

Magnetic Resonance Current Density Imaging

a BioTechMed-Graz Lab Rotation Report

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1 Summary

This lab rotation project was concerned with the development, exploration and implementation of a new optimisation approach for the reconstruction of current density / conductivity of tissue within a Magnetic Resonance Imaging (MRI) setup.

The project was split into two steps: First, solving the forward problem, that is simulating the magnetic field modulation given a current density field and testing the forward solver on a *phantom*, cf. Section 3.1. And second, solving the inverse problem, that is reconstructing current density (including conductivity) using a novel optimisation model, based on measurements of the z-component of the magnetic field, with insight gained from the forward procedure, cf. Section 3.2.

The implementation was built on top of KomaMRI (Castillo-Passi et al. 2022) and is freely available on GitHub (<https://github.com/MrP01/CurrentDensityImaging>). An exemplary reconstruction is documented in the results Section 4.

2 Introduction

The overall intention of MRI is to turn a source of contrast into an image for clinicians to use, mainly for the identification of different (potentially malignant) tissue types. In the usual setting, this source of contrast is the proton density or one of the T_1 , T_2 or T_2^* material constants of individual magnetic dipoles, such as water. Visualising these material properties in an image therefore allows to differentiate between different types of tissue.

2.1 Brief Introduction to Magnetic Resonance Imaging

What we image specifically are magnetic dipoles, also referred to as *spins* in this context. The general equation governing the behaviour of these *spin* objects is the Bloch equation:

$$\frac{d\mathbf{M}}{dt} = \gamma(\mathbf{M} \times \mathbf{B}) - \frac{(M_3 - M_0)\hat{\mathbf{e}}_3}{T_1} - \frac{M_1\hat{\mathbf{e}}_1 + M_2\hat{\mathbf{e}}_2}{T_2}, \quad (1)$$

where $\mathbf{M} = (M_1, M_2, M_3)^T \in \mathbb{R}^3$ is the magnetisation, $\mathbf{B} \in \mathbb{R}^3$ the magnetic field, $M_0 \in \mathbb{R}$ the equilibrium magnetisation, γ the gyromagnetic ratio, T_1 and $T_2 \in \mathbb{R}$ the (*relaxation time*) material constants mentioned above and $t \in \mathbb{R}$ being time (cf. Nishimura 1996). It describes the precession of the magnetisation \mathbf{M} “around” the magnetic field \mathbf{B} . Supplied with the appropriate initial condition (e.g. $\mathbf{M}(t = t_0) = \mathbf{M}_{t_0}$), this system can be solved explicitly.

Starting from a certain MRI sequence controlling the gradient coils to excite spins inside the object of interest (individually governed by Equation (1)), the imaging process generally travels through a grid of points in \mathbf{k} -space (Fourier space) to later reconstruct an image slice using the inverse Fourier transform. This procedure is repeated for every slice of the object until a full image is reconstructed (cf. Nishimura 1996).

For the purposes of Magnetic Resonance Current Density Imaging (MRCDI), the procedure also yields the z-component of the magnetic field $B_z = B_3 = \{\mathbf{B}\}_3$, which we will use for the reconstruction of current density.

2.2 The Biot-Savart Law

For a given current density $\mathbf{j}(t) : \Omega \rightarrow \mathbb{R}^3$ ⁽¹⁾ in \mathbb{R}^3 , with $\Omega \subseteq \mathbb{R}^3$ denoting the region of interest, at time $t \in \mathbb{R}^+$, Maxwell's fourth equation in differential form,

$$\nabla \times \mathbf{B} = \mu_0 \left(\mathbf{j} + \varepsilon_0 \frac{\partial \mathbf{E}}{\partial t} \right),$$

relates the curl of the magnetic field $\mathbf{B}(t) : \Omega \rightarrow \mathbb{R}^3$ to the current density $\mathbf{j}(t)$ and the temporal rate of change in the *electric* field $\mathbf{E}(t) : \Omega \rightarrow \mathbb{R}^3$ (Jackson and Fox 1999, chapter 6, section 1). The constants ε_0 and μ_0 are the electric permittivity and magnetic permeability of free space, respectively. The current density \mathbf{j} is also connected to the electric field \mathbf{E} through the, also position-dependent, electrical conductivity $\sigma : \Omega \rightarrow \mathbb{R}^+$,

$$\mathbf{j} = \sigma \mathbf{E}.$$

In the **electrostatic case** (when the electric field \mathbf{E} is independent of time t), the last term in the above equation will vanish and we arrive at

$$\nabla \times \mathbf{B} = \mu_0 \mathbf{j}. \quad (2)$$

This equation, together with the freedom of divergence of the magnetic field (Maxwell's second equation), $\nabla \cdot \mathbf{B} = 0$, leads to the **Biot-Savart law**,

$$\mathbf{B}(\mathbf{x}) = \frac{\mu_0}{4\pi} \int_{\Omega} \frac{\mathbf{j}(\mathbf{y}) \times (\mathbf{x} - \mathbf{y})}{\|\mathbf{x} - \mathbf{y}\|_2^3} d\mathbf{y} \quad (3)$$

which provides an explicit expression for the magnetic field contribution $\mathbf{B}(\mathbf{x})$ at position $\mathbf{x} \in \Omega$ given a current density field \mathbf{j} (Jackson and Fox 1999, chapter 5, section 2).

2.3 Current Density Imaging

Using a current source and attaching it to two electrodes connected to the object of interest, while reconstructing the image inside the Magnetic Resonance Tomography (MRT) tube, the *imaging current* will induce a magnetic field modulation on top of the magnetic field already present inside the tube according to the Biot-Savart law, which we can measure (the B_z component), cf. Ider and Onart 2004.

As we only have access to one component of the magnetic field, the resulting system is underdetermined and we therefore need to solve an ill-posed *inverse problem*.

3 Methods

For visualisation and as a general framework, we built upon the **KomaMRI** toolchain and Julia software package (Castillo-Passi et al. 2022). This choice was made mostly due to Koma's feature-richness and extensibility, while harnessing the computational power of Julia (Bezanson et al. 2017).

3.1 FFT-accelerated evaluation of the Biot-Savart law

Starting from the Biot-Savart law introduced above (cf. Equation (3)), we observe that splitting the cross-product (\times) into its three respective components

$$\begin{aligned} B_1(\mathbf{x}) &= \frac{\mu_0}{4\pi} \int_{\Omega} \frac{j_2(\mathbf{y}) \cdot (x_3 - y_3) - j_3(\mathbf{y}) \cdot (x_2 - y_2)}{\|\mathbf{x} - \mathbf{y}\|_2^3} d\mathbf{y}, \\ B_2(\mathbf{x}) &= \frac{\mu_0}{4\pi} \int_{\Omega} \frac{j_3(\mathbf{y}) \cdot (x_1 - y_1) - j_1(\mathbf{y}) \cdot (x_3 - y_3)}{\|\mathbf{x} - \mathbf{y}\|_2^3} d\mathbf{y}, \\ B_3(\mathbf{x}) &= \frac{\mu_0}{4\pi} \int_{\Omega} \frac{j_1(\mathbf{y}) \cdot (x_2 - y_2) - j_2(\mathbf{y}) \cdot (x_1 - y_1)}{\|\mathbf{x} - \mathbf{y}\|_2^3} d\mathbf{y}, \end{aligned}$$

¹The current density $\mathbf{j} = nq\mathbf{v}$ describes the flow of n charges q with velocity \mathbf{v} , it is therefore the amount of current I flowing through an infinitesimal area A , motivating its notation as a *density*.

according to the explicit representation of the cross-product in three spatial dimensions, most notably allows us to express B_1 , B_2 and B_3 as convolution integrals with $g_1, g_2, g_3 : \Omega \rightarrow \mathbb{R}$,

$$B_1(\mathbf{x}) = \frac{\mu_0}{4\pi} \int_{\Omega} [j_2(\mathbf{y})g_3(\mathbf{x} - \mathbf{y}) - j_3(\mathbf{y})g_2(\mathbf{x} - \mathbf{y})] d\mathbf{y} = \frac{\mu_0}{4\pi} [(j_2 *_{\Omega} g_3) - (j_3 *_{\Omega} g_2)](\mathbf{x}),$$

$$\text{where } g_1(\mathbf{x}) := \mathbf{x}_1 / \|\mathbf{x}\|_2^3, \quad g_2(\mathbf{x}) := \mathbf{x}_2 / \|\mathbf{x}\|_2^3 \quad \text{and} \quad g_3(\mathbf{x}) := \mathbf{x}_3 / \|\mathbf{x}\|_2^3$$

and the respective analogs for B_2 and B_3 (Yazdanian et al. 2020).

Using the convolution theorem for two functions f and g

$$\mathcal{F}[f * g](\mathbf{k}) = \mathcal{F}[f](\mathbf{k}) \cdot \mathcal{F}[g](\mathbf{k}) \quad \text{where} \quad \mathcal{F}[f](\mathbf{k}) := \int_{\mathbb{R}^3} f(\mathbf{x}) e^{-i\mathbf{k}^T \mathbf{x}} d\mathbf{x},$$

the respective convolution integrals may be evaluated computationally efficiently using the Fast Fourier Transform (FFT) and IFFT (for details on this relationship, we refer the reader to Herman 2022). More precisely, this speedup is due to the computational complexity cost reduction from $\mathcal{O}(n^2)$ when evaluating the Biot-Savart law conventionally to $\mathcal{O}(n \log n)$ when using the FFT, where n would model the number of voxels the system operates on.

The first component of the magnetic field B_1 may then be expressed as

$$B_1(\mathbf{x}) = \frac{\mu_0}{4\pi} \mathcal{F}^{-1} [\mathcal{F}(j_2) \cdot \mathcal{F}(g_3) - \mathcal{F}(j_3) \cdot \mathcal{F}(g_2)](\mathbf{x}),$$

and analogous expressions may be found for the second and third component B_2 and B_3 , respectively. Importantly, we can explicitly evaluate

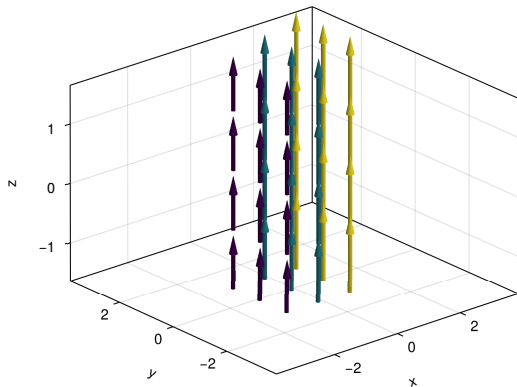
$$\mathcal{F}[g_n](\mathbf{k}) = \mathcal{F} \left[\mathbf{x} \mapsto \frac{x_n}{\|\mathbf{x}\|_2^3} \right](\mathbf{k}) = -4\pi i \frac{k_n}{\|\mathbf{k}\|_2^2}, \quad \text{for } n = 1, 2, 3,$$

where i is the imaginary unit, which allows for a direct evaluation of the above using only the FFT, IFFT, addition and multiplication (Yazdanian et al. 2020).

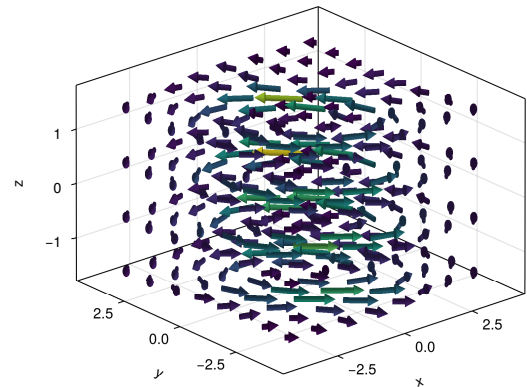
In code, this may be implemented like so:

```
1 function calculate_magnetic_field(cdp::CurrentDensityPhantom)::VectorField
2   # [obtain M_range, calculate g1, g2, g3 on the entire frequency domain]
3   c1, c2, c3 = cross(fft(pad_jx), fft(pad_jy), fft(pad_jz), g1, g2, g3)
4   B1, B2, B3 = real(ifft(c1)), real(ifft(c2)), real(ifft(c3))
5   return mu_0 .* (B1[M_range...], B2[M_range...], B3[M_range...])
6 end
```

For a current density field as given in Figure 1a, the resulting \mathbf{B} field looks as presented in Figure 1b.



(a) Current density $\mathbf{j}(\mathbf{x}) = j_0(-1, -1, 2x_1)^T \mathbf{1}_{\mathbf{x} \in \Omega_s}$ on a sample rectangular domain.



(b) Magnetic field $\mathbf{B}(\mathbf{x})$ resulting from the current density field depicted in Figure 1a.

Figure 1: Current density field \mathbf{j} and the corresponding magnetic field \mathbf{B} as obtained using the Biot-Savart law.

3.1.1 Gridded Data Format

In order to use KomaMRI together with the FFT-based evaluation of the Biot-Savart law (described above) on a phantom, we had to adapt the storage format from Koma's flattened vector storage of coordinates, density and other characteristic constants to a grid-based approach. The flattened storage format allows for higher efficiency in terms of memory usage, especially for less dense phantoms. However, some calculations on the data require neighbouring spins' characteristics in all six directions, for which a three-dimensional tensor storage format is more suitable, as it is provided in, for example, the JEMRIS file format (arrays are stored in `.h5` files). This is especially true for the FFT, where data must be provided in gridded format.

3.2 The Optimisation Approach

In order to be able to reconstruct the current density vector field \mathbf{j} and conductivity scalar field σ , we employed the following optimisation model:

$$\begin{aligned} \mathbf{B}^*, \sigma^* &:= \arg \min_{\mathbf{B}, \sigma} && \frac{1}{2} \int_{\Omega} \|\{\mathbf{B}\}_3 - \mathbf{B}_3^0\|_2^2 \, d\mathbf{x} + \frac{\alpha}{2} \int_{\Omega} \frac{\|\nabla \times \mathbf{B}(\mathbf{x})\|_2^2}{\sigma(\mathbf{x})} \, d\mathbf{x} + R(\sigma), \\ &\text{subject to} && \nabla \cdot \mathbf{B} = \text{div } \mathbf{B} = 0, \\ &\text{and} && \sigma \in [\sigma_0, \sigma_1], \end{aligned}$$

As a regulariser R for the conductivity σ , we used the Total Variation (TV), 1-norm of the gradient, of σ :

$$R(\sigma) := \beta \int_{\Omega} \|\text{grad } \sigma(\mathbf{x})\|_1 \, d\mathbf{x} = \beta \int_{\Omega} \left(\left| \frac{\partial \sigma}{\partial x_1} \right| + \left| \frac{\partial \sigma}{\partial x_2} \right| + \left| \frac{\partial \sigma}{\partial x_3} \right| \right) \, dx_1 dx_2 dx_3$$

One can include the divergence constraint as a penalty term to the objective function:

$$p_{\text{div}} := \gamma \int_{\Omega} |\text{div } \mathbf{B}(\mathbf{x})|^2 \, d\mathbf{x} = \gamma \int_{\Omega} \left| \frac{\partial B_1}{\partial x_1} + \frac{\partial B_2}{\partial x_2} + \frac{\partial B_3}{\partial x_3} \right|^2 \, dx_1 dx_2 dx_3$$

with all constants $\alpha, \beta, \gamma \in \mathbb{R}^+$ sufficiently large. Both the divergence and gradient were implemented using a numerical central difference scheme.

Using Maxwell's second equation, one can obtain the corresponding current density \mathbf{j}^* from Maxwell's fourth equation in the electrostatic case (Equation (2)),

$$\mathbf{j}^* := \frac{1}{\mu_0} (\nabla \times \mathbf{B}^*) = \frac{1}{\mu_0} \text{curl}(\mathbf{B}^*),$$

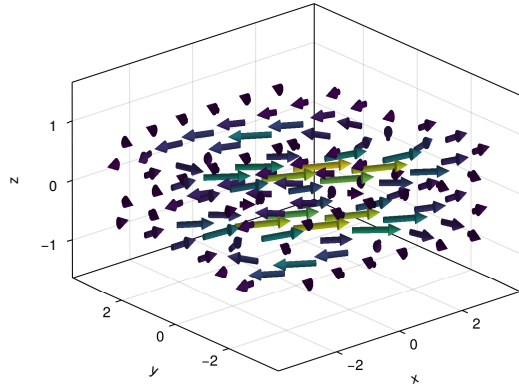
where we evaluated the `curl` using a central finite difference scheme (setting boundary values to 0).

Dimensional Analysis In the objective function, note the unit of the integral kernel $\|\nabla \times \mathbf{B}(\mathbf{x})\|_2^2 / \sigma(\mathbf{x})$: the numerator is the squared norm of the current density, therefore in A^2/m^4 . Together with the corresponding denominator (in units of S m^{-1}), the integral kernel has units of $\text{A}^2 \Omega / \text{m}^3$, becoming $\text{A}^2 \Omega$, which can be simplified to V A , which is a unit of **power** ($\text{V A} = \text{W}$). The entire term may then be interpreted as a form of power dissipation and should therefore be minimised.

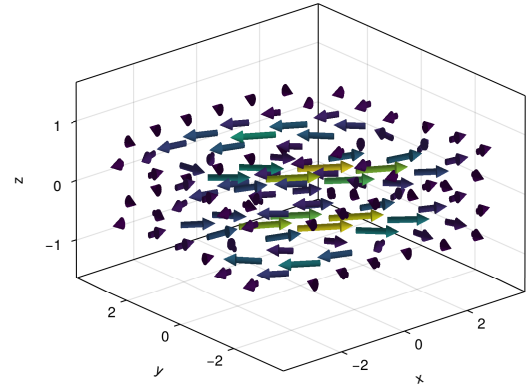
The iterates were produced using the Optim.jl optimisation library (Mogensen and Riseth 2018) using the LBFGS routine (Liu and Nocedal 1989).

4 Results

Starting from an exemplary current density phantom (cf. Figure 2a), the reconstruction routine converges after 40 LBFGS iterations, taking 20s in total. The reconstructed current density may be found to its right, in Figure 2b. The corresponding reconstructed conductivity and magnetic field, along with their originals can be found in Figures 3 and 4.

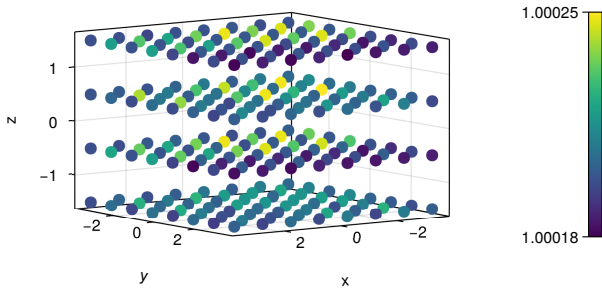


(a) Exemplary current density field $\mathbf{j}(\mathbf{x})$ on a sample rectangular domain.

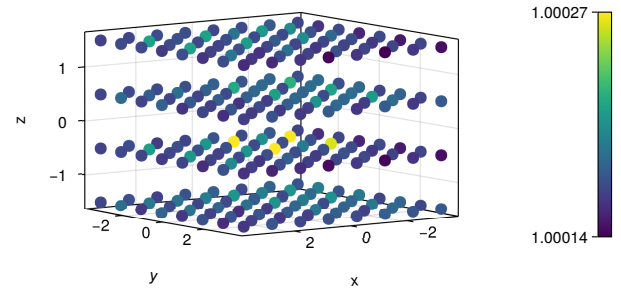


(b) The *reconstructed* current density field $\mathbf{j}^*(\mathbf{x})$ using the method described in Section 3.2.

Figure 2: Original current density field \mathbf{j} (left) and its reconstruction \mathbf{j}^* (right).

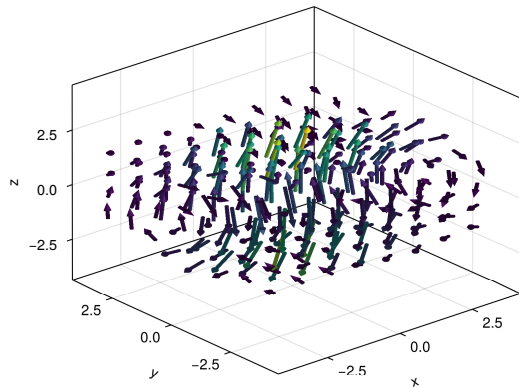


(a) Conductivity σ of the original, obtained as the best possible combination with \mathbf{j} and \mathbf{B} using the optimisation described in Section 3.2.

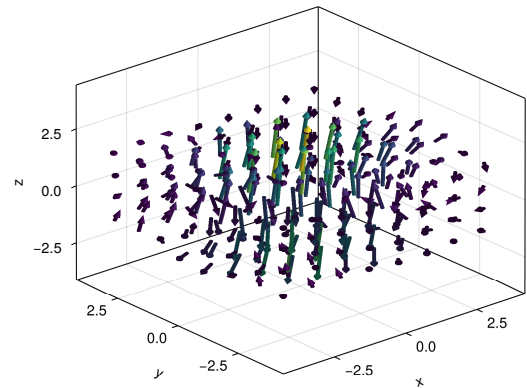


(b) Reconstructed conductivity field σ , marking the solution of the optimisation problem together with Figures 2b and 4b.

Figure 3: Original conductivity σ (left) and its reconstruction σ^* (right).



(a) Magnetic field corresponding to the current density field in Figure 2a, as obtained through the Biot-Savart law.



(b) Magnetic field corresponding to the *reconstructed* current density Figure 2b, as obtained through the Biot-Savart law.

Figure 4: Original magnetic field \mathbf{B} (left) and its reconstruction \mathbf{B}^* (right).

Within the scope of this initial exploration, the current density reconstruction *from only the z-component of the magnetic field* matches the phantom well, along with the conductivity and associated magnetic field (the actual variables of the optimisation routine). The reconstructed conductivity also aligns well with its original, varying only slightly in magnitude while retaining most spatial features.

Further experiments (which this report is too short to contain) suggest it extends well to other current density fields, as long as they are solenoidal (divergence free). Starting from a general vector field, its solenoidal part may be obtained using the Helmholtz decomposition, as long as it decays rapidly enough towards the boundaries of the domain.

4.1 Usage as a Julia Package

The package provided alongside this report offers two abstractions of Koma's `phantom` class:

1. `PhantomOnAGrid`, implementing the conversion and interfacing to the gridded data format.
2. `CurrentDensityPhantom`, extending `PhantomOnAGrid` with the current density vector field (each spatial component is stored in its own respective tensor).

In order to install and try out the package, open up a Julia live interpreter and type:

```
1 import Pkg; Pkg.add(url="https://github.com/MrP01/CurrentDensityImaging");
2 import CurrentDensityImaging as CDI # import the package
3 cdp = CDI.generateDemoCDP() # generate exemplary current density phantom
```

With the constructed current density phantom (`cdp`), one can generate plots for \mathbf{j} , \mathbf{B} and σ using `CDI.plot_current_density`, `CDI.plot_magnetic_field` and `CDI.plot_conductivity`, respectively. It is also possible to load a phantom in the JEMRIS format, for more information we refer the reader to the [code repository](#), specifically the `src/grid.jl` file.

The forward operator (FFT-accelerated evaluation of the Biot-Savart equation) can be invoked as follows:

```
1 B1, B2, B3 = CDI.calculate_magnetic_field(cdp); # obtain the x, y and z components of B
```

A reconstruction from the z-component of the magnetic field B_z can be obtained using

```
1 cdpr, sigmar = CDI.solve(B3); # reconstruction of j and sigma
```

where `cdpr` contains the reconstructed current density phantom and `σr` is the reconstructed conductivity.

5 Future Perspectives

MRCDI is a promising approach toward the better diagnosis of pathologies in patients using imaging techniques. The given model can be extended in various ways, for example one could modify the regularisation imposed on the conductivity, or add further terms to generally improve the reconstruction of \mathbf{j} . It might also be possible to relax the divergence freeness condition in the optimisation procedure, by reducing redundancy in the optimised variable space.

Signatures

This project was carried out as part of a [BioTechMed-Graz Lab Rotation](#). The aim of the programme is to give interested graduates a chance to broaden their scientific interests beyond the topic of their master's thesis.

(Peter Waldert)

Place and Date

(Prof. Kristian Bredies)

Place and Date

References

- Castillo-Passi, Carlos, Ronal Coronado, Gabriel Varela-Mattatall, Carlos Alberola-López, René Botnar and Pablo Irarrazaval (2022). ‘KomaMRI.jl: An open-source framework for general MRI simulations with GPU acceleration’. In: *Magnetic Resonance in Medicine*. DOI: [10.1002/mrm.29635](https://doi.org/10.1002/mrm.29635).
- Nishimura, Dwight George (1996). *Principles of Magnetic Resonance Imaging*. Stanford, CA, USA: Stanford University. URL: https://books.google.at/books/about/Principles_of_magnetic_resonance_imaging.html?id=uz9BAQAAIAAJ&redir_esc=y.
- Jackson, Jon David and Ronald Forrest Fox (1999). ‘Classical Electrodynamics, 3rd ed.’ In: *American Journal of Physics* 67, pp. 841–842. URL: <https://api.semanticscholar.org/CorpusID:123039708>.
- Ider, Y. Ziya and Serkan Onart (Feb. 2004). ‘Algebraic reconstruction for 3D magnetic resonance–electrical impedance tomography (MREIT) using one component of magnetic flux density’. In: *Physiol. Meas.* 25.1, p. 281. ISSN: 0967-3334. DOI: [10.1088/0967-3334/25/1/032](https://doi.org/10.1088/0967-3334/25/1/032).
- Bezanson, Jeff, Alan Edelman, Stefan Karpinski and Viral B Shah (2017). ‘Julia: A fresh approach to numerical computing’. In: *SIAM Review* 59.1, pp. 65–98. DOI: [10.1137/141000671](https://doi.org/10.1137/141000671).
- Yazdani, Hassan, Guilherme B. Saturnino, Axel Thielscher and Kim Knudsen (June 2020). ‘Fast evaluation of the Biot-Savart integral using FFT for electrical conductivity imaging’. In: *J. Comput. Phys.* 411, p. 109408. ISSN: 0021-9991. DOI: [10.1016/j.jcp.2020.109408](https://doi.org/10.1016/j.jcp.2020.109408).
- Herman, Russell (July 2022). ‘9.9: The Convolution Theorem’. In: *Mathematics LibreTexts*. URL: [https://math.libretexts.org/Bookshelves/Differential_Equations/Introduction_to_Partial_Differential_Equations_\(Herman\)/09%3A_Transform_Techniques_in_Physics/9.09%3A_The_Convolution_Theorem](https://math.libretexts.org/Bookshelves/Differential_Equations/Introduction_to_Partial_Differential_Equations_(Herman)/09%3A_Transform_Techniques_in_Physics/9.09%3A_The_Convolution_Theorem).
- Mogensen, Patrick K. and Asbjørn N. Riseth (Apr. 2018). ‘Optim: A mathematical optimization package for Julia’. In: *Journal of Open Source Software* 3.24, p. 615. ISSN: 2475-9066. DOI: [10.21105/joss.00615](https://doi.org/10.21105/joss.00615).
- Liu, Dong C. and Jorge Nocedal (Aug. 1989). ‘On the limited memory BFGS method for large scale optimization’. In: *Math. Program.* 45.1, pp. 503–528. ISSN: 1436-4646. DOI: [10.1007/BF01589116](https://doi.org/10.1007/BF01589116).
- Danisch, Simon and Julius Krumbiegel (2021). ‘Makie.jl: Flexible high-performance data visualization for Julia’. In: *Journal of Open Source Software* 6.65, p. 3349. DOI: [10.21105/joss.03349](https://doi.org/10.21105/joss.03349).

Acronyms

FFT	Fast Fourier Transform	3
IFFT	Inverse Fast Fourier Transform	3
MRCDI	Magnetic Resonance Current Density Imaging	1
MRI	Magnetic Resonance Imaging	1
MRT	Magnetic Resonance Tomography	2
TV	Total Variation	4
UI	User Interface	9

A Appendix

The figures in this report were generated using Makie (Danisch and Krumbiegel 2021).

A.1 User Interface of KomaMRI

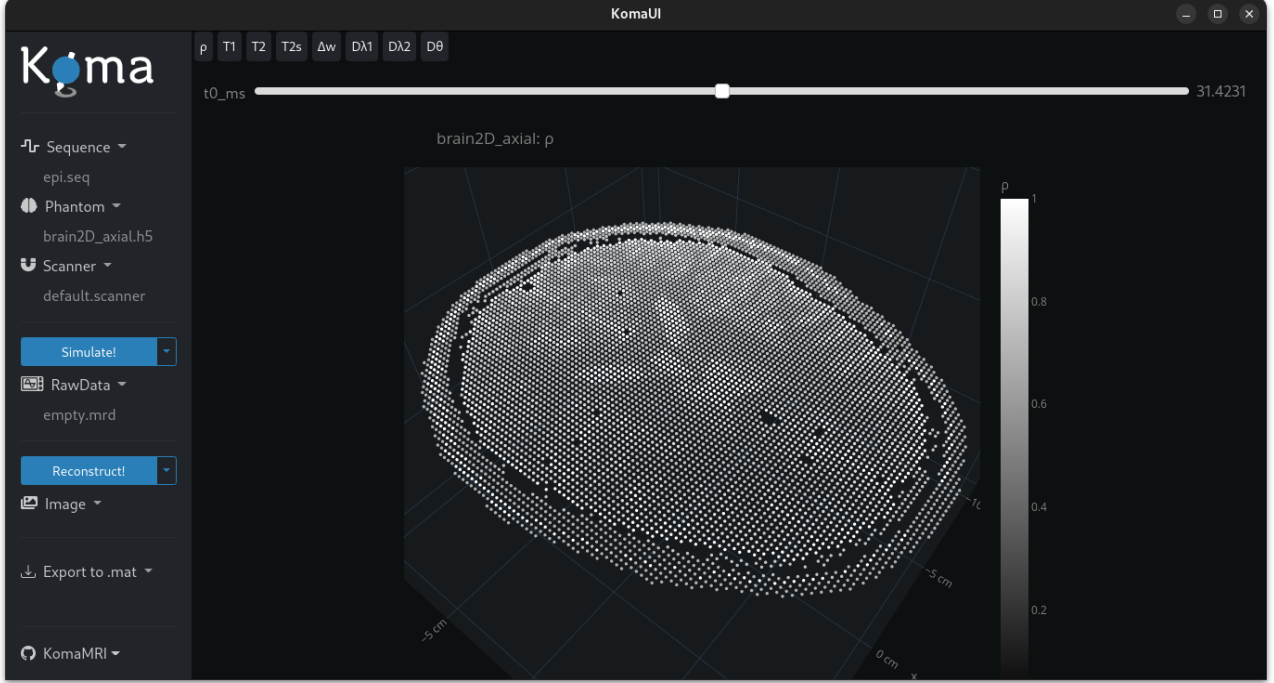


Figure 5: The User Interface (UI) packaged alongside Koma for users to explore the different functionality in a no-code environment. This interface can also be used to visualise results from the reconstruction mechanisms introduced in this report.

A.2 Reconstructed Magnetic Field

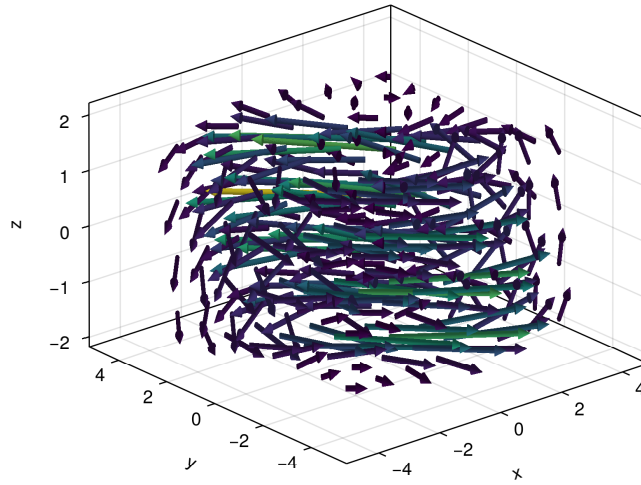


Figure 6: Magnetic field obtained from the Biot-Savart procedure of the [curl](#) of the B -field in Figure 1b.