

## Examples: ABJM Theory

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There's not a whole lot to say until we write the formula:

$$\int \left( \prod_i e^{ik\pi(\theta_i^2 - \phi^2)} \right) \times \frac{\prod_{i \neq j} (2 \sinh \pi(\theta_i - \theta_j) 2 \sinh \pi(\phi_i - \phi_j))}{\prod_{i,j} (2 \cosh \pi(\theta_i - \phi_j))^2}$$

This integral could possibly be over  $[0, 2\pi]^2$  but I am expecting it's over the real numbers:

$$(\theta, \phi) \in \mathbb{R}^N \times \mathbb{R}^N$$

since this the domain of integration of the reals. Here is the simplest one:

$$\int_{-\infty}^{\infty} d\theta e^{ik\pi\theta^2} = \sqrt{\frac{\pi}{8}}(1 + i)$$

I have not done that yet. If we set  $N = 1$  or  $N = 2$  and this integral is totally feasible.

ABJM theory could refer to two things:

- The  $U(N) \times U(N)$  Chern-Simons-matter theory with  $\mathcal{N} = 6$  superconformal symmetry. The first  $U(N)$  is at level  $k$  and the second  $U(N)$  is at level  $-k$ .
- The  $U(N) \times U(N)$  Gaussian matrix integral the CS-theory localizes to.

Both of these are called **ABJM theory**. This is rather confusing.

Here is another Chern-Simons formula:

$$Z_{\text{CS}}(N, k) = \frac{1}{\sqrt{N+k}} \prod_{\alpha > 0} 2 \sin \frac{\pi \alpha \cdot \rho}{k+N}$$

where  $\alpha_{ij} = e_i - e_j \in \mathbb{C}^N$  are vectors<sup>1</sup> and

$$\rho = \frac{1}{2} \sum_{\alpha > 0} \alpha = \sum_{i=1}^N \left( \frac{N+1}{2} - i \right) e_i$$

and also  $k \in \mathbb{Z}$  is a positive integer, and so is  $n \in \mathbb{Z}$ .

Therefore our Chern-Simons partition function  $Z_{\text{CS}}$  should be a number.

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<sup>1</sup>This is called the **Cartan subalgebra** a term which should mean nothing right now.

My understanding is that any localization result of a supersymmetric gauge theory that is not **flat** is indebted to Vasily Pestun<sup>2</sup>.

Flat spaces are things like  $\mathbb{R}^4$  and even  $\mathbb{R}_{\epsilon_1, \epsilon_2}^4$  which are **distorted** versions of flat 4-dimensional space.

The curved spaces being considered are remarkably simple  $\mathbf{S}^3$  the 3-sphere and  $\mathbf{S}^4$  the 4-sphere<sup>3</sup>

We get hints from Pestun's original paper:

equivariant Euler class of the infinite-dimension  
normal bundle to the localization locus

I don't know what this means. Maybe some complicated infinite dimensional object has been erected over our 4-sphere,  $S^4$  ?

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<sup>2</sup>This is just looking at the citations. The 4-sphere is a curved four-dimensional space, I guess.

<sup>3</sup>In either case the procedure is the same:

- Do a very complicated algebra
- Claim a localization result
- Solve the integral

I will have a little bit to say about each of these steps. Here are some old ideas that may help:

- Invariant Theory, Spherical Harmonics
- Symplectic Geometry, ODE
- Integral Geometry

The Chern-Simon's matter action is involved:

$$S = \frac{k}{4\pi} \int \text{Tr}(A \wedge dA + \frac{2}{3}A^3 - \bar{\chi}\chi + 2D\sigma) + W$$

called a **multiplet** and a **superpotential**

$$W = \frac{2\pi}{k} \epsilon^{ab} \epsilon_{\dot{a}\dot{b}} \text{Tr}(A_a B_{\dot{a}} A_b B_{\dot{b}})$$

and we're somehow going to localize this.

## Why Path Integrals Localize to Gaussians<sup>4</sup>

An action  $S$  depends on a function. By calculus of variations:

$$S(x + \delta x) = S(x) + \delta x S'(x) + \frac{(\delta x)^2}{2} S''(x)$$

Our job should be to set  $S'(x) = 0$  this is our critical point. Feynman writes the path integral:

$$Z = \sum e^{S(x)}$$

However our logic is the same, this sum should concentrate *near* the classical trajectory:

$$Z \approx e^{x_{\text{cl}}} \sum_{x=x_{\text{cl}}+t\delta x} e^{\frac{1}{2}(t\delta x)^2 S''(x_{\text{cl}})}$$

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<sup>4</sup>As soon as you hear the logic two things should happen:

- You should remember this has appeared in every single quantum mechanics textbook.
- Realize this is entirely bull-shit.

Pestun cites the **Duistermaat-Heckman** theorem<sup>5</sup>

$$\int_M \frac{\omega^n}{n!} e^{-\mu} = \sum_i \frac{e^{-\mu(x_i)}}{e(x_i)}$$

Here are some notations we might need:

- $M$  is a symplectic compact manifold
- $\omega$  is a symplectic form
- $\dim M = 2n$  the number of dimensions
- $\mu$  is the moment map of the  $U(1)$  action<sup>6</sup>
- $x_i$  are the fixed points of the rotation
- $e(x_i)$  is the product of the **weight** of the  $U(1)$  action on the tangent space at  $x_i$

$$\int_M \alpha = \sum_i \frac{\pi^n \alpha_0(x_i)}{\sqrt{\det(\partial_\mu V^\nu(x_i))}}$$

This **Berline-Vergne-Atiyah-Bott** formula is quite robust. The low-end version of this formula is:

$$\int e^{t f(x)} dx \approx \frac{e^{t f(x_0)}}{\sqrt{f''(x_0)}}$$

Because  $f(x) \approx f(x_0) + \frac{t^2}{2} f''(x_0)$  near  $x \approx x_0$ .

<sup>5</sup>This is a formula from **Symplectic Geometry** or closely related to **Hamiltonian Mechanics**.

<sup>6</sup>It is a **rotation**.

## Some Hamiltonian Mechanics

Consider  $n$  particles with mass  $m_i$  and with phase space  $X = \mathbb{R}^{2n}$  and total angular momentum:

$$\omega = \sum_{i=1}^n dx^i \wedge dy_i$$

and consider the rotation in phase space:

$$\begin{pmatrix} x \\ y \end{pmatrix} \mapsto \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix}$$

with  $\theta \in [0, 2\pi]$ . This are  $n$  independent harmonic oscillators. The total energy of the  $n$  particles is:

$$\mu = \sum m_i(x_i^2 + y_i^2)$$

and mysteriously the Duistermaat Heckman formula produces a Gaussian integral:

$$\frac{1}{n!} \int_{\mathbb{R}^n} (dx^1 \wedge dy_1) \wedge \cdots \wedge (dx^n \wedge dy_n) e^{-\epsilon \sum (x_i^2 + y_i^2)}$$

The pair  $X = \mathbb{R}^{2n}$  and  $\omega = \sum (dx \wedge dy)$  could be the phase space of  $n$  harmonic oscillators.

And we achieve a Gaussian distribution.

$$= \frac{(2\pi)^n \exp(-\epsilon)}{\prod_{i=1}^n w_i}$$

The geometric object used here is the Euler class, rotation each particle  $(x^i, y_i)$  by an angle  $\theta_i$ .

The Hamiltonian is  $H = \mu$ . The DH formula “quantizes” this rotation into  $n$  random walks (or  $n$  copies of the **free particle**).

## References

- (1) Anton Kapustin, Brian Willett, Itamar Yaakov. **Exact Results for Wilson Loops in Superconformal Chern-Simons Theories with Matter**  
<https://terrytao.wordpress.com/2009/08/23/determinantal-processes/>
- (2) Ofer Aharony, Oren Bergman, Daniel Louis Jafferis, Juan Maldacena **N=6 superconformal Chern-Simons-matter theories, M2-branes and their gravity duals** arXiv:0806.1218v4
- (3) Vasily Pestun, Maxim Zabzine. **Introduction to Localization in Quantum Field Theory.** arXiv:1608.02953
- (4) Richard Feynman **Path Integrals and Quantum Mechanics** Dover, 2010.
- (5) Vasily Pestun **Review of Localization in Geometry** arXiv:1608.02954