随机信号处理mooc(下)

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参考:《随机信号处理-西安电子科技大学-赵国庆》

https://www.bilibili.com/video/BV16s411p7iX

因为太长了,做个分篇,这里是41-72集

白噪声通过线性系统

白噪声均值为0,功率谱在无穷区间均匀分布

$$G_{x}(\omega) = \frac{N_{0}}{2}, \quad -\infty < \omega < \infty$$

$$R_{x}(\tau) = \frac{N_{0}}{2}\delta(\tau) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \frac{N_{0}}{2} e^{j\omega\tau} d\omega$$

$$\int_{0}^{\infty} \delta(\tau) d\tau = \frac{1}{2}$$

一般关系式

冲击响应函数/传递函数

h(t) H(w)

$$G_y(\omega) = G_x(\omega) |H(\omega)|^2 = \frac{N_0}{2} |H(\omega)|^2$$
 $R_y(\tau) = R_x(\tau) \otimes h(\tau) \otimes h(-\tau) = \frac{N_0}{2} \frac{1}{2\pi} \int_{-\infty}^{\infty} |H(\omega)|^2 e^{j\omega\tau} d\omega = \frac{N_0}{2} h(\tau) \otimes h(-\tau)$ 前者一般用来计算输出过程功率谱,后者用来计算输出过程相关函数

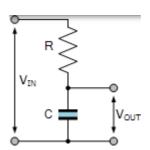
等效噪声通频带 Δf_n

说明输出后的等效宽度有多少

$$\Delta f_n = \frac{\int_0^\infty G(\omega)d\omega}{2\pi \left|G(\omega_0)\right|} \left|G(\omega_0)\right| = \max \left|G(\omega)\right|$$

反映了能量在频谱上的集中程度, 越大代表分布越宽, 频谱占据越宽变化越大

通过RC低频滤波器(积分器)



$$\begin{split} G_{y}(\omega) &= G_{x}(\omega) \big| H(\omega) \big|^{2} = \frac{N_{0}}{2} \frac{a^{2}}{\omega^{2} + a^{2}} \\ R_{y}(\tau) &= \frac{N_{0}a^{2}}{2} \frac{2}{2\pi} \int_{0}^{\infty} \frac{1}{\omega^{2} + a^{2}} cos(\omega \tau) d\omega = \frac{N_{0}}{2} \frac{1}{2} a e^{-a|\tau|} = \frac{N_{0}}{4} a e^{-a|\tau|} \end{split}$$

等效噪声通频带

$$\Delta f_n = \frac{\frac{N_0}{2} \int_0^\infty \frac{a^2}{\omega^2 + a^2} d\omega}{2\pi \frac{N_0}{2}} = \frac{1}{2\pi} \int_0^\infty \frac{a^2}{\omega^2 + a^2} d\omega = \frac{a}{4} = \frac{1}{4RC}$$

RC越大,过滤的高频分量越多,频带越小

相关系数

$$r_y(\tau) = \frac{C_y(\tau)}{C_y(0)} = \frac{R_y(\tau)}{R_y(0)} = \frac{\frac{N_0}{4}ae^{-a|\tau|}}{\frac{N_0}{4}a} = e^{-a|\tau|}$$

相关时间

$$\tau_0 = \int_0^\infty a^{-a\tau} d\tau = \frac{1}{a} = RC = \frac{1}{4\Delta f_n}$$

相关时间和等效噪声通频带关系

几乎所有的 $\tau_0 \propto \frac{1}{\Delta f_n}$ 都成立

通过理想的低通滤波器



$$\begin{aligned} \left| H(\omega) \right| &= \begin{cases} K_0 \ |\omega| \le \Delta \Omega \\ 0 \ others \end{cases} \\ G_y(\omega) &= \frac{N_0}{2} K_0^2 \ |\omega| \le \Delta \Omega \\ R_y(\tau) &= \frac{N_0}{2} K_0^2 \frac{2}{2\pi} \int_0^\infty \cos(\omega \tau) \, d\omega = \frac{N_0}{2} K_0^2 \frac{\sin \Delta \Omega \tau}{\pi \tau} \end{aligned}$$

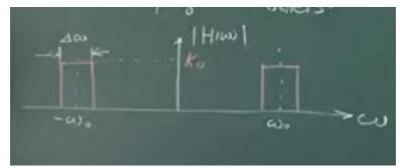
频域矩形谱, 时域为sin函数

$$\Delta f_n = \frac{\frac{N_0}{2} K_0^2 \Delta \Omega}{2\pi \frac{N_0}{2} K_0^2} = \frac{\Delta \Omega}{2\pi}$$

即为矩形谱的宽度ΔΩ

$$r_{y}(\tau) = \frac{\frac{N_{0}}{2}K_{0}^{2}\frac{\sin\Delta\Omega\tau}{\pi\tau}}{\frac{N_{0}}{2}K_{0}^{2}\frac{\Delta\Omega}{\pi}} = \frac{\sin\Delta\Omega\tau}{\Delta\Omega\tau}$$
$$\tau_{0} = \int_{0}^{\infty} \frac{\sin\Delta\Omega\tau}{\Delta\Omega\tau} d\tau = \frac{\pi}{2\Delta\Omega} = \frac{1}{4\Delta f_{n}}$$

通过理想的带通滤波器



$$\begin{split} \left| H(\omega) \right| &= \begin{cases} K_0 \left| \omega \pm \omega_0 \right| \leq \Delta \omega / 2 \\ 0 \text{ others} \end{cases} \\ G_y(\omega) &= \frac{N_0}{2} K_0^2 \left| \omega \pm \omega_0 \right| \leq \frac{\Delta \omega}{2} \\ R_y(\tau) &= \frac{N_0}{2} K_0^2 \frac{2}{2\pi} \int_{\omega_0 - \frac{\Delta \omega}{2}}^{\omega_0 + \frac{\Delta \omega}{2}} \cos(\omega \tau) \, d\omega = \frac{N_0}{2} K_0^2 \frac{1}{2\pi \tau} \cos\omega_0 \tau \sin\frac{\Delta \omega \tau}{2} \end{split}$$

在sin函数的基础上加了一个载波

$$\Delta f_n = \frac{\frac{N_0}{2} K_0^2 \Delta \omega}{2\pi \frac{N_0}{2} K_0^2} = \frac{\Delta \omega}{2\pi} = \Delta f$$

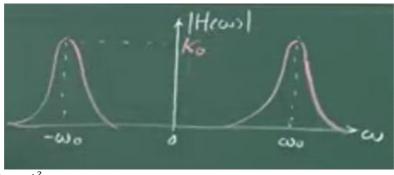
角频率转为物理频率

$$r_{y}(\tau) = \frac{\frac{N_{0}}{2}{K_{0}}^{2} \frac{1}{2\pi\tau} cos\omega_{0}\tau sin\frac{\Delta\omega\tau}{2}}{\frac{N_{0}}{2}{K_{0}}^{2} \frac{1}{2\pi}\frac{\Delta\omega}{2}} = \frac{2cos\omega_{0}\tau sin\frac{\Delta\omega\tau}{2}}{\Delta\omega\tau}$$

相关时间不看载波,只看包络

$$\tau_0 = \int_0^\infty \frac{2}{\Delta\omega\tau} \sin\frac{\Delta\omega\tau}{2} d\tau = \frac{\pi}{\Delta\omega} = \frac{1}{2\Delta f_n}$$

通过高斯型带通滤波器



$$\begin{split} |H(\omega)| &= K_0 \, e^{-\frac{\left(\omega \pm \omega_0\right)^2}{2\beta^2}} - \infty < \omega < \infty \\ G_y(\omega) &= \frac{N_0}{2} K_0^2 e^{-\frac{\left(\omega \pm \omega_0\right)^2}{\beta^2}} - \infty < \omega < \infty \\ R_y(\tau) &= \frac{N_0}{2} K_0^2 \frac{2}{2\pi} \int_{-\infty}^{\infty} e^{-\frac{\left(\omega - \omega_0\right)^2}{\beta^2}} \cos \omega \tau d\omega \, \, \Leftrightarrow \, \omega - \omega_0 = \omega' \\ &= \frac{N_0}{2} K_0^2 \frac{2}{2\pi} \int_{-\infty}^{\infty} e^{-\frac{\omega'^2}{\beta^2}} \cos \left(\omega_0 + \omega'\right) \tau d\omega' = \frac{N_0}{2} K_0^2 \frac{2}{2\pi} \int_{-\infty}^{\infty} e^{-\frac{\omega'^2}{\beta^2}} \cos \omega' \tau d\omega' \cos \omega_0 \tau \\ &= \frac{N_0}{2} K_0^2 \frac{\beta}{\sqrt{\pi}} \cos \omega_0 \tau e^{-\frac{\tau^2 \beta^2}{4}} \end{split}$$

$$\Delta f_n = \frac{\frac{N_0}{2} K_0^2 \int_{-\infty}^{\infty} e^{-\frac{(\omega - \omega_0)^2}{\beta^2}} d\omega}{2\pi \frac{N_0}{2} K_0^2} = \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{-\frac{(\omega - \omega_0)^2}{\beta^2}} d\omega = \frac{\beta}{2\sqrt{\pi}}$$

$$r_{y}(\tau) = \frac{\frac{N_{0}}{2}K_{0}^{2} \frac{\beta}{\sqrt{\pi}} cos\omega_{0}\tau e^{-\frac{\tau^{2}\beta^{2}}{4}}}{\frac{N_{0}}{2}K_{0}^{2} \frac{\beta}{\sqrt{\pi}}} = cos\omega_{0}\tau e^{-\frac{\tau^{2}\beta^{2}}{4}}$$

载波项依然存在

载波项依然不参加积分

$$\tau_0 = \int_0^\infty e^{-\frac{\tau^2 \beta^2}{4}} d\tau = \frac{\sqrt{\pi}}{\beta} = \frac{1}{2\Delta f_n}$$

随机过程线性变换后的概率分布