

# Hermitian operators in quantum mechanics

they allow us to define physical observables, and their eigenvalues are real numbers. and the eigenstate of a Hermitian operator form a basis in the state space.

$\hat{A} = \hat{A}^+$  an operator which is equal to its adjoint.

$$\hat{A}|\psi\rangle = |\psi'\rangle \quad \longleftrightarrow \quad \langle\psi|A^+ = \langle\psi'|$$
$$\longleftrightarrow \quad \langle\psi|\hat{A} = \langle\psi'|$$

① Eigenvalues of a Hermitian operator are real numbers since they represent physical observable and hence should have real eigenvalues.

② Eigenstates of an Hermitian operator can be chosen to form an orthonormal set

we can always write an arbitrary quantum state in the representation given by the eigenstate of a physical observable.

$$A|\psi\rangle = \lambda|\psi\rangle \quad \sim \text{Eigenvalue Equation}$$

$$A|\psi\rangle = \lambda|\psi\rangle \implies \langle\psi|A|\psi\rangle = \lambda\langle\psi|\psi\rangle$$
$$\langle\psi| \xrightarrow{\phantom{x}} \lambda$$

$$\langle\psi|A^+ = \lambda^*\langle\psi|$$

$$\langle\psi|\hat{A} = \lambda^*\langle\psi|$$

we just need to show that  $\lambda^* = \lambda$  and hence eigenvalues are  $\mathbb{R}$ .

$$\langle \psi | \hat{A} | \psi \rangle = \underbrace{\lambda^* \langle \varphi | \psi \rangle}_{\geq 0} \xrightarrow{\lambda = \lambda^*} \lambda \in \mathbb{R}$$

$$\lambda \langle \psi | \psi \rangle = \lambda^* \langle \psi | \psi \rangle$$

$\equiv 0$  iff  $|\psi\rangle = 0$

Eigenstates

$$\hat{A} |\psi\rangle = \lambda |\psi\rangle$$

$$\lambda \neq \mu$$

$$\hat{A} |\psi\rangle = \lambda |\psi\rangle \quad \Rightarrow \quad \boxed{\langle \psi | \hat{A} | \psi \rangle = \lambda \langle \psi | \psi \rangle}$$

$$A |\psi\rangle = \mu |\psi\rangle \quad \exists \langle \psi | A = \mu \langle \psi |$$

$$\langle \psi | A = \mu \langle \psi | \xrightarrow[\in |\psi\rangle]{\hat{A} = A^+} \boxed{\langle \psi | \hat{A} | \psi \rangle = \mu \langle \psi | \psi \rangle}$$

$$\text{so } \lambda \langle \psi | \psi \rangle = \mu \langle \psi | \psi \rangle$$

$$(\lambda - \mu) \langle \psi | \psi \rangle = 0$$

$\neq 0 \quad \xrightarrow{=} 0^-$

$$\langle \psi | \psi \rangle = 0$$

$\therefore$  eigenstate of Hermitian operator, corresponding to different eigenvalues form an orthonormal pair.

### Degenerate Eigenvalues

What happens to the eigenstates then?

$$\textcircled{1} \quad \hat{A} |\psi^i\rangle = \lambda |\psi^i\rangle \quad i=1, 2, \dots, n$$

the  $n$ -eigenstates that share the same eigenvalues  $\lambda$ .

$$|\Psi\rangle = \sum c_i |\psi^i\rangle$$

Linear combination of these also have the same eigenvalue

$$\hat{A} |\Psi\rangle = \lambda |\Psi\rangle$$

For an  $n$ -fold degenerate eigenvalue, any ket that lives in the  $n$ -dim subspace of state space spanned by the eigenkets  $|\psi^i\rangle$  is also a valid eigenstate.

$\textcircled{2}$  For  $n$ -fold degenerate eigenvalues, there are  $n$  linearly independent eigenstates.

$n$ -dim state space

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$n$  linearly independent eigenstates  
(may not be mutually orthogonal)

How to make all of them mutually orthogonal?

Gram-Schmidt orthonormalization

$$\sum |\psi^i\rangle \rightarrow \sum |\psi^i\rangle$$

set of linearly independent  $\rightarrow$  mutually orthogonal  
but not necessarily  
orthogonal

Orthonormal states

$$\langle \psi^i | \psi^j \rangle = \delta_{ij}$$

$$\textcircled{6} \quad |\psi^P\rangle = \frac{|\psi^i\rangle}{\sqrt{\langle \psi^i | \psi^i \rangle}}$$

and sum the states

$\sqrt{\langle \psi' | \psi' \rangle} \rightarrow$  to make sure the basis is normalized.

②  $|\chi^2\rangle = |\psi^2\rangle + \alpha |\psi'\rangle \rightarrow$  we want  $|\chi^2\rangle$  to be orthogonal to  $|\psi'\rangle$

$$\langle \psi' | \chi^2 \rangle = 0$$

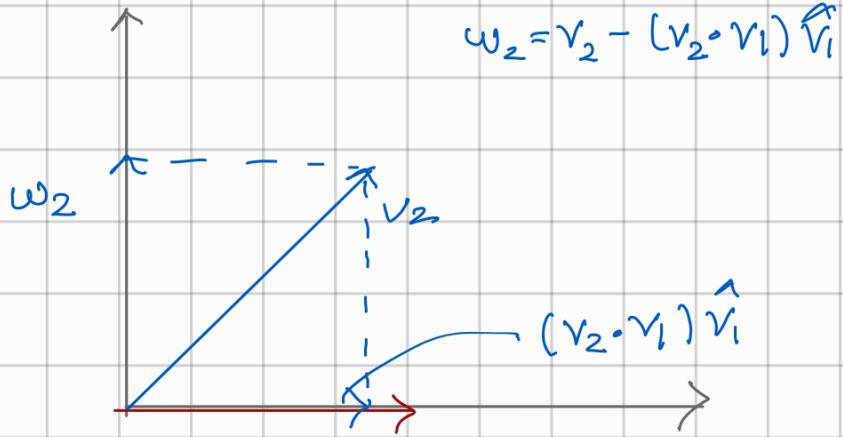
$$\Rightarrow \langle \psi' | \psi^2 \rangle + \alpha \langle \psi' | \psi' \rangle = 0$$

$$\text{if we choose } \alpha = -\langle \psi' | \psi^2 \rangle$$

Replace this in ②

$$|\chi^2\rangle = |\psi^2\rangle - |\psi'| \times \langle \psi' | \psi^2 \rangle$$

$$|\psi^2\rangle = \frac{|\chi^2\rangle}{\sqrt{\langle \chi^2 | \chi^2 \rangle}} \quad |\psi^2\rangle \text{ is parallel to } |\chi^2\rangle$$



Suppose  $v_1$  and  $v_2$  are linearly independent, but not orthogonal. So keeping  $v_1$  as reference, we want to make a slate which is orthonormalized.

Repeat the same procedure.

Matrix formulation  
~ transpose conjugate matrix

$$\hat{A} = A_{ij}^*$$

$$\hat{A}^+ = A_{ji}^*$$

$$A_{ij}^* = A_{ji}^{*\dagger}$$

$$\begin{pmatrix} A_{11} & A_{12} & A_{13} & \cdots \\ A_{21} & A_{22} & A_{23} & \cdots \\ A_{31} & A_{32} & A_{33} & \cdots \\ \vdots & \vdots & \vdots & \ddots \end{pmatrix}$$

$$i=j \Rightarrow A_{ii} = A_{ii}^* \Rightarrow A_{ii} \in \mathbb{R}$$

$$i \neq j \Rightarrow A_{ij}^* = A_{ji}^T$$

$$\downarrow$$

$$\begin{pmatrix} A_{11} & A_{12} & A_{13} & \cdots \\ A_{12}^* & A_{22} & A_{23} & \cdots \\ A_{13}^* & A_{23}^* & A_{33} & \cdots \\ \vdots & \vdots & \vdots & \ddots \end{pmatrix}$$

$$\hat{A} |\psi_m\rangle = \lambda_m |\psi_m\rangle, \langle \psi_m | \psi_m \rangle = \delta_{mm}$$

$$\lambda_{mm} = \langle \psi_m | \hat{A} | \psi_m \rangle \rightarrow \lambda_m \langle \psi_m | \psi_m \rangle$$

—————  
 $\lambda_m |\psi_m\rangle \Rightarrow \lambda_m \delta_{mm}$

matrix is diagonal  
 and the diagonal values are just the eigenvalues

$$\begin{pmatrix} \lambda_1 & 0 & 0 & 0 & \cdots \\ 0 & \lambda_2 & 0 & 0 & \cdots \\ 0 & 0 & \lambda_3 & 0 & \cdots \\ \vdots & \vdots & \vdots & \ddots & \vdots \end{pmatrix}$$

$\lambda_i \in \mathbb{R}$

$\lambda_m \in M$ .

$$\begin{pmatrix} \lambda_2 & 0 & 0 & \cdots \\ 0 & \lambda_3 & 0 & \cdots \\ 0 & 0 & \ddots & \lambda_4 \\ \vdots & \vdots & & \ddots \end{pmatrix}$$

the matrix associate with an operator  $\hat{A}$  is diagonal when written in the basis spanned by the eigenstate of  $\hat{A}$ , and the diagonal entries are the eigenvalues of the operator.