

ChPol

Calculation of ESP charges consistently with an MM-Pol model

Stefano Caprasecca¹

ChPol

ChPol is a Molecolab² tool. It performs a polarizable ESP fitting of a given molecule, according to one of the MMPol models available.³

The next Sections describe how to use the tool and the theory behind it.

Usage

Generation of input files

Four files are needed:

`xxx.gesp` a `gesp` file used to get the QM potential to be fit (generated by Gaussian);

`xxx.mol2` a `mol2` file used to get connectivity and QM dipole information (generated by Antechamber);

`xxx.pol` a file with the values of atomic isotropic polarizabilities, in a.u.³, to be edited by hand.

`xxx.cns` a file with constraints, to be edited by hand;

The way to obtain or write them is detailed below. Please follow the instructions in the correct order.

Generating the `gesp` file This file contains the QM electrostatic potential on a grid of points around the molecule. It is produced by Gaussian from the following input file `xxx.com`:

```
#p b3lyp/6-311+G(d,p) pop=ESP IOp(6/50=1)
```

Calculate ES potential and print it to output file `xxx.gesp`

```
O 1
O 0.0000 0.0000 0.0000
H 1.0000 0.0000 0.0000
H 0.0000 1.0000 0.0000
```

`xxx.gesp`

¹stefano.caprasecca@for.unipi.it

²www.dcci.unipi.it/molecolab — For help please write to molecolab.pisa@gmail.com.

³Only Amber12 (Wang) implemented so far.

Additional options one may wish to use are the following:

I0p(6/41= n) use n layers in the ESP fit (default is 4)

I0p(6/42= n) density of points per unit area in the ESP fit (default is 1)

File `xxx.gesp` is then produced. It contains information on the atom positions, the QM dipole and the ES potential on the grid points.

Generating the mol2 file This file is generated by calling antechamber, which reads the `gesp` file produced in the previous step.

Call antechamber:

```
antechamber -i xxx.gesp -fi gesp -o xxx.mol2 -fo mol2 -c resp -pf yes
```

The `mol2` file generated contains, among other things, the connectivity information. The structure of the connectivity is the following:

```
@<TRIPOS>MOLECULE
```

```
xxx
```

```
60    59    1    0    0
```

```
[...]
```

```
@<TRIPOS>BOND
```

```
1      1      2  1
```

```
2      1     21  1
```

```
3      1     22  1
```

```
4      1     60  1
```

```
5      2      3  2
```

```
[...]
```

The first piece of information, at the beginning of the file, and starting with `@<TRIPOS>MOLECULE`, lists the number of atoms (60) and the number of bonds (59). Other bits of information follow.

In the second piece of information, starting with `@<TRIPOS>BOND`, lists all (59) bonds; in the example, the 1st bond is between atoms 1 and 2, and is single; 2nd bond is between atoms 1 and 21 and is single; and so on.

In some cases you may wish to change the connectivity information. To do so, modify the list of bonds as you wish, and also update the number of bonds in the first part, as appropriate.

No further changes are needed.

Polarizability file The file `xxx.pol` must be edited by listing all the values of atomic polarizabilities (in a.u.³), in the same order the atoms appear in the `xxx.com` input file. The values can be taken from reference papers. When using Amber12 (Wang) model, refer to Table 3, AL model, in Wang et al., *J. Phys. Chem. B* **115**, 3091 (2011). Other MMPol models are not yet implemented.

Constraint file The `xxx.cns` file contains the constraints you wish to apply to your fit. Note that the program perform both a standard, non-polarizable ESP fit, and a polarizable one. Both fits are constrained as requested, including the constraint on the total dipole moment.

The `xxx.cns` file *must* contain, *on the first line*, the total charge of the molecule. Additional constraints are optional.

total charge ⁴The total charge constraint is *not optional* and must be entered on the *first line* of the file.

fragment charge ⁵It is possible to specify that a certain fragment of the molecule must sum to a certain charge \bar{c} . Use the keyword `fragm`, then on the next line specify the number of atoms belonging to the fragment and the total charge, and on the following line the list of the atoms of the fragment, for instance:

```
fragm
4 1.0
9 15 17 19
```

This requires that the fragment made up by the 4 atoms 9,15,17 and 19 sums up to charge +1.0.

equivalence ⁵The equivalence between atoms is specified using the keyword `equiv`, followed on the next line by the number of equivalent atoms and on the next line by the list of such atoms, as in the example:

```
equiv
3
9 10 11
```

This requires that the three atoms 9, 10 and 11 have the same charge.

total dipole ⁵⁶Constraining the total dipole (i.e. the sum of the fixed molecular dipole given by the charge distribution, plus all the induced dipoles) to a certain value can be done by using the keyword `dipole`. Subkeywords are: `qm` (uses the QM molecular dipole from Gaussian calculation), `esp` (uses the dipole computed from the ESP charges computed by Gaussian) and `read` (reads the dipole moment, in a.u., from the `xxx.cns` file (on the next line). See examples below.

dipole	dipole	dipole
qm	esp	read
		0.250 -0.375 0.000

There is no limit on the number of constraints imposed, as long as the conditions are linearly independent. Only one dipole constraint may be imposed. The order the conditions are written in the file is irrelevant, except for the total charge, specified on the first line. See an example of a constraint file below:

⁴Required.

⁵Optional.

⁶Only applied for polarizable ESP fitting.

```

0.00
dipole
read
0.250  -0.375  0.000
equiv
3
1 2 3
fragm
2 1.0
5 6
equiv
2
15 16
equiv
2
18 19

```

Running the program

Compilation Use the compile.csh script to compile the program. Just run

```
./compile
```

from the shell. An executable file named chpol.exe will be created.

What the script does is the following:

```

rm *.o *.exe
pgf90 chpol.f90 -o chpol.exe -llapack -lblas

```

Running Once the program is compiled and the four input files are present, the programme can be executed directly. The input files must be specified when calling the executable, using the following options:

- g (required) Followed by the gesp file name
 - m (required) Followed by the mol2 file name
 - p (required) Followed by the polarizability file name
 - c (required) Followed by the constraints file name
 - h (optional) Print out a help message and quit
 - s (optional) Run in silent mode (minimum printout)
 - d (optional) Run in debug mode (extra printout)
- if neither -s or -d are used, the program will run in normal mode (standard printout)

For instance:

```
./chpol.exe -g xxx.gesp -m xxx.mol2 -p xxx.pol -c xxx.cns -d
```

Theory

The function to minimize is the squared difference between the QM potential \bar{V} calculated on a grid and that generated by a set of charges and dipoles induced by those charges:

$$J = \sum_i^m [\bar{V}_i - V^{\text{chg}} - V^{\text{pol}}]^2 + J_c = \sum_i^m \left[\bar{V}_i - \sum_j^n \frac{q_j}{r_{ij}} - \sum_j^n \frac{\boldsymbol{\mu}_j \cdot \mathbf{r}_{ij}}{r_{ij}^3} \right]^2 + J_c \quad (1)$$

where i runs over the m grid points, j over the n MM atoms, $\mathbf{r}_{ij} \equiv \mathbf{r}_i - \mathbf{r}_j$ and J_c is an additional term eventually including the constraints (see later).

Main function

The minimum of this function is obtained by setting to zero the derivative with respect to the charges:

$$\frac{\partial J}{\partial q_x} = - \sum_i^m 2 [\bar{V}_i - V^{\text{chg}} - V^{\text{pol}}] \left[\frac{\partial}{\partial q_x} \left(\sum_j^n \frac{q_j}{r_{ij}} \right) + \frac{\partial}{\partial q_x} \left(\sum_j^n \frac{\boldsymbol{\mu}_j \cdot \mathbf{r}_{ij}}{r_{ij}^3} \right) \right] \quad (2)$$

Some definitions:

$$A_{ij} \equiv \frac{1}{r_{ij}} \quad (3)$$

$$\mathbf{B}_{ij} \equiv \frac{\mathbf{r}_{ij}}{r_{ij}^3} \quad (4)$$

$$\mathbf{C}_{ij} \equiv \sum_k^n D_{ik}^{-1} \mathbf{B}_{kj} \quad (5)$$

$$E_{ij} \equiv \sum_j^n \mathbf{B}_{ik} \mathbf{C}_{kj} \quad (6)$$

$$F_{ij} \equiv A_{ij} + E_{ij} \quad (7)$$

$$G_{ij} \equiv \sum_k^m F_{ik}^\dagger F_{kj} \quad (8)$$

$$H_i = \sum_j^m F_{ij}^\dagger \bar{V}_j \quad (9)$$

So:

$$\sum_j^n \frac{q_j}{r_{ij}} = \sum_j^n A_{ij} q_j \quad (10)$$

$$\Phi_k = \sum_l^n \frac{q_l \mathbf{r}_{kl}}{r_{kl}^3} = \sum_l^n \mathbf{B}_{kl} q_l \quad (11)$$

$$\mu_j = \sum_k^n \mathbf{D}_{jk}^{-1} \Phi_k = \sum_k^n \sum_l^n \mathbf{D}_{jk}^{-1} \mathbf{B}_{kl} q_l = \sum_l^n \mathbf{C}_{jl} q_l \quad (12)$$

$$\sum_j^n \frac{\mu_j \cdot \mathbf{r}_{ij}}{r_{ij}^3} = \sum_j^n \mathbf{B}_{ij} \mu_j = \sum_j^n \sum_l^n \mathbf{B}_{ij} \mathbf{C}_{jl} q_l = \sum_l^n E_{il} q_l \quad (13)$$

and the derivatives become:

$$\frac{\partial}{\partial q_x} \left(\sum_j^n \frac{q_j}{r_{ij}} \right) = \frac{\partial}{\partial q_x} \left(\sum_j^n A_{ij} q_j \right) = A_{ix} \quad (14)$$

$$\frac{\partial}{\partial q_x} \left(\sum_j^n \frac{\mu_j \cdot \mathbf{r}_{ij}}{r_{ij}^3} \right) = \frac{\partial}{\partial q_x} \left(\sum_j^n E_{ij} q_j \right) = E_{ix} \quad (15)$$

Therefore equation 2 becomes:

$$\begin{aligned} \frac{\partial J}{\partial q_x} &= -2 \sum_i^m \left[\bar{V}_i - \sum_j^n A_{ij} q_j - \sum_j^n E_{ij} q_j \right] [A_{ix} + E_{ix}] = \\ &= -2 \sum_i^m \left[\bar{V}_i - \sum_j^n F_{ij} q_j \right] F_{ix} \end{aligned} \quad (16)$$

Setting this to zero to obtain the minimizing charges gives:

$$\sum_i^m \left[\bar{V}_i - \sum_j^n F_{ij} q_j \right] F_{ix} = 0 \quad (17)$$

$$\sum_i^m \sum_j^n F_{xi}^\dagger F_{ij} q_j = \sum_i^m F_{ix} \bar{V}_i \quad (18)$$

$$\sum_j^n G_{xj} q_j = H_x \quad (19)$$

This can be rewritten in matrix form as:

$$\mathbf{G} \mathbf{q} = \mathbf{H} \quad (20)$$

and the system is solved by:

$$\mathbf{q} = \mathbf{G}^{-1} \mathbf{H} \quad (21)$$

Constraints

Constraints are introduced in function J_c . These are of three different kinds: total charge of a fragment, equivalence between n atoms, and conservation of the total dipole moment. The conservation of the total molecular charge is a particular case of the first kind.

Charge of a fragment The total charge of a fragment F is set to \bar{q} . Therefore the constraint function is:

$$J_c = \lambda \left(\sum_{j \in F} q_j - \bar{q} \right) \quad (22)$$

It follows that:

$$\frac{\partial}{\partial q_x} J_c = \lambda q_x \quad (23)$$

$$\frac{\partial}{\partial \lambda} J_c = \sum_{j \in F} q_j - \bar{q} \quad (24)$$

If the fragment comprises atoms from r to s , the new matrices are:

$$\mathbf{G}_c = \begin{bmatrix} G_{1,1} & G_{1,2} & \cdots & G_{1,r} & \cdots & G_{1,s} & \cdots & G_{1,n} & 0 \\ G_{2,1} & G_{2,2} & \cdots & G_{2,r} & \cdots & G_{2,s} & \cdots & G_{2,n} & 0 \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots \\ G_{r,1} & G_{r,2} & \cdots & G_{r,r} & \cdots & G_{r,s} & \cdots & G_{r,n} & 1 \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots \\ G_{s,1} & G_{s,2} & \cdots & G_{s,r} & \cdots & G_{s,s} & \cdots & G_{s,n} & 1 \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots \\ G_{n,1} & G_{n,2} & \cdots & G_{n,r} & \cdots & G_{n,s} & \cdots & G_{n,n} & 0 \\ 0 & 0 & 0 & 1 & 1 & 1 & 0 & 0 & 0 \end{bmatrix} \quad \mathbf{H}_c = \begin{bmatrix} H_1 \\ H_2 \\ \vdots \\ H_r \\ \vdots \\ H_s \\ \vdots \\ H_n \\ \bar{q} \end{bmatrix} \quad (25)$$

The same holds even in the general case when the fragment comprises non-sequential atoms.

Equivalence of atoms If condition $q_i = q_j = q_k \dots$ holds, this can be split into a system of equations:

$$J_{c1} = \lambda_1 (q_i - q_j)$$

$$J_{c2} = \lambda_2 (q_j - q_k)$$

and so on. For each condition:

$$\frac{\partial}{\partial q_x} J_{c1} = \lambda (\delta_{i,x} - \delta_{j,x}) \quad (26)$$

$$\frac{\partial}{\partial \lambda_1} J_{c1} = q_i - q_j \quad (27)$$

and the new matrices are:

$$\mathbf{G}_c = \begin{bmatrix} G_{1,1} & G_{1,2} & \cdots & G_{1,i} & \cdots & G_{1,j} & \cdots & G_{1,n} & 0 \\ G_{2,1} & G_{2,2} & \cdots & G_{2,i} & \cdots & G_{2,j} & \cdots & G_{2,n} & 0 \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots \\ G_{i,1} & G_{i,2} & \cdots & G_{i,i} & \cdots & G_{i,j} & \cdots & G_{i,n} & 1 \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & 0 \\ G_{j,1} & G_{j,2} & \cdots & G_{j,i} & \cdots & G_{j,j} & \cdots & G_{j,n} & -1 \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & 0 \\ G_{n,1} & G_{n,2} & \cdots & G_{n,i} & \cdots & G_{n,j} & \cdots & G_{n,n} & 0 \\ 0 & 0 & 0 & 1 & 0 & -1 & 0 & 0 & 0 \end{bmatrix} \quad \mathbf{H}_c = \begin{bmatrix} H_1 \\ H_2 \\ \vdots \\ H_i \\ \vdots \\ H_j \\ \vdots \\ H_n \\ 0 \end{bmatrix} \quad (28)$$

Conservation of the total dipole moment If the total dipole moment must be equal to $\bar{\mu}$, this corresponds to three conditions:

$$J_c = \lambda \left(\sum_j \mu_j + \mu^{\text{chg}} - \bar{\mu} \right) \quad (29)$$

where μ_j are the induced atomic dipole moments and μ^{chg} is the fixed dipole given by the charges. Note that when non-polarizable ESP charges are being calculated, the induced dipole moments are assumed to be zero. In such case, the equations below still hold, with $\mathbf{C}_{jk} = 0$ and $\tilde{\mathbf{C}}_k = 0$, and therefore with $\mathbf{K}_k = \mathbf{r}_k$. Equation 29 can be rewritten as:

$$\begin{aligned} J_c &= \lambda \left(\sum_j \left(\sum_k \mathbf{C}_{jk} q_k + \mathbf{r}_j q_j \right) - \bar{\mu} \right) = \\ &= \lambda \left(\sum_k \left(\tilde{\mathbf{C}}_k q_k + \mathbf{r}_k q_k \right) - \bar{\mu} \right) = \lambda \left(\sum_k \mathbf{K}_k q_k - \bar{\mu} \right) \end{aligned} \quad (30)$$

where $\tilde{\mathbf{C}}_k \equiv \sum_j \mathbf{C}_{jk}$ and $\mathbf{K}_k \equiv \tilde{\mathbf{C}}_k + \mathbf{r}_k$. The derivatives are:

$$\frac{\partial}{\partial q_x} J_c = \lambda \mathbf{K}_x \quad (31)$$

$$\frac{\partial}{\partial \lambda} J_c = \sum_k \mathbf{K}_k q_k - \bar{\mu} \quad (32)$$

The new matrices are:

$$\mathbf{G}_c = \begin{bmatrix} G_{1,1} & G_{1,2} & \cdots & G_{1,n} & K_1^{(x)} & K_1^{(y)} & K_1^{(z)} \\ G_{2,1} & G_{2,2} & \cdots & G_{2,n} & K_2^{(x)} & K_2^{(y)} & K_2^{(z)} \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots \\ G_{n,1} & G_{n,2} & \cdots & G_{n,n} & K_n^{(x)} & K_n^{(y)} & K_n^{(z)} \\ K_1^{(x)} & K_2^{(x)} & \cdots & K_n^{(x)} & 0 & 0 & 0 \\ K_1^{(y)} & K_2^{(y)} & \cdots & K_n^{(y)} & 0 & 0 & 0 \\ K_1^{(z)} & K_2^{(z)} & \cdots & K_n^{(z)} & 0 & 0 & 0 \end{bmatrix} \quad \mathbf{H}_c = \begin{bmatrix} H_1 \\ H_2 \\ \vdots \\ H_n \\ \bar{\mu}^{(x)} \\ \bar{\mu}^{(y)} \\ \bar{\mu}^{(z)} \end{bmatrix} \quad (33)$$