

University of California  
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**Constraints on the Higgs boson total decay width  
from off-shell production with decay into pairs of  
Z-bosons in simulated data**

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of the requirements for the degree

Bachelor of Science  
in  
Physics

by

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May 2020

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## Abstract

Constraints on the Higgs boson total decay width from off-shell production with decay into pairs of Z-bosons in simulated data

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The discovery of the Standard Model (SM) Higgs boson at the Large Hadron Collider (LHC) was a major achievement of the experimental particle physics community in the 21st century. Though a fair portion of the physics analysis focus has shifted to Supersymmetry (SUSY) physics, the direct search of SUSY models has yield null result so far. Meanwhile, many properties of the Higgs boson are still to be measured.

In this analysis, we present constraints on the total decay width of Higgs boson,  $\Gamma_H$ , by using the on-shell and off-shell decay rates of Higgs boson to a pair of Z bosons with one Z-boson decaying into a pair of electrons or muons and the other into neutrinos. A total number of  $2$  ( $ee$  or  $\mu\mu$ )  $\times$   $4$  ( $N_{\text{jets}} = 0, 1, 2, 3+$ )  $= 8$  channels are considered and used for limits fitting. The result represent the expected constraints using the physics events and CMS experiment detector response simulated data via Monte Carlo methods (MC). The data and expected results correspond to data collected in 2016, 2017, and 2018, which has an integrated luminosity of  $137.15 \text{ fb}^{-1}$  at a center-of-mass energy of 13 TeV. The expected constraint based on simulation is  $\Gamma_H < 16.4 \text{ MeV}$  95% CL.

# Contents

<b>Abstract</b>	<b>v</b>
<b>1 Introduction</b>	<b>1</b>
1.1 Physics at the LHC (and elsewhere) . . . . .	1
1.2 The CMS Detector . . . . .	2
1.3 The Higgs boson and off-shell methods . . . . .	5
1.4 Background and signal simulation . . . . .	7
1.5 Uncertainties . . . . .	10
<b>2 Methods</b>	<b>11</b>
2.1 Event selection and physical variables . . . . .	11
2.2 Signal samples re-weighting . . . . .	15
2.3 Strategy in variable selection and binning and systematical uncertainties	19
<b>3 Results and interpretation</b>	<b>22</b>
3.1 Limits on Higgs decay width . . . . .	23
<b>A Weights Table for Higgs Sample</b>	<b>25</b>
A.1 Windowed re-weighting results . . . . .	25
A.2 Sample mass factors for ggZZ and VBF_ZZ samples . . . . .	26
<b>B Additional Figures</b>	<b>28</b>
B.1 Signal Sample Fitting Template Histogram of Background and Signal . .	29
<b>Bibliography</b>	<b>33</b>

# Chapter 1

## Introduction

This chapter contains a brief introduction of the experimental aspect of the Large Hadron Collider (LHC) physics, the CMS detector, and the physics behind the off-shell methods for constraining the Higgs decay width. Finally, a very brief technical account for the Monte Carlo (MC) events production is also presented.

~~The work presented here is part of a wider endeavor to measure the decay width as well as test for (the strength of) BSM contributions by studying the off-shell Higgs production. A pure MC based analysis is conducted here but with many of the same tools and studies needed for the final result, as well as the expected value of the final constraint (for the entire Run 2).~~

### 1.1 Physics at the LHC (and elsewhere)

After the discovery of standard model (SM) Higgs boson in 2012, the LHC went through a series of upgrades during 2013–2015 and finally restarted with center-of-mass energy reaching 13 TeV. Though the last particle promised by the SM has been discovered, there remains ‘known’ mysteries in the high energy physics.

The importance of beyond-the-standard model (BSM) physics is self-evident, though the SM is one of the most precise theories in physics we ever had, the model requires many ‘inputs’ from experiment for parametrization, and, it cannot account for phenomena such as neutrino oscillation with its original form. One proposed (and admittedly the most popular one) solution is called Supersymmetry (SUSY), which predicts the existence of more particles than what we have identified so far in the SM. Because these predicted particles have a ‘mirror’ extension/association to the SM particles, they are often addressed as SUSY partners (as each fundamental fermion in the SM would have a bosonic SUSY partner and vice versa). While the direct searches for these SUSY partners in the past few years have all yielded null results, they have not discouraged physicists from other attempts of probing the BSM physics in other ways.

Higgs boson has a special role in the process mainly due to the Higgs field, from which the particle originates from, couples to all massive particles. Thus making Higgs boson a window to probe the BSM physics, as discussed in Section 1.3.

The LHC has been operating at the same energy for the past 3 years (2016-2018) with the delivered luminosity ramping up in each year[1]. In this analysis, we will be using mostly simulation that were produced corresponding to Run period 2018, but scaled to  $137.15\text{ fb}^{-1}$  to simulate the expected result by using the entire Run 2 data, for the Compact Muon Solenoid (CMS) detector.

## 1.2 The CMS Detector

The CMS is one of the two general purpose detectors built around LHC (the other one is ATLAS). It is ‘compact’ only when comparing to the ATLAS detector and is in fact heavier than the latter. Though we are not using the data accumulated by the CMS detector, the simulated events are reconstructed based on the material interaction and

electronics responses of the real detector using the GEANT4 package[2].

The defining feature of the CMS detector is the solenoid (as in its name) around the beam line, in the middle of the detector[3]. This superconducting solenoid produces a 4 T magnetic field and is 12.5 m long free bore[3]. This feature allows the tracking of charged particles by bending their trajectory when hitting various layers of the detector.

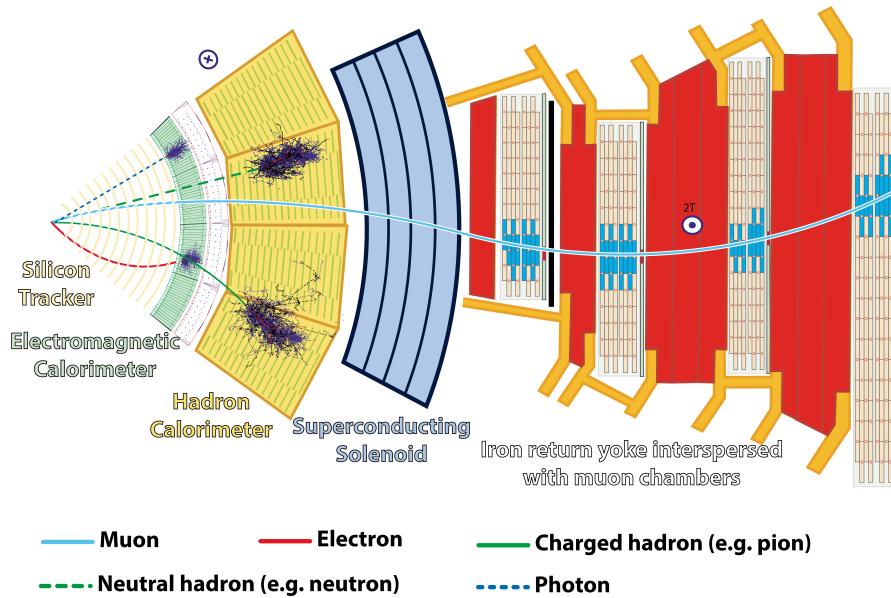


Figure 1: A section slice of the CMS detector where we can see how particles with different physical properties interact differently with the detector.<sup>1</sup>

One important kind of ‘final-state’ particles used in this thesis is charged lepton, which includes electrons and muons (the half-life of tau-on is too short to interact with the detector directly); ‘final-state’ here means they are the end products of a decay chain and are recorded by the detector components directly. If we concentrate on the red and blue lines in Fig. 1, we can see that **for the** electrons (red) leave a few hits that resemble a curved track within the silicon tracker and are stopped at the electromagnetic calorimeter (ECAL).

<sup>1</sup><https://cds.cern.ch/record/2120661>

For the muons (blue line in Fig. 1), they largely go through all the interior layers unhinged due to their higher mass [and the fact that they only interact electromagnetically \(unlike a proton, for example\)](#); pay close attention to the ‘S’-shaped curve, this is due to the opposite magnetic field outside of the superconducting solenoid. The iron return yoke also stops particles other than muons coming through (except neutrinos, they always travel freely). Although information such as momentum relies on hits on the silicon trackers, the hits in the muon chambers and a matching trajectory are needed for the object to be reconstructed as a muon. Notice how the muon chambers occupy more than half of the detector by size, such design allows the CMS detector to excel in muon measurements and is the primary reason for the overall design.

Zooming out from the individual components of the detector, the complete (physics object) reconstruction process is complex and sometimes requires information from multiple parts of the detector, this algorithm is called the Particle Flow (PF)[4]. While the energy of simple and neutral objects such as photon is directly measured by the ECAL (with correction such as zero-suppression, where minimal activities are treated as zero value during readout), electron measurements need the information from the inner tracker (to determine charge and momentum), energy at the ECAL, and also sum of the energy of photons produced by bremsstrahlung (when electron interacts with the detector material) compatible with the electron’s track.

For gluons and quarks that come out of the interaction vertex, due to quark confinement, the detector is only be able to ‘see’ a narrow ‘spray’ of final state (stable) particles that, subsequently, are clustered into objects called ‘jets’. There are different algorithms and criteria to combine a collection of measurements into a single jet, in the CMS at the moment, the anti- $k_t$  clustering algorithm [5, 6] is used. In short, the algorithm would cluster objects in a roughly cone-like space originated from the vertex according to some parameters. The algorithm is resilient to QCD effects that would cause jets to split, such

as shown in Fig. 2

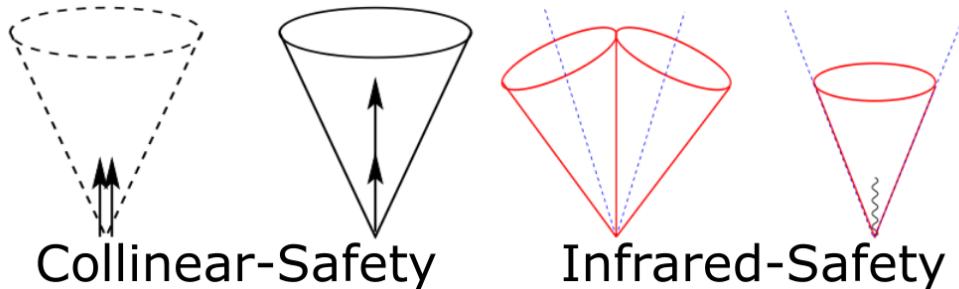


Figure 2: A illustration of two kinds of QCD effects in jets<sup>2</sup>

In this analysis, we use the AK4 jets which means that the algorithm is given a  $R = 0.4$  parameter, another common choice is  $R = 0.8$  which is preferred in some cases for boosted topologies due to a larger ‘opening angle’.

### 1.3 The Higgs boson and off-shell methods

As mentioned above, the Higgs boson has a special place as it can be seen as a ‘bridge’ between the SM and the BSM (in some models) as some SUSY particles can decay into it or decayed from it, or simply have a BSM production Feynman diagram that leads to more a higher Higgs bosons production rate than the SM expects. In any of these cases, the basic properties of the SM Higgs boson remain important.

The importance of studying the Higgs also comes from the fact that it is an ‘add-on’ to the SM as a mass generation mechanism for gauge bosons. The Higgs sector of the SM is not well tested, and, in theory could be more complicated than the minimal SM suggests, for example, multiple Higgs bosons being present, or the observed Higgs being a composite particle. Interestingly, if not composite, the observed Higgs boson would be

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<sup>2</sup>[https://twiki.cern.ch/twiki/bin/viewfile/Sandbox/Lecture?rev=1;filename=Philipp\\_Schieferdeckers\\_Lecture.pdf#page=3](https://twiki.cern.ch/twiki/bin/viewfile/Sandbox/Lecture?rev=1;filename=Philipp_Schieferdeckers_Lecture.pdf#page=3)

the only spin zero fundamental particle we know.

Many properties are well measured, such as its mass and spin [7]; others, however, have not entered the realm of ‘precision’ physics so far. One of them is the total decay width of the Higgs boson, which is of course associated to the particle’s half-life  $\tau$  as  $\Gamma = \frac{h}{\tau}$ . The SM predicts that the Higgs boson to have a decay width of 4.07 MeV. The problem is that the energy resolution of the CMS detector, which is around  $\mathcal{O}(1)$  GeV for the di-photon or 4-leptons final states, is not even remotely close to that prediction.

One way suggested by the theorists so far is to indirectly measure this decay width using the off-shell method [8, 9]. When a Higgs boson decays into two massive vector bosons (VV) (in the scope of this thesis, two Z bosons), one of the them has to be off-shell, or the Higgs boson is off-shell, this is because  $m_H \approx 125$  GeV is smaller than the mass of two W-bosons ( $\approx 160$  GeV) or of two Z-bosons ( $\approx 182$  GeV). The Higgs boson decay branching ratio is coupled to the daughter particles’ mass, and thus this cross section  $\sigma_{H \rightarrow VV}$  is enhanced as the mass of Higgs boson gets closer to the on-shell VV mass.

In fact, the production of Higgs boson from a pair of vector boson is also related to the decay width of Higgs boson  $\Gamma_H$  via the ‘propagator’[10]. In terms of the differential cross section:

$$\frac{d\sigma_{vv \rightarrow H \rightarrow VV}}{dq_H^2} \sim \frac{g_{vvH}^2 g_{HVV}^2}{(q_H^2 - m_H^2)^2 + m_H^2 \Gamma_H^2} \quad (1)$$

where the two  $g$  on the right-hand side are the couplings for production from vv and decay to VV respectively. Integrating this equation near the on-shell mass of the SM Higgs boson or in the tail region (above the mass of VV), one can translate the differential

cross section into event rate that can be (in theory) measured in experiments:

$$\begin{aligned} N_{\text{vv} \rightarrow \text{H} \rightarrow \text{VV}^*}^{\text{on-shell}} &\sim \frac{g_{\text{vvH}}^2 g_{\text{HVV}}^2}{m_{\text{H}} \Gamma_{\text{H}}} \sim \mu_{\text{vvH}} \\ N_{\text{vv} \rightarrow \text{H}^* \rightarrow \text{VV}}^{\text{off-shell}} &\sim \frac{g_{\text{vvH}}^2 g_{\text{HVV}}^2}{(q_{\text{H}}^2 - m_{\text{H}}^2)^2} \sim \mu_{\text{vvH}} \cdot \Gamma_{\text{H}} \end{aligned} \quad (2)$$

The key takeaway is that, up to some correction factors, the event rate of off-shell Higgs scales linearly respect to the Higgs decay width  $\Gamma_{\text{H}}$  which allows indirect measurement of the width itself.

In this thesis, we will focus on the  $\text{H} \rightarrow \text{ZZ} \rightarrow 2\ell 2\nu$  channel at the CMS detector using simulated (MC) events. The work conducted here is part of the ongoing analysis within a group formed by the UCSB HEP group, Université Libre de Bruxelles, and Beihang University under the CMS Collaboration<sup>3</sup>. The analysis items and methods included in this thesis are a subset of what will be in the official analysis and is a part of the final measurement. More detail on what approximation has been taken in order to obtain a preliminary expected result is discussed in the next few sections.

## 1.4 Background and signal simulation

The following list contains the MC samples (by physical process) used in this thesis, the first two contain (off-shell) Higgs boson in the intermediate state:

- ggZZ offshell: Gluon fusion  $gg \rightarrow \text{H} \rightarrow \text{ZZ}$
- VVZZ offshell: Vector boson fusion (VBF) into Higgs boson
- qqZZ, qqWZ, qqWW
- DY: Drell-Yan process

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<sup>3</sup>Internally, CMS AN-20-081

- TT:  $t\bar{t}$ , including samples with additional vector boson (TTW/TTZ) or photon + jets (TTGJets).

The MC samples for all processes except DY are produced in RunIIAutumn18MiniAOD-102X, DY sample is produced in RunIISummer16MiniAODv3 94X and scaled appropriately afterwards. As their names suggest, each year in the Run 2 (2016-2018) has their separate MC samples since the running condition (the LHC) and the condition of the CMS can change from year to year, though the physics remains the same.

Various programs are used in the long chain of simulated events production. POWHEG v2 [11], JHU GENERATOR AND MELA [12], are used for the signal simulation. MADGRAPH [13] is used to generate NLO background samples, PYTHIA [14] for parton showering and NNPDF 3.1[15] sets are used for the parton distribution functions.

Before diving into the procedure in which the signal samples are generated separately and subsequently combined with via a reweighting, it would be appropriate to give an account for the general idea behind the ‘event weight’ and its significance.

As described at the beginning of this section, events that correspond to different physical processes are generated in different MC configurations at different ‘order’ of the QCD/QED physics. Naively, one would imagine a process where the MC can directly simulate the physics at the LHC at a given center-of-mass energy. Unfortunately this is neither efficient nor possible: not possible because some physical processes (especially QCD ones) are non perturbative and stepped approach is taken to gradually build up to the final event.

It is also impossibly inefficient because the processes an analysis concerns (for example, in all SUSY searches) usually have a tiny (if not 0) cross sections compare to other common ones that can be found at  $\sqrt{s} = 13$  TeV at the LHC. And it would be a waste of computing resources to generate the common processes over and over again.

Example of some important weights		
Name	abbr.	Description
Generator weight	GEN wgt	Given by MC event generator
Pile-up weight	PU wgt	Correction for the pile-up effect
Matrix element weight	ME wgt	From generator that uses ME Likelihood approach
K-factor	Kfactor	Correction for LO cross section of QCD processes

Table 1: An incomplete list of weights used in the MC events used.

In reality, Monte Carlo events are each given many ‘weights’ (Tab. 1), so that we don’t have to generate uninteresting processes, and at the same time, for the events that lack in number (results in poor statistics in the distribution), one can optionally generate extension events set for it. In practice, for example as shown in Fig. 3, you can have  $1 \times 10^6$  raw MC entries, but only 100 expected events (yield) in the integration of the histogram, due to the physics process being very rarely seen. But, these  $1 \times 10^6$  would form a much smoother distribution than if you only had 100 entries to work with. Also, this enables the generation of ‘unknown’ processes which can be used to constrain possible new physics in a likelihood fit (against null hypothesis).

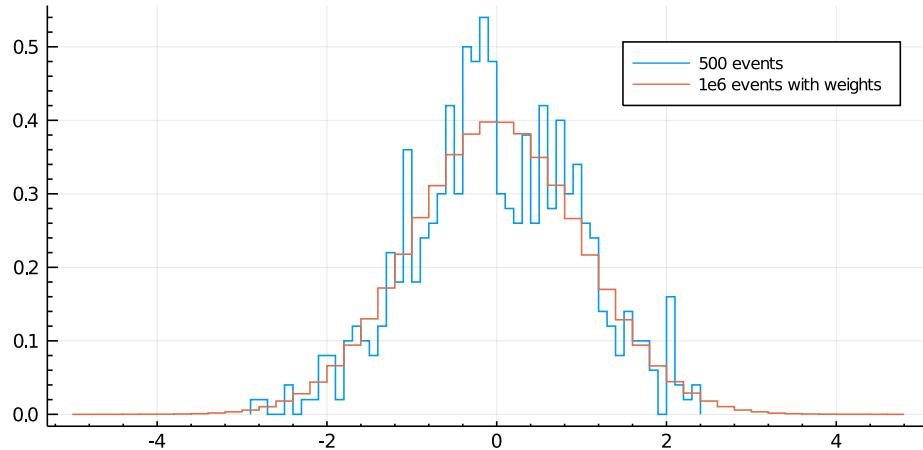


Figure 3: Illustrative plot to show how weights can make distribution smoother without altering the total number of expected events.

In this thesis, we explicitly use K-factors, and variation of it to obtain Electroweak

systematical uncertainty in qqZZ/qqWZ/qqWW backgrounds. For the signal processes involving the Higgs boson, a special treatment is given to merge and obtain high statistics sample from multiple ‘raw’ samples with different ‘true’ Higgs mass corresponding to the  $m_H$  term in the denominator on the right hand side of Eq. 1. This approach is also desirable for generation of off-shell (Higgs) decays. The procedures used and the resultant combined signal samples are discussed in the Sec. 2.2.

## 1.5 Uncertainties

A limited number of experimental and theoretical uncertainties in both signal and background processes are discussed here. Although dedicated to MC-only analysis, the leading theoretical uncertainties are considered:

Source	Uncertainty	Affected processes
Integrated luminosity	2.5%	GGH, VBF,qqZZ,qqWZ
Non-resonant background estimation	10%	TT, qqWW
Electroweak	$1\sigma$	qqZZ, qqWZ
Higgs branching ratio	2%	GGH, VBF
gg background	10%	ggZZ

Table 2: Summary of systematic uncertainties considered in this thesis and their magnitude as well as processes affected by them.

Most of the uncertainties are simply experimental, for example the luminosity. The Electroweak uncertainties are obtained by scaling up and down on the corresponding samples and non-resonant background comes from separate studies that examines  $e/\mu$  events in the data, which are limited by the number of events within the Z-boson mass window. Here we approximate this uncertainty as  $\pm 10\%$ . The gg background uncertainty is parametric respect to the  $M_T^{ZZ}$  defined below [16, 17, 17].

# Chapter 2

## Methods

In this chapter, a description of event selection is given, as well as definitions of key physical variables and how they are used to select events. Then, a procedure regarding how the signal sample is manipulated to produce a high statistics off-shell Higgs sample is presented. Finally, the binning of variables used to obtain the results is determined and defined.

### 2.1 Event selection and physical variables

Proton bunches cross each at a rate of about 400 MHz in the beam line of the LHC, naturally, not all of these crossings are recorded due to both technical limitation of the electronics as well as the fact that the vast majority of these crossings don't produce inelastic collision that is energetic enough to be interesting.

After the selection of Level 1 (L1) trigger and the higher level trigger (HLT), less than 1000 events per second are permanently recorded and would go to off-line, full reconstruction. Among these, we only select the ones that passes certain HLT trigger, for 2018 data set:

- HLT\_Mu17\_TrkIsoVVL\_Mu8\_TrkIsoVVL\_DZ\_Mass3p8
- HLT\_IsoMu24
- HLT\_Ele23\_Ele12\_CaloIdL\_TrackIdL\_IsoVL
- HLT\_DoubleEle25\_CaloIdL\_MW
- HLT\_DoublePhoton70
- HLT\_Ele32\_WPTight\_Gsf
- HLT\_Photon200

Most of the trigger names are self-explanatory, for example, the numbers are the momentum threshold (in GeV) to trigger them. The HLT is the starting point of the analysis. The purpose of the triggers is to ensure the events in the corresponding samples would contain the physics we wanted (since run data samples are divided according to what triggers they passed).

The next step is to ‘add’ composite variables that is significant to the physics of interests as well as making base-line cuts on the events. The jets are all AK4 jets unless mentioned otherwise, and a special type of jets, called b-tagged jets are **produced** identified using the DeepFlavour algorithm.[18] These are jets likely originated from relatively long-lived bottom quarks judging from a combination of many measurements including the displacement of vertexes. We shall also define a few of the uncommon variables in the list, and physical motivations are given in the following paragraphs.

After mandating passage of the set of triggers above, based on more delicate physics reasons, the baseline cuts are:

- $E_T^{\text{miss}} > 125\text{GeV}$ : the signal process creates true  $E_T^{\text{miss}}$  with neutrinos, this cut also reduce bkg such as DY.

- Both leptons have  $p_T > 25 \text{ GeV}$
- $|\Delta\phi_{\ell\ell\_E_T^{\text{miss}}}| > 1.0$
- $|\Delta\phi_{\ell\ell\text{Jets}\_E_T^{\text{miss}}}| > 2.5$
- $|m_{\ell\ell} - 91.2 \text{ GeV}| < 15 \text{ GeV}$ : the signal process consists of  $Z \rightarrow \ell\ell$ , we require the di-lepton system has a mass that is consistent within the Z mass peak.
- $p_T^\ell > 55 \text{ GeV}$ : the Drell–Yan (DY) process creates a lot of backgrounds events, but their di-leptons go back-to-back with expected value of this variable close to 0.
- No ak4-jet b-tagged jet ( $\Delta R = 0.4$ , loose b-tagging working point (WP))
- $\min |\Delta\phi_{j\_E_T^{\text{miss}}}| > 0.25$

In addition, to reduce the detector effects at high pseudo rapidity angels:

- $\eta_\mu < 2.4$
- $\eta_e < 2.5$

First and most importantly, the  $E_T^{\text{miss}}$  cut reduces background from DY significantly since they don't havd real  $E_T^{\text{miss}}$ . Notice many of the cuts are related to angles between various physical objects presented in the reconstruction. The reason is simple: the signal events, where Higgs goes to ZZ and one Z goes to 2 charged leptons and the other goes to 2 neutrinos, ideally would have the two Z's ‘back-to-back’ in Higgs boson's rest frame, leading to a large angel in the transverse plane.

Furthermore, in background events which do not contain this kinematic feature, the correlation in the directions of  $\vec{E}_T^{\text{missing}}$  and the transverse momentum of the leptons is weaker.

To use this kinematic feature to increase signal to background ratio, we define  $|\Delta\phi_{\ell\ell\_E_T^{\text{miss}}}|$  as the azimuthal angel (perpendicular to the beam line) between the di-lepton system and the transverse missing energy. In the signal events that produce 0 jet, this variable should be  $\pi$ . The cut is lowered to 1.0 due to the finding that in the (not so rare) cases where there are jet(s) recoiling against the ZZ system, this variable can go quite low.

This leads to the next variable on the list,  $|\Delta\phi_{\ell\ell\text{Jets\_}E_T^{\text{miss}}}|$ , which is almost the same except that we add all jets' momentum into the di-lepton system to account for the events that have produced jets, which in turn would cause the angel be lower in a multi-body final state.

[Also, a veto on the b-tagged jet is placed to reduce the background events originated from a pair of top quarts \( \$t\bar{t}\$ \) since top quark always decays to a bottom quark.](#)

Finally,  $\min |\Delta\phi_{j\_E_T^{\text{miss}}}|$  is the minimum azimuthal angel difference between any of the jet (that passes cuts) and the  $E_T^{\text{miss}}$ , it exist because jets are one of the most difficult physical objects measure, they often create so-called instrumental  $E_T^{\text{miss}}$  due to jet mis-measurements and it can be quite large in magnitude. However, such mis-measurements often yield large  $E_T^{\text{miss}}$  in the direction of the original jet. This cuts requires angular separation since in signal process, the jet recoils against ZZ system.

We also define variables that are not cut on, but are used for fitting:

$$D_{jj}^{\text{VBF}} = \frac{P_{SM}^{\text{VBF}}(\vec{\Omega})}{P_{SM}^{gg}(\vec{\Omega}) + P_{SM}^{\text{VBF}}(\vec{\Omega})} \quad (3)$$

as first introduced in another CMS paper [19]. In short, this variable (discriminator) is sensitive to the VBF physics and the correlation between (angels and mass of) the outgoing jets resulted from the VBF topology. At the same time:

$$M_T^{\text{ZZ}} = \left[ \sqrt{p_{T,\ell\ell}^2 + m_{2\ell}^2} + \sqrt{{E_T^{\text{miss}}}^2 + m_Z^2} \right]^2 - \left[ \vec{p}_{\ell\ell} + \vec{E}_T^{\text{miss}} \right]^2 \quad (4)$$

is defined based on the hypothesis that the  $E_T^{\text{miss}}$  is comprised of mainly the two neutrinos from one of the Z bosons. We shall see the usefulness of this variable in channels where not enough jets are present to construct the  $D_{jj}^{VBF}$  variable.

On top of the cuts stated in the beginning, several  $E_T^{\text{miss}}$  filters exist to remove detector anomalies. Lepton isolated tracks veto is implemented in the underlying analysis framework to ~~reject events with pathological reconstruction~~, remove events with additional leptons that are likely backgrounds, such as those can be found in WZ.

## 2.2 Signal samples re-weighting

Extra attention was given to the off-shell Higgs sample used in this thesis and two different kinds of re-weighting of the simulated events are applied in to produced a MC sample with wide mass spectrum way beyond the mass of Higgs ( $\approx 125 \text{ GeV}$ ). We use the gluon fusion Higgs (ggH) sample to illustrate the procedures, the same procedures are applied to the VBF samples as well.

We start by generating separate samples with different Higgs mass (LHECandMass). This is the mass of Higgs terms appears in the propagator on the R.H.S of Eq. 1 as mentioned before. The raw distribution of the true mass in different samples (without any weight) is shown in Fig. 4 (left). As expected, the peak of the distribution moves to the right as the mass of the sample becomes larger, at the same time, the ‘peak’ of samples with very large mass are wider because the width of the Higgs increases. We also see that for some lower mass samples (200, 300, 400 etc.), they have a cut-off (due to insufficnet statistics) beyond  $M \approx 2500$  which means they have 0 statistics beyond that mass range. After applying the GEN, PU, and ME weights given by their individual MC process and JHUGen MELA, as shown in Fig. 4, we see that they are consistent with each others’ line shape. However, it is clear that:

- (i) Lower mass samples have insufficient statistics in the tail region
- (ii) Samples have poor statistics in mass windows that are far from their true mass (as listed in the legend).

The second point is best illustrated by the wide spikes of lower mass samples near their ‘cut-offs’, as well as the visible fluctuations of high-mass samples in the mass region (don’t let the visual mislead you, the plots are in semi-log scale).

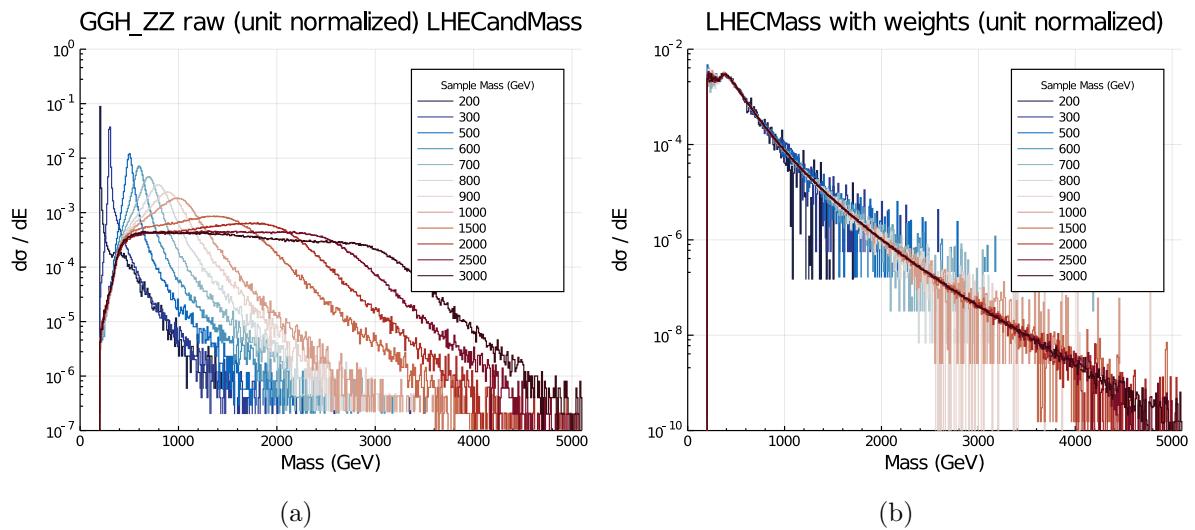


Figure 4: Unit normalized distributions of LHECandMass before (left) and after (right) applying the weights. Together they show a need to combine samples for a wide-range, high statistics signal sample. Bin size = 10 GeV.

The goal of the combination of samples is to use all the events, but with a correction weight such that each sample has a higher weight in the region where they pose good statistics while keeping the overall normalization stays unchanged. To do this, we pick a list of ‘mass windows’ with edges sitting on the true masses of the samples, and we define effective number of events  $N_{\text{eff}} = \frac{(\sum \text{wgts})^2}{\sum (\text{wgts}^2)}$  within each mass window. Here, the wgts corresponds to the product of PU wgt, GEN wgt, K-factor, and ME weight for the GGH sample in consideration. For a specific GGH sample  $i_0$  and its events fall in a mass

window  $j$ ,  $N_{\text{eff}}^{i_0 j}$  is first obtained and a re-weighting factor can be computed:

$$\text{wgt}_{\text{window}}^{i_0 j} = \frac{N_{\text{eff}}^{i_0 j}}{\sum_i N_{\text{eff}}^{ij}} \quad (5)$$

This factor is applied to all events from sample  $i_0$  within the window  $j$ . Conceptually, the effective number of events ensures the weight is not skewed by the difference in the overall normalization of samples, and in each of the mass windows, samples with more concentrated statistics in that window are given a higher weight. In Fig. 5 (a), a clear diagonal pattern can be seen, physically it means that samples with higher true mass are given a higher weight in tail mass windows—consistent with the expected outcome. Complete table of weights can be found in Appendix A.1.

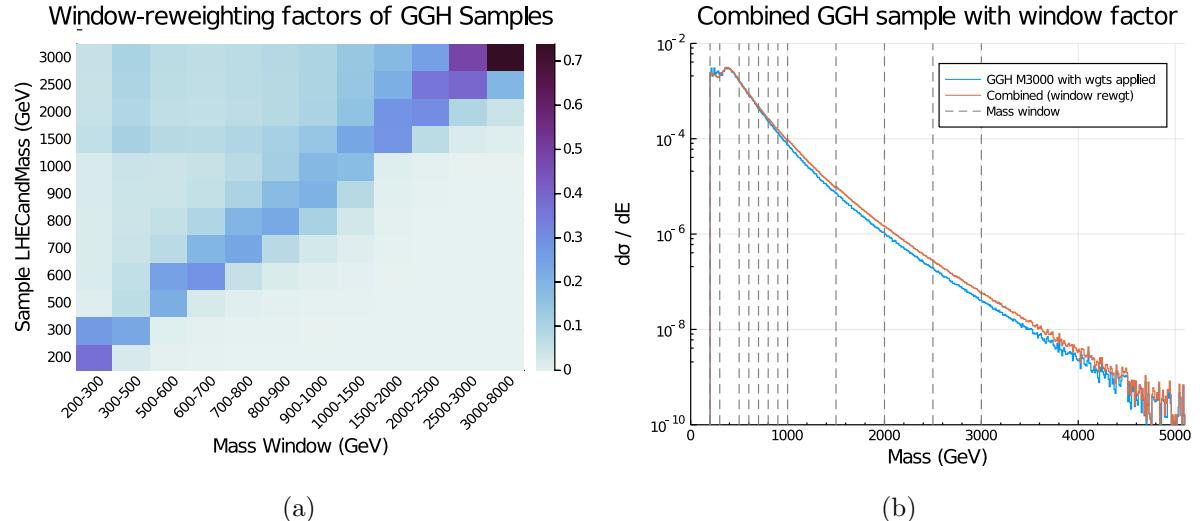


Figure 5: Heatmap of window re-weight factors of different samples and mass windows (left); effects of applying window factors for the combined sample (right)

However, even with the unit normalization, there are still inconsistencies in the shape as shown in Fig. 5 (b). This is likely because the finite number of events and non-infinitesimal mass window size used. We introduce another correction factor for this

small artifacts. Iteratively going through every sample, between the previous and the next one, derive a sample mass factor based on:

$$\text{wgt}_{\text{mass}}^{i,i+1} = \frac{\sum \text{wgt}_i}{\sum \text{wgt}_{i+1}}, \text{ for events that has Mass between sample mass of } i \text{ and } i + 1 \quad (6)$$

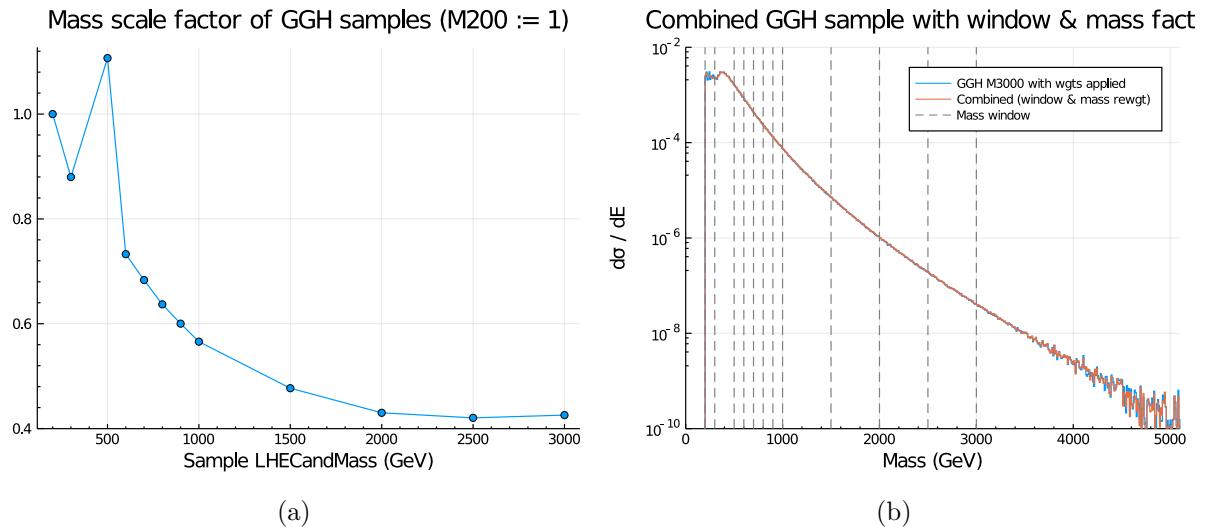


Figure 6: Iterative sample mass factors obtained (left) and the final combined sample (right)

This factor corrects the high variations of overall normalizations between samples, the factors and result are shown in Fig. 6. As expected, high mass samples need a down correction (not by a lot) to eliminate the deviated trend before. Complete tables of this weights for both samples can be found in Appendix A.2.

Finally, 1.10 is multiplied to the weights of all ggZZ processes (all of BKG, SIG, BSI of GGH sample) as a K factor for Next-to-next-to-leading-order (NNLO) → Next-to-next-to-next-to-leading-order (N3LO) QCD.

Although we used the particular matrix element weights for one of the signal hypothesis, these two correction factors apply to all hypothesis and a plot of them without unit area normalization are shown in Fig. 7. As expected, background exceeds signal by more

than 100% which is partially why the constrain is hard to obtain.

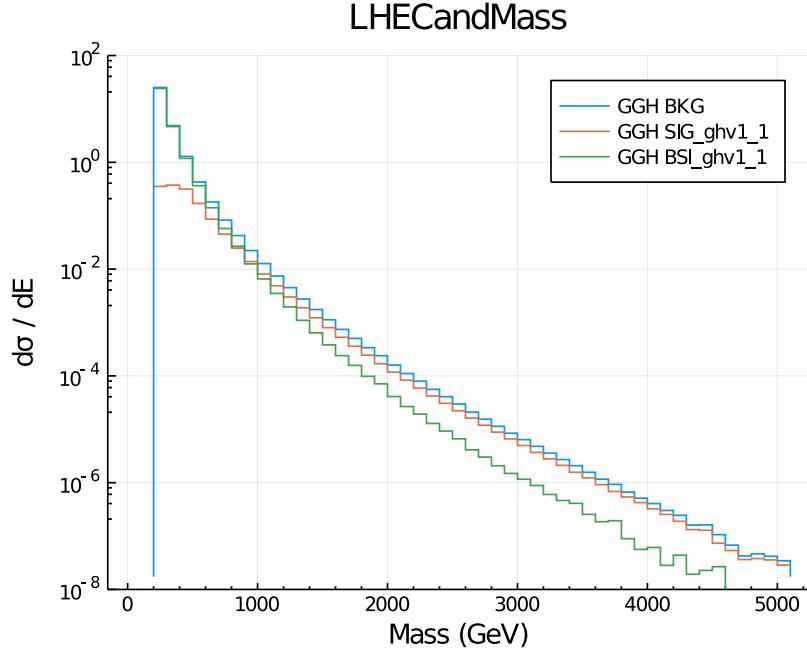


Figure 7: Distributions of background, signal, and background signal interaction

## 2.3 Strategy in variable selection and binning and systematical uncertainties

After preparing the signal samples and decided on the event selection criteria, we move on to decide on the variables and their (2D histogram, as shown in Fig. 8) ‘binning’ before a combined-limits fit can be applied. As discussed in earlier sections, one of the more inventive variables newly introduced specifically for the analysis is the DJJVBF discriminator. However, it is clear that this variable is undefined for events with  $N_{\text{jets}} < 2$ . To not ‘waste’ any statistical significance, we use  $E_T^{\text{miss}}$  in its place for the  $N_{\text{jets}} = 0, 1$  categories. In total, we have  $2$  ( $ee$  or  $\mu\mu$ )  $\times 4$  ( $N_{\text{jets}} = 0, 1, 2, 3+$ )  $= 8$  channels to consider

when making histogram templates. Bin edges for different categories in number of jets are as the following:

- $N_{\text{jets}} \geq 2$ 
  - $M_T^{\text{ZZ}} = 150, 300, 400, 600, 800, 1000, 13000$  (GeV)
  - $\text{KD1} = \text{DJJVBF} = 0, 0.2, 0.4, 0.6, 0.8, 1$
- $N_{\text{jets}} < 2$ 
  - $M_T^{\text{ZZ}} = 150, 300, 400, 600, 800, 1000, 13000$  (GeV)
  - $\text{KD1} = E_T^{\text{miss}} = 125, 200, 280, 420, 500, 800, 13000$  (GeV)

The higher mass ( $M_T^{\text{ZZ}}$ ) bins are wider because samples have difficulty filling them due to physical reasons (especially for backgrounds) and the cuts being applied. As shown in Fig. 8, the relative error (err/counts) of each bin in the histogram templates are displayed, a balance between significance and uncertainty is obtained. We use the above binning for all samples and 8 channels that are considered in this thesis. See Appendix B.1 for a compilation of template histograms of GGH sample.

For low-yield background samples, due to the nature of NLO samples, bins sometimes would have negative content. To mitigate it's effect on the likelihood fitting (pathological), we replace the bin content by (Integral of the histogram)  $\times 10^{-5}$ .

Systematical uncertainties are also included in the fitting:

- Luminosities GGH\_ZZ, VBF\_ZZ, qqZZ, qqWZ
- NRB (non-resonant background) Estimation: TT, WW
- Branching Ratio of Higgs to ZZ to 4l: GGH\_ZZ, VBF\_ZZ
- K-fatcors of backgroun gluongluon parameter

See Table. 1.5 for respective range of these uncertainties.

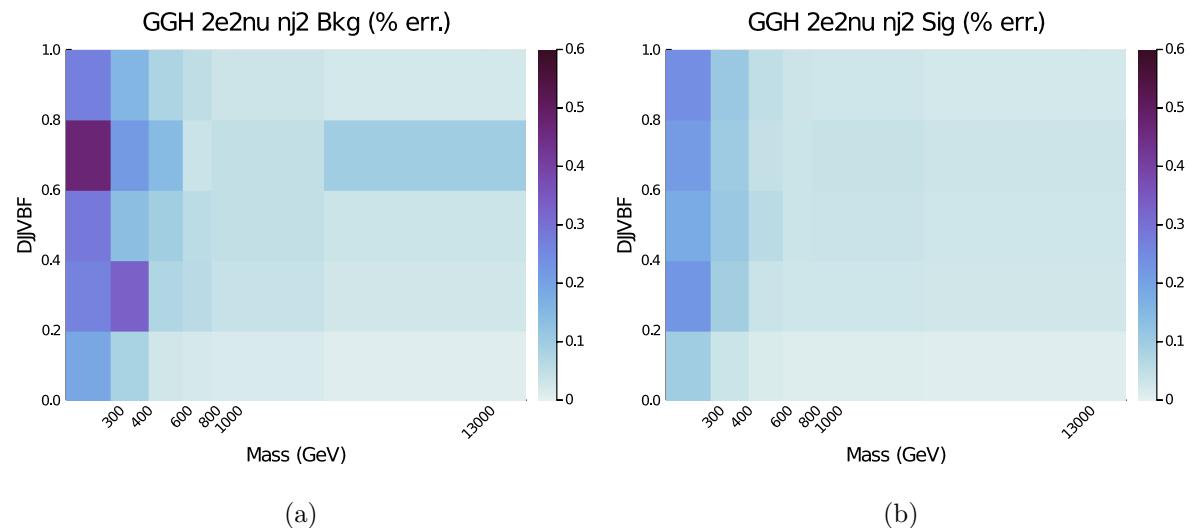


Figure 8: Background (left) and Signal (right) histogram templates for  $2e2\nu$  with  $N_{\text{jets}} = 2$ , showing percentage (statistical) error.

# **Chapter 3**

## **Results and interpretation**

Final results of this thesis are presented. As the result acts as ‘expected’ limits, the ongoing work and potential interpretation are discussed.

### 3.1 Limits on Higgs decay width

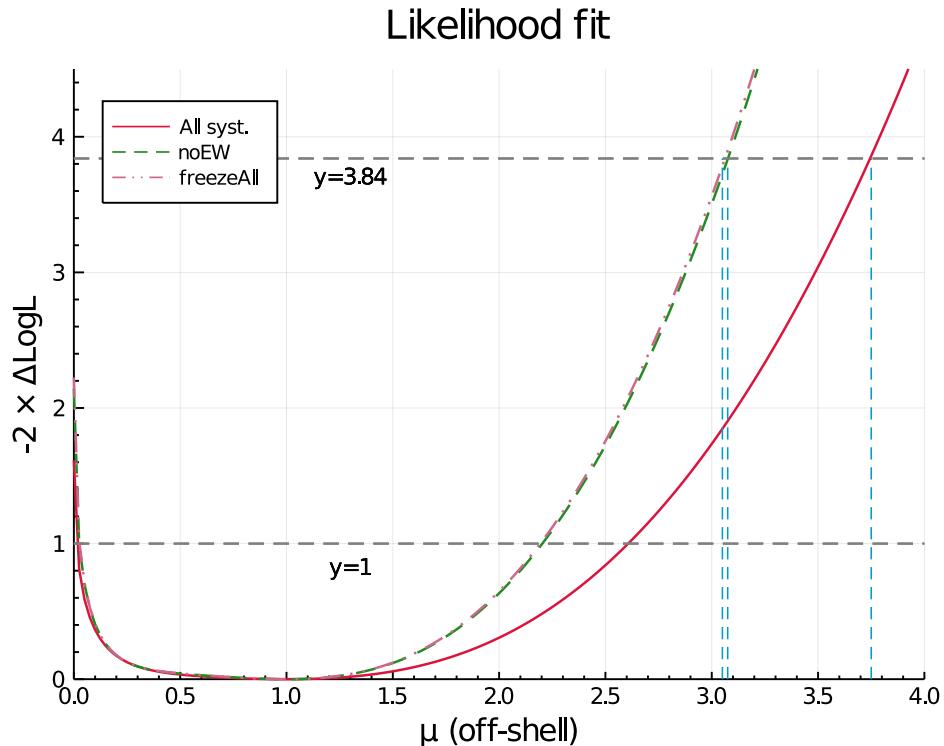


Figure 9: Maximum likelihood fit of  $\mu_{\text{off-shell}}^{\text{H}}$  (off-shell rate ratio). For all systematics (red), no Electroweak syst. (green), 0 syst. (orange):  $y$ -intersect= $\{1.61, 2.14, 2.22\}$ ,  $1\sigma$  lower limits= $\{0.025, 0.025, 0.025\}$ ,  $1\sigma$  higher limits= $\{2.6, 2.2, 2.2\}$ , 95% CL limits= $\{3.75, 3.08, 3.05\}$ , respectively.

After running through Combined Limited tool [20, 21, 22] for likelihood fitting, we first extract the significance of the off-shell rate. (Fig. 9) The y-axis is understood to be  $\sigma^2$  in terms of significance, thus the intersection with the y-axis is the signal sensitivity, or in other words, rejection of the 0 width hypothesis (no off-shell), which has a significance of  $\sqrt{1.61} \approx 1.26\sigma$  in this fit with all the systematic uncertainties included. As the systematics are turned off', the constraint becomes tighter, producing an error band for the expected final result in the upcoming official analysis. The electroweak uncertainty is displayed individually as it encompasses most of the systematic uncertainties.

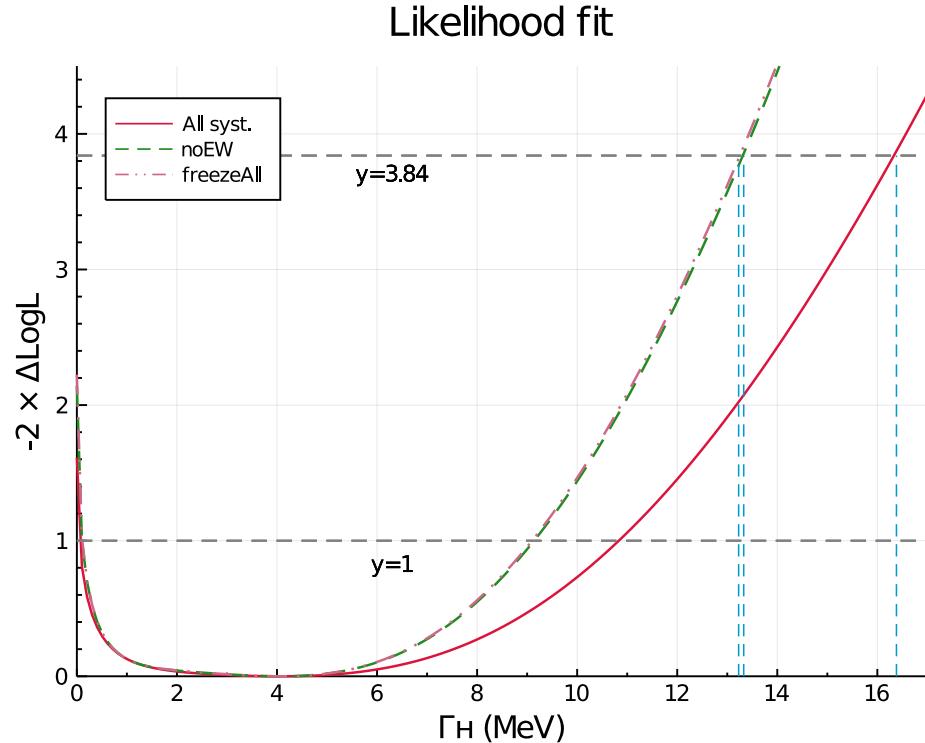


Figure 10: Maximum likelihood fit of Higgs decay width. For all systematics (red), no Electroweak syst. (green), 0 syst. (orange):  $y$ -intersect= $\{1.61, 2.14, 2.22\}$ ,  $1\sigma$  lower limits= $\{0.10, 0.10, 0.10\}$  MeV,  $1\sigma$  higher limits= $\{10.78, 9.16, 9.06\}$  MeV, 95% CL limits= $\{16.38, 13.33, 13.23\}$  MeV, respectively.

Furthermore, by un-constraining the  $\mu_F$  and  $\mu_V$  which are the production rate of Higgs from fermion fusion vs. massive boson fusion, as mentioned in the Eqn. 1. We adopt the range suggested [23]. The constraint on the decay width of Higgs  $\Gamma_H$  is shown in Fig. 10. The minimal (max likelihood) falls on 4.07 MeV, consistent with the Standard Model hypothesis being used. Again,  $1\sigma$  and 95% CL are marked respectively. And a final result of  $\Gamma_H < 16.38$  MeV can be quoted.

# Appendix A

## Weights Table for Higgs Sample

### A.1 Windowed re-weighting results

(100 GeV)	2-3	3-5	5-6	6-7	7-8	8-9	9-10	10-15	15-20	20-25	25-30	30-80
30	0.0572	0.1082	0.0737	0.0686	0.0742	0.0857	0.1015	0.1431	0.1832	0.2573	0.4859	0.74
25	0.0531	0.1011	0.071	0.0659	0.0734	0.083	0.1001	0.1461	0.2169	0.3597	0.3971	0.1964
20	0.0545	0.0913	0.0659	0.0633	0.0713	0.084	0.1021	0.1599	0.2889	0.3043	0.0933	0.0481
15	0.0643	0.1118	0.0829	0.0812	0.0948	0.113	0.1447	0.2395	0.2847	0.0732	0.022	0.0141
10	0.0242	0.0493	0.044	0.0529	0.0752	0.1179	0.1896	0.1752	0.0141	0.0028	0.0008	0.0008
9	0.0238	0.0468	0.0465	0.0602	0.1028	0.1791	0.2064	0.0832	0.0066	0.0013	0.0004	0.0003
8	0.0205	0.045	0.0592	0.0941	0.1991	0.2324	0.1099	0.0341	0.0031	0.0007	0.0002	0.0001
7	0.0207	0.0478	0.076	0.1965	0.2409	0.0787	0.0324	0.012	0.0014	0.0003	0.0001	0.0001
6	0.0212	0.0637	0.2502	0.2875	0.0573	0.0205	0.0098	0.0048	0.0007	0.0002	0.0	0.0001
5	0.014	0.0705	0.217	0.0251	0.0083	0.0041	0.0023	0.0014	0.0002	0.0	0.0	0.0
3	0.2708	0.2386	0.0111	0.0037	0.002	0.0012	0.0008	0.0006	0.0001	0.0	0.0	0.0
2	0.3757	0.0259	0.0026	0.001	0.0006	0.0004	0.0003	0.0001	0.0	0.0	0.0	0.0

Table 3: Windowed re-weighting result for ggZZ samples.

(100 GeV)	2-3	3-4	4-5	5-6	6-7	7-8	8-9	9-10	10-15	15-20	20-25	25-30	30-80
30	0.0013	0.0024	0.0021	0.0035	0.003	0.005	0.01	0.0112	0.0265	0.0764	0.1583	0.3873	0.6233
25	0.0017	0.0028	0.0023	0.0046	0.004	0.0063	0.0121	0.0149	0.039	0.1193	0.2795	0.3844	0.2218
20	0.0015	0.0044	0.0038	0.0067	0.0065	0.0096	0.0211	0.0256	0.0709	0.2562	0.3658	0.1491	0.0904
15	0.0045	0.0088	0.0075	0.015	0.0131	0.0218	0.0448	0.06	0.1727	0.4011	0.1484	0.0569	0.0434
10	0.0052	0.0077	0.0092	0.0184	0.0208	0.0464	0.0152	0.0085	0.2471	0.0459	0.0136	0.0058	0.0055
9	0.0054	0.0142	0.0167	0.0378	0.0449	0.0975	0.332	0.4022	0.2029	0.0389	0.0122	0.0058	0.0052
8	0.0107	0.0196	0.0252	0.0623	0.0998	0.3037	0.187	0.2898	0.1195	0.0255	0.0085	0.0036	0.0042
7	0.0087	0.0226	0.0299	0.1043	0.2057	0.2715	0.2067	0.0905	0.0481	0.0125	0.0047	0.0023	0.0022
6	0.0176	0.0482	0.0882	0.0557	0.4805	0.157	0.1003	0.0499	0.0348	0.0108	0.0039	0.002	0.0017
5	0.029	0.1134	0.2939	0.5247	0.1041	0.0499	0.0397	0.025	0.0188	0.0065	0.0026	0.0015	0.0013
4	0.0634	0.336	0.4353	0.1231	0.001	0.0183	0.0174	0.0129	0.0111	0.0042	0.0015	0.0008	0.0007
3	0.1082	0.3288	0.0725	0.0305	0.0114	0.0086	0.009	0.0064	0.0058	0.0022	0.0009	0.0005	0.0004
2	0.7428	0.0912	0.0134	0.0132	0.005	0.0042	0.0047	0.0031	0.0028	0.0004	0.0	0.0	0.0

Table 4: Windowed re-weighting result for VBF samples.

## A.2 Sample mass factors for ggZZ and VBF\_ZZ samples

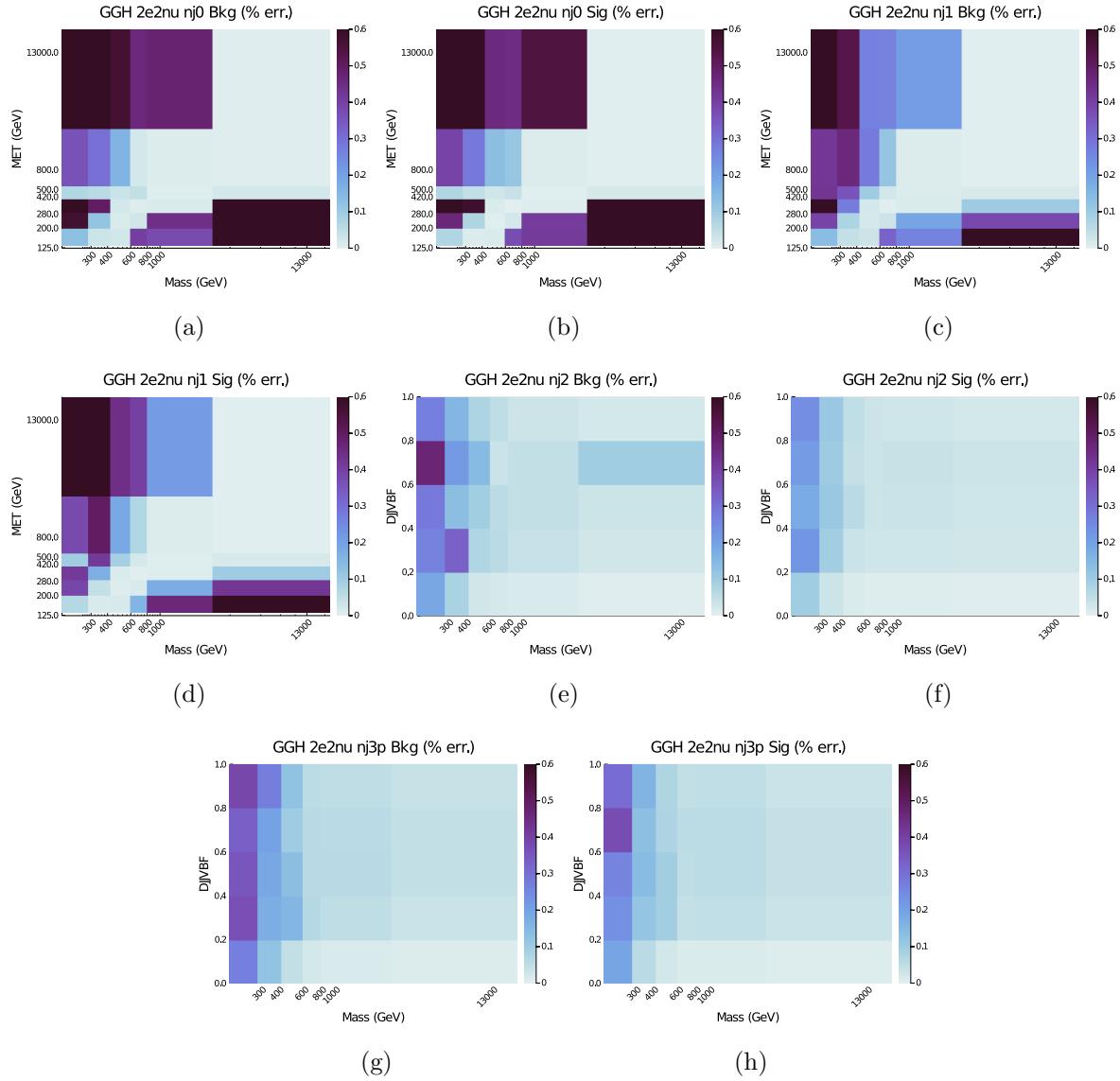
ggZZ		VBF_ZZ	
Mass (GeV)	Factor	Mass (GeV)	Factor
200	1	200	1
300	0.8798	300	0.8731
500	1.1068	400	0.7873
600	0.7327	500	0.7106
700	0.6833	600	0.6361
800	0.6370	700	0.5754
900	0.6001	800	0.5217
1000	0.5658	900	0.4818
1500	0.4769	1000	0.4257
2000	0.4300	1500	0.3255
2500	0.4204	2000	0.2968
3000	0.4256	2500	0.2898
		3000	0.2916

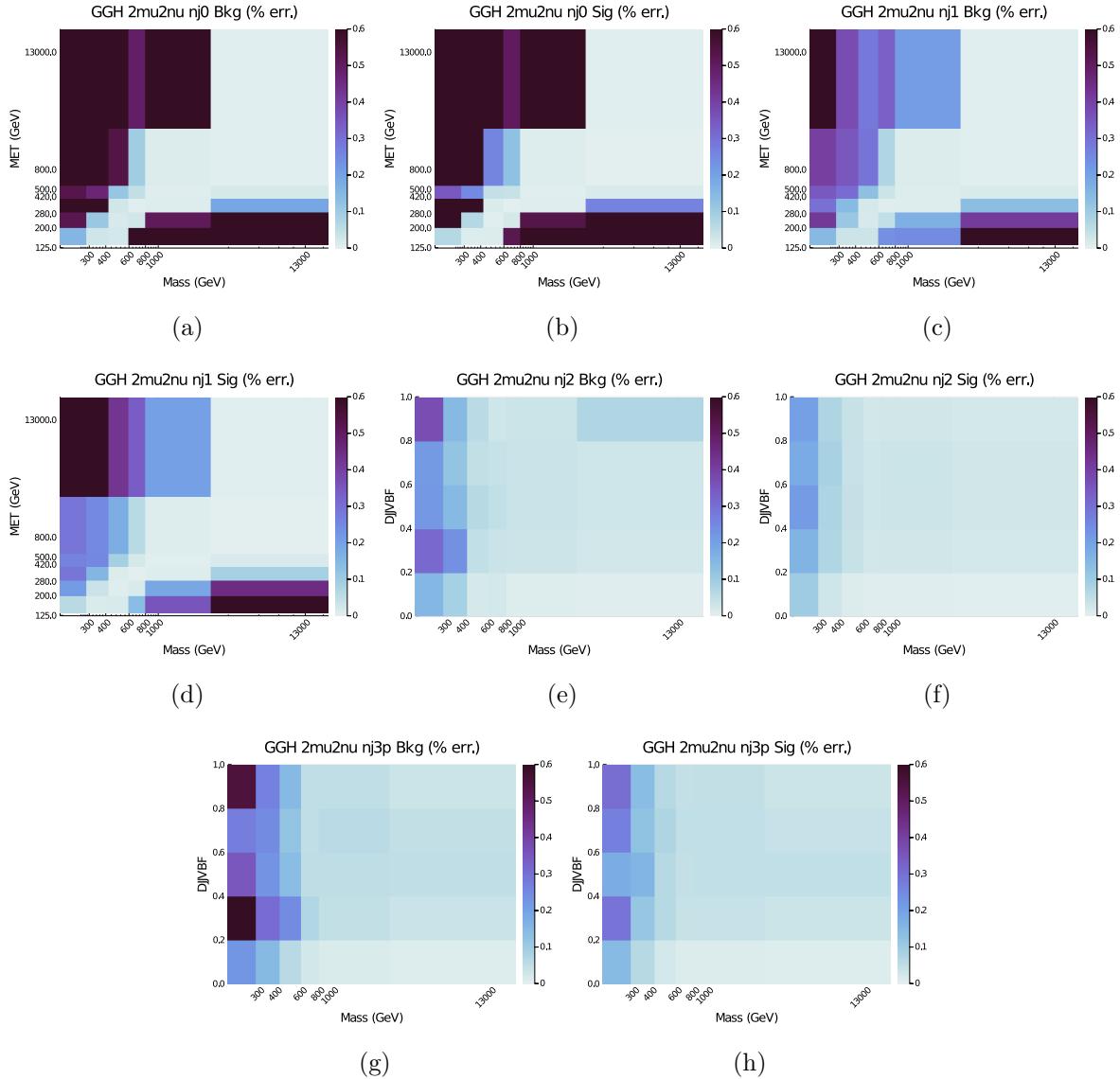
Table 5: Sample mass re-weighting factors for both signal samples.

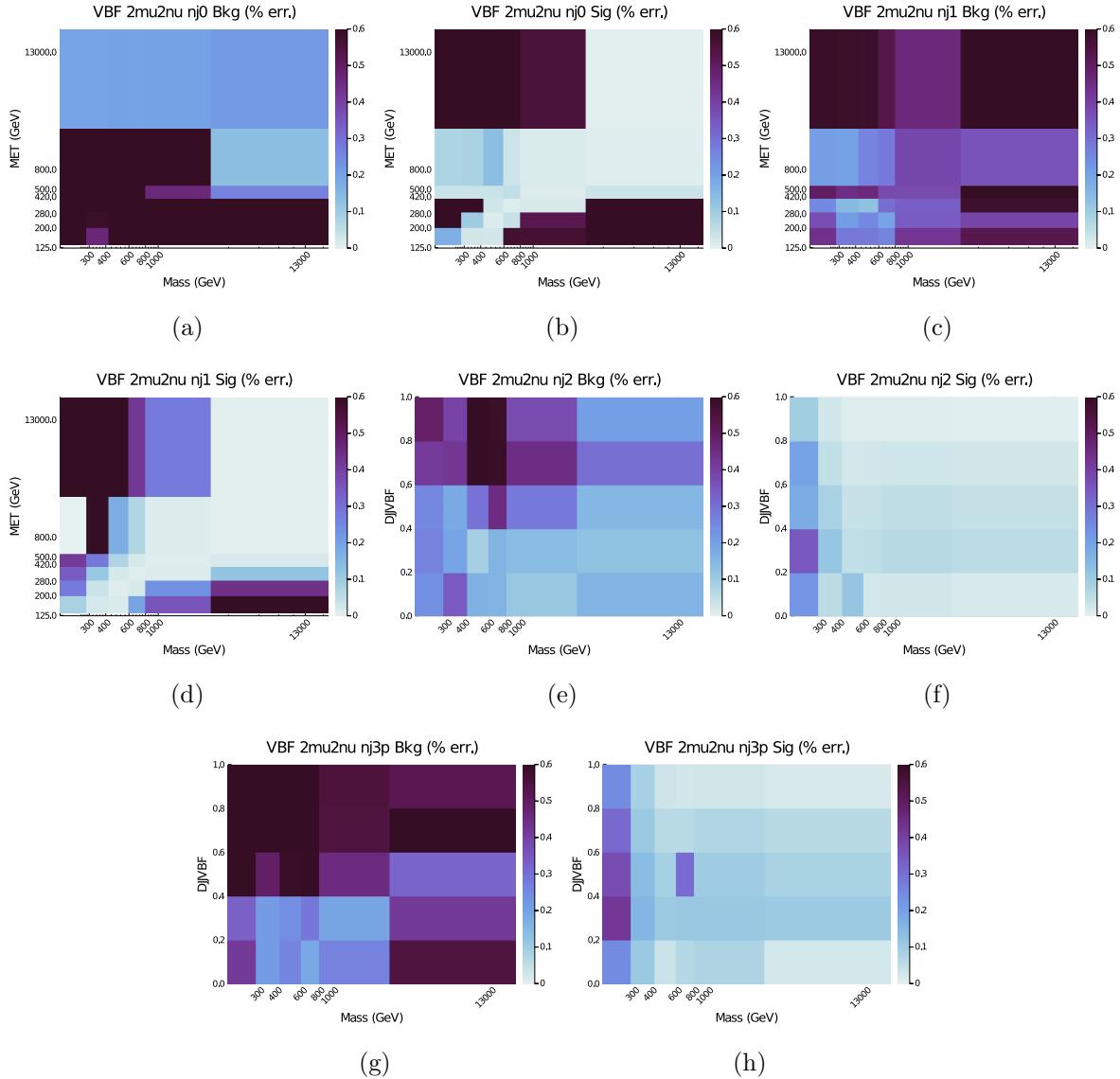
# **Appendix B**

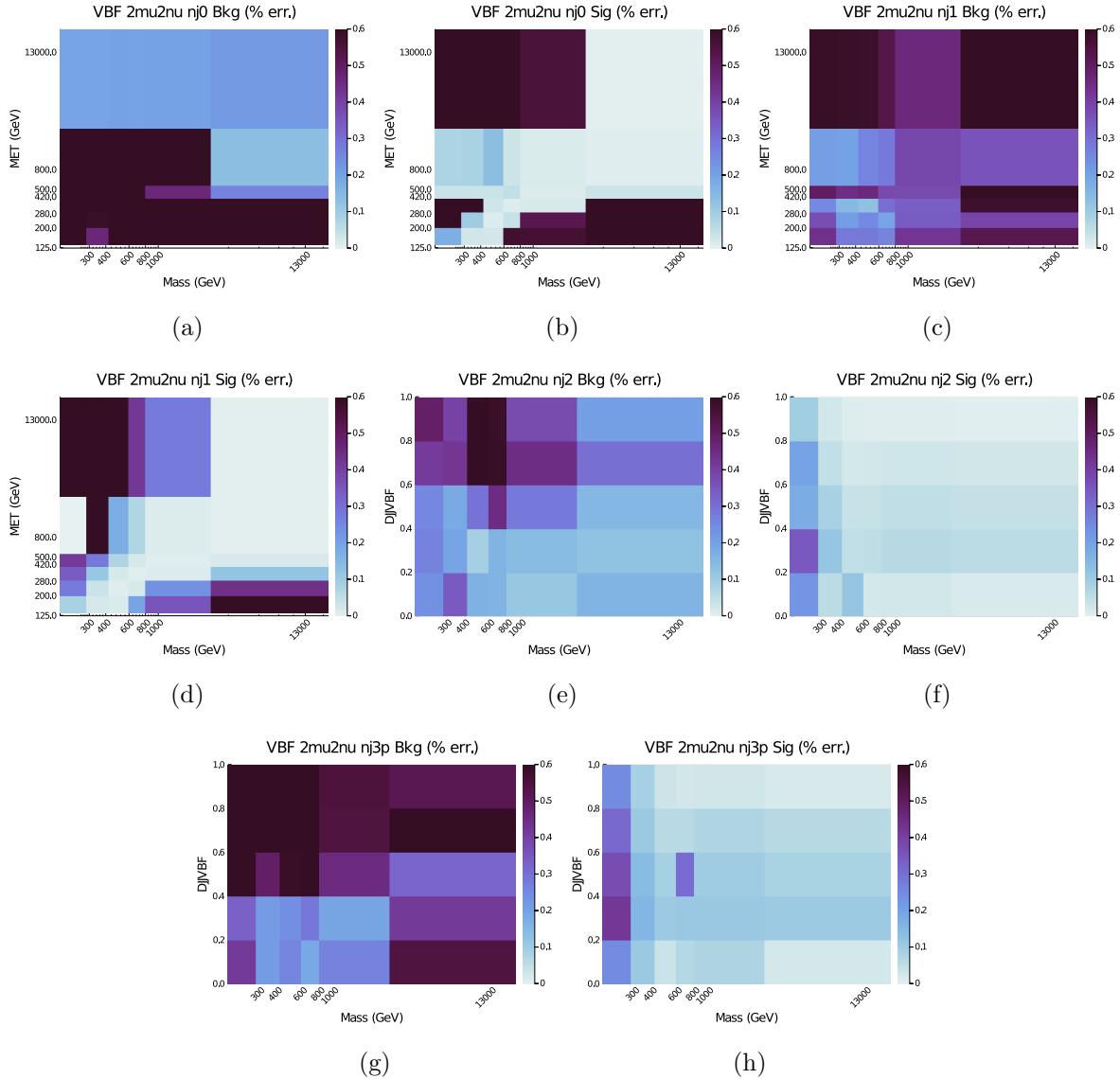
## **Additional Figures**

## B.1 Signal Sample Fitting Template Histogram of Background and Signal









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