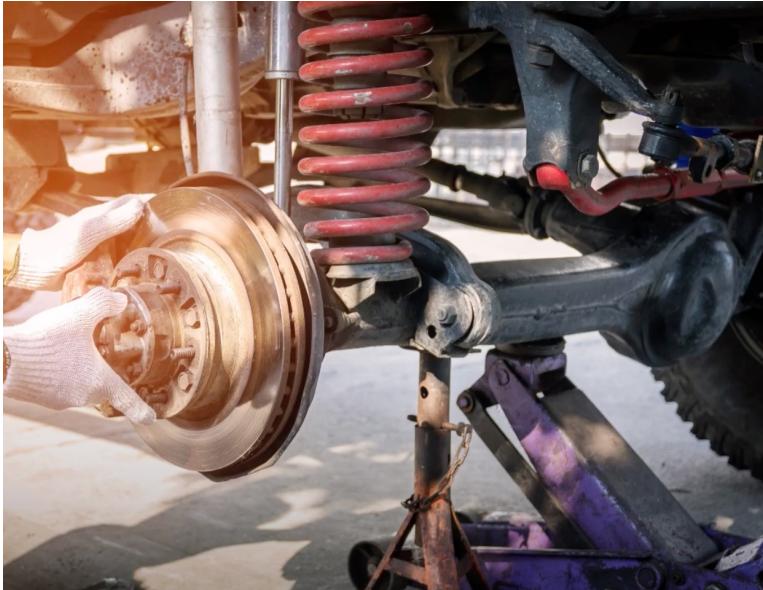


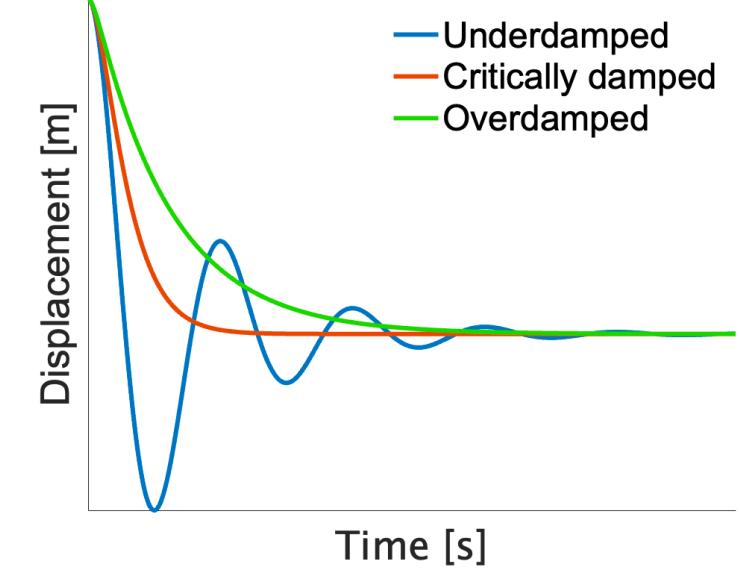
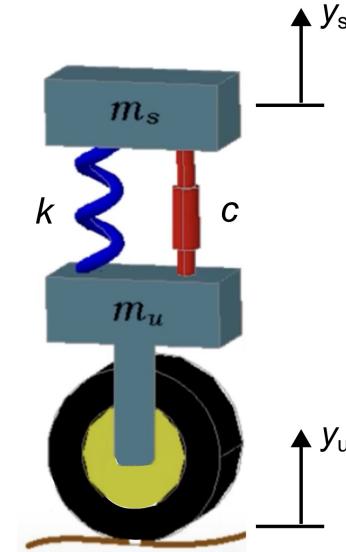
# **Analysis of Single Degree of Freedom Mass–Spring–Damper System: Free Vibration**

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## Automobile Suspension System

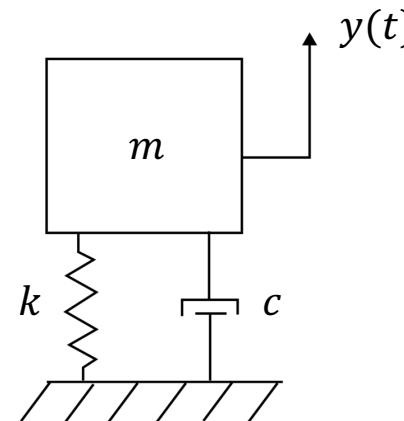


**Mass + Spring + Damper**



- ✓ **Critically damped:** Minimise oscillations & guarantee responsiveness (quick return to rest position).
- ✗ **Underdamped:** undesired oscillations.
- ✗ **Overdamped:** unresponsive suspension system.

## Single Degree of Freedom Spring-Mass-Damper System:



with:

- coordinate  $y(t)$  with time derivatives  $\dot{y}(t)$  and  $\ddot{y}(t)$ ,

- inertia , mass  $m$  [ $N \text{ m}^{-1}\text{s}^2$ ],

- linear spring , coefficient  $k$  [ $N \text{ m}^{-1}$ ],

- viscous damper , coefficient  $c$  [ $N \text{ m}^{-1} \text{ s}$ ],

## Hypotheses:

- single coordinate describes motion;
- lumped mass system;
- linear elasticity for spring;
- linear viscous damping;
- time invariant parameters;
- no additional forces (e.g., friction).

## Find:

Free motion of the mass with initial conditions  
 $y_0$  and  $\dot{y}_0 \rightarrow$  homogeneous solution.

**Homogeneous solution:**  
 right side  $f(t) = 0 \Rightarrow y_h(t)$

**Particular solution:**  
 right side  $f(t) \neq 0 \Rightarrow y_p(t)$

Solution of ODE:

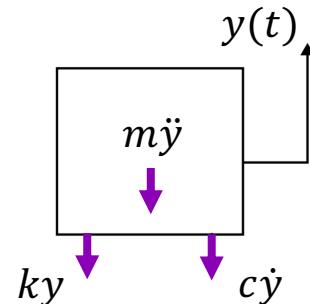
**Total solution:**

$$y_{\text{tot}}(t) = y_h(t) + y_p(t)$$

a) Equation of motion: D'Alembert's principle

- Newton's second law,

$$m \ddot{y}(t) + c \dot{y}(t) + k y(t) = 0$$



$m \ddot{y}(t)$ : inertia force (opposes acceleration)

$c \dot{y}(t)$ : damping force (dissipates energy)

$k y(t)$ : spring force (restores position)

b) Initial conditions:

$$y_0 = y(t_0); \dot{y}_0 = \dot{y}(t_0)$$

Normalised form:

$$\ddot{y}(t) + 2\zeta\omega_n \dot{y}(t) + \omega_n^2 y(t) = 0$$

- undamped natural frequency,

$$\omega_n = \sqrt{k/m} \quad [\text{rad s}^{-1}]$$

- damping ratio,

$$\zeta = \frac{c}{c_{\text{crit}}} \quad [-]$$

$$c_{\text{crit}} = 2\sqrt{m k} \quad [\text{N m}^{-1} \text{s}]$$

1. Assume solution, substitute:  $y(t) = C e^{\lambda t}; \quad \dot{y}(t) = C \lambda e^{\lambda t}; \quad \ddot{y}(t) = C \lambda^2 e^{\lambda t}$

2. Characteristic equation:  $\lambda^2 + 2\zeta\omega_n \lambda + \omega_n^2 = 0 \rightarrow \lambda_{1,2} = -\zeta\omega_n \pm \omega_n\sqrt{\zeta^2 - 1}$ .

3. Two roots ( $\lambda_{1,2}$ ), two solutions:  $y_1(t) = C_1 e^{\lambda_1 t}; \quad y_2(t) = C_2 e^{\lambda_2 t}$

4. General solution:  $y(t) = C_1 e^{[-\zeta\omega_n + \omega_n\sqrt{\zeta^2 - 1}] t} + C_2 e^{[-\zeta\omega_n - \omega_n\sqrt{\zeta^2 - 1}] t}$

$C_1; C_2$

constants from initial conditions.

- No damping:  $\zeta = 0$

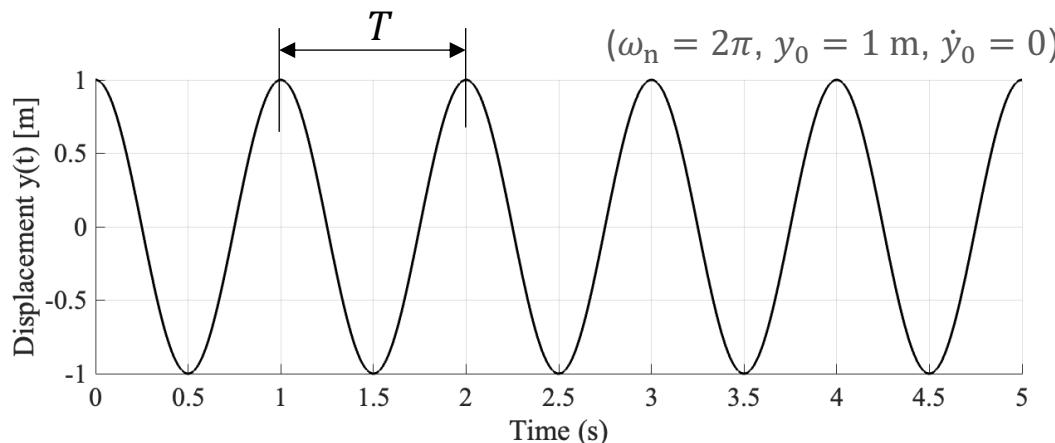
- Solution for complex conjugate roots without real part:

$$y(t) = C_1 e^{[-\zeta \omega_n + i \omega_n \sqrt{1-\zeta^2}] t} + C_2 e^{[-\zeta \omega_n - i \omega_n \sqrt{1-\zeta^2}] t}$$

- Express as harmonic oscillation (Euler's equation\*):

$$y_h(t) = C'_1 \cos(\omega_n t) + C'_2 \sin(\omega_n t)$$

$$\begin{aligned} y_0 &= y(t_0); & \rightarrow C'_1 &= y_0; \\ \dot{y}_0 &= \dot{y}(t_0) & C'_2 &= \frac{\dot{y}_0}{\omega_n} \end{aligned}$$



Note(\*):  $e^{i\alpha} = \cos \alpha + i \sin \alpha$

- Natural frequency without damping:

$$\omega_n = \sqrt{k/m}$$

$$T = \frac{2\pi}{\omega_n}$$

$$\uparrow k \Rightarrow \uparrow \omega_n$$

$$\uparrow m \Rightarrow \downarrow \omega_n$$

Application:

- baseline for more complex systems  
(with damping and forcing).

- Underdamped case:  $0 < \zeta < 1$

$$c < c_{\text{crit}} = 2\sqrt{m k}$$

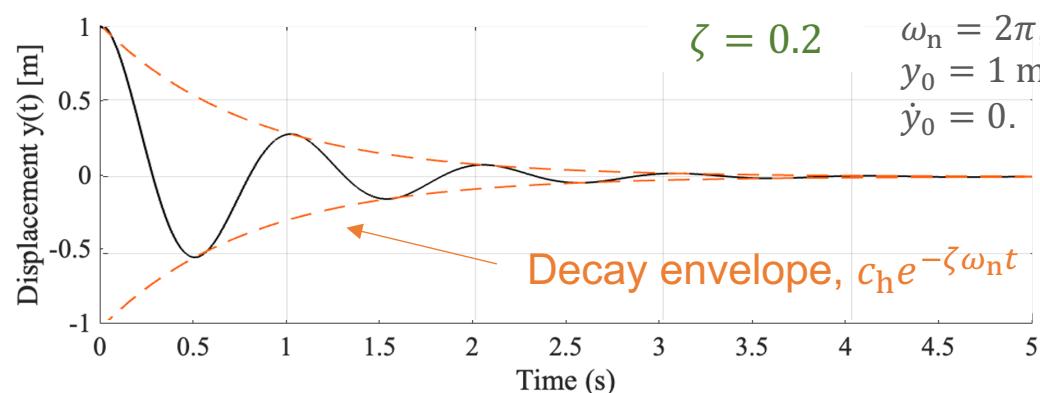
- Solution for complex conjugate roots, negative real part:

$$y(t) = C_1 e^{[-\zeta \omega_n + i \omega_n \sqrt{1-\zeta^2}] t} + C_2 e^{[-\zeta \omega_n - i \omega_n \sqrt{1-\zeta^2}] t}$$

- Express as harmonic oscillation (Euler's equation):

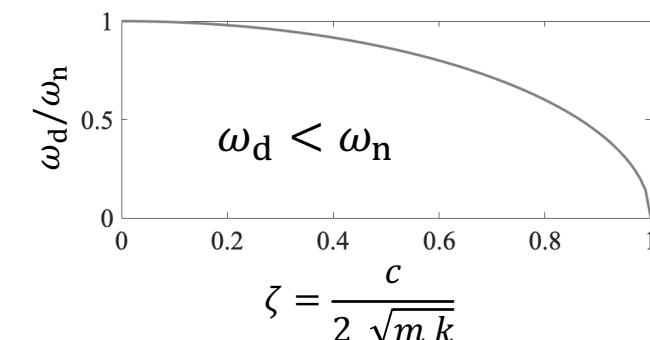
$$y(t) = e^{-\zeta \omega_n t} [C'_1 \cos(\omega_d t) + C'_2 \sin(\omega_d t)]$$

$$\begin{aligned} y_0 &= y(t_0); & C'_1 &= y_0; \\ \dot{y}_0 &= \dot{y}(t_0) & C'_2 &= \frac{\dot{y}_0 + \zeta \omega_n y_0}{\omega_d} \end{aligned}$$



- Frequency of damped vibration:

$$\omega_d = \omega_n \sqrt{1 - \zeta^2}$$



- Damping → lowers frequency, reduces amplitude.

## Applications:

- most automobile suspension,
- vibrating mechanical systems,
- musical instruments.

- Critically damped case:  $\zeta = 1$

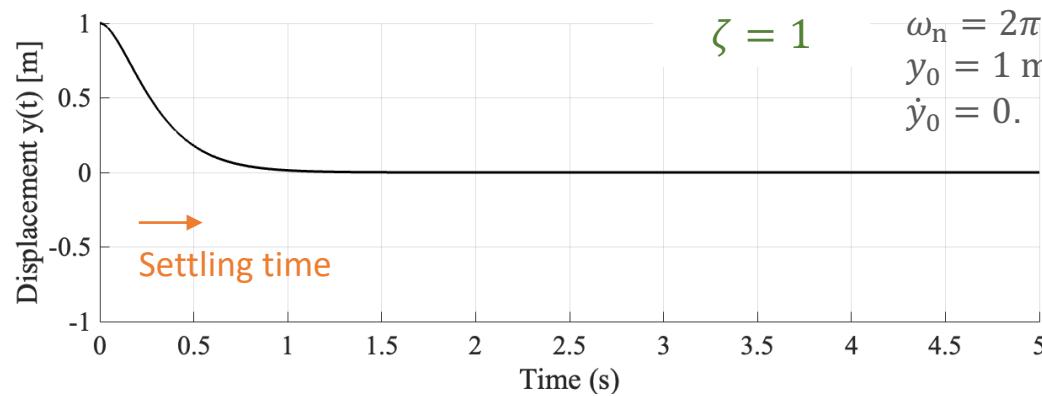
$$c = c_{\text{crit}} = 2\sqrt{m k}$$

- Solution for two real equal roots:

$$y(t) = C'_1 e^{-\omega_n t} + C'_2 t e^{-\omega_n t}$$

$$\begin{aligned} y_0 &= y(t_0); \\ \dot{y}_0 &= \dot{y}(t_0) \end{aligned} \rightarrow C'_1 = y_0; \quad C'_2 = \dot{y}_0 + \omega_n y_0$$

- Exponential decay:



- Minimum damping to avoid oscillations.
- Mass returns to rest position in shortest time (responsiveness) without oscillations.

#### Applications:

- automatic / hydraulic doors,
- some automobile suspensions,
- safety systems in brakes or shock absorbers.

- Overdamped case:  $\zeta > 1$

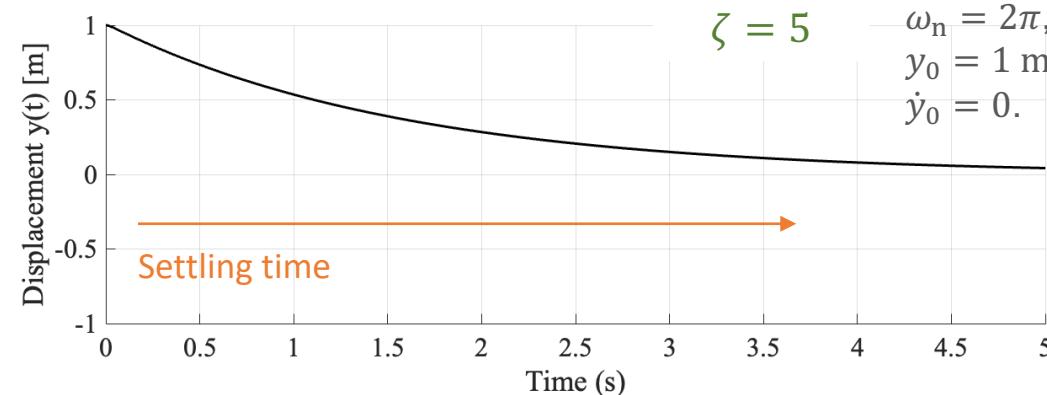
$$c > c_{\text{crit}} = 2\sqrt{m k}$$

- Solution for two real distinct roots:

$$y(t) = C_1 e^{[-\zeta \omega_n + \omega_n \sqrt{\zeta^2 - 1}] t} + C_2 e^{[-\zeta \omega_n - \omega_n \sqrt{\zeta^2 - 1}] t}$$

$$\begin{aligned} y_0 &= y(t_0); & C_1 &= \frac{y_0 \omega_n [\zeta + \sqrt{\zeta^2 - 1}] + \dot{y}_0}{2 \omega_n \sqrt{\zeta^2 - 1}}; \\ \dot{y}_0 &= \dot{y}(t_0) & C_2 &= \frac{-y_0 \omega_n [\zeta - \sqrt{\zeta^2 - 1}] - \dot{y}_0}{2 \omega_n \sqrt{\zeta^2 - 1}}. \end{aligned}$$

- Exponential decay:



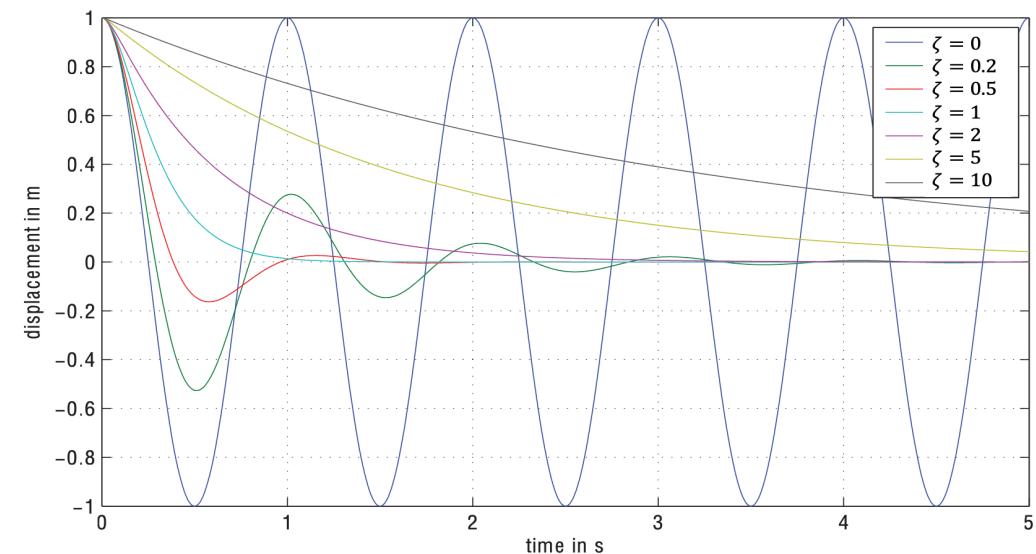
- Equilibrium reached without oscillations, slower than in critically damped case.

#### Applications:

- systems where oscillations are unacceptable,
- seismic dampers,
- aircraft landing gear.

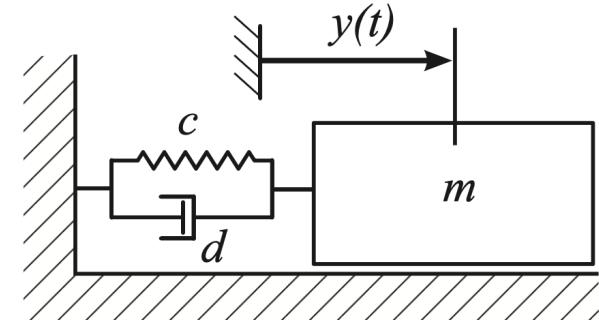
damping:	Solution free vibration:	oscillation:	settling time:	applications
no damping: $\zeta = 0$	$y(t) = C_1 \cos(\omega_n t) + C_2 \sin(\omega_n t)$	Yes	No	Baseline
underdamped: $0 < \zeta < 1$	$y(t) = e^{-\zeta \omega_n t} [C'_1 \cos(\omega_d t) + C'_2 \sin(\omega_d t)]$	Yes	Medium	Suspensions, machinery, musical instruments
critically damped: $\zeta = 1$	$y(t) = C'_1 e^{-\omega_n t} + C'_2 t e^{-\omega_n t}$	No	Minimum possible	Suspensions, measuring instruments, hydraulic doors
overdamped: $\zeta > 1$	$y(t) = C_1 e^{-\zeta \omega_n + \omega_n \sqrt{\zeta^2 - 1} t} + C_2 e^{-\zeta \omega_n - \omega_n \sqrt{\zeta^2 - 1} t}$	No	Slower than critical	Seismic dampers, aircraft landings

- Damping → reduces amplitude, lowers frequency with respect to undamped free vibration.



Statement: Derive the differential equation for the depicted spring-damper-mass system, determine the homogeneous solution  $y_h(t)$  and adapt it to the given initial conditions  $y_0$  and  $\dot{y}_0$ . Consider the four different values for the damping  $c$ .

given:  $c = 0 \text{ N s mm}^{-1} | 0.1 \text{ N s mm}^{-1} | 0.2 \text{ N s mm}^{-1} | 0.3 \text{ N s mm}^{-1}$   
 $k = 10 \text{ N cm}^{-1}; m = 10 \text{ kg}; t_0 = 0 \text{ s}; y_0 = 1 \text{ mm}; \dot{y}_0 = -3 \text{ mm s}^{-1}$



- Book: SS Rao. *Mechanical vibrations*. ISBN 978-0-13-212819-3.
- Moreno-Mateos, M.A. (2025). CD-computdynamics. GitHub repository: <https://github.com/yourusername/comp-dyn-labs>
- Moreno-Mateos, M.A. (2025). MorenoMiguelES/CD-computdynamics: Elastodynamics Computer Laboratory, Moreno-Mateos (v1.1). Zenodo. (<https://doi.org/10.5281/zenodo.15771587>).