#### Simple Pendulum

A simple pendulum also exhibits periodic motion.

It consists of a particle-like bob of mass *m* suspended by a light string of length *L*.

The motion occurs in the vertical plane and is driven by gravitational force.

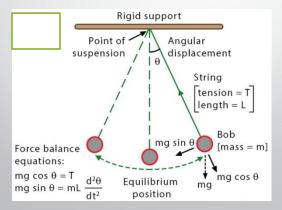
The motion is very close to that of the SHM oscillator if the angle is <10°

#### Simple Pendulum, 2

The forces acting on the bob are the tension and the weight.

- 1)  $\vec{T}$  is the force exerted on the bob by the string.
- 2)  $\mathbf{m}\vec{g}$  is the gravitational force.

The tangential component of the gravitational force is a restoring force.



# Simple Pendulum, 3

Projection along x-axis, 
$$F_x = ma_x$$
  
 $x = L\theta, \quad \dot{x} = L\dot{\theta}, \quad \ddot{x} = L\ddot{\theta}$   
 $-mg\sin\theta = mL\ddot{\theta}$ 

Projection along y-axis  $mg \cos \theta = T$ 

The length L of the pendulum is constant, and for small values of  $\theta$ .  $-mq\theta = mL\ddot{\theta}$ 

$$mL\ddot{\theta} + ma\theta = 0$$

divided by mL, we obtain:

$$\ddot{\theta} + \frac{g}{I}\theta = 0$$

This confirms the mathematical form of the motion is the same as for SHM.

# Simple Pendulum, 4

• The function  $\theta$  can be written as

$$\theta = \theta_m \cos(\omega t + \varphi)$$

The angular frequency is:

$$\omega = \sqrt{\frac{g}{L}}$$

The period is:

$$T = \frac{2\pi}{\omega} = 2\pi \sqrt{\frac{L}{g}}$$

The frequency is:

$$f=rac{1}{T}=rac{1}{2\pi}\sqrt{rac{g}{L}}$$
 Or  $f=rac{\omega}{2\pi}=rac{1}{2\pi}\sqrt{rac{g}{L}}$ 

### Simple Pendulum, Summary

- The period and frequency of a simple pendulum depend only on the length of the string and the acceleration due to gravity.
- The period is independent of the mass.

All simple pendula that are of equal length and are at the same location oscillate with thesame period.

### **Damped Oscillations**

In many real systems, non-conservative forces are present.

- 1. This is no longer an ideal system (the type we have dealt with so far).
- 2. Friction and air resistance are common non-conservative forces.

In this case, the mechanical energy of the system diminishes in time, the motion is said to be *damped*.

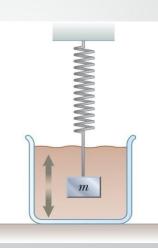
#### **Damped Oscillation, Example**

One example of damped motion occurs when an object is attached to a spring and submerged in a viscous liquid.

The retarding force can be expressed as

$$\vec{R} = -b\vec{v}$$

b is a constantb is called the damping



#### **Damped Oscillations, Graph**

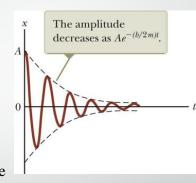
A graph for a damped oscillation.

The amplitude decreases with time.

The blue dashed lines represent the *envelope* of the motion.

Use the active figure to vary the mass and the damping constant and observe the effect on the damped motion.

The restoring force is -kx.



### **Damped Oscillations, Equations**

From Newton's Second Law

$$\sum \vec{F} = m\vec{a}$$

$$\vec{F_s} + \vec{R} = m\vec{a}$$

$$-kx - bv = ma_x$$

or 
$$a_{x} = \ddot{x}$$
 and  $v = \dot{x}$ 

$$so m\ddot{x} + b\dot{x} + kx = 0$$

divided by m, we obtain:

$$\ddot{x} + \frac{b}{m}\dot{x} + \frac{k}{m}x = 0$$

When the retarding force is small compared to the maximum restoring force we can determine the expression for *x*.

• This occurs when *b* is small.

The position can be described by

$$x(t) = Ae^{\frac{-bt}{2m}}\cos(\omega t + \varphi)$$

The angular frequency will be

$$\omega = \sqrt{\frac{k}{m} - \frac{b^2}{4m^2}}$$

## Damped Oscillations, Natural Frequency

- When the retarding force is small, the oscillatory character of the motion is preserved, but the amplitude decreases exponentially with time.
- The motion ultimately ceases.

## **Types of Damping**

If the restoring force is such that  $\frac{b}{2m} < \omega_0$ , the system is said to be **underdamped.** 

When b reaches a critical value b such that  $\frac{b}{2m} = \omega_0$ , the system will not oscillate. The system is said to be <u>critically damped</u>. If the restoring force is such that  $\frac{b}{2m} > \omega_0$ , the system is said to

be <u>overdamped</u>.

