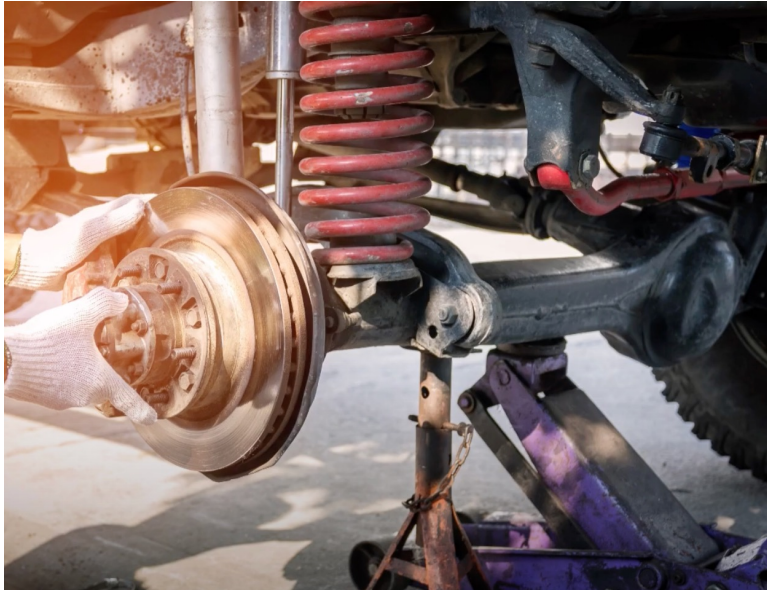




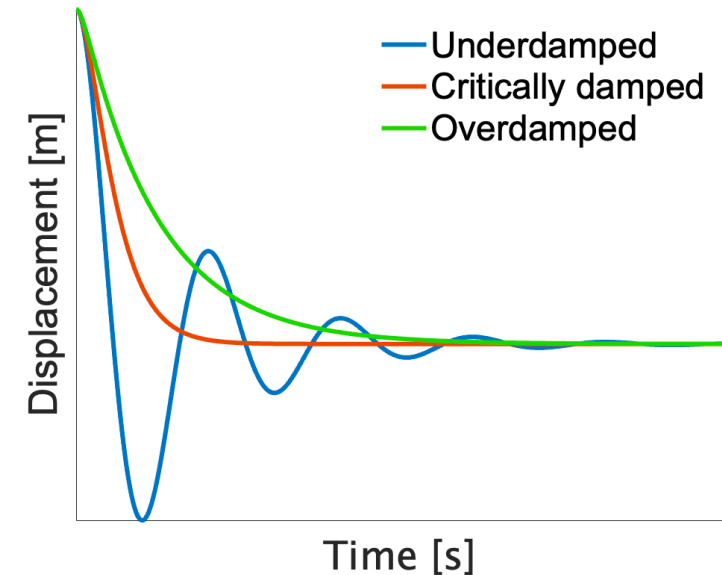
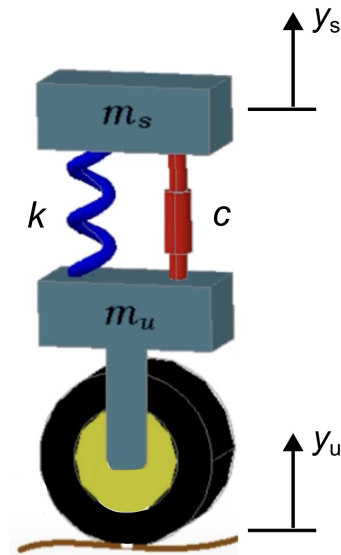
Analysis of Single Degree of Freedom Mass–Spring–Damper System: Free Vibration

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Automobile Suspension System

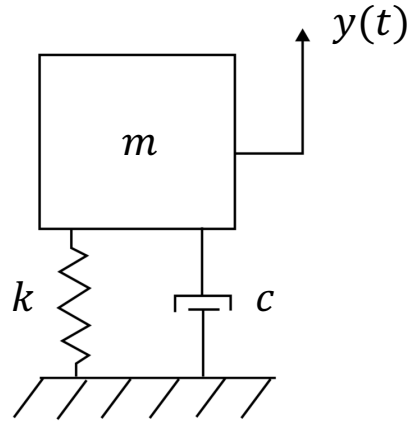


Mass + Spring + Damper



- ✓ **Critically damped:** Minimise oscillations & guarantee responsiveness (quick return to rest position).
- ✗ **Underdamped:** undesired oscillations.
- ✗ **Overdamped:** unresponsive suspension system.

Single Degree of Freedom Spring-Mass-Damper System:



with:

- coordinate $y(t)$ with time derivatives $\dot{y}(t)$ and $\ddot{y}(t)$,

- inertia , mass m [$\text{N m}^{-1}\text{s}^2$],

- linear spring , coefficient k [N m^{-1}],

- viscous damper , coefficient c [$\text{N m}^{-1}\text{s}$],

Hypotheses:

- single coordinate describes motion;
- lumped mass system;
- linear elasticity for spring;
- linear viscous damping;
- time invariant parameters;
- no additional forces (e.g., friction).

Find:

Free motion of the mass with initial conditions y_0 and $\dot{y}_0 \rightarrow$ homogeneous solution.

Homogeneous solution:

right side $f(t) = 0 \Rightarrow y_h(t)$

Particular solution:

right side $f(t) \neq 0 \Rightarrow y_p(t)$

Solution of ODE:

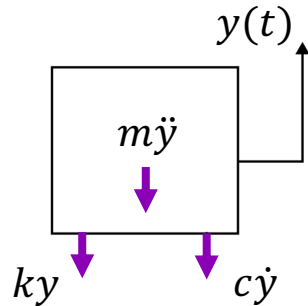
Total solution:

$$y_{\text{tot}}(t) = y_h(t) + y_p(t)$$

a) Equation of motion: D'Alembert's principle

- Newton's second law,

$$m \ddot{y}(t) + c \dot{y}(t) + k y(t) = 0$$



$m \ddot{y}(t)$: inertia force (opposes acceleration)

$c \dot{y}(t)$: damping force (dissipates energy)

$k y(t)$: spring force (restores position)

b) Initial conditions:

$$y_0 = y(t_0); \dot{y}_0 = \dot{y}(t_0)$$

Normalised form:

$$\ddot{y}(t) + 2\zeta\omega_n \dot{y}(t) + \omega_n^2 y(t) = 0$$

- undamped natural frequency,

$$\omega_n = \sqrt{k/m} \quad [\text{rad s}^{-1}]$$

- damping ratio,

$$\zeta = \frac{c}{c_{\text{crit}}} \quad [-]$$

$$c_{\text{crit}} = 2\sqrt{mk} \quad [\text{N m}^{-1} \text{ s}]$$

1. Assume solution, substitute: $y(t) = Ce^{\lambda t}; \quad \dot{y}(t) = C\lambda e^{\lambda t}; \quad \ddot{y}(t) = C\lambda^2 e^{\lambda t}$

2. Characteristic equation: $\lambda^2 + 2\zeta\omega_n \lambda + \omega_n^2 = 0 \rightarrow \lambda_{1,2} = -\zeta\omega_n \pm \omega_n\sqrt{\zeta^2 - 1}.$

3. Two roots ($\lambda_{1,2}$), two solutions: $y_1(t) = C_1 e^{\lambda_1 t}; \quad y_2(t) = C_2 e^{\lambda_2 t}$

4. General solution:

$$y(t) = C_1 e^{[-\zeta\omega_n + \omega_n\sqrt{\zeta^2 - 1}]t} + C_2 e^{[-\zeta\omega_n - \omega_n\sqrt{\zeta^2 - 1}]t}$$

$C_1; C_2$

constants from initial conditions.

- No damping: $\zeta = 0$

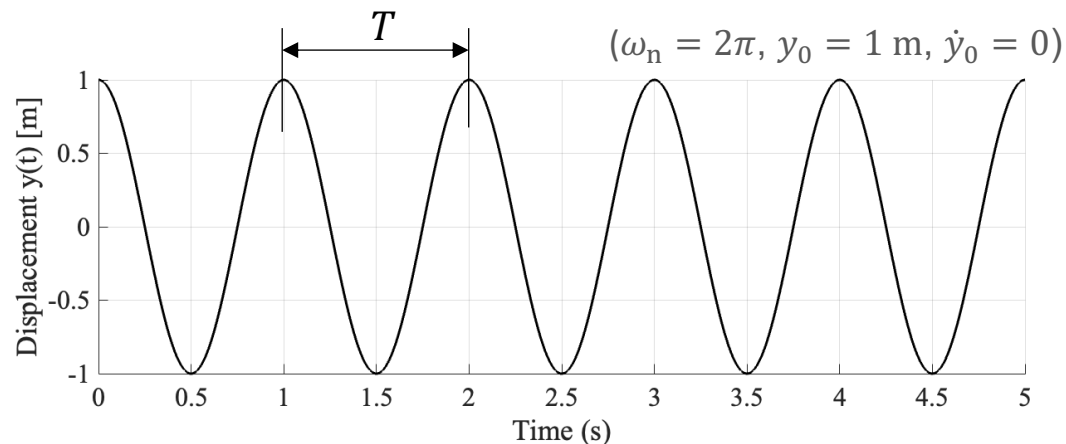
- Solution for complex conjugate roots without real part:

$$y(t) = C_1 e^{[-\zeta\omega_n + i\omega_n\sqrt{1-\zeta^2}]t} + C_2 e^{[-\zeta\omega_n - i\omega_n\sqrt{1-\zeta^2}]t}$$

- Express as harmonic oscillation (Euler's equation*):

$$y_h(t) = C'_1 \cos(\omega_n t) + C'_2 \sin(\omega_n t)$$

$$\begin{aligned} y_0 &= y(t_0); & \dot{y}_0 &= \dot{y}(t_0) \end{aligned} \rightarrow C'_1 = y_0; C'_2 = \frac{\dot{y}_0}{\omega_n}$$



Note(*): $e^{i\alpha} = \cos \alpha + i \sin \alpha$

- Natural frequency without damping:

$$\omega_n = \sqrt{k/m} \quad T = \frac{2\pi}{\omega_n}$$

$$\uparrow k \Rightarrow \uparrow \omega_n$$

$$\uparrow m \Rightarrow \downarrow \omega_n$$

Application:

- baseline for more complex systems (with damping and forcing).

- Underdamped case: $0 < \zeta < 1$ $c < c_{\text{crit}} = 2\sqrt{m k}$

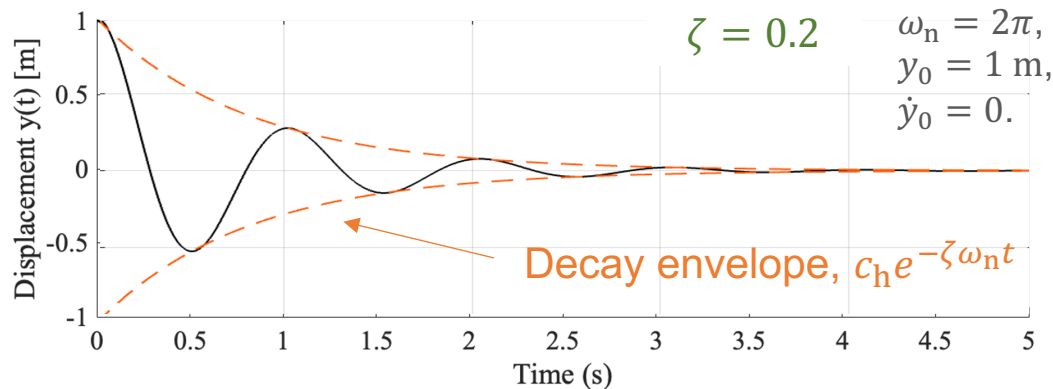
- Solution for complex conjugate roots, negative real part:

$$y(t) = C_1 e^{[-\zeta\omega_n + i\omega_n\sqrt{1-\zeta^2}]t} + C_2 e^{[-\zeta\omega_n - i\omega_n\sqrt{1-\zeta^2}]t}$$

- Express as harmonic oscillation (Euler's equation):

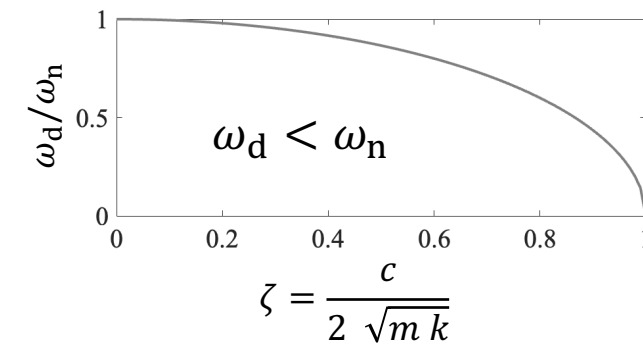
$$y(t) = e^{-\zeta\omega_n t} [C'_1 \cos(\omega_d t) + C'_2 \sin(\omega_d t)]$$

$$\begin{aligned} y_0 &= y(t_0); & \dot{y}_0 &= \dot{y}(t_0) \end{aligned} \rightarrow C'_1 = y_0; \quad C'_2 = \frac{\dot{y}_0 + \zeta\omega_n y_0}{\omega_d}$$



- Frequency of damped vibration:

$$\omega_d = \omega_n \sqrt{1 - \zeta^2}$$



- Damping \rightarrow lowers frequency, reduces amplitude.

Applications:

- most automobile suspension,
- vibrating mechanical systems,
- musical instruments.

- Critically damped case: $\zeta = 1$

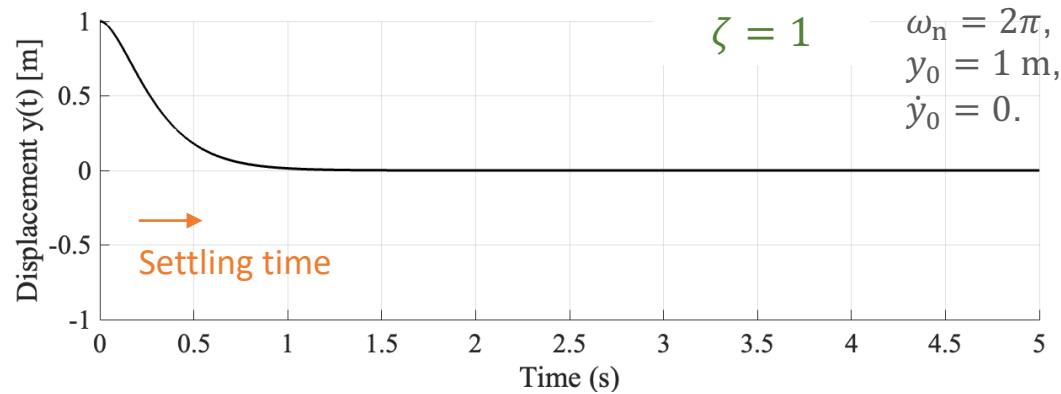
$$c = c_{\text{crit}} = 2\sqrt{m k}$$

- Solution for two real equal roots:

$$y(t) = C_1' e^{-\omega_n t} + C_2' t e^{-\omega_n t}$$

$$\begin{aligned} y_0 &= y(t_0); \\ \dot{y}_0 &= \dot{y}(t_0) \end{aligned} \rightarrow C_1' = y_0; C_2' = \dot{y}_0 + \omega_n y_0$$

- Exponential decay:



- Minimum damping to avoid oscillations.
- Mass returns to rest position in shortest time (responsiveness) without oscillations.

Applications:

- automatic / hydraulic doors,
- some automobile suspensions,
- safety systems in brakes or shock absorbers.

- Overdamped case: $\zeta > 1$

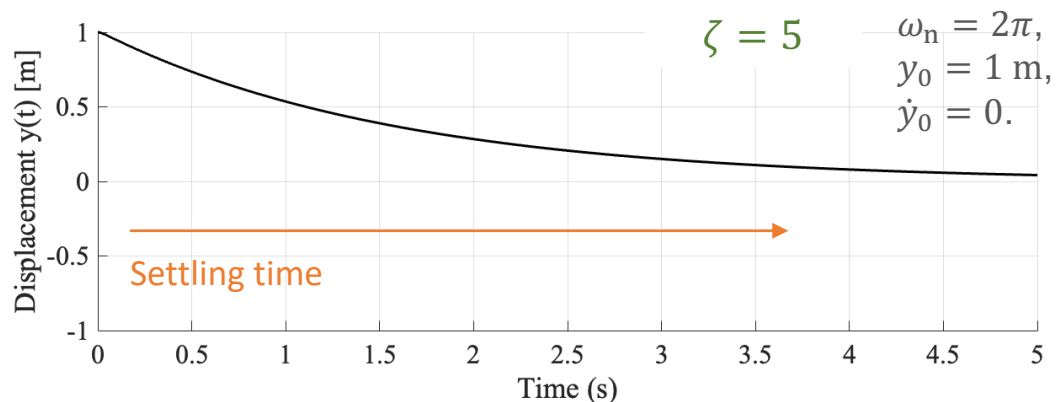
$$c > c_{\text{crit}} = 2\sqrt{m k}$$

- Solution for two real distinct roots:

$$y(t) = C_1 e^{[-\zeta\omega_n + \omega_n\sqrt{\zeta^2 - 1}]t} + C_2 e^{[-\zeta\omega_n - \omega_n\sqrt{\zeta^2 - 1}]t}$$

$$\begin{aligned} y_0 = y(t_0); \quad \dot{y}_0 = \dot{y}(t_0) &\rightarrow C_1 = \frac{y_0\omega_n[\zeta + \sqrt{\zeta^2 - 1}] + \dot{y}_0}{2\omega_n\sqrt{\zeta^2 - 1}}; \\ C_2 &= \frac{-y_0\omega_n[\zeta - \sqrt{\zeta^2 - 1}] - \dot{y}_0}{2\omega_n\sqrt{\zeta^2 - 1}}. \end{aligned}$$

- Exponential decay:



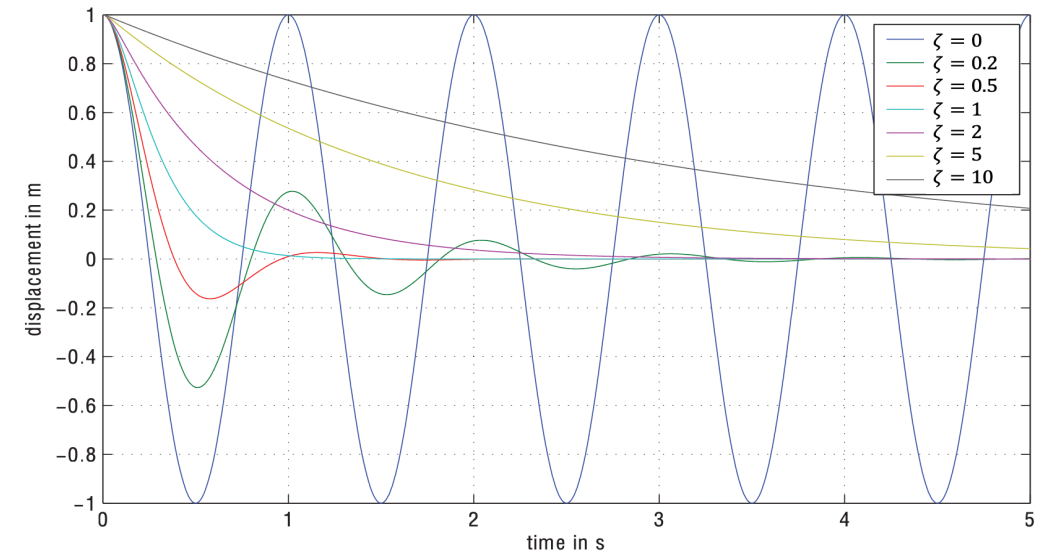
- Equilibrium reached without oscillations, slower than in critically damped case.

Applications:

- systems where oscillations are unacceptable,
- seismic dampers,
- aircraft landing gear.

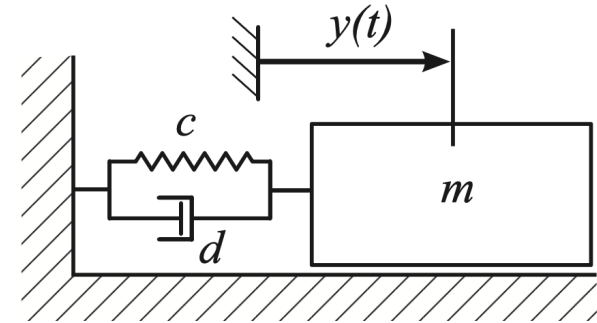
damping:	Solution free vibration:	oscillation:	settling time:	applications
no damping: $\zeta = 0$	$y(t) = C_1 \cos(\omega_n t) + C_2 \sin(\omega_n t)$	Yes	No	Baseline
underdamped: $0 < \zeta < 1$	$y(t) = e^{-\zeta \omega_n t} [C'_1 \cos(\omega_d t) + C'_2 \sin(\omega_d t)]$	Yes	Medium	Suspensions, machinery, musical instruments
critically damped: $\zeta = 1$	$y(t) = C'_1 e^{-\omega_n t} + C'_2 t e^{-\omega_n t}$	No	Minimum possible	Suspensions, measuring instruments, hydraulic doors
overdamped: $\zeta > 1$	$y(t) = C_1 e^{-\zeta \omega_n + \omega_n \sqrt{\zeta^2 - 1} t} + C_2 e^{-\zeta \omega_n - \omega_n \sqrt{\zeta^2 - 1} t}$	No	Slower than critical	Seismic dampers, aircraft landings

- Damping \rightarrow reduces amplitude, lowers frequency with respect to undamped free vibration.



Statement: Derive the differential equation for the depicted spring-damper-mass system, determine the homogeneous solution $y_h(t)$ and adapt it to the given initial conditions y_0 and \dot{y}_0 . Consider the four different values for the damping c .

given: $c = 0 \text{ N s mm}^{-1} \mid 0.1 \text{ N s mm}^{-1} \mid 0.2 \text{ N s mm}^{-1} \mid 0.3 \text{ N s mm}^{-1}$
 $k = 10 \text{ N cm}^{-1}$; $m = 10 \text{ kg}$; $t_0 = 0 \text{ s}$; $y_0 = 1 \text{ mm}$; $\dot{y}_0 = -3 \text{ mm s}^{-1}$



- Book: SS Rao. *Mechanical vibrations*. ISBN 978-0-13-212819-3.
- Moreno-Mateos, M.A. (2025). CD-computodynamics. GitHub repository: <https://github.com/yourusername/comp-dyn-labs>
- Moreno-Mateos, M.A. (2025). MorenoMiguelES/CD-computodynamics: Elastodynamics Computer Laboratory, Moreno-Mateos (v1.1). Zenodo. (<https://doi.org/10.5281/zenodo.15771587>).