

PROBLEMA 1

$$V(r) = \alpha \delta(r-a)$$

$$f(\theta) = \frac{1}{k} \sum_l (2l+1) e^{i\delta_l} \sin \delta_l P_l(\cos \theta)$$

Como solo sobrevive $l=0$ ya que $ka \ll 1 \Rightarrow f(\theta) \cong \frac{1}{k} e^{i\delta_0} \sin \delta_0$

Resolvemos la parte radial de la ecuación de Schrödinger por $l=0$:

$$-\frac{\hbar^2}{2m} \frac{d^2 u}{dr^2} + \alpha \delta(r-a) u = E u$$

Para determinar condiciones de frontera integramos esta ecuación alrededor de a :

$$-\frac{\hbar^2}{2m} \int_{a-\epsilon}^{a+\epsilon} \left(\frac{d^2 u}{dr^2} \right) r^2 dr + \alpha \int_{a-\epsilon}^{a+\epsilon} r^2 \delta(r-a) u dr = E \int_{a-\epsilon}^{a+\epsilon} r^2 u dr \quad E \rightarrow 0$$

Por partes

$$-\frac{\hbar^2}{2m} \left(r^2 \frac{du}{dr} \right) \Big|_{a-\epsilon}^{a+\epsilon} - 2 \int_{a-\epsilon}^{a+\epsilon} r \frac{du}{dr} dr + \alpha a^2 u(a) = 0$$

$$-\frac{\hbar^2}{2m} \left(a^2 (u'(a^+) - u'(a^-)) - 2 \underbrace{U r \Big|_{a-\epsilon}^{a+\epsilon}}_0 + 2 \int_{a-\epsilon}^{a+\epsilon} V dr \right) + \alpha a^2 u(a) = 0$$

Por continuidad de $U(r)$

$$\Rightarrow -\frac{\hbar^2 a^2}{2m} \left(\frac{du}{dr} \Big|_{r=a^+} - \frac{du}{dr} \Big|_{r=a^-} \right) = -\alpha a^2 u(a)$$

$$\Rightarrow \frac{du}{dr} \Big|_{r=a^+} - \frac{du}{dr} \Big|_{r=a^-} = \frac{2m\alpha}{\hbar^2} u(a)$$

$$\text{Sea } \phi = \frac{2m\alpha a}{\hbar^2} \Rightarrow \boxed{\frac{du}{dr} \Big|_{r=a^+} - \frac{du}{dr} \Big|_{r=a^-} = \phi \frac{u(a)}{a}} \quad (*)$$

Para $r \neq a$ $\frac{d^2 U}{dr^2} + K^2 U = 0$ como $K^2 = \frac{2mE}{\hbar^2}$

$$\Rightarrow \begin{cases} U(r) = A \sin(Kr) + A_2 \cos(Kr) & r < a \\ U(r) = B_1 \sin(Kr) + B_2 \cos(Kr) & r > a \end{cases}$$

$$R(r) = \frac{U(r)}{r}$$

$R(0)$ debe ser finita \Rightarrow requiere $A_2 = 0$

Continuidad de $U(r)$ en $r = a$:

$$\textcircled{1} A_1 \sin(Ka) = B_1 \sin(Ka) + B_2 \cos(Ka)$$

Discontinuidad de la derivada ec. $\textcircled{2}$

$$\textcircled{2} -K A_1 \cos(Ka) + K B_1 \cos(Ka) - K B_2 \sin(Ka) = 0 \quad A_1 \frac{\sin(Ka)}{a}$$

Hallamos B_1 y B_2 en términos de A_1 :

$$\text{De } \textcircled{1} : B_1 = A_1 - B_2 \cotan(Ka)$$

$$\text{en } \textcircled{2} : -K A_1 \cos(Ka) + K A_1 \cos(Ka) - K B_2 \cos(Ka) \cotan(Ka) - K B_2 \sin(Ka)$$

$$= 0 \quad A_1 \frac{\sin(Ka)}{a}$$

$$\Rightarrow -K B_2 \left[\sin(Ka) + \frac{\cos^2(Ka)}{\sin(Ka)} \right] = A_1 \left[\frac{\sin(Ka)}{a} \right]$$

$$\Rightarrow -K B_2 = A_1 \frac{\sin^2(Ka)}{a} \Rightarrow \boxed{B_2 = -A_1 \frac{\sin^2(Ka)}{Ka}}$$

$$\Rightarrow B_1 = A_1 \left[1 + \frac{\sin(Ka) \cos(Ka)}{Ka} \right]$$

\Rightarrow Para $r > a$:

$$U(r) = A_1 \left[1 + \frac{\sin(Ka) \cos(Ka)}{Ka} \right] \sin(Kr) - \frac{A_1 \sin^2(Ka) \cos(Kr)}{Ka}$$

y $R(r) = U(r)/r$ se puede escribir como:

$$R(r) = \frac{C}{r} \sin(Kr + \phi_0)$$

$$C = A_1 \left\{ \left(1 + \frac{\sin(Ka) \cos(Ka)}{Ka} \right)^2 + \left(\frac{\sin^2(Ka)}{Ka} \right)^2 \right\}^{1/2}$$

$$-\phi_0 = \tan^{-1} \left\{ \frac{\frac{\sin^2(Ka)}{Ka}}{1 + \frac{\sin(Ka) \cos(Ka)}{Ka}} \right\} = \tan^{-1} \left\{ \frac{1}{\frac{Ka}{\sin^2(Ka)} + \cotan(Ka)} \right\}$$

$$Ka \ll 1 \Rightarrow \sin^2(Ka) \approx K^2 a^2$$

$$\cotan(Ka) \approx \frac{1}{Ka}$$

$$\Rightarrow -\rho_0 = \tan^{-1} \left\{ \frac{1}{\frac{Ka}{\phi K^2 a^2} + \frac{1}{Ka}} \right\} \approx \tan^{-1} \left\{ \frac{\phi Ka}{1+\phi} \right\}$$

$$\sin(\tan^{-1} u) \Rightarrow \tan u = \frac{\sin u}{\cos u} = \frac{\sin u}{\sqrt{1-\sin^2 u}} \Rightarrow \sin^2 u = \tan^2 u (1-\sin^2 u)$$

$$\Rightarrow \sin^2 u = \frac{\tan^2 u}{1+\tan^2 u}$$

$$\sin^2(\tan^{-1} u) = \frac{u^2}{1+u^2} \cos(\tan^{-1} u) = \sqrt{1 - \frac{u^2}{1+u^2}} = \frac{1}{\sqrt{1+u^2}}$$

$$f(\theta) \approx \frac{1}{K} [\cos \rho_0 \sin \rho_0 + i \sin^2 \rho_0]$$

$$\approx \frac{1}{K} \left[\frac{i \phi^2 K^2 a^2 / (1+\phi)^2}{1 + \frac{\phi^2 K^2 a^2}{(1+\phi)^2}} - \frac{\phi Ka / (1+\phi)}{1 + \frac{\phi^2 K^2 a^2}{(1+\phi)^2}} \right]$$

$$f(\theta) \approx \frac{1}{K} \frac{\phi Ka}{(1+\phi)^2 + \phi^2 K^2 a^2} [-(1+\phi) + i \phi Ka]$$

$$\sigma_T = \frac{4\pi}{K^2} \sin^2 \rho_0 = \frac{4\pi}{K^2} \left\{ \frac{\phi^2 K^2 a^2 / (1+\phi)^2}{1 + \frac{\phi^2 K^2 a^2}{(1+\phi)^2}} \right\}$$

$$\sigma_T = 4\pi a^2 \left\{ \frac{\phi^2}{(1+\phi)^2 + \phi^2 K^2 a^2} \right\} \text{ para } \phi^2 K^2 a^2 \ll \phi^2$$

$$\Rightarrow \sigma_T \approx \frac{4\pi a^2 \phi^2}{(1+\phi^2)}$$

PROBLEMA 2

$$\boxed{2} \quad f_0 = -\frac{\alpha K [1 - i\alpha K]}{K [1 + \alpha^2 K^2]}$$

para $f(\theta) = \frac{1}{K} \sum_l (2l+1) e^{i\delta_l} \sin \delta_l P_l(\cos \theta)$

$$\Rightarrow f_0 = \frac{1}{K} e^{i\delta_0} \sin \delta_0 = -\frac{\alpha [1 - i\alpha K]}{1 + \alpha^2 K^2}$$

igualando partes Re e Im :

$$\left\{ \begin{array}{l} \textcircled{1} \quad \frac{1}{K} \cos \delta_0 \sin \delta_0 = -\frac{\alpha}{1 + \alpha^2 K^2} \\ \textcircled{2} \quad \frac{1}{K} \sin^2 \delta_0 = \frac{\alpha^2 K}{1 + \alpha^2 K^2} \end{array} \right\} \begin{array}{l} \textcircled{2}/\textcircled{1} : \\ \Rightarrow \tan \delta_0 = -\alpha K \\ \Rightarrow \boxed{\delta_0 = \tan^{-1} [-\alpha K]} \end{array}$$

PROBLEMA 3

$$\boxed{3} \quad S_0 = k(a + bk^2)$$

$$S_1 = ck^3$$

$$f_i = 0 \quad i > 1$$

$$\boxed{a} \quad \text{Dado que } f(\theta) = \sum_l (2l+1) \frac{e^{i\delta_l} \sin \delta_l}{k} P_l(\cos \theta)$$

Para verificar que $\delta_1 = 0$ para $k > 1$ medimos

$\frac{d\sigma}{d\Omega}$ en función de θ y hacemos un ajuste

a $\frac{d\sigma}{d\Omega}$ con un polinomio en $\cos \theta$ hasta orden 2,

esto es porque si solo hay contribución hasta $l=1$:

$$f(\theta) = \left(\frac{e^{i\delta_0} \sin \delta_0}{k} + e^{i\delta_1} \frac{\sin \delta_1}{k} \cos \theta \right)$$

$$\frac{d\sigma}{d\Omega} = |f(\theta)|^2 \quad \text{solo puede tener términos:}$$

$$\text{constante} + \text{constante} \cos \theta + \text{constante} \cos^2 \theta$$

estas mediciones se hacen a un valor específico de k .

Para medir las constantes se hace un fit a

$\frac{d\sigma}{d\Omega}$ vs k , manteniendo θ constante.

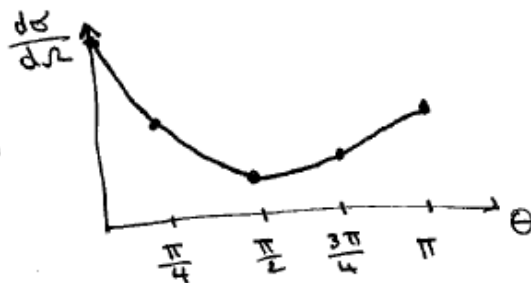
$$\boxed{b} \quad a=b=c=\pi/4, \quad k=1$$

$$\Rightarrow \frac{d\sigma}{d\Omega} = |f(\theta)|^2 = \frac{1}{k^2} \left| (\cos \delta_0 \sin \delta_0 + 3 \cos \delta_1 \sin \delta_1 \cos \theta)^2 + (\sin^2 \delta_0 + 3 \sin^2 \delta_1 \cos^2 \theta) \right|$$

$$= \frac{1}{k^2} (\sin^2 \delta_0 + 9 \cos^2 \theta \sin^2 \delta_1 + 6 \cos \theta \sin \delta_0 \sin \delta_1 \cos(\delta_1 - \delta_0))$$

$$\delta_0 = \pi/2, \quad \delta_1 = \pi/4$$

$$\frac{d\sigma}{d\Omega} = 1 + 3 \cos \theta + \frac{9}{2} \cos^2 \theta$$



PROBLEMA 4

$$\boxed{4} \quad V(r) = \begin{cases} -V_0, & r \leq R \\ 0, & r > R \end{cases}$$

De acuerdo a la primera aproximación de Born:

$$f(\theta) = -\frac{2m}{\hbar^2 q} \int_0^\infty r' V(r') \operatorname{Sen}(qr') dr'$$

$$= \frac{2m V_0}{\hbar^2 q} \int_0^R r' \operatorname{Sen}(qr') dr'$$

$$= \frac{2m V_0}{\hbar^2 q^3} \int_0^{qR} u \operatorname{Sen}(u) du$$

$$= \frac{2m V_0}{\hbar^2 q^3} \left[\operatorname{Sen}(u) - u \cos(u) \right] \Big|_0^{qR}$$

$$f(\theta) = \frac{2m V_0}{\hbar^2 q^3} \left[\operatorname{Sen}(qR) - qR \cos(qR) \right]$$

PROBLEMA 5

Usando la primera aproximación de Born .

$$a) V(r) = V_0 \exp(-\alpha r)$$

$$f(\theta) = \frac{-2m}{\hbar^2 K} \int_0^\infty r V(r) \sin(Kr) dr \quad \text{porque el potencial es esfericamente simétrico}$$

$$\Rightarrow f(\theta) = -\frac{m V_0}{\hbar^2 K \sin(\frac{\theta}{2})} \int_0^\infty r e^{-\alpha r} \sin(kr) dr$$

$$= -\frac{m V_0}{\hbar^2 K \sin(\frac{\theta}{2})} \frac{1}{2i} \left\{ \int_0^\infty r e^{-(\alpha - ik)r} dr - \int_0^\infty r e^{-(\alpha + ik)r} dr \right\}$$

$$u_1 = (\alpha - ik)r \Rightarrow du_1 = (\alpha - ik) dr$$

$$= -\frac{m V_0}{\hbar^2 K \sin(\frac{\theta}{2})} \frac{1}{2i} \left\{ \frac{1}{(\alpha - ik)^2} \underbrace{\int_0^\infty u e^{-u} du}_{\Gamma(2)} - \frac{1}{(\alpha + ik)^2} \underbrace{\int_0^\infty u e^{-u} du}_{\Gamma(2)} \right\}$$

$$\Rightarrow f(\theta) = -\frac{m V_0}{2i \hbar^2 K \sin(\frac{\theta}{2})} \left\{ \frac{1}{(\alpha - ik)^2} - \frac{1}{(\alpha + ik)^2} \right\} \quad \Gamma(2) = 1$$

$$= \frac{m V_0}{2i \hbar^2 K \sin(\frac{\theta}{2})} \left\{ \frac{\alpha^2 - k^2 + 2i\alpha k - \alpha^2 - k^2 - 2i\alpha k}{(\alpha^2 + k^2)^2} \right\}$$

$$\Rightarrow f(\theta) = -\frac{4m V_0 \alpha}{\hbar^2} \frac{1}{(\alpha^2 + k^2)^2}$$

$$\Rightarrow \boxed{f(\theta) = -\frac{4m V_0 \alpha}{\hbar^2} \frac{1}{(\alpha^2 + 4k^2 \sin^2(\frac{\theta}{2}))^2}}$$

$$\frac{d\sigma}{d\Omega} = \frac{16 m^2 V_0^2 \alpha^2}{\hbar^4} \frac{1}{[\alpha^2 + 4K^2 \sin^2(\frac{\theta}{2})]^4}$$

$$\sigma_T = \int_0^{2\pi} \int_0^\pi \left(\frac{d\sigma}{d\Omega} \right) \sin \theta d\theta d\phi$$

⑥ Potencial Gaussiano

$$V(r) = V_0 e^{-\alpha^2 r^2}$$

$$\Rightarrow f(\theta) = -\frac{2m}{\hbar^2 k} \int_0^\infty r e^{-\alpha^2 r^2} \sin(kr) dr$$

$$= -\frac{m}{i\hbar^2 k} \int_0^\infty r \left[e^{-\alpha^2 r^2 + ikr} - e^{-\alpha^2 r^2 - ikr} \right] dr$$

Completamos cuadrados en los argumentos de las exponenciales:

$$e^{-\alpha^2 r^2 + ikr} = e^{-(\alpha r - \frac{ik}{2\alpha})^2 - \frac{k^2}{4\alpha^2}}; \quad e^{-\alpha^2 r^2 - ikr} = e^{-(\alpha r + \frac{ik}{2\alpha})^2 - \frac{k^2}{4\alpha^2}}$$

$$\Rightarrow f(\theta) = -\frac{m}{i\hbar^2 k} e^{-k^2/4\alpha^2} \left\{ \int_0^\infty r e^{-(\alpha r - \frac{ik}{2\alpha})^2} dr - \int_0^\infty r e^{-(\alpha r + \frac{ik}{2\alpha})^2} dr \right\}$$

$$\text{sea } u = \alpha r - \frac{ik}{2\alpha} \Rightarrow dr = \frac{1}{\alpha} du$$

$$f(\theta) = -\frac{m}{i\hbar^2 k \alpha^2} e^{-k^2/4\alpha^2} \left\{ \int_0^\infty u e^{-u^2} du + \frac{ik}{2\alpha} \int_0^\infty e^{-u^2} du - \int_0^\infty u e^{-u^2} du + \frac{ik}{2\alpha} \int_0^\infty e^{-u^2} du \right\}$$

$$f(\theta) = -\frac{m\sqrt{\pi} e^{-k^2/4\alpha^2}}{2\hbar^2 \alpha^3}$$

$$\frac{d\sigma}{d\Omega} = \frac{m^2 \pi}{4\hbar^4 \alpha^6} \exp\left(-\frac{2\sin^2(\theta/2)}{\alpha^2}\right)$$

⑦ $V(r) = \frac{V_0}{(r^2 + d^2)^2}$

$$f(\theta) = -\frac{2mV_0}{\hbar^2 k} \int_0^\infty r \frac{\sin(kr)}{(r^2 + d^2)^2} dr = -\frac{2mV_0}{\hbar^2 k} \frac{1}{4d} e^{-kd} \pi k$$

$$\Rightarrow -\frac{\pi m V_0}{2\hbar^2 d} e^{-kd} \quad k = 2K \sin\left(\frac{\theta}{2}\right)$$

$$\frac{d\sigma}{d\Omega} = |f(\theta)|^2 = \frac{\pi^2 m^2 V_0^2}{4\hbar^4 d^2} e^{-2kd}$$

$$\sigma_T = \frac{\pi^3 m^2 V_0^2}{2\hbar^4 d^2} \int_0^\pi \sin\theta e^{-4Kd \sin(\frac{\theta}{2})} d\theta$$

$$\Rightarrow \frac{\pi^3 m^2 V_0^2}{2\hbar^4 d^2} \frac{1}{4d^2 k^2} \left(1 - (1 + 4dk) e^{-4dk}\right)$$

$$\sigma_T = \frac{\pi^3 m^2 V_0^2}{8\hbar^4 d^4 k^2} \left[1 - (1 + 4dk) e^{-4dk}\right]$$