Numerics of Partial Differential Equations: Stationary Problems $_{\rm Lecture\ Notes}$

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Chapter 1

Introduction

1.1 Strong Form and Variational Form

The finite element method is a scheme for the numerical solution of partial differential equations. In this chapter, we introduce the basic concepts for elliptic problems in the frame of the Riesz theorem. To that end, we consider the most standard example, namely the Poisson equation with mixed Dirichlet-Neumann boundary conditions. We aim to solve

$$-\Delta u = f \quad \text{in } \Omega,$$

$$u = 0 \quad \text{on } \Gamma_D,$$

$$\partial u/\partial n = \phi \quad \text{on } \Gamma_N,$$

$$(1.1)$$

which is said to be the **strong form** of the boundary value problem. Here, Ω denotes a domain in \mathbb{R}^d , d=2,3. The boundary $\Gamma:=\partial\Omega$ is split into the Dirichlet boundary Γ_D and the Neumann boundary Γ_N , respectively. To be more precise, we assume that Γ_D and Γ_N are (relatively) open subsets of Γ with $\Gamma_D \cap \Gamma_N = \emptyset$ and $\Gamma = \overline{\Gamma}_D \cup \overline{\Gamma}_N$. The source term $f:\Omega \to \mathbb{R}$ as well as the Neumann data $\phi:\Gamma_N \to \mathbb{R}$ are given, and $u:\Omega \to \mathbb{R}$ is the unknown solution. Moreover,

$$\Delta u(x) := \sum_{j=1}^{d} \frac{\partial^2 u}{\partial x_j^2}(x) \tag{1.2}$$

denotes the Laplace operator, which is defined in the classical sense for a function $u \in C^2(\overline{\Omega})$, where $C^k(\overline{\Omega}) := \{w|_{\overline{\Omega}} \mid w \in C^k(\mathbb{R}^d)\}$. If $u \in C^2(\overline{\Omega})$ solves (1.1), u is said to be a **strong solution** of the mixed boundary value problem.

Throughout the lecture, we shall assume that Ω is a **Lipschitz domain** in \mathbb{R}^d , i.e.,

- Ω is a bounded, open, and connected subset of \mathbb{R}^d ,
- Ω is locally on one side of Γ ,
- \bullet Γ can locally be parametrized by Lipschitz continuous functions.

An important consequence of this assumption is the validity of the integration by parts formula

$$\int_{\Omega} \frac{\partial u}{\partial x_j} v \, dx + \int_{\Omega} u \, \frac{\partial v}{\partial x_j} \, dx = \int_{\Gamma} u v n_j \, ds \quad \text{for all } u, v \in C^1(\overline{\Omega}), \tag{1.3}$$

where n_j denotes the j-th component of the outer normal vector of Ω on Γ and where ds denotes the surface measure on Γ . For a precise definition and details, we refer, e.g., to [McL].

Let $u \in C^2(\overline{\Omega})$ be a strong solution of (1.1) and $v \in C_D^1(\overline{\Omega}) := \{w \in C^1(\overline{\Omega}) \mid w|_{\Gamma_D} = 0\}$. Multiplication of $-\Delta u = f$ by v, integration over Ω , and integration by parts yield that

$$\int_{\Omega} f v \, dx = -\int_{\Omega} (\Delta u) v \, dx = -\sum_{j=1}^d \int_{\Omega} \frac{\partial^2 u}{\partial x_j^2} v \, dx = \sum_{j=1}^d \left[\int_{\Omega} \frac{\partial u}{\partial x_j} \, \frac{\partial v}{\partial x_j} \, dx - \int_{\Gamma} \frac{\partial u}{\partial x_j} \, v n_j \, ds \right].$$

With $x \cdot y = \sum_{j=1}^d x_j y_j$ the usual scalar product in \mathbb{R}^d , we obtain the first Green formula

$$\int_{\Omega} f v \, dx = \int_{\Omega} \nabla u \cdot \nabla v \, dx - \int_{\Gamma} \frac{\partial u}{\partial n} v \, ds, \tag{1.4}$$

where we have used $\nabla u \cdot n = \partial u/\partial n$. Together with $v|_{\Gamma_D} = 0$ and $\Gamma_N = \Gamma \setminus \overline{\Gamma}_D$, we may plug-in the Neumann data to see that

$$\int_{\Omega} f v \, dx = \int_{\Omega} \nabla u \cdot \nabla v \, dx - \int_{\Gamma_N} \frac{\partial u}{\partial n} \, v \, ds = \int_{\Omega} \nabla u \cdot \nabla v \, dx - \int_{\Gamma_N} \phi v \, ds.$$

Altogether we thus have proven the following proposition:

Proposition 1.1. Let $u \in C^2(\overline{\Omega})$ solve the strong form (1.1). Then, it holds that

$$\int_{\Omega} \nabla u \cdot \nabla v \, dx = \int_{\Omega} f v \, dx + \int_{\Gamma_N} \phi v \, ds \quad \text{for all } v \in C_D^1(\overline{\Omega}), \tag{1.5}$$

which is the **variational form** of the boundary value problem (1.1).

This proposition gives a necessary condition for a function u to solve the strong form (1.1). We stress that any strong solution belongs to $C_D^1(\overline{\Omega})$ and that the variational form (1.5) can be understood for $u \in C_D^1(\overline{\Omega})$. This leads to a symmetric variational formulation: Find $u \in C_D^1(\overline{\Omega})$ such that (1.5) holds.

Exercise 1. Prove the following well-known integral formulae:

• For $f \in C^1(\Omega)^d$, let div $f := \sum_{j=1}^d \frac{\partial f_j}{\partial x_j}$ denote the divergence operators. Then, there holds the **Gauss divergence theorem**

$$\int_{\Omega} \operatorname{div} f \, dx = \int_{\Gamma} f \cdot n \, ds \quad \text{for all } f \in C^{1}(\overline{\Omega})^{d}. \tag{1.6}$$

• Besides the first Green formula, there holds the **second Green formula**

$$\int_{\Omega} (-\Delta u)v \, dx + \int_{\Gamma} \frac{\partial u}{\partial n} \, v \, ds = \int_{\Omega} u(-\Delta v) \, dx + \int_{\Gamma} u \, \frac{\partial v}{\partial n} \, ds \quad \text{for all } u, v \in C^2(\overline{\Omega}). \quad (1.7)$$

Both are easily obtained from the integration by parts formula.

1.2 Solvability of Variational Form

To look for solutions of the weak form (1.5), we will employ the following Riesz theorem.

Theorem 1.2 (Riesz). For a Hilbert space H (over \mathbb{R}), the mapping

$$I_H: H \to H^*, \quad I_H(u) := (u; \cdot)_H$$
 (1.8)

is linear, isometric, and bijective, i.e., for any $F \in H^*$ there is a unique $u \in H$ such that

$$(u;v)_H = F(v) \quad \text{for all } v \in H.$$
 (1.9)

Moreover, it holds that $||u||_H = ||F||_{H^*}$.

The proof of this theorem can be found in each textbook of functional analysis.

First, we observe that the left-hand side

$$(u;v) := \int_{\Omega} \nabla u \cdot \nabla v \, dx$$

of the variational form (1.5) defines a scalar product on $C_D^1(\overline{\Omega})$, provided the Dirichlet boundary Γ_D is nontrivial: Clearly, $(u \ ; v)$ is a symmetric bilinear form on $C_D^1(\overline{\Omega})$. It thus only remains to prove definiteness. Note that $0 = (u \ ; u) = \|\nabla u\|_{L^2(\Omega)}^2$ implies $\nabla u = 0$, whence u is constant in Ω . Together with $u|_{\Gamma_D} = 0$, this proves u = 0. Moreover, the right-hand side

$$F(v) := \int_{\Omega} f v \, dx + \int_{\Gamma_N} \phi v \, ds$$

defines a linear functional on $C_D^1(\overline{\Omega})$ which is continuous with respect to the induced norm $||v|| := (v ; v)^{1/2}$. We prove this claim only in the special situation $\Gamma = \Gamma_D$ and postpone the abstract proof to a subsequent section.

Lemma 1.3 (Friedrichs' inequality). Suppose that $\Omega = [a, b] \times [c, d] \subset \mathbb{R}^2$ and $\Gamma_D = \partial \Omega$. Then, it holds that $\|v\|_{L^2(\Omega)} \leq \operatorname{diam}(\Omega) \|\nabla v\|_{L^2(\Omega)}$ for all $v \in C_D^1(\overline{\Omega})$.

Proof. For $x = (x_1, x_2) \in \Omega$, it holds that $v(x_1, c) = 0$. Therefore, the fundamental theorem of calculus yields that

$$v(x) = \int_{c}^{x_2} \partial_2 v(x_1, t) dt.$$

The Hölder inequality yields that

$$|v(x)| \le |d - c|^{1/2} \left(\int_c^{x_2} |\partial_2 v(x_1, t)|^2 dt \right)^{1/2}.$$

Integration over Ω gives

$$||v||_{L^{2}(\Omega)}^{2} = \int_{\Omega} |v(x)|^{2} dx \le |d - c| \int_{\Omega} \int_{c}^{x_{2}} |\partial_{2}v(x_{1}, t)|^{2} dt dx$$

$$= |d - c| \int_{c}^{d} \int_{a}^{b} \int_{c}^{x_{2}} |\partial_{2}v(x_{1}, t)|^{2} dt dx_{1} dx_{2}$$

$$\le |d - c| \int_{c}^{d} ||\partial_{2}v||_{L^{2}(\Omega)}^{2} dx_{2}$$

$$= |d - c|^{2} ||\partial_{2}v||_{L^{2}(\Omega)}^{2}.$$

This results in $||v||_{L^2(\Omega)} \le |d-c| ||\partial_2 v||_{L^2(\Omega)} \le \operatorname{diam}(\Omega) ||\nabla v||_{L^2(\Omega)}$.

According to the Hölder and the Friedrichs inequality, we obtain that

$$|F(v)| \le ||f||_{L^2(\Omega)} ||v||_{L^2(\Omega)} \le \operatorname{diam}(\Omega) ||f||_{L^2(\Omega)} ||\nabla v||_{L^2(\Omega)} = \operatorname{diam}(\Omega) ||f||_{L^2(\Omega)} ||v||.$$

Therefore, the linear functional F is continuous with respect to $\|\cdot\| := \|\nabla(\cdot)\|_{L^2(\Omega)}$ with operator norm $\|F\|_* \leq \operatorname{diam}(\Omega)\|f\|_{L^2(\Omega)}$. If $C_D^1(\overline{\Omega})$ associated with the norm $\|\cdot\|$ were a Hilbert space, the Riesz theorem would therefore imply the unique solvability of the variational form (1.5). However, $C_D^1(\overline{\Omega})$ is not complete and therefore the Riesz theorem does not apply.

The remedy is to consider the (unique) completion of $C_D^1(\overline{\Omega})$ with respect to $\|\cdot\|$. This leads to a so-called **Sobolev space** $H_D^1(\Omega)$, which is —by definition— complete and hence a Hilbert space. Density arguments then lead to an extended variational form: Find $u \in H_D^1(\Omega)$ such that

$$\int_{\Omega} \nabla u \cdot \nabla v \, dx = \int_{\Omega} f v \, dx + \int_{\Gamma_N} \phi v \, ds \quad \text{for all } v \in H_D^1(\Omega), \tag{1.10}$$

which is the **weak form** of the boundary value problem (1.1). Now, the Riesz theorem applies and proves the unique existence of a **weak solution** $u \in H_D^1(\Omega)$ of (1.10). Later on, we are going to show that

- each strong solution $u \in C^2(\overline{\Omega})$ of (1.1) belongs to $H_D^1(\Omega)$ and is also the unique weak solution of (1.10).
- provided the weak solution $u \in H_D^1(\Omega)$ is smooth, i.e., $u \in C^2(\overline{\Omega})$, the weak solution also solves the strong form (1.1).

In this sense, the strong form (1.1) and the weak form (1.10) are equivalent.

1.3 Finite Element Method

The finite element method for (1.10) essentially consists of replacing the (infinite dimensional) Sobolev space $H_D^1(\Omega)$ by a finite dimensional subspace $X_h \subset H_D^1(\Omega)$: Find $u_h \in X_h$ such that

$$\int_{\Omega} \nabla u_h \cdot \nabla v_h \, dx = \int_{\Omega} f v_h \, dx + \int_{\Gamma_N} \phi v_h \, ds \quad \text{for all } v_h \in X_h.$$
 (1.11)

This problem is equivalent to the solution of a system of linear equations $\mathbf{A}\mathbf{x} = \mathbf{b}$, where the system matrix \mathbf{A} is symmetric and positive definite. Of course, the question of convergence depends on the choice of X_h . Thus, there remain some topics for mathematical discussions later on.

The finite element method is a special **Galerkin scheme**. In this section, we collect the most simple properties of Galerkin schemes. Throughout, H is a (real) Hilbert space, and $\langle \cdot ; \cdot \rangle$ is an equivalent scalar product on H, i.e., there are constants $\alpha, \beta > 0$ such that

$$\alpha \|v\|_H \le \|v\| \le \beta \|v\|_H \quad \text{for all } v \in H, \tag{1.12}$$

where $||v|| := \langle v; v \rangle^{1/2}$ denotes the induced norm. We stress that $\langle \cdot; \cdot \rangle$ and $||\cdot||$ are often called **energy scalar product** and **energy norm**, respectively (see also Exercise 5).

Remark. In the following, we state all results with respect to the norm $\|\cdot\|_H$, which involves the constants $\alpha, \beta > 0$. Analogously, one may state the results with respect to the energy norm $\|\cdot\| = \|\cdot\|_H$, which corresponds to $\alpha = \beta = 1$.

For given $F \in H^*$, the Riesz theorem proves the existence and uniqueness of a solution $u \in H$ of

$$\langle u; v \rangle = F(v) \quad \text{for all } v \in H,$$
 (1.13)

for what we use the short-hand notation

$$\langle\!\langle u ; \cdot \rangle\!\rangle = F \in H^* \tag{1.14}$$

to implicitly indicate that this equation holds (pointwise) for all $v \in H$. Now, the Galerkin method simply consists in replacing the continuous space H by some finite dimensional subspace: Let X_h be a finite-dimensional (and hence closed) subspace of H. Since the Riesz theorem applies to the Hilbert space X_h as well, there is a unique **Galerkin solution** $u_h := \mathbb{G}_h u \in X_h$ such that

$$\langle\!\langle \mathbb{G}_h u ; \cdot \rangle\!\rangle = F \in X_h^*. \tag{1.15}$$

For $u \in H$ and the corresponding functional $\langle u : \cdot \rangle \in H^*$, this defines the Galerkin projection

$$\mathbb{G}_h: H \to X_h \quad \text{where } \mathbb{G}_h u \in X_h \text{ solves} \quad \langle\!\langle \mathbb{G}_h u ; \cdot \rangle\!\rangle = \langle\!\langle u ; \cdot \rangle\!\rangle \in X_h^*.$$
 (1.16)

Note that $\mathbb{G}_h u \in X_h$ is characterized by the **Galerkin orthogonality**

$$\langle u - \mathbb{G}_h u ; v_h \rangle = 0 \quad \text{for all } v_h \in X_h.$$
 (1.17)

Before we proceed with the theoretical analysis of Galerkin schemes, we treat an implementational issue. The following theorem is the fundamental observation: Usually, only the scalar product $\langle \cdot ; \cdot \rangle$ and the right-hand side $F \in H^*$ are known, while the exact solution $u \in H$ of (1.13) is unknown. Then, the Galerkin solution $\mathbb{G}_h u \in X_h$ can be computed by solving a linear system of equations — without knowledge of u.

Theorem 1.4. Let $\{\phi_1, \ldots, \phi_N\}$ be a basis of X_h . We define the Galerkin matrix $A \in \mathbb{R}^{N \times N}$ and the vector $b \in \mathbb{R}^N$ by

$$A_{jk} := \langle \! \langle \phi_k ; \phi_j \rangle \! \rangle \quad and \quad b_j := F(\phi_j). \tag{1.18}$$

Then, A is symmetric and positive definite and, in particular, a regular matrix. Moreover, there holds $\mathbb{G}_h u = \sum_{j=1}^N x_j \phi_j$, where the vector $x \in \mathbb{R}^N$ solves Ax = b.

Proof. 1. step. Symmetry of A clearly follows from the symmetry of $\langle \cdot ; \cdot \rangle$.

2. step. For any $x \in \mathbb{R}^N$ and $v_h := \sum_{i=1}^N x_i \phi_i$, it holds that

$$|||v_h|||^2 = \langle v_h ; v_h \rangle = \sum_{j,k=1}^N x_j x_k \langle \phi_j ; \phi_k \rangle = x \cdot Ax.$$

This proves $Ax \cdot x > 0$ for all $x \neq 0$. By definition, A is positive definite and hence regular.

3. step. Determine Galerkin solution: Let $x \in \mathbb{R}^n$ be the unique solution of the linear Galerkin system Ax = b. We use the basis representation $\mathbb{G}_h u = \sum_{j=1}^N y_j \phi_j$ of the Galerkin solution with some coefficient vector $y \in \mathbb{R}^n$. By use of the linearity of $\langle \cdot, \cdot \rangle$, equation (1.15) becomes

$$b_k = F(\phi_k) = \langle \langle \mathbb{G}_h u ; \phi_k \rangle \rangle = \sum_{j=1}^N y_j \langle \langle \phi_j ; \phi_k \rangle \rangle = (Ay)_k \text{ for all } k = 1, \dots, N.$$

Therefore, the coefficient vector $y \in \mathbb{R}^N$ satisfies Ay = b. This proves x = y, i.e., we obtain $\mathbb{G}_h u$ by solving Ax = b.

Remark. We just remark that Theorem 1.4 can be applied for any orthogonal-type projection, e.g., the L^2 -orthogonal projection onto a discrete space.

We now proceed with the abstract analysis of Galerkin schemes. The following two lemmata provide elementary properties of the Galerkin projection. The first lemma proves stability of the method with respect to changes of the right-hand side F.

Lemma 1.5. The Galerkin projection \mathbb{G}_h is a linear and continuous projection onto X_h with

$$\|\mathbb{G}_h u\|_H \le \frac{\beta}{\alpha} \|u\|_H \quad \text{for all } u \in H, \tag{1.19}$$

where $\alpha, \beta > 0$ are the norm equivalence constants from (1.12). Moreover, \mathbb{G}_h is the orthogonal projection onto X_h with respect to the energy scalar product $\langle \cdot, \cdot \rangle$.

Proof. For $u_h \in X_h$, the Galerkin orthogonality (1.17) implies $\mathbb{G}_h u_h = u_h$. Therefore \mathbb{G}_h is a projection onto X_h . Also the linearity of \mathbb{G}_h follows from the Galerkin orthogonality (1.17). To see the continuity of \mathbb{G}_h , it remains to estimate the operator norm: For $u \in H$ holds

$$\|\mathbb{G}_h u\|^2 = \langle\!\langle \mathbb{G}_h u ; \mathbb{G}_h u \rangle\!\rangle = \langle\!\langle u ; \mathbb{G}_h u \rangle\!\rangle \le \|\|u\|\|\mathbb{G}_h u\|\|,$$

whence $\|\mathbb{G}_h u\| \leq \|u\|$ and

$$\alpha \| \mathbb{G}_h u \|_H \le \| \mathbb{G}_h u \| \le \| u \| \le \beta \| u \|_H$$

where we have used the norm equivalence (1.12) on H as well as the Cauchy inequality for the scalar product $\langle\!\langle \cdot ; \cdot \rangle\!\rangle$. This proves that $\|\mathbb{G}_h u\|_H \leq (\alpha/\beta)\|u\|_H$ and thus continuity of \mathbb{G}_h . Finally,

we remark that the *unique* orthogonal projection with respect to $\langle \cdot ; \cdot \rangle$, is characterized by the orthogonality relation (1.17).

The following Céa lemma states that the **Galerkin error** $||u - \mathbb{G}_h u||_H$ is quasi-optimal, i.e., it behaves like the best approximation error up to multiplicative constants, which depend only on the continuous setting but not on X_h .

Lemma 1.6 (Céa). The Galerkin error is quasi-optimal, i.e.,

$$||u - \mathbb{G}_h u||_H \le \frac{\beta}{\alpha} \min_{v_h \in X_h} ||u - v_h||_H \quad \text{for all } u \in H,$$
 (1.20)

where $\alpha, \beta > 0$ are the norm equivalence constants from (1.12). With respect to the energy norm, it holds that

$$|||u - \mathbb{G}_h u|| = \min_{v_h \in X_h} |||u - v_h|| \quad \text{for all } u \in H,$$
 (1.21)

i.e., the Galerkin solution $\mathbb{G}_h u$ is the best approximation of u with respect to the energy norm.

Proof. For arbitrary $v_h \in X_h$, the Galerkin orthogonality (1.17) proves that

$$||u - \mathbb{G}_h u||^2 = \langle u - \mathbb{G}_h u ; u - v_h \rangle \le ||u - \mathbb{G}_h u|| ||u - v_h||,$$

which yields (1.21) with an infimum on the right-hand side. Of course, the minimum in (1.21) is attained for $v_h = \mathbb{G}_h u$. With the same arguments as in the proof of the last lemma, we even see that

$$\alpha \|u - \mathbb{G}_h u\|_H \le \|u - \mathbb{G}_h u\| \le \|u - v_h\| \le \beta \|u - v_h\|_H$$

which implies (1.20) with an infimum on the right-hand side. This minimum is attained for $v_h = \Pi_h u$ with $\Pi_h : X \to X_h$ being the orthogonal projection onto X_h with respect to $\|\cdot\|_H$.

Exercise 2. Let X be a normed vector space over \mathbb{R} and $X_h \subseteq X$ be a finite dimensional subspace of X. Then, for any $x \in X$, there exists some (not necessarily unique) $x_h \in X_h$ such that

$$||x - x_h||_X = \min_{v_h \in X_h} ||x - v_h||_X,$$

i.e., best approximation errors on finite dimensional spaces as in (1.20) are always attained. Prove that the set of minimizers is convex, closed and bounded (and hence even compact). \Box

A major advantage of Galerkin methods is that one can prove convergence for any exact solution $u \in H$ if one knows that smooth functions can be approximated well. In the following, think of the subscript h > 0 as a mesh-size parameter with corresponding finite dimensional spaces X_h :

Proposition 1.7. For all h > 0, let X_h be a finite-dimensional subspace of H. We assume that there is a dense subspace D of H with approximation property, namely

$$\lim_{h \to 0} \min_{v_h \in X_h} \|v - v_h\|_H = 0 \quad \text{for all } v \in D.$$
 (1.22)

Then, for any $u \in H$, it holds that

$$\lim_{h \to 0} \|u - \mathbb{G}_h u\|_H = 0, \tag{1.23}$$

i.e., the sequence of Galerkin solutions converges to the exact solution u.

Proof. For $v \in D$, the quasi-optimality (1.20) yields that

$$||u - \mathbb{G}_h u||_H \le \frac{\beta}{\alpha} \min_{v_h \in X_h} ||u - v_h||_H \le \frac{\beta}{\alpha} (||u - v||_H + \min_{v_h \in X_h} ||v - v_h||_H).$$

We have to show that

$$\exists C > 0 \,\forall \varepsilon > 0 \,\exists h_0 > 0 \,\forall h \in (0, h_0) \quad \|u - \mathbb{G}_h u\|_H \le C \,\varepsilon.$$

For $\varepsilon > 0$, let $v \in D$ with $||u - v||_H \le \varepsilon$. Choose $h_0 > 0$ according to the approximation assumption (1.22) so that $\min_{v_h \in X_h} ||v - v_h||_H \le \varepsilon$ for all $h \in (0, h_0)$. We thus finally obtain $||u - \mathbb{G}_h u||_H \le 2\beta \varepsilon / \alpha$, which concludes the proof.

Although the result of the preceding lemma seems to be very attractive, we stress, however, that the convergence of a Galerkin scheme can be arbitrarily slow. We argue in the abstract setting: If H is a separable Hilbert space, e.g., H is a Sobolev space, there is a countable orthonormal basis $\{\phi_j \mid j \in \mathbb{N}\}$. Any $u \in H$ can be written as $u = \sum_{j=1}^{\infty} x_j \phi_j$ with coefficients $(x_n) \in \ell_2$. If we define $X_j := \operatorname{span}\{\phi_1, \ldots, \phi_j\}$, it holds that

$$\min_{v_h \in X_h} \|u - v_h\|_H^2 = \sum_{j=k+1}^{\infty} x_j^2.$$

Finally, the decay of the right-hand side can be very slow. One may think of, e.g., $x_j^2 = j^{-(1+\varepsilon)}$ for any $\varepsilon > 0$, so that the series converges but is — in the beginning — almost the divergent harmonic series.

The following exercise shows that the approximation property (1.22) in particular implies that the Hilbert space H has to be separable.

Exercise 3. Suppose that X is a normed space with finite dimensional subspaces $X_{\ell} \subseteq X_{\ell+1} \subseteq X$ for all $\ell \in \mathbb{N}$. Suppose that $\mathcal{D} \subseteq X$ is a dense subspace such that, for all $x \in X$,

$$\lim_{\ell \to \infty} \min_{x_{\ell} \in X_{\ell}} \|x - x_{\ell}\|_{X} = 0.$$
 (1.24)

Then, X is separable, i.e., there is a countable and dense subset $M \subseteq X$.

Exercise 4. Let $X = \ell_{\infty}$ and $X_{\ell} := \{(x_n) \in \ell_{\infty} \mid x_j = 0 \text{ for all } j \geq \ell \}$. Prove that (1.24) fails to hold for any dense subspace \mathcal{D} . Note that this also follows if one proves that ℓ_{∞} is not separable.

Remark. All foregoing results of this section hold (in a slightly modified form) in case that $\langle \cdot ; \cdot \rangle$ only is a continuous and elliptic bilinear form on the Hilbert space H, i.e., in all proofs, one can avoid to use the symmetry of $\langle \cdot ; \cdot \rangle$.

The following exercise explains why $\|\cdot\|$ is called energy norm. In many situations, the function $J(\cdot)$ has the interpretation of a physical energy.

Exercise 5. Let $\langle \! \langle \cdot ; \cdot \rangle \! \rangle$ be a scalar product on the Hilbert space H such that the norm $\| \! | \cdot | \! |$ is equivalent to $\| \cdot \|_H$. Let $F \in H^*$ and $u \in H$. Then, the following assertions are equivalent:

- $\langle u ; \cdot \rangle = F \in H^*;$
- $J(u) = \min_{v \in H} J(v)$, where $J(v) := \frac{1}{2} \langle v ; v \rangle F(v)$.

In particular, the variational formulation is equivalent to energy minimization, and this result also covers the discrete setting. Derive a formula for the energy error $J(\mathbb{G}_h u) - J(u)$, where $\mathbb{G}_h: H \to X_h$ denotes the Galerkin projection.

Finally, we comment on an extension of the concept of Galerkin schemes to some nonlinear problems. We note that this framework does, in particular, cover the frame of the Lax–Milgram lemma.

Exercise 6 (Main Theorem on Strongly Monotone Operators (Zarantonello '60)). Let H be a Hilbert space and $A: H \to H^*$ be a Lipschitz continuous and strongly monotone operator, i.e.,

$$||Au - Av||_{H^*} \le L||u - v||_H$$
 and $\langle Au - Av ; u - v \rangle_{H^* \times H} \ge M||u - v||_H^2$ for all $u, v \in H$

with constants L, M > 0 that only depend on A. Then, A is bijective. **Hint:** Injectivity of A follows from the monotonicity of A. To prove surjectivity, we apply a fixed point argument: Let $I_H: H \to H^*$, $I_H(u) := (u ; \cdot)_H$ denote the Riesz mapping. For given $F \in H^*$ and a certain choice of C > 0, the mapping $\Phi(u) := u - CI_H^{-1}(Au - F)$ is a contraction on H. Therefore, the Banach contraction theorem applies and provides a unique $u \in H$ with $u = \Phi(u)$.

Exercise 7 (Lemma of Lax–Milgram). Use Exercise 6 to derive the Lemma of Lax–Milgram: Let H be a Hilbert space and $a(\cdot,\cdot)$ be a continuous and elliptic bilinear form on H, i.e.,

$$a(u, v) \le L \|u\|_H \|v\|_H$$
 and $a(u, u) \ge M \|u\|_H^2$ for all $u, v \in H$,

where the constants L, M > 0 depend only on $a(\cdot, \cdot)$. Then, given a right-hand side $F \in H^*$,

$CHAPTER\ 1.\ \ INTRODUCTION$

there is a unique $u \in H$ with $a(u, \cdot) = F \in H^*$.

Chapter 2

Sobolev Spaces and Poisson Problem

2.1 Sobolev Spaces on Domains

This section briefly recalls the definition of Sobolev spaces $H^m(\Omega)$, for integer order $m \in \mathbb{N}_0$, on domains $\Omega \subseteq \mathbb{R}^d$. While this section requires Ω only to be open and connected, the following sections will implicitly assume that Ω is a bounded Lipschitz domain.

Definition. A function $u \in L^1_{loc}(\Omega) := \{w : \Omega \to \mathbb{R} \text{ measurable } | \forall K \subset \Omega \text{ compact } w \in L^1(K) \}$ has a **weak partial derivative** $\partial_j u \in L^1_{loc}(\Omega)$, if the pair $(u, \partial_j u)$ satisfies the integration by parts formula with smooth test functions that vanish on the boundary, i.e., it holds that

$$\int_{\Omega} u(\partial_j v) \, dx = -\int_{\Omega} (\partial_j u) v \, dx \quad \text{for all } v \in \mathcal{D}(\Omega) := C_c^{\infty}(\Omega). \tag{2.1}$$

Note that $\partial_j u$ is (so far) only a symbol, whereas $\partial_j v := \partial v / \partial x_j$ is the classical j-th derivative of $v \in \mathcal{D}(\Omega)$. We say that $u \in L^1_{loc}(\Omega)$ is **weakly differentiable with weak gradient** $\nabla u \in L^1_{loc}(\Omega)$, if all weak derivatives $\partial_j u$, for $j = 1, \ldots, d$, exist.

The following main theorem of calculus, which will not be proven in this lecture, we infer that the weak derivative is unique, if it exists. Moreover, the weak derivative and the classical derivative coincide, if the classical derivative exists.

Theorem 2.1 (Fundamental Theorem of Calculus of Variations). Let
$$f \in L^1_{loc}(\Omega)$$
 satisfy $\int_{\Omega} fv \, dx = 0$ for all $v \in \mathcal{D}(\Omega)$. Then, it holds that $f = 0$ almost everywhere in Ω .

Remark. Note that $C(\Omega) \subset L^1_{loc}(\Omega)$. For $f \in C(\Omega)$, the fundamental theorem of calculus of variations can be proven by elementary calculus: Note that for any $x \in \mathbb{R}^d$ and any radius $\varepsilon > 0$, there is a function $\psi \in \mathcal{D}(\mathbb{R}^d)$ such that $\{y \in \mathbb{R}^d \mid \psi(y) > 0\} = U(x,\varepsilon) := \{y \in \mathbb{R}^d \mid |x-y| < \varepsilon\}$; see the following Exercise 8. Provided $f \in C(\Omega)$ with $f(x) \neq 0$ for some $x \in \Omega$, we may assume f(x) > 0. By continuity, there is a small radius $\varepsilon > 0$ such that $U(x,\varepsilon) \subset \Omega$ and that f(y) > 0 for all $y \in U(x,\varepsilon)$. With the associated function $\psi \in \mathcal{D}(\Omega)$, we thus see that $\int_{\Omega} f \psi \, dx > 0$. Note that this argument provides the (logically equivalent) contraposition of the fundamental theorem of calculus of variations in the case of a continuous function f.

Exercise 8. (i) Show that the following definition provides $\phi \in C^{\infty}(\mathbb{R})$ with supp $(\phi) = [-1, 1]$:

$$\phi(t) := \begin{cases} \exp\left(-1/(1-t^2)\right), & \text{for } |t| < 1, \\ 0 & \text{else.} \end{cases}$$

(ii) For $\varepsilon > 0$ and $x \in \mathbb{R}^d$, define the function $\psi_{x,\varepsilon}(y) := \phi(|x-y|^2/\varepsilon)$. Show that $\psi_{x,\varepsilon} \in C^{\infty}(\mathbb{R}^d)$ with $\operatorname{supp}(\psi_{x,\varepsilon}) = \left\{ y \in \mathbb{R}^d \, \middle| \, |x-y| \le \varepsilon \right\}$ and $\psi_{x,\varepsilon}(y) > 0$ for all $y \in \left\{ y \in \mathbb{R}^d \, \middle| \, |x-y| < \varepsilon \right\}$. \square

Corollary 2.2. (i) The weak derivative $\partial_j u$ is unique, if it exists: If $\partial_j u$, $\widetilde{\partial_j u} \in L^1_{\ell oc}(\Omega)$ satisfy (2.1), it holds that $\partial_j u = \widetilde{\partial_j u}$ almost everywhere in Ω .

(ii) A function $u \in C^1(\Omega)$ is weakly differentiable, and the weak derivative coincides with the classical derivative.

Proof. (i) It holds that $\int_{\Omega} (\partial_j u - \widetilde{\partial_j u}) v \, dx = 0$ for all $v \in \mathcal{D}(\Omega)$ and thus $\partial_j u - \widetilde{\partial_j u} = 0$ almost everywhere in Ω . (ii) follows from (i) and the integration by parts formula.

A deeper result is the following, which is somehow, nevertheless, quite natural and expected.

Theorem 2.3. If $u \in L^1_{loc}(\Omega)$ is weakly differentiable with $\nabla u = 0$, then the function u is constant, i.e., there is a constant $c \in \mathbb{R}$ such that u = c almost everywhere in Ω .

Definition. For m=0, we define $H^0(\Omega):=L^2(\Omega)$ as the classical Lebesgue space of square integrable functions. For m=1, the **Sobolev space** $H^1(\Omega)$ is defined by

$$H^1(\Omega) := \{ u \in L^2(\Omega) \mid u \text{ weakly differentiable, } \nabla u \in L^2(\Omega) \}$$
 (2.2)

and associated with the graph norm

$$||u||_{H^1(\Omega)} := (||u||_{L^2(\Omega)}^2 + ||\nabla u||_{L^2(\Omega)}^2)^{1/2}.$$
 (2.3)

Higher-order Sobolev spaces of integer order $m \in \mathbb{N}$ may be defined inductively by

$$H^{m}(\Omega) := \left\{ u \in L^{2}(\Omega) \mid u \text{ weakly differentiable, } \nabla u \in H^{m-1}(\Omega) \right\}, \tag{2.4}$$

with associated norm

$$||u||_{H^{m}(\Omega)} := (||u||_{L^{2}(\Omega)}^{2} + ||\nabla u||_{H^{m-1}(\Omega)}^{2})^{1/2}.$$
(2.5)

Remark. Clearly, $C^1(\overline{\Omega}) \subseteq H^1(\Omega)$ and we note below that $C^1(\overline{\Omega})$ is even dense in $H^1(\overline{\Omega})$.

Theorem 2.4. For all $m \in \mathbb{N}_0$, the Sobolev space $H^m(\Omega)$ is a Hilbert space.

Proof. The proof uses the (hopefully) well-known fact that $H^0(\Omega) = L^2(\Omega)$ is a Hilbert space. We shall proceed by induction on m. However, we explicitly consider the case m = 1 first: Obviously, the H^1 -norm is induced by the scalar product

$$(u; v)_{H^1(\Omega)} := (u; v)_{L^2(\Omega)} + (\nabla u; \nabla v)_{L^2(\Omega)}$$
 for all $u, v \in H^1(\Omega)$,

i.e., $||u||_{H^1(\Omega)}^2 = (u; u)_{H^1(\Omega)}$. Therefore, it only remains to prove the completeness of $H^1(\Omega)$. Let (u_n) be a Cauchy sequence in $H^1(\Omega)$. Note that, by definition of the H^1 -norm, (u_n) as well as (∇u_n) are Cauchy sequences in $L^2(\Omega)$. Since $L^2(\Omega)$ is complete, there are unique $u \in L^2(\Omega)$ and $g \in L^2(\Omega)^d$ such that

$$\lim_{n \to \infty} ||u - u_n||_{L^2(\Omega)} = 0 = \lim_{n \to \infty} ||g - \nabla u_n||_{L^2(\Omega)}.$$

By definition of $H^1(\Omega)$, it thus only remains to prove that u is weakly differentiable with gradient $\nabla u = g$. Let $v \in \mathcal{D}(\Omega)$ be an arbitrary test function. From the weak differentiability of each u_n and L^2 -convergence, we obtain that

$$(u \; ; \; \partial_j v)_{L^2(\Omega)} = \lim_{n \to \infty} (u_n \; ; \; \partial_j v)_{L^2(\Omega)} = -\lim_{n \to \infty} (\partial_j u_n \; ; \; v)_{L^2(\Omega)} = -(g_j \; ; \; v)_{L^2(\Omega)}.$$

Therefore, g_j is the j-th weak derivative of u and consequently $g = \nabla u$. This concludes the case m = 1. The induction step for $H^m(\Omega)$ is left to the reader, but obviously follows from the same arguments, where we replace $g \in L^2(\Omega)^d$ by $g \in H^{m-1}(\Omega)^d$.

2.2 Main Theorems on Sobolev Spaces

From now on, it will be important and thus assumed that $\Omega \subset \mathbb{R}^d$ is a bounded Lipschitz domain. By definition of the Sobolev spaces $H^m(\Omega)$, there holds $H^m(\Omega) \subset H^{m-1}(\Omega)$ with $\|u\|_{H^{m-1}(\Omega)} \leq \|u\|_{H^m(\Omega)}$. In other words, the identity operator $id: H^m(\Omega) \to H^{m-1}(\Omega)$ is well-defined and continuous. The following Rellich theorem states that it is also compact. This is a pretty strong result. The impact of which will become clear in our proofs of the Poincaré inequality and the Friedrichs inequality.

Theorem 2.5 (Rellich Compactness Theorem). For any integer order
$$m \in \mathbb{N}$$
, the embedding $H^m(\Omega) \subseteq H^{m-1}(\Omega)$ is compact.

We recall that an operator $A \in L(X;Y)$ between normed spaces X and Y is compact, if and only if each bounded set $S \subseteq X$ is mapped to a pre-compact set $A(S) \subseteq Y$, i.e., $\overline{A(S)} \subseteq Y$ is compact.

Before the statement and the proof of the Poincaré inequality, we need a further technical lemma. The result is rather standard in the analysis of variational problems.

Lemma 2.6. A continuous and convex functional $f: X \to \mathbb{R}$ on a normed space X is weakly lower semicontinuous, i.e., for each weakly convergent sequence (x_n) in X with $x_n \rightharpoonup x \in X$,

it holds that

$$f(x) \le \liminf_{n \in \mathbb{N}} f(x_n). \tag{2.6}$$

Proof. 1. step. We prove that the epigraph $G := \{(x, \alpha) \in X \times \mathbb{R} \mid f(x) \leq \alpha\}$ is convex: For $(x, \alpha), (y, \beta) \in G$ and $0 \leq \theta \leq 1$, the convexity of f proves that

$$f(\theta x + (1 - \theta)y) \le \theta f(x) + (1 - \theta)f(y) \le \theta \alpha + (1 - \theta)\beta,$$

whence $\theta(x,\alpha) + (1-\theta)(y,\beta) \in G$, i.e., $G \subseteq X \times \mathbb{R}$ is convex.

2. step. We use the continuity of f to prove that G is also closed: Let (x_n, α_n) be a convergent sequence in G, i.e., it holds that $x_n \to x \in X$ and $\alpha_n \to \alpha \in \mathbb{R}$. We prove that $(x, \alpha) \in G$, which follows from

$$f(x) = \lim_{n \to \infty} f(x_n) \le \lim_{n \to \infty} \alpha_n = \alpha.$$

3. step. The following step in the proof is known as $Mazur's\ lemma$: We prove that the closed and convex set G is also weakly closed in $X \times \mathbb{R} =: Y$, i.e., closed with respect to the weak topology on Y. We argue by contradiction and assume that G is not weakly closed. Then, there is an element $y \in \overline{G}^{\sigma} \setminus G$, where \overline{G}^{σ} denotes the weak closure of G. According to the Hahn-Banach separation theorem, there is a functional $\phi \in Y^*$ and a scalar $\lambda \in \mathbb{R}$ such that $\phi(y) < \lambda \leq \inf \phi(G)$. Therefore $U := \phi^{-1}(-\infty, \lambda)$ is weakly open with $y \in U$ and $U \cap G = \emptyset$. This contradicts topologically that $y \in G$ is in the weak closure of G. Hence, $G = \overline{G}^{\sigma}$ is weakly closed, and we may proceed with the proof of G.

4. step. We show the weak lower semicontinuity of f: Suppose that $x_n \to x \in X$. For $\alpha := \liminf_n f(x_n) = \infty$, (2.6) is trivial. We thus may assume $\alpha < \infty$. Let $\beta > \alpha$ and define $\alpha_n := \max\{\beta, f(x_n)\} \to \beta$. Clearly, $(x_n, \alpha_n) \in G$. Moreover, this sequence is weakly convergent $(x_n, \alpha_n) \to (x, \beta)$. We deduce $(x, \beta) \in G$. Thus, $f(x) \leq \beta$ for all $\beta > \alpha$ and therefore finally $f(x) \leq \alpha = \lim_{n \to \infty} f(x_n)$.

A first consequence of the preceding abstract results is that one can easily construct equivalent norms on the Sobolev space $H^1(\Omega)$.

Proposition 2.7. Let $|\cdot|_{H^1}$ be a continuous seminorm on $H^1(\Omega)$ which is definite on the constant functions, i.e., $|c|_{H^1} = 0$ implies c = 0 for all $c \in \mathbb{R}$. Then, there are constants $C_1, C_2 > 0$ such that

 $|v|_{H^1} \le C_1 \|v\|_{H^1(\Omega)}$ as well as $C_2^{-1} \|v\|_{L^2(\Omega)} \le \|v\| := \|\nabla v\|_{L^2(\Omega)} + |v|_{H^1}$ for all $v \in H^1(\Omega)$.

In particular, $\|\cdot\|$ defines an equivalent norm on $H^1(\Omega)$, i.e.,

$$(1+C_1)^{-1} |||v||| \le ||v||_{H^1(\Omega)} \le (1+C_2) |||v||| \quad \text{for all } v \in H^1(\Omega).$$

Proof. 1. step. Existence of C_1 : By definition of continuity, there exists an open neighborhood $O \subseteq H^1(\Omega)$ of 0 such that $|v|_{H^1} \le 1$ for all $v \in O$. Without loss of generality, we may choose a

radius r > 0 sufficiently small such that $\overline{B_r(0)} \subset O \subset H^1(\Omega)$ for the closed ball with radius r and center zero. This implies

$$|v|_{H^1} = \frac{1}{r} ||v||_{H^1(\Omega)} |r \frac{v}{||v||_{H^1(\Omega)}}|_{H^1} \le \frac{1}{r} ||v||_{H^1(\Omega)}.$$

This proves existence of $C_1 := 1/r$.

2. step. Existence of C_2 : We assume that there is no constant $C_2 > 0$ such that $||v||_{L^2(\Omega)} \le C_2 ||v||$ for all $v \in H^1(\Omega)$. Therefore, there exists a sequence (v_n) in $H^1(\Omega)$ such that

$$\frac{1}{n} \|v_n\|_{L^2(\Omega)} > \|v_n\| = \|\nabla v_n\|_{L^2(\Omega)} + |v_n|_{H^1}$$

The definition of $w_n := v_n/\|v_n\|_{L^2(\Omega)}$ leads to to a sequence (w_n) in $H^1(\Omega)$ such that

$$||w_n||_{L^2(\Omega)} = 1$$
, $||\nabla w_n||_{L^2(\Omega)} \le 1/n$, $||w_n||_{H^1} \le 1/n$.

Therefore, (w_n) is a bounded sequence in the Hilbert space $H^1(\Omega)$. A Hilbert space is reflexive. By virtue of the Banach-Alaoglou theorem, each bounded sequence thus has a weakly convergent subsequence. Therefore, we may assume that $w_n \rightharpoonup w \in H^1(\Omega)$. An application of Lemma 2.6 proves that

$$\|\nabla w\|_{L^2(\Omega)} \le \liminf_{n \to \infty} \|\nabla w_n\|_{L^2(\Omega)} = 0,$$

whence the weak limit w is constant. Another application of Lemma 2.6 proves that

$$|w|_{H^1} \le \liminf_{n \to \infty} |w_n|_{H^1} = 0$$

since a seminorm is always convex. Therefore, w=0. On the other hand, the Rellich theorem states the strong convergence $w_n \to w \in L^2(\Omega)$ and thus $\|w\|_{L^2(\Omega)} = \lim_{n \to \infty} \|w_n\|_{L^2(\Omega)} = 1$. This contradiction concludes the existence of C_2 . In particular, we hence observe $\|v\|_{H^1(\Omega)} \leq \|v\|_{L^2(\Omega)} + \|\nabla v\|_{L^2(\Omega)} \leq (C_2+1) \|v\|$.

Corollary 2.8 (Poincaré Inequality). It holds that

$$||v||_{L^2(\Omega)} \le \widetilde{C}_P \left(||\nabla v||_{L^2(\Omega)} + \left| \int_{\Omega} v \, dx \right| \right) \quad \text{for all } v \in H^1(\Omega), \tag{2.7}$$

where the constant $\widetilde{C}_P > 0$ depends only on Ω . Moreover, $||v|| := ||\nabla v||_{L^2(\Omega)} + |\int_{\Omega} v \, dx|$ defines even an equivalent norm on $H^1(\Omega)$.

Proof. According to Proposition 2.7, it only remains to show that

$$|v|_{H^1} := \left| \int_{\Omega} v \, dx \right| \quad \text{for } v \in H^1(\Omega)$$

defines a continuous seminorm on $H^1(\Omega)$ which is definite on the constant functions. The equality $|c|_{H^1} = |\Omega| |c|$ for $c \in \mathbb{R}$ verifies the definiteness. Lipschitz continuity follows from

$$\left| |v|_{H^1} - |w|_{H^1} \right| \le \left| \int_{\Omega} v - w \, dx \right| \le |\Omega|^{1/2} ||v - w||_{L^2(\Omega)} \le |\Omega|^{1/2} ||v - w||_{H^1(\Omega)}$$

and from the boundedness of Ω .

Corollary 2.9 (Poincaré Inequality). There is a constant $C_P > 0$, which depends only on the shape of Ω but not on its diameter, such that

$$||v||_{L^{2}(\Omega)} \le C_{P} \operatorname{diam}(\Omega) ||\nabla v||_{L^{2}(\Omega)} \quad \text{for all } v \in H^{1}_{*}(\Omega) := \{ w \in H^{1}(\Omega) \mid \int_{\Omega} w \, dx = 0 \}, \quad (2.8)$$

where $\operatorname{diam}(\Omega) := \sup \{|x - y| \mid x, y \in \Omega\}$ denotes the diameter of Ω .

Proof. The proof is a so-called **scaling argument**: We define $\lambda := \operatorname{diam}(\Omega)$ and $\widetilde{\Omega} := \lambda^{-1}\Omega$. Note that the scaled domain $\widetilde{\Omega}$ satisfies $\operatorname{diam}(\widetilde{\Omega}) = 1$ and depends only on the shape of Ω . We consider the affine bijection $\Phi : \Omega \to \widetilde{\Omega}$, $\Phi(x) := \lambda^{-1}x$. Recall the transformation theorem, which holds for arbitrary diffeomorphisms $\Phi : \Omega \to \widetilde{\Omega}$ and states that

$$\int_{\widetilde{\Omega}} \widetilde{f} \, dy = \int_{\Omega} \widetilde{f}(\Phi(x)) \, | \det D\Phi(x) | \, dx \quad \text{for all } \widetilde{f} \in L^1(\widetilde{\Omega}).$$

Note that $\det D\Phi(x) = \lambda^{-d}$ since $D\Phi = \lambda^{-1}\boldsymbol{I}$ in our case. For $v \in H^1(\Omega)$, we define $\widetilde{v} := v \circ \Phi^{-1} \in H^1(\widetilde{\Omega})$. Then,

$$\|\widetilde{v}\|_{L^2(\widetilde{\Omega})}^2 = \int_{\widetilde{\Omega}} |\widetilde{v}|^2 dy = \lambda^{-d} \int_{\Omega} |v|^2 dx = \lambda^{-d} \|v\|_{L^2(\Omega)}^2.$$

According to the chain rule, it holds that $\nabla \tilde{v} = \lambda (\nabla v) \circ \Phi^{-1}$ and consequently that

$$\|\nabla \widetilde{v}\|_{L^2(\widetilde{\Omega})}^2 = \lambda^{2-d} \|\nabla v\|_{L^2(\Omega)}^2.$$

With $\widetilde{C}_P > 0$ the Poincaré constant from (2.7) for $\widetilde{\Omega}$, we thus infer

$$\|v\|_{L^2(\Omega)}^2 = \lambda^d \, \|\widetilde{v}\|_{L^2(\widetilde{\Omega})}^2 \leq \lambda^d \, \widetilde{C}_P^{\, 2} \, \|\nabla \widetilde{v}\|_{L^2(\widetilde{\Omega})}^2 = \lambda^2 \, \widetilde{C}_P^{\, 2} \, \|\nabla v\|_{L^2(\Omega)}^2.$$

Note that \widetilde{C}_P depends only on $\widetilde{\Omega}$ und thus only on the shape of Ω . This concludes the proof.

Remark. We stress that $Iv := \int_{\Omega} v \, dx$ defines a linear and continuous functional on $H^1(\Omega)$. In particular, $H^1_*(\Omega) = \ker(I)$ is a closed subspace of $H^1(\Omega)$ and hence a Hilbert space. According to the Poincaré inequality, it holds that $\|\nabla v\|_{L^2(\Omega)} \le \|v\|_{H^1(\Omega)} \le (1 + \widetilde{C}_P^2)^{1/2} \|\nabla v\|_{L^2(\Omega)}$ for all $v \in H^1_*(\Omega)$. In particular, $\|\nabla v\|_{L^2(\Omega)}$ defines an equivalent Hilbert norm on $H^1_*(\Omega)$ with associated scalar product $(\nabla u : \nabla v)_{L^2(\Omega)}$.

Theorem 2.10 (Meyers-Serrin). For each integer order $m \in \mathbb{N}$, $C^{\infty}(\overline{\Omega})$ and, in particular, $C^{\infty}(\Omega) \cap H^m(\Omega)$ are dense subspaces of $H^m(\Omega)$.

Theorem 2.11 (Trace Operator). There is a unique operator $\gamma \in L(H^1(\Omega); L^2(\Gamma))$ such that $\gamma v = v|_{\Gamma}$ for all $v \in C^1(\overline{\Omega})$, i.e., γ extends the classical trace defined as restriction $v|_{\Gamma}$ on the boundary for smooth functions v.

As a first corollary to Theorem 2.11, we can prove that the integration by parts formula also holds for Sobolev functions $u, v \in H^1(\Omega)$.

Corollary 2.12 (Integration by Parts). For all $u, v \in H^1(\Omega)$, it holds that

$$\int_{\Omega} u \, \frac{\partial v}{\partial x_j} \, dx + \int_{\Omega} \frac{\partial u}{\partial x_j} \, v \, dx = \int_{\Gamma} \gamma u \, \gamma v \, n_j \, ds. \tag{2.9}$$

Proof. The formula (2.9) holds for $u, v \in C^1(\overline{\Omega})$. All three terms define continuous bilinear forms on $H^1(\Omega) \times H^1(\Omega)$. Therefore (2.9) follows, for arbitrary $u, v \in H^1(\Omega)$ from the density of $C^1(\overline{\Omega})$ in $H^1(\Omega)$: Given $u, v \in H^1(\Omega)$, there are sequences (u_n) and (v_n) in $C^1(\overline{\Omega})$ which converge to u resp. v in $H^1(\Omega)$. Therefore, if $a(\cdot, \cdot) : H^1(\Omega) \times H^1(\Omega) \to \mathbb{R}$ is continuous, then it holds that $\lim_{n \to \infty} a(u_n, v_n) = a(u, v)$. This concludes the proof.

The analytical treatment of the Dirichlet problem makes use of the so-called Friedrichs inequality, whereas the analytical treatment of the Neumann problem uses the previously proven Poincaré inequality.

Corollary 2.13 (Friedrichs Inequality). Assume that the Dirichlet boundary $\Gamma_D \subseteq \Gamma$ has positive surface measure $|\Gamma_D| > 0$. Then, it holds that

$$||v||_{L^2(\Omega)} \le \widetilde{C}_F \left(||\nabla v||_{L^2(\Omega)} + ||\gamma v||_{L^2(\Gamma_D)} \right) \quad \text{for all } v \in H^1(\Omega)$$
 (2.10)

with a constant $\widetilde{C}_F > 0$, which depends only on Ω and Γ_D . Moreover, the right-hand side $\|v\| := \|\nabla v\|_{L^2(\Omega)} + \|\gamma v\|_{L^2(\Gamma_D)}$ even defines an equivalent norm on $H^1(\Omega)$.

Proof. We again apply Proposition 2.7. It only remains to show that

$$|v|_{H^1} := \|\gamma v\|_{L^2(\Gamma_D)} \quad \text{for } v \in H^1(\Omega)$$

defines a continuous seminorm on $H^1(\Omega)$ which is definite on the constant functions. The definiteness is again easily obtained from $|c|_{H^1} = |\Gamma_D|^{1/2}|c|$ for $c \in \mathbb{R}$. Lipschitz continuity follows from

$$||v|_{H^1} - |w|_{H^1}| \le ||\gamma v - \gamma w||_{L^2(\Gamma_D)} = ||\gamma (v - w)||_{L^2(\Gamma_D)} \le C ||v - w||_{H^1(\Omega)}$$

according to the continuity of the trace operator $\gamma \in L(H^1(\Omega); L^2(\Gamma))$.

Definition. We define $H_0^1(\Omega) := \overline{\mathcal{D}(\Omega)}^{\|\cdot\|_{H^1}}$ and $H_D^1(\Omega) := \overline{C_D^1(\overline{\Omega})}^{\|\cdot\|_{H^1}}$, where the subscript D indicates the Dirichlet boundary Γ_D . By definition, $H_0^1(\Omega)$ as well as $H_D^1(\Omega)$ are closed subspaces of $H^1(\Omega)$ and thus Hilbert spaces. In particular, it holds that $H_0^1(\Omega) \subseteq H_D^1(\Omega)$.

The same scaling argument as for the Poincaré inequality proves the following variant of the Friedrichs inequality, where we note that continuity of the trace operator γ proves that $\gamma v = 0$, for $v \in H_0^1(\Omega)$, as well as $(\gamma v)|_{\Gamma_D} = 0$, for $v \in H_D^1(\Omega)$.

$$||v||_{L^2(\Omega)} \le C_F \operatorname{diam}(\Omega) ||\nabla v||_{L^2(\Omega)} \quad \text{for all } v \in H^1_D(\Omega)$$
(2.11)

with a constant $C_F > 0$ that depends only on the shape of Ω and Γ_D .

We finally note the relation between $H_D^1(\Gamma)$ and the trace operator, cf. the Theorem of Meyers-Serrin.

Theorem 2.15. There holds
$$H_0^1(\Omega) = \ker(\gamma)$$
 with $\gamma \in L(H^1(\Omega); L^2(\Gamma))$ the trace operator. Moreover, $H_D^1(\Omega) = \{v \in H^1(\Omega) \mid (\gamma v)|_{\Gamma_D} = 0\}$.

Exercise 9. Usually, one defines the range of the trace operator as $H^{1/2}(\Gamma) := \operatorname{range}(\gamma) \subseteq L^2(\Gamma)$. This space is associated with the norm $\|v\|_{H^{1/2}(\Gamma)} := \inf \left\{ \|\widehat{v}\|_{H^1(\Omega)} \mid \widehat{v} \in H^1(\Omega) \text{ with } \gamma \widehat{v} = v \right\}$. Prove that $H^{1/2}(\Gamma)$ associated with this norm is a Hilbert space with continuous inclusion $H^{1/2}(\Gamma) \subseteq L^2(\Gamma)$. **Hint:** Recall the definition and the standard results on quotient spaces and the associated quotient norm!

For $X=H^1(\Omega)$ and $Y=L^2(\Omega)$, the following exercise shows that the L^2 -scalar products $(f;\cdot)_{L^2(\Omega)}$ for $f\in L^2(\Omega)$ give (up to density) all linear and continuous functionals on $H^1(\Omega)$, i.e., the embedding $L^2(\Omega)\to H^1(\Omega)^*$, $f\mapsto (f;\cdot)_{L^2(\Omega)}$ is well-defined, linear, continuous, and injective with dense image.

Exercise 10. Let X and Y be Hilbert spaces with continuous embedding $X \subseteq Y$. Show that the mapping $I: Y^* \to X^*, Iy^* := y^*|_X$ is well-defined, linear, and continuous. Prove that $I(Y^*) \subseteq X^*$ is a dense subspace. Moreover, if $X \subseteq Y$ is dense with respect to $\|\cdot\|_Y$, then the embedding I is even injective.

2.3 Weak Form of Laplace Problem

2.3.1 Dirichlet Problem

In this section, we generalize the variational form derived in the introductory section to our Hilbert space setting. We start with the homogeneous Dirichlet problem

$$-\Delta u = f \quad \text{in } \Omega,$$

$$u = 0 \quad \text{on } \Gamma.$$
 (2.12)

Recall that this formulation is called the **strong form** of the boundary value problem. The following proposition provides the — in some sense — equivalent and always uniquely solvable weak form of the boundary value problem.

Proposition 2.16. (i) Provided that $u \in C^2(\overline{\Omega})$ solves (2.12) for a given source term $f \in C(\overline{\Omega})$, it holds that $u \in H^1_0(\Omega)$ as well as

$$(\nabla u ; \nabla v)_{L^2(\Omega)} = (f ; v)_{L^2(\Omega)} \quad \text{for all } v \in H_0^1(\Omega).$$
 (2.13)

(ii) Given $f \in L^2(\Omega)$, the **weak form** (2.13) has a unique solution $u \in H_0^1(\Omega)$. It holds that

$$||u||_{H^{1}(\Omega)} \le C \sup_{v \in H_{0}^{1}(\Omega) \setminus \{0\}} \frac{(f; v)_{L^{2}(\Omega)}}{||v||_{H^{1}(\Omega)}} \le C ||f||_{L^{2}(\Omega)}, \tag{2.14}$$

where the constant C > 0 depends only on Ω .

(iii) Provided that $f \in C(\overline{\Omega})$ and that the weak solution $u \in H_0^1(\Omega)$ of (2.13) additionally satisfies $u \in C^2(\overline{\Omega})$, then u even solves the strong form (2.12).

Proof. (i) We have already seen before that a strong solution $u \in C^2(\overline{\Omega})$ solves the variational form (2.13) for test functions $v \in C_0^1(\overline{\Omega}) := \{w \in C^1(\overline{\Omega}) \mid w \mid_{\Gamma} = 0\}$ replacing $H_0^1(\Omega)$; see Proposition 1.1. If we keep u fixed, the left-hand side as well as the right-hand side of (2.13) define continuous and linear functionals on $H^1(\Omega)$. Note that the closure of $C_0^1(\overline{\Omega})$ with respect to the H^1 -norm leads to the Hilbert space $H_0^1(\Omega)$. Therefore, standard density arguments prove (2.13).

(ii) According to the Friedrichs inequality, it holds that

$$\|\nabla v\|_{L^2(\Omega)}^2 \le \|v\|_{H^1(\Omega)}^2 \le (1 + \widetilde{C}_F^2) \|\nabla v\|_{L^2(\Omega)}^2$$
 for all $v \in H_0^1(\Omega)$.

Therefore, the left-hand side of (2.13) defines an equivalent scalar product on $H_0^1(\Omega)$. The Riesz theorem thus provides a unique weak solution $u \in H_0^1(\Omega)$ of (2.13). Plugging-in $u = v \in H_0^1(\Omega)$, the weak form yields that

$$(1 + \widetilde{C}_F^2)^{-1} \|u\|_{H^1(\Omega)}^2 \le \|\nabla u\|_{L^2(\Omega)}^2 = (f \; ; \; u)_{L^2(\Omega)} \le \sup_{v \in H_0^1(\Omega) \setminus \{0\}} \frac{(f \; ; \; v)_{L^2(\Omega)}}{\|v\|_{H^1(\Omega)}} \|u\|_{H^1(\Omega)}$$

which results in the first estimate of (2.14). The second estimate follows from the Cauchy inequality

$$(f;v)_{L^2(\Omega)} \le ||f||_{L^2(\Omega)} ||v||_{L^2(\Omega)} \le ||f||_{L^2(\Omega)} ||v||_{H^1(\Omega)}.$$

(iii) Since the weak solution u is smooth, we may use integration by parts to see that

$$(\nabla u ; \nabla v)_{L^2(\Omega)} = (-\Delta u ; v)_{L^2(\Omega)}$$
 for all $v \in H_0^1(\Omega)$.

The difference with the weak form (2.13) thus yields that

$$0 = (f + \Delta u; v)_{L^2(\Omega)} \quad \text{for all } v \in H_0^1(\Omega).$$

Note that $F := f + \Delta u \in C(\overline{\Omega})$. With $\mathcal{D}(\Omega) \subseteq H_0^1(\Omega)$, Theorem 2.1 proves F = 0; see also the remark right after Theorem 2.1. Consequently, it holds that $-\Delta u = f$ in Ω . The Dirichlet boundary conditions (in the strong form) follow from $0 = \gamma u = u|_{\Gamma}$. Altogether, u solves (2.12)

2.3.2 Mixed Boundary Value Problem

Second, we consider the mixed boundary value problem

$$-\Delta u = f \quad \text{in } \Omega,$$

$$u = 0 \quad \text{on } \Gamma_D,$$

$$\partial u/\partial n = \phi \quad \text{on } \Gamma_N,$$
(2.15)

with $\Gamma = \overline{\Gamma}_D \cup \overline{\Gamma}_N$, $\Gamma_D \cap \Gamma_N = \emptyset$, and $|\Gamma_D| > 0$. The limit case $|\Gamma_D| = 0$ corresponds to the Neumann problem which is treated in Section 2.3.3. Recall the trace norm $\|\cdot\|_{H^{1/2}(\Gamma)}$ from Exercise 9. Then, the main proposition reads as follows:

Proposition 2.17. (i) Suppose that Γ_N is smooth, i.e., the outer normal vector depends continuously on $x \in \Gamma_N$. Provided that $u \in C^2(\overline{\Omega})$ solves the **strong form** (2.15) for a given source term $f \in C(\overline{\Omega})$ and Neumann data $\phi \in C(\overline{\Gamma}_N)$, it holds that $u \in H^1_D(\Omega)$ as well as

$$(\nabla u; \nabla v)_{L^2(\Omega)} = (f; v)_{L^2(\Omega)} + (\phi; \gamma v)_{L^2(\Gamma_N)} \quad \text{for all } v \in H^1_D(\Omega).$$
 (2.16)

(ii) Given $f \in L^2(\Omega)$ and $\phi \in L^2(\Gamma_N)$, the **weak form** (2.16) has a unique solution $u \in H^1_D(\Omega)$. It holds that

$$||u||_{H^{1}(\Omega)} \leq C_{1} \left(\sup_{v \in H^{1}_{D}(\Omega) \setminus \{0\}} \frac{(f ; v)_{L^{2}(\Omega)}}{||v||_{H^{1}(\Omega)}} + \sup_{w \in H^{1/2}(\Gamma) \setminus \{0\}} \frac{(\phi ; w)_{L^{2}(\Gamma_{N})}}{||w||_{H^{1/2}(\Gamma)}} \right)$$

$$\leq C_{2} \left(||f||_{L^{2}(\Omega)} + ||\phi||_{L^{2}(\Gamma_{N})} \right)$$

$$(2.17)$$

where the constants $C_1, C_2 > 0$ depend only on Ω and Γ_D .

(iii) Provided that $f \in C(\overline{\Omega})$ and $\phi \in C(\overline{\Gamma}_N)$ and that the weak solution $u \in H_D^1(\Omega)$ of (2.16) additionally satisfies $u \in C^2(\overline{\Omega})$, then u even solves the strong form (2.15).

Proof is done in the exercises.

2.3.3 Neumann Problem

Finally, we consider the Neumann problem

$$-\Delta u = f \quad \text{in } \Omega,$$

$$\partial u/\partial n = \phi \quad \text{on } \Gamma.$$
 (2.18)

Note that the solution u of (2.18) cannot be unique: If $u \in C^2(\overline{\Omega})$ solves the **strong form** (2.18), also u + c solves (2.18), for all $c \in \mathbb{R}$. To fix the additive constant, we seek a solution which additionally satisfies, e.g., that

$$\int_{\Omega} u \, dx = 0. \tag{2.19}$$

Moreover, the Gauss divergence theorem shows

$$-\int_{\Omega} f \, dx = \int_{\Omega} \Delta u \, dx = \int_{\Omega} \operatorname{div}(\nabla u) \, dx = \int_{\Gamma} \frac{\partial u}{\partial n} \, ds = \int_{\Gamma} \phi \, ds.$$

Therefore, the data f and ϕ have to satisfy the compatibility condition

$$\int_{\Omega} f \, dx + \int_{\Gamma} \phi \, ds = 0 \tag{2.20}$$

to allow for the existence of (strong) solutions. Recall the trace norm $\|\cdot\|_{H^{1/2}(\Gamma)}$ from Exercise 9.

Proposition 2.18. (i) Suppose that Γ is smooth, i.e., the outer normal vector depends continuously on $x \in \Gamma$. Provided that $u \in C^2(\overline{\Omega})$ solves (2.18) for a given source term $f \in C(\overline{\Omega})$ and Neumann data $\phi \in C(\Gamma)$, it holds that $u \in H^1(\Omega)$ and

$$(\nabla u \; ; \nabla v)_{L^{2}(\Omega)} = (f \; ; v)_{L^{2}(\Omega)} + (\phi \; ; \gamma v)_{L^{2}(\Gamma)} \quad \text{for all } v \in H^{1}(\Omega).$$
 (2.21)

(ii) Given $f \in L^2(\Omega)$ and $\phi \in L^2(\Gamma)$, the variational formulation

$$(\nabla u; \nabla v)_{L^2(\Omega)} = (f; v)_{L^2(\Omega)} + (\phi; \gamma v)_{L^2(\Gamma)} \quad \text{for all } v \in H^1_*(\Omega)$$
 (2.22)

has a unique solution $u \in H^1_*(\Omega) := \{v \in H^1(\Omega) \mid \int_{\Omega} v \, dx = 0\}.$

(iii) Provided that the data $f \in L^2(\Omega)$ and $\phi \in L^2(\Gamma)$ satisfy (2.20), the unique solution $u \in H^1_*(\Omega)$ of (2.22) even solves the **weak form** (2.21). Moreover, it holds that

$$||u||_{H^{1}(\Omega)} \leq C_{1} \left(\sup_{v \in H^{1}(\Omega) \setminus \{0\}} \frac{(f ; v)_{L^{2}(\Omega)}}{||v||_{H^{1}(\Omega)}} + \sup_{w \in H^{1/2}(\Gamma) \setminus \{0\}} \frac{(\phi ; w)_{L^{2}(\Gamma)}}{||w||_{H^{1/2}(\Gamma)}} \right)$$

$$\leq C_{2} \left(||f||_{L^{2}(\Omega)} + ||\phi||_{L^{2}(\Gamma)} \right)$$

$$(2.23)$$

where the constants $C_1, C_2 > 0$ depend only on Ω .

(iv) Provided that $f \in C(\overline{\Omega})$ and $\phi \in C(\Gamma)$ satisfy (2.20) and that the weak solution $u \in H^1_*(\Omega)$ of (2.21) resp. (2.22) additionally satisfies $u \in C^2(\overline{\Omega})$, then u even solves the strong form (2.18).

Proof. (i) The variational form (2.21) holds for test functions $v \in C^1(\overline{\Omega})$ according to integration by parts. For fixed u, the left-hand as well as the right-hand side define continuous linear functionals on $H^1(\Omega)$. Thus, (2.21) follows for $v \in H^1(\Omega)$ by density arguments. (ii) According to the Poincaré inequality, it holds that

$$\|\nabla v\|_{L^2(\Omega)}^2 \le \|v\|_{H^1(\Omega)}^2 \le (1 + \widetilde{C}_P^2) \|\nabla v\|_{L^2(\Omega)}^2$$
 for all $v \in H^1_*(\Omega)$.

Therefore, the left-hand side of (2.22) defines an equivalent scalar product on $H^1_*(\Omega)$. Note that $H^1_*(\Omega)$ is a closed subspace of $H^1(\Omega)$ and hence a Hilbert space. Therefore, (2.22) follows from the Riesz theorem. (iii) For a function $v \in H^1(\Omega)$, we define $\tilde{v} := v - v_{\Omega} \in H^1_*(\Omega)$, where $v_{\Omega} \in \mathbb{R}$ denotes the integral mean $v_{\Omega} := (1/|\Omega|) \int_{\Omega} v \, dx \in \mathbb{R}$. Note that (2.20) implies that

$$(f; v_{\Omega})_{L^{2}(\Omega)} + (\phi; v_{\Omega})_{L^{2}(\Gamma)} = 0.$$

Thus, (2.22) proves that

$$(\nabla u \; ; \; \nabla v)_{L^{2}(\Omega)} = (\nabla u \; ; \; \nabla \widetilde{v})_{L^{2}(\Omega)} = (f \; ; \; \widetilde{v})_{L^{2}(\Omega)} + (\phi \; ; \; \gamma \widetilde{v})_{L^{2}(\Gamma)} = (f \; ; \; v)_{L^{2}(\Omega)} + (\phi \; ; \; \gamma v)_{L^{2}(\Gamma)},$$

i.e., u even solves (2.21). Plugging-in u = v, we see that

$$\|\nabla u\|_{L^{2}(\Omega)}^{2} \leq \sup_{v \in H^{1}(\Omega) \setminus \{0\}} \frac{(f ; v)_{L^{2}(\Omega)}}{\|v\|_{H^{1}(\Omega)}} \|u\|_{H^{1}(\Omega)} + \sup_{w \in H^{1/2}(\Gamma) \setminus \{0\}} \frac{(\phi ; w)_{L^{2}(\Gamma)}}{\|w\|_{H^{1/2}(\Gamma)}} \|\gamma u\|_{H^{1/2}(\Gamma)},$$

where we have used that $H^{1/2}(\Gamma) = \operatorname{range}(\gamma)$. Note that the $H^{1/2}$ -norm is defined in such a way that $\gamma \in L(H^1(\Omega); H^{1/2}(\Gamma))$ with $\|\gamma u\|_{H^{1/2}(\Gamma)} \leq \|u\|_{H^1(\Omega)}$. Therefore,

$$\|\nabla u\|_{L^{2}(\Omega)}^{2} \leq \|u\|_{H^{1}(\Omega)} \left(\sup_{v \in H^{1}(\Omega) \setminus \{0\}} \frac{(f ; v)_{L^{2}(\Omega)}}{\|v\|_{H^{1}(\Omega)}} + \sup_{w \in H^{1/2}(\Gamma) \setminus \{0\}} \frac{(\phi ; w)_{L^{2}(\Gamma)}}{\|w\|_{H^{1/2}(\Gamma)}} \right).$$

Together with $(1+\widetilde{C}_P^2)^{-1}\|u\|_{H^1(\Omega)}^2 \leq \|\nabla u\|_{L^2(\Omega)}^2$, this proves the first estimate in (2.23). As above, the first supremum may be estimated by $\|f\|_{L^2(\Omega)}$. With the continuous embedding $H^{1/2}(\Gamma) \subset L^2(\Gamma)$, the numerator of the second supremum can be dominated by

$$(\phi; w)_{L^{2}(\Gamma)} \leq \|\phi\|_{L^{2}(\Gamma)} \|w\|_{L^{2}(\Gamma)} \leq \widetilde{C} \|\phi\|_{L^{2}(\Gamma)} \|w\|_{H^{1/2}(\Gamma)}.$$

This provides the upper bound $\widetilde{C} \|\phi\|_{L^2(\Gamma)}$ for the second supremum. (iv) As above, we may use integration by parts to see that

$$(f + \Delta u; v)_{L^2(\Omega)} + (\phi - \partial u/\partial n; \gamma v)_{L^2(\Gamma)} = 0$$
 for all $v \in H^1(\Omega)$.

From this, we first conclude $f = -\Delta u$ by use of Theorem 2.1 for test functions $v \in \mathcal{D}(\Omega) \subset H_0^1(\Omega) \subset H^1(\Omega)$. To prove $\phi = \partial u/\partial n$, one proceeds analogously to the remark right after Theorem 2.1.

Chapter 3

A Priori Analysis

3.1 P1-Finite Element Method in 2D

A set $T \subset \mathbb{R}^2$ is called a **non-degenerate triangle** provided that there are nodes $x_T, y_T, z_T \in \mathbb{R}^2$ with $T = \text{conv}\{x_T, y_T, z_T\}$ and provided that |T| > 0, i.e., T has positive measure. We note that T is in particular bounded and closed, whence compact. We denote by

$$\mathcal{K}_T := \{x_T, y_T, z_T\} \tag{3.1}$$

the **set of nodes** of T and by

$$\mathcal{E}_T := \{ \text{conv}\{x_T, y_T\}, \text{conv}\{y_T, z_T\}, \text{conv}\{z_T, x_T\} \}$$
 (3.2)

the **set of edges** of T. The **diameter** of T is denoted by

$$h_T := \operatorname{diam}(T) := \max\{|x - y| \mid x, y \in T\}.$$
 (3.3)

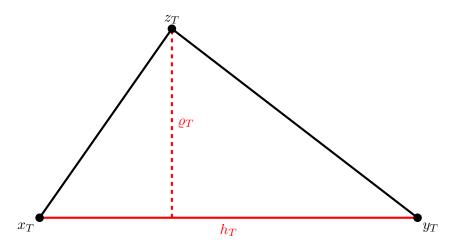


FIGURE 3.1. The diameter h_T of the triangle T is the length of the longest edge (possibly non unique). The quantity ϱ_T denotes the corresponding height.

Moreover, we define the **edge length**

$$h_E := \operatorname{diam}(E) := \max\{|x - y| \mid x, y \in E\}$$
 (3.4)

for all edges $E \in \mathcal{E}_T$. Clearly, the diameter h_T of a triangle is the length of the longest edge (possibly non unique), i.e., there is some $E \in \mathcal{E}_T$ with $h_T = h_E$. The **height** over the longest edge E of T is denoted by ϱ_T , cf. Figure 3.1. Recall that the measure of the triangle reads

$$|T| = \frac{h_T \rho_T}{2}. (3.5)$$

The most important example is the **reference triangle**

$$T_{\text{ref}} := \text{conv}\{(0,0), (1,0), (0,1)\}$$
 (3.6)

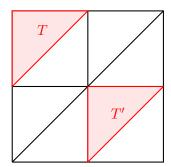
which has measure $|T_{\text{ref}}| = 1/2$.

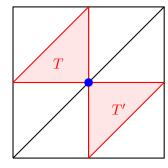
Exercise 11. Give a formal proof that the diameter of a triangle T is the length of one longest edge, i.e., $h_T = \max_{E \in \mathcal{E}_T} h_E$. Hint: Use that the convex hull $\operatorname{conv}(M) := \bigcap \{\widehat{M} \subseteq \mathbb{R}^d \mid \widehat{M} \text{ is convex with } M \subseteq \widehat{M} \}$ of a set $M \subseteq \mathbb{R}^d$ is also characterized by $\operatorname{conv}(M) = \{\sum_{j=1}^N \lambda_j x_j \mid N \in \mathbb{N}, x_j \in M, \lambda_j \geq 0 \text{ with } \sum_{j=1}^N \lambda_j = 1\}$. The proof then directly applies to general simplices in \mathbb{R}^d , i.e., $T = \operatorname{conv}\{x_0, \dots, x_d\} \subset \mathbb{R}^d$.

Definition. A set \mathcal{T} is a **triangulation** of Ω (consisting of triangles) if and only if

- \bullet \mathcal{T} is a finite set of non-degenerate triangles,
- the closure of Ω is covered by \mathcal{T} , i.e., $\overline{\Omega} = \bigcup \mathcal{T}$,
- for all $T, T' \in \mathcal{T}$ with $T \neq T'$, it holds that $|T \cap T'| = 0$, i.e., the overlap is a set of measure zero.

By $\mathcal{K} := \bigcup \{x \in \mathcal{K}_T \mid T \in \mathcal{T}\}$, we then denote the **set of nodes** of the triangulation \mathcal{T} and by $\mathcal{E} := \bigcup \{E \in \mathcal{E}_T \mid T \in \mathcal{T}\}$ the **set of edges** of the triangulation \mathcal{T} . A triangulation of Ω is called **conforming** or **regular** (in the sense of Ciarlet) provided that the intersection of two elements $T, T' \in \mathcal{T}$ with $T \neq T'$ is





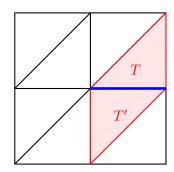


FIGURE 3.2. For a regular triangulation \mathcal{T} , the intersection of two elements $T \neq T'$ is either empty, a joint node, or a joint edge.

- either empty,
- or a joint node, i.e., $T \cap T' = \{z\} = \mathcal{K}_T \cap \mathcal{K}_{T'}$,
- or a joint edge, i.e., $E := T \cap T' \in \mathcal{E}_T \cap \mathcal{E}_{T'}$,

cf. Figure 3.2. According to this regularity assumption, an edge $E \in \mathcal{E}$ with surface measure $|E \cap \Gamma| > 0$ automatically satisfies $E \subseteq \Gamma$, i.e., an edge E is either a boundary edge or an interior edge. Additionally, we always assume that a regular triangulation resolves the boundary conditions: If $\Gamma = \partial \Omega$ is partitioned into Dirichlet and Neumann boundary Γ_D and Γ_N , respectively, each boundary edge $E \in \mathcal{E}$ with $E \subseteq \Gamma$ satisfies

- either $E \subseteq \overline{\Gamma}_D$
- or $E \subseteq \overline{\Gamma}_N$.

With this assumption, we define the (disjoint) sets of boundary edges

$$\mathcal{E}_D := \left\{ E \in \mathcal{E} \mid E \subseteq \overline{\Gamma}_D \right\} \quad \text{and} \quad \mathcal{E}_N := \left\{ E \in \mathcal{E} \mid E \subseteq \overline{\Gamma}_N \right\} \tag{3.7}$$

as well as the set of all interior edges

$$\mathcal{E}_{\Omega} := \mathcal{E} \setminus (\mathcal{E}_D \cup \mathcal{E}_N). \tag{3.8}$$

We finally note that, for each $E \in \mathcal{E}_{\Omega}$, there are two elements $T, T' \in \mathcal{T}$ with $E = T \cap T'$.

Exercise 12. Let \mathcal{T} be a regular triangulation of Ω and $v:\Omega\to\mathbb{R}$ such that $v|_T\in C^1(T)$ for all $T\in\mathcal{T}$. Prove that $v\in H^1(\Omega)$ if and only if $v\in C(\Omega)$.

The following proposition essentially follows from the regularity of the triangulation \mathcal{T} .

Proposition 3.1. For a regular triangulation \mathcal{T} of Ω , we define the discrete space

$$S^{1}(\mathcal{T}) := \{ v_{h} \in C(\Omega) \mid \forall T \in \mathcal{T} \quad v_{h}|_{T} \text{ affine} \}$$
(3.9)

of all T-piecewise affine and globally continuous functions. Then, there holds the following:

- (i) $S^1(\mathcal{T})$ is an N-dimensional subspace of $H^1(\Omega)$ with $N = \#\mathcal{K}$ the number of nodes.
- (ii) For each node $z \in \mathcal{K}$, there is a unique **hat function**

$$\zeta_z \in \mathcal{S}^1(\mathcal{T}) \quad with \quad \zeta_z(z') = \delta_{zz'} \quad for \ all \ z' \in \mathcal{K}.$$
 (3.10)

(iii) The set $\mathcal{B} := \{ \zeta_z \mid z \in \mathcal{K} \}$ is a basis of $\mathcal{S}^1(\mathcal{T})$, the so-called **nodal basis**.

Proof. 1. step. According to the regularity of \mathcal{T} , hat functions ζ_z are automatically continuous on Ω : For each element $T \in \mathcal{T}$, an affine function $v_h : T \to \mathbb{R}$ is uniquely determined by the nodal values $v_h(z)$ for $z \in \mathcal{K}_T$. Therefore, the \mathcal{T} -piecewise affine hat function ζ_z defined by $\zeta_z(z') = \delta_{zz'}$ is uniquely defined. We now show that $\zeta_z \in C(\Omega)$: If $T, T' \in \mathcal{T}$ are elements with $T \cap T' \neq \emptyset$, regularity of \mathcal{T} implies that either T = T' or $\{z'\} = T \cap T'$ is a joint point or $E = T \cap T'$ is a joint edge. In the latter case, note that the trace on E of the affine function $\zeta_z|_T$ as well as of $\zeta_z|_{T'}$ is

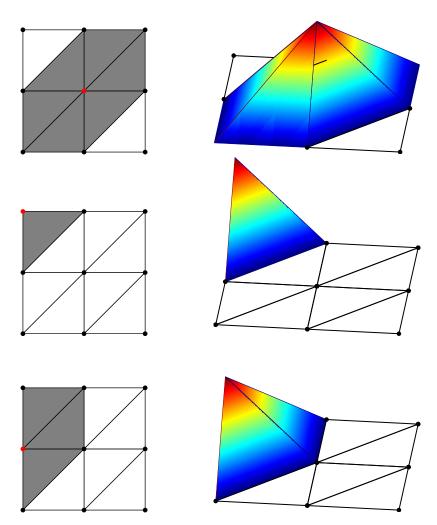


FIGURE 3.3. Examples of P1 hat functions ζ_z : The left figures show the mesh as well as the support $\operatorname{supp}(\zeta_z)$ in grey, where the corresponding node $z \in \mathcal{K}$ is indicated in red. The right figures show the plots of the hat functions. Triangles $T \in \mathcal{T}$ with $\zeta_z|_{T} = 0$ are filled with white.

uniquely defined on the edge E by the nodal values $\zeta_z(x_E)$ and $\zeta_z(y_E)$, where $E = \text{conv}\{x_E, y_E\}$. Therefore the traces of $\zeta_z|_T$ and $\zeta_z|_{T'}$ on E coincide, i.e., ζ_z is continuous on each interior edge.

2. step. The nodal basis \mathcal{B} is a basis of $\mathcal{S}^1(\mathcal{T})$ and dim $\mathcal{S}^1(\mathcal{T}) = \#\mathcal{K}$: Clearly, the hat functions are linearly independent, $\mathcal{B} \subseteq \mathcal{S}^1(\mathcal{T})$, and $\#\mathcal{B} = \#\mathcal{K}$. Moreover, each function $v_h \in \mathcal{S}^1(\mathcal{T})$ is uniquely defined by the nodal values $v_h(z)$ for $z \in \mathcal{K}$ and can thus be written as the linear combination of the hat functions, i.e., $\mathcal{S}^1(\mathcal{T}) \subseteq \operatorname{span}(\mathcal{B})$.

3. step. The inclusion $S^1(\mathcal{T}) \subset H^1(\Omega)$ follows from Exercise 12.

Remark. Examples for hat functions ζ_z are shown in Figure 3.3. Note that the support supp (ζ_z) is always local. This leads to a sparse Galerkin matrix A, i.e., most of the entries of A are zero. \Box

For a given Dirichlet boundary $\Gamma_D \subseteq \Gamma$, we use the discrete space $\mathcal{S}_D^1(\mathcal{T})$ to discretize the weak form of the mixed boundary value problem. In case of $\Gamma_D = \Gamma$, we consider the space $\mathcal{S}_0^1(\mathcal{T})$.

Corollary 3.2. Let \mathcal{T} be a regular triangulation of Ω . Then, the space

$$S_D^1(\mathcal{T}) := \left\{ v_h \in S^1(\mathcal{T}) \mid \forall z \in \mathcal{K} \cap \overline{\Gamma}_D \quad v_h(z) = 0 \right\}$$
(3.11)

is a finite dimensional subspace of $H_D^1(\Omega)$ of dimension $\#\{z \in \mathcal{K} \mid z \notin \overline{\Gamma}_D\}$. The space

$$S_0^1(\mathcal{T}) := \left\{ v_h \in S^1(\mathcal{T}) \middle| \forall z \in \mathcal{K} \cap \Gamma \quad v_h(z) = 0 \right\}$$
(3.12)

is a finite dimensional subspace of $H_0^1(\Omega)$ of dimension $\#\{z \in \mathcal{K} \mid z \notin \Gamma\}$.

Proof. We only need to show that $v_h|_{\Gamma_D} = 0$ for $v_h \in \mathcal{S}_D^1(\mathcal{T})$. Let $x \in \Gamma_D$. According to the regularity of \mathcal{T} , there is an edge $E \in \mathcal{E}_D$ such that $x \in E$. Since the trace $v_h|_E$ is affine, it is uniquely determined by the nodal values $v_h(x_T) = 0 = v_h(y_T)$, where $E = \text{conv}\{x_T, y_T\}$. Consequently, $v_h|_E = 0$ for all $E \in \mathcal{E}_D$ and hence $v_h \in H_D^1(\Omega)$. In particular, we obtain the claim for $\mathcal{S}_0^1(\mathcal{T})$ in case of $\Gamma_D = \Gamma$.

For the discretization of the Neumann problem, we are dealing with $\mathcal{S}^1_*(\mathcal{T})$.

Corollary 3.3. For a regular triangulation \mathcal{T} of Ω , the space

$$S_*^1(\mathcal{T}) := \left\{ v_h \in S^1(\mathcal{T}) \mid \int_{\Omega} v_h \, dx = 0 \right\} \tag{3.13}$$

is a finite dimensional subspace of $H^1_*(\Omega)$ of dimension $\#\mathcal{K}-1$.

Proof. Clearly, it holds that $\mathcal{S}^1_*(\mathcal{T}) \subseteq H^1_*(\Omega)$. Note that $I(v_h) := \int_{\Omega} v_h \, dx$ is a linear functional on $\mathcal{S}^1(\mathcal{T})$ with kernel $\mathcal{S}^1_*(\mathcal{T}) = \ker(I)$. Since $\operatorname{rank}(I) = 1$, Linear Algebra yields that $\dim \mathcal{S}^1_*(\mathcal{T}) = \dim \mathcal{S}^1(\mathcal{T}) - 1$.

The **P1 Finite Element Method** now consists of using the Galerkin method with the discrete spaces $S_0^1(\mathcal{T})$, $S_D^1(\mathcal{T})$, and $S_*^1(\mathcal{T})$ to approximate the weak solution of the Dirichlet problem, the mixed boundary value problem, and the Neumann problem, respectively. From now on, we shall assume that \mathcal{T} is a regular triangulation of Ω . We start with the **Dirichlet problem**

$$-\Delta u = f \quad \text{in } \Omega,$$

$$u = 0 \quad \text{on } \Gamma,$$

for given data $f \in L^2(\Omega)$. The P1-FEM then reads: Find $u_h \in \mathcal{S}_0^1(\mathcal{T})$ such that

$$(\nabla u_h ; \nabla v_h)_{L^2(\Omega)} = (f ; v_h)_{L^2(\Omega)} \quad \text{for all } v_h \in \mathcal{S}_0^1(\mathcal{T}). \tag{3.14}$$

Second, the mixed boundary value problem reads

$$\begin{split} -\Delta u &= f & \text{in } \Omega, \\ u &= 0 & \text{on } \Gamma_D, \\ \partial u / \partial n &= \phi & \text{on } \Gamma_N, \end{split}$$

with $\Gamma = \overline{\Gamma}_D \cup \overline{\Gamma}_N$, $\Gamma_D \cap \Gamma_N = \emptyset$, and $|\Gamma_D| > 0$. The data satisfy $f \in L^2(\Omega)$ and $\phi \in L^2(\Gamma_N)$. The P1-FEM for the mixed BVP reads: Find $u_h \in \mathcal{S}_D^1(\mathcal{T})$ such that

$$(\nabla u_h ; \nabla v_h)_{L^2(\Omega)} = (f ; v_h)_{L^2(\Omega)} + (\phi ; v_h)_{L^2(\Gamma_N)} \quad \text{for all } v_h \in \mathcal{S}_D^1(\mathcal{T}). \tag{3.15}$$

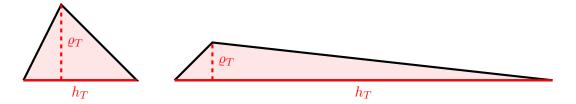


FIGURE 3.4. two triangles with small (left) and large (right) shape regularity constant $\sigma(T) := h_T/\varrho_T$

Finally, we consider the **Neumann problem**

$$\begin{split} -\Delta u &= f & \text{ in } \Omega, \\ \partial u / \partial n &= \phi & \text{ on } \Gamma, \end{split}$$

where the data $f \in L^2(\Omega)$ and $\phi \in L^2(\Gamma)$ are assumed to satisfy $\int_{\Omega} f \, dx + \int_{\Gamma} \phi \, ds = 0$. The P1-FEM for the Neumann problem reads: Find $u_h \in \mathcal{S}^1_*(\mathcal{T})$ such that

$$(\nabla u_h ; \nabla v_h)_{L^2(\Omega)} = (f ; v_h)_{L^2(\Omega)} + (\phi ; v_h)_{L^2(\Gamma)} \quad \text{for all } v_h \in \mathcal{S}^1_*(\mathcal{T}). \tag{3.16}$$

3.2 Approximation Theorem and Bramble-Hilbert Lemma

3.2.1 Uniform Mesh-Refinement and Shape Regularity

Let $h \in L^{\infty}(\Omega)$ and $\rho \in L^{\infty}(\Omega)$ denote the **local mesh-width** functions which are defined by

$$h|_T := h_T = \operatorname{diam}(T) \quad \text{and} \quad \rho|_T := \rho_T \quad \text{for all } T \in \mathcal{T}.$$
 (3.17)

Moreover, the quantities

$$\sigma(T) := \frac{h_T}{\varrho_T} \quad \text{and} \quad \sigma(T) := \|h/\varrho\|_{L^{\infty}(\Omega)} = \max_{T \in \mathcal{T}} \frac{h_T}{\varrho_T} \ge 1$$
 (3.18)

denote the **shape regularity constant** of an element $T \in \mathcal{T}$ resp. the triangulation \mathcal{T} , see Fig. 3.4. Note that $|T| = h_T \varrho_T / 2$ so that $2 h_T / \varrho_T = h_T^2 / |T|$. The shape regularity constant will affect all error estimates, so that mesh-refinement has to avoid a blow-up of $\sigma(\mathcal{T})$. We say that a regular mesh \mathcal{T} is γ -shape regular, if $\sigma(\mathcal{T}) \leq \gamma < \infty$.

For this section, we stick with the so-called **uniform mesh-refinement**: Given a regular triangulation $\mathcal{T}^{(\text{old})}$, we obtain a new triangulation $\mathcal{T}^{(\text{new})}$ as follows: Each element $T \in \mathcal{T}^{(\text{old})}$ is split into 4 similar triangles $T_1, \ldots, T_4 \in \mathcal{T}^{(\text{new})}$, cf. Figure 3.5. Therefore, each node $z \in \mathcal{K}^{(\text{new})}$ either belongs to $\mathcal{K}^{(\text{old})}$ or is the midpoint of an edge $E \in \mathcal{E}^{(\text{old})}$. We stress some simple observations:

- The new triangulation $\mathcal{T}^{(\text{new})}$ is also regular.
- The local mesh-width functions satisfy $h^{\text{(new)}} = h^{\text{(old)}}/2$ and $\varrho^{\text{(new)}} = \varrho^{\text{(old)}}/2$.
- In particular, the shape regularity constant satisfies that $\sigma(\mathcal{T}^{(\text{old})}) = \sigma(\mathcal{T}^{(\text{new})})$.

Further mesh-refinement strategies are discussed in the following section.

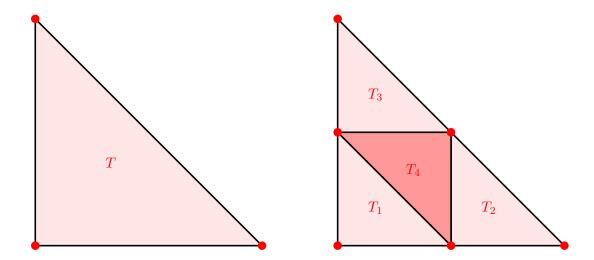


FIGURE 3.5. Red-refinement refines the element $T \in \mathcal{T}^{(\text{old})}$ into 4 similar elements $T_1, \ldots, T_4 \in \mathcal{T}^{(\text{new})}$. The new nodes $\mathcal{K}^{(\text{new})} \setminus \mathcal{K}^{(\text{old})}$ are just the edge midpoints for all edges $E \in \mathcal{E}^{(\text{old})}$. In particular, regularity of $\mathcal{T}^{(\text{old})}$ implies regularity of $\mathcal{T}^{(\text{new})}$.

Exercise 13. Let $T = \text{conv}\{z_1, z_2, z_3\}$ be a non-degenerate triangle in \mathbb{R}^2 . Prove that the shape regularity constant h_T/ϱ_T tends to infinity if and only if the smallest angle in T tends to zero.

Exercise 14. Often, the shape regularity constant is defined as the maximal quotient h_T/r_T , where $r_T > 0$ denotes the maximal radius of a ball $B(x, r_T) := \{y \in \mathbb{R}^2 \mid |x - y| \le r_T\}$ inscribed in T, i.e., $B(x, r_T) \subseteq T$. Let $T = \text{conv}\{z_1, z_2, z_3\}$ be a non-degenerate triangle in \mathbb{R}^2 . What is the relation between ρ_T and r_T ?

3.2.2 Statement and Interpretation of Approximation Theorem

To state our first main result in this section, we need to know that certain Sobolev functions are at least continuous.

Theorem 3.4 (Sobolev). Let Ω be a Lipschitz domain in \mathbb{R}^d and m > d/2. Then, there holds the continuous inclusion $H^m(\Omega) \subseteq C(\overline{\Omega})$.

In particular, for d=2,3, each Sobolev function $u \in H^2(\Omega)$ is continuous so that evaluation of u at the nodes $z \in \mathcal{K}$ is well-defined. Throughout the remaining section, we assume that \mathcal{T} is a regular triangulation of a bounded Lipschitz domain $\Omega \subset \mathbb{R}^2$. We stress, however, that the same results — even with the same proofs — hold for d=3 as well. As in the previous section, the nodal basis function corresponding to a node $z \in \mathcal{K}$ is denoted by $\zeta_z \in \mathcal{S}^1(\mathcal{T})$.

Theorem 3.5 (Approximation Theorem). For $u \in H^2(\Omega)$, the nodal interpolant reads

$$I_h u := \sum_{z \in \mathcal{K}} u(z) \zeta_z \in \mathcal{S}^1(\mathcal{T}). \tag{3.19}$$

For all $T \in \mathcal{T}$, there hold the elementwise error estimates

$$||u - I_h u||_{L^2(T)} \le C ||h^2 D^2 u||_{L^2(T)}$$
(3.20)

and

$$\|\nabla(u - I_h u)\|_{L^2(T)} \le C \,\sigma(T) \,\|h D^2 u\|_{L^2(T)},\tag{3.21}$$

where the generic constant C > 0 is independent of u, \mathcal{T} , and Ω , but depends only on the reference triangle. In particular, this proves for all $\alpha \in \mathbb{R}$ the global error estimates

$$||h^{\alpha}(u - I_h u)||_{L^2(\Omega)} \le C ||h^{2+\alpha} D^2 u||_{L^2(\Omega)}$$
(3.22)

and

$$||h^{\alpha}\nabla(u - I_h u)||_{L^2(\Omega)} \le C \,\sigma(\mathcal{T}) \,||h^{1+\alpha}D^2 u||_{L^2(\Omega)}.$$
 (3.23)

This theorem will be shown later. First, we discuss the following immediate consequence:

Corollary 3.6. For $u \in H^2(\Omega) \cap H^1_D(\Omega)$, it holds that $I_h u \in \mathcal{S}^1_D(\mathcal{T})$ and thus

$$\min_{v_h \in \mathcal{S}_D^1(\mathcal{T})} \|u - v_h\|_{H^1(\Omega)} \le \|u - I_h u\|_{H^1(\Omega)} \le C \,\sigma(\mathcal{T}) \,\|h D^2 u\|_{L^2(\Omega)}. \tag{3.24}$$

For $u \in H^2(\Omega) \cap H^1_*(\Omega)$, it holds that

$$\min_{v_h \in \mathcal{S}^1_*(\mathcal{T})} \|u - v_h\|_{H^1(\Omega)} = \min_{v_h \in \mathcal{S}^1(\mathcal{T})} \|u - v_h\|_{H^1(\Omega)} \le \|u - I_h u\|_{H^1(\Omega)}
\le C \,\sigma(\mathcal{T}) \,\|h D^2 u\|_{L^2(\Omega)}.$$
(3.25)

In either case, the constant C > 0 depends only on diam(Ω).

Proof. Let $C_{\text{apx}} > 0$ denote the constant from the approximation theorem. Then,

$$||u - I_h u||_{H^1(\Omega)}^2 = ||u - I_h u||_{L^2(\Omega)}^2 + ||\nabla (u - I_h u)||_{L^2(\Omega)}^2 \le C_{\text{apx}}^2(\text{diam}(\Omega)^2 + \sigma(\mathcal{T})^2)||hD^2 u||_{L^2(\Omega)}^2.$$

Since $\sigma(\mathcal{T}) > 1$, we obtain that

$$||u - I_h u||_{H^1(\Omega)} \le C_{\text{apx}} \sigma(\mathcal{T}) (\text{diam}(\Omega)^2 + 1)^{1/2} ||h D^2 u||_{L^2(\Omega)}.$$

For $u \in H^2(\Omega) \cap H^1_D(\Omega)$, it holds that u(z) = 0 for all $z \in \overline{\Gamma}_D$. This implies that $I_h u \in \mathcal{S}^1_D(\mathcal{T})$ and hence (3.24). Before we prove (3.25), note that $I_h u \in \mathcal{S}^1(\mathcal{T})$ does not belong to $\mathcal{S}^1_*(\mathcal{T})$ in general. However, let $\mathbb{P}_h : H^1(\Omega) \to \mathcal{S}^1(\mathcal{T})$ denote the H^1 -orthogonal projection onto $\mathcal{S}^1(\mathcal{T})$. Since $1 \in \mathcal{S}^1(\mathcal{T})$, it holds that

$$0 = \int_{\Omega} u \, dx = (u \, ; \, 1)_{H^{1}(\Omega)} = (\mathbb{P}_{h} u \, ; \, 1)_{H^{1}(\Omega)} = \int_{\Omega} \mathbb{P}_{h} u \, dx \quad \text{for all } u \in H^{1}_{*}(\Omega).$$

Therefore, $\mathbb{P}_h u \in \mathcal{S}^1_*(\mathcal{T})$, and the best approximation property of the orthogonal projection \mathbb{P}_h thus implies that

$$||u - \mathbb{P}_h u||_{H^1(\Omega)} = \min_{v_h \in \mathcal{S}^1(\mathcal{T})} ||u - v_h||_{H^1(\Omega)} \le \min_{v_h \in \mathcal{S}^1_*(\mathcal{T})} ||u - v_h||_{H^1(\Omega)} \le ||u - \mathbb{P}_h u||_{H^1(\Omega)}$$

and hence equality. As before, this proves (3.25).

Remark. Corollary 3.6 has two important consequences: First, according to Céa's lemma, the Galerkin error is up to a constant the best approximation error. For a smooth exact solution $u \in H^2(\Omega)$, the P1-FEM thus leads (at least and in fact even) to a convergence order $\mathcal{O}(h)$. Second, $C_D^{\infty}(\overline{\Omega})$ is dense in $H_D^1(\Omega)$ and $C_*^{\infty}(\overline{\Omega}) := \{v \in C^{\infty}(\overline{\Omega}) \mid \int_{\Omega} v \, dx = 0\}$ is dense in $H_*^1(\Omega)$. Corollary 3.6 therefore implies convergence of the Galerkin scheme on a dense subspace. The abstract framework provides convergence of the P1-FEM even without any regularity assumptions on u, cf. Proposition 1.7.

Exercise 15. Use the Poincaré inequality and the Meyers-Serrin theorem to prove that $C_*^{\infty}(\overline{\Omega})$ is dense in $H_*^1(\Omega)$.

3.2.3 Bramble-Hilbert Lemma

It now remains to prove the Approximation Theorem 3.5. The proof of which needs three lemmata. The first two lemmata provide the basis for general scaling arguments. We therefore state the results even in a slightly generalized setting.

Definition. For a multiindex $\alpha \in \mathbb{N}_0^d$ and $x \in \mathbb{R}^d$, we define the **monomial** $x^{\alpha} := \prod_{j=1}^d x_j^{\alpha_j}$, where $|\alpha| := \sum_{j=1}^d \alpha_j$ is the **(total) degree** of α . For a Lipschitz domain $T \subseteq \mathbb{R}^d$, we define

$$\mathcal{P}^{m}(T) := \left\{ v : T \to \mathbb{R} \mid v \text{ is linear combination of monomials of degree } \le m \right\}$$
 (3.26)

the space that consists of all **polynomials** of degree less than or equal to $m \in \mathbb{N}$.

Lemma 3.7 (Bramble-Hilbert). For a Lipschitz domain $T \subset \mathbb{R}^d$ and a normed space X, let $A \in L(H^{m+1}(T); X)$ be a linear and continuous operator with $\mathcal{P}^m(T) \subseteq \ker(A)$. Besides the classical continuity estimate

$$||Av||_X \le ||A|| ||v||_{H^{m+1}(T)} \quad \text{for all } v \in H^{m+1}(T),$$
 (3.27)

it holds that

$$||Av||_X \le C ||A|| ||D^{m+1}v||_{L^2(T)} \quad for \ all \ v \in H^{m+1}(T),$$
 (3.28)

where the constant C > 0 depends only on m and T.

Proof. 1. step. Construct an equivalent norm on $H^{m+1}(T)$: Note that $\mathcal{P}^m(T)$ is a finite dimensional space. Let $\Pi: L^2(T) \to \mathcal{P}^m(T)$ denote the L^2 -orthogonal projection onto $\mathcal{P}^m(T)$. We define

$$||v|| := ||D^{m+1}v||_{L^2(T)} + ||\Pi v||_{L^2(T)}$$
 for $v \in H^{m+1}(T)$.

From $\|\Pi v\|_{L^2(T)} \leq \|v\|_{L^2(T)}$, we infer that

$$|||v||| \le ||D^{m+1}v||_{L^2(T)} + ||v||_{L^2(T)} \le \sqrt{2} ||v||_{H^{m+1}(T)}.$$

Next, we prove the converse inequality, i.e., there exists a constant C > 0 such that

$$||v||_{H^{m+1}(T)} \le C |||v||| \text{ for all } v \in H^{m+1}(T).$$

As above, we use the Rellich theorem and argue by contradiction: If the claim is wrong, we find $v_n \in H^{m+1}(T)$ such that $||v_n||_{H^{m+1}(T)} > n|||v_n|||$. We define $w_n := v_n/||v_n||_{H^{m+1}(T)}$. Note that

$$||w_n||_{H^{m+1}(T)} = 1$$
 as well as $|||w_n||| \le \frac{1}{n}$.

According to reflexivity, we may thus assume that $w_n \to w \in H^{m+1}(T)$. According to Lemma 2.6, convexity and continuity of $\|\cdot\|$ imply that $\|w\| = 0$. Therefore, it holds that $D^{m+1}w = 0$ as well as $\Pi w = 0$. With the help of Exercise 16, we deduce that $w \in \mathcal{P}^m(T)$ and consequently $\|w\|_{L^2(T)} = \|\Pi w\|_{L^2(T)} = 0$. According to Rellich's theorem, we have $w_n \to w = 0 \in H^m(T)$. Since $D^{m+1}w_n \to 0 \in L^2(T)$, we even conclude that $w_n \to 0 = w \in H^{m+1}(T)$. This however, contradicts $\|w_n\|_{H^{m+1}(T)} = 1$. Altogether, we have shown that $\|\cdot\|$ is an equivalent norm on $H^{m+1}(T)$.

2. step. With the norm equivalence constant C > 0 of step 1, it holds that

$$||Av||_X = ||A(v - \Pi v)||_X \le ||A|| ||v - \Pi v||_{H^{m+1}(T)} \le C ||A|| |||v - \Pi v||| = C ||A|| ||D^{m+1}v||_{L^2(T)}$$

for all $v \in H^{m+1}(T)$.

Exercise 16. Prove that a function $v \in H^{m+1}(T)$ on a bounded Lipschitz domain $T \subset \mathbb{R}^d$ satisfies $D^{m+1}v = 0$ if and only if $v \in \mathcal{P}^m(T)$. **Hint:** You should use the case m = 0 without a proof, cf. Theorem 2.3.

3.2.4 Scaling Argument and Proof of Approximation Theorem

Lemma 3.8 (Transformation Formula). Let $T, \widehat{T} \subset \mathbb{R}^d$ be Lipschitz domains. Let $\Phi(x) := Bx + y$ with regular matrix $B \in \mathbb{R}^{d \times d}$ and vector $y \in \mathbb{R}^d$ be an affine diffeomorphism with $\Phi(\widehat{T}) = T$. For $u \in H^m(T)$, it holds that $u \circ \Phi \in H^m(\widehat{T})$ with

$$||D^{m}(u \circ \Phi)||_{L^{2}(\widehat{T})} \le |\det B|^{-1/2} ||B||_{F}^{m} ||D^{m}u||_{L^{2}(T)}, \tag{3.29}$$

where $||B||_F$ denotes the Frobenius norm of B. Moreover, for m=0, there even holds equality.

Proof. 1. step. The case m = 0: According to the transformation theorem and $D\Phi(x) = B$, it holds that

$$\|u\|_{L^2(T)}^2 = \int_T u^2 \, dy = \int_{\widehat{T}} (u \circ \Phi)^2 \, |\det D\Phi| \, dx = |\det B| \, \|u \circ \Phi\|_{L^2(\widehat{T})}^2.$$

2. step. To treat the higher-order case for smooth functions $u \in C^{\infty}(\overline{T})$, we first prove by induction on m that for all $j_{\ell} \in \{1, \ldots, d\}$, it holds that

$$\partial_{j_1} \cdots \partial_{j_m} (u \circ \Phi)(x) = \sum_{k_1=1}^d \cdots \sum_{k_m=1}^d \partial_{k_1} \cdots \partial_{k_m} u(\Phi(x)) \prod_{\ell=1}^m B_{k_\ell j_\ell}, \tag{3.30}$$

which is the special case of the Faà di Bruno formula (chain rule for partial derivatives): The case m=1 follows from the chain rule $D(u \circ \Phi)(x) = Du(\Phi(x))D\Phi(x) = Du(\Phi(x))B$, where, e.g., $Du(y) = (\partial_1 u, \ldots, \partial_d u)(y)$. Therefore,

$$\partial_j(u \circ \Phi)(x) = \sum_{k=1}^d \partial_k u(\Phi(x)) B_{kj}.$$

Assuming that (3.30) holds up to $m \in \mathbb{N}$, we now prove the equality for m + 1:

$$\partial_{j_{1}} \cdots \partial_{j_{m+1}} (u \circ \Phi)(x) \stackrel{!}{=} \partial_{j_{1}} \left(\sum_{k_{2}=1}^{d} \cdots \sum_{k_{m+1}=1}^{d} \partial_{k_{2}} \cdots \partial_{k_{m+1}} u(\Phi(x)) \prod_{\ell=2}^{m+1} B_{k_{\ell} j_{\ell}} \right)$$

$$= \sum_{k_{2}=1}^{d} \cdots \sum_{k_{m+1}=1}^{d} \partial_{j_{1}} \left(\partial_{k_{2}} \cdots \partial_{k_{m+1}} u(\Phi(x)) \right) \prod_{\ell=2}^{m+1} B_{k_{\ell} j_{\ell}}$$

$$\stackrel{!}{=} \sum_{k_{2}=1}^{d} \cdots \sum_{k_{m+1}=1}^{d} \sum_{k_{1}=1}^{d} \partial_{k_{1}} \partial_{k_{2}} \cdots \partial_{k_{m+1}} u(\Phi(x)) B_{k_{1} j_{1}} \prod_{\ell=2}^{m+1} B_{k_{\ell} j_{\ell}}$$

$$= \sum_{k_{1}=1}^{d} \cdots \sum_{k_{m+1}=1}^{d} \partial_{k_{1}} \partial_{k_{2}} \cdots \partial_{k_{m+1}} u(\Phi(x)) \prod_{\ell=1}^{m+1} B_{k_{\ell} j_{\ell}},$$

where we have used the induction hypothesis for m and the initial step m = 1. This verifies (3.30).

3. step. We apply the Cauchy inequality to (3.30) to see that

$$\begin{aligned} \left| \partial_{j_{1}} \cdots \partial_{j_{m}} (u \circ \Phi)(x) \right|^{2} &\leq \left(\sum_{k_{1}=1}^{d} \cdots \sum_{k_{m}=1}^{d} \left| \partial_{k_{1}} \cdots \partial_{k_{m}} u(\Phi(x)) \right|^{2} \right) \left(\sum_{k_{1}=1}^{d} \cdots \sum_{k_{m}=1}^{d} \left| \prod_{\ell=1}^{m} B_{k_{\ell} j_{\ell}} \right|^{2} \right) \\ &= \left(\sum_{k_{1}=1}^{d} \cdots \sum_{k_{m}=1}^{d} \left| \partial_{k_{1}} \cdots \partial_{k_{m}} u(\Phi(x)) \right|^{2} \right) \left(\sum_{k_{1}=1}^{d} \cdots \sum_{k_{m}=1}^{d} \prod_{\ell=1}^{m} B_{k_{\ell} j_{\ell}}^{2} \right) \\ &\stackrel{!}{=} \left(\sum_{k_{1}=1}^{d} \cdots \sum_{k_{m}=1}^{d} \left| \partial_{k_{1}} \cdots \partial_{k_{m}} u(\Phi(x)) \right|^{2} \right) \left(\prod_{\ell=1}^{m} \sum_{k_{\ell}=1}^{d} B_{k_{\ell} j_{\ell}}^{2} \right), \end{aligned}$$

where the last equality follows from another simple induction argument.

4. step. We prove the transformation formula (3.29) for $u \in C^{\infty}(\overline{T})$:

$$|\det B| \|D^{m}(u \circ \Phi)\|_{L^{2}(\widehat{T})}^{2} = \int_{\widehat{T}} \sum_{j_{1}=1}^{d} \cdots \sum_{j_{m}=1}^{d} |\partial_{j_{1}} \cdots \partial_{j_{m}}(u \circ \Phi)(x)|^{2} |\det D\Phi(x)| dx$$

$$\leq \underbrace{\left(\sum_{j_{1}=1}^{d} \cdots \sum_{j_{m}=1}^{d} \prod_{\ell=1}^{m} \sum_{k_{\ell}=1}^{d} B_{k_{\ell}j_{\ell}}^{2}\right)}_{=\prod_{\ell=1}^{m} \sum_{j_{\ell}=1}^{d} \sum_{k_{\ell}=1}^{d} B_{k_{\ell}j_{\ell}}^{2}} \underbrace{\left(\int_{\widehat{T}} \sum_{k_{1}=1}^{d} \cdots \sum_{k_{m}=1}^{d} |\partial_{k_{1}} \cdots \partial_{k_{m}} u(\Phi(x))|^{2} |\det D\Phi(x)| dx\right)}_{=\|D^{m}u\|_{L^{2}(T)}^{2}}$$

$$= \|B\|_{F}^{2m} \|D^{m}u\|_{L^{2}(T)}^{2}.$$

5. step. We prove the transformation formula (3.29) for general $u \in H^m(T)$: According to the Meyers-Serrin theorem, $C^{\infty}(\overline{T})$ is a dense subspace of $H^m(T)$. Note that (3.29) implies for $u \in C^{\infty}(\overline{T})$ the estimate $\|u \circ \Phi\|_{H^m(\widehat{T})} \leq C \|u\|_{H^m(T)}$, where C > 0 depends only on m and B. Hence, $\Psi u := u \circ \Phi$ extends uniquely to a linear and continuous mapping $\Psi : H^m(T) \to H^m(\widehat{T})$. For $u \in H^m(T)$, choose $(u_n) \subset C^{\infty}(\overline{T})$ with $u_n \to u \in H^m(T)$. By continuity of Ψ , it holds that $u_n \circ \Phi = \Psi u_n \to \Psi u$ in $H^m(\widehat{T})$. Moreover, according to step 1, it holds that $u_n \circ \Phi \to u \circ \Phi \in L^2(\widehat{T})$. This implies that $u \circ \Phi = \Psi u \in H^m(\widehat{T})$, i.e., the (unique) extension of Ψ from $C^{\infty}(\overline{T})$ to $H^m(T)$ is, in fact, the composition. Moreover, the left-hand side and the right-hand side of (3.29) depend continuously (with respect to $H^m(T)$) on u. This and (3.29) for $u_n \in C^{\infty}(\overline{T})$ prove that

$$||D^{m}(u \circ \Phi)||_{L^{2}(\widehat{T})} = \lim_{n \to \infty} ||D^{m}(u_{n} \circ \Phi)||_{L^{2}(\widehat{T})} \le \lim_{n \to \infty} |\det B|^{-1/2} ||B||_{F}^{m} ||D^{m}u_{n}||_{L^{2}(T)}$$
$$= |\det B|^{-1/2} ||B||_{F}^{m} ||D^{m}u||_{L^{2}(T)}$$

and conclude the proof.

Lemma 3.9. For $\widehat{T} = T_{\text{ref}}$ the reference element and $T = \text{conv}\{z_1, z_2, z_3\} \subset \mathbb{R}^2$ being a non-degenerate triangle, we define

$$\Phi_T : T_{\text{ref}} \to T, \quad \Phi_T(s, t) := z_1 + B \begin{pmatrix} s \\ t \end{pmatrix}, \quad \text{where } B := (z_2 - z_1 \quad z_3 - z_1) \in \mathbb{R}^{2 \times 2}.$$
 (3.31)

Then, it holds that $|\det B| = 2|T|$ and

$$h_T/\sqrt{2} \le ||B||_F \le \sqrt{2} h_T$$
 as well as $\varrho_T^{-1}/\sqrt{2} \le ||B^{-1}||_F \le \sqrt{2} \varrho_T^{-1}$. (3.32)

Proof. It holds that

$$||B||_F^2 = |z_2 - z_1|^2 + |z_3 - z_1|^2 \le 2h_T^2.$$

Moreover,

$$|z_3 - z_2| \le |z_3 - z_1| + |z_2 - z_1| \le \sqrt{2} (|z_3 - z_1|^2 + |z_2 - z_1|^2)^{1/2} \le \sqrt{2} \|B\|_F.$$

In particular, $h_T = \max\{|z_2 - z_1|, |z_3 - z_1|, |z_3 - z_2|\} \le \sqrt{2} \|B\|_F$. The transformation theorem gives

$$\frac{1}{2} |\det B| = |T_{\text{ref}}| |\det B| = \int_{T_{\text{ref}}} |\det D\Phi_T| \, dx = \int_T \, dx = |T| > 0.$$

Hence, $0 < |\det B| = 2|T| = h_T \varrho_T$. In particular, B^{-1} as well as ϱ_T^{-1} are well-defined. It holds that

$$B^{-1} = \frac{1}{\det B} \begin{pmatrix} b_{22} & -b_{12} \\ -b_{21} & b_{11} \end{pmatrix}$$
 for $B = \begin{pmatrix} b_{11} & b_{12} \\ b_{21} & b_{22} \end{pmatrix}$.

In particular, this proves that

$$||B^{-1}||_F = \frac{||B||_F}{|\det B|} = \frac{||B||_F}{h_T \varrho_T},$$

and the second estimate in (3.32) follows from the first.

Proof of Approximation Theorem 3.5. 1. step. Estimate on the reference element T_{ref} : Let $I_h^{\text{ref}}: H^2(T_{\text{ref}}) \to \mathcal{P}^1(T_{\text{ref}})$ denote the nodal interpolation operator on the reference element. We consider the operator

$$A := 1 - I_h^{\text{ref}} : H^2(T_{\text{ref}}) \to H^k(T_{\text{ref}}) \quad \text{for } k = 0, 1$$

and observe that $\mathcal{P}^1(T_{\text{ref}}) \subseteq \ker(A)$. To see that A is continuous, we estimate

$$||Av||_{H^k(T_{\text{ref}})} \le ||v||_{H^2(T_{\text{ref}})} + ||I_h^{\text{ref}}v||_{H^k(T_{\text{ref}})}.$$

Let z_1, z_2, z_3 denote the nodes of the reference element. Since all norms on the finite dimensional space $\mathcal{P}^1(T_{\text{ref}})$ are equivalent, we use the Sobolev inequality to see that

$$\|I_h^{\mathrm{ref}}v\|_{H^k(T_{\mathrm{ref}})} \leq C_{\mathrm{norm}} \max_{j=1,\dots,3} |I_h^{\mathrm{ref}}v(z_j)| \leq C_{\mathrm{norm}}\|v\|_{\infty,T_{\mathrm{ref}}} \leq C_{\mathrm{norm}}C_{\mathrm{sobolev}}\|v\|_{H^2(T_{\mathrm{ref}})}.$$

Altogether, we obtain that $||Av||_{H^k(T_{ref})} \leq (1 + C_{norm}C_{sobolev})||v||_{H^2(T_{ref})}$, whence continuity of the operator A. Consequently, the Bramble-Hilbert lemma provides a constant $C_{ref} > 0$ that depends only on T_{ref} with

$$||v - I_h^{\text{ref}}v||_{H^k(T_{\text{ref}})} \le C_{\text{ref}}||D^2v||_{L^2(T_{\text{ref}})}$$
 for all $v \in H^2(T_{\text{ref}})$ and $k = 0, 1$.

2. step. Scaling arguments provide the estimate on each element T: Let $\Phi = \Phi_T$ denote the affine diffeomorphism from Lemma 3.9. Note that $I_h^{\text{ref}}(u \circ \Phi) = (I_h u) \circ \Phi$. Define $v := u \circ \Phi$ and observe that $(u - I_h u) \circ \Phi = (1 - I_h^{\text{ref}})v$. First, we apply the transformation formula to Φ^{-1} ,

$$\begin{split} \|D^k(u-I_hu)\|_{L^2(T)} &= \|D^k((v-I_h^{\mathrm{ref}}v)\circ\Phi^{-1})\|_{L^2(T)} \\ &\leq |\det B^{-1}|^{-1/2}\|B^{-1}\|_F^k\|D^k(v-I_h^{\mathrm{ref}}v)\|_{L^2(T_{\mathrm{ref}})} \\ &\leq C_{\mathrm{ref}}|\det B|^{1/2}\|B^{-1}\|_F^k\|D^2v\|_{L^2(T_{\mathrm{ref}})}. \end{split}$$

Second, we plug-in $v = u \circ \Phi$ and apply the transformation formula to Φ ,

$$||D^2v||_{L^2(T_{\mathrm{ref}})} = ||D^2(u \circ \Phi)||_{L^2(T_{\mathrm{ref}})} \le |\det B|^{-1/2} ||B||_F^2 ||D^2u||_{L^2(T)}.$$

The combination of the last two estimates proves that

$$||D^k(u - I_h u)||_{L^2(T)} \le C_{\text{ref}} ||B^{-1}||_F^k ||B||_F^2 ||D^2 u||_{L^2(T)} \le C_{\text{ref}} 2^{(k+2)/2} h_T^2 \varrho_T^{-k} ||D^2 u||_{L^2(T)},$$

where we have used the geometric interpretation of $||B||_F$ and $||B^{-1}||_F$. This proves that

$$||u - I_h u||_{L^2(T)} \le 2C_{\text{ref}} ||h^2 D^2 u||_{L^2(T)} \text{ and } ||\nabla (u - I_h u)||_{L^2(T)} \le 2^{3/2} C_{\text{ref}} \sigma(\mathcal{T}) ||h D^2 u||_{L^2(T)}.$$

and thus concludes the proof.

Remark. The proof of Theorem 3.5 shows that it is enough to assume $u \in C(\overline{\Omega}) \cap H^2(\mathcal{T})$, where $H^k(\mathcal{T}) := \{u \in L^2(\Omega) \mid \forall T \in \mathcal{T} \mid u \in H^k(T)\}$ for $k \geq 1$. According to the Sobolev inequality, it holds that $H^2(\Omega) \subseteq C(\overline{\Omega}) \cap H^2(\mathcal{T})$. For the *broken Sobolev spaces* $H^k(\mathcal{T})$, we write $D_h^k v$ for the \mathcal{T} -piecewise k-th derivative of v and, in particular, $\nabla_h v = D_h^1 v$ for the \mathcal{T} -piecewise gradient. \square

Remark. We recall the procedure of a scaling argument for proving an estimate. To that end, let $\Phi_T: T_{\text{ref}} \to T$ be the affine diffeomorphism with linear part B.

• First, transfer the left-hand side from T to T_{ref} :

$$||D^{k}v||_{L^{2}(T)} = ||D^{k}(v \circ \Phi_{T} \circ \Phi_{T}^{-1})||_{L^{2}(T)} \le |\det B^{-1}|^{-1/2} ||B^{-1}||_{F}^{k} ||D^{k}(v \circ \Phi_{T})||_{L^{2}(T_{ref})}$$
$$\simeq |T| \varrho_{T}^{-k} ||D^{k}(v \circ \Phi_{T})||_{L^{2}(T_{ref})},$$

i.e., derivative on the left-hand side give rise to negative powers of ρ_T .

- Second, prove an appropriate estimate on the reference element $T_{\rm ref}$.
- Third, transfer the right-hand side from T_{ref} to T:

$$||D^{\ell}(w \circ \Phi_T)||_{L^2(T_{\text{ref}})} \le |\det B|^{-1/2} ||B||_F^{\ell} ||D^{\ell}w||_{L^2(T)} \simeq |T|^{-1/2} h_T^{\ell} ||D^{\ell}w||_{L^2(T)},$$

i.e., derivatives on the right-hand side give rise to positive powers of h_T .

Plugging everything together, proves the desired estimate.

Note that the heart of the proof of the approximation theorem is the Rellich theorem and thus a compactness argument. The following exercise shows that approximation results are necessarily proved by use of compactness.

Exercise 17. Let X be a Banach space and Y be a normed space with continuous inclusion $Y \subseteq X$. For $h \to 0$, let X_h be finite dimensional subspaces of X and $I_h \in L(Y; X_h)$ be a continuous and linear operator with

$$||u - I_h u||_X \le Ch^\alpha ||u||_Y$$
 for all $u \in Y$,

where the constants $C, \alpha > 0$ are independent of u and h. Then, the continuous inclusion $Y \subseteq X$ is already compact.

Chapter 4

A Posteriori Analysis

4.1 Introduction

We consider the model problem

$$-\Delta u = f \quad \text{in } \Omega,$$

$$u = 0 \quad \text{on } \Gamma_D,$$

$$\partial u/\partial n = \phi \quad \text{on } \Gamma_N.$$
(4.1)

Let $u \in H_D^1(\Omega)$ be the weak solution of (4.1) and $u_h \in \mathcal{S}_D^1(\mathcal{T})$ be the P1-FEM approximation of u. In the previous chapter, we aimed to control the error $||u-u_h||_{H^1(\Omega)}$ by a priori knowledge, e.g., regularity of the given data and the exact solution (but essentially independent of the discrete solution u_h). Since u is unknown, in general, the **a priori analysis** provides a qualitative understanding of the FEM, e.g., convergence with certain rates, but the derived bounds are non-computable in practice. In this chapter, we aim to derive numerically computable bounds $\eta = \eta(u_h, f, \phi, \mathcal{T})$ for the error $||u-u_h||_{H^1(\Omega)}$, which may depend on u_h , the triangulation \mathcal{T} , and the given data f and ϕ (but not on the exact solution u). The quantity η is referred to as (a posteriori) error estimator, and emphasis is laid on the fact that η can be computed algorithmically as soon as the discrete solution $u_h \in \mathcal{S}_D^1(\mathcal{T})$ has been computed. An error estimator η is called reliable provided that

$$||u - u_h||_{H^1(\Omega)} \le C_{\text{rel}} \eta. \tag{4.2}$$

Usually, the information η provides, is used to steer a mesh-refinement that leads to a sequence \mathcal{T}_{ℓ} of regular meshes with nested spaces $\mathcal{S}_D^1(\mathcal{T}_{\ell}) \subseteq \mathcal{S}_D^1(\mathcal{T}_{\ell+1})$, i.e., $\mathcal{T}_{\ell+1}$ is a certain refinement of \mathcal{T}_{ℓ} . If η is reliable the (numerically or algorithmically observed) decrease of η to zero implies the convergence of u_h towards u. However, it might (formally) occur that u_h tends to u, while η does not tend to zero. Therefore, an error estimator η is called **efficient** provided that

$$C_{\text{eff}} \eta \le \|u - u_h\|_{H^1(\Omega)}. \tag{4.3}$$

For an efficient error estimator η , the convergence of u_h to u necessarily implies the convergence of η to zero. Finally, if η is reliable and efficient, we observe for η the same order of convergence as for $||u - u_h||_{H^1(\Omega)}$.

The aim of a posteriori error estimates is twofold:

- We want to control the accuracy $||u u_h||_{H^1(\Omega)}$ of a discrete solution u_h and stop the computation if u_h is sufficiently accurate.
- The mesh-refinement should be steered automatically by the algorithm so that we are led to the highest possible accuracy with the lowest number of degrees of freedom.

Remark. Throughout, we allow the cases $\Gamma_D = \Gamma$ as well as $\Gamma_D = \emptyset$. In the latter case, (4.1) becomes the Neumann problem, for which we have to assume the compatability condition $\int_{\Omega} f \, dx + \int_{\Gamma} \phi \, ds = 0$. Then, $u \in H^1_*(\Omega)$ and, even more important, the test space $H^1_*(\Omega)$ in the weak formulation can equivalently be replaced by the entire space $H^1(\Omega)$. The same holds for the P1-FEM, where $u_h \in \mathcal{S}^1_*(\mathcal{T})$ and where the discrete test is $\mathcal{S}^1_*(\mathcal{T})$ or equivalently $\mathcal{S}^1(\mathcal{T})$.

4.2 Scott-Zhang Projection

Since H^1 -functions are in general not continuous, nodal interpolation requires additional regularity assumptions. In this section, we aim to provide some quasi-interpolation operator which is well-defined for all $u \in H^1(\Omega)$ and also has the projection property. We start with the following elementary lemma

Lemma 4.1. For $z \in \mathcal{K}$, choose an edge $E_z \in \mathcal{E}$ with $z \in E_z$. Then, there is a unique dual function $\psi_z \in \mathcal{P}^1(E_z)$ such that

$$\int_{E_z} \psi_z \zeta_{z'} \, ds = \delta_{zz'} \quad \text{for all } z' \in \mathcal{K}. \tag{4.4}$$

Moreover, it holds that $\|\psi_z\|_{L^{\infty}(E_z)} \leq C |E_z|^{-1}$ for some generic constant C > 0, which is in particular independent of z and \mathcal{T} .

Proof. According to the Riesz theorem, there is a unique function $\widehat{\psi} \in \mathcal{P}^1[0,1]$ such that

$$\int_0^1 \widehat{\psi} \, \widehat{\phi} \, dt = \widehat{\phi}(0) \quad \text{for all } \widehat{\phi} \in \mathcal{P}^1[0,1].$$

Let $\Phi_z:[0,1]\to E_z$ be an affine parametrization of the edge E_z with $\Phi_z(0)=z$. We define

$$\psi_z := \frac{1}{|E_z|} \, \widehat{\psi} \circ \Phi_z^{-1} \in \mathcal{P}^1(E_z).$$

Clearly, $\|\psi_z\|_{L^{\infty}(E_z)} \leq \|\widehat{\psi}\|_{L^{\infty}(0,1)} |E_z|^{-1}$. Note that $|\Phi'_z| = |E_z|$ and hence

$$\int_{E_z} \psi_z \zeta_{z'} \, ds = \int_0^1 (\psi_z \circ \Phi_z) \, (\zeta_{z'} \circ \Phi_z) \, |\Phi_z'| \, dt = \int_0^1 \widehat{\psi}(t) \, (\zeta_{z'} \circ \Phi_z)(t) \, dt = \zeta_{z'} \big(\Phi_z(0)\big) = \zeta_{z'}(z).$$

This concludes the proof.

Definition. For each node $z \in \mathcal{K}$ of \mathcal{T} , we choose an edge $E_z \in \mathcal{E}$ such that

•
$$E_z \subseteq \overline{\Gamma}_D$$
 for $z \in \overline{\Gamma}_D$,

- $E_z \subseteq \Gamma$ for $z \in \Gamma$,
- E_z arbitrary for $z \in \Omega$.

Note that the precise choice is immaterial for the following analysis. For $z \in \mathcal{K}$, let $\psi_z \in \mathcal{P}^1(E_z)$ be the corresponding dual function. Then, the **Scott-Zhang projection** is defined by

$$J_h v := \sum_{z \in \mathcal{K}} \left(\int_{E_z} \psi_z v \, ds \right) \zeta_z. \tag{4.5}$$

Clearly, $J_h: H^1(\Omega) \to \mathcal{S}^1(\mathcal{T})$ is well-defined and linear. Our first proposition states that J_h is in fact a projection which preserves discrete boundary data.

Proposition 4.2. For $v \in H^1(\Omega)$ and $v_h \in S^1(\mathcal{T})$, the following properties (i)–(iii) are true:

- (i) $J_h v_h = v_h$.
- (ii) $(J_h v)|_{\omega}$ depends only on the trace $v|_{\omega}$ for $\omega \in \{\Gamma, \Gamma_D\}$.
- (iii) $v|_{\omega} = v_h|_{\omega}$ implies that $(J_h v)|_{\omega} = v|_{\omega}$ for $\omega \in \{\Gamma, \Gamma_D\}$.

Proof. (i) Note that $v_h = \sum_{z' \in \mathcal{K}} v_h(z') \zeta_{z'}$. By choice of ψ_z , this shows

$$\int_{E_z} \psi_z v_h \, ds = \sum_{z' \in \mathcal{K}} v_h(z') \int_{E_z} \psi_z \zeta_{z'} \, ds = v_h(z).$$

With this, we deduce

$$J_h v_h = \sum_{z \in \mathcal{K}} \Big(\int_{E_z} \psi_z v_h \, ds \Big) \zeta_z = \sum_{z \in \mathcal{K}} v_h(z) \, \zeta_z = v_h.$$

(ii) follows from the choice of the edges E_z . (iii) We consider only $\omega = \Gamma_D$. We first note that

$$(J_h u)|_{\Gamma_D} = \sum_{z \in \mathcal{K}} \Big(\int_{E_z} \psi_z u \, ds \Big) \zeta_z|_{\Gamma_D} = \sum_{z \in \mathcal{K} \cap \overline{\Gamma}_D} \Big(\int_{E_z} \psi_z u \, ds \Big) \zeta_z|_{\Gamma_D} \quad \text{for all } u \in H^1(\Omega).$$

For $z \in \mathcal{K} \cap \overline{\Gamma}_D$, it holds that $E_z \subseteq \overline{\Gamma}_D$ and hence $\int_{E_z} \psi_z v_h \, ds = \int_{E_z} \psi_z v \, ds$. Together with the last equation and the projection property, we obtain that

$$(J_h v)|_{\Gamma_D} = (J_h v_h)|_{\Gamma_D} = v_h|_{\Gamma_D}.$$

This concludes the proof.

Exercise 18. Show that Lemma 4.1 holds for any dimension
$$d \geq 2$$
.

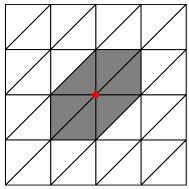
Note that the Scott-Zhang projection $J_h v$ is not defined for general L^2 -functions, since $L^2(T)$ does not provide traces. However, one can define an appropriate variant as follows:

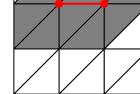
Exercise 19. Construct a linear projection $P_h: L^2(\Omega) \to \mathcal{S}^1(\mathcal{T})$ which satisfies

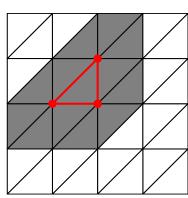
- $||P_h v||_{L^2(\Omega)} \le C ||v||_{L^2(\Omega)}$ for all $v \in L^2(\Omega)$.
- $P_h v_h = v_h$ for all $v_h \in \mathcal{S}^1(\mathcal{T})$.

The constant C > 0 may only depend on $\sigma(\mathcal{T})$. **Hint.** Proceed as for the standard Scott-Zhang projection. Instead of an edge E_z , associate with each node $z \in \mathcal{K}$ an arbitrary element $T_z \in \mathcal{T}$ with $z \in T_z$.

Next, we aim to show that the Scott-Zhang projection has local stability and approximation properties. Unlike nodal interpolation, this will require appropriate patches.







Patch Ω_z of a node $z \in \mathcal{K}$

Patch Ω_E of an edge $E \in \mathcal{E}$

Patch Ω_T of an element $T \in \mathcal{T}$

FIGURE 4.1. For the a posteriori analysis, we need three types of patches $\omega \subseteq \Omega$, namely patches of nodes, edges, and elements, respectively. Note that the patch of an edge (or of an element) just is the union of the patches of its nodes.

Definition. For the a posteriori analysis, we need certain unions of elements, called **patches**, cf. Figure 4.1: For a node $z \in \mathcal{K}$, we define

$$\widetilde{\Omega}_z := \left\{ T \in \mathcal{T} \,\middle|\, z \in \mathcal{K}_T \right\} \quad \text{as well as} \quad \Omega_z := \bigcup \widetilde{\Omega}_z := \left\{ x \in \mathbb{R}^2 \,\middle|\, \exists T \in \widetilde{\Omega}_z \quad x \in T \right\}. \tag{4.6}$$

For an edge $E \in \mathcal{E}$, we define

$$\widetilde{\Omega}_E := \left\{ T \in \mathcal{T} \,\middle|\, \mathcal{K}_E \cap T \neq \emptyset \right\} = \left\{ T \in \widetilde{\Omega}_z \,\middle|\, z \in \mathcal{K}_E \right\} \quad \text{as well as} \quad \Omega_E := \bigcup \widetilde{\Omega}_E. \tag{4.7}$$

Finally, for an element $T \in \mathcal{T}$, we define

$$\widetilde{\Omega}_T := \left\{ T' \in \mathcal{T} \,\middle|\, \mathcal{K}_T \cap T' \neq \emptyset \right\} = \left\{ T' \in \widetilde{\Omega}_z \,\middle|\, z \in \mathcal{K}_T \right\} \quad \text{as well as} \quad \Omega_T := \bigcup \widetilde{\Omega}_T. \tag{4.8}$$

The patches Ω_z , Ω_E , and Ω_T are visualized in Figure 4.1.

Lemma 4.3. There is a constant C > 0 which depends only on $\sigma(\mathcal{T})$, such that

- $\#\widetilde{\Omega}_z \leq C \text{ for all } z \in \mathcal{K},$
- $\#\widetilde{\Omega}_E \leq C \text{ for all } E \in \mathcal{E},$

• $\#\widetilde{\Omega}_T \leq C \text{ for all } T \in \mathcal{T},$

i.e., the number of elements per patch is uniformly bounded. Moreover,

• $\#\{T' \in \mathcal{T} \mid T \in \widetilde{\Omega}_{T'}\} \leq C \text{ for all } T \in \mathcal{T},$

i.e., an element $T \in \mathcal{T}$ belongs only to finitely many patches.

Proof. Note that $\sigma(\mathcal{T})$ provides a bound for the minimal interior angle of all elements $T \in \mathcal{T}$; see Exercise 13. Consequently, there is a maximal number C > 0 of elements in $\widetilde{\Omega}_z$, for all nodes $u \in \mathcal{K}$. By definition, there follows $\#\widetilde{\Omega}_E \leq 2C$ as well as $\#\widetilde{\Omega}_T \leq 3C$.

An essential consequence of Lemma 4.3 is that

$$||v||_{L^2(\Omega)} \le \left(\sum_{T \in \mathcal{T}} ||v||_{L^2(\Omega_T)}^2\right)^{1/2} \le C_{\text{patch}} ||v||_{L^2(\Omega)} \quad \text{for all } v \in L^2(\Omega),$$

where $C_{\text{patch}} > 0$ depends only on $\sigma(\mathcal{T})$. Another consequence of Lemma 4.3 is that the diameter $\operatorname{diam}(\Omega_T)$ of a patch is proportional to $h_T = \operatorname{diam}(T)$. This is stated in the following lemma.

Lemma 4.4. For a regular triangulation, it holds that

- diam $(\Omega_z) \leq C h_T$ for all $z \in \mathcal{K}$ and $T \in \widetilde{\Omega}_z$,
- diam $(\Omega_E) \le C h_E \le C h_T$ for all $E \in \mathcal{E}$ and $T \in \widetilde{\Omega}_E$,
- diam $(\Omega_{T'}) \leq C h_T$ for all $T' \in \mathcal{T}$ and $T \in \widetilde{\Omega}_{T'}$.

The constant C > 0 depends only on $\sigma(\mathcal{T})$.

Proof. 1. step. Note that $h_T \leq \sigma(\mathcal{T})\varrho_T \leq \sigma(\mathcal{T})h_E$ for all $T \in \mathcal{T}$ and all edges $E \in \mathcal{E}_T$.

2. step. Patch of a node $z \in \mathcal{K}$: For $\widetilde{\Omega}_z = \{T_1, \dots, T_n\}$, we may choose a numbering such that T_{j-1}, T_j are neighbours, i.e., $T_{j-1} \cap T_j \in \mathcal{E}$. From step 1, we derive $h_{T_{j-1}} \leq \sigma(\mathcal{T})h_{T_j}$, whence $h_{T'} \leq \sigma(\mathcal{T})^{n-1}h_T$ for all $T, T' \in \widetilde{\Omega}_z$. This yields that

$$\operatorname{diam}(\Omega_z) \leq 2 \max_{T' \in \widetilde{\Omega}_z} h_{T'} \leq 2\sigma(\mathcal{T})^{n-1} h_T \quad \text{for all } T \in \widetilde{\Omega}_z.$$

3. step. Patch of an edge $E \in \mathcal{E}$: With $E = \text{conv}\{z_1, z_2\}$ for some $z_1, z_2 \in \mathcal{K}$, it holds that $\widetilde{\Omega}_E = \widetilde{\Omega}_{z_1} \cup \widetilde{\Omega}_{z_2}$ as well as $\widetilde{\Omega}_{z_1} \cap \widetilde{\Omega}_{z_2} \neq \emptyset$. Let $T \in \widetilde{\Omega}_E$ and $n := \max\{\#\widetilde{\Omega}_{z_1}, \#\widetilde{\Omega}_{z_2}\}$. Without loss of generality, we may assume $T \in \widetilde{\Omega}_{z_1}$. Choose $T' \in \widetilde{\Omega}_{z_1} \cap \widetilde{\Omega}_{z_2}$. Then,

$$\operatorname{diam}(\Omega_E) \leq \operatorname{diam}(\Omega_{z_1}) + \operatorname{diam}(\Omega_{z_2}) \leq 2\sigma(\mathcal{T})^{n-1}(h_T + h_{T'}) \leq 2\sigma(\mathcal{T})^{n-1}(1 + \sigma(\mathcal{T})^{n-1})h_{T'}$$

4. step. Patch of an element $T \in \mathcal{T}$: Simply use the same arguments as in step 3.

The Scott-Zhang projection is locally H^1 -stable and has a local first-order approximation property.

Proposition 4.5. For all $T \in \mathcal{T}$, it holds that

$$||v - J_h v||_{L^2(T)} + h_T ||\nabla J_h v||_{L^2(T)} \le C h_T ||\nabla v||_{L^2(\Omega_T)} \quad \text{for all } v \in H^1(\Omega).$$
(4.9)

The constant C > 0 depends only on γ -shape regularity of \mathcal{T} .

The proof requires the following technical lemmata which are also valid in any dimension $d \ge 2$ as the proofs reveal.

Theorem 4.6 (Trace Inequality). Let $T = \text{conv}\{z_0, \ldots, z_d\} \subset \mathbb{R}^d$ be a simplex in \mathbb{R}^d with |T| > 0 and diameter $h_T := \text{diam}(T)$. Let $E = \text{conv}\{z_1, \ldots, z_d\}$ denote one particular side of the simplex. Then, for $v \in H^1(T)$, it holds

$$||v||_{L^{2}(E)}^{2} \leq \frac{|E|}{|T|} \left(||v||_{L^{2}(T)}^{2} + \frac{2}{d} h_{T} ||v\nabla v||_{L^{1}(T)} \right) \leq \frac{|E|}{|T|} \left(1 + 2 h_{T}/d \right) ||v||_{H^{1}(T)}^{2}. \tag{4.10}$$

With the integral means $v_T := |T|^{-1} \int_T v \, dx$ and $v_E := h_E^{-1} \int_E v \, ds$, it holds that

$$||v - v_E||_{L^2(E)}^2 \le ||v - v_T||_{L^2(E)}^2 \le C \frac{|E| h_T^2}{|T|} ||\nabla v||_{L^2(T)}^2, \tag{4.11}$$

where C>0 depends only on the reference element $T_{\rm ref}$ and the dimension d.

The proof of the trace inequalities (4.10)–(4.11) is done with the help of the following lemma. In particular, we shall see that both estimates are sharp. Note that, for d=2, it holds that $|E|/|T| \le 2\varrho_T^{-1}$ and $|E| h_T^2/|T| \le 2\sigma(T)h_T$.

Lemma 4.7 (Trace Identity). Let $T = \text{conv}\{z_0, \dots, z_d\} \subset \mathbb{R}^d$ be a simplex in \mathbb{R}^d with |T| > 0. Let $E = \text{conv}\{z_1, \dots, z_d\}$ denote one particular side of the simplex. Then,

$$\frac{1}{|E|} \int_{E} w \, ds = \frac{1}{|T|} \int_{T} w \, dx + \frac{1}{d|T|} \int_{T} (x - z_{0}) \cdot \nabla w \, dx \tag{4.12}$$

for all $w \in C^1(\overline{T})$.

Proof. We apply the Gauss Divergence Theorem to the function $f(x) := w(x)(x - z_0)$. With $\operatorname{div} f(x) = \nabla w(x) \cdot (x - z_0) + dw(x)$, we obtain that

$$d\int_T w \, dx + \int_T (x - z_0) \cdot \nabla w(x) \, dx = \int_T \operatorname{div} f \, dx = \int_{\partial T} f \cdot n \, ds.$$

Note that $(x - z_0) \cdot n(x) = 0$ for $x \in \partial T \setminus E$, whereas $(x - z_0) \cdot n(x) = \operatorname{dist}(z_0, H)$, where $H \subset \mathbb{R}^d$ denotes the hyperplane with $E \subseteq H$. Therefore, the boundary integral simplifies to $\int_{\partial T} f \cdot n \, ds = \operatorname{dist}(z_0, H) \int_E w \, ds$ and the latter equality reads

$$\frac{1}{|T|} \int_T w \, dx + \frac{1}{d|T|} \int_T (x - z_0) \cdot \nabla w \, dx = \frac{\operatorname{dist}(z_0, H)|E|}{d|T|} \frac{1}{|E|} \int_E w \, ds,$$

which holds for any $w \in C^1(\overline{T})$. The special choice w=1 can be used to determine the w-independent constant $\frac{\operatorname{dist}(z_0,H)|E|}{d|T|}=1$. This concludes the proof.

Remark. Note that Lemma 4.7 holds for any $w \in W^{1,1}(T) := \{w \in L^1(T) \text{ weakly differentiable } | \nabla w \in L^1(T)^d \}$, even with the same proof.

Proof of Theorem 4.6. According to standard density arguments, it suffices to consider $v \in C^1(\overline{T})$. Plugging $w := v^2 \in C^1(\overline{T})$ into the trace identity (4.12), we see that

$$\frac{1}{|E|} \int_E v^2 \, ds = \frac{1}{|T|} \int_T v^2 \, dx + \frac{1}{d|T|} \int_T (x - z_0) \cdot (2v \nabla v) \, dx.$$

This is rewritten in the form

$$\frac{|T|}{|E|} \|v\|_{L^{2}(E)}^{2} = \|v\|_{L^{2}(T)}^{2} + \frac{2}{d} \int_{T} (x - z_{0}) \cdot (v\nabla v) \, dx \le \|v\|_{L^{2}(T)}^{2} + \frac{2}{d} h_{T} \|v\nabla v\|_{L^{1}(T)} \\
\le (1 + 2h_{T}/d) \|v\|_{H^{1}(T)}^{2}$$

which proves (4.10). For the proof of (4.11), we simply replace v by $v - v_T$ and apply the Poincaré inequality. This leads to

$$||v - v_T||_{L^2(E)}^2 \le \frac{|E|}{|T|} (||v - v_T||_{L^2(T)}^2 + \frac{2}{d} h_T ||v - v_T||_{L^2(T)} ||\nabla v||_{L^2(T)})$$

$$\le \frac{|E|}{|T|} (C_P^2 h_T^2 ||\nabla v||_{L^2(T)}^2 + \frac{2}{d} C_P h_T^2 ||\nabla v||_{L^2(T)}^2)$$

$$= (C_P^2 + 2C_P/d) \frac{|E| h_T^2}{|T|} ||\nabla v||_{L^2(T)}^2.$$

The remaining estimate $||v - v_E||_{L^2(E)} \le ||v - v_T||_{L^2(E)}$ follows from the L^2 -best approximation property of the integral mean.

Lemma 4.8 (Generalized Poincaré-Friedrichs inequality). Let $v \in H^1(\Omega)$, $T \in \mathcal{T}$, $T' \in \widetilde{\Omega}_T$, and $E' \in \mathcal{E}_{T'}$. Define the integral means $v_T := (1/|T|) \int_T v \, dx$, $v_{T'} := (1/|T'|) \int_{T'} v \, dx$, and $v_{E'} := (1/|E'|) \int_{E'} v \, ds$. Then,

$$||v_T - v_{T'}||_{L^2(T)} + ||v_T - v_{E'}||_{L^2(T)} \le C h_T ||\nabla v||_{L^2(\Omega_T)}.$$
(4.13)

In particular, this implies that

$$||v - v_{T'}||_{L^2(\Omega_T)} + ||v - v_{E'}||_{L^2(\Omega_T)} \le C h_T ||\nabla v||_{L^2(\Omega_T)}.$$
(4.14)

In either estimate, the constant C > 0 depends only on γ -shape regularity of \mathcal{T} , but is independent of Ω and the shape of Ω_T .

Proof. To ease the notation, let $v_E := (1/|E|) \int_E v \, ds$ also denote the integral mean over edges.

1. step. For any $w \in H^1(T)$, it holds that $|w_T - w_E| \le C \|\nabla w\|_{L^2(T)}$, where C > 0 depends only on shape regularity: With the trace inequality (4.11) on T, we see that

$$|w_T - w_E| \le |E|^{-1} \|w - w_T\|_{L^1(E)} \le |E|^{-1/2} \|w - w_T\|_{L^2(E)}$$

$$\le C \frac{h_T}{|T|^{1/2}} \|\nabla w\|_{L^2(T)} =: C \|\nabla w\|_{L^2(T)}$$

2. step. For all $T, T' \in \mathcal{T}$ with $E := T \cap T' \in \mathcal{E}$, shape regularity and the triangle inequality yield that

$$||v_T - v_{T'}||_{L^2(T)} = |T|^{1/2} |v_T - v_{T'}| \lesssim |T|^{1/2} |v_T - v_E| + |T'|^{1/2} |v_E - v_{T'}|$$

$$= ||v_T - v_E||_{L^2(T)} + ||v_{T'} - v_E||_{L^2(T')}.$$

With Step 1, we thus see that

$$||v_T - v_{T'}||_{L^2(T)} \lesssim h_T ||\nabla v||_{L^2(T)} + h_{T'} ||\nabla v||_{L^2(T')} \lesssim h_T ||\nabla v||_{L^2(T \cup T')},$$

where the hidden constant now depends on C > 0 from step 1 and from shape regularity of \mathcal{T} .

- **3.** step. If $T \cap T' \neq \emptyset$, there is a minimal $n \in \mathbb{N}$ and elements $T_0, \ldots, T_n \in \mathcal{T}$ with $T_0 = T$, $T_j \cap T_{j-1} \in \mathcal{E}$ and $T_j \subseteq \Omega_T$ for all $j = 1, \ldots, n$, and $T_n = T'$. Note that n is uniformly bounded in terms of the γ -shape regularity of \mathcal{T} . Iterating the argument from Step 2, we conclude (4.13) with $\bigcup_{j=0}^n T_j \subseteq \Omega_T$. The overall constant then depends on C > 0 and γ .
 - **4. step.** For each element $T'' \in \widetilde{\Omega}_T$, the Poincaré inequality and (4.13) show

$$\begin{split} &\|v-v_{T'}\|_{L^{2}(T'')} + \|v-v_{E'}\|_{L^{2}(T'')} \\ &\lesssim \|v-v_{T''}\|_{L^{2}(T'')} + \|v_{T}-v_{T'}\|_{L^{2}(T'')} + \|v_{T}-v_{T''}\|_{L^{2}(T'')} + \|v_{T}-v_{E'}\|_{L^{2}(T'')} \\ &\simeq \|v-v_{T''}\|_{L^{2}(T'')} + \|v_{T}-v_{T'}\|_{L^{2}(T)} + \|v_{T}-v_{T''}\|_{L^{2}(T)} + \|v_{T}-v_{E'}\|_{L^{2}(T)} \\ &\lesssim h_{T''} \|\nabla v\|_{L^{2}(T'')} + h_{T} \|\nabla v\|_{L^{2}(\Omega_{T})} \\ &\lesssim h_{T} \|\nabla v\|_{L^{2}(\Omega_{T})}. \end{split}$$

Summing this estimate over all $T'' \in \widetilde{\Omega}_T$, we obtain that

$$||v - v_{T'}||_{L^2(\Omega_T)} + ||v - v_{E'}||_{L^2(\Omega_T)} \lesssim h_T ||\nabla v||_{L^2(\Omega_T)},$$

where the hidden constants depends only on γ -shape regularity of \mathcal{T} .

Proof of Proposition 4.5 (H¹-stability). For $z \in \mathcal{K}$, let $E_z \subset T_z \in \mathcal{T}$ and $h_z := \operatorname{diam}(T_z)$. Note that $T_z \subseteq \Omega_T$ for $z \in T$. The trace inequality (4.10) yields that

$$\|v\|_{L^2(E_z)}^2 \lesssim h_z^{-1} \left(\|v\|_{L^2(T_z)}^2 + h_z \, \|v\|_{L^2(T_z)} \|\nabla v\|_{L^2(T_z)} \right) \lesssim h_z^{-1} \left(\|v\|_{L^2(T_z)}^2 + h_z^2 \, \|\nabla v\|_{L^2(T_z)}^2 \right)$$

With this and Lemma 4.1, we see that

$$\left| \int_{E_z} \psi_z v \, ds \right| \le \|\psi_z\|_{L^{\infty}(E_z)} \|v\|_{L^1(E_z)} \lesssim |E_z|^{-1/2} \|v\|_{L^2(E_z)}$$

$$\lesssim |E_z|^{-1/2} h_z^{-1/2} \left(\|v\|_{L^2(T_z)} + h_z \|\nabla v\|_{L^2(T_z)} \right).$$

For any hat function, an inverse estimate shows

$$\|\nabla \zeta_z\|_{L^2(T)} \lesssim h_T^{-1} \|\zeta_z\|_{L^2(T)} \leq |T|^{1/2} h_T^{-1}.$$

Together with $|E_z| h_z \simeq |T_z| \simeq |T|$ and $h_z \simeq h_T$, we therefore obtain that, for all $v \in H^1(\Omega)$,

$$\|\nabla J_h v\|_{L^2(T)} \le \sum_{z \in \mathcal{K} \cap T} \left| \int_{E_z} \psi_z v \, ds \right| \|\nabla \zeta_z\|_{L^2(T)} \lesssim \sum_{z \in \mathcal{K} \cap T} \left(h_z^{-1} \, \|v\|_{L^2(T_z)} + \|\nabla v\|_{L^2(T_z)} \right). \tag{4.15}$$

With the integral mean $v_T := (1/|T|) \int_T v \, dx$ and the projection property $J_h v_T = v_T$, we apply the last estimate for $w := v - v_T$ and see that

$$\|\nabla J_h v\|_{L^2(T)} = \|\nabla J_h(v - v_T)\|_{L^2(T)} \lesssim \sum_{z \in \mathcal{K} \cap T} \left(h_z^{-1} \|v - v_T\|_{L^2(T_z)} + \|\nabla v\|_{L^2(T_z)}\right).$$

According to the Poincaré inequality and Lemma 4.8, it holds that for all $z \in \mathcal{K} \cap T$,

$$||v - v_T||_{L^2(T_z)} \le ||v - v_{T_z}||_{L^2(T_z)} + ||v_{T_z} - v_T||_{L^2(T_z)} \lesssim h_z ||\nabla v||_{L^2(\Omega_T)}.$$

$$(4.16)$$

Combining the last two estimates, we thus conclude $\|\nabla J_h v\|_{L^2(T)} \lesssim \|\nabla v\|_{L^2(\Omega_T)}$.

Proof of Proposition 4.5 (approximation property). We adopt the notation from the proof of local H^1 -stability. Arguing as for (4.15), we see that

$$||J_h v||_{L^2(T)} \le \sum_{z \in \mathcal{K} \cap T} \left| \int_{E_z} \psi_z v \, ds \right| ||\zeta_z||_{L^2(T)} \lesssim \sum_{z \in \mathcal{K} \cap T} \left(||v||_{L^2(T_z)} + h_z \, ||\nabla v||_{L^2(T_z)} \right). \tag{4.17}$$

With the integral mean $v_T := (1/|T|) \int_T v \, dx$ and the projection property $J_h v_T = v_T$, we apply the last estimate for $w := v - v_T$ and see that

$$||v - J_h v||_{L^2(T)} = ||(v - v_T) - J_h(v - v_T)||_{L^2(T)}$$

$$\leq ||v - v_T||_{L^2(T)} + ||J_h(v - v_T)||_{L^2(T)}$$

$$\lesssim h_T ||\nabla v||_{L^2(T)} + \sum_{z \in \mathcal{K} \cap T} (||v - v_T||_{L^2(T_z)} + h_z ||\nabla v||_{L^2(T_z)})$$

Finally, we employ (4.16) and $h_z \simeq h_T$ to conclude $||v - J_h v||_{L^2(T)} \lesssim h_T ||\nabla v||_{L^2(\Omega_T)}$.

The following theorem concludes the main properties of the Scott-Zhang projection:

Theorem 4.9. The Scott-Zhang projection $J_h: H^1(\Omega) \to \mathcal{S}^1(\mathcal{T})$ has the following properties (i)–(vii):

(i) J_h is linear and continuous with respect to the H^1 -norm, i.e.,

$$||J_h v||_{H^1(\Omega)} \le C (1 + \operatorname{diam}(\Omega)) ||v||_{H^1(\Omega)} \quad \text{for all } v \in H^1(\Omega).$$
 (4.18)

(ii) J_h is a projection onto $S^1(\mathcal{T})$, i.e.,

$$J_h v_h = v_h \quad \text{for all } v_h \in \mathcal{S}^1(\mathcal{T}).$$
 (4.19)

(iii) J_h preserves discrete boundary data, i.e., for $\omega \in \{\Gamma_D, \Gamma\}$ it holds that

$$(J_h v)|_{\omega} = v|_{\omega} \quad \text{for all } v \in H^1(\Omega) \text{ with } v|_{\omega} \in \mathcal{S}^1(\mathcal{T}|_{\omega}).$$
 (4.20)

(iv) J_h is locally H^1 -stable, i.e.,

$$\|\nabla J_h v\|_{L^2(T)} \le C \|\nabla v\|_{L^2(\Omega_T)} \quad \text{for all } v \in H^1(\Omega) \text{ and } T \in \mathcal{T}.$$

$$\tag{4.21}$$

(v) J_h has a local first-order approximation property, i.e.,

$$\|(1 - J_h)v\|_{L^2(T)} \le Ch_T \|\nabla v\|_{L^2(\Omega_T)} \quad \text{for all } v \in H^1(\Omega) \text{ and } T \in \mathcal{T}.$$
 (4.22)

(vi) J_h is quasi-optimal in the sense of the Céa lemma, i.e.,

$$\|(1 - J_h)v\|_{H^1(\Omega)} \le C(1 + \operatorname{diam}(\Omega)) \min_{v_h \in \mathcal{S}^1(\mathcal{T})} \|v - v_h\|_{H^1(\Omega)} \quad \text{for all } v \in H^1(\Omega).$$
 (4.23)

(vii) For all $\alpha \in \mathbb{R}$, J_h is quasi-optimal in the sense of

$$||h^{\alpha}\nabla(1-J_h)v||_{L^2(\Omega)} \le C \min_{v_h \in S^1(\mathcal{T})} ||h^{\alpha}\nabla(v-v_h)||_{L^2(\Omega)}.$$
 (4.24)

The constant C > 0 in (i)-(vii) depends only on γ -shape regularity of \mathcal{T} .

Proof. (ii)–(v) have already been shown, and (i) is a direct consequence of (vi) and the triangle inequality. (vii) Let $v_h \in \mathcal{S}^1(\mathcal{T})$. With the projection property of J_h and (iv), we see that, for all $T \in \mathcal{T}$,

$$\|\nabla(1-J_h)v\|_{L^2(T)} = \|\nabla(1-J_h)(v-v_h)\|_{L^2(T)} \lesssim \|\nabla(v-v_h)\|_{L^2(\Omega_T)}.$$

With γ -shape regularity and hence $h_T \simeq h_{T'}$ for all $T' \subseteq \Omega_T$, we infer

$$||h^{\alpha}\nabla(1-J_h)v||_{L^2(T)} \lesssim ||h^{\alpha}\nabla(v-v_h)||_{L^2(\Omega_T)}.$$

Using the γ -shape regularity again, this results in

$$||h^{\alpha}\nabla(1-J_h)v||_{L^{2}(\Omega)}^{2} = \sum_{T\in\mathcal{T}} ||h^{\alpha}\nabla(1-J_h)v||_{L^{2}(T)}^{2} \lesssim \sum_{T\in\mathcal{T}} ||h^{\alpha}\nabla(v-v_h)||_{L^{2}(\Omega_T)}^{2}$$
$$\lesssim ||h^{\alpha}\nabla(v-v_h)||_{L^{2}(\Omega)}^{2}.$$

This proves (vii) with an infimum on the right-hand side. Due to finite dimension, this infimum is, in fact, attained. To prove (vi), it remains to estimate the L^2 -part and use $\alpha = 0$ in (vii). With the projection property of J_h and (v), shape regularity yields that

$$\begin{aligned} \|(1 - J_h)v\|_{L^2(\Omega)}^2 &= \sum_{T \in \mathcal{T}} \|(1 - J_h)(v - v_h)\|_{L^2(T)}^2 \lesssim \sum_{T \in \mathcal{T}} h_T^2 \|\nabla(v - v_h)\|_{L^2(\Omega_T)}^2 \\ &\lesssim \operatorname{diam}(\Omega)^2 \|\nabla(v - v_h)\|_{L^2(\Omega)}^2. \end{aligned}$$

Altogether, we thus see that

$$\|(1-J_h)v\|_{H^1(\Omega)}^2 \lesssim (1+\operatorname{diam}(\Omega)^2) \|\nabla(v-v_h)\|_{L^2(\Omega)}^2 \lesssim (1+\operatorname{diam}(\Omega))^2 \|v-v_h\|_{H^1(\Omega)}^2$$

This concludes the proof of (vi).

Remark. Theorem 4.9 holds for any dimension $d \ge 2$ and for any fixed polynomial degree $p \ge 1$.

One drawback of the Scott-Zhang projection is that it is not positivity conserving, i.e., $v \ge 0$ does not necessarily imply that $J_h v \ge 0$.

Exercise 20. Suppose that \mathcal{T} is a regular triangulation of $\Omega := [0,1]^2$ into 2 triangles. Find an example of a function $v \in H^1(\Omega)$ with $v \geq 0$ such that there exists some $x \in \Omega$ with $J_h v < 0$. **Hint.** Compute the function $\widehat{\psi} \in \mathcal{P}^1(0,1)$ from Lemma 4.1 explicitly.

Exercise 21. Extend the approach of Exercise 19 and construct an operator $P_h: L^2(\Omega) \to \mathcal{S}_D^1(\mathcal{T})$ with the following properties:

(i) $P_h: L^2(\Omega) \to \mathcal{S}_D^1(\mathcal{T})$ is a well-defined linear projection,

$$P_h v_h = v_h$$
 for all $v_h \in \mathcal{S}_D^1(\mathcal{T})$.

(ii) P_h is locally L^2 -stable, i.e., for all $T \in \mathcal{T}$, it holds that

$$\|(1-P_h)v\|_{L^2(T)} \le C \|v\|_{L^2(\Omega_T)}$$
 for all $v \in L^2(\Omega)$.

(iii) P_h is locally H^1_D -stable, i.e., for all $T \in \mathcal{T}$, it holds that

$$\|\nabla(1-P_h)v\|_{L^2(T)} \le C \|\nabla v\|_{L^2(\Omega_T)}$$
 for all $v \in H_D^1(\Omega)$.

(iv) P_h has a local first-order approximation property

$$\|(1-P_h)v\|_{L^2(T)} \le Ch_T \|\nabla v\|_{L^2(\Omega_T)}$$
 for all $v \in H_D^1(\Omega)$.

- (v) $P_h: L^2(\Omega) \to L^2(\Omega)$ as well as $P_h: H^1_D(\Omega) \to H^1_D(\Omega)$ are bounded linear operators.
- (vi) P_h is quasi-optimal in the sense of the Céa lemma, i.e.,

$$\|(1-P_h)v\|_{H^1(\Omega)} \le C \min_{v_h \in \mathcal{S}_D^1(\mathcal{T})} \|v-v_h\|_{H^1(\Omega)} \text{ for all } v \in H_D^1(\Omega).$$

(vii) For all $\alpha \in \mathbb{R}, P_h$ is quasi-optimal in the sense of

$$||h^{\alpha}(1-P_h)v||_{L^2(\Omega)} \le C \min_{v_h \in \mathcal{S}_D^1(\mathcal{T})} ||h^{\alpha}(v-v_h)||_{L^2(\Omega)} \text{ for all } v \in L^2(\Omega).$$

(viii) For all $\alpha \in \mathbb{R}$, P_h is quasi-optimal in the sense of

$$\|h^{\alpha}\nabla(1-P_h)v\|_{L^2(\Omega)} \leq C \min_{v_h \in \mathcal{S}_D^1(\mathcal{T})} \|h^{\alpha}\nabla(v-v_h)\|_{L^2(\Omega)} \quad \text{for all } v \in H_D^1(\Omega).$$

The constant C > 0 in (i)–(viii) depends only on γ -shape regularity of \mathcal{T} . **Hint.** Let $\mathcal{K}_F := \mathcal{K} \setminus \Gamma_D$ denote the free nodes, where possibly $\mathcal{K}_F = \mathcal{K}$ for $\Gamma_D = \emptyset$. Then, P_h can be chosen as

$$P_h v = \sum_{z \in \mathcal{K}_F} \left(\int_{T_z} v \psi_z \, dx \right) \zeta_z$$

with appropriate $T_z \in \mathcal{T}$ and $\psi_z \in \mathcal{P}^1(T_z)$.

Definition. The Scott-Zhang projection is just a special example of a Clément-type quasi-interpolation operator: We say that an operator $J_h: H^1_D(\Omega) \to \mathcal{S}^1_D(\mathcal{T})$ is a **Clément-type** quasi-interpolation operator if, for all $v \in H^1_D(\Omega)$ and all $T \in \mathcal{T}$, it holds that

• it is locally H^1 -stable

$$\|\nabla(1 - J_h)v\|_{L^2(T)} \le C \|\nabla v\|_{L^2(\Omega_T)},\tag{4.25}$$

• and has a local first-order approximation property

$$||(1 - J_h)v||_{L^2(T)} \le Ch_T ||\nabla v||_{L^2(\Omega_T)}. \tag{4.26}$$

The constant C > 0 may only depend on γ -shape regularity of \mathcal{T} (and possibly the shapes of possible patches in \mathcal{T}).

For the a posteriori error analysis, we shall need the following simple consequence.

Lemma 4.10. Suppose that $J_h: H_D^1(\Omega) \to \mathcal{S}_D^1(\mathcal{T})$ is a Clément-type operator, i.e., (4.25)–(4.26) hold. Let $T \in \mathcal{T}$ and $E \in \mathcal{E}_T$. Then, it holds that

$$\|(1 - J_h)v\|_{L^2(E)} \le Ch_E^{1/2} \|\nabla v\|_{L^2(\Omega_T)} \quad \text{for all } v \in H_D^1(\Omega). \tag{4.27}$$

The constant C > 0 depends only on γ -shape regularity of \mathcal{T} .

Proof. We apply the trace inequality

$$||w||_{L^{2}(E)}^{2} \lesssim h_{T}^{-1}(||w||_{L^{2}(T)}^{2} + h_{T} ||w||_{L^{2}(T)} ||\nabla w||_{L^{2}(T)})$$

for $w:=(1-J_h)v\in H^1_D(\Omega)$. With the Clément properties (4.25)–(4.26), this yields that

$$||(1-J_h)v||_{L^2(E)}^2 \lesssim h_T ||\nabla v||_{L^2(\Omega_T)}^2$$

Shape regularity and hence $h_T \simeq h_E$ concludes the proof.

The following example is one further *classical* example of a Clément-type operator. The analysis will be left to the reader, but requires the following simple observation:

Exercise 22. Use a scaling argument to show that

$$C^{-1}h_T \|\nabla v_h\|_{L^{\infty}(T)} \le \|\nabla v_h\|_{L^2(T)} \le \frac{h_T}{\sqrt{2}} \|\nabla v_h\|_{L^{\infty}(T)}$$
 for all $v_h \in \mathcal{P}^m(T)$,

where the constant C > 0 only depends on $\sigma(\mathcal{T})$ and the polynomial degree $m \in \mathbb{N}_0$.

Exercise 23. Let $\mathcal{K}_F := \mathcal{K} \setminus \overline{\Gamma}_D$ denote the free nodes (where possibly $\mathcal{K}_F = \mathcal{K}$ if $\Gamma_D = \emptyset$). Define

$$J_h v := \sum_{z \in \mathcal{K}_F} v_z \zeta_z \quad \text{with} \quad v_z := \frac{1}{|\Omega_z|} \int_{\Omega_z} v \, dx, \tag{4.28}$$

where $\Omega_z \subseteq \overline{\Omega}$ denotes the patch of a node $z \in \mathcal{K}$. Prove that J_h satisfies the following properties:

- (i) $J_h: L^2(\Omega) \to \mathcal{S}_D^1(\mathcal{T})$ is a well-defined linear operator.
- (ii) J_h is locally L^2 -stable, i.e., for all $T \in \mathcal{T}$, it holds that

$$||(1-J_h)v||_{L^2(T)} \le C ||v||_{L^2(\Omega_T)}$$
 for all $v \in L^2(\Omega)$.

(iii) J_h is locally H_D^1 -stable, i.e., for all $T \in \mathcal{T}$, it holds that

$$\|\nabla(1 - J_h)v\|_{L^2(T)} \le C \|\nabla v\|_{L^2(\Omega_T)}$$
 for all $v \in H_D^1(\Omega)$.

(iv) J_h has a local first-order approximation property

$$\|(1-J_h)v\|_{L^2(T)} \le Ch_T \|\nabla v\|_{L^2(\Omega_T)}$$
 for all $v \in H_D^1(\Omega)$.

- (v) $J_h:L^2(\Omega)\to L^2(\Omega)$ as well as $J_h:H^1_D(\Omega)\to H^1_D(\Omega)$ are bounded linear operators.
- (vi) J_h is positivity preserving, i.e., $J_h v \geq 0$ for all $v \in L^2(\Omega)$ with $v \geq 0$.
- (vii) With $\Pi_h: L^2(\Omega) \to \mathcal{P}^0(\mathcal{T})$ the L^2 -orthogonal projection onto $\mathcal{P}^0(\mathcal{T})$, it holds that $J_h\Pi_h = J_h$.

The constant C > 0 depends only on γ -shape regularity of \mathcal{T} .

Exercise 24. Find a counter example which shows that the operator J_h from Exercise 23 is no projection, i.e., it holds that $J_h v_h \neq v_h$ for some $v_h \in \mathcal{S}_D^1(\mathcal{T})$.

4.3 Residual-Based Error Estimator

Residual-based a posteriori error estimates follow a general strategy. Recall that the weak solution $u \in H_D^1(\Omega)$ of (4.1) solves the variational form

$$(\nabla u \,; \nabla v)_{L^{2}(\Omega)} = (f \,; v)_{L^{2}(\Omega)} + (\phi \,; v)_{L^{2}(\Gamma_{N})} \quad \text{for all } v \in H^{1}_{D}(\Omega). \tag{4.29}$$

For an approximation $u_h \in \mathcal{S}_D^1(\mathcal{T})$ it is thus natural to define the **residual** $R_h \in H_D^1(\Omega)^*$ by

$$R_h(v) := (f; v)_{L^2(\Omega)} + (\phi; v)_{L^2(\Gamma_N)} - (\nabla u_h; \nabla v)_{L^2(\Omega)}, \tag{4.30}$$

i.e., $R_h = 0$ if and only if $u_h = u$. Let $||w|| := ||\nabla w||_{L^2(\Omega)}$ denote the energy norm on $H_D^1(\Omega)$ and

$$||\!| \Phi |\!|\!|_* := \sup_{w \in H^1_D(\Omega) \setminus \{0\}} \frac{\Phi(w)}{|\!|\!| w |\!|\!|}$$

the induced operator norm on $H_D^1(\Omega)^*$, where we stress that both are equivalent norms on $H_D^1(\Omega)$ and its dual space, respectively. Then, the Riesz theorem and $R_h(v) = (\nabla (u - u_h); \nabla v)_{L^2(\Omega)}$ yield

$$|||R_h|||_* = |||u - u_h|||.$$

To derive a reliable error estimator η , we thus need to prove an estimate of the type

$$R_h(v) \le \widetilde{C_{\text{rel}}} \eta ||v|| \quad \text{for all } v \in H_D^1(\Omega).$$
 (4.31)

To derive an efficient error estimator η , we need to show

$$R_h(v) \ge \widetilde{C_{\text{eff}}} \eta |\!|\!| v |\!|\!| \quad \text{for some } v \in H_D^1(\Omega) \setminus \{0\},$$
 (4.32)

where this $v \in H_D^1(\Omega)$ has to be constructed appropriately.

Exercise 25. Prove that reliability (4.2) of an error estimator η is, in fact, equivalent to (4.31). Prove that efficiency (4.3) of η holds if and only if (4.32) holds.

So far, our observations did not use that we are dealing with Galerkin schemes. We stress that the Galerkin orthogonality reads

$$R_h(v_h) = 0 \quad \text{for all } v_h \in \mathcal{S}_D^1(\mathcal{T})$$
 (4.33)

with respect to the residual R_h . To provide a reliable (and residual-based) error estimator η , we will use some Clément-type operator $J_h: H^1(\Omega) \to \mathcal{S}_D^1(\Omega)$ in connection with the Galerkin orthogonality (4.33).

Before introducing a first a posteriori error estimator, we introduce the following notational conventions. We define the \mathcal{T} -piecewise resp. \mathcal{E} -piecewise constant mesh-width functions

$$h_{\mathcal{T}}|_T := h_T$$
 and $h_{\mathcal{E}}|_E := h_E$

for elements $T \in \mathcal{T}$ and edges $E \in \mathcal{E}$, respectively. Moreover, we write

$$||h_{\mathcal{E}}^{1/2}\psi||_{L^{2}(\mathcal{E}_{*})} := \left(\sum_{E\in\mathcal{E}_{*}} h_{E}||\psi||_{L^{2}(E)}^{2}\right)^{1/2}$$

for any set $\mathcal{E}_* \subseteq \mathcal{E}$ of edges and any function ψ which belongs to $L^2(E)$ for all $E \in \mathcal{E}_*$. Recall that \mathcal{E}_D and \mathcal{E}_N denote the Dirichlet and Neumann edges of \mathcal{T} , respectively. Moreover, let \mathcal{E}_Ω denote the set of all **interior edges**, i.e., for $E \in \mathcal{E}_\Omega$, there are unique elements $T_E^+, T_E^- \in \mathcal{T}$ with $E = T_E^+ \cap T_E^-$. Finally, for $E \in \mathcal{E}_\Omega$, we define the **jump of the normal derivative** by

$$[\![\partial_n u_h]\!]_E := \frac{\partial u_h}{\partial n_E^+} + \frac{\partial u_h}{\partial n_E^-} \in \mathbb{R}, \tag{4.34}$$

where n_E^{\pm} denote the outer normal vectors of the elements T_E^{\pm} on the edge E. Note that $n_E^+ = -n_E^-$ so that the sum in the definition is, in fact, a difference.

Theorem 4.11. The error estimator

$$\eta := \left(\|h_{\mathcal{T}} f\|_{L^{2}(\Omega)}^{2} + \|h_{\mathcal{E}}^{1/2} [\![\partial_{n} u_{h}]\!]\|_{L^{2}(\mathcal{E}_{\Omega})}^{2} + \|h_{\mathcal{E}}^{1/2} (\phi - \partial_{n} u_{h})\|_{L^{2}(\mathcal{E}_{N})}^{2} \right)^{1/2} \tag{4.35}$$

satisfies the reliability estimate

$$||u - u_h||_{H^1(\Omega)} \le C \,\eta,\tag{4.36}$$

where the constant C > 0 depends only on γ -shape regularity of \mathcal{T} .

Proof. For all $w \in H_D^1(\Omega)$, elementwise integration by parts proves

$$\begin{split} R_h(w) &= (f \; ; w)_{L^2(\Omega)} + (\phi \; ; w)_{L^2(\Gamma_N)} - \sum_{T \in \mathcal{T}} (\nabla u_h \; ; \nabla w)_{L^2(T)} \\ &= (f \; ; w)_{L^2(\Omega)} + \sum_{E \in \mathcal{E}_N} (\phi \; ; w)_{L^2(E)} - \sum_{T \in \mathcal{T}} (\partial_n u_h \; ; w)_{L^2(\partial T)} \\ &= \sum_{T \in \mathcal{T}} (f \; ; w)_{L^2(T)} + \sum_{E \in \mathcal{E}_N} (\phi - \partial_n u_h \; ; w)_{L^2(E)} - \sum_{E \in \mathcal{E}_\Omega} ([\![\partial_n u_h]\!] \; ; w)_{L^2(E)} \\ &\leq \sum_{T \in \mathcal{T}} \|f\|_{L^2(T)} \|w\|_{L^2(T)} + \sum_{E \in \mathcal{E}_N} \|\phi - \partial_n u_h\|_{L^2(E)} \|w\|_{L^2(E)} + \sum_{E \in \mathcal{E}_\Omega} \|[\![\partial_n u_h]\!]\|_{L^2(E)} \|w\|_{L^2(E)}. \end{split}$$

For arbitrary $v \in H_D^1(\Omega)$, we now choose $w = v - J_h v$ and note that $R_h(v) = R_h(w)$ according to the Galerkin orthogonality. Then, we estimate the three sums separately. The approximation property of the Clément-type operator J_h and Lemma 4.3 imply

$$\sum_{T \in \mathcal{T}} \|f\|_{L^{2}(T)} \|v - J_{h}v\|_{L^{2}(T)} \lesssim \left(\sum_{T \in \mathcal{T}} \|h_{\mathcal{T}}f\|_{L^{2}(T)}^{2}\right)^{1/2} \left(\sum_{T \in \mathcal{T}} \|\nabla v\|_{L^{2}(\Omega_{T})}^{2}\right)^{1/2}
\lesssim \left(\sum_{T \in \mathcal{T}} \|h_{\mathcal{T}}f\|_{L^{2}(T)}^{2}\right)^{1/2} \left(\sum_{T \in \mathcal{T}} \|\nabla v\|_{L^{2}(T)}^{2}\right)^{1/2}
= \|h_{\mathcal{T}}f\|_{L^{2}(\Omega)} \|\nabla v\|_{L^{2}(\Omega)}.$$

For each edge $E \in \mathcal{E}$, we choose an arbitrary element $T_E \in \mathcal{T}$ with $E \in \mathcal{E}_{T_E}$. Let $\mathcal{E}_* \subset \mathcal{E}$ and $\psi \in L^2(E)$ for all $E \in \mathcal{E}_*$. Recall that $\|(1 - J_h)v\|_{L^2(E)} \lesssim h_E^{1/2} \|\nabla v\|_{L^2(\Omega_{T_E})}$. Therefore, the same arguments as before prove

$$\sum_{E \in \mathcal{E}_*} \|\psi\|_{L^2(E)} \|v - J_h v\|_{L^2(E)} \lesssim \left(\sum_{E \in \mathcal{E}_*} \|h_{\mathcal{E}}^{1/2} \psi\|_{L^2(E)}^2\right)^{1/2} \left(\sum_{E \in \mathcal{E}_*} \|\nabla v\|_{L^2(\Omega_{T_E})}^2\right)^{1/2}$$

$$\lesssim \|h_{\mathcal{E}}^{1/2} \psi\|_{L^2(\mathcal{E}_*)} \|\nabla v\|_{L^2(\Omega)},$$

where we note that an element $T \in \mathcal{T}$ may satisfy $T = T_E$ for at most three edges. Altogether, we now see

$$R_{h}(v) \lesssim \|\nabla v\|_{L^{2}(\Omega)} \left(\|h_{\mathcal{T}}f\|_{L^{2}(\Omega)} + \|h_{\mathcal{E}}^{1/2} [\![\partial_{n}u_{h}]\!]\|_{L^{2}(\mathcal{E}_{\Omega})} + \|h_{\mathcal{E}}^{1/2} (\phi - \partial_{n}u_{h})\|_{L^{2}(\mathcal{E}_{N})} \right)$$

$$\leq \sqrt{3} \|\nabla v\|_{L^{2}(\Omega)} \eta.$$

The hidden constant C depends only (on the Clément operator J_h and) on γ -shape regularity of \mathcal{T} .

Remark. Note that we have used $u_h \in \mathcal{S}^1(\mathcal{T})$ in the sense that the elementwise Laplacian satisfies $\Delta u_h|_{\mathcal{T}} = 0$ for all $T \in \mathcal{T}$. For general \mathcal{T} -piecewise polynomials, the same proof applies with $\|h_{\mathcal{T}}f\|_{L^2(\Omega)}$ replaced by $\|h_{\mathcal{T}}(f + \Delta u_h)\|_{L^2(\Omega)}$.

Exercise 26. We consider the mixed boundary value problem

$$-\Delta u = f \quad \text{in } \Omega,$$

$$u = u_D \quad \text{on } \Gamma_D,$$

$$\partial_n u = \phi \quad \text{on } \Gamma_N.$$

with inhomogeneous Dirichlet data $u_D \in H^{1/2}(\Gamma_D)$. Let $u \in H^1(\Omega)$ denote the weak solution and $u_h \in \mathcal{S}^1(\mathcal{T}_h)$ the P1-FEM solution for discrete Dirichlet data $u_{Dh} := \widehat{u}_{Dh}|_{\Gamma_D}$ with $\widehat{u}_{Dh} \in \mathcal{S}^1(\mathcal{T}_h)$. Use the additional problem

$$-\Delta w = 0 \qquad \text{in } \Omega,$$

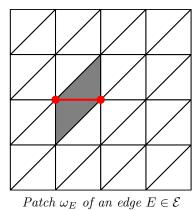
$$w = u_D - u_{Dh} \quad \text{on } \Gamma_D,$$

$$\partial_n u = 0 \quad \text{on } \Gamma_N.$$

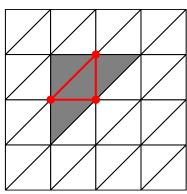
with weak solution $w \in H^1(\Omega)$ to derive a reliable error estimator for $||u - u_h||_{H^1(\Omega)}$.

Hint. Prove that $||w||_{H^1(\Omega)} \simeq ||u_D - u_{Dh}||_{H^{1/2}(\Gamma_D)}$, where the right-hand side is already an a posteriori term. Then, consider the residual $\widetilde{R}_h \in H^1_D(\Omega)^*$ corresponding to the function $(u - u_h) - w \in H^1_D(\Omega)$.

Next, we prove the efficiency of the residual-based error estimator η from (4.35) — at least up to terms of higher order. The efficiency estimate even holds locally with refined patches ω_E and ω_T shown in Figure 4.2: For an interior edge $E \in \mathcal{E}_{\Omega}$, let $T_E^+, T_E^- \in \mathcal{T}$ be the unique elements with







Patch ω_T of an element $T \in \mathcal{T}$

FIGURE 4.2. To prove the efficiency estimate, it suffices to consider smaller patches $\omega_E \subseteq \Omega_E$ and $\omega_T \subseteq \Omega_T$, for edges $E \in \mathcal{E}$ and elements $T \in \mathcal{T}$, respectively. For comparison with the larger patches Ω_E and Ω_T , the reader may consider Figure 4.1 on page 40.

 $E = T_E^+ \cap T_E^-$. For a boundary edge $E \in \mathcal{E}_{\Gamma}$, there is a unique element $T_E \in \mathcal{T}$ with $E \in \mathcal{E}_{T_E}$. We define the refined patch of an edge $E \in \mathcal{E}$ by

$$\omega_E := \begin{cases} T_E^+ \cup T_E^- & \text{for } E \in \mathcal{E}_{\Omega}, \\ T_E & \text{for } E \in \mathcal{E}_{\Gamma}. \end{cases}$$

$$(4.37)$$

Moreover, we define the refined patch of an element $T \in \mathcal{T}$ by

$$\omega_T := \bigcup \left\{ \omega_E \mid E \in \mathcal{E}_T \right\}. \tag{4.38}$$

Note that $\omega_E \subseteq \Omega_E$ and $\omega_T \subseteq \Omega_T$, so that Lemma 4.3 and Lemma 4.4 even hold for the refined patches.

Usually, one is interested in error estimators which are localized with respect to the elements or the edges of \mathcal{T} , respectively. For instance, one considers the **element-based residual error** estimator

$$\eta_{\mathcal{T}} := \left(\sum_{T \in \mathcal{T}} \eta_T^2\right)^{1/2},\tag{4.39}$$

where

$$\eta_T = \left(h_T^2 \|f\|_{L^2(T)}^2 + h_T \|[[\partial_n u_h]]\|_{L^2(\partial T \cap \Omega)}^2 + h_T \|\phi - \partial_n u_h\|_{L^2(\partial T \cap \Gamma_N)}^2\right)^{1/2}$$
(4.40)

or the edge-based residual error estimator

$$\eta_{\mathcal{E}} := \left(\sum_{E \in \mathcal{E}} \eta_E^2\right)^{1/2},\tag{4.41}$$

where

$$\eta_{E} = \begin{cases}
\left(h_{E}^{2} \|f\|_{L^{2}(\omega_{E})}^{2} + h_{E} \| [\partial_{n} u_{h}] \|_{L^{2}(E)}^{2}\right)^{1/2} & \text{for } E \in \mathcal{E}_{\Omega}, \\
\left(h_{E}^{2} \|f\|_{L^{2}(\omega_{E})}^{2} + h_{E} \|\phi - \partial_{n} u_{h}\|_{L^{2}(E)}^{2}\right)^{1/2} & \text{for } E \in \mathcal{E}_{N}, \\
0 & \text{for } E \in \mathcal{E}_{D}.
\end{cases}$$
(4.42)

Alternatively, one could also define

$$\eta_{\mathcal{T}\cup\mathcal{E}} := \left(\sum_{T\in\mathcal{T}} \tilde{\eta}_T^2 + \sum_{E\in\mathcal{E}} \tilde{\eta}_E^2\right)^{1/2},\tag{4.43}$$

where

$$\tilde{\eta}_T = h_T \|f\|_{L^2(T)},\tag{4.44a}$$

$$\tilde{\eta}_{E} = \begin{cases} h_{E}^{1/2} \| [\![\partial_{n} u_{h}]\!] \|_{L^{2}(E)} & \text{for } E \in \mathcal{E}_{\Omega}, \\ h_{E}^{1/2} \| \phi - \partial_{n} u_{h} \|_{L^{2}(E)} & \text{for } E \in \mathcal{E}_{N}, \\ 0 & \text{for } E \in \mathcal{E}_{D}. \end{cases}$$
(4.44b)

Note that $\eta_{\mathcal{T}}$ as well as $\eta_{\mathcal{E}}$ are equivalent to the error estimator η from (4.35): There holds

$$\eta = \eta_{\mathcal{T} \cup \mathcal{E}} \le \eta_{\mathcal{T}} \le \sqrt{2} \, \sigma(\mathcal{T}_h)^{1/2} \, \eta$$
 as well as $\sigma(\mathcal{T})^{-1} \, \eta \le \eta_{\mathcal{E}} \le \sqrt{3} \, \eta$,

since η_T adds the contributions of interior edges twice and $h_E \leq h_T \leq \sigma(\mathcal{T}_h) h_E$ for each edge $E \in \mathcal{E}_T$, whereas $\eta_{\mathcal{E}}$ adds the element contribution at most three times. The local quantities η_T and η_E can be used to steer an adaptive mesh-refining algorithm. They are therefore called **refinement indicators**. We are going to discuss adaptive mesh-refinement below.

Theorem 4.12 (inverse estimate). For all polynomial degrees $m \in \mathbb{N}$ and $k, r \in \mathbb{N}$ with k > r, there exists a constant C > 0 such that

$$||D^k v_h||_{L^2(T)} \le C\sigma(\mathcal{T}) h_T^{r-k} ||D^r v_h||_{L^2(T)} \quad \text{for all } v_h \in \mathcal{P}^m(\mathcal{T}) \text{ and all } T \in \mathcal{T},$$

$$(4.45)$$

where
$$\mathcal{P}^m(\mathcal{T}) := \{ v_h : \Omega \to \mathbb{R} \mid \forall T \in \mathcal{T} \quad v_h|_T \in \mathcal{P}^m(T) \}.$$

Proof. The proof is done \mathcal{T} -elementwise and follows from a scaling argument. We start with the estimate on the reference element.

1. step. We prove

$$||D^k w_h||_{L^2(T_{\text{ref}})} \le C_{\text{ref}} ||D^r w_h||_{L^2(T_{\text{ref}})} \quad \text{for all } w_h \in \mathcal{P}^m(T_{\text{ref}}).$$
 (4.46)

To this end we introduce an interpolation operator $I: \mathcal{P}^m(T_{\text{ref}}) \to \mathcal{P}^{r-1}(T_{\text{ref}})$ via

$$(Iv)(z_j) = v(z_j), \qquad j = 1, \dots, N_r,$$

for $N_r = \dim \mathcal{P}^{r-1}(T_{\text{ref}})$ piecewise different points $z_j \in T_{\text{ref}}$. It is straightforward to see that

$$|\!|\!|\!|v|\!|\!|\!| := |\!|\!|D^r v|\!|\!|_{L^2(T_{\mathrm{ref}})} + \sum_{j=1}^{N_r} |v(z_j)|$$

is a norm on $\mathcal{P}^m(T_{\text{ref}})$. Moreoever, since $\mathcal{P}^m(T_{\text{ref}})$ is finite dimensional, it is equivalent to the norm $\|\cdot\|_{H^k(T_{\text{ref}})}$. Due to $D^k v_h = D^r v_h = 0$ for all $v_h \in \mathcal{P}^{r-1}(T_{\text{ref}})$ and $(Iw_h)(z_j) = w_h(z_j)$ for all $w_h \in \mathcal{P}^m(T_{\text{ref}})$, we have

$$\|D^k w_h\|_{L^2(T_{\text{ref}})} = \|D^k (w_h - I w_h)\|_{L^2(T_{\text{ref}})} \le \|w_h - I w_h\|_{H^k(T_{\text{ref}})} \le C_{\text{ref}} \|w_h - I w_h\| = C_{\text{ref}} \|D^r w_h\|_{L^2(T_{\text{ref}})}.$$

2. step. Proof of the statement: Let $\Phi: T_{\text{ref}} \to T$ be an affine diffeomorphism and $B \in \mathbb{R}^{2 \times 2}$ its linear part. We apply the transformation formula to Φ^{-1} to see that

$$||D^k v_h||_{L^2(T)} \le |\det B^{-1}|^{-1/2} ||B^{-1}||_F^k ||D^k (v_h \circ \Phi)||_{L^2(T_{\text{ref}})}.$$

Note that the L^2 -norm can be estimated by step 1 since $v_h \circ \Phi \in \mathcal{P}^m(T_{\text{ref}})$. The application of the transformation formula to Φ proves that

$$||D^r(v_h \circ \Phi)||_{L^2(T_{ref})} \le |\det B|^{-1/2} ||B||_F^r ||D^r v_h||_{L^2(T)}.$$

By definition of the shape regularity constant $\sigma(\mathcal{T})$, we obtain that

$$||D^{k}v_{h}||_{L^{2}(T)} \leq C_{\text{ref}} ||B^{-1}||_{F}^{k} ||B||_{F}^{r} ||D^{r}v_{h}||_{L^{2}(T)} \leq \sqrt{2} C_{\text{ref}} \varrho_{T}^{-k} h_{T}^{r} ||D^{r}v_{h}||_{L^{2}(T)}$$

$$\leq \sqrt{2} C_{\text{ref}} \sigma(\mathcal{T}) h_{T}^{r-k} ||D^{r}v_{h}||_{L^{2}(T)},$$

where we have used that $||B^{-1}||_F \leq \sqrt{2} \, \varrho_T^{-1}$. This concludes the proof.

Theorem 4.13. We define $f_{\mathcal{T}} \in \mathcal{P}^0(\mathcal{T})$ by $f_{\mathcal{T}}|_T := |T|^{-1} \int_T f \, dx$ and $\phi_{\mathcal{E}} \in \mathcal{P}^0(\mathcal{E}_N)$ by $\phi_{\mathcal{E}}|_E := h_E^{-1} \int_E \phi \, ds$. For each element $T \in \mathcal{T}$, the refinement indicator η_T from (4.40) satisfies

$$\eta_T \le C \left(\|\nabla(u - u_h)\|_{L^2(\omega_T)}^2 + \|h_{\mathcal{T}}(f - f_{\mathcal{T}})\|_{L^2(\omega_T)}^2 + \|h_{\mathcal{E}}^{1/2}(\phi - \phi_{\mathcal{E}})\|_{L^2(\partial T \cap \Gamma_N)}^2 \right)^{1/2}. \tag{4.47}$$

Moreover, the error estimator η from (4.35) is efficient in the sense that

$$\eta \le C \left(\|u - u_h\|_{H^1(\Omega)} + \|h_{\mathcal{T}}(f - f_{\mathcal{T}})\|_{L^2(\Omega)} + \|h_{\mathcal{E}}^{1/2}(\phi - \phi_{\mathcal{E}})\|_{L^2(\Gamma_N)} \right). \tag{4.48}$$

The constant C > 0 only depends on the shape regularity constant $\sigma(\mathcal{T})$.

Remark. For $f \in H^1(\mathcal{T})$ holds $||h_{\mathcal{T}}(f - f_{\mathcal{T}})||_{L^2(\Omega)} = \mathcal{O}(h^2)$. For $\phi \in C^1(\mathcal{E}_N)$ holds $||h_{\mathcal{E}}^{1/2}(\phi - \phi_{\mathcal{E}})||_{L^2(\Gamma_N)} = \mathcal{O}(h^{3/2})$. Even for $u \in H^2(\Omega)$, the error as well as the error estimator η only satisfy $||u - u_h||_{H^1(\Omega)} = \mathcal{O}(h) = \eta$. Therefore, the two terms on the right-hand side are of higher order. \square

Proof of Theorem 4.13. 1. step. Estimate (4.48) is a consequence of (4.47) since

$$\eta \leq \eta_{\mathcal{T}} = \left(\sum_{T \in \mathcal{T}} \eta_T^2\right)^{1/2} \leq 2C \left(\|u - u_h\|_{H^1(\Omega)} + \|h_{\mathcal{T}}(f - f_{\mathcal{T}})\|_{L^2(\Omega)} + \|h_{\mathcal{E}}^{1/2}(\phi - \phi_{\mathcal{E}})\|_{L^2(\Gamma_N)} \right).$$

Here, the factor $2 = 4^{1/2}$ appears since each element $T \in \mathcal{T}$ belongs at most to four patches $\omega_{T'}$.

The proof of (4.47) is split into three steps, where we consider each of the three contributions of η_T separately.

2. step. There holds

$$||h_{\mathcal{T}}f||_{L^{2}(T)} \le C\left(||\nabla(u - u_{h})||_{L^{2}(T)} + ||h_{\mathcal{T}}(f - f_{\mathcal{T}})||_{L^{2}(T)}\right): \tag{4.49}$$

For $T \in \mathcal{T}$, we define the **element bubble function**

$$b_T := \prod_{z \in \mathcal{K}_T} \zeta_z \in H_0^1(T) \cap \mathcal{P}^3(T)$$

as product of all three hat functions. It is essential to observe the following estimate

$$||f_{\mathcal{T}}b_T||_{L^2(T)} \le ||f_{\mathcal{T}}b_T^{1/2}||_{L^2(T)} \le ||f_{\mathcal{T}}||_{L^2(T)} \le C_{\text{ref}}||f_{\mathcal{T}}b_T||_{L^2(T)},\tag{4.50}$$

where the existence of an independent constant $C_{\text{ref}} > 0$ follows from a scaling argument. We stress, however, that $||f_{\mathcal{T}}b_T||_{L^2(T)}$ and hence C_{ref} —since $f_{\mathcal{T}}$ is constant on T—can explicitly be computed. In the following, the main idea is to use integration by parts for $v := f_{\mathcal{T}}b_T \in H_0^1(T)$ to show

$$C_{\text{ref}}^{-2} \|f_{\mathcal{T}}\|_{L^{2}(T)}^{2} \leq \|f_{\mathcal{T}}b_{T}^{1/2}\|_{L^{2}(T)}^{2} = (f_{\mathcal{T}}; v)_{L^{2}(T)} = (f_{\mathcal{T}} - f; v)_{L^{2}(T)} + (f - \Delta u_{h}; v)_{L^{2}(T)}$$
$$= (f_{\mathcal{T}} - f; v)_{L^{2}(T)} + (\nabla (u - u_{h}); \nabla v)_{L^{2}(T)}.$$

Now, we estimate each of the two scalar products on the right-hand side by use of the Cauchy inequality. Together with $v = f_T b_T \in \mathcal{P}^3(\mathcal{T})$ we observe

$$(f_{\mathcal{T}} - f; v)_{L^{2}(T)} \le ||f_{\mathcal{T}} - f||_{L^{2}(T)} ||f_{\mathcal{T}} b_{T}||_{L^{2}(T)} \le ||f_{\mathcal{T}} - f||_{L^{2}(T)} ||f_{\mathcal{T}}||_{L^{2}(T)}$$

as well as

$$\begin{split} (\nabla(u - u_h) \; ; \; \nabla v)_{L^2(T)} &\leq \|\nabla(u - u_h)\|_{L^2(T)} \|\nabla(f_T b_T)\|_{L^2(T)} \\ &\leq C_{\text{inv}} h_T^{-1} \|\nabla(u - u_h)\|_{L^2(T)} \|f_T b_T\|_{L^2(T)} \\ &\leq C_{\text{inv}} h_T^{-1} \|\nabla(u - u_h)\|_{L^2(T)} \|f_T\|_{L^2(T)}. \end{split}$$

Altogether, we see

$$h_T \| f_T \|_{L^2(T)} \le C_{\text{ref}}^2 (h_T \| f_T - f \|_{L^2(T)} + C_{\text{inv}} \| \nabla (u - u_h) \|_{L^2(T)}),$$

which finally results in

$$h_T \|f\|_{L^2(T)} \le (1 + C_{\text{ref}}^2) (h_T \|f_T - f\|_{L^2(T)} + C_{\text{inv}} \|\nabla(u - u_h)\|_{L^2(T)}).$$

3. step. For an interior edge $E \in \mathcal{E}_{\Omega}$, there holds

$$h_E^{1/2} \| [\![\partial_n u_h]\!] \|_{L^2(E)} \le C \left(\| \nabla (u - u_h) \|_{L^2(\omega_E)} + \| h_{\mathcal{T}}(f - f_{\mathcal{T}}) \|_{L^2(\omega_E)} \right) : \tag{4.51}$$

To prove this estimate, we define the edge bubble function

$$b_E := \prod_{z \in \mathcal{K}_E} \zeta_z \in H_0^1(\omega_E) \cap \mathcal{P}^2(\mathcal{T}).$$

The essential estimate reads

$$||b_E||_{L^2(E)} \le ||b_E^{1/2}||_{L^2(E)} \le h_E^{1/2} \le C_{\text{ref}} ||b_E||_{L^2(E)},$$
 (4.52)

where the constant $C_{\text{ref}} > 0$ is independent of E. In particular, this provides

$$C_{\text{ref}}^{-2} \| \llbracket \partial_n u_h \rrbracket \|_{L^2(E)}^2 \le \| \llbracket \partial_n u_h \rrbracket b_E^{1/2} \|_{L^2(E)}^2 = (\llbracket \partial_n u_h \rrbracket \; ; \; \llbracket \partial_n u_h \rrbracket b_E)_{L^2(E)}.$$

Let $T_E^+, T_E^- \in \mathcal{T}$ be the unique elements with $T_E^+ \cap T_E^- = E$ and $\omega_E = T_E^+ \cup T_E^-$. Note that $v := [\![\partial_n u_h]\!] b_E \in \mathcal{P}^2(T_E^\pm)$ satisfies $v|_{\partial T_E^\pm \setminus E} = 0$. Therefore, integration by parts on T_E^\pm proves

$$\begin{split} (\llbracket \partial_n u_h \rrbracket \; ; \, v)_{L^2(E)} &= (\partial_n u_h \; ; \, v)_{L^2(\partial T_E^+)} + (\partial_n u_h \; ; \, v)_{L^2(\partial T_E^-)} \\ &= (\nabla u_h \; ; \, \nabla v)_{L^2(\omega_E)} \\ &= (\nabla (u_h - u) \; ; \, \nabla v)_{L^2(\omega_E)} + (f \; ; \, v)_{L^2(\omega_E)} \\ &\leq \left(C_{\mathrm{inv}} \| \nabla (u_h - u) \|_{L^2(\omega_E)} + \| h_{\mathcal{T}} f \|_{L^2(\omega_E)} \right) \| h_{\mathcal{T}}^{-1} v \|_{L^2(\omega_E)}, \end{split}$$

where we have applied the Cauchy inequality and an inverse estimate for $v \in \mathcal{P}^2(\mathcal{T})$. For $T \in \{T_E^+, T_E^-\}$ holds

$$||v||_{L^{2}(T)} = |[[\partial_{n}u_{h}]]_{E}||b_{E}||_{L^{2}(T)} \leq |T|^{1/2}|[[\partial_{n}u_{h}]]_{E}| \leq \frac{h_{T}^{1/2}}{\sqrt{2}}||[[\partial_{n}u_{h}]]||_{L^{2}(E)},$$

since $|T| \leq \frac{1}{2}h_T h_E$. From this, we infer

$$h_E^{1/2} \|h_{\mathcal{T}}^{-1}v\|_{L^2(\omega_E)} \le \|h_{\mathcal{T}}^{-1/2}v\|_{L^2(\omega_E)} \le \|[\![\partial_n u_h]\!]\|_{L^2(E)}.$$

This finally proves

$$h_E^{1/2} \| [\![\partial_n u_h]\!] \|_{L^2(E)}^2 \le C_{\text{ref}}^2 \big(C_{\text{inv}} \| \nabla (u_h - u) \|_{L^2(\omega_E)} + \| h_{\mathcal{T}} f \|_{L^2(\omega_E)} \big) \| [\![\partial_n u_h]\!] \|_{L^2(E)}$$

and we may conclude this step by use of step 2 to dominate $||h_{\mathcal{T}}f||_{L^2(\omega_E)}$.

4. step. For $T \in \mathcal{T}$ and a Neumann edge $E \in \mathcal{E}_N \cap \mathcal{E}_T$, it holds

$$h_{E}^{1/2} \|\phi - \partial_{n} u_{h}\|_{L^{2}(E)} \leq C \left(\|h_{\mathcal{E}}^{1/2} (\phi - \phi_{\mathcal{E}})\|_{L^{2}(E)} + \|\nabla (u - u_{h})\|_{L^{2}(T)} + \|h_{\mathcal{T}} (f - f_{\mathcal{T}})\|_{L^{2}(T)} \right) :$$

$$(4.53)$$

We consider again the edge bubble function $b_E \in \mathcal{P}^2(T)$ and note that $b_E|_{\partial T \setminus E} = 0$. With $v := (\phi_{\mathcal{E}} - \partial_n u_h)b_E \in \mathcal{P}^2(T)$, we proceed as in step 3 and obtain

$$C_{\text{ref}}^{-2} \|\phi_{\mathcal{E}} - \partial_n u_h\|_{L^2(E)}^2 \le (\phi_{\mathcal{E}} - \partial_n u_h ; v)_{L^2(E)} = (\phi_{\mathcal{E}} - \phi ; v)_{L^2(E)} + (\phi - \partial_n u_h ; v)_{L^2(E)}.$$

For the second term, we employ integration by parts to see

$$(\phi - \partial_n u_h ; v)_{L^2(E)} = (\partial_n (u - u_h) ; v)_{L^2(\partial T)}$$

$$= (\nabla (u - u_h) ; \nabla v)_{L^2(T)} - (f ; v)_{L^2(T)}$$

$$\leq (C_{\text{inv}} \|\nabla (u - u_h)\|_{L^2(T)} + \|h_{\mathcal{T}} f\|_{L^2(T)}) h_F^{-1/2} \|\phi_{\mathcal{E}} - \partial_n u_h\|_{L^2(E)}.$$

The first term is estimated by the Cauchy inequality directly

$$(\phi_{\mathcal{E}} - \phi; v)_{L^{2}(E)} \le \|h_{\mathcal{E}}^{1/2}(\phi_{\mathcal{E}} - \phi)\|_{L^{2}(E)} \|h_{\mathcal{E}}^{-1/2}v\|_{L^{2}(E)}.$$

There holds

$$||v||_{L^{2}(E)} = |(\phi_{\mathcal{E}} - \partial_{n}u_{h})|_{E}|||b_{E}||_{L^{2}(E)} \le h_{E}^{1/2}|(\phi_{\mathcal{E}} - \partial_{n}u_{h})|_{E}| = ||\phi_{\mathcal{E}} - \partial_{n}u_{h}||_{L^{2}(E)}.$$

Altogether, we thus have shown

$$h_{E}^{1/2} \|\phi_{\mathcal{E}} - \partial_{n} u_{h}\|_{L^{2}(E)}^{2}$$

$$\leq C_{\text{ref}}^{2} \left(C_{\text{inv}} \|\nabla(u - u_{h})\|_{L^{2}(T)} + \|h_{\mathcal{T}} f\|_{L^{2}(T)} + \|h_{\mathcal{E}}^{1/2} (\phi_{\mathcal{E}} - \phi)\|_{L^{2}(E)} \right) \|\phi_{\mathcal{E}} - \partial_{n} u_{h}\|_{L^{2}(E)}.$$

Here, $||h_{\mathcal{T}}f||_{L^2(T)}$ is estimated by step 2, and $\phi_{\mathcal{E}}$ on the left-hand side is replaced by ϕ with the help of the triangle inequality.

Exercise 27. Prove that $f_{\mathcal{T}}$ in Theorem 4.13 can be replaced by an arbitrary \mathcal{T} -elementwise polynomial $f_{\mathcal{T}} \in \mathcal{P}^m(\mathcal{T})$. The constant C > 0 in (4.47)–(4.48) then additionally depends on the polynomial degree $m \in \mathbb{N}_0$.

Remark. With the help of a so-called extension operator that extends a polynomial $p: E \to \mathbb{R}$ to a polynomial $F_{\text{ext}}p: T \to \mathbb{R}$, one can show that $\phi_{\mathcal{E}}$ in Theorem 4.13 can be replaced by an arbitrary \mathcal{E}_N -edgewise polynomial (with respect to the arclength).

Actually, the volume residual contribution $||h_{\mathcal{T}}f||_{L^2(\Omega)} = \mathcal{O}(h)$ to η can be improved. This is done in the following exercise, where this term is replaced by some higher-order term $\mathcal{O}(h^2)$.

Exercise 28. Let $\Omega_z = \text{supp}(\zeta_z)$ denote the node patch of $z \in \mathcal{K}$. For $f \in L^2(\Omega)$, let $f_z := |\Omega_z|^{-1} \int_{\Omega_z} f \, dx$ denote the corresponding integral mean. Prove the following claims:

(i) For all inner nodes $z \in \mathcal{K} \backslash \Gamma$, it holds

$$\int_{\Omega_z} f f_z \zeta_z \, dx \le C \Big(\sum_{\substack{E \in \mathcal{E}_{\Omega} \\ z \in E}} \| [\![\partial_n u_h]\!] \|_{L^2(E)}^2 \Big)^{1/2} \| h_{\mathcal{T}}^{-1/2} f_z \|_{L^2(\Omega_z)}.$$

(ii) For all inner nodes $z \in \mathcal{K} \setminus \Gamma$ and elements $T \in \mathcal{T}$ with $z \in T$, it holds

$$C^{-1} \|h_{\mathcal{T}}f\|_{L^{2}(T)}^{2} \leq \|h_{\mathcal{T}}(f - f_{z})\|_{L^{2}(\Omega_{z})}^{2} + \sum_{\substack{E \in \mathcal{E}_{\Omega} \\ z \in E}} \|h_{\mathcal{E}}^{1/2} [\![\partial_{n} u_{h}]\!]\|_{L^{2}(E)}^{2}.$$

(iii) Derive the equivalence

$$C^{-1}\eta^{2} \leq \widetilde{\eta}^{2} := \|h_{\mathcal{E}}^{1/2} [\![\partial_{n} u_{h}]\!]\|_{L^{2}(\mathcal{E}_{\Omega})}^{2} + \|h_{\mathcal{E}}^{1/2} (\phi - \partial_{n} u_{h})\|_{L^{2}(\mathcal{E}_{N})}^{2} + \sum_{z \in \mathcal{K} \setminus \Omega} \|h_{\mathcal{T}} (f - f_{z})\|_{L^{2}(\Omega_{z})}^{2} \leq C \eta^{2}.$$

(iv) Conclude that the improved error estimator $\tilde{\eta}$ is reliable and efficient.

(v) In what sense is the error estimator $\tilde{\eta}$ improved when compared to η .

The constant C > 0 in (i)–(iii) depends only on γ -shape regularity of \mathcal{T} .

4.4 Adaptive Mesh-Refining Algorithm

Usually, a posteriori error estimates are not only used to estimate the (unknown) error $\|\nabla(u - u_h)\|_{L^2(\Omega)}$ but even to steer the local mesh-refinement. Let

$$\eta := \left(\sum_{T \in \mathcal{T}} \eta(T)^2\right)^{1/2}$$

be an a posteriori error estimator, where the quantities $\eta(T) := \eta_T$ reflect —at least heuristically—the (unknown) local error $\|\nabla(u - u_h)\|_{L^2(T)}$ for all $T \in \mathcal{T}$. We then aim to refine only the elements $T \in \mathcal{T}$, where $\eta(T)$ is large. Therefore, the quantities $\eta(T)$ are usually called **refinement indicators** (or error indicators). To state our version of an adaptive algorithm, we introduce some additional notation which will be used from now on.

- the index $\ell \in \mathbb{N}_0$ denotes the step of the adaptive algorithm,
- \mathcal{T}_{ℓ} is the mesh in the ℓ -th step of the adaptive algorithm.
- \mathcal{N}_{ℓ} and \mathcal{E}_{ℓ} denote the associated sets of nodes and edges, respectively.
- $U_{\ell} \in \mathcal{X}_{\ell} := \mathcal{S}_{D}^{1}(\mathcal{T}_{\ell})$ denotes the Galerkin solution in the ℓ -th step.
- $h_{\ell} \in \mathcal{P}^0(\mathcal{T}_{\ell}), h_{\ell}|_T := \operatorname{diam}(T)$ is the local mesh-side function.

With this notation, one common strategy is the following: Let $\theta \in (0,1)$ be the parameter for the adaptive algorithm.

Algorithm 4.14 (Adaptive Mesh-Refinement). Input: Initial triangulation \mathcal{T}_0 , tolerance $\tau > 0$, adaptivity parameter $\theta \in (0,1)$, counter $\ell := 0$.

- (i) Compute discrete solution U_{ℓ} .
- (ii) Compute refinement indicators $\eta_{\ell}(T)$ and error estimator $\eta_{\ell} = \left(\sum_{T \in \mathcal{T}_{\ell}} \eta_{\ell}(T)^2\right)^{1/2}$.
- (iii) Stop computation provided that $\eta_{\ell} \leq \tau$
- (iv) Choose the minimal set $\mathcal{M}_{\ell} \subseteq \mathcal{T}_{\ell}$ of marked elements such that

$$\theta \,\eta_{\ell}^2 = \theta \, \sum_{T \in \mathcal{T}_{\ell}} \eta_{\ell}(T)^2 \le \sum_{T \in \mathcal{M}_{\ell}} \eta_{\ell}(T)^2. \tag{4.54}$$

- (v) Generate a new regular mesh $\mathcal{T}_{\ell+1}$, where at least all marked elements have been refined.
- (vi) Update $\ell \mapsto \ell + 1$ and goto (i).

Output: Finite sequence of discrete solutions U_{ℓ} and corresponding error estimators η_{ℓ} .

Remark. Clearly, the stopping criterion (iii) is only meaningful if η_{ℓ} is reliable and if the reliability constant in $\|\nabla(u-U_{\ell})\|_{L^2(\Omega)} \leq C_{\text{rel}} \eta_{\ell}$ is known. In practice, runtime and storage requirements are the limiting quantities for a numerical simulation. Usually, one thus uses rather a maximal runtime or a maximal storage requirement, e.g., the maximal number of elements, as a stopping criterion. Adaptivity is then used to obtain an —in some sense— optimal approximation with respect to these side constraints.

Remark. The marking criterion (4.54) was introduced by DÖRFLER (1996). It will be crucial to prove convergence of U_{ℓ} to the exact solution $u \in H_D^1(\Omega)$ of (4.1). Note that the choice $\theta \to 0$ in (4.54) leads to highly adapted meshes, whereas $\theta \to 1$ corresponds to (almost) uniform mesh-refinement. However, for small θ , only a few elements are refined per step. This may result in too many steps in the sense that usually the assembly of the Galerkin data is the most time consuming part of the algorithm. In practice, a good compromise between sufficient mesh-adaption and as few steps in the loop as possible appears to be $\theta \approx 0.25$.

Remark. In the beginning of the analysis of adaptive FEM, Babuška proposed the following marking criterion: An element $T \in \mathcal{T}$ is marked for refinement if and only if

$$\eta_T \ge \theta \, \max \big\{ \eta_{T'} \, \big| \, T' \in \mathcal{T} \big\}, \tag{4.55}$$

which is called **bulk criterion** in the literature. Convergence (but *not* optimality) of this version of adaptive FEM was proven by MORIN, SIEBERT & VEESER (2008). Very recently, DIENING, KREUZER & STEVENSON (2014) proved the so-called *instance optimality* of the adaptive algorithm for some extended bulk criterion.

4.4.1 Red-Green-Blue Refinement

It now remains to discuss the mesh-refinement. Recall that all error estimates are affected by the shape regularity $\sigma(\mathcal{T}_{\ell})$ in the sense that the involved constants become unbounded for $\sigma(\mathcal{T}_{\ell}) \xrightarrow{\ell \to \infty} \infty$. Therefore, the mesh-refinement has to take care of the interior angles of the elements $T \in \mathcal{T}_{\ell}$ since $\sigma(\mathcal{T}_{\ell})$ tends to infinity if and only if the minimal interior angle of the triangulation tends to zero. We follow the so-called **red-green-blue strategy** (or **RGB-refinement**): This refinement strategy is based on edge-refinement. First, we thus use the following marking rule:

• If an element $T \in \mathcal{T}_{\ell}$ is marked for refinement, we mark all edges $E \in \mathcal{E}_T$ for refinement.

We now proceed recursively as follows:

• For each element $T \in \mathcal{T}_{\ell}$, we mark its longest edge $E \in \mathcal{E}_T$ for refinement provided that \mathcal{E}_T contains a marked edge.

Each marked edge will be halved, i.e., the midpoint m_E of a marked edge belongs to the new set $\mathcal{K}_{\ell+1}$ of nodes. Finally, we have the following refinement rules, for all $T \in \mathcal{T}_{\ell}$:

- If no edge in \mathcal{E}_T is marked for refinement, T is not refined, i.e., $T \in \mathcal{T}_{\ell+1}$.
- If all edges in \mathcal{E}_T are marked, we use a **red-refinement** of T, i.e., T is refined uniformly into four similar triangles, cf. Figure 4.3.
- If one edge in \mathcal{E}_T is marked (and hence the longest edge), we use a **green-refinement**, i.e., T is refined into two triangles, cf. Figure 4.4.
- If two edges in \mathcal{E}_T are marked one of which is, according to the marking rule, the longest edge of T —, we use a **blue-refinement**, i.e., T is split into three triangles, cf. Figure 4.5.

In Figure 4.6, we visualize a simple example for an RGB-refined mesh.



FIGURE 4.3. Red-refinement: If all edges of a triangle $T \in \mathcal{T}_{\ell}$ are marked (left), T is refined into four similar triangles $T_1, T_2, T_3, T_4 \in \mathcal{T}_{\ell+1}$ (right).



FIGURE 4.4. Green-refinement: If only the longest edge of a triangle $T \in \mathcal{T}_{\ell}$ is marked (left), T is refined into two new triangles $T_1, T_2 \in \mathcal{T}_{\ell+1}$ (right).

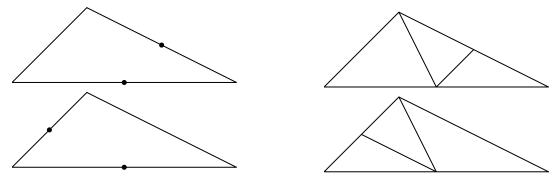
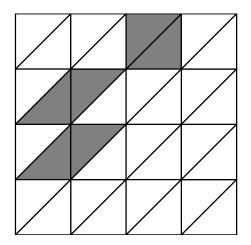


FIGURE 4.5. Blue-refinement: If besides the longest edge of a triangle $T \in \mathcal{T}_{\ell}$ just one other edge is marked for refinement (left), T is refined into three new triangles $T_1, T_2, T_3 \in \mathcal{T}_{\ell+1}$ (right).

We state the following elementary but important theorem without a proof.



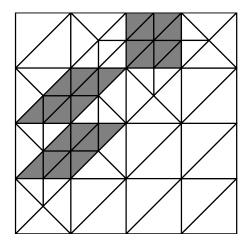


FIGURE 4.6. The left plot shows an initial mesh \mathcal{T}_{ℓ} with marked elements coloured in grey. The right plot shows the mesh $\mathcal{T}_{\ell+1}$ obtained by RGB-refinement of the marked elements. The grey elements are obtained by uniform refinement of a marked element $T \in \mathcal{T}_0$.

Theorem 4.15. Let \mathcal{T}_0 be a regular triangulation such that $\varepsilon > 0$ is a lower bound for the smallest angle of a triangle $T \in \mathcal{T}_0$. Let \mathcal{T}_ℓ be a sequence of meshes, where \mathcal{T}_ℓ is obtained by RGB-refinement of the mesh $\mathcal{T}_{\ell-1}$ and where the set $\mathcal{M}_{\ell-1} \subseteq \mathcal{T}_{\ell-1}$ of marked elements is arbitrary. Then, \mathcal{T}_ℓ is regular and the smallest angle of all triangles $T \in \mathcal{T}_\ell$ is bounded from below by $\varepsilon/2$. In particular, there holds

$$\sup_{\ell \in \mathbb{N}} \sigma(\mathcal{T}_{\ell}) < \infty, \tag{4.56}$$

which is an equivalent formulation for the fact that the smallest angles of the triangulations \mathcal{T}_{ℓ} do not tend to zero.

4.5 Convergence of Adaptive FEM

In the following, we aim to show that Algorithm 4.14 creates a sequence U_{ℓ} of discrete solutions which converges to the exact solution $u \in H := H_D^1(\Omega)$. The adaptive algorithm generates a sequence $\mathcal{X}_{\ell} = \mathcal{S}_D^1(\mathcal{T}_{\ell})$ of finite dimensional nested subspaces of H, i.e., $\mathcal{X}_{\ell} \subsetneq \mathcal{X}_{\ell+1}$ for all $\ell \in \mathbb{N}_0$, since $\mathcal{T}_{\ell+1}$ is some refinement of \mathcal{T}_{ℓ} . We first stress that the sequence U_{ℓ} is always convergent to some limit $U_{\infty} \in H$. However, we even stress that one may in general expect that $U_{\infty} \neq u$.

Exercise 29. Let \mathcal{X}_{ℓ} be nested subspaces of a Hilbert space H, i.e., $\mathcal{X}_{\ell} \subseteq \mathcal{X}_{\ell+1}$ for all $\ell \in \mathbb{N}_0$. Let $\langle \cdot : \cdot \rangle$ be an elliptic and continuous bilinear form on H with corresponding Galerkin solutions $U_{\ell} \in \mathcal{X}_{\ell}$. Prove that the limit $U_{\infty} := \lim_{\ell \to \infty} U_{\ell}$ exists in H. Hint: Define \mathcal{X}_{∞} as the closure of $\bigcup_{\ell=0}^{\infty} \mathcal{X}_{\ell}$ in H. Let $U_{\infty} \in \mathcal{X}_{\infty}$ be the corresponding Galerkin solution, and prove that U_{∞} is the limit of the sequence U_{ℓ} .

Exercise 30. Let $H = H_D^1(\Omega)$ and $\mathcal{X}_{\ell} = \mathcal{S}_D^1(\mathcal{T}_{\ell})$, where the regular initial mesh \mathcal{T}_0 is given and where \mathcal{T}_{ℓ} is obtained iteratively by uniform refinement of $\mathcal{T}_{\ell-1}$. Prove that $\mathcal{X}_{\infty} = H$ for the space \mathcal{X}_{∞} from Exercise 29.

The interpretation of the last exercises is the following: For uniform mesh-refinement, there usually holds $\mathcal{X}_{\infty} = H$ and thus $u = U_{\infty}$, i.e., we have convergence of the sequence of discrete solutions U_{ℓ} towards the unique solution u. However, adaptive mesh-refinement may lead to $\mathcal{X}_{\infty} \subsetneq H$. Consequently, the question arrises whether the adaptive algorithm guarantees $U_{\infty} = u$ or not. This will be discussed in the following sections.

Throughout the subsequent section, we use the following notation, which is now collected for the convenience of the reader:

- $U_{\ell} \in \mathcal{X}_{\ell} := \mathcal{S}_{D}^{1}(\mathcal{T}_{\ell})$ denotes the Galerkin solution.
- For $T \in \mathcal{T}_{\ell}$ and some $V \in \mathcal{S}_D^1(\mathcal{T}_{\ell})$, $\eta_{\ell}(T,V)$ denotes the associated refinement indicator, e.g.,

$$\eta_{\ell}(T, V)^{2} = h_{T}^{2} \|f\|_{L^{2}(T)}^{2} + h_{T} \| [\![\partial_{n} V]\!] \|_{L^{2}(\partial T \cap \Omega)}^{2} + h_{T} \| \phi - \partial_{n} V \|_{L^{2}(\partial T \cap \Gamma_{N})}^{2}. \tag{4.57}$$

- For some subset $\mathcal{M} \subseteq \mathcal{T}_{\ell}$ and $V \in \mathcal{S}_D^1(\mathcal{T}_{\ell})$, let $\eta_{\ell}(\mathcal{M}, V) := \left(\sum_{T \in \mathcal{M}} \eta_{\ell}(T, V)^2\right)^{1/2}$.
- We abbreviate $\eta_{\ell}(\mathcal{M}) = \eta_{\ell}(\mathcal{M}, U_{\ell})$ and $\eta_{\ell} = \eta_{\ell}(\mathcal{T}_{\ell})$.

Note that in case of (4.57), η_{ℓ} is the residual a posteriori error estimator discussed in Section 4.3. We recall some technical terms, proven above for the residual error estimator η_{ℓ} .

• η_{ℓ} is **reliable** if

$$||u - U_\ell||_H \le C_{\text{rel}} \eta_\ell. \tag{4.58}$$

• η_{ℓ} is **efficient** (up to oscillation terms which depend only on \mathcal{T}_{ℓ}), if

$$\eta_{\ell} \le C_{\text{eff}} \left(\|u - U_{\ell}\|_{H} + \text{osc}_{\ell} \right), \tag{4.59}$$

where $\operatorname{osc}_{\ell} := \operatorname{osc}_{\ell}(\mathcal{T}_{\ell}), \operatorname{osc}_{\ell}(\mathcal{M}) := \left(\sum_{T \in \mathcal{M}} \operatorname{osc}_{\ell}(T)^{2}\right)^{1/2} \text{ for } \mathcal{M} \subseteq \mathcal{T}_{\ell}, \text{ and}$

$$\operatorname{osc}_{\ell}(T)^{2} := h_{T}^{2} \|f - f_{T}\|_{L^{2}(T)} + h_{T} \|\phi - \phi_{\mathcal{E}}\|_{L^{2}(\partial T \cap \Gamma_{N})}^{2}. \tag{4.60}$$

• The set $\mathcal{M}_{\ell} \subseteq \mathcal{T}_{\ell}$ of marked elements is usually assumed to satisfy the **Dörfler marking**

$$\theta \, \eta_{\ell} \le \eta_{\ell}(\mathcal{M}_{\ell}) \tag{4.61}$$

for some fixed parameter $\theta \in (0,1)$.

Exercise 31. Prove that $||u - U_{\ell}||_H$ as well as $\operatorname{osc}_{\ell}$ are monotonously decreasing for $\ell \to \infty$. Prove that in case of the residual-based indicators (4.57), there holds $\operatorname{osc}_{\ell}(T) \leq \eta_{\ell}(T)$ for all $T \in \mathcal{T}_{\ell}$, i.e., the error estimator dominates the oscillation terms.

The following convergence theorem is a result of CASCÓN, KREUZER, NOCHETTO & SIEBERT from 2008, where it is proven that the combined error quantity, which consists of error and error estimator, has a contraction property. We stress two important observations:

• For their analysis, Cascón, Kreuzer, Nochetto, and Siebert re-define the mesh width

$$h_T := |T|^{1/2} \quad \text{for } T \in \mathcal{T}_{\ell}, \tag{4.62}$$

whereas we considered diam(T) before. Note that, however, $|T| \leq \text{diam}(T)^2 \leq 2\sigma(T)|T|$ so that both definition are equivalent for shape regular meshes, and we shall use the new definition in what follows.

• If $T \in \mathcal{T}_{\ell}$ is refined, each son $T' \in \mathcal{T}_{\ell+1}$ satisfies at least $|T'| \leq |T|/2$, which now results in a strict reduction $h_{T'} \leq h_T/\sqrt{2}$ of the local mesh-width (which fails, in general, for the usual definition $h_T = \text{diam}(T)$). This observation is used in step 2 of the proof of the following theorem.

We note that the analysis holds for general symmetric problems. For non-symmetric problems, the correspoding result has been open until Feischl, Führer & Praetorius (2014).

Theorem 4.16 (Cascón, Kreuzer, Nochetto & Siebert '08). Suppose that the set of marked elements \mathcal{M}_{ℓ} satisfies the Dörfler marking for some fixed $\theta \in (0,1)$. Then, there are constants $\kappa > 0$ and $q \in (0,1)$ which depend only on θ and uniform γ -shape regularity of \mathcal{T}_{ℓ} for all $\ell \in \mathbb{N}_0$, such that

$$\left(\|\nabla(u - U_{\ell+1})\|_{L^2(\Omega)}^2 + \kappa \,\eta_{\ell+1}^2\right)^{1/2} \le q \left(\|\nabla(u - U_{\ell})\|_{L^2(\Omega)}^2 + \kappa \,\eta_{\ell}^2\right)^{1/2} \text{ for all } \ell \in \mathbb{N}.$$
 (4.63)

In particular, this implies convergence $\lim_{\ell \to \infty} \|\nabla(u - U_{\ell})\|_{L^{2}(\Omega)} = 0 = \lim_{\ell \to \infty} \eta_{\ell}$.

Proof. 1. step. There holds the following quasi-triangle inequality for the error estimator

$$\eta_{\ell}(V) \le \eta_{\ell}(W) + C_{\Delta} \|\nabla(V - W)\|_{L^{2}(\Omega)} \quad \text{for all } V, W \in \mathcal{S}_{D}^{1}(\mathcal{T}_{\ell})$$
 (4.64)

with some constant $C_{\Delta} > 0$ which depends only on $\sigma(\mathcal{T}_{\ell})$: From the triangle inequalities in ℓ_2 and L^2 , we infer

$$\eta_{\ell}(V) = \left[\|h_{\ell}f\|_{L^{2}(\Omega)}^{2} + \sum_{T \in \mathcal{T}_{\ell}} h_{T} \left(\| [\partial_{n}V] \|_{L^{2}(\partial T \cap \Omega)}^{2} + \|\phi - \partial_{n}V\|_{L^{2}(\partial T \cap \Gamma_{N})}^{2} \right) \right]^{1/2} \\
\leq \left[\|h_{\ell}f\|_{L^{2}(\Omega)}^{2} + \sum_{T \in \mathcal{T}_{\ell}} h_{T} \left(\| [\partial_{n}W] \|_{L^{2}(\partial T \cap \Omega)}^{2} + \|\phi - \partial_{n}W\|_{L^{2}(\partial T \cap \Gamma_{N})}^{2} \right) \right]^{1/2} \\
+ \left[\sum_{T \in \mathcal{T}_{\ell}} h_{T} \left(\| [\partial_{n}(V - W)] \|_{L^{2}(\partial T \cap \Omega)}^{2} + \|\partial_{n}(V - W) \|_{L^{2}(\partial T \cap \Gamma_{N})}^{2} \right) \right]^{1/2}.$$

For fixed $T \in \mathcal{T}_{\ell}$ and $E \in \mathcal{E}_T$, a scaling argument proves

$$h_T(\|[[\partial_n(V-W)]]\|_{L^2(E\cap\Omega)}^2 + \|\partial_n(V-W)\|_{L^2(E\cap\Gamma_N)}^2) \lesssim \|\nabla(V-W)\|_{L^2(\omega_E)}^2$$

where the constant depends only on $\sigma(\mathcal{T}_{\ell})$. Consequently, we end up with (4.64).

2. step. There holds an estimator reduction in the sense that there is a constant $\varrho \in (0,1)$ with

$$\eta_{\ell+1}^2 \le (1+\delta)\varrho \,\eta_{\ell}^2 + C_\delta \|\nabla (U_{\ell+1} - U_{\ell})\|_{L^2(\Omega)}^2 \quad \text{for all } \delta > 0,$$
 (4.65)

where $C_{\delta} > 0$ depends only on δ and $C_{\Delta} > 0$. The constant ϱ depends only on θ and the reduction of the mesh-side on marked elements: Let $\Omega_* := \bigcup_{T \in \mathcal{M}_{\ell}} T$ denote the subdomain of Ω , where the elements are marked. Recall that $h_{T'} \leq h_T/\sqrt{2}$ for all sons $T' \in \mathcal{T}_{\ell+1}$ of a marked element $T \in \mathcal{M}_{\ell}$. The crucial step is to observe that the error indicators

$$\eta_{\ell}(T, V)^{2} = h_{T}^{2} \|f\|_{L^{2}(T)}^{2} + h_{T} \| [\![\partial_{n} V]\!] \|_{L^{2}(\partial T \cap \Omega)}^{2} + h_{T} \| \phi - \partial_{n} V \|_{L^{2}(\partial T \cap \Gamma_{N})}^{2}.$$

for $h_T = |T|^{1/2}$ lead to

$$\eta_{\ell+1}(U_{\ell})^{2} = \sum_{\substack{T' \in \mathcal{T}_{\ell+1} \\ T' \subseteq \Omega_{*}}} \eta_{\ell+1}(T', U_{\ell})^{2} + \sum_{\substack{T' \in \mathcal{T}_{\ell+1} \\ T' \subseteq \overline{\Omega} \setminus \Omega_{*}}} \eta_{\ell+1}(T', U_{\ell})^{2} \\
\leq \frac{1}{\sqrt{2}} \sum_{\substack{T \in \mathcal{T}_{\ell} \\ T \subseteq \Omega_{*}}} \eta_{\ell}(T, U_{\ell})^{2} + \sum_{\substack{T \in \mathcal{T}_{\ell} \\ T \subseteq \overline{\Omega} \setminus \Omega_{*}}} \eta_{\ell}(T, U_{\ell})^{2} \\
= 2^{-1/2} \eta_{\ell}(\mathcal{M}_{\ell})^{2} + \eta_{\ell}(\mathcal{T}_{\ell} \setminus \mathcal{M}_{\ell})^{2} \\
= (2^{-1/2} - 1) \eta_{\ell}(\mathcal{M}_{\ell})^{2} + \eta_{\ell}^{2}.$$

By use of the Dörfler marking $\theta \eta_{\ell}^2 \leq \eta_{\ell}(\mathcal{M}_{\ell})^2$, we thus obtain

$$\eta_{\ell+1}^2(U_\ell) \le \eta_\ell^2 - (1 - 2^{-1/2})\eta_\ell(\mathcal{M}_\ell)^2 \le \varrho \, \eta_\ell^2 \quad \text{with} \quad \varrho := (1 - \theta(1 - 2^{-1/2})).$$

Now, Young's inequality and step 1 conclude

$$\eta_{\ell+1}^2 \le (1+\delta)\eta_{\ell+1}(U_{\ell})^2 + (1+\delta^{-1})C_{\Delta}^2 \|\nabla(U_{\ell+1} - U_{\ell})\|_{L^2(\Omega)}^2$$

$$\le (1+\delta)\varrho \,\eta_{\ell}^2 + (1+\delta^{-1})C_{\Delta}^2 \|\nabla(U_{\ell+1} - U_{\ell})\|_{L^2(\Omega)}^2.$$

3. step. Proof of contraction property (4.63): Let $\kappa, \delta, \beta > 0$ be constants which are fixed later. Let $\varrho \in (0,1)$ be the given constant from step 2. We recall the Galerkin orthogonality

$$\|\nabla(u - U_{\ell})\|_{L^{2}(\Omega)}^{2} = \|\nabla(u - U_{\ell+1})\|_{L^{2}(\Omega)}^{2} + \|\nabla(U_{\ell+1} - U_{\ell})\|_{L^{2}(\Omega)}^{2}$$

This and the estimator reduction imply

$$\|\nabla(u - U_{\ell+1})\|_{L^{2}(\Omega)}^{2} + \kappa \eta_{\ell+1}^{2} = \|\nabla(u - U_{\ell})\|_{L^{2}(\Omega)}^{2} - \|\nabla(U_{\ell+1} - U_{\ell})\|_{L^{2}(\Omega)}^{2} + \kappa \eta_{\ell+1}^{2}$$

$$\leq \|\nabla(u - U_{\ell})\|_{L^{2}(\Omega)}^{2} + (\kappa C_{\delta} - 1) \|\nabla(U_{\ell+1} - U_{\ell})\|_{L^{2}(\Omega)}^{2} + \kappa (1 + \delta)\varrho \eta_{\ell}^{2}.$$

Provided that $\kappa C_{\delta} \leq 1$, we infer

$$\begin{split} \|\nabla(u - U_{\ell+1})\|_{L^{2}(\Omega)}^{2} + \kappa \,\eta_{\ell+1}^{2} &\leq \|\nabla(u - U_{\ell})\|_{L^{2}(\Omega)}^{2} + \kappa(1+\delta)\varrho \,\eta_{\ell}^{2} \\ &= \|\nabla(u - U_{\ell})\|_{L^{2}(\Omega)}^{2} - \kappa\beta \,\eta_{\ell}^{2} + \kappa((1+\delta)\varrho + \beta) \,\eta_{\ell}^{2} \end{split}$$

Reliability $\|\nabla(u-U_{\ell})\|_{L^{2}(\Omega)} \leq C_{\mathrm{rel}}\eta_{\ell}$ finally leads to

$$\|\nabla(u - U_{\ell+1})\|_{L^{2}(\Omega)}^{2} + \kappa \,\eta_{\ell+1}^{2} \le (1 - \kappa \beta C_{\text{rel}}^{-2}) \|\nabla(u - U_{\ell})\|_{L^{2}(\Omega)}^{2} + \kappa \left((1 + \delta)\varrho + \beta\right) \,\eta_{\ell}^{2}$$

$$\le \max\left\{1 - \kappa \beta C_{\text{rel}}^{-2}, \, (1 + \delta)\varrho + \beta\right\} \left(\|\nabla(u - U_{\ell})\|_{L^{2}(\Omega)}^{2} + \kappa \,\eta_{\ell}^{2}\right).$$

It remains to choose the constants κ, δ, β so that $q^2 := \max\left\{1 - \kappa \beta C_{\mathrm{rel}}^{-2}, \, (1+\delta)\varrho + \beta\right\} \in (0,1)$:

- $\rho := (1 \theta(1 2^{-1/2})) \in (0, 1).$
- Choose $\delta > 0$ such that $(1 + \delta)\rho < 1$.
- Choose $\kappa > 0$ such that $\kappa C_{\delta} \leq 1$ with $C_{\delta} := (1 + \delta^{-1})C_{\Lambda}^2$.
- Choose $\beta > 0$ such that $(1 + \delta)\rho + \beta < 1$.

This implies $q \in (0,1)$ and concludes the proof.

Remark. We again collect the main arguments of the preceding proof, namely a certain quasitriangle inequality of the estimator (4.64) and a strict reduction $\eta_{\ell+1}(\mathsf{sons}(\mathcal{M}_\ell), U_\ell) \leq \kappa \, \eta_\ell(\mathcal{M}_\ell, U_\ell)$ with some $\kappa \in (0,1)$ based on the strict reduction of the local mesh-width for marked elements. Besides this, only Galerkin orthogonality and Dörfler marking are used. Therefore, the proof works for a quite general class of symmetric problems and a variety of error estimators. The original work of Cascón, Kreuzer, Nochetto & Siebert (2008) considers linear second order symmetric and elliptic problems in divergence form and H^1 -conforming finite element spaces with fixed polynomial degree. Finally, we stress that the proof also works for higher dimensions $d \geq 2$, where $h_T = |T|^{-1/d}$. For 2D, the usual definition $h_T := \operatorname{diam}(T)$ is sufficient if marked elements are refined, e.g., by red-refinement or bisec(3), since then all edges are bisected.

Exercise 32. Prove the following variants of Young's inequality, for all $a, b \in \mathbb{R}$ and $\delta > 0$,

•
$$ab \leq \frac{a^2}{2\delta} + \frac{\delta b^2}{2}$$

•
$$ab \le \frac{a^2}{2\delta} + \frac{\delta b^2}{2}$$
.
• $(a+b)^2 \le (1+\delta^{-1})a^2 + (1+\delta)b^2$.

Chapter 5

A Priori Analysis II

5.1 FEM with Data Approximation

Now that we have realized that the P1-FEM is of order $\mathcal{O}(h)$, we need to show that the quadrature rules used for implementation are sufficiently accurate. For P1-FEM we will show that it is sufficient to approximate the right-hand side of the exact P1-FEM

$$F(v) := \int_{\Omega} fv \, dx + \int_{\Gamma_N} \phi v \, ds \quad \text{for } v \in H^1(\Omega)$$
 (5.1)

by

$$F_h(v_h) := \sum_{T \in \mathcal{T}} |T| f(s_T) v_h(s_T) + \sum_{E \in \mathcal{E}_N} h_E \phi(m_E) v_h(m_E) \quad \text{for } v_h \in \mathcal{S}^1(\mathcal{T}), \tag{5.2}$$

where s_T denotes the center of mass of an element $T \in \mathcal{T}$ and where m_E denotes the midpoint of a Neumann edge $E \in \mathcal{E}_N$. Such an approximation can be considered as a perturbed P1-FEM and we need to study the convergence of this perturbed scheme.

5.1.1 First Strang Lemma

In this section, we go back to the abstract formulation of Galerkin schemes: Let H be a real Hilbert space with norm $\|\cdot\|_H$. Let $\langle\cdot,\cdot\rangle$ be a bilinear form which is assumed to be elliptic and continuous, i.e., it holds that

$$\alpha \|v\|_H^2 \le \langle \langle v; v \rangle \quad \text{as well as} \quad \langle \langle v; w \rangle \rangle \le \beta \|v\|_H \|w\|_H \quad \text{for all } v, w \in H. \tag{5.3}$$

Let $F \in H^*$ be a given right-hand side. Then, the Lax-Milgram lemma applies and yields the existence and uniqueness of the solution $u \in H$ of

$$\langle\!\langle u ; \cdot \rangle\!\rangle = F \in H^*. \tag{5.4}$$

For a discretization parameter h > 0, let X_h be a finite dimensional subspace of H. It is an important property of a Galerkin scheme that it is stable with respect to certain perturbations of the scalar product $\langle \cdot ; \cdot \rangle$ or the right-hand side F. — For the interpretation, recall that usually the right-hand side $F \in H^*$ as well as the scalar product $\langle \cdot ; \cdot \rangle$ involve integrals, which are computed

numerically by quadrature rules. For a fixed discrete space X_h , this leads to perturbations $F_h \in X_h^*$ and $\langle \cdot ; \cdot \rangle_h$ of F and $\langle \cdot ; \cdot \rangle_h$, respectively. In particular, this gives rise to additional **consistency** errors

$$||F - F_h||_{X_h^*}$$
 and $\sup_{v_h \in X_h \setminus \{0\}} \frac{||\langle v_h ; \cdot \rangle - \langle v_h ; \cdot \rangle_h||_{X_h^*}}{||v_h||_H}$.

In practice, the **best approximation error** (or **discretization error**) behaves like

$$\min_{v_h \in X_h} \|u - v_h\|_H = \mathcal{O}(h^{\alpha}) \quad \text{for } h \to 0,$$

where the **convergence order** $\alpha > 0$ usually depends on the regularity of the exact solution u. Then, the Céa lemma proves that

$$||u - \mathbb{G}_h u||_H = \mathcal{O}(h^{\alpha}).$$

The following result due to Strang shows that the consistency errors should be at least of the same order, i.e., one needs a sufficiently large order for the quadrature rules. Then, the perturbed Galerkin scheme

$$\langle \langle u_h ; v_h \rangle \rangle_h = F_h(v_h) \quad \text{for all } v_h \in X_h$$
 (5.5)

still allows for a unique solution $u_h \in X_h$. Moreover the approximation error still satisfies

$$||u - u_h||_H = \mathcal{O}(h^\alpha).$$

However, the consequence of Strang's lemma even works the other way around: You should avoid to compute integrals exactly (or with high accuracy quadrature rules) since this is usually computationally expensive and since this expense does not pay in the sense of an increased order of convergence. Finally, we note that analytic computation of integrals via antiderivatives, i.e., $\int_a^b f \, dx = F(b) - F(a)$ for the simple 1D case, necessarily leads to cancellation for small mesh-sizes. These are, however, avoided for numerical integration via Gaussian quadrature rules, since the Gaussian quadrature weights are all positive. In explicit terms, this implies that approximate computation will be numerically more accurate than analytic computation, if the quadrature rules are deliberately chosen.

Proposition 5.1 (First Strang Lemma). Assume that $\langle \cdot ; \cdot \rangle_h$ is a bilinear form on X_h and that $F_h: X_h \to \mathbb{R}$ is linear. Then, there holds the following:

(i) Assume convergence of $\langle \cdot ; \cdot \rangle_h$ to $\langle \cdot ; \cdot \rangle$, i.e.,

$$\lim_{h \to 0} E_h = 0 \quad with \quad E_h := \sup_{v_h, w_h \in X_h \setminus \{0\}} \frac{|\langle v_h ; w_h \rangle - \langle v_h ; w_h \rangle_h|}{\|v_h\|_H \|w_h\|_H}.$$
 (5.6)

Then, the bilinear forms are uniformly elliptic for small h, i.e.,

$$\exists \alpha_0 > 0 \,\exists h_0 > 0 \,\forall h \in (0, h_0) \,\forall v_h \in X_h \quad \alpha_0 \|v_h\|_H^2 \le \langle\!\langle v_h \,; \, v_h \rangle\!\rangle_h. \tag{5.7}$$

In particular, there exist unique $u_h \in X_h$ with $\langle u_h ; \cdot \rangle_h = F_h \in X_h^*$ for sufficiently small h > 0. (ii) Provided (5.7), there holds the Céa type estimate

$$C^{-1} \|u - u_h\|_H \le \inf_{v_h \in X_h} (\|u - v_h\|_H + \|\langle\langle v_h ; \cdot \rangle\rangle - \langle\langle v_h ; \cdot \rangle\rangle_h \|_{X_h^*}) + \|F - F_h\|_{X_h^*}$$

$$\le (1 + E_h) \min_{v_h \in X_h} \|u - v_h\|_H + E_h \|u\|_H + \|F - F_h\|_{X_h^*}$$
(5.8)

with u being the exact solution of (5.4). The constant C > 0 depends only on $\langle \cdot ; \cdot \rangle$.

Proof. Let $0 < \varepsilon < \alpha$ and $h_0 > 0$ such that

$$\forall h \in (0, h_0) \quad \sup_{v_h \in X_h \setminus \{0\}} \frac{|\langle \langle v_h ; v_h \rangle - \langle \langle v_h ; v_h \rangle \rangle_h|}{\|v_h\|_H^2} \le \varepsilon.$$

Then, $\alpha \|v_h\|_H^2 \leq \langle v_h ; v_h \rangle \leq \langle v_h ; v_h \rangle_h + |\langle v_h ; v_h \rangle_h - \langle v_h ; v_h \rangle_h| \leq \langle v_h ; v_h \rangle_h + \varepsilon \|v_h\|_H^2$, whence

$$(\alpha - \varepsilon) \|v_h\|_H^2 \le \langle\!\langle v_h ; v_h \rangle\!\rangle_h,$$

i.e., $\langle \cdot : \cdot \rangle_h$ is an elliptic bilinear form on X_h for $h < h_0$. This concludes the proof of (i) with $\alpha_0 := \alpha - \varepsilon > 0$. To prove (ii), let $v_h \in X_h$ be arbitrary. Then,

$$\alpha_0 \|v_h - u_h\|_H^2 \le \langle v_h - u_h ; v_h - u_h \rangle_h = \langle v_h ; v_h - u_h \rangle_h - F_h(v_h - u_h).$$

Together with

$$\langle (u - v_h ; v_h - u_h) \rangle = F(v_h - u_h) - \langle (v_h ; v_h - u_h) \rangle$$

we obtain that

$$\alpha_{0} \|v_{h} - u_{h}\|_{H}^{2} \leq \left[F(v_{h} - u_{h}) - F_{h}(v_{h} - u_{h}) \right] + \left[\langle v_{h} ; v_{h} - u_{h} \rangle_{h} - \langle v_{h} ; v_{h} - u_{h} \rangle_{h} \right] - \langle u - v_{h} ; v_{h} - u_{h} \rangle$$

$$\leq \|v_{h} - u_{h}\|_{H} \left[\|F - F_{h}\|_{X_{h}^{*}} + \|\langle v_{h} ; \cdot \rangle_{h} - \langle v_{h} ; \cdot \rangle_{\|X_{h}^{*}} + \beta \|u - v_{h}\|_{H} \right].$$

Finally, the combination with a triangle inequality yields that

$$||u - u_h||_H \le ||u - v_h||_H + ||v_h - u_h||_H$$

$$\le C[||F_h - F||_{X_h^*} + ||\langle v_h ; \cdot \rangle| - \langle v_h ; \cdot \rangle_h ||_{X_h^*} + ||u - v_h||_H]$$

for any $v_h \in X_h$ with $C = 1 + \beta/\alpha_0$. This proves the first estimate in (5.8). To see the second estimate, note that

$$\|\langle\langle v_h ; \cdot \rangle\rangle - \langle\langle v_h ; \cdot \rangle\rangle_h\|_{X_h^*} \le E_h \|v_h\|_H \le E_h \|u\|_H + E_h \|u - v_h\|_H.$$

This concludes the proof.

Under the assumptions of the Strang lemma, one can even show convergence of the perturbed Galerkin scheme.

Exercise 33. Assume that $\langle \cdot ; \cdot \rangle_h$ is a symmetric bilinear form on X_h^* and that $F_h \in X_h^*$. We assume convergence of the data in the sense that

$$\lim_{h \to 0} E_h = 0 = \lim_{h \to 0} \|F - F_h\|_{X_h^*} \quad \text{with} \quad E_h := \sup_{v_h, w_h \in X_h \setminus \{0\}} \frac{|\langle v_h; w_h \rangle - \langle v_h; w_h \rangle_h|}{\|v_h\|_H \|w_h\|_H}.$$
(5.9)

For sufficiently small h > 0, let $u_h \in X_h$ be the unique solutions of the perturbed Galerkin scheme (5.5). Under the approximation assumption

$$\lim_{h \to 0} \min_{v_h \in X_h} \|v - v_h\|_H = 0 \quad \text{for all } v \in D$$
 (5.10)

for some dense subspace D of H, there holds convergence

$$\lim_{h \to 0} \|u - u_h\|_H = 0$$

with u being the exact solution of (5.4).

5.1.2 Approximation of Volume Forces

For the Poisson problem and a P1-FEM, the bilinear form $\langle v_h ; w_h \rangle$ can be computed analytically. Nevertheless, for the right-hand side in most cases quadrature formulas are needed. Let F and F_h be given by (5.1)–(5.2), respectively. According to the first Strang lemma 5.1, we only need to show that

$$||F - F_h||_{\mathcal{S}^1(\mathcal{T})^*} = \mathcal{O}(h)$$

to guarantee that the perturbed P1-FEM is also of order $\mathcal{O}(h)$. We consider the two contributions of the right-hand side separately.

Proposition 5.2. Let $f \in H^2(\mathcal{T})$ and $F(v) := \int_{\Omega} fv \, dx$ for $v \in H^1(\Omega)$. Let $F_h(v_h) := \sum_{T \in \mathcal{T}} |T| f(s_T) v_h(s_T)$ for $v_h \in \mathcal{S}^1(\mathcal{T})$, where $s_T \in \mathbb{R}^2$ denotes the center of mass of an element $T \in \mathcal{T}$. Then, it holds that

$$||F - F_h||_{\mathcal{S}^1(\mathcal{T})^*} \le C ||h^2 \nabla f||_{H^1(\mathcal{T})},$$
 (5.11)

where the constant C > 0 depends only on T_{ref} , but not on Ω , \mathcal{T} , or f.

Proof. The proof is done elementwise. For $T \in \mathcal{T}$ and $w \in H^1(T)$, we define the integral mean $w_T := |T|^{-1} \int_T w \, dx$. According to the Poincaré inequality, it holds that $||w - w_T||_{L^2(T)} \le C_P h_T ||\nabla w||_{L^2(T)}$, where the constant $C_P > 0$ is independent of T and w. Moreover, $w \mapsto w_T$ is the L^2 -orthogonal projection onto $\mathcal{P}^0(T)$.

1. step. It holds that

$$\left| \int_{T} f v_h \, dx - |T| f(s_T) v_h(s_T) \right| \leq C_P^2 h_T^2 \|\nabla f\|_{L^2(T)} \|\nabla v_h\|_{L^2(T)} + \|f_T - f(s_T)\|_{L^2(T)} \|v_h\|_{L^2(T)} :$$

From $\int_T v_h dx = |T|v_h(s_T)$, we infer that

$$\int_{T} f v_{h} dx - |T| f(s_{T}) v_{h}(s_{T}) = (f - f(s_{T}) ; v_{h})_{L^{2}(T)}
= (f - f(s_{T}) ; v_{h} - v_{hT})_{L^{2}(T)} + (f - f(s_{T}) ; v_{hT})_{L^{2}(T)}
= (f - f_{T} ; v_{h} - v_{hT})_{L^{2}(T)} + (f_{T} - f(s_{T}) ; v_{h})_{L^{2}(T)}
\leq ||f - f_{T}||_{L^{2}(T)} ||v_{h} - v_{hT}||_{L^{2}(T)} + ||f_{T} - f(s_{T})||_{L^{2}(T)} ||v_{h}||_{L^{2}(T)}.$$

where we have used orthogonality of $(\cdot)_T$ in the last but one step. The Poincaré inequality concludes the proof of step 1.

2. step. It holds that $||f_T - f(s_T)||_{L^2(T)} \le 2C_{\text{ref}}h_T^2||D^2f||_{L^2(T)}$ with an independent constant $C_{\text{ref}} > 0$, which is obtained from a scaling argument: Let $\Phi: T_{\text{ref}} \to T$ denote an affine diffeomorphism with linear part $B \in \mathbb{R}^{2\times 2}$. Note that

$$f_T = \frac{1}{|T|} \int f \, dx = \frac{|\det B|}{|T|} \int_{T_{\text{ref}}} f \circ \Phi \, dx = 2 \int_{T_{\text{ref}}} f \circ \Phi \, dx = \frac{1}{|T_{\text{ref}}|} \int_{T_{\text{ref}}} f \circ \Phi \, dx = (f \circ \Phi)_{T_{\text{ref}}}.$$

Together with $f(s_T) = (f \circ \Phi)(s_{T_{ref}})$, this yields that

$$||f_T - f(s_T)||_{L^2(T)} = |\det B^{-1}|^{-1/2} ||(f \circ \Phi)_{T_{\text{ref}}} - (f \circ \Phi)(s_{T_{\text{ref}}})||_{L^2(T_{\text{ref}})}.$$

We define $g := f \circ \Phi \in H^2(T_{\text{ref}})$ and consider the operator $A : H^2(T_{\text{ref}}) \to L^2(T_{\text{ref}})$ defined by $Ag := g_{T_{\text{ref}}} - g(s_{T_{\text{ref}}})$. Then, $\mathcal{P}^1(T_{\text{ref}}) \subseteq \ker A$ and continuity of A follows from the Sobolev inequality

$$||Ag||_{L^{2}(T_{\text{ref}})} \leq ||g_{T_{\text{ref}}}||_{L^{2}(T_{\text{ref}})} + |T_{\text{ref}}|^{1/2}|g(s_{T_{\text{ref}}})| \leq ||g||_{L^{2}(T_{\text{ref}})} + |T_{\text{ref}}|^{1/2}||g||_{\infty,T_{\text{ref}}}$$
$$\leq (1 + C_{\text{Sobolev}}|T_{\text{ref}}|^{1/2})||g||_{H^{2}(T_{\text{ref}})}$$

Therefore, the Bramble-Hilbert lemma provides a constant $C_{\text{ref}} > 0$ with $||Ag||_{L^2(T_{\text{ref}})} \le C_{\text{ref}} ||D^2 g||_{L^2(T_{\text{ref}})}$. We conclude the scaling argument by

$$C_{\text{ref}}^{-1} \| (f \circ \Phi)_{T_{\text{ref}}} - (f \circ \Phi)(s_{T_{\text{ref}}}) \|_{L^{2}(T_{\text{ref}})} \le \| D^{2}(f \circ \Phi) \|_{L^{2}(T_{\text{ref}})} \le |\det B|^{-1/2} \| B \|_{F}^{2} \| D^{2}f \|_{L^{2}(T)},$$

which finally leads to

$$||f_T - f(s_T)||_{L^2(T)} \le 2C_{\text{ref}} h_T^2 ||D^2 f||_{L^2(T)}.$$

3. step. It holds that $\left| \int_T f v_h dx - |T| f(s_T) v_h(s_T) \right| \le \max\{C_P^2, 2C_{\text{ref}}\} h_T^2 \|\nabla f\|_{H^1(T)} \|v_h\|_{H^1(T)}$: The combination of step 1 and step 2 proves that

$$\left| \int_{T} f v_{h} dx - |T| f(s_{T}) v_{h}(s_{T}) \right|$$

$$\leq \max\{C_{P}^{2}, 2C_{\text{ref}}\} h_{T}^{2} (\|\nabla f\|_{L^{2}(T)} \|\nabla v_{h}\|_{L^{2}(T)} + \|D^{2} f\|_{L^{2}(T)} \|v_{h}\|_{L^{2}(T)}).$$

Note that the brackets contain an \mathbb{R}^2 -scalar product which is estimated with the help of the Cauchy inequality $ab + cd \leq (a^2 + c^2)^{1/2}(b^2 + d^2)^{1/2}$. This concludes the proof of step 3.

4. step. With $C := \max\{C_P^2, 2C_{\text{ref}}\}$, we finally sum over all elements $T \in \mathcal{T}$ to obtain that

$$|F(v_h) - F_h(v_h)| \le \sum_{T \in \mathcal{T}} \left| \int_T f v_h \, dx - |T| f(s_T) v_h(s_T) \right|$$

$$\le C \sum_{T \in \mathcal{T}} \|h^2 \nabla f\|_{H^1(T)} \|v_h\|_{H^1(T)}$$

$$\le C \left(\sum_{T \in \mathcal{T}} \|h^2 \nabla f\|_{H^1(T)}^2 \right)^{1/2} \left(\sum_{T \in \mathcal{T}} \|v_h\|_{H^1(T)}^2 \right)^{1/2}$$

$$= C \|h^2 \nabla f\|_{H^1(\mathcal{T})} \|v_h\|_{H^1(\Omega)}$$

by use of the Cauchy inequality. This concludes the proof

We stress that the proof does not work for $f \in H^1(\mathcal{T})$ since H^1 -functions are in general discontinuous so that the evaluation of f at s_T is not well-defined. However, for $f \in C^1(\mathcal{T})$, everything works well.

Exercise 34. For $f \in C^1(\mathcal{T})$, define $F \in H^1(\Omega)^*$ and $F_h \in \mathcal{S}^1(\mathcal{T})^*$ as in Proposition 5.2. Then, there holds

$$||F - F_h||_{\mathcal{S}^1(\mathcal{T})^*} \le C ||h\nabla f||_{L^{\infty}(\Omega)}, \tag{5.12}$$

where the constant C > 0 does neither depend on Ω nor \mathcal{T} or f.

However, if the volume force only satisfies $f \in H^1(\mathcal{T})$, one can proceed as follows:

Exercise 35. For $f \in H^1(\mathcal{T})$, define $F \in H^1(\Omega)^*$ as in Proposition 5.2 and $F_h \in \mathcal{S}^1(\mathcal{T})^*$ by $F_h(v_h) := \sum_{T \in \mathcal{T}} |T| f_T v_h(s_T)$, where $f_T := |T|^{-1} \int_T f \, dx$ denotes the integral mean. Then,

$$||F - F_h||_{\mathcal{S}^1(\mathcal{T})^*} \le C ||h^2 \nabla f||_{L^2(\Omega)},$$
 (5.13)

where the constant C > 0 does neither depend on Ω nor \mathcal{T} or f.

5.1.3 Approximation of Neumann Data

Finally, we consider the approximation of the Neumann contribution.

Proposition 5.3. Let $\phi \in C^2(\mathcal{E}_N) := \{ \psi \in L^2(\Gamma_N) \mid \forall E \in \mathcal{E}_N \quad \psi|_E \in C^2(E) \}$ and $F(v) := \int_{\Gamma_N} \phi v \, ds$ for $v \in H^1(\Omega)$. Let $F_h(v_h) := \sum_{E \in \mathcal{E}_N} h_E \phi(m_E) v_h(m_E)$ for $v_h \in \mathcal{S}^1(\mathcal{T})$, where $m_E \in \mathbb{R}^2$ denotes the midpoint of a Neumann edge $E \in \mathcal{E}_N$. With the mesh-size function $h \in L^{\infty}(\Gamma_N)$, $h|_E := h_E = \operatorname{diam}(E)$, it then holds

$$||F - F_h||_{\mathcal{S}^1(\mathcal{T})^*} \le C ||h^{3/2}\phi'||_{C^1(\mathcal{E}_N)} := \max_{E \in \mathcal{E}_N} \left(h_E^{3/2} \max\{||\phi'||_{L^{\infty}(E)}, ||\phi''||_{L^{\infty}(E)}\} \right)$$
 (5.14)

where the constant C > 0 depends only on $\sigma(\mathcal{T})$ and $|\Gamma_N|$.

Proof. We aim to follow the lines of the proof of Proposition 5.2. For a Neumann edge $E \in \mathcal{E}_N$ and $w \in L^2(E)$, let $w_E := h_E^{-1} \int_E w \, ds$ denote the integral mean.

1. step. From $\int_E v_h ds = h_E v_h(m_E)$, we infer that

$$\int_{E} \phi v_{h} ds - h_{E} \phi(m_{E}) v_{h}(m_{E}) = (\phi - \phi(m_{E}) ; v_{h})_{L^{2}(E)}$$

$$= (\phi - \phi(m_{E}) ; v_{h} - v_{hE})_{L^{2}(E)} + (\phi - \phi(m_{E}) ; v_{hE})_{L^{2}(E)}$$

$$= (\phi - \phi_{E} ; v_{h} - v_{hE})_{L^{2}(E)} + (\phi_{E} - \phi(m_{E}) ; v_{h})_{L^{2}(E)}$$

$$\leq \|\phi - \phi_{E}\|_{L^{2}(E)} \|v_{h} - v_{hE}\|_{L^{2}(E)} + \|\phi_{E} - \phi(m_{E})\|_{L^{2}(E)} \|v_{h}\|_{L^{2}(E)}.$$

where we have simply used orthogonality of $(\cdot)_E$. Therefore, the trace inequalities (4.10)–(4.11) vield that

$$\left| \int_{E} \phi v_h \, ds - h_E \phi(m_E) v_h(m_E) \right| \leq C \left(h_E^{1/2} \|\phi - \phi_E\|_{L^2(E)} + h_E^{-1/2} \|\phi_E - \phi(m_E)\|_{L^2(E)} \right) \|v_h\|_{H^1(T)},$$

where $T \in \mathcal{T}$ is an arbitrary element with $E \in \mathcal{E}_T$. The constant C > 0 depends only on $\sigma(\mathcal{T})$ and on $|\Gamma_N|$.

2. step. It holds that $\|\phi - \phi_E\|_{L^2(E)} \le h_E^{3/2} \|\phi'\|_{L^{\infty}(E)}$: Note that $w := \phi - \phi_E \in C^1(E)$ has necessarily a zero $\zeta \in E$. Therefore, the fundamental theorem of calculus proves that

$$|w(x)| = \left| \int_{\zeta}^{x} w' \, ds \right| \le h_E^{1/2} ||w'||_{L^2(E)}.$$

Integration over E thus yields that

$$\|\phi - \phi_E\|_{L^2(E)}^2 = \|w\|_{L^2(E)}^2 = \int_E |w(x)|^2 ds_x \le h_E^2 \|w'\|_{L^2(E)}^2 = h_E^2 \|\phi'\|_{L^2(E)}^2 \le h_E^3 \|\phi'\|_{L^\infty(E)}^2.$$

3. step. It holds that $\|\phi_E - \phi(m_E)\|_{L^2(E)} \le (1/2) h_E^{5/2} \|\phi''\|_{L^{\infty}(E)}$: Let $p \in \mathcal{P}^1(E)$ be a polynomial on E (with respect to the arclength) such that $\phi(m_E) = p(m_E)$ and $\phi'(m_E) = p'(m_E)$. Then,

$$\|\phi_E - \phi(m_E)\|_{L^2(E)} = h_E^{1/2} |\phi_E - \phi(m_E)| = h_E^{-1/2} \left| \int_E \phi \, ds - h_E p(m_E) \right| = h_E^{-1/2} \left| \int_E (\phi - p) \, ds \right|$$

With $w := \phi - p$ and hence $w'' = \phi''$, this implies that

$$\|\phi_E - \phi(m_E)\|_{L^2(E)} \le h_E^{-1/2} \|w\|_{L^1(E)} \le \|w\|_{L^2(E)}.$$

Note that w as well as w' have zeros at the edge midpoint m_E . Therefore, the same arguments as in step 2 (with the zero $\zeta = m_E$ and hence integration along a segment of length $h_E/2$) prove that

$$\|w\|_{L^{2}(E)}^{2} \le \frac{h_{E}^{2}}{2} \|w'\|_{L^{2}(E)}^{2}$$
 as well as $\|w'\|_{L^{2}(E)}^{2} \le \frac{h_{E}^{2}}{2} \|w''\|_{L^{2}(E)}^{2} = \frac{h_{E}^{2}}{2} \|\phi''\|_{L^{2}(E)}^{2}$.

Altogether, we see that

$$\|\phi_E - \phi(m_E)\|_{L^2(E)}^2 \le \|w\|_{L^2(E)}^2 \le \frac{h_E^4}{4} \|\phi''\|_{L^2(E)}^2 \le \frac{h_E^5}{4} \|\phi''\|_{L^\infty(E)}^2.$$

4. step. The combination of the preceding steps proves that

$$\left| \int_{E} \phi v_{h} ds - h_{E} \phi(m_{E}) v_{h}(m_{E}) \right| \leq C h_{E}^{2} (\|\phi'\|_{L^{\infty}(E)} + \|\phi''\|_{L^{\infty}(E)}) \|v_{h}\|_{H^{1}(T)}$$

$$\leq 2C h_{E}^{2} \|\phi'\|_{C^{1}(E)} \|v_{h}\|_{H^{1}(T)}$$

$$\leq 2C h_{E}^{1/2} \|h^{3/2} \phi'\|_{C^{1}(\mathcal{E}_{N})} \|v_{h}\|_{H^{1}(T)}$$

by definition of $||w||_{C^1(E)} := \max\{||w||_{L^{\infty}(E)}, ||w'||_{L^{\infty}(E)}\}.$

5. step. We obtain the final result by summing over all Neumann edges $E \in \mathcal{E}_N$: For each $E \in \mathcal{E}_N$ we choose an element $T_E \in \mathcal{T}$ with $E \in \mathcal{E}_T$. Note that the element T_E can arise at most 3 times. Therefore,

$$|F(v_h) - F_h(v_h)| \le 2C \|h^{3/2}\phi'\|_{C^1(\mathcal{E}_N)} \sum_{E \in \mathcal{E}_N} h_E^{1/2} \|v_h\|_{H^1(T_E)}$$

$$\le 2C \|h^{3/2}\phi'\|_{C^1(\mathcal{E}_N)} \Big(\sum_{E \in \mathcal{E}_N} h_E\Big)^{1/2} \Big(\sum_{E \in \mathcal{E}_N} \|v_h\|_{H^1(T_E)}^2\Big)^{1/2}$$

$$\le 2\sqrt{3} C |\Gamma_N|^{1/2} \|h^{3/2}\phi'\|_{C^1(\mathcal{E}_N)} \Big(\sum_{T \in \mathcal{T}} \|v_h\|_{H^1(T)}^2\Big)^{1/2}$$

$$= 2\sqrt{3} C |\Gamma_N|^{1/2} \|h^{3/2}\phi'\|_{C^1(\mathcal{E}_N)} \|v_h\|_{H^1(\Omega)}.$$

This concludes the proof.

Exercise 36. Consider the homogenous Dirichlet problems

$$-\Delta u = 1 \quad \text{in } \Omega,$$

$$u = 0 \quad \text{on } \Gamma = \partial \Omega,$$

with Ω being either the square $\Omega = (-1,1)^2$ or the *L*-shaped domain $\Omega = (-1,1)^2 \setminus [0,1]^2$. Which experimental convergence rates $||u - u_h|| = \mathcal{O}(h^{\alpha})$ are observed? Do you expect that the solutions belong to $H^2(\Omega)$? **Hint:** For a convergent sequence $(x_j)_{j \in \mathbb{N}}$, the Δ^2 -sequence reads

$$y_j = x_j - \frac{(x_{j+1} - x_j)^2}{x_{j+2} - 2x_{j+1} + x_j}.$$

Under certain assumptions on $(x_j)_{j\in\mathbb{N}}$ the sequence $(y_j)_{j\in\mathbb{N}}$ then converges faster to $\lim_{j\to\infty} x_j$. \square

5.2 Inhomogeneous Dirichlet Data

Under the usual assumptions of the mixed boundary value problem of Section 2.3.2, we consider the boundary value problem

$$-\Delta u = f \quad \text{in } \Omega,$$

$$u = u_D \quad \text{on } \Gamma_D,$$

$$\partial_n u = \phi \quad \text{on } \Gamma_N.$$
(5.15)

The only difference to the problem treated above is the fact, that the Dirichlet data u_D might be nontrivial. A function $u \in C^2(\overline{\Omega})$ that solves (5.15) is called **strong solution** of (5.15), and the formulation (5.15) is called the **strong form** of the boundary value problem. A function $u \in H^1(\Omega)$ is **weak solution** of (5.15) provided that

$$\gamma u|_{\Gamma_D} = u_D \tag{5.16a}$$

$$(\nabla u; \nabla v)_{L^2(\Omega)} = (f; v)_{L^2(\Omega)} + (\phi; \gamma v)_{L^2(\Gamma_N)} \quad \text{for all } v \in H^1_D(\Omega).$$
 (5.16b)

These two equations are referred to as the **weak form** of the boundary value problem (5.15). Note that the variational part (5.16b) of the weak form is the same as for the mixed boundary value problem with homogeneous Dirichlet conditions $u_D = 0$.

The following proposition shows that (5.15) and (5.16) are essentially equivalent and that the weak solution is unique. The unique solvability, however, needs certain assumptions on the Dirichlet data: If (5.16) has a solution $u \in H^1(\Omega)$, then it holds that $\gamma u|_{\Gamma_D} = u_D$, i.e., u_D can be extended from Γ_D to a function $\widehat{u}_D \in H^1(\Omega)$. With the same arguments as above, cf. Exercise 9 on page 18, one shows that

$$H^{1/2}(\Gamma_D) := \left\{ \gamma u |_{\Gamma_D} \, \middle| \, u \in H^1(\Omega) \right\} \quad \text{with norm} \quad \|v\|_{H^{1/2}(\Gamma_D)} = \inf \left\{ \|\widehat{v}\|_{H^1(\Omega)} \, \middle| \, \gamma \widehat{v}|_{\Gamma_D} = v \right\}$$

is a Hilbert space. Moreover, $H^{1/2}(\Gamma_D)$ is continuously embedded into $L^2(\Gamma_D)$, and the restriction operator $(\cdot)|_{\Gamma_D}: H^{1/2}(\Gamma) \to H^{1/2}(\Gamma_D)$ is well-defined and continuous.

Proposition 5.4. (i) Provided that $u \in C^2(\overline{\Omega})$ solves the strong form (5.15), u solves also the weak form (5.16).

- (ii) Provided that $f \in C(\overline{\Omega})$, $\phi \in C(\overline{\Gamma}_N)$, and $u_D \in C(\overline{\Gamma}_D)$ and that the weak solution $u \in H^1(\Omega)$ of (5.16) additionally satisfies $u \in C^2(\overline{\Omega})$, then u even solves the strong form (5.15).
- (iii) Let $\widehat{u}_D \in H^1(\Omega)$ be an arbitrary extension of the Dirichlet data $u_D \in H^{1/2}(\Gamma)$. Given $f \in L^2(\Omega)$ and $\phi \in L^2(\Gamma_N)$, there exists a unique $u_0 \in H^1_D(\Omega)$ such that

$$(\nabla u_0 ; \nabla v)_{L^2(\Omega)} = (f ; v)_{L^2(\Omega)} - (\nabla \widehat{u}_D ; \nabla v)_{L^2(\Omega)} + (\phi ; \gamma v)_{L^2(\Gamma_N)} \quad \text{for all } v \in H_D^1(\Omega).$$
(5.17)

- (iv) Under the assumptions of (iii), a function $u \in H^1(\Omega)$ with $\gamma u|_{\Gamma_D} = u_D$ solves the weak form (5.16), if and only if $u_0 := u \widehat{u}_D \in H^1_D(\Omega)$ solves (5.17).
- (v) Under the assumptions of (iii), there exists a unique weak solution $u \in H^1(\Omega)$ of (5.16). Contrary to $u_0 \in H^1_D(\Omega)$, however, the function $u \in H^1(\Omega)$ does not depend on the special choice of \widehat{u}_D .
- (vi) The weak solution $u \in H^1(\Omega)$ satisfies

$$||u||_{H^{1}(\Omega)} \leq C_{1} \left(\sup_{v \in H^{1}_{D}(\Omega) \setminus \{0\}} \frac{(f ; v)_{L^{2}(\Omega)}}{||v||_{H^{1}(\Omega)}} + \sup_{w \in H^{1/2}(\Gamma_{N}) \setminus \{0\}} \frac{(\phi ; w)_{L^{2}(\Gamma_{N})}}{||w||_{H^{1/2}(\Gamma_{N})}} + ||u_{D}||_{H^{1/2}(\Gamma_{D})} \right)$$

$$\leq C_{2} \left(||f||_{L^{2}(\Omega)} + ||\phi||_{L^{2}(\Gamma_{N})} + ||u_{D}||_{H^{1/2}(\Gamma_{D})} \right)$$

$$(5.18)$$

where the constants $C_1, C_2 > 0$ only depend on Ω and Γ_D .

Proof. Note that the variational form (5.16b) does not consider whether u_D is zero or not. Therefore, the same proofs as for the mixed boundary value problem with homogeneous Dirichlet data apply to prove (i) and (ii). To verify (iii), simply note that the left-hand side of (5.17) defines an equivalent scalar product on the Hilbert space $H_D^1(\Omega)$. The right-hand side is linear and continuous on $H_D^1(\Omega)$. Therefore, existence and uniqueness of u_0 follows from the Riesz theorem. (iv) is obvious, and (v) thus an immediate consequence of (iii) and (iv). To prove the stability estimate, we argue as for the homogeneous Dirichlet conditions. With the Friedrichs inequality, we see that

$$\begin{split} C_F^{-2} \|u_0\|_{H^1(\Omega)}^2 &\leq \|\nabla u_0\|_{L^2(\Omega)}^2 \\ &= (\nabla u_0 \; ; \; \nabla u_0)_{L^2(\Omega)} \\ &= (f \; ; \; \nabla u_0)_{L^2(\Omega)} + (\phi \; ; \; \gamma u_0)_{L^2(\Gamma_N)} - (\nabla \widehat{u}_D \; ; \; \nabla u_0)_{L^2(\Omega)} \\ &\leq \|u_0\|_{H^1(\Omega)} \left(\sup_{v \in H_D^1(\Omega) \setminus \{0\}} \frac{(f \; ; \; v)_{L^2(\Omega)}}{\|v\|_{H^1(\Omega)}} + \sup_{w \in H^{1/2}(\Gamma_N) \setminus \{0\}} \frac{(\phi \; ; \; w)_{L^2(\Gamma_N)}}{\|w\|_{H^{1/2}(\Gamma_N)}} + \|\widehat{u}_D\|_{H^1(\Omega)} \right) \end{split}$$

Second, the triangle inequality gives

$$||u||_{H^{1}(\Omega)} \leq ||\widehat{u}_{D}||_{H^{1}(\Omega)} + ||u_{0}||_{H^{1}(\Omega)}$$

$$\leq (1 + C_{F}^{2}) \left(\sup_{v \in H_{D}^{1}(\Omega) \setminus \{0\}} \frac{(f ; v)_{L^{2}(\Omega)}}{||v||_{H^{1}(\Omega)}} + \sup_{w \in H^{1/2}(\Gamma_{N}) \setminus \{0\}} \frac{(\phi ; w)_{L^{2}(\Gamma_{N})}}{||w||_{H^{1/2}(\Gamma_{N})}} + ||\widehat{u}_{D}||_{H^{1}(\Omega)} \right).$$

Taking the infimum over all \hat{u}_D , we conclude the stability estimate (5.18).

Remark. A first idea for the numerical approximation of the weak solution $u \in H^1(\Omega)$ of (5.15) might be the following:

- Construct an extension $\widehat{u}_D \in H^1(\Omega)$ of the Dirichlet data.
- Discretize the variational form (5.17) by P1-FEM to obtain an approximation $u_{0h} \in \mathcal{S}_D^1(\mathcal{T})$ of $u_0 \in H_D^1(\Omega)$.
- Compute $u_h := u_{0h} + \widehat{u}_D$ to obtain an approximation of u.

We stress, however, that then $u_h \notin S^1(\mathcal{T})$ so that a postprocessing or evaluation of u_h is nontrivial. Moreover, we have to compute the scalar product $(\nabla \widehat{u}_D ; \nabla v_h)$ for discrete functions to build the load vector of the P1-FEM for u_0 . This leads to additional quadrature errors. Finally and most important, it might be hard to compute \widehat{u}_D unless the Dirichlet data u_D are rather simple.

To overcome the difficulties mentioned in the previous remark, one uses the following approach in practice, which is then called P1-FEM of the weak form (5.16):

- Discretize Dirichlet data $u_D \in H^{1/2}(\Gamma_D)$ by some $u_{Dh} \in \mathcal{S}^1(\mathcal{T}|_{\Gamma_D}) := \{v_h|_{\Gamma_D} \mid v_h \in \mathcal{S}^1(\mathcal{T})\}.$
- Construct extension $\widehat{u}_{Dh} \in \mathcal{S}^1(\mathcal{T})$ with $\widehat{u}_{Dh}|_{\Gamma_D} = u_{Dh}$.
- With \widehat{u}_{Dh} replacing \widehat{u}_D , compute P1-FEM approximation $u_{0h} \in \mathcal{S}_D^1(\mathcal{T})$, cf. (5.17).

• Finally, define $u_h := u_{0h} + \widehat{u}_{Dh} \in \mathcal{S}^1(\mathcal{T})$ as approximation of the weak solution $u \in H^1(\Omega)$.

Note that the discrete solution $u_h \in \mathcal{S}^1(\mathcal{T})$ then belongs to the affine space $\widehat{u}_{Dh} + \mathcal{S}_D^1(\mathcal{T})$. The following result is the corresponding Céa-type lemma:

Lemma 5.5 (Céa lemma, first version). Let $u \in H^1(\Omega)$ be the weak solution of (5.15). Let $\widehat{u}_{Dh} \in \mathcal{S}^1(\mathcal{T})$ be the approximate Dirichlet data and $u_{Dh} := \widehat{u}_{Dh}|_{\Gamma_D}$. Let $u_h \in \mathcal{S}^1(\mathcal{T})$ be the unique solution of

$$u_h|_{\Gamma_D} = u_{Dh}$$

$$(\nabla u_h ; \nabla v_h)_{L^2(\Omega)} = (f ; v_h)_{L^2(\Omega)} + (\phi ; v_h)_{L^2(\Gamma_N)} \quad \text{for all } v_h \in \mathcal{S}_D^1(\mathcal{T}).$$

$$(5.19)$$

Then, u_h is quasioptimal in the sense that there exists a constant C>0 such that

$$C^{-1} \|u - u_h\|_{H^1(\Omega)} \le \min_{v_h \in \mathcal{S}_D^1(\mathcal{T})} \|u - (v_h + \widehat{u}_{Dh})\|_{H^1(\Omega)} = \min_{\substack{w_h \in \mathcal{S}^1(\mathcal{T}) \\ w_h|_{\Gamma_D} = u_D}} \|u - w_h\|_{H^1(\Omega)}.$$
 (5.20)

The constant C > 0 only depends on Ω and Γ_D .

Proof. Note that the variational formulations (5.16) and (5.19) imply the Galerkin orthogonality

$$(\nabla(u-u_h); \nabla v_h)_{L^2(\Omega)} = 0$$
 for all $v_h \in \mathcal{S}_D^1(\mathcal{T})$.

We define $u_{0h} := u_h - \widehat{u}_{Dh} \in \mathcal{S}_D^1(\mathcal{T})$ and observe that

$$\begin{split} \|\nabla(u - u_h)\|_{L^2(\Omega)}^2 &= (\nabla(u - u_h) \; ; \; \nabla(u - [u_{0h} + \widehat{u}_{Dh}]))_{L^2(\Omega)} \\ &= (\nabla(u - u_h) \; ; \; \nabla(u - [v_h + \widehat{u}_{Dh}]))_{L^2(\Omega)} \\ &\leq \|\nabla(u - u_h)\|_{L^2(\Omega)} \|\nabla(u - [v_h + \widehat{u}_{Dh}])\|_{L^2(\Omega)} \end{split}$$

for each $v_h \in \mathcal{S}_D^1(\mathcal{T})$. Next, recall that $|||v||| := ||\nabla v||_{L^2(\Omega)} + ||\gamma v||_{L^2(\Gamma_D)}$ provides an equivalent norm on $H^1(\Omega)$, i.e., there are constants $C_1, C_2 > 0$ such that $C_1^{-1} |||v||| \le ||v||_{H^1(\Omega)} \le C_2 |||v|||$ for all $v \in H^1(\Omega)$. Consequently,

$$C_{2}^{-1}\|u - u_{h}\|_{H^{1}(\Omega)} \leq \|\nabla(u - u_{h})\|_{L^{2}(\Omega)} + \|\gamma(u - u_{h})\|_{L^{2}(\Gamma_{D})}$$

$$= \|\nabla(u - u_{h})\|_{L^{2}(\Omega)} + \|u_{D} - u_{Dh}\|_{L^{2}(\Gamma_{D})}$$

$$\leq \|\nabla(u - [v_{h} + \widehat{u}_{Dh}])\|_{L^{2}(\Omega)} + \|\gamma(u - [v_{h} + \widehat{u}_{Dh}])\|_{L^{2}(\Gamma_{D})}$$

$$\leq C_{1}\|u - (v_{h} + \widehat{u}_{Dh})\|_{H^{1}(\Omega)}$$

for all $v_h \in \mathcal{S}_D^1(\mathcal{T})$. This proves (5.20) with an infimum on the right-hand side. Standard arguments show that this infimum is, in fact, attained.

Exercise 37. Proof that (5.19) has a unique solution
$$u_h \in \mathcal{S}^1(\mathcal{T})$$
.

Remark. Note that Lemma 5.5 is independent of how the Dirichlet data are actually discretized, but the discretization enters the right-hand side, since it constraints the affine space for the minimum in (5.20). Later on, we shall see that appropriate discretization $u_{Dh} = J_h u_D$ by means of the

Scott-Zhang projection J_h even guarantees that

$$||u - u_h||_{H^1(\Omega)} \le C \min_{w_h \in S^1(\mathcal{T})} ||u - w_h||_{H^1(\Omega)},$$

where the right-hand side is independent of how u_D is actually discretized; see also Exercise 39–40 below.

Remark. If the Dirichlet data u_D have an extension $\widehat{u}_D \in H^2(\Omega)$ with $\gamma \widehat{u}_D|_{\Gamma_D} = u_D$, then u_D is continuous. We define $\widehat{u}_{Dh} \in \mathcal{S}^1(\mathcal{T})$ nodewise by

$$\widehat{u}_{Dh}(z) = \begin{cases} u_D(z) & \text{for } z \in \overline{\Gamma}_D, \\ 0 & \text{else,} \end{cases}$$

for $z \in \mathcal{K}$. Let $u \in H^1(\Omega)$ denote the weak solution of (5.15) and $u_0 := u - \widehat{u}_D \in H^1_D(\Omega)$. We additionally define $\widetilde{u}_{Dh} \in \mathcal{S}^1_D(\mathcal{T})$ nodewise by

$$\widetilde{u}_{Dh}(z) = \begin{cases} 0 & \text{for } z \in \overline{\Gamma}_D, \\ \widehat{u}_D(z) & \text{else,} \end{cases}$$

for $z \in \mathcal{K}$. Note that the nodal interpolant of \widehat{u}_D reads $I_h \widehat{u}_D = \widehat{u}_{Dh} + \widetilde{u}_{Dh}$ and that $\|\widehat{u}_D - I_h \widehat{u}_D\|_{H^1(\Omega)} = \mathcal{O}(h)$ decays with optimal order. Consequently, we may plug-in $u = \widehat{u}_D + u_0$ into Céa's lemma to observe that

$$C^{-1} \| u - u_h \|_{H^1(\Omega)} \leq \min_{v_h \in \mathcal{S}_D^1(\mathcal{T})} \| u - (\widehat{u}_{Dh} + v_h) \|_{H^1(\Omega)}$$

$$= \min_{v_h \in \mathcal{S}_D^1(\mathcal{T})} \| (\widehat{u}_D - I_h \widehat{u}_D) + (u_0 - v_h + \widetilde{u}_{Dh}) \|_{H^1(\Omega)}$$

$$= \min_{w_h \in \mathcal{S}_D^1(\mathcal{T})} \| (\widehat{u}_D - I_h \widehat{u}_D) + (u_0 - w_h) \|_{H^1(\Omega)}$$

$$\leq \| \widehat{u}_D - I_h \widehat{u}_D \|_{H^1(\Omega)} + \min_{w_h \in \mathcal{S}_D^1(\mathcal{T})} \| u_0 - w_h \|_{H^1(\Omega)}.$$

Conversely, it holds that

$$\min_{w_h \in \mathcal{S}_D^1(\mathcal{T})} \|u_0 - w_h\|_{H^1(\Omega)} = \min_{w_h \in \mathcal{S}_D^1(\mathcal{T})} \|(u - \widehat{u}_D) - w_h\|_{H^1(\Omega)}
= \min_{w_h \in \mathcal{S}_D^1(\mathcal{T})} \|u - (w_h + I_h \widehat{u}_D) - (\widehat{u}_D - I_h \widehat{u}_D)\|_{H^1(\Omega)}
\leq \min_{w_h \in \mathcal{S}_D^1(\mathcal{T})} \|u - (w_h + I_h \widehat{u}_D)\|_{H^1(\Omega)} + \|\widehat{u}_D - I_h \widehat{u}_D\|_{H^1(\Omega)}
\leq \|u - u_h\|_{H^1(\Omega)} + \|\widehat{u}_D - I_h \widehat{u}_D\|_{H^1(\Omega)}.$$

Therefore, the proposed P1-FEM for the approximation of $u \in H^1(\Omega)$ converges with the same order as the P1-FEM for the approximation of $u_0 \in H^1_D(\Omega)$.

The inhomogeneous Dirichlet problem allows the proof that the trace operator has a right inverse \mathcal{L} . This inverse is called *lifting operator*.

Exercise 38. Let $\gamma \in L(H^1(\Omega); H^{1/2}(\Gamma))$ denote the trace operator. Prove that there exists a lifting operator $\mathcal{L} \in L(H^{1/2}(\Gamma); H^1(\Omega))$ such that $\gamma \mathcal{L}v = v$ for all $v \in H^{1/2}(\Gamma)$. **Hint.** Consider an appropriate Dirichlet-Problem with inhomogeneous Dirichlet data $v \in H^{1/2}(\Gamma)$ and let $u := \mathcal{L}v \in H^1(\Omega)$ denote the unique solution.

The assumptions of the following exercise will be satisfied for the Scott-Zhang projection.

Exercise 39 (Céa lemma, second version). Suppose that there exists a linear projection $P_h: H^1(\Omega) \to \mathcal{S}^1(\mathcal{T})$ with the following properties

- (i) $||P_h v||_{H^1(\Omega)} \le C_{\text{stab}} ||v||_{H^1(\Omega)}$ for all $v \in H^1(\Omega)$
- (ii) $P_h v_h = v_h$ for all $v_h \in \mathcal{S}^1(\mathcal{T})$
- (iii) $(P_h v)|_{\omega} = v|_{\omega}$ for all $v \in H^1(\Omega)$ with $v|_{\omega} \in \mathcal{S}^1(\mathcal{T}|_{\omega})$ and $\omega \in \{\Gamma, \Gamma_D\}$
- (iv) $(P_h v)|_{\omega}$ depends only on the trace $v|_{\omega}$ for all $v \in H^1(\Omega)$ and $\omega \in \{\Gamma, \Gamma_D\}$

Then, for $u \in H^1(\Omega)$ being the solution of (5.15) and $u_{Dh} := (P_h u)|_{\Gamma_D}$, it holds that

$$\min_{\substack{v_h \in \mathcal{S}^1(\mathcal{T}) \\ v_h|_{\Gamma_D} = u_{Dh}}} \|u - v_h\|_{H^1(\Omega)} \le C \min_{\substack{w_h \in \mathcal{S}^1(\mathcal{T})}} \|u - w_h\|_{H^1(\Omega)},$$

where C > 0 depends only on the stability constant C_{stab} , Ω , and Γ_D . In particular, this implies an unconstrained Céa lemma for the mixed boundary value problem with inhomogeneous Dirichlet data, i.e., under the assumptions of Lemma 5.5 and with $u_{Dh} = (P_h u)|_{\Gamma_D}$, it holds

$$||u - u_h||_{H^1(\Omega)} \le C \min_{w_h \in S^1(\mathcal{T})} ||u - w_h||_{H^1(\Omega)}.$$

Hint. Let $w \in H^1(\Omega)$ be the weak solution of $\Delta w = 0$ in Ω subject to the boundary conditions $w = u - u_{Dh}$ on Γ_D and $\partial_n w = 0$ on Γ_N . Define $u_0 := u - w$. Prove that $u_0 = u_{Dh}$ on Γ_D and $\|u - u_0\|_{H^1(\Omega)} \simeq \|u - u_{Dh}\|_{H^{1/2}(\Gamma)}$. Choose $v_h := P_h u_0$.

The existence of a Scott-Zhang-type projection is essentially equivalent to the validity of the Céa lemma.

Exercise 40. (a) Suppose that $P_h: H^1(\Omega) \to \mathcal{S}^1(\mathcal{T})$ satisfies the properties (i)–(iii) of Exercise 39 for $\omega = \Gamma$ only. Then, for all $u \in H^1(\Omega)$ and all $u_{Dh} \in \mathcal{S}^1(\mathcal{T})$, it holds that

$$\min_{\substack{v_h \in \mathcal{S}^1(\mathcal{T}) \\ v_h \mid_{\Gamma} = u_{Dh}}} \|u - v_h\|_{H^1(\Omega)} \le C \left[\min_{\substack{w_h \in \mathcal{S}^1(\mathcal{T}) \\ v_h \mid_{\Gamma} = u_{Dh}}} \|u - w_h\|_{H^1(\Omega)} + \|u - u_{Dh}\|_{H^{1/2}(\Gamma)} \right], \tag{5.21}$$

where C > 0 depends only on Ω and the stability constant C_{stab} . Hint. Argue along the lines of Exercise 39.

(b) Suppose that (5.21) holds true. Then, there exists a linear projection $P_h: H^1(\Omega) \to \mathcal{S}^1(\mathcal{T})$ which satisfies the properties (i)–(iii) of Exercise 39. **Hint.** For given $u \in H^1(\Omega)$, let $u_{Dh} \in \mathcal{S}^1(\mathcal{T}|_{\Gamma})$ be the $H^{1/2}(\Gamma)$ -best approximation of $u|_{\Gamma} \in H^{1/2}(\Gamma)$. With this, let

 $P_h u := u_h \in \mathcal{S}^1(\mathcal{T})$ be the FEM solution of the inhomogeneous Dirichlet problem with discrete Dirichlet data u_{Dh} .

Remark. Note that, for inhomogeneous Dirichlet data, it holds that

$$|||u - u_h||^2 \neq ||u||^2 - ||u_h||^2$$

in general. Therefore, we cannot proceed as in Exercise 36 to approximate the error. Instead, in academic examples, where u is known, one has to compute

$$|||u - u_h||^2 = \sum_{T \in \mathcal{T}} ||\nabla u - \nabla u_h||^2_{L^2(T)}$$

by \mathcal{T} -piecewise numerical quadrature.

5.3 Higher Dimensions

A set $T \subset \mathbb{R}^d$ is called **non-degenerate simplex** provided that there are nodes $z_0, \ldots, z_d \in \mathbb{R}^d$ with $T = \text{conv}\{z_0, \ldots, z_d\}$ and provided that |T| > 0, i.e., T has positive measure. We note that T is in particular bounded and closed, whence compact. For d = 2, this definition describes non-degenerate triangles; for d = 3, this definition describes non-degenerate tetrahedra.

The most important example is the **reference simplex**

$$T_{\text{ref}} := \text{conv}\{0, \mathbf{e}_1, \dots, \mathbf{e}_d\},\tag{5.22}$$

where \mathbf{e}_{j} is the j-th unit vector. There holds $|T_{\text{ref}}| = 1/d!$

The **diameter** of T is denoted by

$$h_T := \operatorname{diam}(T) := \max\{|x - y| \mid x, y \in T\}.$$
 (5.23)

Moreover, ρ_T denotes the radius of the largest ball inscribed of T, i.e.,

$$\rho_T := \sup \left\{ \rho > 0 \mid \exists x \in T \quad B(x, \rho) \subseteq T \right\}. \tag{5.24}$$

By $\mathcal{K}_T = \{z_0, \dots, z_d\}$, we denote the set of nodes of T. By \mathcal{E}_T , we denote the set of faces of T, i.e., $\mathcal{E}_T := \{\operatorname{conv}(M) \mid M \subseteq \mathcal{K}_T \text{ with } \# M = d\}$. Note that $E \in \mathcal{E}_T$ is a hyper-simplex of dimension d-1, e.g., the faces of a tetrahedron are 2-dimensional surface triangles.

Definition. Let Ω be a bounded Lipschitz domain in \mathbb{R}^d , $d \geq 2$. A set \mathcal{T} is a **triangulation** of Ω (consisting of simplices) if and only if

- T is a finite set of non-degenerate simplices,
- the closure of Ω is covered by \mathcal{T} , i.e., $\overline{\Omega} = \bigcup \mathcal{T}$,
- for all $T, T' \in \mathcal{T}$ with $T \neq T'$ holds $|T \cap T'| = 0$, i.e., the overlap is a set of measure zero.

By $\mathcal{K} := \bigcup \{x \in \mathcal{K}_T \mid T \in \mathcal{T}\}$, we then denote the **set of nodes** of the triangulation \mathcal{T} and by $\mathcal{E} := \bigcup \{E \in \mathcal{E}_T \mid T \in \mathcal{T}\}$ the **set of faces** of the triangulation \mathcal{T} . A triangulation of Ω is called **conforming** or **regular** (in the sense of Ciarlet) provided that the intersection of two elements $T, T' \in \mathcal{T}$ with $T \neq T'$ is

- either empty,
- or a joint k-dimensional hyper-simplex of both T and T', i.e., $T \cap T' = \text{conv}(M)$ with $M \subseteq \mathcal{K}_T \cap \mathcal{K}_{T'}$ and $\#M = k \leq d-1$.

According to this regularity assumption, a face $E \in \mathcal{E}$ with surface measure $|E \cap \Gamma| > 0$ automatically satisfies $E \subseteq \Gamma$, i.e., a face E is either a boundary face or an interior face. Additionally, we always assume that a regular triangulation resolves the boundary conditions: If $\Gamma = \partial \Omega$ is partitioned into Dirichlet and Neumann boundary Γ_D and Γ_N , respectively, each boundary face $E \in \mathcal{E}$ with $E \subseteq \Gamma$ satisfies

- either $E \subseteq \overline{\Gamma}_D$
- or $E \subseteq \overline{\Gamma}_N$.

With this assumption, we define the (disjoint) sets of boundary faces

$$\mathcal{E}_D := \left\{ E \in \mathcal{E} \mid E \subseteq \overline{\Gamma}_D \right\} \quad \text{and} \quad \mathcal{E}_N := \left\{ E \in \mathcal{E} \mid E \subseteq \overline{\Gamma}_N \right\}$$
 (5.25)

as well as the set of all interior faces

$$\mathcal{E}_{\Omega} := \mathcal{E} \setminus (\mathcal{E}_D \cup \mathcal{E}_N). \tag{5.26}$$

We finally note that, for each $E \in \mathcal{E}_{\Omega}$, there are two elements $T, T' \in \mathcal{T}$ with $E = T \cap T'$.

For a regular triangulation \mathcal{T} , the hat functions provide a basis of $\mathcal{S}^1(\mathcal{T})$, and all results of Section 3.1 hold accordingly.

5.3.1 Shape Regularity & Scaling Arguments

A regular triangulation \mathcal{T} is γ -shape regular if

$$\sigma(\mathcal{T}) := \max_{T \in \mathcal{T}} \frac{h_T}{\rho_T} \le \gamma < \infty. \tag{5.27}$$

According to Exercise 14, this new definition is (up to some generic constant) equivalent to the definition given in Section 3.2.

For a non-degenerate simplex $T = \text{conv}\{z_0, \dots, z_d\} \subset \mathbb{R}^d$, we define

$$\Phi_T : T_{\text{ref}} \to T, \quad \Phi_T v := z_0 + B_T v, \quad \text{where } B_T := (z_1 - z_0 \quad z_2 - z_0 \quad \dots \quad z_d - z_0) \in \mathbb{R}^{d \times d}.$$

Arguing as in Lemma 3.9, we see that $||B_T||_F \simeq h_T$, since the diameter of a simplex is its longest edge. To employ scaling arguments, it remains to prove $||B_T^{-1}||_F \lesssim \rho_T$. This is done with the help of the following lemma.

Lemma 5.6. Let $T_1, T_2 \subset \mathbb{R}^d$ be compact sets with $B(x_j, \rho_j) \subseteq T_j \subseteq B(y_j, r_j)$ for some $x_j, y_j \in T_j$ and $\rho_j, r_j > 0$. Let $\Phi : T_1 \to T_2$ be affine with $\Phi(v) := Bv + w$ and $B \in \mathbb{R}^{d \times d}$. Then, it holds $\|B\|_2 \leq r_2/\rho_1$ for the Euclidean operator norm.

Proof. 1. step. For $x \in \mathbb{R}^d$ with $|x| \leq 2\rho_1$, it holds $|Bx| \leq 2r_2$: Since $B(x_1, \rho_1) \subseteq T_1$, we find $y, z \in T_1$ with x = y - z. Then, $\Phi(y), \Phi(z) \in T_2$. Since $T_2 \subseteq B(y_2, r_2)$, it follows $2r_2 \geq |\Phi(y) - \Phi(z)| = |B(y - z)| = |Bx|$.

2. step. For $x \in \mathbb{R}^d$, it holds $|Bx| \le (r_2/\rho_1)|x|$: Let $x \in \mathbb{R}^d \setminus \{0\}$. Define $v := (2\rho_1/|x|)x$. From $|v| = 2\rho_1$, we obtain $(2\rho_1/|x|)|Bx| = |Bv| \le 2r_2$. This concludes the proof.

Corollary 5.7. With the above notation, the matrix $B_T \in \mathbb{R}^{d \times d}$ is invertible with $|\det B_T| \simeq |T|$ and $||B_T^{-1}||_F \lesssim \rho_T^{-1}$.

Proof. As for 2D, one obtains $|\det B_T| \simeq |T| > 0$, and hence Φ_T are invertible. Note that B_T^{-1} is the linear part of the affine mapping Φ_T^{-1} . Hence, Lemma 5.6 gives $||B_T^{-1}||_2 \le h_{\text{ref}}/\rho_T \le \rho_T^{-1}$. Norm equivalence on $\mathbb{R}^{d \times d}$ concludes $||B_T^{-1}||_F \simeq ||B_T^{-1}||_2 \le \rho_T^{-1}$.

The analysis of the previous chapters transfers from d=2 to general dimension $d\geq 2$.

- The whole Chapter 2 is stated for $\Omega \subset \mathbb{R}^d$, $d \geq 2$.
- All results of Section 3.1 now hold verbatim for $d \geq 2$.
- The Approximation Theorem 3.5 holds for d = 2, 3. For $d \ge 4$, one requires higher smoothness of u to ensure continuity (cf. the Sobolev Theorem 3.4).
- Bramble-Hilbert Lemma 3.7 and transformation formula (Lemma 3.8) have already been formulated for $d \ge 2$.
- The inverse estimate and its applications hold verbatim.
- The data approximation analysis of Section 5.1 in the frame of the first Strang lemma applies for d = 2, 3. For $d \ge 4$, it requires only higher regularity assumptions on f to ensure continuity.
- Technical auxiliary results like the trace inequality remain valid for general $d \geq 2$.
- The a posteriori analysis of Chapter 4 remains valid. Only the proof of Lemma 4.1 which provides the dual basis functions to define the Scott-Zhang projection, has to be adapted.
- Finally, the adaptive convergence analysis requires $h_T := |T|^{1/d} \simeq \operatorname{diam}(T)$. All arguments remain valid.

5.4 Higher-order FEM

Due to the Céa lemma, we observe that it can be advantageous to consider higher-order polynomial discrete spaces for FEM. For example if the exact solution is smooth, higher-order spaces will achieve a better rate of convergence of the FEM error.

5.4.1 Higher-order elements in 1D

We consider a "triangulation" \mathcal{T} with nodes x_i , i = 0, ..., M on an interval $\Omega \subset \mathbb{R}$. Instead of piecewise linear functions, we may also use higher-order polynomials to construct our discrete spaces, i.e.

$$S^{p}(\mathcal{T}) := \{ u \in H^{1}(\Omega) \mid u \mid_{T} \circ \Phi_{T} \in \mathcal{P}^{p}(T_{ref}) \quad \forall T \in \mathcal{T} \},$$
 (5.28a)

$$S_0^p(\mathcal{T}) := S^p(\mathcal{T}) \cap H_0^1(\Omega) \tag{5.28b}$$

Here, the we use the mappings from the reference element $T_{\text{ref}} = [-1, 1], \, \Phi_T : T_{\text{ref}} \to T$ as defined above.

Remark. Since Φ_T is affine, $u|_T \circ \Phi_T$ is a polynomial of degree p if and only if $u|_T$ is a polynomial of degree p. This means, the definition above is equivalent to $\mathcal{S}^p(\mathcal{T}) = \{u \in H^1(\Omega) \mid u|_T \in \mathcal{P}_p \forall T \in \mathcal{T}\}$. However, for non-affine maps Φ_T (e.g., for curved elements) the Definition (5.28) is still valid, while the second definition above does not generalize. (Remember that the scaling arguments and inverse estimates in the previous chapters required $u|_T \circ \Phi_T$ to be polynomial.)

We construct a basis of $S^p(\mathcal{T})$ on the reference element. We choose a basis $\{N_i | i = 1, \dots, p+1\}$ of the polynomial space $\mathcal{P}^p(T_{\text{ref}})$ such that

$$N_1(\xi) = \frac{1}{2}(1-\xi), \quad N_2(\xi) = \frac{1}{2}(1+\xi), \quad N_i(\pm 1) = 0 \quad i \ge 3.$$

Remark. The functions N_i , $i \geq 3$ can be chosen quite freely. The simplest possibility is $N_i(\xi) = (1 - \xi^2)\xi^{i-3}$ for all $i \in \{3, \dots, p+1\}$. For small p = 2, 3, 4, this choice is fine. However, for higher p, the choice leads to very badly conditioned stiffness matrices and hence to numerical instabilities. It is better to choose more "orthogonal" basis functions as for example:

$$N_i(\xi) = \int_{-1}^{\xi} L_{i-2}(t) dt, \qquad (5.29)$$

where $L_i \in \mathcal{P}_i$ is the *i*-th Legendre polynomial. Due to the orthogonality properties of Legendre polynomials, we have $N_i(\pm 1) = 0$ for $i \geq 3$. For the practical implementation, it is important to be able to quickly evaluate the basis functions. On one hand, there holds $(2i + 1) \int_{-1}^{\xi} L_i(t) dt = L_{i+1}(\xi) - L_{i-1}(\xi)$ and on the other hand, the Legendre polynomials can be computed very efficiently via three-term recurrences.

Since the basis functions vanish for $i \geq 3$, it is easy to construct a basis of $\S S^p(\mathcal{T})$ from these local definitions, i.e.,

$$\mathcal{B} = \mathcal{B}^{lin} \cup \left(\cup_{T \in \mathcal{T}} \mathcal{B}^T \right), \tag{5.30}$$

where $\mathcal{B}^{lin} = \{\varphi_i | i = 0, ..., M\}$ are the hat-functions corresponding to $x_i, i = 0, ..., M$ and $\mathcal{B}^T = \{\varphi_{T,i} | i = 3, ..., p+1\}$ with

$$\varphi_{T,i}(x) = \begin{cases} N_i(\Phi_T^{-1}(x)) & x \in T \\ 0 & x \in \Omega \setminus T \end{cases}$$

We note that the construction of the basis above followed a typical recipe in FEM: The local basis function (the form functions) are associated with geometrical objects, e.g., the hat-functions are associated with nodes, whereas the bubble functions $\varphi_{T,i}$ are associated with elements $T \in \mathcal{T}$. Moreover, we observe that $\varphi|_T \circ \Phi_T \in \{0, N_1, \dots, N_{p+1}\}$, i.e., a basis function vanishes on an element, or it is exactly one of the local basis functions N_i .

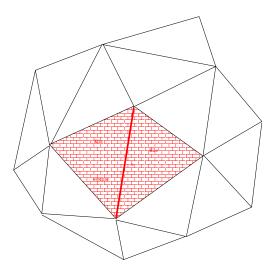


FIGURE 5.1. The edge E with elements T_E^1 , T_E^2 and $\Omega_E = T_E^1 \cup T_E^2$.

5.4.2 Higher-order elements in 2D

Analogously to the 1D case, we may define higher-order basis functions in 2D. Let \mathcal{T} denote a regular triangulation and define

$$\mathcal{S}^p(\mathcal{T}) := \{ u \in H^1(\Omega) \mid u_{|K} \circ \Phi_T \in \mathcal{P}^p(T_{ref}) \},$$

and
$$\mathcal{S}_0^p(\mathcal{T}) := \mathcal{S}^p(\mathcal{T}) \cap H_0^1(\Omega)$$
.

When constructing the basis functions for the FEM-spaces, we implicitly obeyed the following rules:

- 1. The basis functions $\varphi \in \mathcal{B}$ have a simple structure on the reference element T_{ref} , i.e., for all $T \in \mathcal{T}$ the function satisfies $\varphi|_T \circ \Phi_T \in \{0, N_1, N_2, \dots, \}$, where $\{0, N_1, \dots, \}$ is known explicitly.
- 2. The support supp φ of the basis functions $\varphi \in \mathcal{B}$ is small. This leads to sparse stiffness matrices and hence more efficient solvers.
- 3. For implementation, it is often advantageous to associate certain basis functions with geometrical objects, e.g., nodes, edges, elements, ...

The case p=2

Idea: Construct \mathcal{B} as a union of hat functions \mathcal{B}^{lin} and edges-bubble functions. The latter functions φ_E are supported on Ω_E (see Figure 5.1). On the edge E, φ_E is a quadratic function as shown in Figure 5.2.

In engineering literature, the basis functions N_i are often illustrated with a diagram in which each dot represents a form-function:

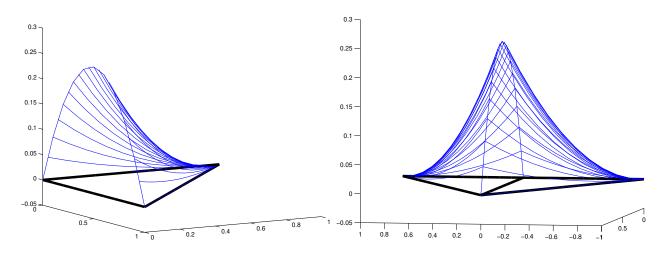
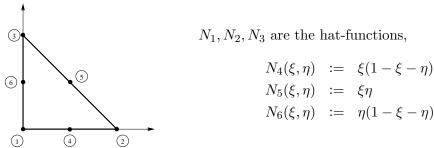


FIGURE 5.2. Left: N_4 on T_{ref} . Right: φ_E .



We note that the edge bubbles N_4, \ldots, N_6 are chosen such that they vanish on two edges of T_{ref} . Hence, we may associate each of those functions with one edge where it is non-zero. We write the basis \mathcal{B} of $\mathcal{S}^2(\mathcal{T})$ as

$$\mathcal{B} = \mathcal{B}^{lin} \cup \left(\cup_{e \in \mathcal{E}} \mathcal{B}^E \right),$$

where \mathcal{B}^{lin} is again the set of hat-functions associated with the nodes \mathcal{N} . The one-element sets $\mathcal{B}^E = \{\varphi_E\}, E \in \mathcal{E}$ contain the edge bubble functions, which are characterized as follows:

$$\varphi_E \in H^1(\Omega), \quad \sup \varphi_E \subset \overline{\Omega_E}, \quad \varphi_E|_T \circ \Phi_T \in \{N_4, N_5, N_6\} \quad \forall T \in \Omega_E.$$
 (5.31)

Remark. If we restrict them to an edge of $T_{\rm ref}$, the functions N_i $(i \in \{4,5,6\})$ are symmetric with respect to the midpoint of the edge. Hence, the above definition of φ_E leads to a continuous basis function. To see this, let $E \in \mathcal{E}$ with two elements T_E^1 , $T_E^2 \in \Omega_E$. Let $\Gamma_4 = \{(x,0) \mid x \in (0,1)\}$, $\Gamma_5 = \{(x,y) \mid x \in (0,1) \mid 1-x-y=0\}$, $\Gamma_6 = \{(0,y) \mid y \in (0,1)\}$ denote the three edges of the reference element $T_{\rm ref}$. Let $i, j \in \{4,5,6\}$ denote the edge numbers corresponding to E, i.e., $\Phi_{T_E^1}(\Gamma_i) = E$ and $\Phi_{T_E^2}(\Gamma_j) = E$. Then, the above definition of φ_E is equivalent to

$$\varphi_E(x) := \begin{cases} N_i \circ \Phi_{T_E^1}^{-1}(x) & x \in \overline{T_E^1} \\ N_j \circ \Phi_{T_E^2}^{-1}(x) & x \in \overline{T_E^2} \\ 0 & else \end{cases}$$

The symmetry of N_i on the edges shows that this is well-defined since the two cases coincide on the edge E.

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Appendix A

Some Facts from Functional Analysis

I this appendix we collect some results from introductory functional analysis courses which are used throughout. We stick with the case of vector spaces over \mathbb{R} .

A.1 Main Theorems from Functional Analysis

Theorem A.1 (Hahn-Banach Extension Theorem). Let $p: X \to \mathbb{R}$ be a sublinear functional on a linear space X, i.e. $p(x+y) \le p(x) + p(y)$ and $p(\lambda x) = \lambda p(x)$ for all $x, y \in X$ and $\lambda \ge 0$. If Y is a subspace of X and $f: Y \to \mathbb{R}$ is a linear functional with $f \le p$ on Y, there is a linear extension $F: X \to \mathbb{R}$ with $F|_Y = f$ and $F \le p$ on X.

If X is a normed space and $f \in Y^*$, one may choose $p(x) = ||x||_X ||f||_{X^*}$ to prove the extension theorem for continuous linear functionals.

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Corollary A.2. If Y is the subspace of a normed space X and f \in Y^*, there is an extension F \in X^* with F|_Y = f and ||F||_{X^*} = ||f||_{Y^*}.
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One then considers the subspace $Y := \operatorname{span}\{x\}$ and $f(\lambda x) = \lambda ||x||_X$ to derive the following corollary:

Corollary A.3. If X is a normed space and
$$x \in X$$
, there is a linear functional $f \in X^*$ with $||f||_{X^*} = 1$ and $f(x) = ||x||_X = \sup_{\|f\|_{X^*} = 1} |f(x)|$.

Theorem A.4 (Hahn-Banach Separation Theorem). Let X be a normed space, and let A and B be convex, nonempty subsets of X with $A \cap B = \emptyset$.

- (i) If A is open, there is a linear functional $f \in X^*$ and a scalar $\lambda \in \mathbb{R}$ such that $f(x) < \lambda \le f(y)$ for all $x \in A$ and $y \in B$.
- (ii) If A is compact and B is closed, there is a linear functional $f \in X^*$ and scalars $\lambda_1, \lambda_2 \in \mathbb{R}$ such that $f(x) \leq \lambda_1 < \lambda_2 \leq f(y)$ for all $x \in A$ and $y \in B$.

If Y is a subspace of X, one can use (ii) to characterize the closure \overline{Y} of Y in X. The proof only needs that each bounded linear functional $f \in Y^*$ is trivial, i.e. $f|_Y = 0$.

Corollary A.5. Let Y be a subspace of the normed space X. Then, $x \in X$ satisfies $x \in \overline{Y}$ if and only if f(x) = 0 for all $f \in X^*$ with $f|_Y = 0$.

Proof. For $x \in \overline{Y}$ and $f \in X^*$ with $f|_Y = 0$, continuity yields f(x) = 0. The converse implication is proven by contradiction: We assume that $x \notin \overline{Y}$ and choose $f \in X^*$ such that $f(x) < \lambda \le f(y)$ for all $y \in Y$ and some fixed $\lambda \in \mathbb{R}$. Using that Y is a vector space, we infer that $\lambda \le f(\pm y) = -f(\mp y) \le -\lambda$ and thus $f(y) \in [\lambda, -\lambda]$ for all $y \in Y$. As bounded linear functionals are trivial, we obtain $f|_Y = 0$. According to our assumptions, this implies f(x) = 0 and thus contradicts $f(x) < \lambda \le f(0) = 0$.

The following corollary is an immediate consequence of the last one.

Corollary A.6. Let Y be a subspace of the normed space X. Then, Y is dense in X if and only if each functional $f \in X^*$ with $f|_Y = 0$ is trivial, i.e., $f = 0 \in X^*$.

For an operator $T \in L(X;Y)$, one defines $(T^*y^*)(x) := y^*(Tx)$ for all $y^* \in Y^*$ and $x \in X$. From the continuity of T, we see that $T^*y^* \in X^*$, and obviously $T^*:Y^* \to X^*$ is a linear operator. From the corollary of the Hahn-Banach extension theorem, we derive for the operator norm

$$||T^*|| = \sup_{\|y^*\|_{Y^*}=1} ||T^*y^*||_{X^*} = \sup_{\|y^*\|_{Y^*}=1} \sup_{\|x\|_X=1} (T^*y^*)(x)$$

$$= \sup_{\|x\|_X=1} \sup_{\|y^*\|_{Y^*}=1} (y^*)(Tx) = \sup_{\|x\|_X=1} ||Tx||_Y = ||T||,$$

i.e. there holds $T^* \in L(Y^*; X^*)$ with operator norm $||T^*|| = ||T||$. The operator T^* is called the **adjoint operator** of T.

Theorem A.7 (Banach Closed Range Theorem). For an operator $T \in L(X;Y)$ between Banach spaces X and Y and $T^* \in L(Y^*;X^*)$ its adjoint, the following is pairwise equivalent:

- (i) range(T) is a closed subspace of Y.
- (ii) range $(T) = (\ker T^*)_{\circ} := \{ y \in Y \mid \forall y^* \in \ker(T^*) \mid y^*(y) = 0 \}.$
- (iii) range(T^*) is a closed subspace of X^* .
- (iv) range $(T^*) = (\ker T)^\circ := \{x^* \in X^* \mid \forall x \in \ker(T) \ x^*(x) = 0\}.$

A.2 Hilbert Spaces

A space X is called **Hilbert space** if it is a Banach space whose norm is induced by a scalar product.

Theorem A.8. Let Y be the closed subspace of a Hilbert space X and $Y^{\perp} := \{x \in X \mid \forall y \in Y \mid (x;y)_X = 0\}$ the orthogonal complement. Then, there holds $X = Y \oplus Y^{\perp}$ in the sense of the linear algebra, i.e. every element $x \in X$ has a unique decomposition $x = y + y^{\perp}$ with some $y \in Y$ and $y^{\perp} \in Y^{\perp}$.

With the orthogonal decomposition $X = Y \oplus Y^{\perp}$, one can define a projection $\pi_Y : X \to Y$ by $x = y + y^{\perp} \mapsto y$.

Corollary A.9. Let Y be the closed subspace of a Hilbert space X. Then, there is a unique linear operator $\Pi: X \to Y$ with $\Pi|_Y = id$ and $\ker(\Pi) = Y^{\perp}$, which is called **orthogonal projection** onto Y. This projection is continuous with operator norm $\|\Pi\| = 1$ and symmetric, i.e. $(x; y)_X = (\Pi x; y)_X$ for all $x \in X$ and $y \in Y$. Moreover, the orthogonal projection is the solution operator for the best approximation problem, $\|x - \Pi x\|_X = \min_{y \in Y} \|x - y\|_X$.

The dual space X^* of a Hilbert space X has a straight-forward representation, and one can somehow identify X with X^* .

Theorem A.10 (Riesz). For a Hilbert space X, the Riesz mapping $I_X: X \to X^*$, $I_X x := (x; \cdot)_X \in X^*$, is an isometric isomorphism.