{2\pi}} \left[ - \int_{-\infty}^0 e^{ax+i\xi x} dx + \int_0^{\infty} e^{i\xi x-ax} dx \right] \\
&= \frac{1}{\sqrt{2\pi}} \left[ - \int_{-\infty}^0 e^{x(a+i\xi)} dx + \int_0^{\infty} e^{x(i\xi-a)} dx \right] \\
&= \frac{1}{\sqrt{2\pi}} \left[ - \int_{-\infty}^0 \frac{a+i\xi}{a+i\xi} e^{x(a+i\xi)} dx + \int_0^{\infty} \frac{i\xi-a}{i\xi-a} e^{x(i\xi-a)} dx \right]
\end{aligned}$$

sea  $u = x(a+i\xi)$  y  $v = x(i\xi-a)$ ,  $du = (a+i\xi)dx$ ,  $dv = (i\xi-a)dx$

$$\begin{aligned}
F[f; x \rightarrow \xi] &= \frac{1}{\sqrt{2\pi}} \left[ - \frac{1}{a+i\xi} \int e^u du + \frac{1}{i\xi-a} \int e^v dv \right] \\
&= \frac{1}{\sqrt{2\pi}} \left[ - \frac{1}{a+i\xi} e^u + \frac{1}{i\xi-a} e^v \right] \\
&= \frac{1}{\sqrt{2\pi}} \left[ - \frac{1}{a+i\xi} e^{x(a+i\xi)} \Big|_{-\infty}^0 + \frac{1}{i\xi-a} e^{x(i\xi-a)} \Big|_0^{\infty} \right]
\end{aligned}$$

$$\begin{aligned}
&= \frac{1}{\sqrt{2\pi}} \left[ -\frac{e^{0(a+i\xi)} - \cancel{e^{-\infty(a+i\xi)}}}{a+i\xi} + \frac{e^{\infty(i\xi-a)} - e^{0(i\xi-a)}}{i\xi-a} \right] \\
&= \frac{1}{\sqrt{2\pi}} \left[ -\frac{1}{a+i\xi} + \frac{e^{\infty(i\xi)} \cancel{e^{-a\infty}} - 1}{i\xi-a} \right] \\
&= \frac{1}{\sqrt{2\pi}} \left[ -\frac{1}{a+i\xi} + \frac{-1}{i\xi-a} \right] = \frac{2i\xi}{\sqrt{2\pi}(\xi^2+a^2)} = F(\xi)
\end{aligned}$$

3. La temperatura  $u(x, t)$  de una varilla semiinfinita  $0 \leq x < \infty$  satisface la ecuación diferencial parcial

$$\frac{\partial u}{\partial t} = \kappa \frac{\partial^2 u}{\partial x^2}$$

sujeta a las condiciones

$$u(x, 0) = 0, \quad \frac{\partial u}{\partial x} = \lambda \text{ una constante, cuando } x = 0, t > 0$$

Determina la temperatura  $u(x, t)$  en la varilla

**SOLUCIÓN:**

Apliquemos la transformada coseno de Fourier

$$F_c \left[ \frac{\partial u}{\partial t} \right] = F_c \left[ \kappa \frac{\partial^2 u}{\partial x^2} \right]$$

entonces, como la transformada de coseno de Fourier de una derivada es  $F[f'(t)] = -\xi^2 F(\xi) - f(0)$

$$\frac{\partial U(\xi, t)}{\partial t} = \kappa(-\xi^2 U(\xi, t) - u_x(0, t))$$

donde  $U = F_c(u)$ , y por las condiciones de frontera

$$\frac{\partial U(\xi, t)}{\partial t} = -\kappa(\xi^2 U(\xi, t) + \lambda)$$

$$\frac{-\kappa\xi^2}{-\kappa\xi^2 - \kappa(\xi^2 U(\xi, t) + \lambda)} U' = 1$$

$$\frac{1}{-\kappa\xi^2} \int \frac{-\kappa\xi^2 U'}{-\kappa(\xi^2 U(\xi, t) + \lambda)} dt = \int 1 dt$$

$$\frac{\ln -\kappa(\xi^2 U(\xi, t) + \lambda)}{-\kappa\xi^2} = t + c$$

$$-\kappa(\xi^2 U(\xi, t) + \lambda) = e^{-\kappa \xi^2 t} e^{-\kappa \xi^2 c} = C e^{-\kappa \xi^2 t}$$

$$U(\xi, t) = \frac{1}{-\kappa \xi^2} (C e^{-\kappa \xi^2 t} + \kappa \lambda)$$

por la otra condición inicial se tiene que

$$\frac{1 + \kappa \lambda}{-\kappa \xi} = 0 \quad \rightarrow \quad \kappa = -1/\lambda$$

entonces la solución en el espacio transformado es

$$U(\xi, t) = \frac{\lambda}{\xi^2} (C e^{-\kappa \xi^2 t} - 1)$$

ahora aplicando la transformada inversa coseno de Fourier

$$u(x, t) = \frac{2}{\pi} \int_0^\infty \frac{\lambda}{\xi^2} (C e^{-\kappa \xi^2 t} - 1) \cos \xi x d\xi$$

4. La función seno integral se define como:

$$Si(t) = \int_0^t \frac{\sin u}{u} du$$

mientras que la función coseno integral queda definida por:

$$Ci(t) = - \int_t^\infty \frac{\cos u}{u} du$$

Calcula la transformada de Laplace de  $Si(t)$  y de  $Ci(t)$

**SOLUCIÓN:**

Primero demostremos que  $\frac{d}{dp} L[f(t); t \rightarrow p] = -L[tf(t); t \rightarrow p]$

$$\begin{aligned} \frac{d}{dp} L[f(t); t \rightarrow p] &= \frac{d}{dp} \int_0^\infty f(t) e^{-pt} dt \\ &= \int_0^\infty \frac{d}{dp} (f(t) e^{-pt}) dt = \int_0^\infty f(t) \frac{d}{dp} (e^{-pt}) dt \\ &= \int_0^\infty f(t) (-t) e^{-pt} dt = - \int_0^\infty t f(t) e^{-pt} dt = -L[tf(t); t \rightarrow p] \end{aligned}$$

ahora sí, veamos las transformadas de Laplace, por el teorema fundamental del cálculo

$$Ci'(t) = -\frac{\cos t}{t}$$

$$tCi'(t) = -\cos t$$

$$L[tCi'(t); t \rightarrow p] = L[-\cos t; t \rightarrow u]$$

como la transformada de Laplace es lineal y por lo demostrado anteriormente

$$-\frac{d}{dp}L[Ci'(t); t \rightarrow p] = -L[\cos t; t \rightarrow u]$$

y por el ejemplo 3.5

$$-\frac{d}{dp}L[Ci'(t); t \rightarrow p] = -\frac{u}{u^2 + 1}$$

por el teorema 3.5 de las notas

$$-\frac{d}{dp}(pL[Ci(t); t \rightarrow p] - \cancel{Ci(0)}) = -\frac{u}{u^2 + 1}$$

que integrando

$$pL[Ci(t); t \rightarrow p] = \int_0^p \frac{u}{u^2 + 1} du$$

sea  $w = u^2 + 1$

$$= \frac{1}{2} \int_1^{p^2+1} \frac{dw}{w} = \frac{1}{2}(\ln(p^2 + 1) - \ln 1)$$

$$\therefore L[Ci(t); t \rightarrow p] = \frac{\ln(p^2 + 1)}{2p}$$

ahora para el seno integral

$$Si'(t) = -\frac{\sin t}{t}$$

$$tSi'(t) = -\sin t$$

$$L[tSi'(t); t \rightarrow p] = L[-\sin t; t \rightarrow u]$$

como la transformada de Laplace es lineal y por lo demostrado anteriormente

$$-\frac{d}{dp}L[Si'(t); t \rightarrow p] = -L[\sin t; t \rightarrow u]$$

y por el ejemplo 3.5

$$-\frac{d}{dp}L[Si'(t); t \rightarrow p] = -\frac{1}{u^2 + 1}$$

por el teorema 3.5 de las notas

$$-\frac{d}{dp}(pL[Si(t); t \rightarrow p] - \cancel{Si(0)}) = -\frac{1}{u^2 + 1}$$

que integrando

$$pL[Si(t); t \rightarrow p] = \int_0^p \frac{1}{u^2 + 1} du$$

$$= tg^{-1}(p) - \cancel{tg^{-1}(0)} = tg^{-1}(p)$$

$$\therefore L[Si(t); t \rightarrow p] = \frac{tg^{-1}(p)}{p}$$

5. Una función periódica  $f(t)$  de periodo  $2\pi$  presenta una discontinuidad finita en  $t = \pi$ , está definida por:

$$f(t) = \begin{cases} \sin t & , \quad 0 \leq t < \pi \\ \cos t & , \quad \pi < t \leq 2\pi \end{cases}$$

**Evalúa la transformada de Laplace.**

**SOLUCIÓN:**

Como  $\sin t$  y  $\cos t$  son cotinuas en todo su dominio,  $f(t)$  solo tiene discontinuidades en  $(2n + 1)\pi$ , por lo que es continua a tramos, tambien como esas funciones están acotadas por 1, entonces  $f(t)$  es de orden exponencial con  $M = 2, \sigma = 0$  ( $|f(t)| < Me^{\sigma t} = 2$ ), entonces la transformada de Laplace de  $f(t)$  existe.

Recordando que

$$H(x - a) = \begin{cases} 0 & x - a < 0 \quad ; x < a \\ 1 & x - a > 0 \quad ; x > a \end{cases}$$

$$H(x - b) = \begin{cases} 0 & x - b < 0 \quad ; x < b \\ 1 & x - b > 0 \quad ; x > b \end{cases}$$

entonces, si  $a < b$



$$H(x-a) - H(x-b) = \begin{cases} 0 & x < a \\ 1 & a < x < b \\ 1 - 1 = 0 & b < x \end{cases}$$

con esto, podemos escribir a  $f(t)$  como

$$f(t) = \sin t [H(t - 2n\pi) - H(t - (2n+1)\pi)] + \cos t [H(t - (2n+1)\pi) - H(t - 2(n+1)\pi)]$$

con  $n$  natural ( $0 \in \mathbb{N}$ ), entonces la transformada de Laplace de  $f(t)$  es

$$\begin{aligned} L[f(t); t \rightarrow p] &= \int_0^\infty f(t) e^{-pt} dt \\ &= \int_0^\infty \{ \sin t [H(t - 2n\pi) - H(t - (2n+1)\pi)] + \cos t [H(t - (2n+1)\pi) - H(t - 2(n+1)\pi)] \} e^{-pt} dt \end{aligned}$$

por la linealidad de integrales

$$\begin{aligned} &= \int_0^\infty \sin t [H(t - 2n\pi) - H(t - (2n+1)\pi)] e^{-pt} dt + \int_0^\infty \cos t [H(t - (2n+1)\pi) - H(t - 2(n+1)\pi)] e^{-pt} dt \\ &= \int_0^\infty \sin t H(t - 2n\pi) e^{-pt} dt - \int_0^\infty \sin t H(t - (2n+1)\pi) e^{-pt} dt + \\ &\quad \int_0^\infty \cos t [H(t - (2n+1)\pi) e^{-pt} dt - \int_0^\infty \cos t H(t - 2(n+1)\pi) e^{-pt} dt \end{aligned}$$

y por el ejemplo 3.1 de la notas

$$= \int_{2n\pi}^\infty \sin t e^{-pt} dt - \int_{(2n+1)\pi}^\infty \sin t e^{-pt} dt + \int_{(2n+1)\pi}^\infty \cos t e^{-pt} dt - \int_{2(n+1)\pi}^\infty \cos t e^{-pt} dt$$

ahora resolvamos las 2 integrales necesarias y luego evaluemos los límites

$$\begin{aligned} \int \sin t e^{-pt} dt &= \int \frac{e^{it} - e^{-it}}{2i} e^{-pt} dt \\ &= \int \frac{e^{it-pt} - e^{-it-pt}}{2i} dt = \int \frac{e^{t(i-p)} - e^{-t(i+p)}}{2i} dt \\ &= \frac{1}{2i} \left[ \int e^{t(i-p)} dt - \int e^{-t(i+p)} dt \right] = \frac{1}{2i} \left[ \int \frac{i-p}{i-p} e^{t(i-p)} dt - \int \frac{-(i+p)}{-(i+p)} e^{-t(i+p)} dt \right] \end{aligned}$$

$$= \frac{1}{2i} \left[ \frac{1}{i-p} \int e^{t(i-p)}(i-p)dt - \frac{1}{-(i+p)} \int e^{-t(i+p)}[-(i+p)]dt \right]$$

sea  $u = t(i-p)$  y  $dv = -t(i+p)$ ;  $du = (i-p)dt$   $dv = -(i+p)dt$

$$= \frac{1}{2i} \left[ \frac{1}{i-p} \int e^u du - \frac{1}{-(i+p)} \int e^v dv \right] = \frac{1}{2i} \left[ \frac{e^u}{i-p} - \frac{e^v}{-(i+p)} \right]$$

$$= \frac{1}{2i} \left[ \frac{-(i+p)e^u - (i-p)e^v}{-(i-p)(i+p)} \right] = \frac{(i+p)e^{t(i-p)} + (i-p)e^{-t(i+p)}}{-2i(1+p^2)}$$

$$\int \cos te^{-pt} dt = \int \frac{e^{it} + e^{-it}}{2} e^{-pt} dt$$

$$= \int \frac{e^{it-pt} + e^{-it-pt}}{2} dt = \int \frac{e^{t(i-p)} + e^{-t(i+p)}}{2} dt$$

$$= \frac{1}{2} \left[ \int e^{t(i-p)} dt + \int e^{-t(i+p)} dt \right] = \frac{1}{2} \left[ \int \frac{i-p}{i-p} e^{t(i-p)} dt + \int \frac{-(i+p)}{-(i+p)} e^{-t(i+p)} dt \right]$$

$$= \frac{1}{2} \left[ \frac{1}{i-p} \int e^{t(i-p)}(i-p)dt + \frac{1}{-(i+p)} \int e^{-t(i+p)}[-(i+p)]dt \right]$$

sea  $u = t(i-p)$  y  $dv = -t(i+p)$ ;  $du = (i-p)dt$   $dv = -(i+p)dt$

$$= \frac{1}{2} \left[ \frac{1}{i-p} \int e^u du + \frac{1}{-(i+p)} \int e^v dv \right] = \frac{1}{2} \left[ \frac{e^u}{i-p} + \frac{e^v}{-(i+p)} \right]$$

$$= \frac{1}{2} \left[ \frac{-(i+p)e^u + (i-p)e^v}{-(i-p)(i+p)} \right] = \frac{(i-p)e^{-t(i+p)} - (i+p)e^{t(i-p)}}{2(1+p^2)}$$

ahora si evaluemosy supongamos que  $\text{Re}(p) > 0$

$$L[f(t); t \rightarrow p] = \int_{2n\pi}^{\infty} \sin te^{-pt} dt - \int_{(2n+1)\pi}^{\infty} \sin te^{-pt} dt + \int_{(2n+1)\pi}^{\infty} \cos te^{-pt} dt - \int_{2(n+1)\pi}^{\infty} \cos te^{-pt} dt$$

$$= \left[ \frac{(i+p)e^{\infty(i-p)} + (i-p)e^{-\infty(i+p)}}{-2i(1+p^2)} - \frac{(i+p)e^{2n\pi(i-p)} + (i-p)e^{-2n\pi(i+p)}}{-2i(1+p^2)} \right] -$$

$$\left[ \frac{(i+p)e^{\infty(i-p)} + (i-p)e^{-\infty(i+p)}}{-2i(1+p^2)} - \frac{(i+p)e^{(2n+1)\pi(i-p)} + (i-p)e^{-(2n+1)\pi(i+p)}}{-2i(1+p^2)} \right] +$$

$$\left[ \frac{(i-p)e^{-\infty(i+p)} - (i+p)e^{\infty(i-p)}}{2(1+p^2)} - \frac{(i-p)e^{-(2n+1)\pi(i+p)} - (i+p)e^{(2n+1)\pi(i-p)}}{2(1+p^2)} \right] -$$

$$\begin{aligned}
& \left[ \frac{(i-p)e^{-\infty(i+p)} - (i+p)e^{\infty(i-p)}}{2(1+p^2)} - \frac{(i-p)e^{-2(n+1)\pi(i+p)} - (i+p)e^{2(n+1)\pi(i-p)}}{2(1+p^2)} \right] \\
&= -\frac{(i+p)e^{2n\pi(i-p)} + (i-p)e^{-2n\pi(i+p)}}{-2i(1+p^2)} - \frac{(i+p)e^{(2n+1)\pi(i-p)} + (i-p)e^{-(2n+1)\pi(i+p)}}{-2i(1+p^2)} - \\
& \frac{(i-p)e^{-(2n+1)\pi(i+p)} - (i+p)e^{(2n+1)\pi(i-p)}}{2(1+p^2)} - \frac{(i-p)e^{-2(n+1)\pi(i+p)} - (i+p)e^{2(n+1)\pi(i-p)}}{2(1+p^2)} \\
& \frac{-1}{(1+p^2)} \left[ \frac{(i+p)e^{\pi(i-p)} + (i-p)e^{-\pi(i+p)}}{2i} - \frac{(i-p)e^{-\pi(i+p)} + (i+p)e^{\pi(i-p)}}{2} \right] \\
&= \frac{(i-1)((i+p)e^{\pi(i-p)} + (i-p)e^{\pi(i+p)})}{2i(1+p^2)}
\end{aligned}$$

también lo hice de otra forma pero creo que la correcta es la anterior, igual la pongo por si acaso

integremos por partes, sea  $u = \sin t$  y  $dv = e^{-pt}dt$ ;  $du = \cos t dt$   $v = \frac{e^{-pt}}{p}$

$$\int \sin t e^{-pt} dt = -\sin t \frac{e^{-pt}}{p} + \int \frac{e^{-pt}}{p} \cos t dt$$

ahora sea  $u' = \cos t$  y mantegamos  $dv$

$$\int \sin t e^{-pt} dt = -\sin t \frac{e^{-pt}}{p} + \frac{1}{p} \left[ -\cos t \frac{e^{-pt}}{p} - \int \frac{e^{-pt}}{p} \sin(t) dt \right]$$

$$\int \sin t e^{-pt} dt = -\sin t \frac{e^{-pt}}{p} + \frac{1}{p^2} \left[ -\cos t e^{-pt} - \int \sin(t) e^{-pt} dt \right]$$

$$\int \sin t e^{-pt} dt + \frac{1}{p^2} \int \sin(t) e^{-pt} dt = -\sin t \frac{e^{-pt}}{p} - \frac{1}{p^2} \cos t e^{-pt}$$

$$\left(\frac{1}{p^2} + 1\right) \int \sin(t) e^{-pt} dt = -\sin t \frac{e^{-pt}}{p} - \frac{1}{p^2} \cos t e^{-pt}$$

$$\frac{p^2+1}{p^2} \int \sin(t)e^{-pt} dt = -\sin t \frac{e^{-pt}}{p} - \frac{1}{p^2} \cos t e^{-pt}$$

$$\int \sin(t)e^{-pt} dt = e^{-pt} \left( -\frac{\sin t}{p} - \frac{\cos t}{p^2} \right) \frac{p^2}{p^2+1}$$

$$\int \sin(t)e^{-pt} dt = -e^{-pt} \frac{p \sin t + \cos t}{p^2+1}$$

Sea  $u = \cos t$  y  $dv = e^{-pt} dt$ ;  $du = -\sin t dt$   $v = \frac{-e^{-pt}}{p}$

$$\int \cos t e^{-pt} dt = \cos t \frac{-e^{-pt}}{p} - \frac{1}{p} \int \sin t e^{-pt} dt$$

ahora sea  $u' = \sin t$  y mantengamos  $dv$

$$\int \cos t e^{-pt} dt = \cos t \frac{-e^{-pt}}{p} - \frac{1}{p} \left[ -\sin t \frac{e^{-pt}}{p} + \int \frac{e^{-pt}}{p} \cos t dt \right]$$

$$\int \cos t e^{-pt} dt = \cos t \frac{-e^{-pt}}{p} + \frac{1}{p^2} \left[ \sin t e^{-pt} - \int \cos t e^{-pt} dt \right]$$

$$\frac{1}{p^2} \int \cos t e^{-pt} dt + \int \cos t e^{-pt} dt = \cos t \frac{-e^{-pt}}{p} + \frac{1}{p^2} \sin t e^{-pt}$$

$$\frac{p^2+1}{p^2} \int \cos t e^{-pt} dt = e^{-pt} \left( \frac{-\cos t}{p} + \frac{\sin t}{p^2} \right)$$

$$\int \cos t e^{-pt} dt = e^{-pt} \left( \frac{-\cos t}{p} + \frac{\sin t}{p^2} \right) \frac{p^2}{p^2+1}$$

$$\int \cos t e^{-pt} dt = e^{-pt} \frac{\sin t - p \cos t}{p^2+1}$$

ahora solo hay que evaluar, supongamos que  $\text{Re}(p) > 0$  (si  $\text{Re}(p) < 0$  entonces no está definida)

$$\begin{aligned} L[f(t); t \rightarrow p] &= \int_{2n\pi}^{\infty} \sin t e^{-pt} dt - \int_{(2n+1)\pi}^{\infty} \sin t e^{-pt} dt + \int_{(2n+1)\pi}^{\infty} \cos t e^{-pt} dt - \int_{2(n+1)\pi}^{\infty} \cos t e^{-pt} dt \\ &= - \left[ \cancel{e^{-p\infty}} \frac{p \sin \infty + \cos \infty}{p^2+1} + e^{-p2n\pi} \frac{\overline{p \sin 2n\pi} + \cos 2n\pi}{p^2+1} \right] \\ &\quad - \left[ \cancel{e^{-p\infty}} \frac{p \sin \infty + \cos \infty}{p^2+1} + e^{-p(2n+1)\pi} \frac{\overline{p \sin (2n+1)\pi} + \cos (2n+1)\pi}{p^2+1} \right] \end{aligned}$$

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$$\begin{aligned}
& + \left[ \cancel{e^{-p\infty}} \frac{\sin \infty - p \cos \infty}{p^2 + 1} - e^{-p(2n+1)\pi} \frac{\sin \cancel{(2n+1)\pi} - p \cos (2n+1)\pi}{p^2 + 1} \right] \\
& - \left[ \cancel{e^{-p\infty}} \frac{\sin \infty - p \cos \infty}{p^2 + 1} - e^{-p2(n+1)\pi} \frac{\sin \cancel{2(n+1)\pi} - p \cos 2(n+1)\pi}{p^2 + 1} \right] \\
& = e^{-p2n\pi} \frac{p}{p^2 + 1} + e^{-p(2n+1)\pi} \frac{p}{p^2 + 1} \cancel{- e^{-p(2n+1)\pi} \frac{p}{p^2 + 1}} + e^{-p2(n+1)\pi} \frac{p}{p^2 + 1} \\
& \qquad \qquad \qquad = \frac{pe^{-p(2n+2(n+1))\pi}}{p^2 + 1} = \frac{pe^{-2p(2n+1)\pi}}{p^2 + 1}
\end{aligned}$$

con  $Re(p) > 0$