

AES 770 Satellite Remote Sensing II Project III

Smoke and Cloud Radiative Forcing

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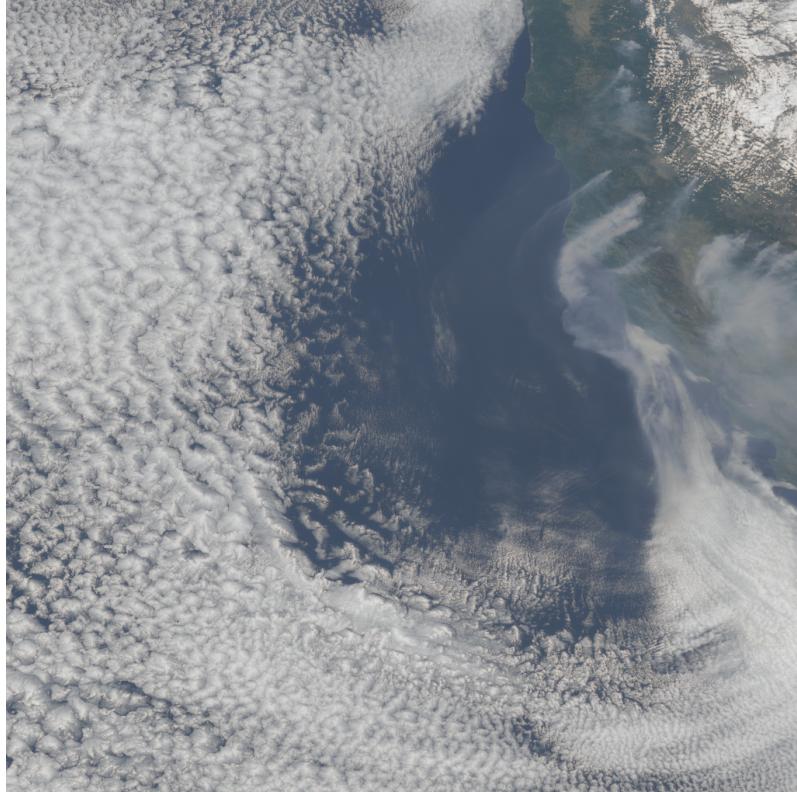


Figure 1: GOES-17 ABI Truecolor RGB of the Region of Interest, captured on August 19, 2021 at 2137z. The scene shows widespread stratocumulus cover as well as thick wildfire smoke over the ocean, which is typical for the observations used in this report.

1 Abstract

In this report, I provide an assessment of the Cloud Radiative Forcing (CRF) and Aerosol Radiative Forcing (ARF) in both the longwave and shortwave spectral ranges using L2 Single-Scan Footprint (SSF) data from the Clouds and Earth Radiant Energy System (CERES) instrument, which includes imager data from the Moderate Resolution Imaging Spectroradiometer (MODIS) instrument on the same satellite platforms. I focus on a Region of Interest (RoI) spanning $135^{\circ}W$ to $120^{\circ}W$ and $30^{\circ}N$ to $45^{\circ}N$, and the data I include consists of daytime overpasses from both the Terra and Aqua satellites between August 15 and August 25, 2021. Over land, the cloud forcing is shown to be about $CRF = -50 \text{ W m}^{-2}$, or about -11 W m^{-2} per unit cloud optical depth (COD), and the total aerosol forcing is approximately $ARF = -10 \text{ W m}^{-2}$, and about the same amount per unit aerosol optical depth (AOD). The forcing effects over the ocean are demonstrated to be much more substantial, with $CRF = -190 \text{ W m}^{-2}$, or -23 W m^{-2} per unit COD, and $ARF = -19 \text{ W m}^{-2}$, or -76 W m^{-2} per unit AOD.

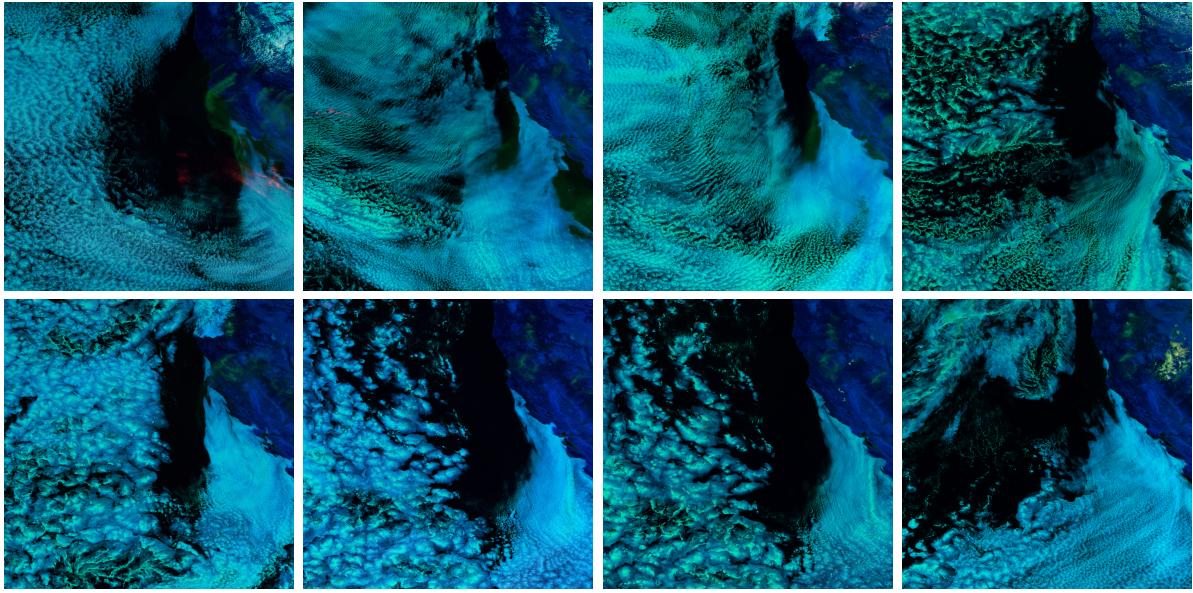


Figure 2: GOES ABI Day cloud Phase RGBs captured at a subset of the Terra and Aqua overpass times used in this report, showing the extent of wildfire smoke and prevalence of marine stratocumulus clouds between August 18 (top left) to August 25 (bottom right). Light blue indicates warm water clouds, green hues suggest fine and uniform smoke particles, and red indicates cirrus cloud cover.

2 Introduction

Clouds and aerosols both play a significant role Earth’s radiation budget, but remain two of the greatest sources of uncertainty in modeling climate sensitivity [1]. Clouds are difficult to characterize for several reasons including that they typically have conflicting longwave and shortwave effects that depend on the type and the temperature of both a cloud and the surface it occludes. Additionally, the feedback mechanisms that affect changes in the global distribution of clouds with respect to changing temperatures remain unclear [2]. Aerosols also pose challenges to modeling because (1) their type and distribution is highly regionally and contextually variable, (2) they can facilitate cloud nucleation (the so-called indirect effect), and (3) because they can enhance the evaporation rate of a juxtaposed cloud layer due to warming by absorption (semi-direct effect) [3]. What’s more is that in both cases there are no reliable counterfactual observations for comparing results to a pre-industrial atmosphere, which makes the anthropogenic contribution to climate forcing difficult to quantify. As such, it’s especially important to employ direct observational data of clear skies, aerosols, and full-sky outgoing fluxes with broadband instruments like the Earth Radiation Budget Experiment (ERBE) constellation, and more recently CERES, which I employ in this report in order to estimate the radiative forcing due to clouds and aerosols over land and ocean surfaces.

The setting for this study is a region covering the coast of central California over a 10-day period during a significant wildfire event. This is an appropriate choice for studying the radiative impact of clouds and smoke aerosols from biomass burning because the characteristic morning sea breeze consistently draws smoke aerosols out over the ocean, and facilitates the development of low-level stratocumulus clouds with a wide variety of optical depths. I restrict the data I use to only the daytime overpasses by Terra and

Aqua, which are around 11:00am (19z) and 1:00pm (21z), respectively.

For context, the 2021 California wildfire season was one of the worst on record, following a persistent drought that lasted throughout 2020 and until at least September of 2021. During this period of time, regional precipitation totals were the lowest since at least 1895 [4], which culminated in 7,396 wildfires burning over 2,569,000 acres according to the California Wildfire Incident Archive [5]. Two of the most significant wildfire events were the Dixie and Caldor wildfires, which burned 963,309 and 221,835 acres respectively, mutually forced the evacuation of more than 29,000 residents, and were both the first known fires to cross the divide of the Sierra Nevada range [6][5]. The Caldor fire began conflagrating the High Sierra just West of Lake Tahoe on August 14 and rapidly grew over the following days, while the Dixie fire had already been burning further North since mid-July at the time of this study. These two forest fire events contributed much of the smoke considered in this study, and I personally witnessed the devastation and subsequent re-growth following the Caldor fire while backpacking the area in June 2023.

3 Background: Satellites and Earth's Radiation Budget

3.1 Early Satellite Retrievals of Outgoing Radiation

The first satellite observations of Earth's radiation budget were acquired by the low-resolution thermistor-based radiometers on board the TIROS-generation satellites between 1962 and 1966, and later by the similar ESSA III platform in 1966 and 1967 [7][8]. The instruments simultaneously directly measured incoming solar irradiance, as well as rough outgoing broad-band radiance in the longwave and shortwave ranges [9][10]. Early measurements from these platforms concluded in an estimated global mean albedo of 30%, and a global average longwave emission of 230.3 W m^{-2} [9][11].

3.2 Earth's Radiation Budget Experiment

In the 1978, NASA launched the Nimbus 7 satellite carrying the first generation of the Earth's Radiation Budget (ERB) instrument, which enabled the development of a new methodology for retrieving accurate outgoing radiation estimates. This new strategy involves the co-location of narrow-FOV Advanced Very High Resolution Radiospectrometer (AVHRR) data with wide-FOV ERB data, and subsequently the use of angular distribution models (ADMs) to characterize surfaces and estimate the full-sky radiative flux of a surface from a radiance observation (with a particular Sun/surface/satellite geometry) [12][13]. The success of this mission was soon followed by the launch of 3 similar instruments, which flew aboard the NOAA-9, NOAA-10, and ERBS satellite platforms, which collectively form the Earth's Radiation Budget Experiment (ERBE) constellation [14].

The advances in observational capacity as well as analysis techniques following the ERBE experiment facilitated the development of regional and seasonal estimates of Earth's radiation budget, which substantially enhanced understanding of the shortwave and longwave global energy balance [15]. In particular, the juxtaposition of data from imagers having finer spatial and spectral resolutions enabled the scenes observed by broadband instruments to be spectrally classified into surface types with common characteristics. This advancement led to the development of surface-specific bidirectional models (ADMs) for more accurate flux retrievals, as well as derivations of the radiative forcing by clouds as opposed to clear-sky [2][16]. The improved observations led to an updated

average planetary albedo estimate of 31%, and an outgoing longwave radiation average of 235 W m^{-2} [17].

3.3 Clouds and the Earth's Radiant Energy System

In 1997, NASA and the Japanese Aerospace Exploration Agency (JAXA) began collecting data with a new generation of broad-band instruments known as Clouds and the Earth's Radiant Energy System (CERES) launched on the Tropical Rainfall Measuring Mission (TRMM). This first instance of the instrument was limited in that TRMM didn't have global coverage, and CERES soon failed in August of 1998. Nonetheless, the NASA launched new versions of the instrument in 1999 and 2002 with the Terra and Aqua satellites as part of the Earth Observing System (EOS), which have proven extremely successful missions that presently continue to contribute valuable science data more than 20 years later.

CERES is a 3-channel broadband radiometer that improves on ERB's measurement accuracy by up to a factor of 3, and offers a reduced sample FoV of about $16 \times 32 \text{ km}$. The three channels supported by CERES include $.2 - 5\mu\text{m}$ shortwave-range, $.2 - 100\mu\text{m}$ full-spectrum, and $8 - 12\mu\text{m}$ infrared window channels. Designers of the instrument neglected to include an independent longwave channel as a means of decreasing observation error due to material shortcomings, but accurate longwave estimates can be obtained by simply subtracting the shortwave instrument observations from those of the full-spectrum. The window channel is useful for estimating clear-sky surface fluxes as well as estimating the greenhouse effects of species with strong absorption outside the atmospheric window [18].

In addition to the CERES instrument's improvements over ERB, the EOS satellite platforms include Moderate Resolution Imaging Spectroradiometer (MODIS) imagers with a wealth of narrow-band channels co-located with the CERES observations. These substantially enhance the quality of spectral classification and material retrievals from within a CERES FoV (hereafter referred to as a footprint). These help supplement CERES data with high-quality information about surface types and the properties/distribution of atmospheric materials. Additionally, Terra and Aqua include a second CERES instrument (in addition to the typical one scanning in a cross-track mode) which spins azimuthally in order to contribute near-simultaneous observations of surfaces from a second viewing angle. Data from the second instrument is critical for producing accurate and empirical ADMs.

4 Converting CERES Radiance to TOA Flux

4.1 Filtered to Unfiltered Radiance

The first step to determining the TOA flux from broadband observations of radiance by the CERES instrument is to remove the effect of the instrument's relative spectral response function (RSR). As with any physically realistic radiometer, the digital counts recorded by the instrument are somewhat spectrally dependent on the wavelength of incident light. Since radiation budget experiments seek to observe the actual magnitude of energy associated with outgoing radiation, as opposed to relative magnitudes that are often sufficient for imager applications, it is critically important to correct for the error caused by the instrument's imperfect spectral sensitivity.

In general, the observed radiances can be thought of as a filter convolution of the RSR function over the actual surface's radiance, which is a fundamentally information-

lossy process. Thus, “undoing” this effect is necessarily an approximation. In the case of CERES, this involves estimating so-called *unfiltered* radiance in each channel in terms of empirically-derived coefficients, which form linear combinations of the filtered radiances. These coefficients depend on each footprint’s surface class as well as the viewing geometry, and are sometimes be periodically updated throughout the instrument’s lifetime [19].

4.2 Angular Distribution Models

As mentioned in the previous section, one of the major advancements that enables accurate characterization of top-of-atmosphere fluxes are angular distribution models, or ADMs, which seek to represent the dependence of the observed radiance of a particular scene type on the viewing geometry. Broad-band imagers on satellite platforms can only ever observe *radiances* in $\text{W m}^{-2} \text{sr}^{-1}$ since they necessarily have a specific position with respect to the surface they are observing. This poses a challenge for the retrieval of the full-sky radiant flux of a surface, because the absorbed or emitted radiation may be strongly anisotropic. This challenge is even more prominent for instruments with a relatively large footprint such as CERES because the incident surface type may be highly variable. As such, there is a need for careful modeling of the dependence of observed radiance on viewing angle.

$$F(\theta_0) = \int_0^{2\pi} \int_0^{\frac{\pi}{2}} I(\theta_0, \theta, \phi) \cos(\theta) \sin(\theta) d\theta d\phi \quad (1)$$

Equation 1 expresses the full-sky TOA flux $F(\theta_0)$ given solar zenith angle θ_0 with respect to radiance I , viewing zenith angle θ , and viewing azimuth angle ϕ . Once a surface type has been identified for a footprint (typically using spectral/spatial classification of co-located imager data), the task of that surface’s ADM is to provide a discrete estimate of the above integral in terms the footprint’s (θ_0, θ, ϕ) configuration. This estimate takes the form of a lookup table value known as an *anisotropic factor* R_s .

$$R_s(\theta_0, \theta, \phi) \approx \pi \frac{\bar{I}_s(\theta_0, \theta, \phi)}{\bar{F}_s(\theta_0)} \Rightarrow F_s(\theta_0) \approx \pi \frac{I_s(\theta_0, \theta, \phi)}{R_s(\theta_0, \theta, \phi)} \quad (2)$$

Equation 2 shows how the anisotropic factor R_s for surface s is estimated, and subsequently how it can be used to approximate the full-sky flux. On the left hand side, the anisotropic factor is approximated as the ratio of the mean isotropic radiance $\pi \bar{I}_s(\theta_0, \theta, \phi)$ within an angular bin to the surface’s average full-sky-flux $\bar{F}_s(\theta_0)$. Since the fundamental problem here is that $\bar{F}_s(\theta_0)$ can’t be directly observed, its value is estimated using large ensembles of observations, which will ideally converge on reliable values per angular bin, if surface classes are selected well. Once lookup tables of anisotropic factors have been empirically determined for a surface, they are applied to scale actual radiance observations $I_s(\theta_0, \theta, \phi)$ into estimates of TOA flux $F_s(\theta_0)$ [20][21][22].

An additional consideration that must be taken into account while constructing angular dependence models is the TOA reference level used when modeling outgoing fluxes. This is not normally a consideration that must be made for imager data, which is supposed to correspond to radiance emerging directly from the observed surface, however for large-footprint flux calculations it becomes important to consider the influence of radiance corresponding to secant “slant paths” through the atmosphere but above the observed surface. Changing the reference height essentially just modifies the surface area over which outgoing radiation is presumed to be distributed, however for CERES-sized

footprints, this can be responsible for an error of a few W m^{-2} of flux associated with the observed surface. The ideal TOA reference height determined for the CERES ADMs is $z_{TOA} = 20 \text{ km}$, as identified by (Loeb, 2002)[23].

- | | | |
|--------------------------------|---------------------------------|-------------------------|
| 1. Evergreen Needleleaf | 7. Open Shrubland | 15. Permanent Snow/Ice |
| 2. Evergreen Broadleaf | 8. Woody Savannas | 16. Exposed Soil/Rock |
| 3. Deciduous Needleleaf | 9. Savannas | 17. Water Bodies |
| 4. Deciduous Broadleaf | 10. Grassland | 18. Tundra |
| 5. Mixed Forest | 11. Permanent Wetland | 19. Fresh Snow |
| 6. Closed Shrubland | 12. Cropland | 20. Sea Ice |
| | 13. Urban/Industrialized | |
| | 14. Cropland Mosaic | |

The ADMs used for retrieving fluxes from CERES data are separated into the 20 surface types shown above, of which only the bolded items are considered in this report. Although there are some footprints in the RoI identified as urban/industrialized or fresh snow in the data I acquired, they did not meet the thresholded requirements for consideration in my energy budget calculations.

4.3 Co-location of Imager Data

$$\bar{x} = \frac{\int_{FOV} P(\delta, \beta) x(\delta, \beta) \cos(\delta) d\beta d\delta}{\int_{FOV} P(\delta, \beta) \cos(\delta) d\beta d\delta} \quad (3)$$

As mentioned in previous sections, an accurate assessment of the TOA flux at CERES' footprint resolution requires careful co-location of spatially and spectrally finer-resolution pixel data from the on-board MODIS radiometer. CERES' sensitivity to radiance within a footprint is not constant throughout the footprint area due to the nonzero exposure time of a pixel, the translational motion of the satellite, as well as inherent instrument properties. Thus, in order to appropriately estimate the influence of the properties retrieved from imager data on the flux observations by CERES, the sensor's point spread function (PSF) $P(\delta, \beta)$ must be considered with respect to orthogonal viewing angles δ and β within a single footprint. Equation 3 is the continuous form of the expression used to estimate the PSF-weighted average value of an imager pixel property \bar{x} within a footprint's *FOV*. This spatial averaging technique is used in the generation of CERES SSF data in order to produce full-footprint averages of cloud/aerosol optical depth, percent coverages of incident surface types, estimates of cloud top temperature/pressure, and other useful metrics derived from MODIS radiances.

5 Forcing Theoretical Basis

$$H = S(1 - \alpha) - F_{LW} = H_{SW} - F_{LW} \quad (4)$$

The climate forcing attributable to an atmospheric component such as clouds or aerosols is a function of the material's influence on incoming shortwave solar radiation, as well as its effect on emitted longwave radiation. As such, in the most general form, the net radiative heating of an atmospheric column can be expressed in terms of the difference between the heat from absorbed shortwave radiation within the column H_{SW} and the column's outgoing longwave flux emission F_{LW} , as demonstrated by Equation 4.

6 Cloud and Aerosol Forcing Results

Since clouds and aerosols are transient phenomena with radiation effects that depend on other properties such as optical depth, single scatter albedo, and effective radius, it's often helpful to express their forcing effect in terms of the variation in the net outgoing flux of an atmospheric column containing the cloud or aerosol from that of the same atmospheric column without the cloud or aerosol. Given a clear-sky outgoing flux F as shown above, and the flux given the presence of a cloud $F^{(cld)}$, the total climate forcing as a consequence of the cloud (i.e. cloud radiative forcing) is determined to be $CRF = F - F^{(cld)}$. Similarly, given the outgoing flux of an aerosol-laden atmospheric column $F^{(aer)}$, the aerosol radiative forcing is $ARF = F - F^{(aer)}$.

The above expressions for CRF and ARF are structured such that positive values are associated with a net heating effect from clouds or aerosols, and negative values indicate net cooling within the atmosphere. It's also critical to note that the clear-sky flux F in both cases is highly dependent on the incident surface, and even the relative heating rates in the presence of clouds or aerosols can depend on the surface type in a nonlinear fashion due to feedback mechanisms. Therefore in this project I calculate the atmospheric forcing effects of clouds and aerosols independently per surface type observed within the RoI.

$$\begin{aligned} CRF &= F - F^{(cld)} = (F_{SW} + F_{LW}) - (F_{SW}^{(cld)} + F_{LW}^{(cld)}) \\ CRF &\approx (F_{SW} + F_{LW}) - \left(F_{SW} + \frac{dF_{SW}}{d\tau} \cdot \tau + F_{LW} + \frac{dF_{LW}}{d\tau} \cdot \tau \right) \\ &= -\tau \cdot \left(\frac{dF_{SW}}{d\tau} + \frac{dF_{LW}}{d\tau} \right) \\ \frac{dCRF}{d\tau} &= \frac{dF_{SW}}{d\tau} + \frac{dF_{LW}}{d\tau} \end{aligned} \tag{5}$$

Since CERES directly observes outgoing fluxes in the shortwave and longwave spectrum, and because the scene includes clouds and aerosols with a wide variety of properties that influence their radiative characteristics, a more appropriate expression for the measured radiative forcing is in terms of the change in shortwave and longwave flux with respect to a determining property. For example, instead of modeling CRF in terms of cloudy and clear-sky heating rates, Equation 5 employs shortwave and longwave components of both cloudy and clear-sky outgoing fluxes, assumes a linear relationship between COD (τ) and the outgoing flux in each waveband, and cancels out the clear-sky fluxes to obtain CRF as a function of the rate of change of F_{SW} and F_{LW} with respect to COD. The same principle applies to aerosol radiative forcing. For the purposes of this project, I will maintain the admittedly weak assumption of this linear relationship, and calculate CRF and ARF in terms of the first-order rates of change with respect to COD and AOD.

6 Cloud and Aerosol Forcing Results

6.1 Methodology

The CERES SSF dataset I used for this assignment contains a wealth of information associated with each footprint, which makes it straightforward to set thresholds based on sub-footprint imager data and satellite geometric configuration, and to independently consider surfaces based on their type class and consistency. Aerosol-cloud interactions and the radiative forcing over highly heterogeneous surfaces are out of the scope of this project, so I restricted the footprints I used to exclude those with aerosols and clouds

6 Cloud and Aerosol Forcing Results

in close proximity, those with any high-level cirrus contamination, and those containing inconsistent surface types.

The following constraints were imposed on the footprints used in each flux forcing calculation, based on the auxilliary data provided alongside the SSF flux data, as defined in (Loeb, 2018) [24].

Cloud Optical Depth Constraints

1. At least 80% of imager pixels identified as the dominant surface class
2. Layer 2 (cirrus) cloud optical depth less than 1
3. Standard deviation of imager pixel cloud optical depth less than 10
4. Upper bound of 2% of imager pixels identified as containing aerosols

Aerosol Optical Depth Constraints

1. At least 80% of imager pixels identified as the dominant surface class.
2. Layer 2 (cirrus) cloud optical depth less than 1
3. At least 98% of pixels marked cloud-free by MOD04 cloud mask.

As an exception to the restrictions on dominant surface class, due to the scarcity of some valid samples and similarity in their observed radiative forcing curves I merged (1) Evergreen Needleleaf and Deciduous Broadleaf classes into a single class for all Forests, (2) Savannas, Woody Savannas, and Grassland classes into a single class for Savanna, and (3) Cropland and Cropland Mosaic classes into a single class for Crops.

6.2 Cloud Radiative Forcing

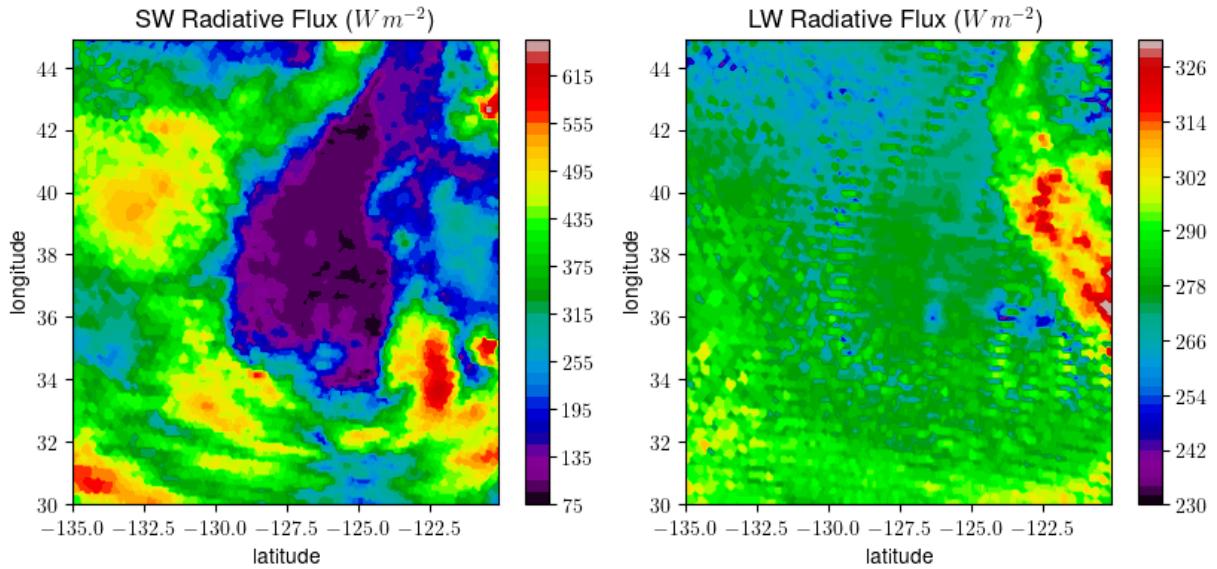


Figure 3: Outgoing Shortwave and Longwave fluxes observed by CERES on August 18, 2021 at 1910z (same swath as Figure 1). The $\sim 16 \times 32$ km footprints are resampled onto a uniform 0.1° grid by nearest-neighbor interpolation.

Clouds have a substantial influence on the radiation budget due to their high reflectivity as well as emissivity. In the shortwave range, enhanced reflection of insolation causes a

6 Cloud and Aerosol Forcing Results

general cooling effect that is most prominent in the mid-latitude summer, and over bodies of water [25]. This effect is clearly demonstrated by the broad-band shortwave flux in Figure 6.3, which shows that outgoing shortwave radiation over the highest-reflectivity water clouds exceeds that of the ocean surface by as much as 540 W m^{-2} .

The longwave impact of clouds is less prominent for liquid water clouds over the ocean since in the observed scenarios both surfaces have a comparable temperature and emissivity, however over land (in the top right of the second image in Figure 6.3), there are clear depressions in outgoing longwave radiation (OLR) associated with low-level water clouds. Furthermore, the thin cirrus clouds associated with red hues in the center-right of the first image in Figure 1 correspond to a decrease in OLR over the same area in Figure 6.3. Both of cases are indicative of the infrared greenhouse effect of clouds which have a net warming effect that is generally less prominent than the SW forcing [2].

Surface	F_{SW}	$F_{SW}^{(cld)}$	F_{LW}	$F_{LW}^{(cld)}$	$\frac{dF_{SW}}{d\tau}$	$\frac{dF_{LW}}{d\tau}$	$\frac{dCRF}{d\tau}$	CRF
Ocean	112.606	301.863	279.217	280.211	23.583	-0.129	-23.454	-190.251
Forest	126.681	190.923	301.541	291.280	18.476	-2.951	-15.525	-53.981
Crop	181.188	249.028	307.006	292.812	16.040	-3.356	-12.684	-53.646
Savanna	180.448	224.008	315.442	309.667	12.353	-1.638	-10.715	-37.785
Shrub	150.317	219.108	321.178	300.303	18.966	-5.755	-13.211	-47.916
Urban	144.108	197.317	302.574	298.587	12.388	-0.928	-11.460	-49.222

Table 1: Cloud radiative forcing results from the regression technique, corresponding to the trend lines for each surface in Figures 5 and 6. $F_{SW}^{(cld)}$ and $F_{LW}^{(cld)}$ are determined using the average flux of all pixels identified as cloudy using the thresholding described in the previous section.

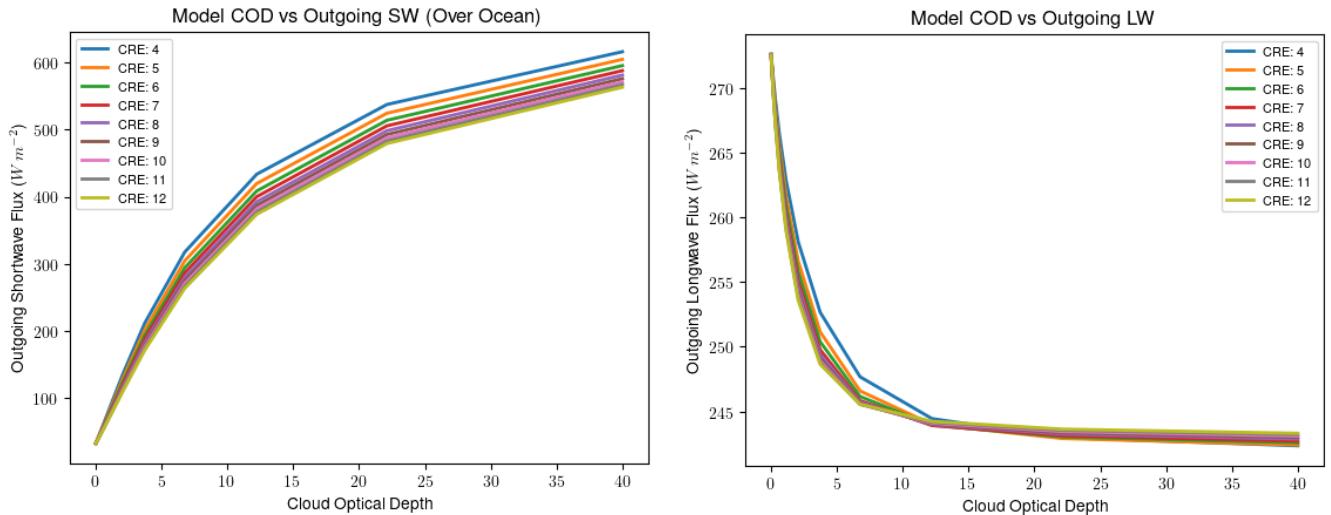


Figure 4: Model-derived dependence of total TOA shortwave and longwave flux with respect to cloud optical depth (COD) and cloud effective radius (CRE) over an ocean surface. Values were calculated using the SBDART radiative transfer model assuming an ocean surface temperature of 300 K, a mid-latitude Summer atmosphere, and a solar zenith angle of 20° . Shortwave flux dependence on solar zenith angle is also modeled in Figure 10.

6 Cloud and Aerosol Forcing Results

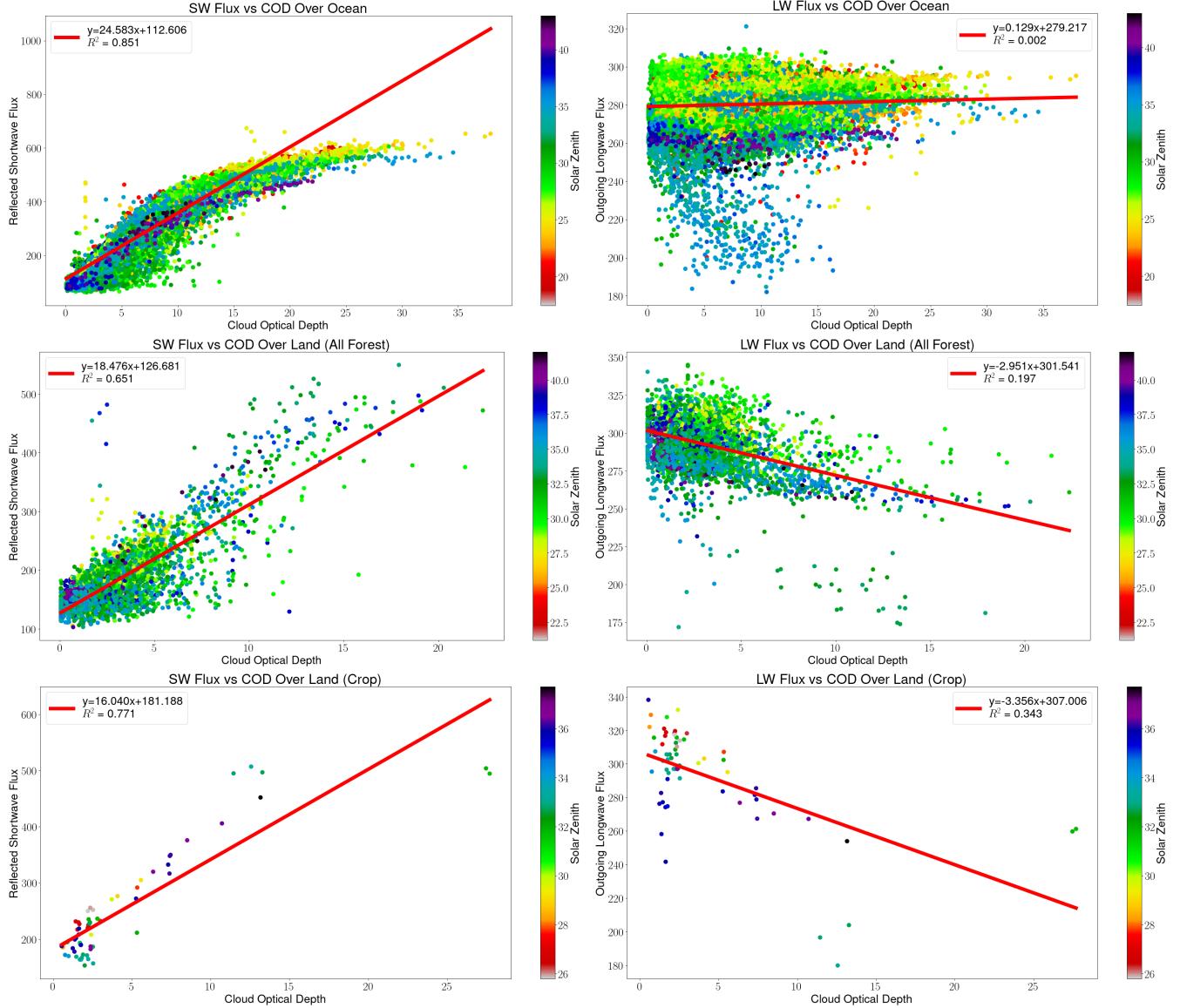


Figure 5: CERES-observed shortwave (left) and longwave (right) dependence of ocean (top), forest (middle), and crop land (bottom) surface classes on cloud optical depth. Includes all Terra and Aqua CERES footprints within the RoI that meet the above constraints between August 15 and August 25, 2021.

6 Cloud and Aerosol Forcing Results

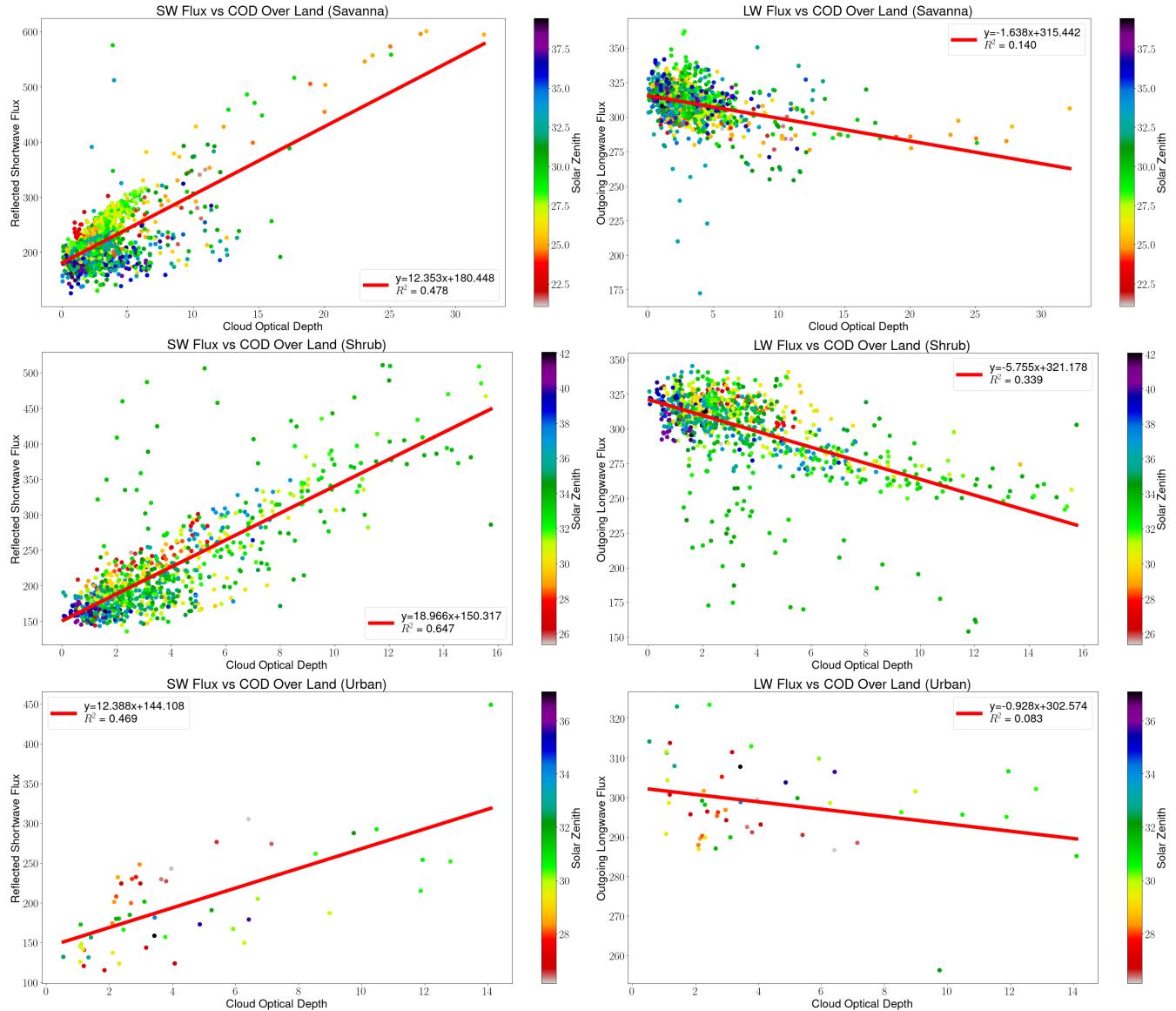


Figure 6: SW (left) and LW (right) dependence of savanna (top), shrub (middle), and urban (bottom) land surface types on cloud optical depth.

6.3 Aerosol Radiative Forcing

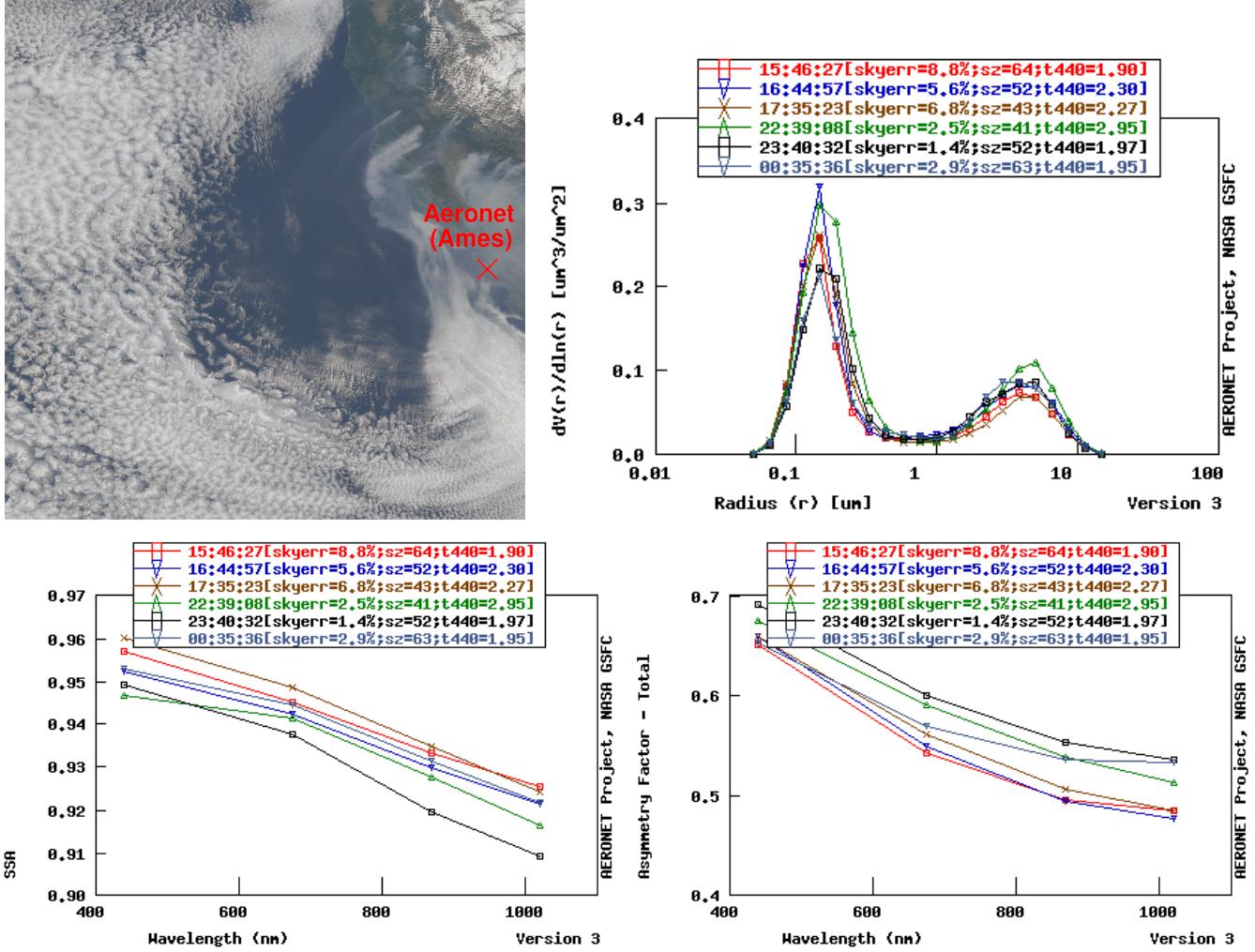


Figure 7: Top: Location of NASA Ames AERONET station, and cotemporal measurements of aerosol effective radius distribution. Bottom: Same AERONET station's observation of aerosol single scatter albedo and assymetry parameter. These observations informed the model calculations in Figure 8.

Surface	F_{SW}	$F_{SW}^{(aer)}$	F_{LW}	$F_{LW}^{(aer)}$	$\frac{dF_{SW}}{d\tau}$	$\frac{dF_{LW}}{d\tau}$	$\frac{dARF}{d\tau}$	ARF
Ocean	75.280	91.156	274.705	278.662	62.307	14.145	-76.452	-19.833
Forest	136.597	146.987	299.052	299.079	23.141	0.312	-23.453	-10.417
Savanna	180.233	190.139	324.041	321.058	15.294	-6.256	-9.038	-6.923
Shrub	156.140	170.537	331.089	327.455	22.228	-13.291	-8.937	-10.763

Table 2: Results for aerosol radiative forcing using the linear regression technique. Since the influence of outgoing longwave radiation is weakly correlated with AOD, only short-wave effects are included in the ARF.

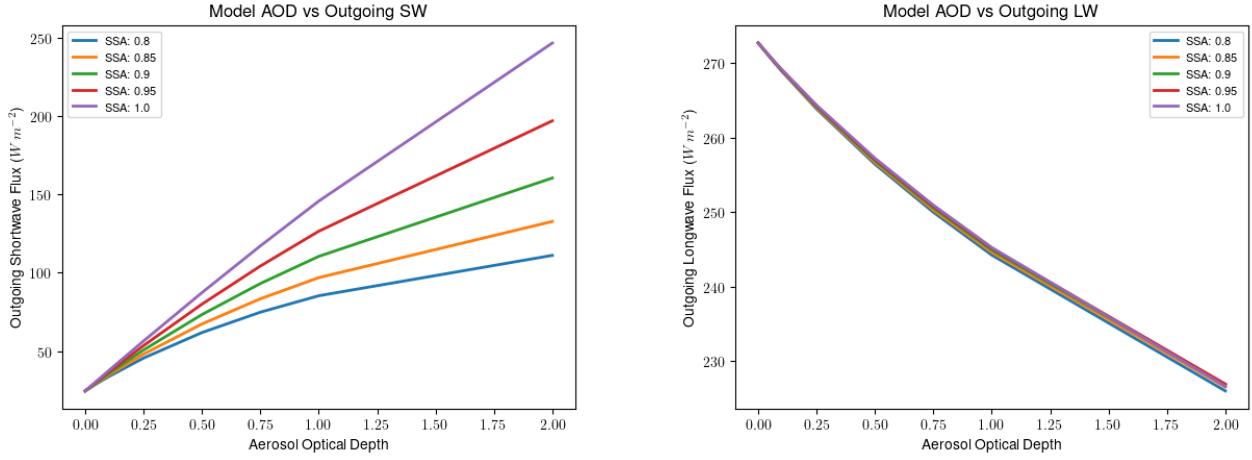


Figure 8: Model-derived dependence of total TOA shortwave and longwave flux with respect to aerosol optical depth (AOD) and single scatter albedo (SSA) over an ocean surface. Values were calculating using the SBDART radiative transfer model assuming an ocean surface temperature of 300 K, a mid-latitude Summer atmosphere, and a solar zenith angle of 20°. Shortwave flux dependence on solar zenith angle is also modeled in Figure 10.

6 Cloud and Aerosol Forcing Results

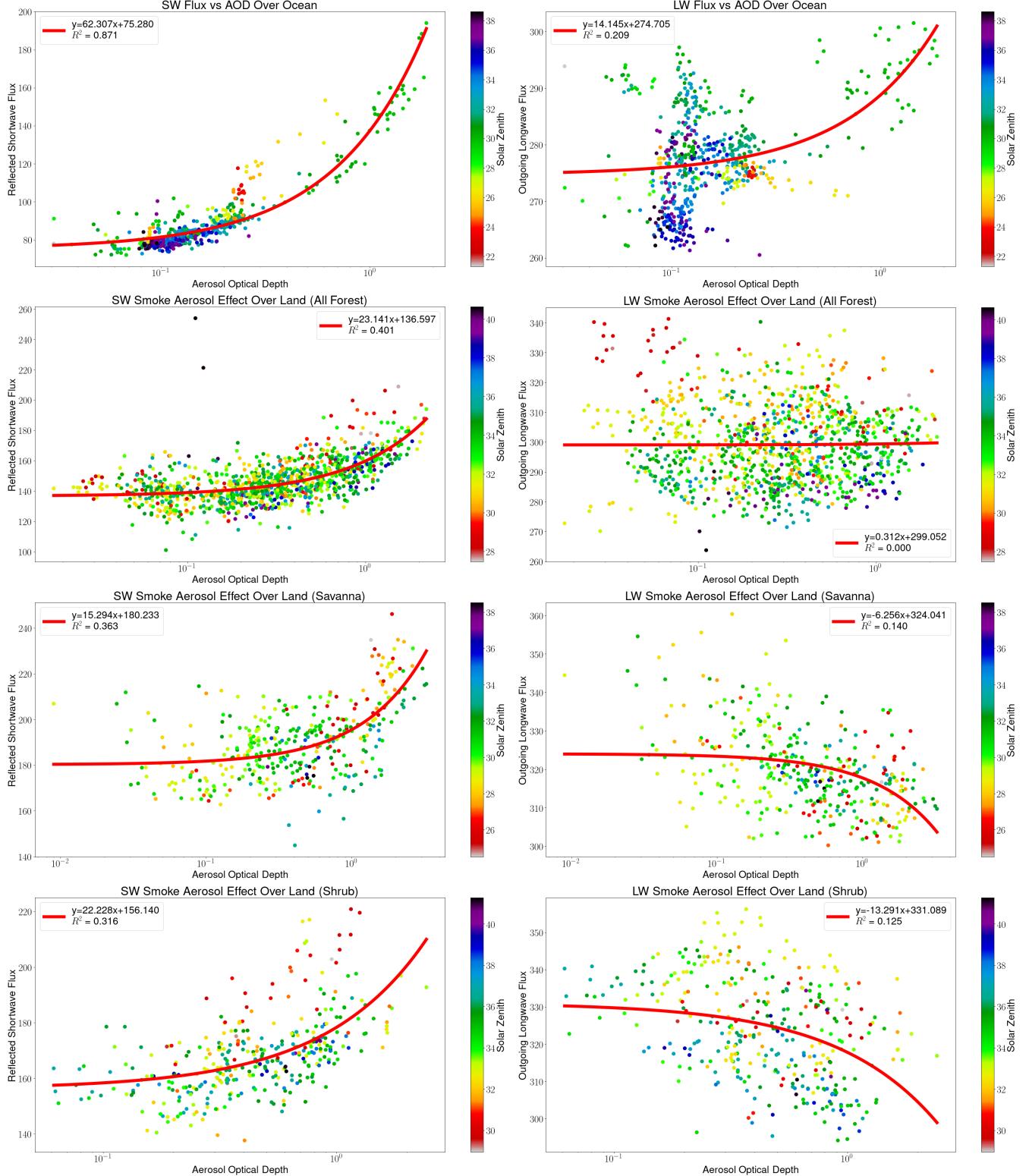


Figure 9: CERES-observed shortwave (left) and longwave (right) dependence of ocean (top), forest, savanna, and shrub land (bottom) surface classes on aerosol optical depth. Note that trend lines only appear nonlinear due to the logarithmically-scaled x-axis. Outgoing longwave radiation, especially over the ocean, exhibits clustering of solar zenith values, which suggests the trend is dominantly the consequence of surface temperature changes between overpass times.

7 Conclusion

8 Appendix

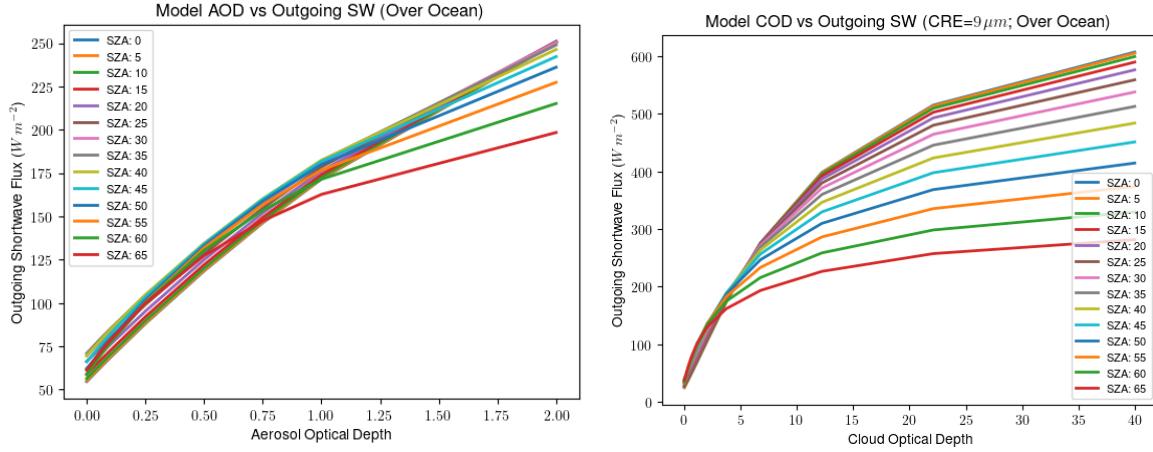


Figure 10: Model-derived dependence on solar zenith angle for shortwave TOA flux with respect to aerosol optical depth (left) and cloud optical depth (right).

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