



Quantum Espresso Overview and Practical Guide

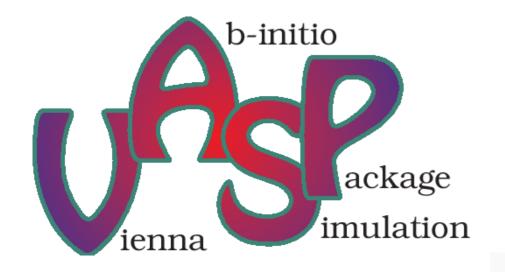
Dr. Dzade Nelson
The Pennsylvania State University

DFT codes/software



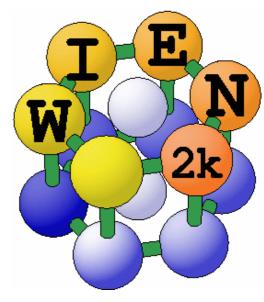








Gaussian 16







Spanish Initiative for Electronic Simulations with Thousands of Atoms

CAmbridge **S**erial **T**otal **E**nergy **P**ackage

Kohn-Sham equations

IN RUN OUT

Model:

unit cell lattice vectors basis

Physical approx:

xc-approximation GGA, LDA, ...

Numerical approx:

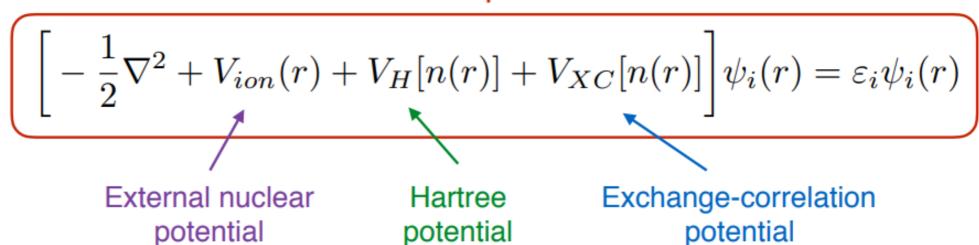
energy cut-off k-points grid SCF procedure



Physical quantities:

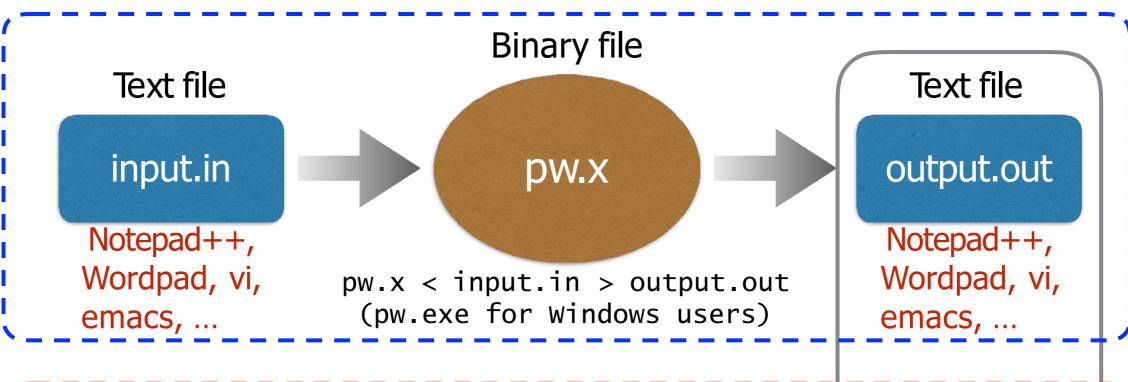
charge density total energy KS wavefunctions KS energies







Structure of QE I/O file



Text & data file

band structure, phonon, charge density

gnuplot, xmgrace, xcrysden, Vesta

Binary file

pw.x, ph.x, pp.x,

pw.x < input.in > output.out
 (pw.exe for Windows users)

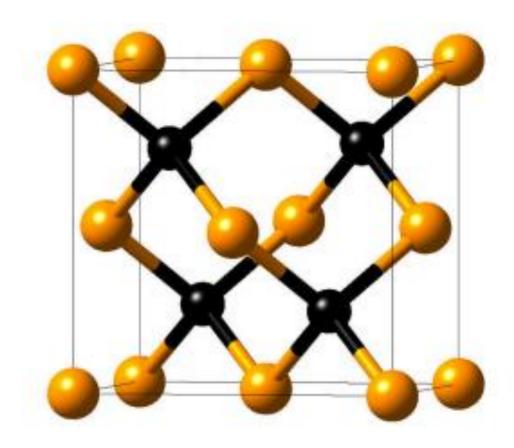
Data file

tmp folder

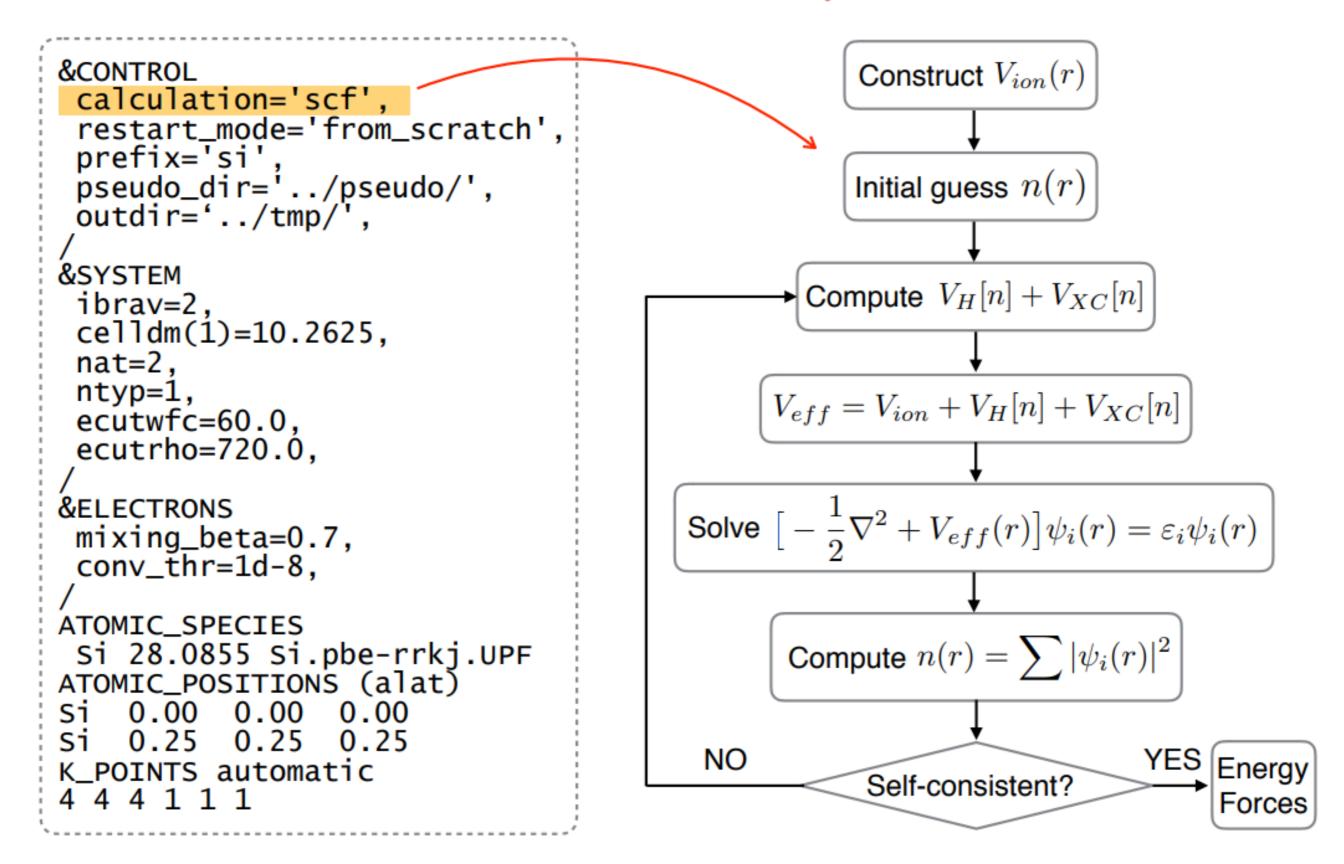
Text file

new input.in

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Si
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                0.25
K_POINTS automatic
```



Silicon



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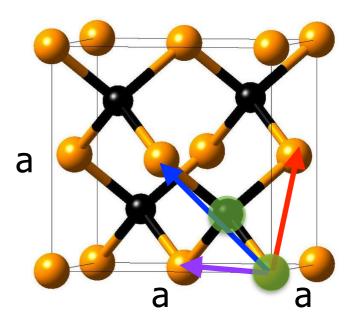
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ibrav = 2 (FCC)
ibrav = 4 (Hexagonal)
simple cubic:
v1 = a(1,0,0)
v2 = a(0,1,0)
v3 = a(0,0,1)
face centered cubic:
v1 = (a/2)(-1,0,1)
v2 = (a/2)(0,1,1)
v3 = (a/2)(-1,1,0)
```

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v2 = (a/2)(0,1,1)
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 Si 28.0855 Si.pbe-rrkj.UPF
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si 0.25 0.25 0.25
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K_POINTS automatic
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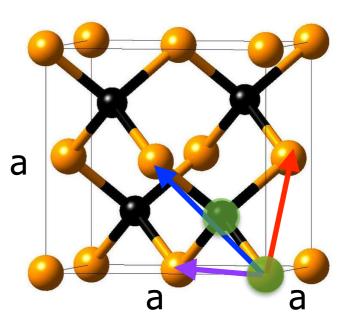
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K_POINTS automatic
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FCC structure



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FCC structure



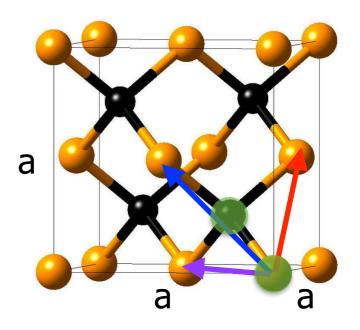
Silicon

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         0.25 0.25
K_POINTS automatic
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K_POINTS automatic
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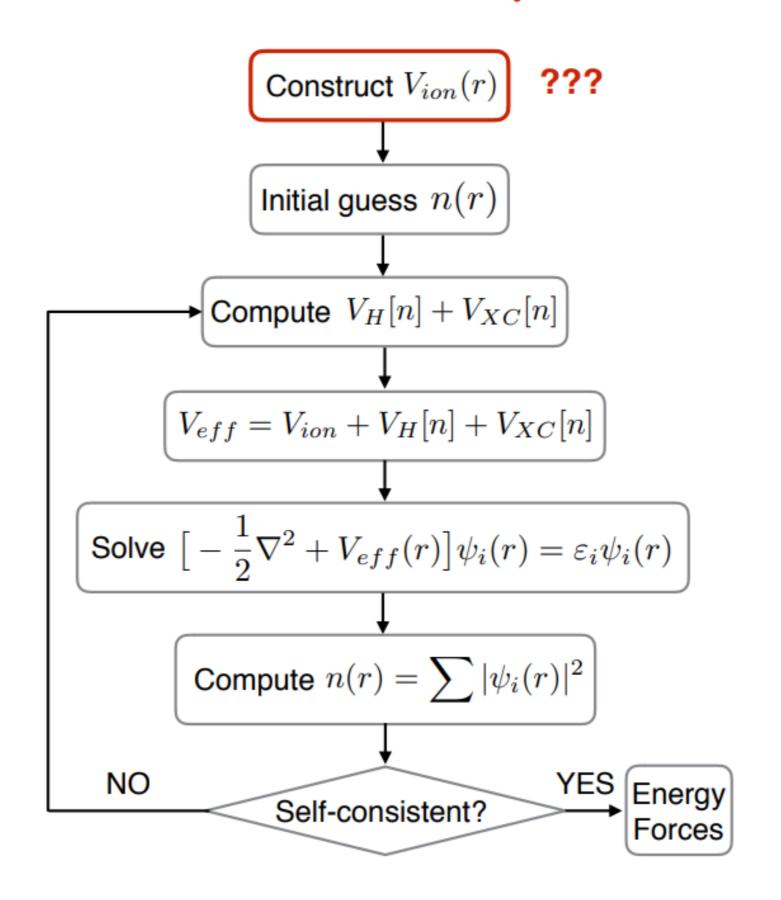
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ATOMIC_POSITIONS (alat) ←
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si 0.25 0.25 0.25
K_POINTS automatic
4 4 4 1 1 1
```

FCC structure



atomic positions are in cartesian coordinates, in units of the lattice parameter.

External nuclear potential



External nuclear potential

Electrons experience a Coulomb potential due to the nuclei with simple form:

$$V_{ion} = -\frac{Z}{r}$$

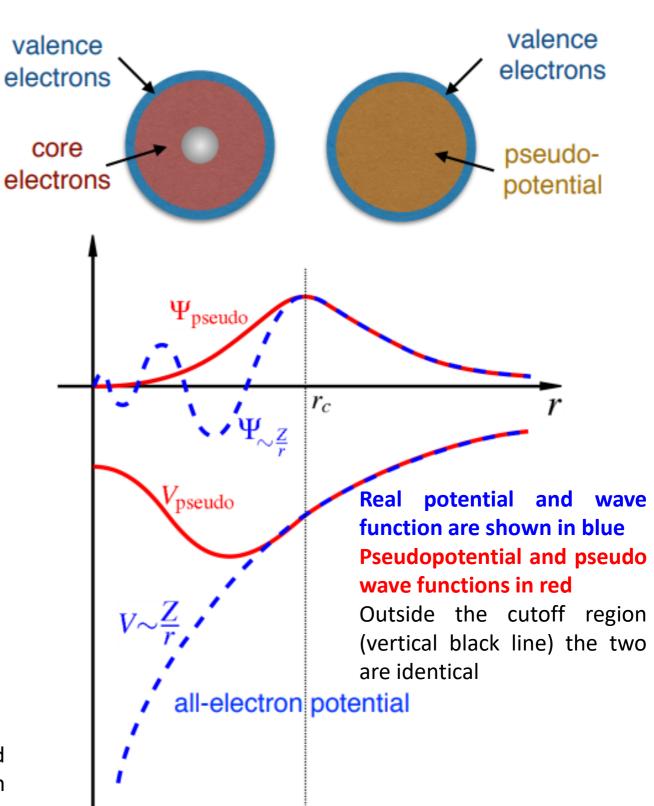
But this leads to computational problems!



Pseudopotential (PP):

Replace the strong Coulomb potential of the nucleus and tightly bound core electrons by an effective ionic potential acting on the valence electrons.

The rapidly changing potential field near the atomic core is replaced by a smoother function that simulates the potential field far from the core very well. By doing so, it requires less number plane wave basis for wavefunction expansion.



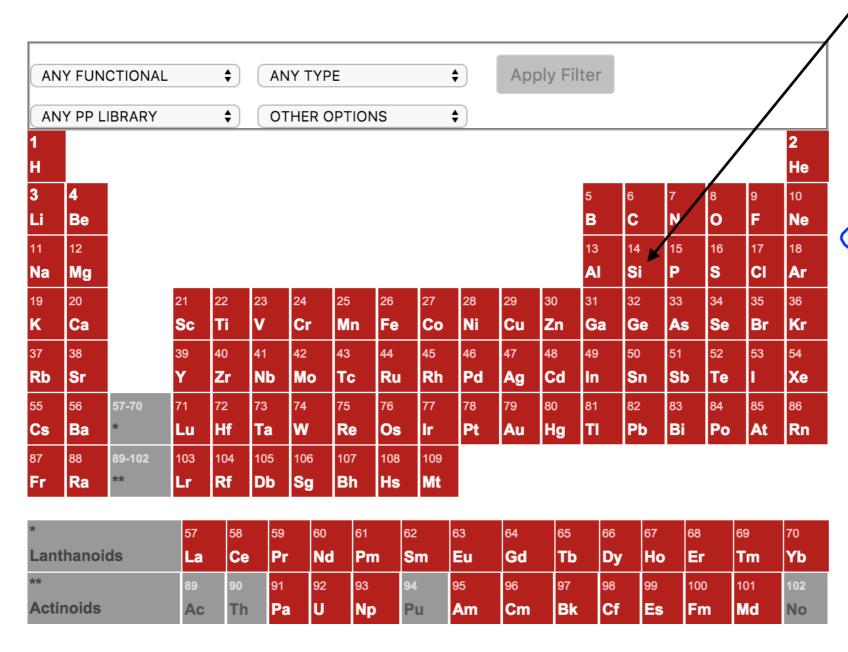
Pseudopotential

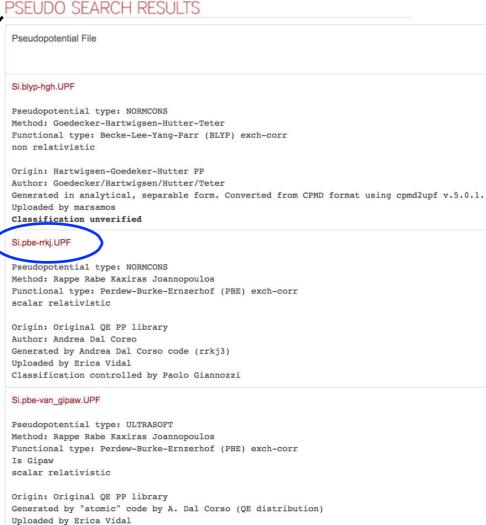
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    0.25 0.25
K_POINTS automatic
```

Name of the pseudopotential file

How to get pseudopotential

http://www.quantum-espresso.org/pseudopotentials/





Si.pbe-rrkj.UPF

Classification controlled by Paolo Giannozzi

- type of exchangecorrelation functional
- type of pseudopotential

Pseudopotential

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    0.25 0.25 0.25
K_POINTS automatic
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Si.pbe-rrkj.UPF

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  2.5000000000E+00
nl pn l occ
                                               Rcut US
                                                                     E pseu
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2 1 2.00
                     2.50000000000
                                         2.60000000000
                                                             0.0000000000
                     2.50000000000
                                         2.60000000000
                                                             0.0000000000
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                        PBE Exchange-Correlation functional
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    4.00000000000
                        z valence
 -7.47480832270 Total energy 0.0000000 0.0000000 Suggested cutoff for wfc and rho
                        Max angular momentum component
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                        Number of points in mesh
                        Number of Wavefunctions, Number of Projectors
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                               occ
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Plane wave expansion

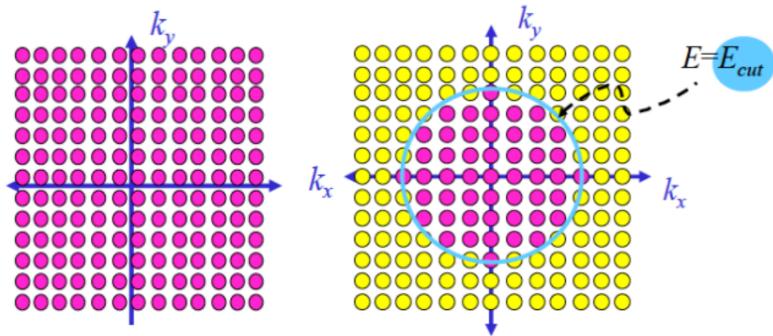
In a periodic system we can write the KS states as a superposition of plane waves:

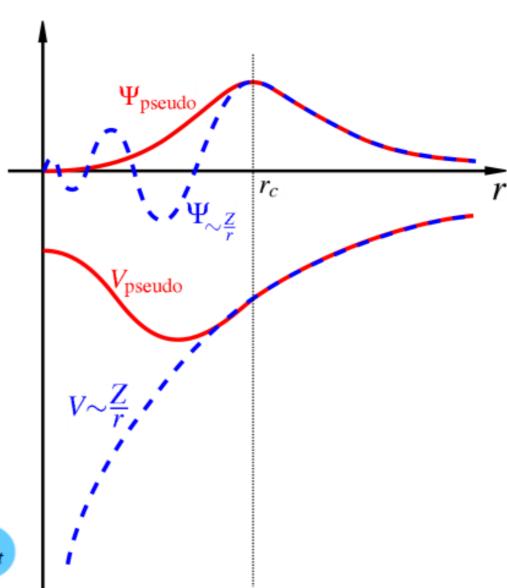
$$\Psi_{k,n}(r) = \frac{1}{\Omega} \sum_{G} c_{k,n}^{G} e^{i(k+G)\cdot r}$$

G are vectors in reciprocal space.

The sum, in principle infinite, can be truncated:

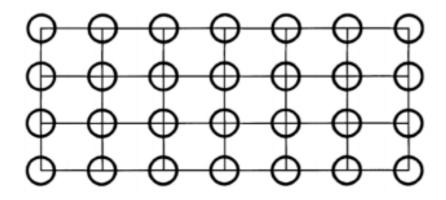
$$\frac{\hbar^2}{2m} |k+G|^2 \le E_{cut}$$





Crystalline solids and plane wave DFT

A crystal: periodic arrangement atoms



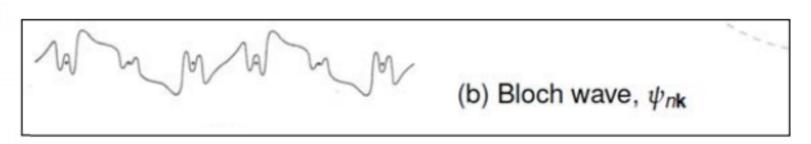
(a) Periodic potential, $U(\mathbf{r})$

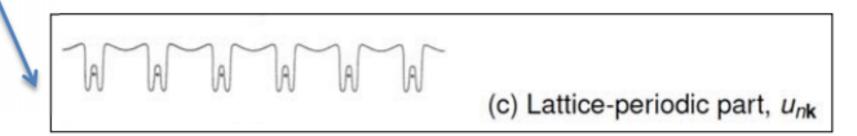
Free electrons: Plane waves:

(d) Plane wave, e^{ik·r}

Electrons in a periodic potential are Bloch waves:

 $\psi_{n\mathbf{k}}(\mathbf{r}) = \exp(i\mathbf{k} \cdot \mathbf{r}) \, \mathbf{u}_{n\mathbf{k}}(\mathbf{r})$ (Perturbed free electrons)



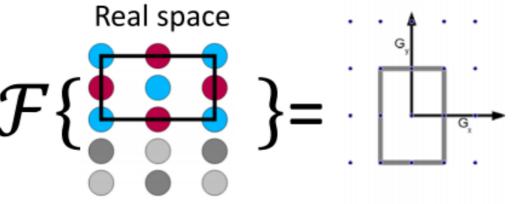


Crystalline solids and plane wave DFT

-Cutoff energy

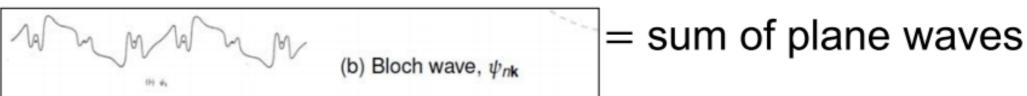
Reciprocal space:

Reciprocal space



All periodic functions can be expanded by a Fourier series:

$$\psi_{n\mathbf{k}}(\mathbf{r}) = \exp(i\mathbf{k} \cdot \mathbf{r}) \, \mathbf{u}_{n\mathbf{k}}(\mathbf{r}) = \exp(i\mathbf{k} \cdot \mathbf{r}) \sum_{\mathbf{G}} c_{\mathbf{k}} \exp(i\mathbf{G} \cdot \mathbf{r})$$



- Each plane wave in the sum have kinetic energy: $E = \frac{\hbar}{2m} |\mathbf{k} + \mathbf{G}|^2$
- Numerically we must define a cutoff energy for the expansion!
- The choice of cutoff energy must be tested with respect to energy convergence in calculations

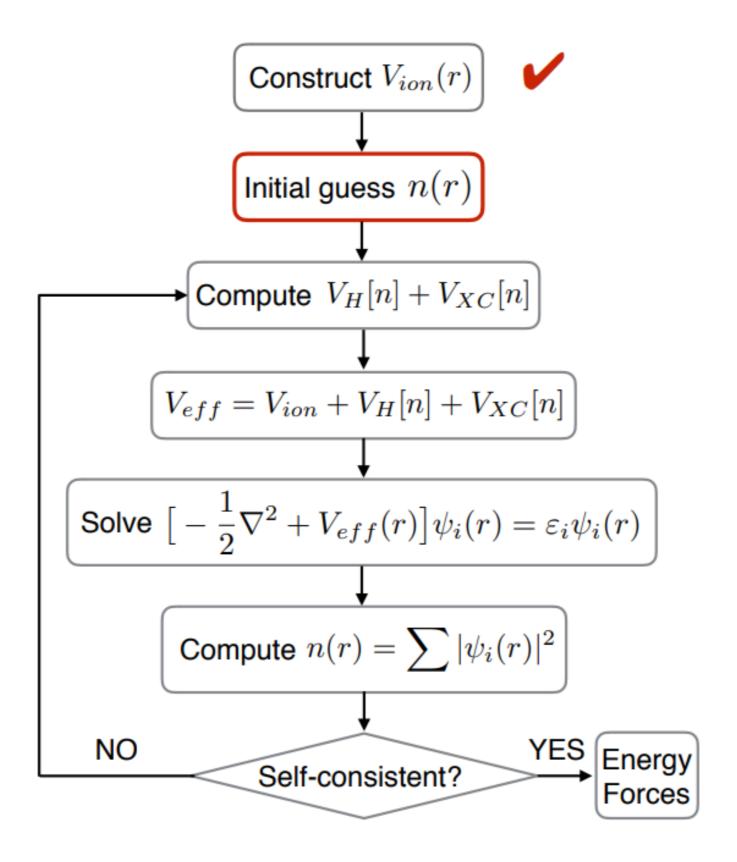
Plane wave expansion

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ATOMIC_POSITIONS (alat)
si 0.00 0.00 0.00
   0.25 0.25 0.25
K_POINTS automatic
4 4 4 1 1 1
```

$$\frac{\hbar^2}{2m}|k+G|^2 \le E_{cut}$$

Units: Ry (1 Ry = 0.5 Ha = 13.6057 eV) For ultrasoft pseudopotentials we have also: ecutrho = usually $8-12 \times \text{ecutwfc}$

Initial n(r)



Initial n(r)

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ATOMIC_POSITIONS (alat)
Si 0.00 0.00 0.00
          0.25
K_POINTS automatic
4 4 4 1 1 1
```

Sampling of the Brillouin zone

Many quantities we need to compute involve an integral over the BZ:

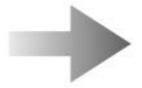
$$\bar{A} = \frac{1}{\Omega_{BZ}} \int_{BZ} A(k) d(k)$$

An example is the electronic density n(r):

$$n(r) = \frac{1}{\Omega_{BZ}} \sum_{i} \int_{BZ} |\Psi_{i,k}(r)|^2 f(\epsilon_{i,k} - \epsilon_F) d(k)$$

In practice the integral is discretized:

$$\frac{1}{\Omega_{BZ}} \int_{BZ} d(k) \to \sum_{k} \omega_{k}$$

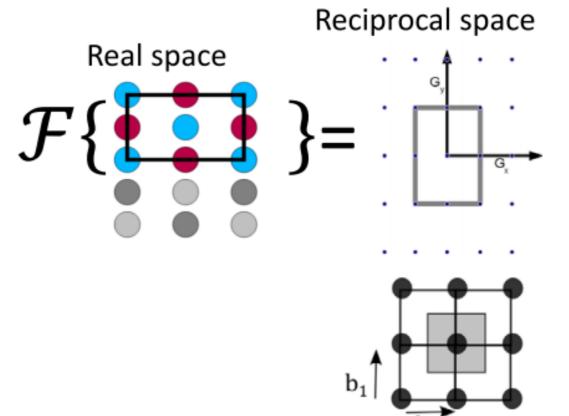


How do we choose the k points to include in the sum???

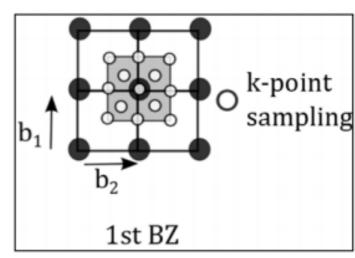
Crystalline solids and plane wave DFT - k-points

- Wavevectors $\mathbf{k} = \frac{2\pi}{\lambda} \left[\frac{1}{m} \right]$ are reciprocal space
- The primitive unit cell in reciprocal space is called the 1st Brillouin Zone (BZ)
- Any k-point that differ by a reciprocal lattice vector G are equivalent: k`=k+G
- Integrals need only be evaluated in the 1st BZ

- Numerically we must chose an appropriate number of k-points to sample the BZ
- The k-point sampling must be tested with respect to energy convergece



1st BZ



Monkhorst and Pack (1976)

Example: square 2D lattice

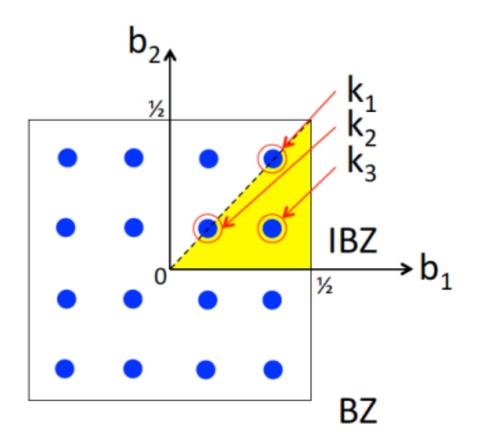
4×4 k-points grid (16 points)

3 inequivalent point (IBZ)

$$4 \times k_1 \to \omega_1 = 4/16 = 1/4$$

$$4 \times k_2 \to \omega_2 = 4/16 = 1/4$$

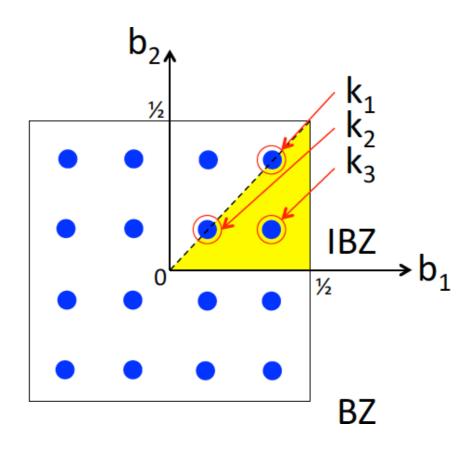
$$8 \times k_3 \to \omega_3 = 8/16 = 1/2$$



$$\frac{1}{\Omega_{BZ}} \int_{BZ} A(k)d(k) \simeq \frac{1}{4} A(k_1) + \frac{1}{4} A(k_2) + \frac{1}{2} A(k_3)$$

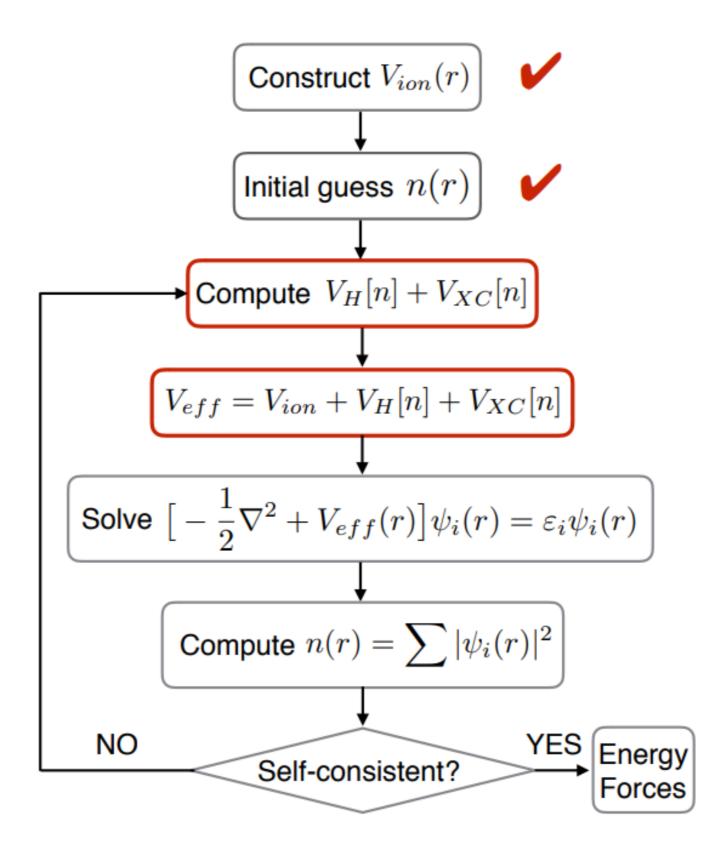
Sampling of the Brillouin zone

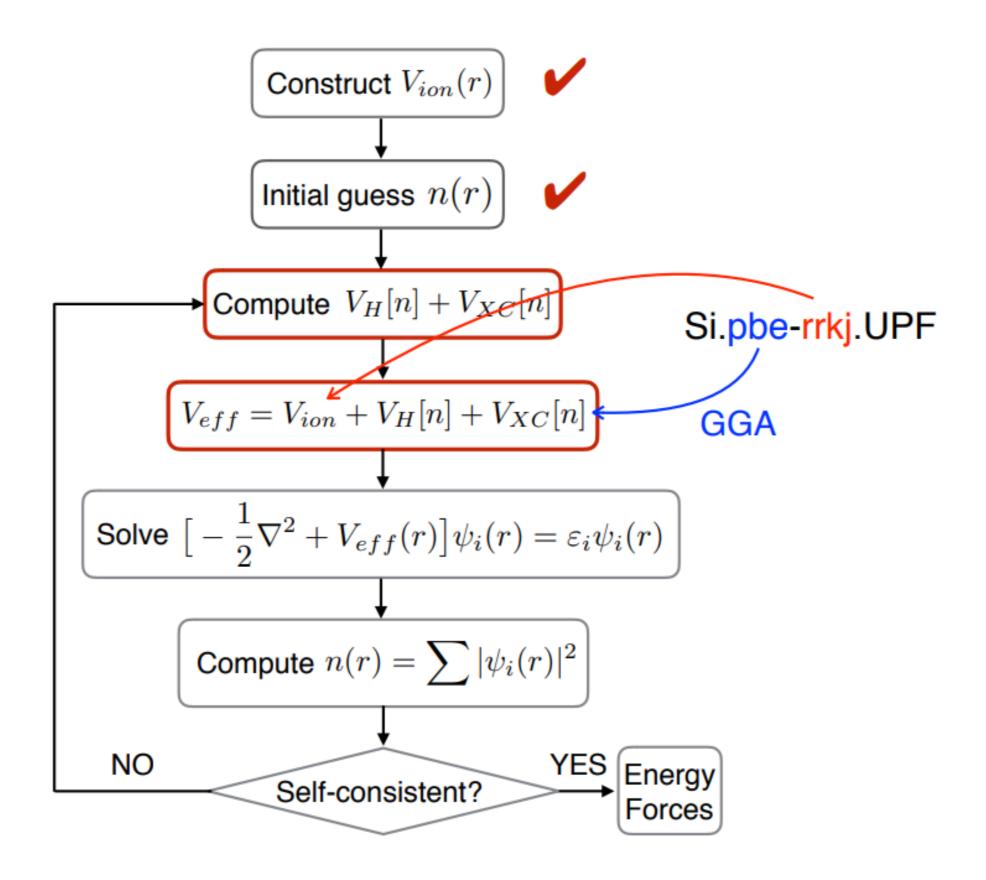
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ATOMIC_POSITIONS (alat)
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         0.25
K_POINTS automatic
4 4 4 1 1 1
```

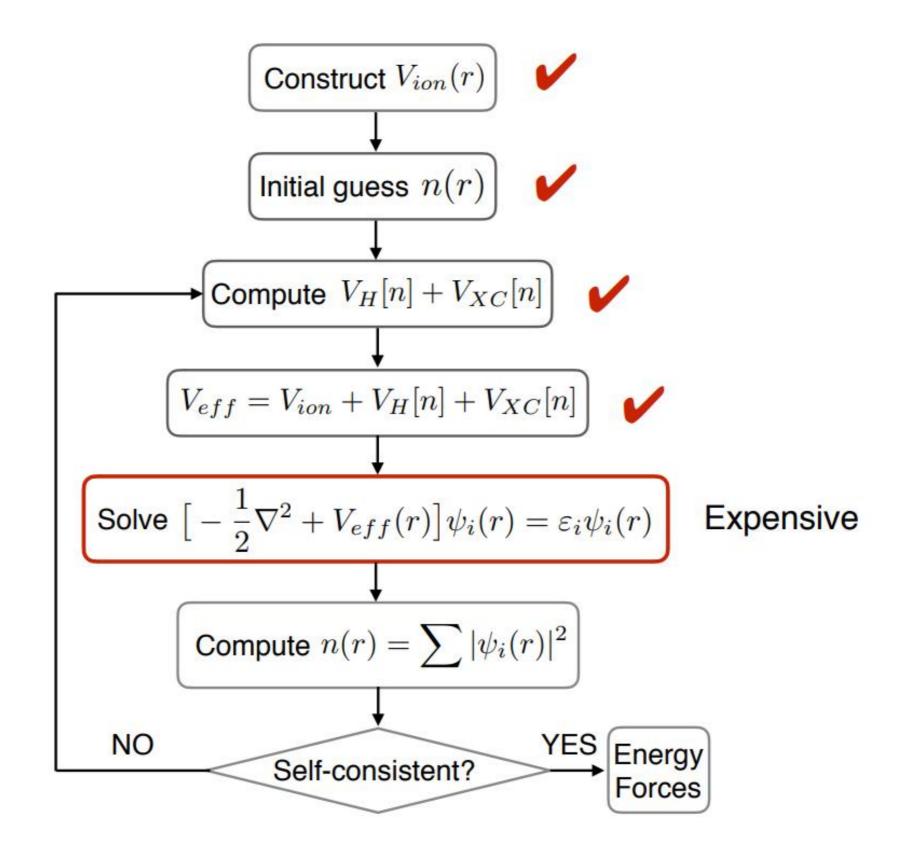


K_POINTS automatic
nk1, nk2, nk3, k1, k2, k3

nk1, nk2, nk3 as in Monkhorst-Pack grids k1, k2, k3 must be 0 (no offset) or 1 (grid displaced by half a grid step in the corresponding direction)







Solve wave equation

```
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ATOMIC_POSITIONS (alat)
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Si
si 0.25 0.25 0.25
K_POINTS automatic
4 4 4 1 1 1
```

Quantum Espresso Input File – diagonalization

```
&CONTROL
    calculation='scf'.
    restart_mode='from_scratch',
    prefix='si',
    pseudo_dir='/opt/pseudo/',
    outdir='tmp',
&SYSTEM
    ibrav=2,
    celldm(1)=10.2625,
    nat=2,
    ntyp=1,
    ecutwfc=60.0,
    ecutrho=720.0,
&ELECTRONS
    diagonalization = 'david'
    mixing_beta=0.7,
    conv_thr=1d-8,
ATOMIC SPECIES
    Si 28.0855 Si.pbe-rrkj.UPF
ATOMIC_POSITIONS (alat)
Si 0.00 0.00 0.00
Si 0.25 0.25 0.25
K POINTS automatic
444111
```

If the Kohn-Sham equations written in the matrix form, then we have

$$H[\rho]\psi_{bk} = \epsilon_{bk}\psi_{bk}$$

The corresponding energy levels of the non- (wavefunctions) for the interacting electrons.

Kohn-Sham orbitals non-interacting electrons.

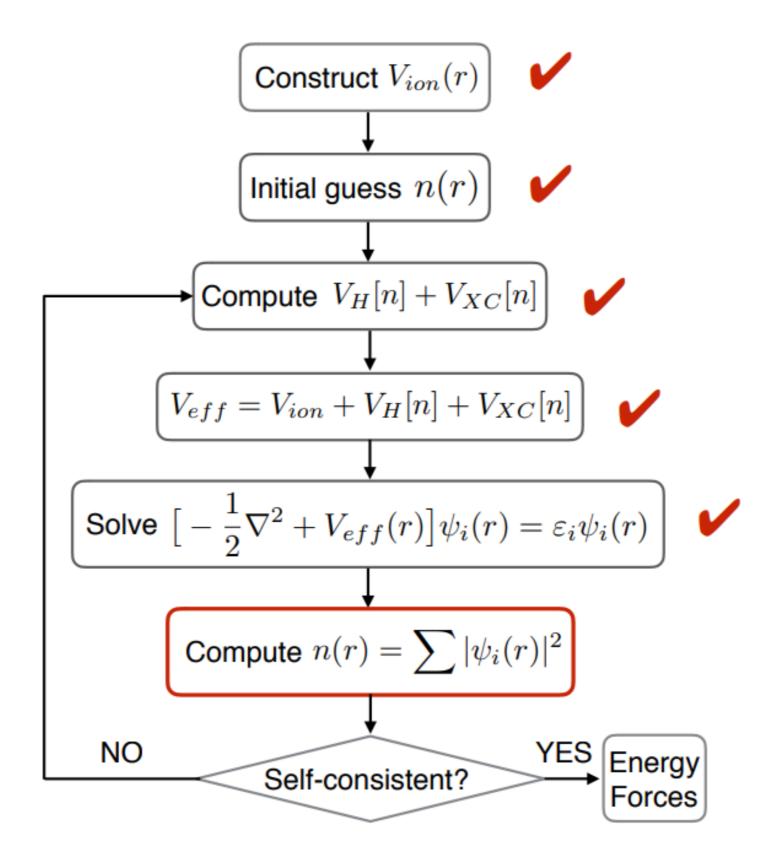
Eigenenergies can be found by the diagonalization of this matrix

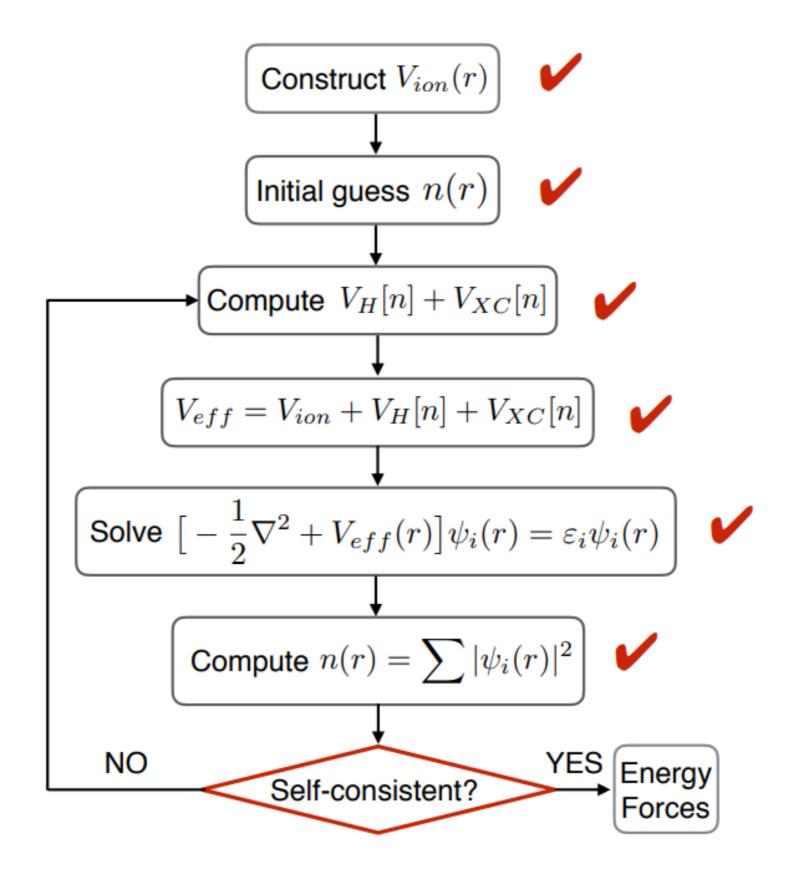
Davidson iterative diagonalization is the **DEFAULT** approach

Quantum Espresso Input File – diagonalization

```
&CONTROL
    calculation='scf',
    restart mode='from scratch',
    prefix='si',
    pseudo_dir='/opt/pseudo/',
    outdir='tmp',
&SYSTEM
    ibrav=2,
    celldm(1)=10.2625,
    nat=2,
    ntyp=1,
    ecutwfc=60.0,
    ecutrho=720.0,
&ELECTRONS
    diagonalization = 'david',
    mixing beta=0.7,
    conv thr=1d-8,
ATOMIC SPECIES
    Si 28.0855 Si.pbe-rrkj.UPF
ATOMIC_POSITIONS (alat)
Si 0.00 0.00 0.00
Si 0.25 0.25 0.25
K POINTS automatic
444111
```

```
Available options are:
  'david' :
       Davidson iterative diagonalization with overlap matrix
       (default). Fast, may in some rare cases fail.
  'cg' :
       Conjugate-gradient-like band-by-band diagonalization.
       MUCH slower than 'david' but uses less memory and is
       (a little bit) more robust.
  'ppcg':
       PPCG iterative diagonalization
  'paro', 'ParO' :
       ParO iterative diagonalization
  'rmm-davidson', 'rmm-paro':
       RMM-DIIS iterative diagonalization.
       To stabilize the SCF loop
       RMM-DIIS is alternated with calls to Davidson or
       ParO solvers depending on the string used.
       Other variables that can be used to tune the behavior of
       RMM-DIIS are: diago rmm ndim and diago rmm conv
```

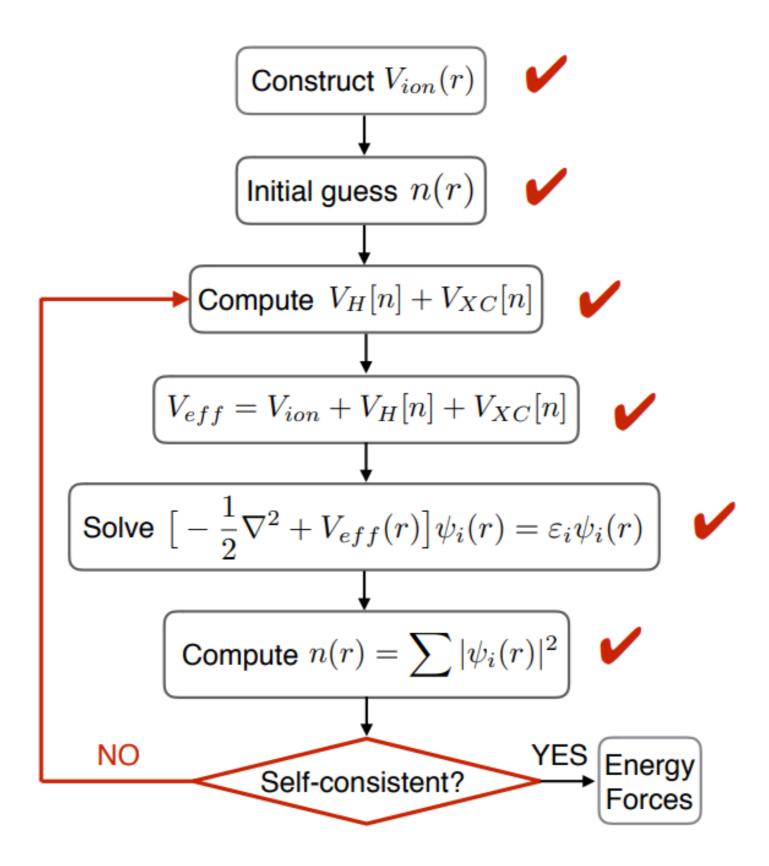




Self-consistency

```
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 prefix='si'
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 outdir='../tmp/',
&SYSTEM
 ibrav=2.
 celldm(1)=10.2625,
 nat=2.
 ntyp=1,
 ecutwfc=60.0,
 ecutrho=720.0,
&FL FCTRONS
 mixing_beta=0.7,
 conv_{thr}=1d-8,
ATOMIC_SPECIES
 Si 28.0855 Si.pbe-rrkj.UPF
ATOMIC_POSITIONS (alat)
    0.00 0.00
    0.25 0.25
K_POINTS automatic
4 4 4 1 1 1
```

Convergence threshold for self-consistency: estimated energy error > conv_thr (NO) or energy error < conv_thr (YES)



new n(r)

```
&CONTROL
 calculation='scf',
 restart_mode='from_scratch',
 prefix='si'
 pseudo_dir='../pseudo/',
 outdir='../tmp/',
&SYSTEM
 ibrav=2,
 celldm(1)=10.2625,
 nat=2.
 ntyp=1,
 ecutwfc=60.0,
 ecutrho=720.0,
&ELECTRONS
 mixing_beta=0.7, ← conv_thr=1d-8,
ATOMIC_SPECIES
 Si 28.0855 Si.pbe-rrkj.UPF
ATOMIC_POSITIONS (alat)
Si 0.00 0.00 0.00
         0.25
K_POINTS automatic
4 4 4 1 1 1
```

Mix new and old density: 0.7 = 70% of the new density and 30% of old density at first step

Quantum Espresso Input File – mixing_mode

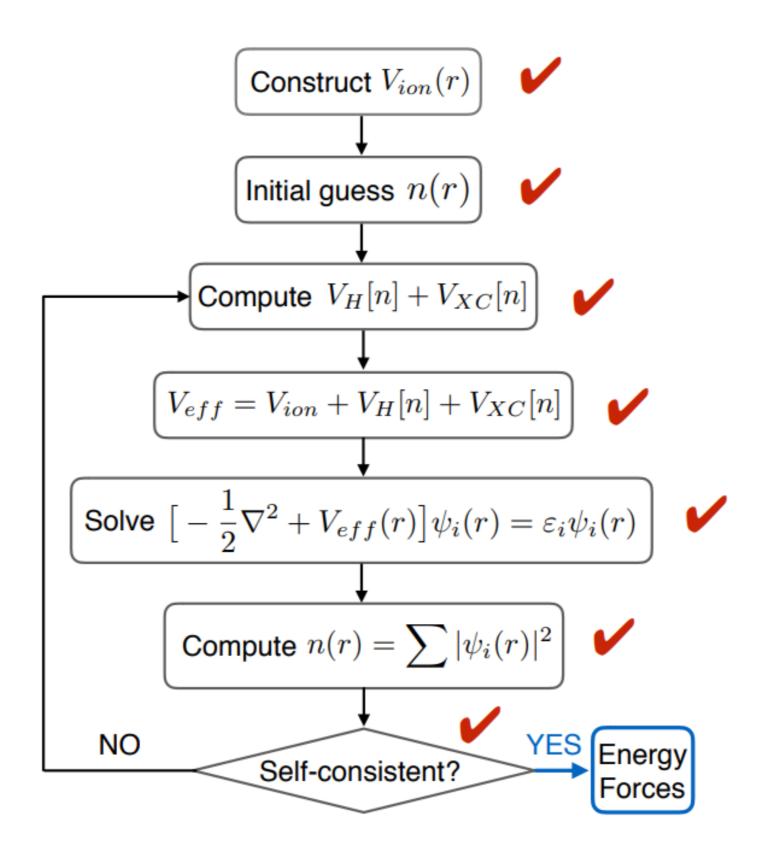
```
&CONTROL
    calculation='scf',
    restart mode='from scratch',
    prefix='si',
    pseudo_dir='/opt/pseudo/',
    outdir='tmp',
&SYSTEM
    ibrav=2,
    celldm(1)=10.2625,
    nat=2.
    ntyp=1,
    ecutwfc=60.0,
    ecutrho=720.0,
&ELECTRONS
    mixing_mode='plain',"
    mixing_beta=0.7,
    conv thr=1d-8,
ATOMIC SPECIES
    Si 28.0855 Si.pbe-rrkj.UPF
ATOMIC POSITIONS (alat)
Si 0.00 0.00 0.00
Si 0.25 0.25 0.25
K POINTS automatic
444111
```

- Allows to specify the mode used for mixing electronic charge density in SCF calculations
- 'plain': (DEFAULT)
 charge density Broyden mixing
- 'TF':

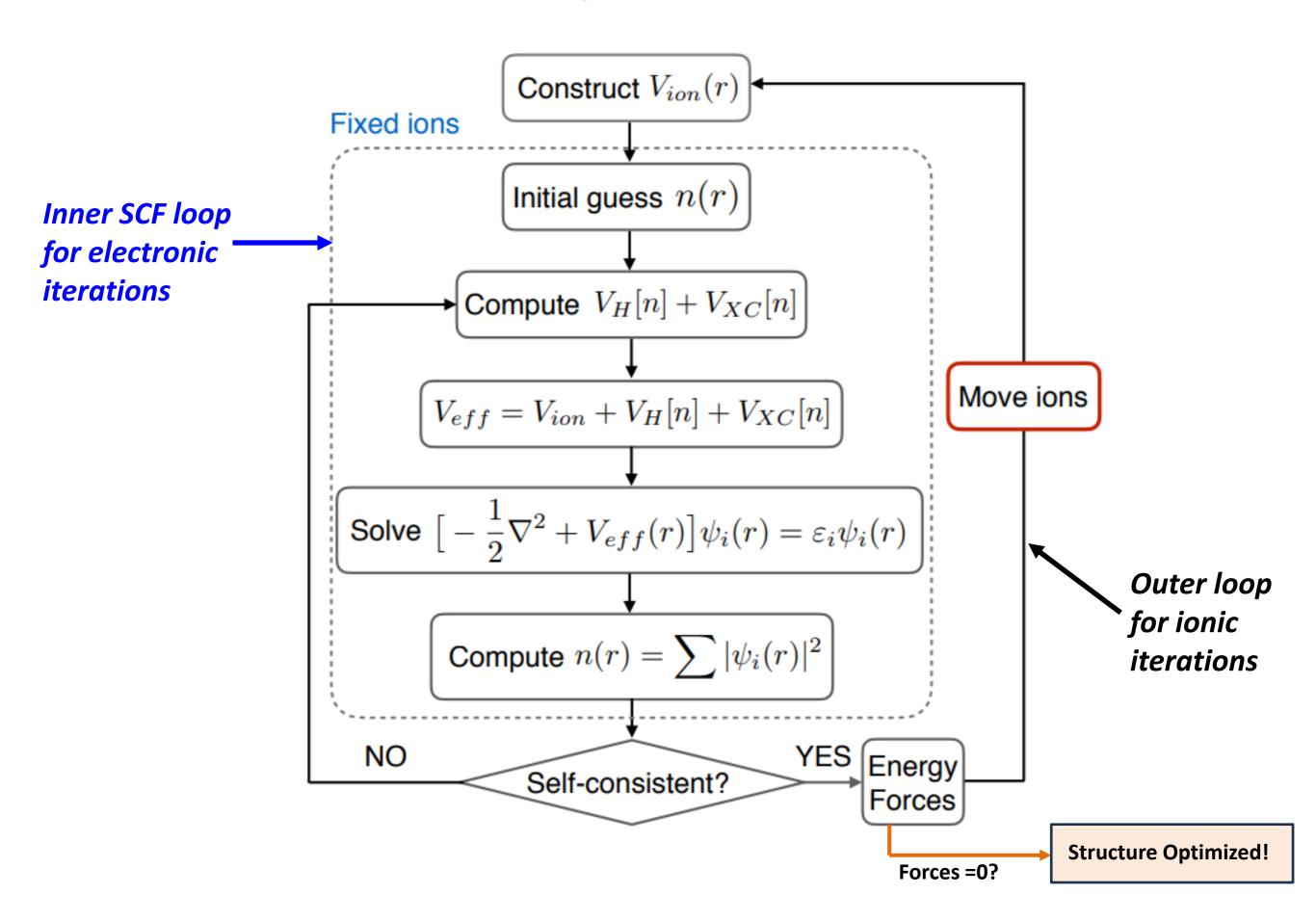
as above, with simple Thomas-Fermi screening (for highly homogeneous systems)

'local-TF':

as above, with local-densitydependent TF screening (for highly inhomogeneous systems)



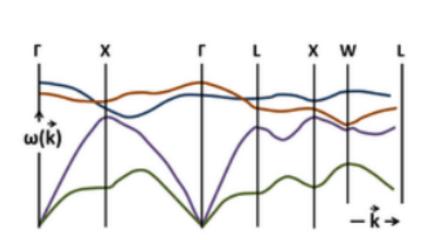
Forces & Geometry Optimization

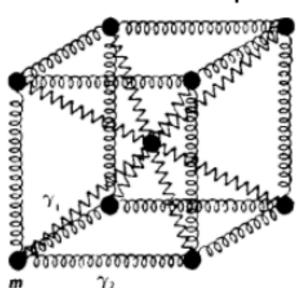


Forces on atoms easily calculated when electron ground state is obtained

$$F_{I} = -\frac{dE}{dr_{I}} = -\left|\psi_{i} \left| \frac{\partial \widehat{H}}{\partial r_{I}} \right| \psi_{i}\right\rangle$$

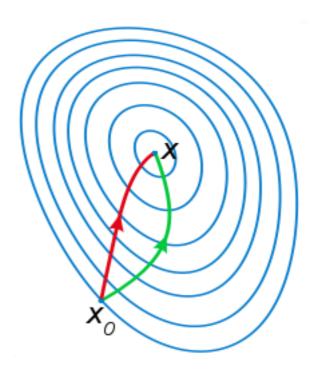
- By moving along the ionic forces (steepest descent) the ionic ground can be calculated
- We can displace ions from the ionic groundstate, and determine the forces on all other ions
 - Effective interatomic force constants and vibrational frequences





Structure optimization

```
&CONTROL
calculation='vc-relax',
 restart_mode='from_scratch'
    prefix='si'
pseudo_dir='../pseudo/',
outdir='../tmp/',
forc_conv_thr=1d-5,
&SYSTEM
ibrav=2.
celldm(1)=10.2625,
nat=2.
ntyp=1.
ecutwfc=60.0.
ecutrho=720.0,
&ELECTRONS
mixing_beta=0.7,
conv_thr=1d-8,
&IONS
ion_dynamics='bfgs', €
&CELL
cell_dynamics='bfgs',
press=0.0,
press_conv_thr=0.5,
ATOMIC_SPECIES
Si 28.0855 Si.pbe-rrkj.UPF
ATOMIC_POSITIONS (alat)
Si 0.00 0.00 0.00
si 0.25 0.25 0.25
K_POINTS automatic
4 4 4 1 1 1
```



Broyden–Fletcher–Goldfarb–Shanno (BFGS) algorithm is an iterative method for solving unconstrained nonlinear optimization problems.

Structure optimization

```
&CONTROL
 calculation='vc-relax',
 restart_mode='from_scratch'
    prefix='si'
 pseudo_dir='../pseudo/',
 outdir='../tmp/',
 forc_conv_thr=1d-5,
&SYSTEM
 ibrav=2,
(celldm(1)=10.2625,
nat=2.
 ntyp=1.
 ecutwfc=60.0,
 ecutrho=720.0,
&ELECTRONS
mixing_beta=0.7,
 conv_thr=1d-8,
&IONS
ion_dynamics='bfgs',
&CELL
cell_dynamics='bfgs',
 press=0.0,
 press_conv_thr=0.5,
ATOMIC SPECIES
Si 28.0855 Si.pbe-rrkj. UPF
ATOMIC_POSITIONS (alat)
si 0.00 0.00
               0.00
   0.25 0.25
                0.25
K_POINTS automatic
4 4 4 1 1 1
```

Cell lattice parameters and free (internal) coordinates of the atoms may be changed by relaxation

Properties Derived from the Ground State Energy

Second derivative of E _{tot} or F wrt	Physical property
d _i , d _i (atomic displacement)	Force spring const: phonons
E, E (E-field)	Dielectric constant
ε, ε (strain)	Elastic constant
Ε, ε	Piezo-electric constant
E, d _i	Born Dynamical charge
ϵ , d_i	Strain-phonon coupling
H, H (magnetic field)	Magnetic susceptibility
E, H	Magneto-electric constant
Η, ε	Piezo-magnetic constant