# CP Violation In and Beyond The Standard Model

Two Higgs Doublet Model Type II Corrections to Flavour Observables

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## The Standard Model

➤ One of the great achievements of the 20th Century, the Standard Model:

$$\mathcal{L} = \underbrace{-\frac{1}{4}F_{\mu\nu}F^{\mu\nu}}_{\text{gauge fields}} \underbrace{+i\bar{\Psi}\cancel{D}\Psi}_{\text{fermions}} \underbrace{+(D_{\mu}\Phi)^{\dagger}(D^{\mu}\Phi) - V(\Phi)}_{\text{Higgs}} \underbrace{-Y_{ij}\bar{\Psi}_{i}\Phi\Psi_{j} + h.c.}_{\text{Yukawa}}$$
(1)

- ➤ A gauge field theory describing matter and its interactions with 25 fundamental particles
- ► Each particle is described by a field transforming under the gauge groups of the Standard Model:  $SU(3)_c \otimes SU(2)_L \otimes U(1)_Y$
- ➤ Has successfully described most particle phenomena we have observed to date



#### 1. Introduction

## Unsolved Problems of the Standard Model

- Quantum gravity; Dark matter; Neutrino masses
- $\blacktriangleright$  Deviations between experiment and theory, e.g.  $\mathcal{R}(\mathit{K}^{(*)})$
- ➤ Sakharov Criteria for Baryogenesis:
  - 1. Baryon Number Violation found in Sphalerons
  - 2. C and CP Violation present but not enough
  - 3. First Order Phase Transition only if  $m_h < 60 \,\text{GeV}$

To answer these questions, we need to consider models to extend our physics Beyond the Standard Model. These models should:

- > preserve predictions in agreement with experiment
- ➤ agree with experimental bounds
- ➤ follow the structures of gauge field theory for a physical model, e.g. renormalisability



# The Two Higgs Doublet Model Type II

In the Standard Model:

➤ One complex Higgs doublet, 4 scalar fields:

$$\Phi_1 = \begin{pmatrix} \phi_1 + i\phi_2 \\ \phi_0 + i\phi_3 \end{pmatrix} \tag{2}$$

- ➤ 3 fields "eaten" by W<sup>±</sup>, Z bosons; 1 real field left, h
- ➤ Introduce the Hermitian conjugate for masses of all fermions

#### In 2HDM:

Add a second doublet, now 8 scalar fields

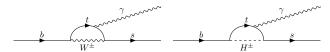
$$\Phi_2 = \begin{pmatrix} \phi_5 + i\phi_6 \\ \phi_4 + i\phi_7 \end{pmatrix} \tag{3}$$

- ➤ Now 5 fields left  $H^{\pm}$ ,  $H^0$ ,  $h^0$ ,  $A^0$
- ➤ No need for Hermitian conjugate
- ▶ In Type II,  $\Phi_1$  couples to down quarks;  $\Phi_2$  to up quarks and charged leptons



# Why Two Higgs Doublet Model?

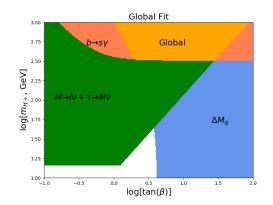
- ➤ Minimal Extension to SM
- ➤ Limited number of new parameters:
  - ightharpoonup Masses of  $H^{\pm}$ ,  $H^0$ ,  $A^0$ ; VEV ratio  $\tan \beta = \frac{v_2}{v_1}$ ; scalar mixing angle
- ➤ Sakharov Criteria:
  - 1. Baryon Number Violation Sphalerons
  - 2. C and CP violation more of it
  - 3. First Order Phase Transition now present
- ➤ Charged weak currents gain additional decay paths, replacing  $W^{\pm}$  with  $H^{\pm}$  allows for easy constraining





# First Inputs

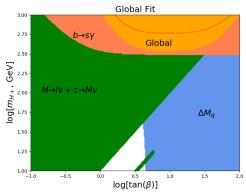
- $\blacktriangleright$  1 $\sigma$  scan
- ➤ Leptonic, mixing, and radiative
- ightharpoonup No real constraint on an eta
- ►  $m_{H^+} > 340 \,\text{GeV}$





## Statistical Fitting of Scans

- Aimed to replicate process of original paper (see below)
- ➤ Scanned at 95% CL
- $ightharpoonup \chi^2$  fit to  $1\sigma$
- ➤ Replicated  $m_{H^+} \gtrsim 316 \, \text{GeV}$ , 95% CL

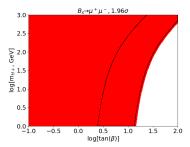


O Deschamps et al, Phys. Rev. D82 (2010) 073012, arxiv:0907.5135 [hep-ph]

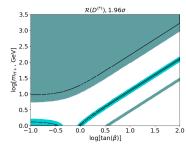


## New Inputs

$$\blacktriangleright$$
  $B_s \rightarrow \mu^+\mu^-$ 



$$\blacktriangleright R(D^{(*)}) = \frac{\Gamma[B \to D^{(*)} \tau \nu]}{\Gamma[B \to D^{(*)} l \nu]}$$

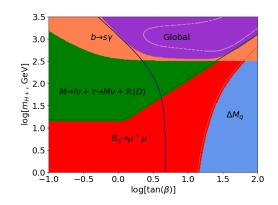


- $\triangleright$  2HDM historically struggles fitting both  $\mathcal{R}(D)$  and  $\mathcal{R}(D^*)$
- ➤ Does this kill the 2HDM?



## Extended Global Fit

- ightharpoonup Added  $B_s 
  ightarrow \mu^+ \mu^-$  and  $\mathcal{R}(D)$
- $ightharpoonup \mathcal{R}(D^*)$  not included
- ➤ 95% CL:  $m_{H^+} > 390 \,\text{GeV}$
- ►  $1\sigma$ :  $m_{H^+} > 530 \,\text{GeV}$
- $\blacktriangleright \tan \beta \gtrsim 2$





#### 2. Global Fits

## CKM Element Modifications

- ➤ CKM Matrix contains information of quark mixing
- ► In SM, a  $3 \times 3$  unitary matrix
- Measurements would be incorrect in 2HDM; e.g. from leptonic decays, modified as

$$V_{ij} = \frac{V_{ij}^{SM}}{(1 + r_H)^2} \tag{4}$$

- ➤ Exclusive measurements from light quark mesons would have negligible change
- ightharpoonup Possibility to improve unitarity constraints, e.g. second row currently sums to >1
- ➤ Space for a fourth generation?



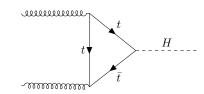
## Four Generations?

### Why SM4?

- ▶  $3 \times 3$  CKM  $\rightarrow 4 \times 4$ ; 1 CP-violating phase to 3
- $\blacktriangleright$  New heavy quarks, t' and b', extra loop diagrams to change decay widths
- ➤ A simple extension, and no reason for 3, so why not 4?

#### Exclusion of Chiral SM4:

- ► Light neutrinos measured precisely as  $N_{\nu} = 3$
- ➤ SM4 gluon fusion  $\approx$  9 times SM3 gluon fusion from heavy quarks



# SM4 with 2HDM Type II

- ➤ A Chiral SM4 not excluded with 2HDM
- ➤ Introduce a "wrong sign" limit to cancel new Higgs couplings

$$\kappa_{u} = -\kappa_{d,l} \tag{5}$$

For  $\tan \beta \gg 2$ , the wrong sign limit yields

$$\cos(\beta - \alpha) = \frac{2}{\tan \beta}, \qquad \cos(\beta - \alpha) = \sin 2\beta$$
 (6)

- ightharpoonup These are equivalent in the large an eta limit
- ➤ Relations allow us to reduce free parameters of SM4×2HDM
- ➤ Aim to extend 2HDM scans to SM4 and constrain model using flavour observables



#### 4. Questions

# Any Questions?

