Engineering Mathematics 2B Module 3: Differentiation

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Motivation:

There are some universal principles underpinning phenomena in electromagnetics, heat transfer, and other fields.

We have already seen that in electrostatics $\mathbf{E} = -\nabla \phi$, and this implies that \mathbf{E} is conservative. In module 2 however we said that if \mathbf{E} is conservative then $\nabla \times \mathbf{E} = 0$.

We also saw that the magnetic field is always solenoidal, that is $\nabla \cdot \mathbf{H} = 0$. Maxwell has shown that $\nabla \times \mathbf{E} = -\mathrm{i}\omega \mu \mathbf{H}$, that is the magnetic field is the curl of the electric field multiplied by a constant. (The reverse also holds).

These aren't mere coincidences. In fact, the curl of a gradient field is always a zero field, and the divergence of a curl field is always zero.

Harmonic fields

Laplace's equation is ubiquitous in engineering and physics. It describes potential in electrostatics, temperature in steady-state heat transfer, and flow in porous media (among others)

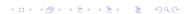
$$\nabla \cdot \nabla f = 0$$

The notation

$$\nabla \cdot \nabla = \nabla^2 = \Delta = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}$$

allows to write this as $\nabla^2 f = 0$, with ∇^2 known as the Laplace operator. Any scalar field f that satisfies this equation is called harmonic.

Recall that we have defined the vector operator $\nabla \doteq (\partial/\partial x, \partial/\partial y)$ so you can verify the above.



¹In 2D. The extension to 3D is trivial.

Harmonic fields

Consider an infinitely large 2D (resp. 3D) object Ω that we heat using a high power, high precision laser beam at point \mathbf{r}' . The temperature f in the object is a scalar field given by

$$\nabla^2 f = -\delta(\mathbf{r} - \mathbf{r}'), \quad \forall \mathbf{r} \in \Omega.$$

thus at any point other than \mathbf{r}' it satisfies Laplace's equation

$$\nabla^2 f = 0, \quad \mathbf{r} \in \Omega \setminus \mathbf{r}'.$$

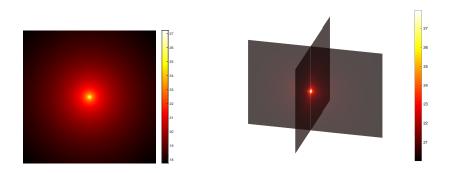
If Ω is 2D then the temperature in the object is given by the harmonic field

$$f_{\mathrm{2D}} = -\frac{1}{2\pi} \log |\mathbf{r} - \mathbf{r}'|,$$

while in 3D

$$f_{3D} = \frac{1}{4\pi |\mathbf{r} - \mathbf{r}'|}$$

Temperature harmonic fields



Plots of $f_{\rm 2D}$ and $f_{\rm 3D}$ assuming ${\bf r'}$ is the origin and the object's initial temperature was 17 C.

Harmonic fields

A graphical plot of the harmonic field is reassuring, but we still need to verify $\nabla^2 f = 0$ mathematically.

To compute the Laplacian of $f = -\frac{1}{2\pi} \log(\sqrt{x^2 + y^2})$ recall the partial derivatives of the 2D distance function $\frac{\partial}{\partial x} \log(\sqrt{x^2 + y^2}) = \frac{x}{x^2 + y^2}$, and $\frac{\partial}{\partial y} \log(\sqrt{x^2 + y^2}) = \frac{y}{x^2 + y^2}$.

$$\nabla^2 f = \nabla \cdot \nabla f = \frac{\partial}{\partial x} \left(\frac{\partial f}{\partial x} \right) + \frac{\partial}{\partial y} \left(\frac{\partial f}{\partial y} \right)$$
$$= \frac{\partial}{\partial x} \frac{x}{x^2 + y^2} + \frac{\partial}{\partial y} \frac{y}{x^2 + y^2}$$
$$= \frac{x^2 + y^2 - 2x^2}{(x^2 + y^2)^2} + \frac{x^2 + y^2 - 2y^2}{(x^2 + y^2)^2} = 0$$

This holds everywhere apart from (0,0).

Try the 3D case for f_{3D} as an exercise.



The curl of a gradient

In module 2 we have established that if the curl of a continuous vector field \mathbf{a} is zero, then there must exist a scalar potential function f such that $\mathbf{a} = \nabla f$. This can be read in two ways:

If
$$\mathbf{a} = \nabla f$$
 then $\nabla \times \mathbf{a} = 0^2$

If $\nabla \times \mathbf{a} = 0$ then there must exists f such that $\mathbf{a} = \nabla f$

This isn't just a piece of mathematical jargon. It asserts for example that in electrostatics the curl of the electric field has to be zero, since we have electric potential.

If we magnetise a non-metallic body (e.g. a human inside an MRI scanner) then there are no currents flowing through. Since Maxwell gives $\mathbf{J} = \nabla \times \mathbf{H}$, then $\mathbf{J} = 0$ means that there must exist a magnetic potential, $\mathbf{H} = -\nabla \phi$.



 $^{^{2}}$ This is a vector field (0,0,0).

The curl of a gradient proof

The identify $\nabla \times \nabla f = 0$ holds true for any continuous f.

$$\nabla \times \nabla \mathbf{f} = \begin{vmatrix} \hat{\mathbf{i}} & \hat{\mathbf{j}} & \hat{\mathbf{k}} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ \frac{\partial \mathbf{f}}{\partial x} & \frac{\partial \mathbf{f}}{\partial y} & \frac{\partial \mathbf{f}}{\partial z} \end{vmatrix}$$

$$= \left(\frac{\partial^2 f}{\partial y \partial z} - \frac{\partial^2 f}{\partial y \partial z} \right) \hat{\mathbf{i}} + \left(\frac{\partial^2 f}{\partial z \partial x} - \frac{\partial^2 f}{\partial z \partial x} \right) \hat{\mathbf{j}} + \left(\frac{\partial^2 f}{\partial x \partial y} - \frac{\partial^2 f}{\partial x \partial y} \right) \hat{\mathbf{k}}$$

$$= 0 \hat{\mathbf{i}} + 0 \hat{\mathbf{j}} + 0 \hat{\mathbf{k}}.$$

All conservative, irrotational, gradient fields have zero curl.

The curl of the wind



Simulation of a wind velocity field courtesy of John Hopkins University.

Wind fields are hardly ever conservative, and that's a good thing when aiming to spin turbine blades.

The divergence of the curl

The divergence of the curl of a vector field is always zero.

This statement can be posed in two ways:

If
$$\mathbf{a} = \nabla \times \mathbf{b}$$
 then $\nabla \cdot \mathbf{a} = 0$

If $\nabla \cdot \mathbf{a} = 0$ then there must exists **b** such that $\mathbf{a} = \nabla \times \mathbf{b}$

Maxwell's state that $\nabla \cdot \mathbf{B} = 0$, where **B** is the magnetic flux density, related to the magnetic field by $\mathbf{B} = \mu \mathbf{H}$. This in turn justifies the co-existence of magnetic and electric fields, since

$$\mathbf{H} = \frac{\mathrm{i}}{\omega \mu} \nabla \times \mathbf{E}.$$

The divergence of a curl proof

The identify $\nabla \cdot \nabla \times \mathbf{a} = 0$ holds true for any continuous \mathbf{a} .

Let $\mathbf{a} = f(x, y, z)\mathbf{\hat{i}} + g(x, y, z)\mathbf{\hat{j}} + h(x, y, z)\mathbf{\hat{k}}$, then

$$\nabla \cdot \nabla \times \mathbf{a} = \nabla \cdot \left(\left(\frac{\partial h}{\partial y} - \frac{\partial g}{\partial z} \right) \hat{\mathbf{i}} + \left(\frac{\partial f}{\partial z} - \frac{\partial h}{\partial x} \right) \hat{\mathbf{j}} + \left(\frac{\partial g}{\partial x} - \frac{\partial f}{\partial y} \right) \hat{\mathbf{k}} \right)$$

$$= \frac{\partial}{\partial x} \left(\frac{\partial h}{\partial y} - \frac{\partial g}{\partial z} \right) + \frac{\partial}{\partial y} \left(\frac{\partial f}{\partial z} - \frac{\partial h}{\partial x} \right) + \frac{\partial}{\partial z} \left(\frac{\partial g}{\partial x} - \frac{\partial f}{\partial y} \right)$$

$$= \frac{\partial^2 h}{\partial x \partial y} - \frac{\partial^2 h}{\partial y \partial x} - \frac{\partial^2 g}{\partial x \partial z} + \frac{\partial^2 g}{\partial z \partial x} + \frac{\partial^2 f}{\partial y \partial z} - \frac{\partial^2 f}{\partial z \partial y} = 0$$

since for any function w(x, y, z)

$$\frac{\partial^2 w}{\partial x \partial y} = \frac{\partial^2 w}{\partial y \partial x}, \quad \frac{\partial^2 w}{\partial x \partial z} = \frac{\partial^2 w}{\partial z \partial x}, \quad \frac{\partial^2 w}{\partial y \partial z} = \frac{\partial^2 w}{\partial z \partial y}$$

Formulas

Let $\mathbf{F} = (f, g, h)$ a continuous vector field.

ightharpoonup The Laplacian of f is

$$\Delta f = \nabla^2 f = \nabla \cdot \nabla f = \frac{\partial^2 f}{\partial x^2} + \frac{\partial^2 f}{\partial y^2} + \frac{\partial^2 f}{\partial z^2}$$

- ▶ Harmonic fields have $\Delta f = 0$
- $\nabla \times \nabla f = 0$
- $\nabla \cdot \nabla \times \mathbf{F} = 0$

Main outcomes of module 3

You MUST know:

- 1. What are the scalar harmonic fields and how they are related to conservative vector fields.
- 2. The curl of a gradient field is always zero.
- 3. The divergence of a curl field is always zero.
- 4. To compute second partial derivatives.
- 5. The notation $\operatorname{curl} \mathbf{a} \equiv \nabla \times \mathbf{a}$, $\operatorname{div} \mathbf{a} \equiv \nabla \cdot \mathbf{a}$, $\operatorname{grad} f \equiv \nabla f$.

Good to know:

Try to find examples of irrotational and solenoidal vector fields relevant to your own discipline.