

Physics 2130: Assignment III

Jeremy Favro

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Problem 1. Consider the driven, damped harmonic oscillator. Its equation of motion is:

$$\ddot{x} + 2\beta\dot{x} + \omega_0^2 = A \cos(\omega t)$$

In the case of an underdamped oscillator, i.e., $\beta < \omega_0^2$ we found that the solution for the equation of motion is:

$$x(t) = x_c(t) + x_p(t)$$

Where:

$$x_c(t)e^{-\beta t} [c_1 e^{i\omega_1 t} + c_2 e^{-i\omega_1 t}] = B e^{-\beta t} \cos(\omega_1 t - \phi) = \Gamma(t)s(t)$$

And where:

$$\omega_1 = \sqrt{\omega_0^2 - \beta^2}$$

$$\Gamma(t) = B e^{-\beta t}$$

$$s(t) = \cos(\omega_1 t - \phi)$$

The particular solution is instead:

$$x_p(t) = D \cos(\omega t - \delta)$$

Where:

$$D = \frac{A}{\sqrt{(\omega_0^2 - \omega^2)^2 + 4\omega^2\beta^2}}$$

$$\delta = \arctan\left(\frac{2\omega\beta}{\omega_0^2 - \omega^2}\right)$$

- Consider an underdamped driven oscillator that starts with the initial conditions $x(t=0) = x_0$ and $\dot{x}(t=0) = 0$. Find the analytical expressions for the unknown coefficients in $x(t)$ using these initial conditions.
- Write a python program that returns (and prints) the values of ω_1 , β , ϕ , D and δ for a given x_0 . This should be coded as a function called `harm_osc_params` that accepts as inputs ω_0 , β , A , ω and x_0 .
- In the same python program, now write a new function, `harm_osc_x_pos`, that calculates the array of positions x of the harmonic oscillator for a given array of times t . This function should receive the array t as an input as well as the values of ω_0 , β , A , ω and x_0 . It should call the previously written function `harm_osc_params` for the calculations of all the oscillation parameters. It should return the array of positions x of the harmonic oscillator for each value of t .
- In the same python program, now write a new function to plot the data, `harm_osc_x_plot_single`. The function receives as inputs the arrays t and x generated at the previous step.
- Pick three values of β (remember of the constraint $\beta < \omega_0$) and plot in the same graph $x(t)$ for the three chosen values. For reference use $\omega_0 = 1\text{s}^{-1}$ and $A = 1\text{s}^{-2}$ (but play with the values of A to see the relative importance of the driver and the damping). Comment on what effect β has on both the transient and steady-state solution.
- To help you distinguish the different effects, now write a function called `harm_osc_damp_drive` that plots in the same graph, $s(t)$, $\Gamma(t)$, $x_c(t)$ and $x_p(t)$. Comment on the results, specifically on the contribution of each component.

- (g) Finally, write a new function called `harm_osc_euler_cromer` that receives has input parameters ω_0 , β , A , ω , x_0 , and the analytical solution $x(t)$. This function should calculate a new $x(t)$, called $x_{EC}(t)$, that uses the Euler-Cromer method to determine the position of the oscillator as a function of time for the same initial conditions. Additionally, the function should plot, on the same graph, the solutions $x(t)$ and $x_{EC}(t)$ obtained with the analytical and Euler-Cromer methods, respectively, and the residuals, i.e., the difference $x(t) - x_{EC}(t)$. Discuss how you choose a value of Δt that gives a sufficiently accurate answer (which means defining “sufficiently accurate”).

Solution 1.

- (a)

$$\begin{aligned} x(t=0) &= B \cos(-\phi) + D \cos(-\delta) = x_0 \\ \Rightarrow x_0 - D \cos(\delta) &= B \cos(\phi) \\ \Rightarrow \frac{x_0 - D \cos(\delta)}{\cos(\phi)} &= B \end{aligned}$$

$$\begin{aligned} \dot{x}(t=0) &= -B\omega_1 \sin(-\phi) - B\beta \cos(-\phi) - D\omega \sin(-\delta) = 0 \\ &= B\omega_1 \sin(\phi) - B\beta \cos(\phi) + D\omega \sin(\delta) \\ &= \frac{x_0 - D \cos(\delta)}{\cos(\phi)} \omega_1 \sin(\phi) - \frac{x_0 - D \cos(\delta)}{\cos(\phi)} \beta \cos(\phi) + D\omega \sin(\delta) \\ &= (x_0 - D \cos(\delta)) \omega_1 \tan(\phi) - (x_0 - D \cos(\delta)) \beta + D\omega \sin(\delta) \\ \Rightarrow -\frac{D\omega \sin(\delta)}{(x_0 - D \cos(\delta))} &= \omega_1 \tan(\phi) - \beta \\ \Rightarrow \arctan\left(\frac{\beta - \frac{D\omega \sin(\delta)}{(x_0 - D \cos(\delta))}}{\omega_1}\right) &= \phi \end{aligned}$$

- (b) `import numpy as np`

```
def harm_osc_params(w_0: float,
                    beta: float,
                    A: float,
                    w: float,
                    x_0: float
                    ) -> tuple[float, float, float, float, float]:

    # "Given"
    w_1 = np.sqrt(w_0**2 - beta**2)
    D = A / (np.sqrt((w_0**2 - w**2)**2 + 4 * (w**2) * (beta**2)))
    delta = np.arctan2((2 * w * beta), (w_0**2 - w**2))

    # Unknowns
    phi = np.arctan2(
        ((beta - (D * w * np.sin(delta))) / (x_0 - D * np.cos(delta))),
        w_1)
    B = (x_0 - D * np.cos(delta)) / (np.cos(phi))

    print(f"w_1: {w_1}", f"B: {B}", f"phi: {phi}",
          f"D: {D}", f"delta: {delta}", sep="\n")

    return w_1, B, phi, D, delta
```

```

(c) import numpy as np
    import q1b

    def harm_osc_x_pos(w_0: float,
                      beta: float,
                      A: float,
                      w: float,
                      x_0: float,
                      t: np.ndarray
                      ) -> list:
        w_1, B, phi, D, delta = q1b.harm_osc_params(w_0,
                                                    beta,
                                                    A,
                                                    w,
                                                    x_0)

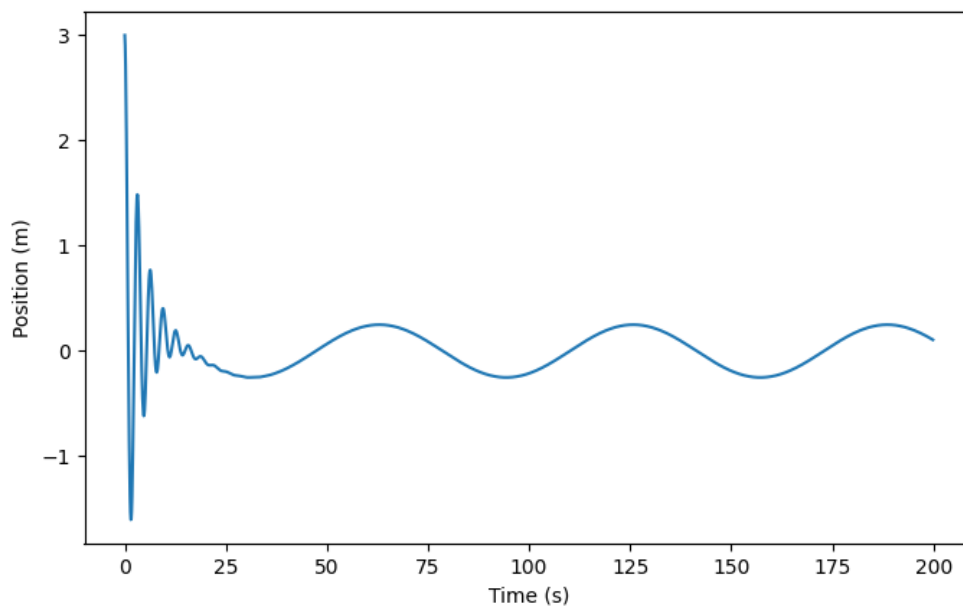
        return B*np.exp(-beta*t)*np.cos(w_1*t-phi)+D*np.cos(w*t-delta)

(d) import matplotlib.pyplot as pyplot
    import numpy as np
    import q1c

    def harm_osc_x_plot_single(x: np.ndarray, t: np.ndarray) -> None:
        pyplot.plot(t, x)
        pyplot.xlabel("Time (s)")
        pyplot.ylabel("Position (m)")
        pyplot.show()

    t_vals = np.arange(0, 200, 0.1)
    x_vals = q1c.harm_osc_x_pos(2, 0.25, 1, 0.1, 3, t_vals)
    harm_osc_x_plot_single(x_vals, t_vals)

```



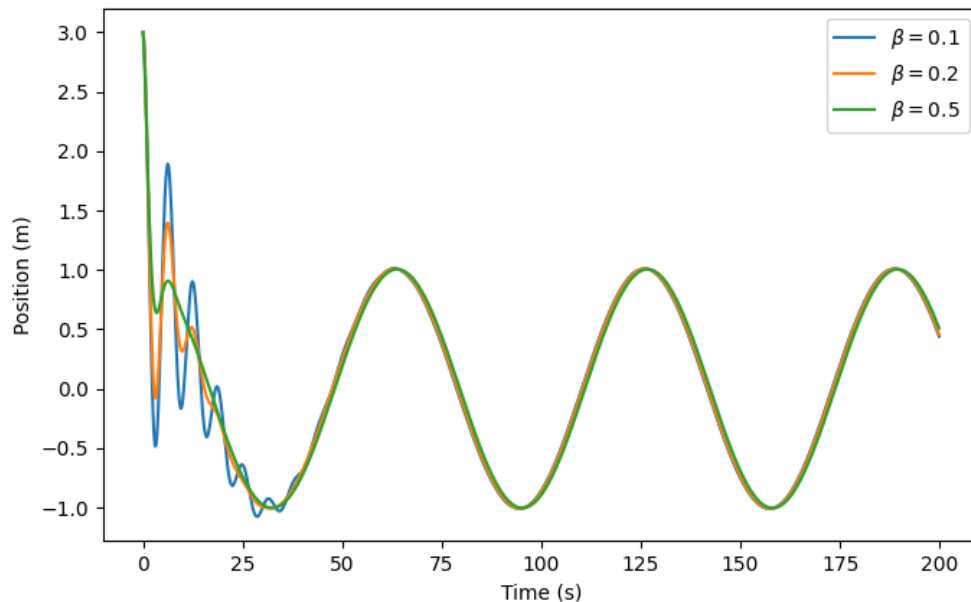
The first regime, visible as the rapid oscillations from $t = 0$ to $t \approx 30$ is dominated by the $Be^{-\beta t} \cos(\omega_1 t - \phi)$ term. Because This term contains an exponential decay factor it will fade away as time progresses leaving the $D \cos(\omega t - \delta)$ term as the only oscillator. The first regime has a period of $\approx 3.2 \text{ rad s}^{-1}$, which is consistent with $\omega_1 = \frac{2\pi}{\sqrt{w_0^2 - \beta^2}} = \frac{2\pi}{\sqrt{2^2 - 0.25^2}} \approx 3.2 \text{ rad s}^{-1}$. The second regime has a period of $\approx 63 \text{ rad s}^{-1}$, which is consistent with $\omega = \frac{2\pi}{63} \approx 0.1$.

```
(e) import numpy as np
import matplotlib.pyplot as pyplot
import q1c

t_vals = np.arange(0, 200, 0.1)

for beta in [0.1, 0.2, 0.5]:
    x_vals = q1c.harm_osc_x_pos(1, beta, 1, 0.1, 3, t_vals)
    pyplot.plot(t_vals, x_vals, label=r"$\beta$={beta}")

pyplot.xlabel("Time (s)")
pyplot.ylabel("Position (m)")
pyplot.legend()
pyplot.show()
```



As is expected here, greater values of β result in quicker die-off of the transient solution. β only slightly effects the steady-state solution, causing slight displacement as the transient solution never actually reaches zero.

```
(f) import numpy as np
import matplotlib.pyplot as pyplot
import q1b

t_vals = np.arange(0, 200, 0.1)

def harm_osc_damp_drive():
```

```

t = np.arange(0, 100, 0.1)
w_0 = 2
beta = 0.25
A = 1
w = 0.1
x_0 = 3
w_1, B, phi, D, delta = q1b.harm_osc_params(w_0,
                                              beta,
                                              A,
                                              w,
                                              x_0)

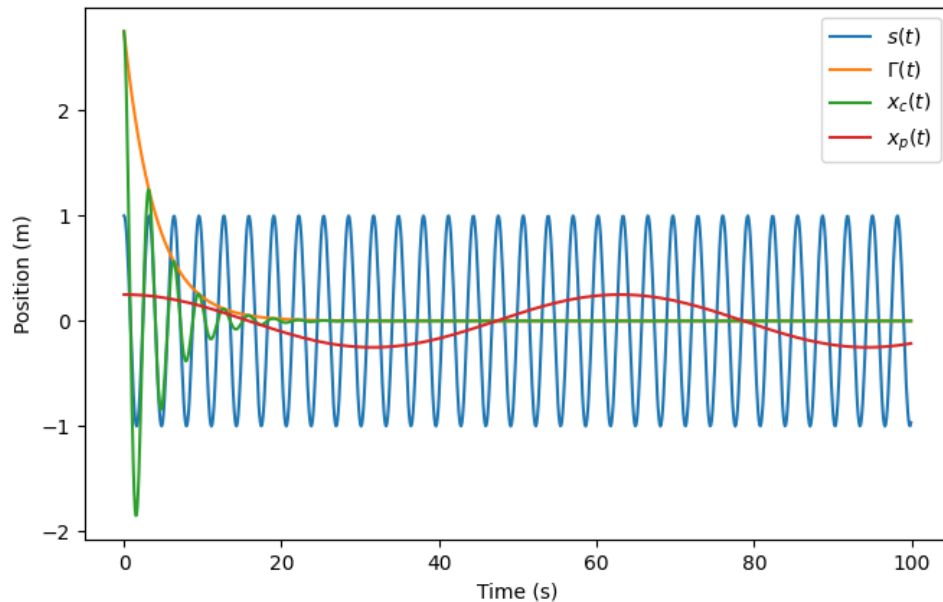
s = np.cos(w_1*t-phi)
gamma = B*np.e**(-beta*t)
x_c = s*gamma
x_p = D*np.cos(w*t-delta)

pyplot.plot(t, s, label=r"$s(t)$")
pyplot.plot(t, gamma, label=r"$\Gamma(t)$")
pyplot.plot(t, x_c, label=r"$x_c(t)$")
pyplot.plot(t, x_p, label=r"$x_p(t)$")

pyplot.xlabel("Time (s)")
pyplot.ylabel("Position (m)")
pyplot.legend()
pyplot.show()

harm_osc_damp_drive()

```



$s(t)$ contributes the oscillation of the transient solution. $\Gamma(t)$ contributes the exponential decay envelope. $x_c(t)$ contributes, being the product of $s(t)$ and $\Gamma(t)$ produces the initial decaying transient solution. $x_p(t)$

contributes the steady-state solution.

```
(g) import numpy as np
import matplotlib.pyplot as pyplot
import q1c

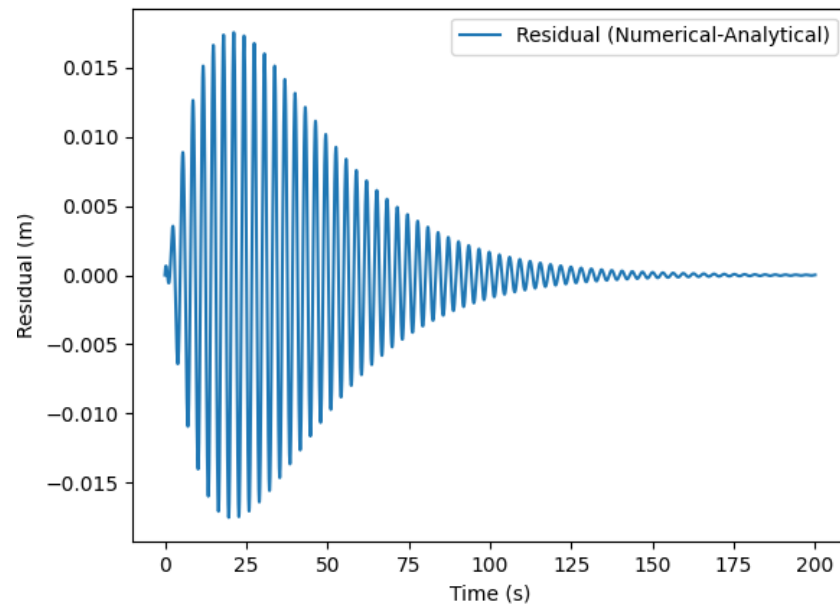
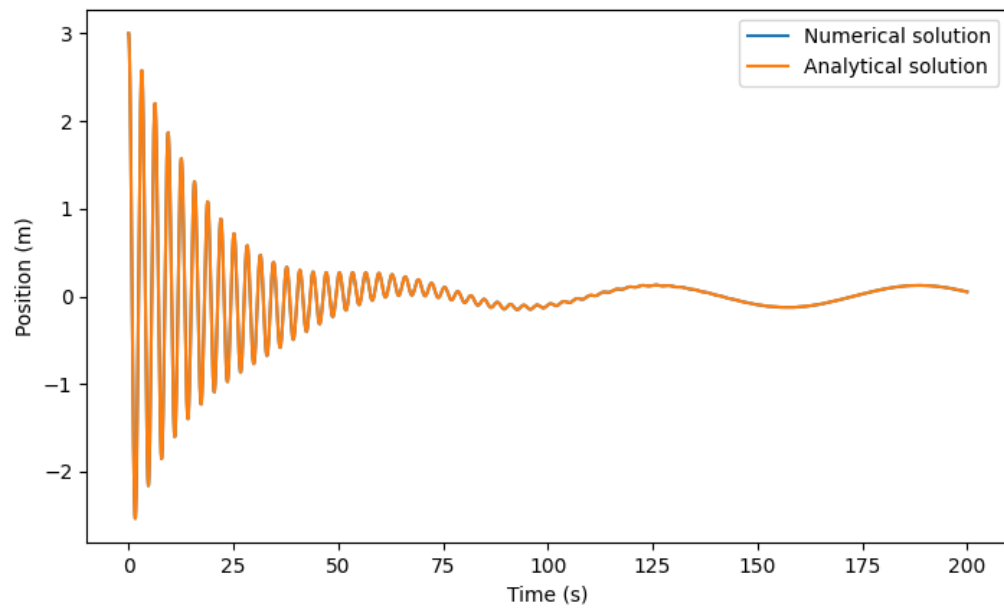
def harm_osc_euler_cromer(w_0: float,
                          beta: float,
                          A: float,
                          w: float,
                          x_0: float,
                          t: np.ndarray
                          ):
    dt = t[1]-t[0]
    x = np.full_like(t, x_0)
    x_r = q1c.harm_osc_x_pos(w_0, beta, A, w, x_0, t)
    v = np.zeros_like(t)
    for t_i in range(len(t) - 1):
        a = A*np.cos(w*t[t_i])-2*beta*v[t_i]-(w_0**2)*x[t_i]

        v[t_i+1] = v[t_i] + a*dt
        x[t_i+1] = x[t_i] + v[t_i+1]*dt

    pyplot.plot(t, x, label="Numerical solution")
    pyplot.plot(t, x_r, label="Analytical solution")
    pyplot.xlabel("Time (s)")
    pyplot.ylabel("Position (m)")
    pyplot.legend()
    pyplot.figure(0)

    pyplot.plot(t, x-x_r, label="Residual (Numerical-Analytical)")
    pyplot.xlabel("Time (s)")
    pyplot.ylabel("Residual (m)")
    pyplot.legend()
    pyplot.show()

t_vals = np.arange(0, 200, 1/64)
harm_osc_euler_cromer(2, 0.05, 1/2, 0.1, 3, t_vals)
```



Note that here I plotted the residuals and numerical-analytical solutions on different graphs due to the difference in amplitude of each

Problem 2. We want to study now the behaviour of the **non-linear (physical) pendulum**. The analytical solutions we generally get for oscillating systems are derived under the small-angle approximation, which allows us to write $\sin(\theta) \approx \theta$. If one drops this hypothesis things change, and the motion now becomes dependent on the amplitude. The equation of motion is now:

$$\frac{d^2\theta}{dt^2} = -\frac{g}{l} \sin(\theta) - q \frac{d\theta}{dt} + F_d \sin(\Omega_D t)$$

Where the only difference from the linear case is that we now have $\sin \theta$ instead of θ . This equation of motion has no analytical solutions; therefore, we need to solve it numerically. We thus need to write the equations for angular acceleration and velocity and use them, e.g., with the Euler-Cromer method:

$$\frac{d\omega}{dt} = -\frac{g}{l} \sin(\theta) - q \frac{d\theta}{dt} + F_d \sin(\Omega_D t)$$

$$\frac{d\theta}{dt} = \omega$$

- (a) Write a python program that calculates numerically $\theta(t)$ using the Euler-Cromer method.
- (b) Write another python program (you can mostly base it on what you wrote in (2a)) that calculates two separate solutions $\theta_1(t)$ and $\theta_2(t)$, where $\theta_1(t)$ starts with an initial angle $\theta_0 = 10^\circ$ while $\theta_2(t)$ starts with an ever so slightly different initial angle $\theta_0 = 10.05^\circ$.
- (c) In the program you wrote in (2b), add a phase-space plot where you now plot ω vs. θ .

Solution 2:

```
import numpy as np
import matplotlib.pyplot as pyplot

g = 9.8
t = np.arange(0, 200, 0.1)

def euler_chromer_nonlinear(theta_0: float,
                             w_0: float,
                             l: float,
                             q: float,
                             F_d: float,
                             W: float,
                             t_vals: np.ndarray) -> tuple[np.ndarray, np.ndarray]:
    theta = np.full_like(t_vals, np.deg2rad(theta_0))
    omega = np.full_like(t_vals, w_0)
    dt = t_vals[1]-t_vals[0]
    for t_i in range(len(t_vals)- 1):
        omega[t_i+1] = omega[t_i] + \
            (-(g/l)*np.sin(theta[t_i])-q *
             omega[t_i] + F_d*np.sin(W*t_vals[t_i]))*dt
        theta[t_i+1] = theta[t_i]+omega[t_i+1]*dt

    return theta

pyplot.figure(0)
pyplot.title(r"$F_D=0.5$")
for i in [10, 25, 45]:
    pyplot.plot(t, euler_chromer_nonlinear(
        i, 0, g, 0.5, 0.5, 2/3, t), label=rf"$\theta={i}^\circ$")
pyplot.xlabel("Time (s)")
pyplot.ylabel("Position (rad)")
pyplot.legend()
```

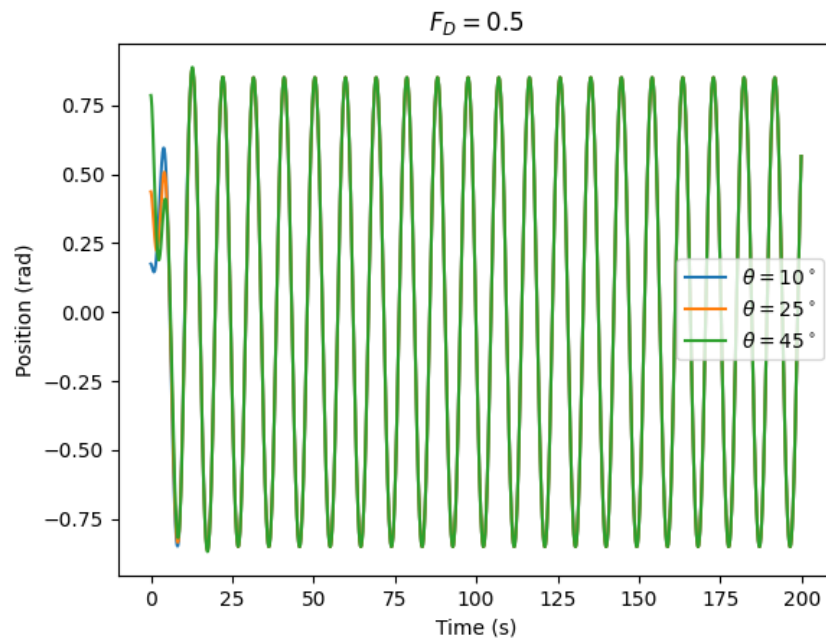
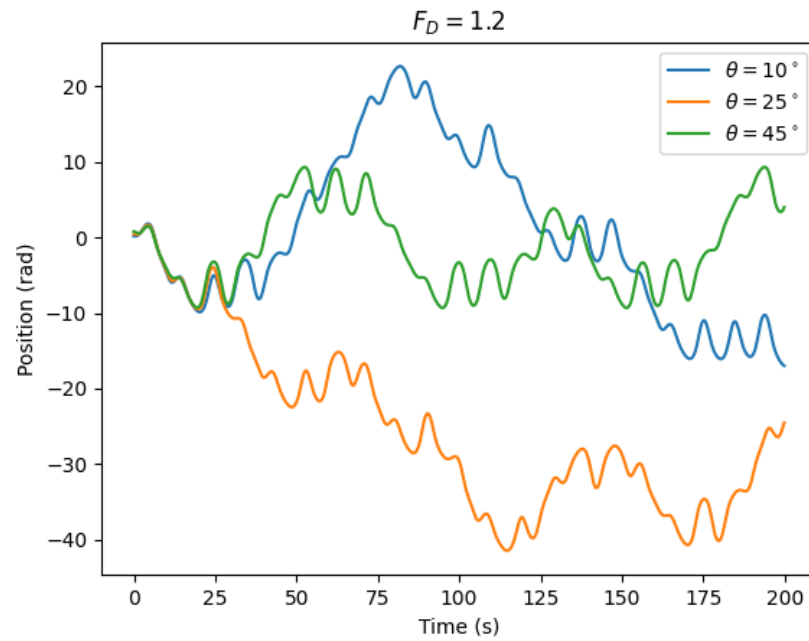


```

pyplot.figure(1)
pyplot.title(r"$F_D=1.2$")
for i in [10, 25, 45]:
    pyplot.plot(t, euler_chromer_nonlinear(
        i, 0, g, 0.5, 1.2, 2/3, t), label=r"$\theta={i}^\circ$")
pyplot.xlabel("Time (s)")
pyplot.ylabel("Position (rad)")
pyplot.legend()

pyplot.show()

```



(b) `import numpy as np`

```

import matplotlib.pyplot as pyplot

g = 9.8

def euler_chromer_nonlinear(theta_0: float,
                             w_0: float,
                             l: float,
                             q: float,
                             F_d: float,
                             W: float,
                             t_vals: np.ndarray) -> tuple[np.ndarray, np.ndarray]:
    theta = np.full_like(t_vals, np.deg2rad(theta_0))
    omega = np.full_like(t_vals, w_0)
    dt = t_vals[1]-t_vals[0]
    for t_i in range(len(t_vals) - 1):
        omega[t_i+1] = omega[t_i] + \
            (-(g/l)*np.sin(theta[t_i])-q *
             omega[t_i] + F_d*np.sin(W*t_vals[t_i]))*dt
        theta[t_i+1] = theta[t_i]+omega[t_i+1]*dt

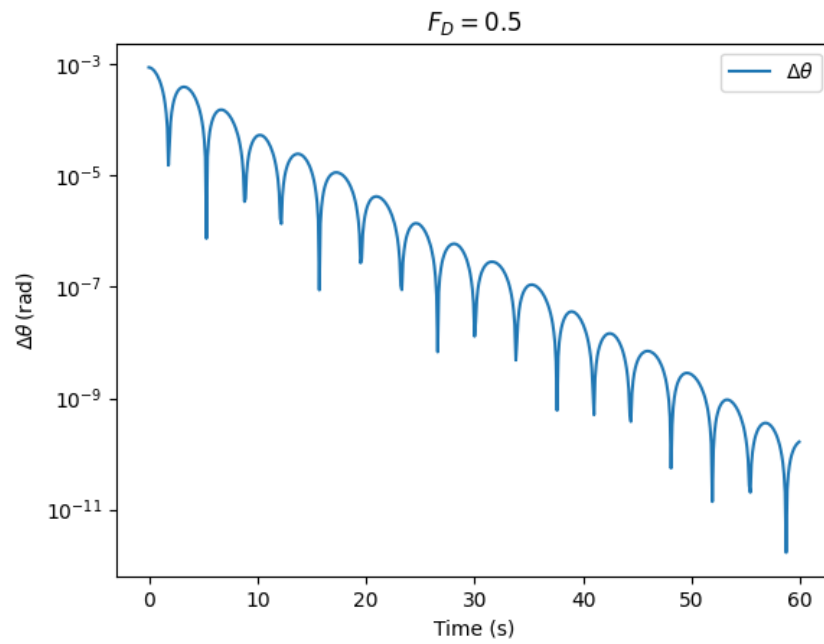
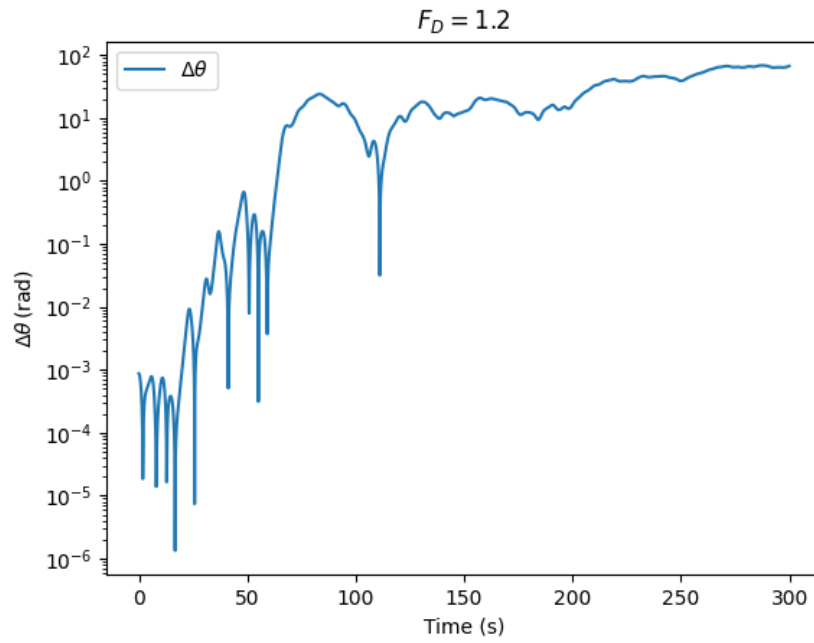
    return theta

t = np.arange(0, 60, 0.1)
pyplot.figure(0)
pyplot.title(r"$F_D=0.5$")
theta_0 = euler_chromer_nonlinear(10, 0, g, 0.5, 0.5, 2/3, t)
theta_1 = euler_chromer_nonlinear(10.05, 0, g, 0.5, 0.5, 2/3, t)
pyplot.yscale("log")
pyplot.plot(t, np.abs(theta_0-theta_1), label=r"$\Delta\theta$")
pyplot.xlabel("Time (s)")
pyplot.ylabel(r"$\Delta\theta$, \text{(rad)}$")
pyplot.legend()

t = np.arange(0, 300, 0.1)
pyplot.figure(1)
pyplot.title(r"$F_D=1.2$")
theta_0 = euler_chromer_nonlinear(10, 0, g, 0.5, 1.2, 2/3, t)
theta_1 = euler_chromer_nonlinear(10.05, 0, g, 0.5, 1.2, 2/3, t)
pyplot.yscale("log")
pyplot.plot(t, np.abs(theta_0-theta_1), label=r"$\Delta\theta$")
pyplot.xlabel("Time (s)")
pyplot.ylabel(r"$\Delta\theta$, \text{(rad)}$")
pyplot.legend()

pyplot.show()

```



```
(c) import numpy as np
import matplotlib.pyplot as pyplot

g = 9.8
t = np.arange(0, 300, 0.1)

def euler_chromer_nonlinear(theta_0: float,
                             w_0: float,
                             l: float,
                             q: float,
                             F_d: float,
                             W: float,
```

```

        t_vals: np.ndarray) -> tuple[np.ndarray, np.ndarray]:
    theta = np.full_like(t_vals, np.deg2rad(theta_0))
    omega = np.full_like(t_vals, w_0)
    dt = t_vals[1]-t_vals[0]
    for t_i in range(len(t_vals) - 1):
        omega[t_i+1] = omega[t_i] + \
            (-(g/l)*np.sin(theta[t_i])-q *
             omega[t_i] + F_d*np.sin(W*t_vals[t_i]))*dt
        theta[t_i+1] = theta[t_i]+omega[t_i+1]*dt

    return theta, omega

pyplot.figure(0)
pyplot.title(r"$F_D=0.5$")
theta_0, omega_0 = euler_chromer_nonlinear(10, 0, g, 0.5, 0.5, 2/3, t)

pyplot.plot(theta_0, omega_0, label=r"$\Delta\theta$")
pyplot.ylabel(r"$\omega \, (s^{-1})$")
pyplot.xlabel(r"$\theta \, (rad)$")
pyplot.legend()

pyplot.figure(1)
pyplot.title(r"$F_D=1.2$")
theta_0, omega_0 = euler_chromer_nonlinear(10, 0, g, 0.5, 1.2, 2/3, t)

pyplot.plot(theta_0, omega_0, label=r"$\Delta\theta$")
pyplot.ylabel(r"$\omega \, (s^{-1})$")
pyplot.xlabel(r"$\theta \, (rad)$")
pyplot.legend()

pyplot.show()

```

