# Physics 3200Y: Assignment I

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**Problem 1.** Let  $\mathbf{z} = \mathbf{r} - \mathbf{r}'$  be the separation between  $\mathbf{r}$  and  $\mathbf{r}'$ , where  $\mathbf{r}' = (x', y', z')$  is a fixed point and  $\mathbf{r} = (x, y, z)$ . Let  $z = |\mathbf{z}|$  be the magnitude of the separation.

- (a) Show that  $\nabla (r^2) = 2\mathbf{r}$ .
- (b) Show that  $\nabla \exp\left(\vec{k} \cdot \vec{\mathbf{z}}\right) = \vec{k} \exp\left(\vec{k} \cdot \vec{\mathbf{z}}\right)$ , where  $\vec{k}$  is a vector constant.
- (c) Show that  $\nabla \exp(kr) = k\hat{\mathbf{i}} \exp(kr)$ .
- (d) Show that  $\nabla (1/r) = -\hat{\mathbf{x}}/r^2$ .

### Solution 1.

(a) Proof.

$$\begin{split} &= \nabla \left( t^2 \right) \\ &= \left[ \frac{\partial}{\partial x} \hat{x} + \frac{\partial}{\partial y} \hat{y} + \frac{\partial}{\partial z} \hat{z} \right] \sqrt{\left( x - x' \right)^2 + \left( y - y' \right)^2 + \left( z - z' \right)^2} \\ &= \left[ \frac{\partial}{\partial x} \hat{x} + \frac{\partial}{\partial y} \hat{y} + \frac{\partial}{\partial z} \hat{z} \right] \left[ \left( x - x' \right)^2 + \left( y - y' \right)^2 + \left( z - z' \right)^2 \right] \\ &= \frac{\partial}{\partial x} \left( x - x' \right)^2 \hat{x} + \frac{\partial}{\partial y} \left( y - y' \right)^2 \hat{y} + \frac{\partial}{\partial z} \left( z - z' \right)^2 \hat{z} \quad \text{Note}^1 \\ &= 2 \left( x - x' \right) \hat{x} + 2 \left( y - y' \right) \hat{y} + 2 \left( z - z' \right) \hat{z} \quad \text{By chain rule} \\ &= 2 \left( \left( x - x' \right), \left( y - y' \right), \left( z - z' \right) \right) = 2 \mathbf{z} \end{split}$$

(b) Proof.

$$\begin{split} &= \nabla \exp\left(\vec{k} \cdot \vec{a}\right) \\ &= \nabla \exp\left(k_x(x - x') + k_y(y - y') + k_z(z - z')\right) \\ &= k_x \exp\exp\left(k_x(x - x') + k_y(y - y') + k_z(z - z')\right) \hat{x} \\ &+ k_y \exp\left(k_x(x - x') + k_y(y - y') + k_z(z - z')\right) \hat{y} \\ &+ k_z \exp\left(k_x(x - x') + k_y(y - y') + k_z(z - z')\right) \hat{z} \\ &= (k_x, k_y, k_z) \exp\left(k_x(x - x') + k_y(y - y') + k_z(z - z')\right) \\ &= \vec{k} \exp\left(\vec{k} \cdot \vec{a}\right) \end{split}$$

<sup>&</sup>lt;sup>1</sup>The partials kill the terms that don't contain their variable of differentiation, omitted for brevity

(c) Proof.

$$\begin{split} &= \nabla \exp{(k x)} \\ &= \nabla \exp{\left(\sqrt{k_x^2 + k_y^2 + k_z^2} \sqrt{(x - x')^2 + (y - y')^2 + (z - z')^2}\right)} \\ &= \nabla \exp{\left(\sqrt{k_x^2 + k_y^2 + k_z^2} \sqrt{(x - x')^2 + (y - y')^2 + (z - z')^2}\right)} \\ &= \exp{\left(\sqrt{k_x^2 + k_y^2 + k_z^2} \sqrt{(x - x')^2 + (y - y')^2 + (z - z')^2}\right)} \\ &= \tilde{k} \nabla \sqrt{(x - x')^2 + (y - y')^2 + (z - z')^2} \quad \text{By chain rule} \\ &= \frac{\sqrt[3]{k} \exp{\left(\sqrt{(x - x')^2 + (y - y')^2 + (z - z')^2}\right)}}{\sqrt{(x - x')^2 + (y - y')^2 + (z - z')^2}} \\ &= \tilde{k} \exp{(k x)} \frac{\left[(x - x') \hat{x} + (y - y') \hat{y} + (z - z') \hat{z}\right]}{\sqrt{(x - x')^2 + (y - y')^2 + (z - z')^2}} \\ &= \tilde{k} \hat{\mathbf{z}} \exp{(k x)} \end{split}$$

(d) Proof.

$$\begin{split} &= \nabla \left( 1/\imath \right) \\ &= \nabla \left( \left( x - x' \right)^2 + \left( y - y' \right)^2 + \left( z - z' \right)^2 \right)^{-1/2} \\ &= -\frac{1}{2} \left( \left( x - x' \right)^2 + \left( y - y' \right)^2 + \left( z - z' \right)^2 \right)^{-3/2} \left( \cancel{2} \left( x - x' \right) \hat{x} + \cancel{2} \left( y - y' \right) \hat{y} + \cancel{2} \left( z - z' \right) \hat{z} \right) \quad \text{By chain rule...} \\ &= -\frac{\mathbf{\hat{z}}}{\imath^3} = -\frac{\hat{\mathbf{\hat{z}}}}{\imath^2} \end{split}$$

**Problem 2.** For each of the vector fields illustrated below, find an equation that *could* describe the field and calculate its divergence and curl.

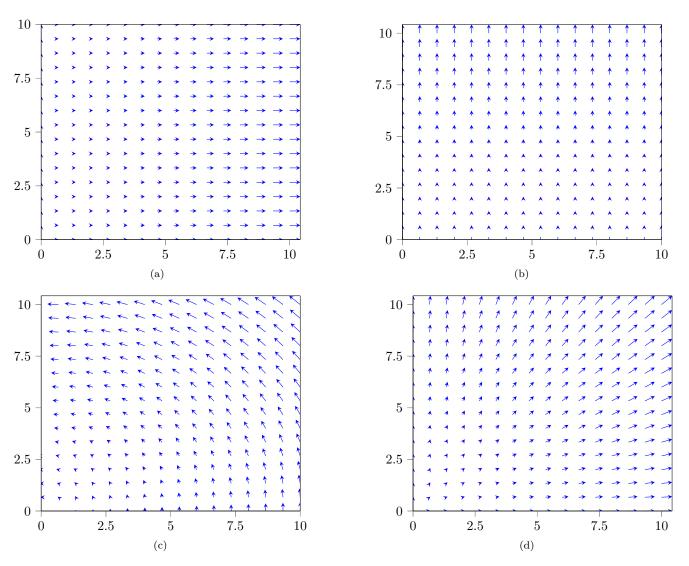


Figure 1: Plots generated with TikZ

Solution 2. (a)  $\mathbf{F} = x\hat{x} + 0\hat{y} + 0\hat{z}$ .  $\operatorname{div}(\mathbf{F}) = \vec{\nabla} \cdot \mathbf{F} = 1$ .

$$\operatorname{curl}(\mathbf{F}) = \vec{\nabla} \times \mathbf{F}$$

$$= \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ x & 0 & 0 \end{vmatrix}$$

$$= 0\hat{x} + 0\hat{y} + 0\hat{z}$$

(b) 
$$\mathbf{F} = 0\hat{x} + y\hat{y} + 0\hat{z}$$
.  $\operatorname{div}(\mathbf{F}) = \vec{\nabla} \cdot \mathbf{F} = 1$ 

$$\operatorname{curl}(\mathbf{F}) = \vec{\nabla} \times \mathbf{F}$$

$$= \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ 0 & y & 0 \end{vmatrix}$$

$$= 0\hat{x} + 0\hat{y} + 0\hat{z}$$

(c) 
$$\mathbf{F} = -y\hat{x} + x\hat{y} + 0\hat{z}$$
.  $\operatorname{div}(\mathbf{F}) = \vec{\nabla} \cdot \mathbf{F} = 0$ 

$$\operatorname{curl}(\mathbf{F}) = \vec{\nabla} \times \mathbf{F}$$

$$= \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ -y & x & 0 \end{vmatrix}$$

$$= 0\hat{x} + 0\hat{y} + 2\hat{z}$$

(d) 
$$\mathbf{F} = x\hat{x} + y\hat{y} + 0\hat{z}$$
.  $\operatorname{div}(\mathbf{F}) = \vec{\nabla} \cdot \mathbf{F} = 2$ 

$$\operatorname{curl}(\mathbf{F}) = \vec{\nabla} \times \mathbf{F}$$

$$= \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ -y & x & 0 \end{vmatrix}$$

$$= 0\hat{x} + 0\hat{y} + 0\hat{z}$$

**Problem 3.** Find the electric field a distance z above the centre of a flat circular disk of radius R that carries a uniform charge density  $\sigma$ . Complete all the steps we discussed in class:

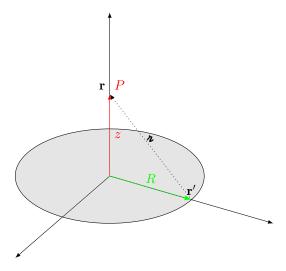
- (a) State which equation you will use and draw a diagram that illustrates the variables used in the equation.
- (b) Find expressions for each of the variables (r, r', etc.) that are needed to evaluate the equation for the field.
- (c) Solve the equation for the field.
  - Be careful to distinguish between z > 0 and z < 0.
  - Check that your answer satisfies the boundary conditions given by Eq. (2.33) in Griffiths.
- (d) Check that your answer makes sense in the limit  $z \gg R$ , where it should look like a point charge.

### Solution 3.

(a) For this we will use the surface charge equation,

$$\mathbf{E}(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \int \frac{\sigma(\mathbf{r}')}{\imath^2} \hat{\mathbf{i}} da'.$$

We will integrate over the disk which will probably require a change to polar coordinates but means that a' represents the area of the circle.



(b) r is the vector to the field "test point", labeled P in the figure. It's fairly obvious that it is (0,0,z). r' is also fairly obviously (x,y,0) as it is every point s.t.  $x^2+y^2\leq R^2$  (this is also our bound of integration until we make a change to polar coordinates). This means that  $\mathbf{a}=(-x,-y,z)$ 

(c) 
$$= \frac{1}{4\pi\epsilon_0} \int \frac{\sigma(\mathbf{r}')}{\imath^2} \hat{\boldsymbol{a}} da'$$

$$= \frac{\sigma}{4\pi\epsilon_0} \int \frac{\boldsymbol{a}}{\imath^3} da'$$

$$= \frac{\sigma}{4\pi\epsilon_0} \int \frac{(-x, -y, z)}{(x^2 + y^2 + z^2)^{3/2}} da'$$

$$= \frac{\sigma}{4\pi\epsilon_0} \int \frac{(-x, -y, z)}{(x^2 + y^2 + z^2)^{3/2}} da' \rightsquigarrow x = r' \cos \theta, y = r' \sin \theta$$

$$= \frac{\sigma}{4\pi\epsilon_0} \int_0^{2\pi} \int_0^R \frac{(-r \cos \theta, -r \sin \theta, z)}{(r'^2 \cos^2 \theta + r'^2 \sin^2 \theta + z^2)^{3/2}} r' dr' d\theta \rightsquigarrow x = r' \cos \theta, y = r' \sin \theta$$

$$= \frac{\sigma}{4\pi\epsilon_0} \int_0^{2\pi} \int_0^R \frac{(-r \cos \theta, -r \sin \theta, z)}{(r'^2 + z^2)^{3/2}} r' dr' d\theta \quad \text{Integrating w.r.t } \theta \text{ kills the sin and cos by symmetry}$$

$$= \frac{\sigma}{4\pi\epsilon_0} \int_{0=r'}^R \frac{(0, 0, 2\pi z)}{2u^{3/2}} du \rightsquigarrow u = r'^2 + z^2 \implies du/2r' = dr'$$

$$= \frac{\sigma}{4\epsilon_0} \int_{0=r'}^R \frac{(0, 0, z)}{u^{3/2}} du$$

$$= -\frac{z\sigma}{2\epsilon_0} \left[ (r'^2 + z^2)^{-1/2} \right]_0^R \hat{z} = \frac{z\sigma}{2\epsilon_0} \left[ \frac{1}{|z|} - \frac{1}{\sqrt{R^2 + z^2}} \right] \hat{z}$$

For boundary conditions

### Why do I have to discard the abs

$$\mathbf{E}_{above} - \mathbf{E}_{below} = \frac{z\sigma}{2\epsilon_0} \left[ \frac{1}{z} - \frac{1}{\sqrt{R^2 + z^2}} + \frac{1}{z} + \frac{1}{\sqrt{R^2 + z^2}} \right] \hat{z}$$

$$= \frac{z\sigma}{2\epsilon_0} \left[ \frac{2}{z} \right] \hat{z}$$

$$= \frac{\sigma}{\epsilon_0} \hat{z}$$

Which satisfies equation 2.33 as  $\hat{z}$  here is our normal vector.

(d) We first begin with some reworking using the binomial approximation given by  $(1+x)^n \approx 1 + xn$ 

$$\begin{split} &= \frac{z\sigma}{2\epsilon_0} \left[ \frac{1}{z} - \frac{1}{\sqrt{R^2 + z^2}} \right] \hat{z} \\ &= \frac{\sigma}{2\epsilon_0} \left[ 1 - 1 - \left[ 1 + \left( \frac{R}{z} \right)^2 \right]^{-1/2} \right] \hat{z} \\ &\approx \frac{\sigma}{2\epsilon_0} \left[ \frac{R^2}{2z^2} \right] \hat{z} \\ &= \frac{\sigma R^2}{4\epsilon_0 z^2} \hat{z} \end{split}$$

We then recognize that  $\sigma = Q/\pi R^2$  because we have a uniform charge distribution of net charge Q spread over a disk of area  $\pi R^2$ . Substituting this in,

$$\frac{QR^2}{4\epsilon_0\pi R^2z^2}\hat{z} = \frac{Q}{4\pi\epsilon_0z^2}\hat{z}$$

Which is the equation for a point charge, as we would hope.

**Problem 4.** For a solid sphere of radius R and charge density  $\rho(\mathbf{r}) = \kappa r$ , use Gauss' law to find

- (i) The electric field.
- (ii) The potential inside and outside the sphere.

#### Solution 4.

(i) Gauss' law states that

$$\oint_{S} \mathbf{E} \cdot d\mathbf{a} = \frac{Q_{enc}}{\epsilon_0}$$

Where  $Q_{enc}$  is given by

$$Q_{enc} = \int_{\mathcal{V}} \rho \, d\tau.$$

Because we have a sphere of charge we can transform this to use spherical coordinates,

$$Q_{enc} = \int_0^R dr \int_0^{\pi} d\theta \int_0^{2\pi} d\phi \, kr^3 \sin\theta$$
$$= 4\pi k \int_0^R r^3 dr$$
$$= \pi k R^4$$

We can now return to our original statement involving  $\mathbf{E}$  and solve. We can also observe that because  $\mathbf{E}$  points radially outwards from our sphere and  $\mathbf{a}$  does as well we can reduce the dot product to a product of their magnitudes.

$$\frac{\pi k R^4}{\epsilon_0} = \oint_S |\mathbf{E}| da$$

Now because  $|\mathbf{E}|$  is constant over the *surface* because the charge density is spherically symmetric, only depending on r, we can extract  $|\mathbf{E}|$  from the integral which then just becomes the area of a sphere,

$$\begin{split} \frac{\pi k R^4}{\epsilon_0} &= \oint_S |\mathbf{E}| da \\ \frac{\pi k R^4}{\epsilon_0} &= |\mathbf{E}| \oint_S da \\ \frac{\pi k R^4}{\epsilon_0} &= 4\pi r^2 |\mathbf{E}| \\ \frac{k R^4}{4\epsilon_0 r^2} &= |\mathbf{E}| \end{split}$$

(For r > R)

(ii) To start here we need the electric field inside the sphere which is given by the same process as above except we integrate from 0 to  $r \leq R$  because we are now inside the sphere. This gives  $Q_e nc = \pi k r^4 \implies \mathbf{E} = \frac{kr^2}{4\epsilon_0}\hat{\mathbf{r}}$ . We can now evaluate for the outside of the sphere,

$$V_{r>R}(r) = -\int_{\infty}^{r} \frac{kR^4}{4\epsilon_0 r'^2} dr'$$
$$= -\int_{\infty}^{r} \frac{kR^4}{4\epsilon_0 r'^2} dr'$$
$$= \frac{kR^4}{4\epsilon_0} \left[ \frac{1}{r'} \right]_{\infty}^{r} = \frac{kR^4}{4\epsilon_0 r}$$

Then for the potential inside the sphere we split the integral.

I'm not sure why we do this but it's done in the text

$$V_{R < r}(r) = -\int_{\infty}^{R} \frac{kR^4}{4\epsilon_0 r'^2} dr' - \int_{R}^{r} \frac{kr^2}{4\epsilon_0} dr'$$
$$= \frac{kR^4}{4\epsilon_0 R} - \frac{k}{12\epsilon_0} \left[r'^3\right]_{R}^{r}$$
$$= \frac{k}{12\epsilon_0} \left(4R^3 - r^3\right)$$

## Check both of these w/ $E = -\nabla V$

**Problem 5.** Consider the electric field

$$\mathbf{E}(\mathbf{r}) = E_0 \left[ y^2 \hat{x} + \left( 2xy + z^2 \right) \hat{y} + 2yz\hat{z} \right].$$

Find, by integrating  $-\int \mathbf{E} \cdot d\boldsymbol{\ell}$ , the potential at an arbitrary point  $\mathbf{r}$ , taking the origin as the reference point. Note that the most direct integration path may not be simplest.

**Solution 5.** Here we have that  $d\ell = dx\hat{x} + dy\hat{y} + dz\hat{z}$  which means that  $\mathbf{E} \cdot d\ell = E_0 \left[ y^2 dx + \left( 2xy + z^2 \right) dy + 2yz dz \right]$  so,

$$V(\mathbf{r}) = \int_0^{\mathbf{r}} E_0 \left[ y^2 dx + (2xy + z^2) dy + 2yz dz \right]$$
  
=  $E_0 \left[ \int_0^{r_x} y^2 dx + \int_0^{r_y} (2xy + z^2) dy + \int_0^{r_z} 2yz dz \right]$ 

Which simplifies nicely along several paths (though it's already fairly simple). I'll go along the x, the y, then the z. This means that for our first integral along x we are moving from  $\langle 0,0,0 \rangle$  to  $\langle r_x,0,0 \rangle$  so y is zero along the whole range which means that the integral just becomes zero as its integrand is  $y^2$ . Next along y we are moving from  $\langle r_x,0,0 \rangle$  to  $\langle r_x,r_y,0 \rangle$  so

$$\int_0^{r_y} \left(2xy + \mathbf{z}^2\right)^0 dy = r_x r_y^2.$$

Finally along z we are moving from  $\langle r_x, r_y, 0 \rangle$  to  $\langle r_x, r_y, r_z \rangle$  so

$$\int_{0}^{r_{z}} 2yzdz = \int_{0}^{r_{z}} 2yzdz = r_{y}r_{z}^{2}.$$

This yields

$$V(\mathbf{r}) = E_0 \left[ r_x r_y^2 + r_y r_z^2 \right].$$